

HIGH FREQUENCY COMMUNICATION SYSTEMS

Lecture 4

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LECTURE OUTLINE

- · Transmission Lines Theory & Analysis
- $\cdot \ \, \text{Load Mismatching}$
- · The Magic of Quarter-wave Transformer

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TRANSMISSION LINE THEORY

- · A technique to bridge the gap between basic circuit analysis and electromagnetic fields theory
 - · Lots of similarities and analogies
- · Commonly used to design microwave devices and circuits
- . We use transmission lines when the electrical length of the device is greater than $\lambda/10$.

TL & EM FIELD THEORY - DUALITY

- In wave scattering problems, we place the co-ordinate system at the boundary
- In transmission lines, we use the two different lines as the boundary
- For open problems, the load is considered to be at infinity.

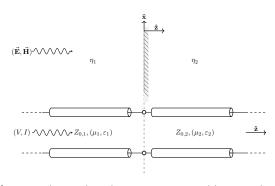


Figure 1: The analogy between wave problems and transmission lines.

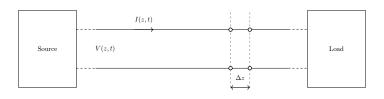
TL & EM FIELD THEORY DUALITY

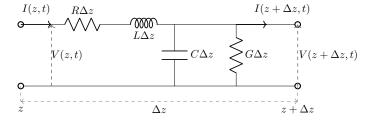
TL Theory	EM Field Theory
$V_{1} = V_{0}e^{-j\gamma_{1}z} \left(1 + \Gamma_{L}e^{j2\gamma_{1}z}\right)$ $I_{1} = \frac{V_{0}}{Z_{0,1}}e^{-j\gamma_{1}z} \left(1 + \Gamma_{L}e^{j2\gamma_{1}z}\right)$ $\Gamma_{L} = \frac{Z_{L}-Z_{0,1}}{Z_{L}+Z_{0,1}} = \frac{Z_{0,2}-Z_{0,1}}{Z_{0,2}+Z_{0,1}}$ $V_{2} = TV_{0}e^{-j\gamma_{2}z}$ $I_{2} = T\frac{V_{0}}{Z_{0,2}}e^{-j\gamma_{2}z}$	$E_{x,1} = E_0 e^{-jk_1 z} \left(1 + \Gamma e^{j2k_1 z} \right)$ $H_{x,1} = \frac{E_0}{\eta_0} e^{-jk_1 z} \left(1 + \Gamma e^{j2k_1 z} \right)$ $\Gamma = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1}$ $E_{x,2} = TE_0 e^{-jk_2 z}$ $H_{x,2} = T\frac{E_0}{\eta_2} e^{-jk_2 z}$
$T = \frac{7Z_{0,2}}{2Z_{0,2}}$ $T = \frac{2Z_{0,2}}{Z_{0,2} + Z_{0,1}}$	$T = \frac{2\eta_2}{\eta_2 + \eta_1}$

TL THEORY AND CIRCUIT ELEMENTS

- · For a transmission line, we use a distributed circuit approach
 - · Energy stored in the magnetic field \rightarrow L
 - · Energy stored in the electric field \rightarrow C
 - · Conductive losses $\rightarrow R$
 - · Dielectric losses $\rightarrow G$
- · All the circuit elements are expressed per unit length
- We can express the voltage and current at any given point z and time t.

THE TRANSMISSION LINE EQUATION





- · Solving the electric circuit leads to a second order differential equation
- · Analogous to the wave equation
 - · Hence the analogies between voltage V and electric field $\vec{\mathbf{E}}$

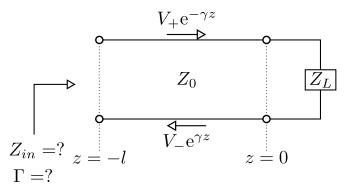
$$\left\{ \frac{\partial^2}{\partial z^2} - \left[(R + j\omega L)(G + j\omega C) \right] \right\} V(z, t) = 0$$

The solution is of the type:

$$V(z) = V_{+}e^{-\gamma z} + V_{-}e^{+\gamma z}$$

where the propagation constant $\gamma = \sqrt{(R+j\omega L)(G+j\omega C)}$. For a lossless case R=G=0.

- · For a transmission line terminated with a load Z_L
 - We get reflections in the line (superposition incoming and reflected wave make up the total voltage)
- The input impedance depends on the Z_L as well as the length l of the transmission line
- · For Transmission line analysis, we shift the origin to the load.



The voltage reflection coefficient Γ can for the line is

$$\Gamma = \frac{V_{ref}}{V_{inc}} = \frac{V_{-}e^{\gamma z}}{V_{+}e^{-\gamma z}}$$
$$= \frac{V_{-,z=0}}{V_{+,z=0}}e^{2\gamma z} = \Gamma_{L}e^{2\gamma z}$$

We can express the voltage at any point on the line as:

$$V(z) = V_{+} \left(e^{-\gamma z} + \Gamma_{L} e^{\gamma z} \right)$$

After finding a similar expression for the current, we can write the impedance as:

$$Z(z) = Z_0 \frac{e^{-\gamma z} + \Gamma_L e^{-\gamma z}}{e^{-\gamma z} - \Gamma_L e^{-\gamma z}}$$

At z = 0, $Z(0) = Z_L$, and we get:

$$Z_{L} = Z_{0} \frac{1 + \Gamma_{L}}{1 - \Gamma_{L}} \implies \Gamma_{L} = \frac{Z_{L} - Z_{0}}{Z_{L} + Z_{0}}$$

At z = -l, the input impedance $Z(-l) = Z_{in}$ is given by:

$$Z(-l) = Z_0 \frac{e^{\gamma l} + e^{-\gamma l} \Gamma_L}{e^{\gamma l} - e^{\gamma l} \Gamma_L}$$

$$= Z_0 \frac{(Z_L + Z_0) e^{\gamma l} + (Z_L - Z_0) e^{-\gamma l}}{(Z_L + Z_0) e^{\gamma l} - (Z_L - Z_0) e^{-\gamma l}}$$

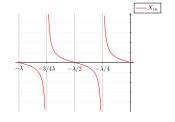
$$= Z_0 \frac{Z_L + Z_0 \tanh \gamma l}{Z_0 + Z_L \tanh \gamma l}$$

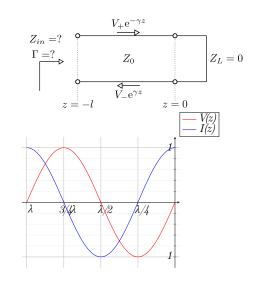
For lossless TLs $\gamma = j\beta$, therefore,

$$Z(-l) = Z_0 \frac{Z_L + jZ_0 \tan \beta l}{Z_0 + jZ_L \tan \beta l}$$

For
$$Z_L=0$$
,
$$Z_{in}=Z_0 j \tan(\beta l)$$

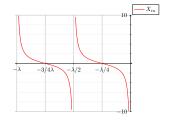
$$\Gamma_{in}=-1$$

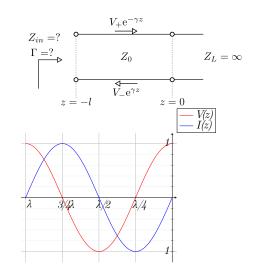




For
$$Z_L = \infty$$
,
$$Z_{in} = -Z_0 j \cot(\beta l)$$

$$\Gamma_{in} = 1$$



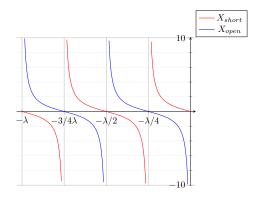


For
$$Z_L = Z_0$$
, $Z_{in} = Z_0$ $Z_{in} = ?$ Z_0 $Z_L = Z_0$ $Z_L = Z_0$

- \cdot Ideally we desire this condition for maximum power transfer.
- · No reflections take place in this case

SPECIAL CASES OF LOAD IMPEDANCE - THE QUARTER-WAVE TRANSFORMER

- · We see that for terminated loads with special lengths that are multiples of quarter wavelengths $(\lambda/4 + n\lambda/2)$ for $n = 1, 2, 3, \cdots$
- \cdot The load impedance is inverted
 - · Open-circuit load becomes short-circuit and vice versa
- For such a line, the input impedance is $Z_{in} = \frac{Z_0^2}{Z_L}$



- · We want maximum power transferred to the load from the transmission line.
 - · For that to happen, we require the impedance matching of the transmission line.
- · The reflections in the TL are undesirable and lead to a standing wave generation ($V = V_+ + V_-$).
- The voltage standing wave ratio (VSWR) measures the quality of impedance matching in TL.

Recall,

$$V(z) = V_{+} \left(e^{-\gamma z} + \Gamma e^{\gamma z} \right)$$

The values of the voltage reflection coefficient range from $-1 \leq |\Gamma| \leq 1.$

The VSWR is the ratio of the maximum and minimum absolute values of the voltage

$$VSWR = \frac{|V_{max}|}{|V_{min}|}$$

In other words,

$$\begin{aligned} |V_{\textit{max}}| &= V_0 \left(1 + |\Gamma| \right) \\ |V_{\textit{min}}| &= V_0 \left(1 - |\Gamma| \right) \\ VSWR &= \frac{1 + |\Gamma|}{1 - \Gamma} \end{aligned}$$

- Often we use the decibels (dB) to express the power levels and performance parameters
- · Here we look into the the terms return loss and insertion loss

For mismatched load, the power *returned* to the generator is the return loss:

$$RL = -20 \log |\Gamma| (dB)$$

• For passive circuits like an antenna, the return loss is always a non-negative number.

- · Like reflection, we also have transmission coefficient
- · This is useful when we have lots of elements in a transmission line.

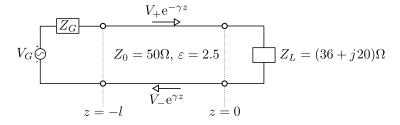
The trasmission coefficient is defined as:

$$T = 1 + \Gamma = \frac{2Z_1}{Z_0 + Z_1}$$

The insertion loss becomes:

$$\mathrm{IL} = -20\log|\mathit{T}|\mathrm{dB}$$

We use 20 in both the definitions as we primarily are dealing with power levels.



For the TL above with a length l of 6.330 m, find the average power delivered to the load as well as the line. The source voltage is $V_G=100\,\mathrm{V}$ operating at a frequency of 200 MHz. The source impedance Z_G is also 50 Ω . First, check whether the circuit is actually a transmission line. Also find the input impedance $Z_{in}=Z(-l)$ of the line.

First, we check the electrical length of the transmission line.

$$\beta = k = \frac{2\pi}{\lambda} = \frac{\omega}{v}$$

$$\implies \lambda = \frac{v}{f} = \frac{c}{\sqrt{\varepsilon_r}f} = \frac{3 \times 10^8}{\sqrt{2.5} \times 20 \times 10^6}$$

$$\lambda = 10.55 \,\mathrm{m} \implies l = \frac{6.33}{10.55} = 0.6\lambda$$

As $l>0.1\lambda$, the given circuit can be treated as a transmission line.

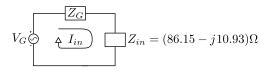
As Z_0 is real-valued, we treat the TL as lossless.

$$Z_{in} = Z(-l) = Z_0 \frac{Z_L + jZ_0 \tan \beta l}{Z_0 + jZ_L \tan \beta l}$$

$$= 50 \times \frac{(36 + j20) + j50 \tan (2\pi \times 0.6)}{50 + j(36 + j20) \tan (2\pi \times 0.6)}$$

$$Z_{in} = (86.15 - j10.93) \Omega$$

The equivalent circuit now looks like:



The current in the circuit is:

$$I_{in} = \frac{V_G}{Z_G + Z_{in}}$$

$$= \frac{100}{50 + 86.15 - j10.93} = (0.7300 + j0.06000) A$$

The average input power is:

$$P_{in} = 1/2 \operatorname{Re} \{V_{in}I_{in}^*\} = 1/2 \operatorname{Re} \{Z_{in}I_{in}I_{in}^*\} = 23.09 \,\mathrm{W}$$

We first calculate the current I(z) in the transmission line.

$$I(z) = \frac{V_{+}}{Z_{0}} \left(e^{-j\beta z} - \Gamma e^{j\beta z} \right)$$

At
$$z = -l$$
, $I(-l) = I_{in}$ and treating $\frac{V_{+}}{Z_{0}} = I_{+}$
$$I(-l) = I_{in} = I_{+} \left(e^{+j\beta l} - \Gamma_{in} e^{-j\beta l} \right)$$

$$I_{+} = I_{in} / \left(e^{+j\beta l} - \Gamma_{in} e^{-j\beta l} \right) = (-0.4700 + j0.8900) \, \text{A}$$

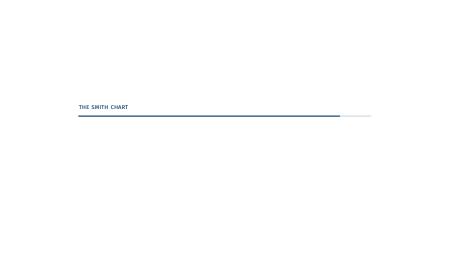
The current at the load is:

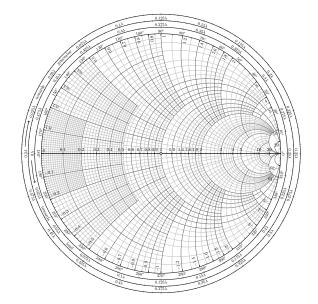
$$I_L = I_+(1 - \Gamma_L)$$

The power at the load is hence:

$$P_{load} = 1/2 \operatorname{Re}\{I_{L}I_{L}^{*}Z_{L}\}$$

= 23.09 W





WHY SMITH CHART

- · Developed by P Smith in 1939
- \cdot To this day, it is an integral part of microwave circuit design
- Provides a tool to visualise the transmission line phenomena such as impedance matching
- \cdot It is simply a polar plot of the reflection coefficient, Γ

NAVIGATING THE SMITH CHART

- · In polar coordinates, $\Gamma = |\Gamma| \mathrm{e}^{\mathrm{j}\theta}$
- · We plot the magnitude as a radius $(|\Gamma| \le 1)$ from the centre
- \cdot The angle θ ranges from -180° to 180°

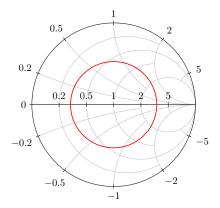


Figure 2: Γ plotted on the Smith chart