

Let us consider Background of Fourier Analysis a 2-D Euclidean space as V.S.

Given by:

$$\vec{x} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \in \mathbb{R}^2 \quad \text{s.t. } \|\vec{x}\| = \sqrt{x_1^2 + x_2^2} < \infty; \langle \vec{x}, \vec{y} \rangle = \vec{y}^T \vec{x}$$

Given by: $\vec{x} = q_1 \vec{u}_1 + q_2 \vec{u}_2 = \begin{bmatrix} q_1 \\ q_2 \end{bmatrix}$

Coordinate of \vec{x} along \vec{u}_1 & \vec{u}_2
or
Coefficient of the linear combination.

Orthogonal basis set.
↳ spans complete \mathbb{R}^2

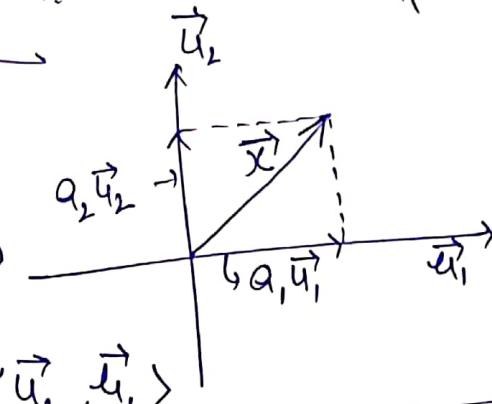
Derivation of q_1 & q_2

Typically, $\vec{u}_1 = \begin{bmatrix} 3 \\ 0 \end{bmatrix}$
 $\vec{u}_2 = \begin{bmatrix} 0 \\ 2 \end{bmatrix}$
 $\vec{x} = \begin{bmatrix} 6 \\ 4 \end{bmatrix}$

i.e. every $\vec{x} \in \mathbb{R}^2$ can be rep. as $\vec{x} = \sum_{k=1}^2 q_k \vec{u}_k$
 $q_k \in \mathbb{R}$.

1. $\langle \vec{x}, \vec{u}_1 \rangle = \cancel{\langle \vec{x}, \vec{u}_1 \rangle}$

$= \langle (q_1 \vec{u}_1 + q_2 \vec{u}_2), \vec{u}_1 \rangle$



$= q_1 \langle \vec{u}_1, \vec{u}_1 \rangle + q_2 \langle \vec{u}_2, \vec{u}_1 \rangle$

\downarrow

$\langle \vec{x}, \vec{u}_1 \rangle = q_1 \|\vec{u}_1\|^2 + 0$

$\downarrow = 0$

$\because \vec{u}_2 \perp \vec{u}_1 \Rightarrow \vec{u}_1^T \vec{u}_2 = [3 \ 0] \begin{bmatrix} 0 \\ 2 \end{bmatrix} = 0$

$q_1 = \frac{\langle \vec{x}, \vec{u}_1 \rangle}{\|\vec{u}_1\|^2} = \frac{\vec{u}_1^T \vec{x}}{\vec{u}_1^T \vec{u}_1} = \frac{[3 \ 0] \begin{bmatrix} 6 \\ 4 \end{bmatrix}}{[3 \ 0] \begin{bmatrix} 3 \\ 0 \end{bmatrix}} = \frac{18}{9} = 2$

Similarly-

$q_2 = \frac{\langle \vec{x}, \vec{u}_2 \rangle}{\|\vec{u}_2\|^2} = \frac{\langle \vec{x}, \vec{u}_2 \rangle}{\|\vec{u}_2\|^2} = \frac{[0 \ 2] \begin{bmatrix} 6 \\ 4 \end{bmatrix}}{[0 \ 2] \begin{bmatrix} 0 \\ 2 \end{bmatrix}} = \frac{8}{4} = 2$

e.g. conducted here.

$\vec{x} = q_1 \vec{u}_1 + q_2 \vec{u}_2 = 2 \vec{u}_1 + 2 \vec{u}_2$

$\begin{bmatrix} 6 \\ 4 \end{bmatrix} = 2 \begin{bmatrix} 3 \\ 0 \end{bmatrix} + 2 \begin{bmatrix} 0 \\ 2 \end{bmatrix}$

Analogously ;

~~Exponential Fourier Series~~

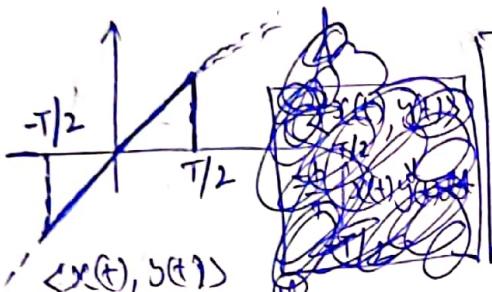
considers

$\mathbb{C} \rightarrow$ complex number

a signal space

$x(t) \in \mathbb{C}$ defines over $-\frac{T}{2} \leq t \leq \frac{T}{2}$

$$\text{s.t. } \|x(t)\|^2 = \langle x(t), x(t) \rangle = \int_{-T/2}^{T/2} x^*(t) x(t) dt$$



$$\|x(t)\|^2 = \int_{-T/2}^{T/2} |x(t)|^2 dt = E_x < \infty$$

Inner product
of $x(t)$ with
 $y(t)$

$$\langle x(t), y(t) \rangle = \int_{-T/2}^{T/2} y^*(t) \cdot x(t) dt = \int_{-T/2}^{T/2} x(t) y(t) dt$$

Such signals can be represented by the linear combination of the following orthogonal basis set.

$$1). \left\{ e^{j k \omega_0 t} ; k \in \mathbb{Z} ; k \in \{-2, -1, 0, 1, 2, \dots\} \right\}$$

$$2). \left\{ 1, \cos(k \omega_0 t), \sin(k \omega_0 t) ; k \in \mathbb{Z} ; k \in \{1, 2, 3, \dots\} \right\}$$

where $\omega_0 = \frac{2\pi}{T}$

Hence
EFS

$$x(t) \approx \sum_{k=-\infty}^{\infty} C_k \cdot e^{j k \omega_0 t} ; \frac{-T}{2} \leq t \leq \frac{T}{2}$$

$$C_k = \frac{\langle x(t), e^{j k \omega_0 t} \rangle}{\|e^{j k \omega_0 t}\|^2} ; k \in \mathbb{Z}$$

Synthesis Eqⁿ

Analysis Eqⁿ

$$\|e^{j k \omega_0 t}\|^2 = \int_{-T/2}^{T/2} |e^{j k \omega_0 t}|^2 dt$$

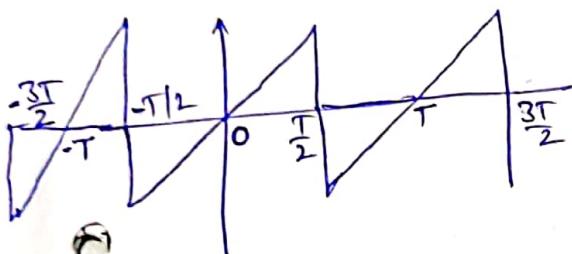
$$= \int_{-T/2}^{T/2} 1 dt = T$$

EFS
Coefficients.

$$C_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) \cdot e^{-j k \omega_0 t} dt$$

PS: Moreover, if $x(t) \Rightarrow x(t+kT)$ is periodic with fundamental time-period, T ; i.e.

then
$$x(t) \Leftarrow \sum c_k e^{jk\omega_0 t} ; t \in \mathbb{R}$$



where $c_k \triangleq \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-j k \omega_0 t} dt$.

Similarly, for $x(t) = x(t+T)$ s.t. $\int_{-T/2}^{T/2} |x(t)|^2 < \infty$.

TF RFS $\Rightarrow x(t) \Leftarrow a_0 \cdot 1 + \sum_{k=1}^{\infty} a_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} b_k \sin(k\omega_0 t) + \dots$

Trigonometric
or
Real FS
only consider even frequencies
i.e. only a_k terms

$$a_0 = \frac{\langle x(t), 1 \rangle / \|1\|^2}{\int_{-T/2}^{T/2} |1|^2 dt} = \frac{\int_{-T/2}^{T/2} x(t) \cdot 1 \cdot dt}{\int_{-T/2}^{T/2} |1|^2 dt} = \frac{\int_{-T/2}^{T/2} x(t) dt}{\int_{-T/2}^{T/2} dt}$$

$$a_0 = \frac{1}{T} \int_{-T/2}^{T/2} x(t) dt$$

$$a_k = \frac{\langle x(t), \cos(k\omega_0 t) \rangle}{\|\cos(k\omega_0 t)\|^2} =$$

$$\begin{aligned} \|\cos(k\omega_0 t)\|^2 &= \int_{-T/2}^{T/2} \cos^2(k\omega_0 t) dt = \int_{-T/2}^{T/2} \left(\frac{1}{2} + \frac{\cos(2k\omega_0 t)}{2} \right) dt \\ &= \int_{-T/2}^{T/2} \frac{1}{2} dt + \int_{-T/2}^{T/2} \cos(2k\omega_0 t) dt = \frac{1}{2} T + \frac{\sin(2k\omega_0 \cdot \frac{T}{2})}{2 \cdot 2k\omega_0} \Big|_{-T/2}^{T/2} \end{aligned}$$

$$a_k = \frac{T}{2} + \frac{\cos\left(k\omega_0 \cdot \frac{T}{2}\right) - \cos\left(-k\omega_0 \cdot \frac{T}{2}\right)}{4\omega_0}$$

$$= \frac{T}{2} + \frac{\cos\left(k \cdot \frac{2\pi}{T} \cdot t\right) - \cos\left(-k \cdot \frac{2\pi}{T} \cdot t\right)}{4\omega_0}$$

$$\left| \left| \cos(k\omega_0 t) \right| \right|^2 = \frac{T}{2} + \frac{1 - 1}{4\omega_0} \Rightarrow \boxed{\left| \left| \cos(k\omega_0 t) \right| \right|^2 = \frac{T}{2}}$$



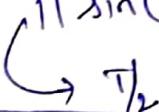


$$a_k = \frac{\langle x(t), \cos(k\omega_0 t) \rangle}{\| \cos(k\omega_0 t) \|^2}$$

$$\boxed{a_k = \frac{2}{T} \int_{-T/2}^{T/2} x(t) \cos(k\omega_0 t) dt}$$

Similarly

$$\boxed{b_k = \frac{\langle x(t), \sin(k\omega_0 t) \rangle}{\| \sin(k\omega_0 t) \|^2} = \frac{2}{T} \int_{-T/2}^{T/2} x(t) \sin(k\omega_0 t) dt}$$



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Relationship 1/2 EFS \wedge TFS \Leftrightarrow

Relationship b/w EFS & TFS.

0

EFS Coefficient of $x(t)$ $\Rightarrow C_K = \frac{a_K}{2} + j \frac{b_K}{2}$ } Holds true even if $x(t)$ is complex valued.

May be real or complex valued.

$$C_{-K} = \frac{a_{-K}}{2} - j \frac{b_{-K}}{2}$$

Reciprocal

EFS Let $x(t) = x(t+T)$;
 $x(t) \leq \sum_{k=-\infty}^{\infty} c_k e^{j k \omega_0 t}; \quad \omega_0 = \frac{2\pi}{T}$

$$c_k = \frac{1}{T} \int_{-T/2}^{T/2} x(t) e^{-jk\omega_0 t} dt; \quad k \in \mathbb{Z}$$

TFS

$$x(t) \leq a_0 + \sum_{k=1}^{\infty} a_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} b_k \sin(k\omega_0 t)$$

$$a_0 = \frac{1}{T} \int_{-T/2}^{T/2} x(t) dt; \quad b_k = \frac{2}{T} \int_{-T/2}^{T/2} x(t) \sin(k\omega_0 t) dt$$

$$a_k = \frac{2}{T} \int_{-T/2}^{T/2} x(t) \cos(k\omega_0 t) dt$$

$T/2$

$-T/2$

$T/2$

$-T/2$

$T/2$

$-T/2$

$T/2$

$-T/2$

$$a_{-k} = a_k \Rightarrow a_k \text{ is always an even function of } k$$

$$b_{-k} = -b_k \Rightarrow b_k \text{ is always an odd func of } k$$

$$c_k = \frac{a_k}{2} + j \frac{b_k}{2} \quad \Rightarrow \quad ②$$

Eq ①, ②, ③ & ④ are true irrespective of $x(t)$ is real or complex. Furthermore ① & ②

BUT

$$c_{-k} = c_k^* \Rightarrow$$

iff a_k & b_k are real.

$$a_k = c_k + c_{-k} \quad ③$$

$$b_k = c_{-k} - c_k \quad ④$$

$x(t)$ is real.

So,

Conclusion

$$x(t) = x(t)^*$$

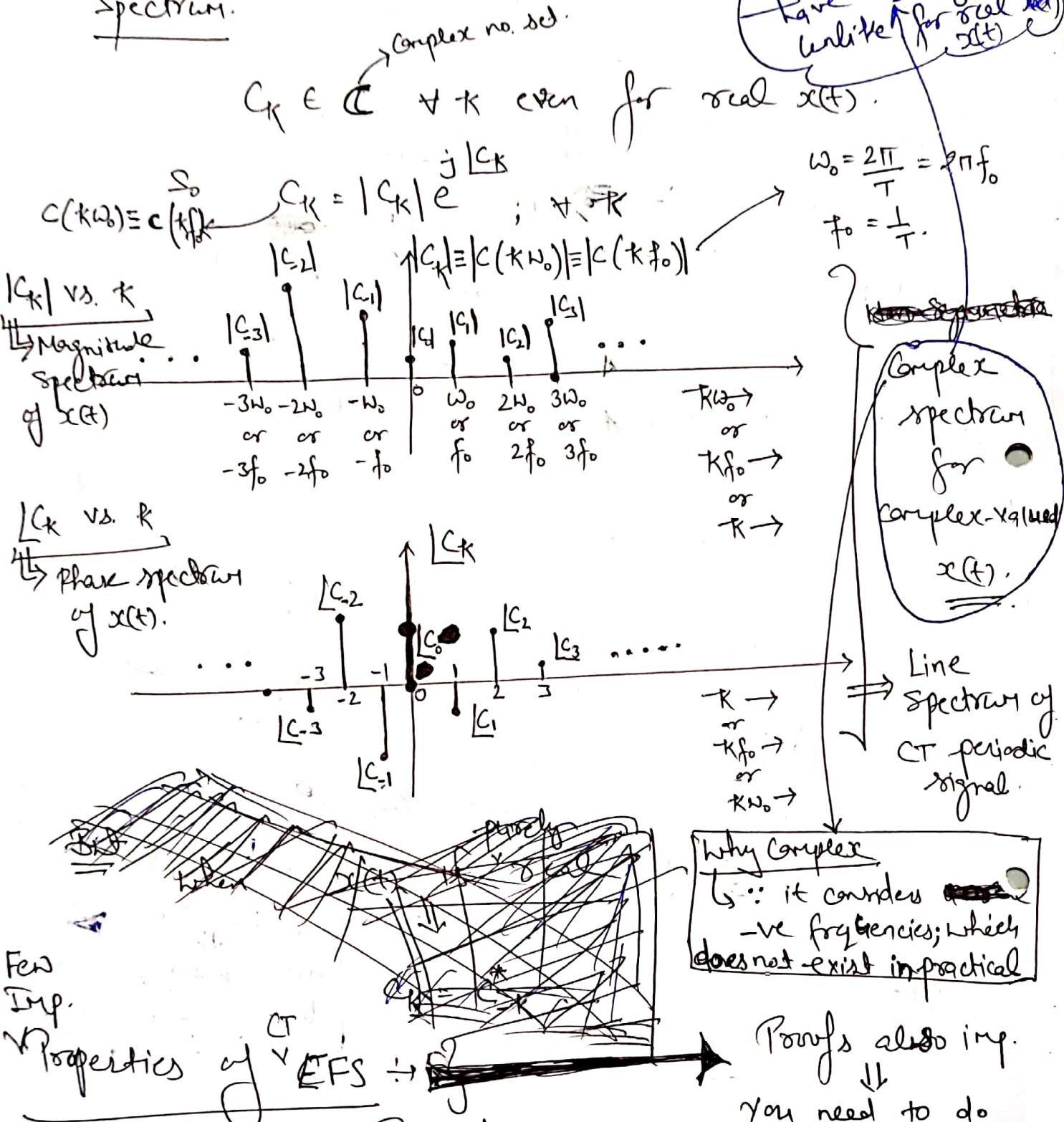
i.e.
real

\Leftrightarrow CTFs

$$c_k = c_{-k}^*$$

Conjugate Symmetric
for real $x(t)$.

Spectrum.



$x(t) \leftrightarrow C_K$

$\text{purely } x^*(t) \leftrightarrow C_{-K}^*$

1.1 If $x(t)$ is "real" $\Rightarrow x(t) = x^*(t) \leftrightarrow C_K = C_{-K}^*$ (Conj Symmetric CS)

1.2 If $x(t)$ is purely imag. $\Rightarrow x(t) = -x^*(t) \leftrightarrow C_K = -C_{-K}^*$ (Anti-Symmetric CAS)

(2)

1.3. If $x(t)$ is neither purely real nor purely imag. $\Rightarrow C_K$ is neither CS nor CAS.

For Case 1.1:

When $x(t)$ is purely Real

$$C_K = \frac{a_K}{2} - j \frac{b_K}{2}$$

$$C_{-K} = \frac{a_K}{2} + j \frac{b_K}{2} = C_K^*$$

$$C_K = C_K^*$$

$$|C_K| = \sqrt{|a_K|^2 + |b_K|^2} = |C_{-K}|$$

$$\angle C_K = -\tan^{-1}\left(\frac{b_K}{a_K}\right) = C_{-K}$$

$$C_K = \pm \tan^{-1}\left(\frac{b_K}{a_K}\right) = C_{-K}$$

$$|C_K| = |C_{-K}| \Rightarrow \text{Even Symmetric}$$

$$C_K = -C_{-K} \Rightarrow \text{Odd Symmetric}$$

For Case 1.2:

When $x(t)$ is purely imag. $\Rightarrow a_K, b_K$ are purely imag.

$$C_K = \frac{a_K}{2} - j \frac{b_K}{2}$$

$$C_{-K} = \frac{a_K}{2} + j \frac{b_K}{2} = -C_K^* = -\frac{a_K}{2} - j \frac{b_K}{2}$$

$$|C_K| = \sqrt{|a_K|^2 + |b_K|^2} = |C_{-K}|$$

$$\angle C_K = +\tan^{-1}\left(\frac{a_K}{b_K}\right) = -C_{-K}$$

Even Symm.

Odd Symmetric

i.e. Either $x(t)$ is purely Real or purely Imag. $\Rightarrow |C_K|$ is even of 'k'. $|C_K|$ is odd of 'k'.

$|C_K|$ is even of 'k'.
 $|C_K|$ is odd of 'k'.

signal is real valued

' ω ' are zero not unreal; ie purely.

(3)

Real Spectrum (Real FS spectrum)

Assuming $x(t)$ is purely real ~~then $C_K = \bar{C}_{-K}$~~

then

$$C_K = |C_K| e^{j\angle C_K} \stackrel{\text{s.t.}}{\Rightarrow} \boxed{|C_K| = |C_{-K}|} \quad \& \quad \boxed{\angle C_K = -\angle C_{-K}}$$

TFS \Rightarrow

$$x(t) = a_0 + \sum_{k=1}^{\infty} a_k \cos(k\omega_0 t) + \sum_{k=1}^{\infty} b_k \sin(k\omega_0 t) \quad \boxed{\text{TFS}}$$

Ans.

$$a \cos \theta + b \sin \theta = r \cos(\theta - \phi)$$

$$= \underbrace{r \cos \phi}{a} \cos \theta + \underbrace{r \sin \phi}{b} \sin \theta \quad \left| \begin{array}{l} r = \sqrt{a^2 + b^2} \\ \phi = \tan^{-1} \frac{b}{a} \end{array} \right.$$

So Assuming

$$\text{P1} \rightarrow a_k \cos(k\omega_0 t) + b_k \sin(k\omega_0 t) = r_k \cos(k\omega_0 t - \phi_k).$$

$$\boxed{a_k = r_k \cos(\phi_k)} \\ \boxed{b_k = r_k \sin(\phi_k)}$$

where

$$\boxed{r_k = \sqrt{a_k^2 + b_k^2} = 2|C_k|} \\ \boxed{\phi_k = \tan^{-1} \left(\frac{b_k}{a_k} \right) = -\angle C_k}$$

Further for $k=0$

$$\text{P2} \rightarrow \text{Assuming via } b_0 = 0 \Rightarrow \boxed{r_0 = a_0} \\ \boxed{\phi_0 = 0}$$

$$\boxed{r_0 = \sqrt{a_0^2 + b_0^2} = a_0} \\ \boxed{\phi_0 = \tan^{-1} \left(\frac{b_0}{a_0} \right) = 0}$$

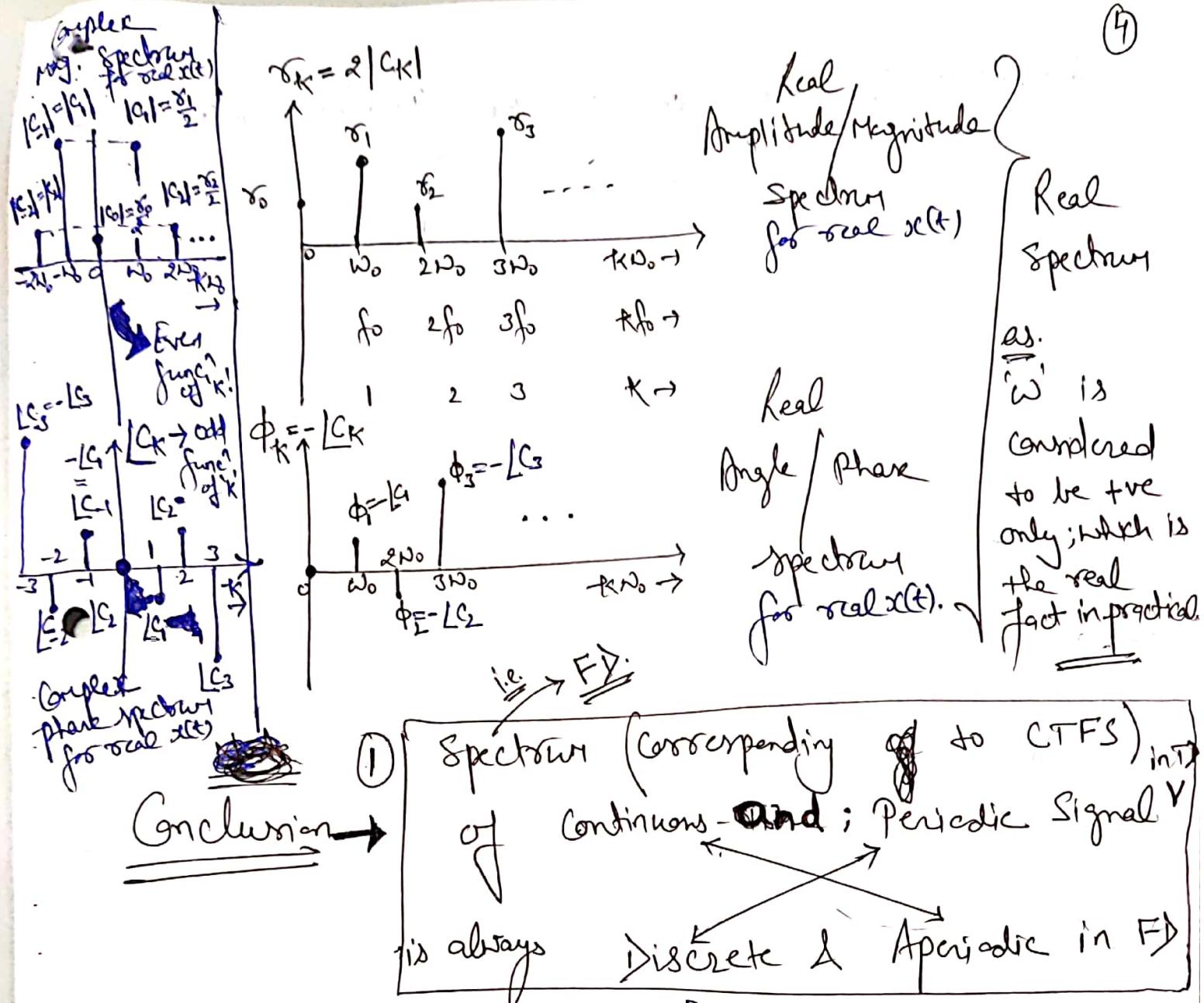
TFS

can be re-written as

$$x(t) \approx \sum_{k=0}^{\infty} r_k \cos(k\omega_0 t - \phi_k) \Rightarrow \begin{array}{l} \text{Harmonic Amplitude} \\ \text{Angle} \end{array}$$

from of CTFs

$r_k \rightarrow$ Amplitude at $k\omega_0$, $\phi_k \rightarrow$ Phase at $k\omega_0$ where $k\omega_0 \rightarrow$ harmonics.



② Real magnitude spectrum of real-valued signal reduces to complex magnitude spectrum by reducing/dividing the amplitude/magnitude at all the freq. equally in complex magnitude spectrum.

③ Key Points to Note while looking at Complex Spectrum:

There's (i) whenever ~~the~~ freq. dependent info (e.g. 'Bandwidth') needs to be addressed; only +ve freq. will be taken care.

(ii) whenever, the magnitude dependent info e.g. Energy/Power needs to be addressed / find; ~~both~~ Both -ve & +ve freq. will be taken care in the magnitude spectrum.

Lec-3.

Existence & Convergence conditions for CTFs of a continuous periodic signal \rightarrow (Developed by P.L. Dirichlet).

Recommended Signal System Reading by A.V. Oppenheim

Dirichlet conditions \rightarrow guarantees that $x(t) = \sum_{k=-\infty}^{\infty} c_k e^{j2\pi k t}$, $\forall t$

are only the sufficient conditions (i.e. not necessary) to have EFS of continuous periodic signals or of periodic signals with finite no. of finite-size discontinuity.

Exception $x(t) = \sum_{k=-\infty}^{\infty} \delta(t - kT) \leftrightarrow c_k = \frac{1}{T}$

except at isolated values of t for which $x(t)$ is discontinuous.

At these values of t , the infinite series, $\sum_{k=-\infty}^{\infty} c_k e^{jk\omega_0 t}$ converges to the

average of the values on either side of the discontinuity.

Condition-1 (Existence Condition of CTFs $\Rightarrow |c_k| < \infty \forall k$).

$$\int_T |x(t)| dt < \infty$$

$$|c_k| = \left| \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt \right| \leq \frac{1}{T} \int_T |x(t)| |e^{-jk\omega_0 t}| dt$$

$x(t) = \frac{1}{t}; 0 < t \leq 1$
violates Condition-1
 $c_k = ??$

$$|c_k| \leq \frac{1}{T} \int_T |x(t)| dt$$

$$\therefore \text{if } \int_T |x(t)| dt < \infty \Rightarrow |c_k| < \infty \forall k.$$

Condition-2 \rightarrow

$x(t)$ must have finite no. of maxima & minima during any suff. single period of the signal.

$x(t) = \sin(\frac{2\pi}{T} t)$, $0 < t \leq 1$ meets condition 1 but not condition-2.

Condition-3 \rightarrow In any finite interval of time, there are only a finite no. of discontinuities. Furthermore, each of these discontinuities is of finite size.

Conclusions →

Case 1 For periodic ~~non~~ signals → Has no discontinuities
 e.g. $e^{j\omega_0 t}$, $\sin(\omega_0 t)$, $\cos(\omega_0 t)$ etc.; $\sum_{K=-\infty}^{\infty} C_K e^{j K \omega_0 t} = x(t) + t$.
 finite-sized

Case 2 For a periodic signal with a finite no. of discontinuities in each period; e.g. periodic square wave,

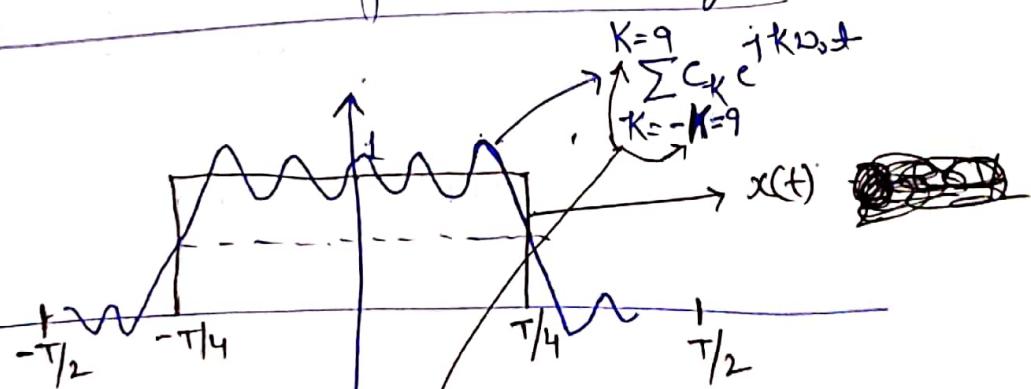
$$\sum_{K=-\infty}^{\infty} C_K e^{j K \omega_0 t} = x(t) + t \text{ except at 't' where discontinuities are located. let's say}$$

$$\sum_{K=-\infty}^{\infty} C_K e^{j K \omega_0 t} = \frac{x(t) + x(t^-)}{2} \quad 't' \text{ where discontinuities are located}$$

Furthermore; $e(t) = x(t) - \sum_{K=-\infty}^{\infty} C_K e^{j K \omega_0 t}$ has zero energy over one 'T'.
 i.e. $\int_T |e(t)|^2 dt = 0$.

Gibbs Phenomenon occurs for case-2 signals ↴

i.e.



Conclusion ↴

Even for large values of 'K' an overshoot occurs at discontinuity but the ripples goes lower & lower as $K \rightarrow \infty$.

Parseval's Relation (CTFS).

$$\int_{-\infty}^{\infty} f(t) \overline{f(t)} dt = \sum_{n=-\infty}^{\infty} C_n \overline{C_n}$$

$$P_c = ?$$

But

$$C_n = \frac{a_n}{2} - j \frac{b_n}{2}$$

$$\frac{1}{T} \int_{-T/2}^{T/2} f(t) f^*(t) dt$$

$$C_0 = a_0 \Rightarrow |C_0|^2 = P_c \Rightarrow DC \text{ Power}$$

$$\frac{1}{T} \int_{-T/2}^{T/2} f(t) \left(\sum_{n=-\infty}^{\infty} C_n e^{j n \omega t} \right) dt$$

$$C_n = \frac{a_n}{2} + j \frac{b_n}{2}$$

$$= \frac{1}{T} \int_{-T/2}^{T/2} f(t) \left(\sum_{n=-\infty}^{\infty} C_n e^{-j n \omega t} \right) dt$$

$$|C_n|^2 = \frac{a_n^2 + b_n^2}{4}$$

$$= \sum_{n=-\infty}^{\infty} C_n \cdot \frac{1}{T} \int_{-T/2}^{T/2} f(t) e^{-j n \omega t} dt$$

$$|C_n|^2 = \frac{a_n^2 + b_n^2}{4}$$

$$= \sum_{n=-\infty}^{\infty} C_n^2$$

$$|C_0|^2 + |C_1|^2 + |C_n|^2 = \frac{a_0^2}{2} + \frac{b_0^2}{2} \Rightarrow AC \text{ Power}$$

$$P_c = \sum_{n=-\infty}^{\infty} |C_n|^2$$

Replace in eq shown on 1st page

$$P = C_0^2 + \sum_{n=1}^{\infty} \{ |C_n|^2 + |C_{-n}|^2 \}$$

$$= C_0^2 + 2|C_1|^2 + \sum_{n=1}^{\infty} |C_n|^2$$

$$|C_1|^2 + |C_2|^2 + |C_3|^2 \dots$$

$$= C_0^2 + 2|C_1|^2 + \sum_{n=1}^{\infty} |C_n|^2$$

$$= C_0^2 + \sum_{n=1}^{\infty} |C_n|^2 + \sum_{n=1}^{\infty} |C_{-n}|^2$$

$$= C_0^2 + \sum_{n=1}^{\infty} |C_n|^2 + \sum_{n=1}^{\infty} |C_n|^2$$

$$(a_n^2 + b_n^2)$$

$$P_{f(t)} = \sum_{n=-\infty}^{\infty} |C_n|^2 \rightarrow \text{Parseval's Power Relation}$$

for periodic power signal.

$$\frac{1}{T} \int_{-T/2}^{T/2} |f(t)|^2 dt = \sum_{n=-\infty}^{\infty} |C_n|^2 \Rightarrow \text{Average Power of Periodic Signal}$$

$$= a_0^2 + \sum_{n=1}^{\infty} a_n^2 + b_n^2$$

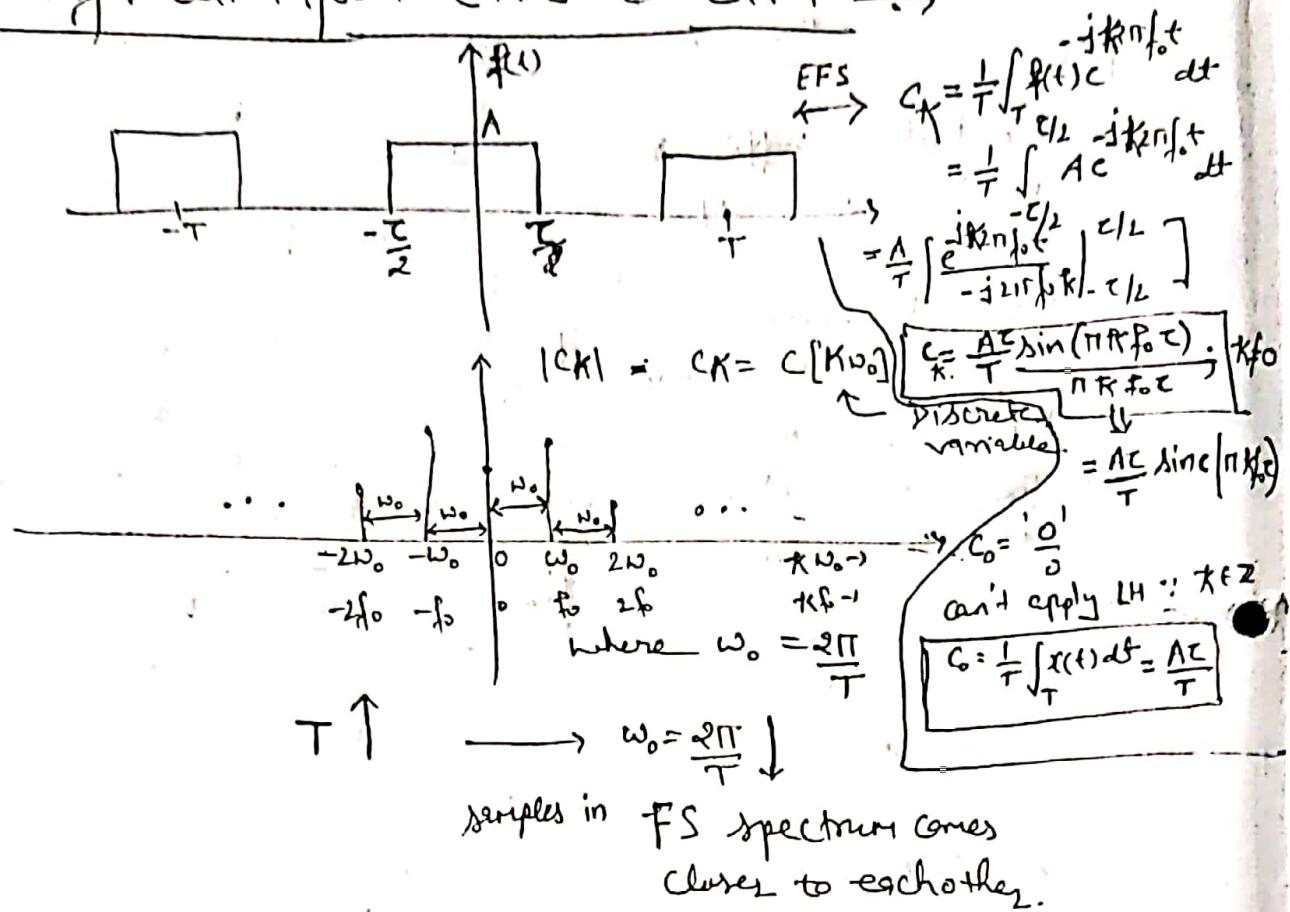
$|C_n|^2$ or Power Spectral density ($P_{f(t)}$)

power spectrum of the given CT periodic signal

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Scanned with CamScanner

Transformation from CTS to CTFT



$T \rightarrow 0$

$$\omega_0 = \frac{2\pi}{T} \rightarrow 0 \equiv \delta\omega \text{ or } d\omega$$

$K\omega_0 \rightarrow \omega \leftarrow \text{continuous variable.}$

$$C\{K\omega_0\} = \frac{1}{T} \int_{-T/2}^{T/2} f(t) e^{-jK\omega_0 t} dt$$

$$T C\{K\omega_0\} = \int_{-T/2}^{T/2} f(t) \cdot e^{-j\omega_0 t} dt$$

$$\text{Let } T C\{K\omega_0\} = F\{K\omega_0\}$$

$$F\{K\omega_0\} = \int_{-T/2}^{T/2} f(t) e^{-jK\omega_0 t} dt$$

$$f(t) = \sum_{K=-\infty}^{\infty} \dots, C_K \cdot e^{jK\omega_0 t}$$

$$= \sum_{K=-\infty}^{\infty} C\{K\omega_0\} e^{jK\omega_0 t} = \sum_{K=-\infty}^{\infty} \frac{F\{K\omega_0\}}{T} e^{jK\omega_0 t}$$

$$f(t) = \sum_{K=-\infty}^{\infty} F\{K\omega_0\} e^{jK\omega_0 t} \cdot \frac{\omega_0}{2\pi}$$

As $T \rightarrow \infty$, $\omega_0 \rightarrow 0$, $n\omega_0 \rightarrow \omega$; $\sum \rightarrow \int$
 or
 $\frac{1}{T} \sum f(t) \rightarrow \frac{1}{2\pi} \int F(\omega) d\omega$.

$f(t)$ periodic $\xrightarrow{T=0}$ Aperiodic & ~~continuous~~
 Discrete Frequency \rightarrow Continuous Frequency Spectrum Spectrum.

Analysis Eq

$$F(\omega) = \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt$$

$F\{K\omega_0\} \rightarrow F(\omega)$.

Fourier Transform

Synthesis Eq

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{j\omega t} d\omega$$

Inverse Fourier Transform

Note

Fourier Series existed only for Periodic signals but Fourier Transform will exist for both Aperiodic as well as periodic signals but FT of periodic signals are Limitedly defined i.e. $FT = \infty$ at one freq.

E.g. Find FT of $e^{-at} u(t)$, $a > 0$



$$A \rightarrow F(\omega) = \int_{-\infty}^{\infty} e^{-at} u(t) e^{-j\omega t} dt$$

$$\begin{aligned} &= \int_0^{\infty} e^{-at} e^{-j\omega t} dt \\ &= \int_0^{\infty} e^{-(a+j\omega)t} dt \\ &= \left[\frac{e^{-(a+j\omega)t}}{-(a+j\omega)} \right]_0^{\infty} \end{aligned}$$

$$F(\omega) = \frac{1}{a+j\omega} \Rightarrow F(\omega) = \frac{1}{a+j\omega} \frac{(a-j\omega)}{(a-j\omega)} = \frac{a-j\omega}{a^2+\omega^2}$$

$$F(\omega) = \frac{a}{a^2+\omega^2} - j \frac{\omega}{a^2+\omega^2}$$

$$\operatorname{Re}[F(\omega)] = \frac{a}{a^2+\omega^2} \Rightarrow \text{Even function of } \omega.$$

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$$\text{Im}\{F(\omega)\} = \frac{-\omega}{q^2 + \omega^2} \rightarrow \text{odd func of } \omega$$

Hence $F(\omega) \Rightarrow \text{Conjugate Symmetric}$

i.e. $F(\omega) = F^*(\omega)$

Note: $\xrightarrow{\text{FT}} \boxed{F^*(-\omega) = F(\omega)}$

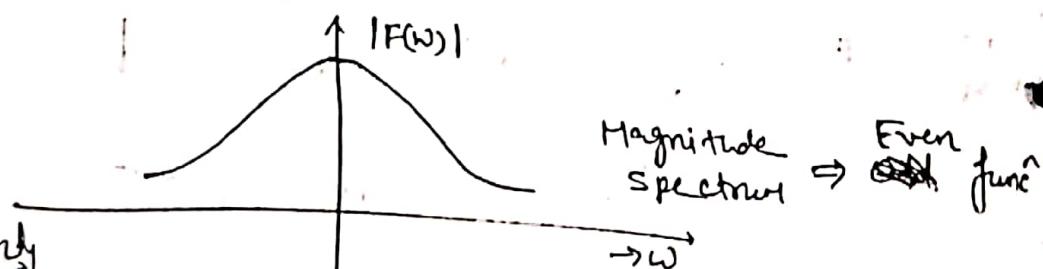
If $f(t)$ is real valued, $F(\omega)$ must be conjugate symmetric in frequency; $|F(\omega)| \rightarrow \text{even}$ & $\boxed{F(\omega) \rightarrow \text{odd}}$

Polar form: $F(\omega) \rightarrow \text{Frequency Response}$

Frequency Response $\leftarrow F(\omega) \rightarrow |F(\omega)| \quad \boxed{|F(\omega)|}$

Magnitude Response \uparrow Phase Response

$$|F(\omega)| = \frac{1}{\sqrt{q^2 + \omega^2}} ; \quad \angle F(\omega) = -\tan^{-1} \frac{\omega}{q}$$



Conjugate Property

of FT

$$\tilde{f}(t) \leftrightarrow F(\omega)$$

$$f^*(t) \leftrightarrow F^*(-\omega)$$

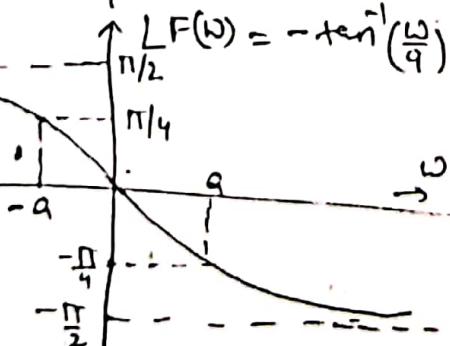
$$\begin{aligned} \text{If } f(t) \rightarrow \text{real} \\ f(t) = \tilde{f}(t) \leftrightarrow F(\omega) = \tilde{F}(\omega) \end{aligned}$$

\rightarrow Conjugate Symmetric

Conj-L

$$f(t) - \tilde{f}(t) \leftrightarrow \boxed{F(\omega) - \tilde{F}(\omega)}$$

\rightarrow Conjugate M.T. Sym.



Phase \Rightarrow odd func

referred to as the Dirichlet conditions, require that:

1. $x(t)$ be absolutely integrable; that is,

$$\int_{-\infty}^{+\infty} |x(t)| dt < \infty. \quad (4.13)$$

2. $x(t)$ have a finite number of maxima and minima within any finite interval.
3. $x(t)$ have a finite number of discontinuities within any finite interval. Furthermore each of these discontinuities must be finite.

Parseval's Energy Relation
(With CTFT)

Spectral Density Property

~~Signal energy = spectral density × time energy input.~~

$$E = \int_{-\infty}^{\infty} |f(t)|^2 dt$$

$$= \int_{-\infty}^{\infty} f(\omega) \cdot f(\omega)^* dt$$

$$= \int_{-\infty}^{\infty} f(\omega) \left(\frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} F(\nu) e^{-j\nu t} d\nu \right)^* dt$$

$$= \int_{-\infty}^{\infty} f(\omega) \left(\int_{-\infty}^{\infty} F(\nu) e^{j\nu t} d\nu \right)^* dt$$

Changing the order of integration

$$= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} F(\nu)^* \left(\int_{-\infty}^{\infty} f(\omega) e^{-j\nu t} d\omega \right) d\nu$$

$$C = a + jb$$

$$C^* = a - jb$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(\nu)^* F(\nu) d\nu$$

$$C \cdot C^* = a^2 + b^2$$

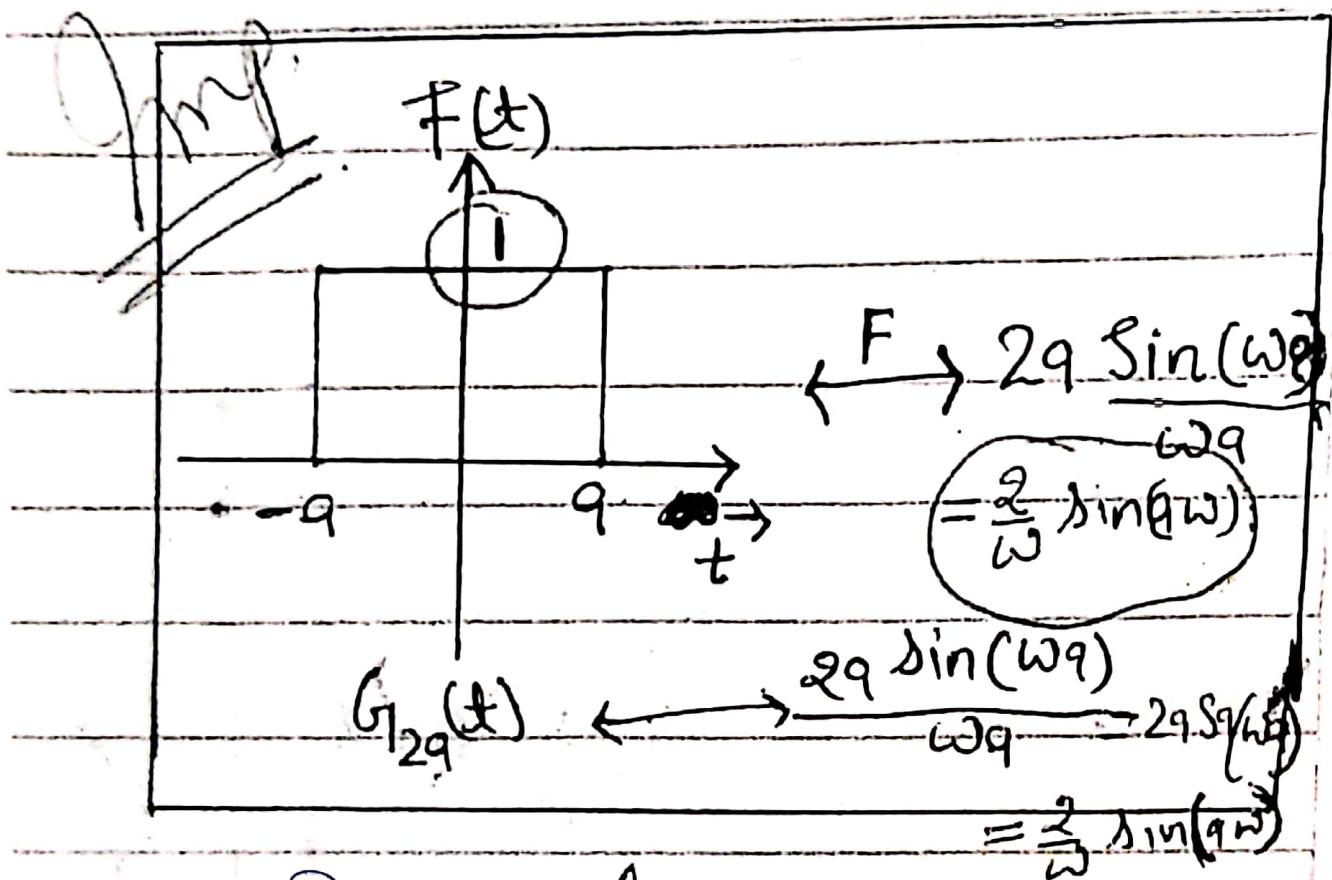
$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(\nu)^* F(\nu) d\nu$$

$$= |C|^2$$

$$E = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (|F(\nu)|^2) d\nu \rightarrow \text{Physical Energy}$$

$\Rightarrow F(\nu)$
Energy
spectral
Density
of $f(t)$.

or
Parseval's energy
theorem.



Duality Property:

$$x(t) \xleftrightarrow{F} X(\omega)$$

$$X(t) \xleftrightarrow{F} 2\pi x(-\omega)$$

CJ

$$\begin{aligned} e^{-2t} u(t) &\xleftrightarrow{F} \frac{1}{2+tj\omega} \\ \frac{1}{2+jt} &\xleftrightarrow{F} 2\pi e^{2\omega} u(-\omega) \end{aligned}$$

$$\begin{aligned} \delta(t) &\xleftrightarrow{F} 1 \\ 1 &\xleftrightarrow{F} 2\pi \delta(-\omega) \\ &= 2\pi \delta'(\omega) \end{aligned}$$

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