

Physics 115C HW 2

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Problem 1

(a)

We have that the first order corrections to the energy is given by

$$E_n^1 = \langle \psi_n | H' | \psi_n \rangle$$

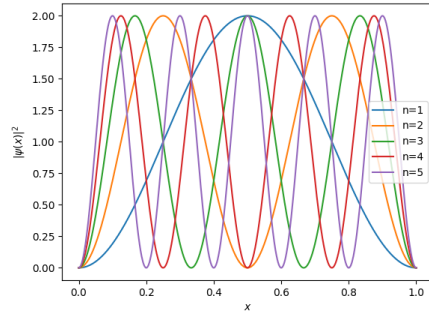
Therefore in our case we have that

$$E_n^1 = \alpha \langle \psi_n | \delta(x - \frac{L}{2}) | \psi_n \rangle$$

Therefore we get that

$$\begin{aligned} E_n^1 &= \alpha \int \psi_n^*(x) \delta(x - \frac{L}{2}) \psi_n(x) dx \\ &= \alpha |\psi(\frac{L}{2})|^2 \\ &= \alpha \frac{2}{L} \sin^2\left(\frac{n\pi}{2}\right) \\ &= \begin{cases} \alpha \frac{2}{L} & \text{if } n \text{ is odd} \\ 0 & \text{if } n \text{ is even} \end{cases} \end{aligned}$$

(b)



As we can see the states that do not exhibit a energy shift of the first order, are states where the probability density is zero at the location of the perubation. Thus the particle will not feel the effect of the perubation.

(c)

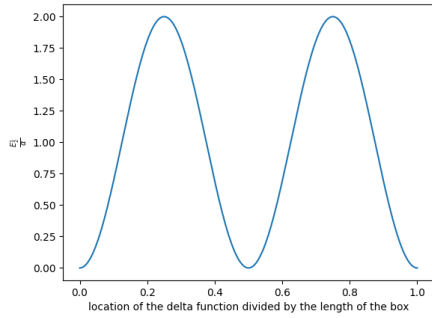
Likewise we can see that the states that do exhibit a energy shift are states where the probability density is not zero at the location of the perubation. Thus the particle will feel the effect of the perubation.

(d)

The delta function acts to effectively "sample" the wavefunction for the first order correction in the energy. Thus for the maximum first order energy corrections, we would want to place the delta function at the location where the probability density is the largest. Therefore for the $n = 2$ state this will be $\delta(x - \frac{L}{4}), \delta(x - \frac{3L}{4})$. For the $n = 4$ state this would be $\delta(x - \frac{L}{8}), \delta(x - \frac{3L}{8}), \delta(x - \frac{5L}{8}), \delta(x - \frac{7L}{8})$. And for the $n = 6$ state this would be $\delta(x - \frac{1}{12}L), \delta(x - \frac{3}{12}L), \delta(x - \frac{5}{12}L), \delta(x - \frac{7}{12}L), \delta(x - \frac{9}{12}L), \delta(x - \frac{11}{12}L)$.

(e)

Since it samples the wavefunction, as we move the delta function, the resulting first order correction will just be the amplitude of the wavefunction at the delta function:



Problem 2

(a)

$H^0 = \omega S_z$, therefore the eigenvalues are $+\omega$ and $-\omega$, and the eigenstates are $|\uparrow_z\rangle = [1, 0]^T$ and $|\downarrow_z\rangle = [0, 1]^T$.

(b)

We have that the characteristic equation is given by

$$x^2 - \omega^2 - (\lambda\Omega)^2 = 0$$

Therefore the eigenvalues are $\pm\sqrt{\omega^2 + (\lambda\Omega)^2}$.

(c)

We have that

$$H |\psi_1\rangle = \left(\omega \cos\left(\frac{\phi}{2}\right) + \lambda\Omega \sin\left(\frac{\phi}{2}\right) \right) |0\rangle + \left(\lambda\Omega \cos\left(\frac{\phi}{2}\right) - \omega \sin\left(\frac{\phi}{2}\right) \right) |1\rangle$$

In order $|\psi_1\rangle$ to be an eigenstate we must have:

$$H |\psi_1\rangle = \sqrt{\omega^2 + (\lambda\Omega)^2} |\psi_1\rangle$$

ie we need:

$$\begin{aligned} \omega \cos\left(\frac{\phi}{2}\right) + \lambda\Omega \sin\left(\frac{\phi}{2}\right) &= \sqrt{\omega^2 + (\lambda\Omega)^2} \cos\left(\frac{\phi}{2}\right) \\ \lambda\Omega \cos\left(\frac{\phi}{2}\right) - \omega \sin\left(\frac{\phi}{2}\right) &= \sqrt{\omega^2 + (\lambda\Omega)^2} \sin\left(\frac{\phi}{2}\right) \end{aligned}$$

We have that

$$\begin{aligned} \omega \cos\left(\frac{\phi}{2}\right) + \lambda\Omega \sin\left(\frac{\phi}{2}\right) &= \sqrt{\omega^2 + (\lambda\Omega)^2} \cos\left(\frac{\phi}{2}\right) \\ \cos(\phi) \cos\left(\frac{\phi}{2}\right) + \sin(\phi) \sin\left(\frac{\phi}{2}\right) &= \cos\left(\frac{\phi}{2}\right) \\ \sin^2(\phi) \sin^2\left(\frac{\phi}{2}\right) &= (1 - \cos(\phi))^2 \cos\left(\frac{\phi}{2}\right) \\ \sin^2(\phi) \frac{1 - \cos(\phi)}{2} &= (1 - \cos(\phi))^2 \frac{1 + \cos(\phi)}{2} \\ (1 - \cos^2(\phi))(1 - \cos(\phi)) &= (1 - \cos(\phi))^2(1 + \cos(\phi)) \\ (1 - \cos(\phi))^2(1 + \cos(\phi)) &= (1 - \cos(\phi))^2(1 + \cos(\phi)) \end{aligned}$$

Likewise we have that

$$\begin{aligned} \lambda\Omega \cos\left(\frac{\phi}{2}\right) - \omega \sin\left(\frac{\phi}{2}\right) &= \sqrt{\omega^2 + (\lambda\Omega)^2} \sin\left(\frac{\phi}{2}\right) \\ \sin(\phi) \cos\left(\frac{\phi}{2}\right) - \cos(\phi) \sin\left(\frac{\phi}{2}\right) &= \sin\left(\frac{\phi}{2}\right) \\ \sin^2(\phi) \cos^2\left(\frac{\phi}{2}\right) &= (1 + \cos(\phi))^2 \sin^2\left(\frac{\phi}{2}\right) \\ \sin^2(\phi)(1 + \cos(\phi)) &= (1 + \cos(\phi))^2(1 - \cos(\phi)) \\ (1 - \cos^2(\phi))(1 + \cos(\phi)) &= (1 + \cos(\phi))^2(1 - \cos(\phi)) \\ (1 + \cos(\phi))^2(1 - \cos(\phi)) &= (1 + \cos(\phi))^2(1 - \cos(\phi)) \end{aligned}$$

Therefore as we can see $|\psi_1\rangle$ is the eigenstate for the energy eigenvalue $\sqrt{\omega^2 + (\lambda\Omega)^2}$.

We can repeat a similar process for $|\psi_2\rangle$, we want to have

$$H|\psi_2\rangle = -\sqrt{\omega^2 + (\lambda\Omega)^2}|\psi_2\rangle$$

ie we need:

$$\begin{aligned} -\omega \sin\left(\frac{\phi}{2}\right) + \lambda\Omega \cos\left(\frac{\phi}{2}\right) &= \sqrt{\omega^2 + (\lambda\Omega)^2} \sin\left(\frac{\phi}{2}\right) \\ -\lambda\Omega \sin\left(\frac{\phi}{2}\right) - \omega \cos\left(\frac{\phi}{2}\right) &= -\sqrt{\omega^2 + (\lambda\Omega)^2} \cos\left(\frac{\phi}{2}\right) \end{aligned}$$

For the first equation we have that

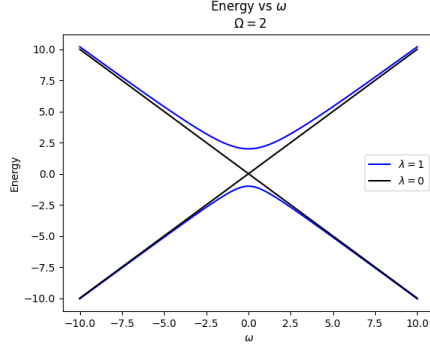
$$\begin{aligned} -\omega \sin\left(\frac{\phi}{2}\right) + \lambda\Omega \cos\left(\frac{\phi}{2}\right) &= \sqrt{\omega^2 + (\lambda\Omega)^2} \sin\left(\frac{\phi}{2}\right) \\ -\cos(\phi) \sin\left(\frac{\phi}{2}\right) + \sin(\phi) \cos\left(\frac{\phi}{2}\right) &= \sin\left(\frac{\phi}{2}\right) \end{aligned}$$

We already proved equality holds for this when we proved that the $|1\rangle$ components of $H|\psi_1\rangle$ and $\sqrt{\omega^2 + (\lambda\Omega)^2}|\psi_1\rangle$ are equal. For the second equation we have

$$\begin{aligned} -\lambda\Omega \sin\left(\frac{\phi}{2}\right) - \omega \cos\left(\frac{\phi}{2}\right) &= -\sqrt{\omega^2 + (\lambda\Omega)^2} \cos\left(\frac{\phi}{2}\right) \\ \cos(\phi) \cos\left(\frac{\phi}{2}\right) + \sin(\phi) \sin\left(\frac{\phi}{2}\right) &= \cos\left(\frac{\phi}{2}\right) \end{aligned}$$

We already proved equality holds for this when we proved that the $|0\rangle$ components of $H|\psi_1\rangle$ and $\sqrt{\omega^2 + (\lambda\Omega)^2}|\psi_1\rangle$ are equal. Thus $|\psi_2\rangle$ is the eigenstate for the energy eigenvalue $-\sqrt{\omega^2 + (\lambda\Omega)^2}$.

(d)



As we can see around $\omega = 0$ the energy eigenvalues become those for only the perturbation term, ie they become $\pm\lambda\Omega$. This is because the unperturbed hamiltonian becomes way smaller than the perturbation.

(e)

We have that:

$$\left. \frac{d\sqrt{\omega^2 + (\lambda\Omega)^2}}{d\lambda} \right|_{\lambda=0} = 0$$

so the first order correction is 0. taking the derivative again we get

$$\left. \frac{d^2\sqrt{\omega^2 + (\lambda\Omega)^2}}{d\lambda^2} \right|_{\lambda=0} = \frac{\Omega^2}{\omega}$$

Therefore we have that the second order correction is $\frac{\Omega^2}{\omega}$.

For the eigenstates we have that

$$\left. \frac{d \cos\left(\frac{\phi}{2}\right)}{d\lambda} \right|_{\lambda=0} = 0$$

$$\left. \frac{d^2 \sin\left(\frac{\phi}{2}\right)}{d\lambda^2} \right|_{\lambda=0} = \frac{\Omega^2}{4\omega^2}$$

and

$$\left. \frac{d \sin\left(\frac{\phi}{2}\right)}{d\lambda} \right|_{\lambda=0} = \frac{\Omega}{2\omega}$$

$$\left. \frac{d^2 \cos\left(\frac{\phi}{2}\right)}{d\lambda^2} \right|_{\lambda=0} = -\frac{\Omega^2}{4\omega^2}$$

Therefore we have that the first and second order corrections to the energy eigenstates are

$$|\psi_1\rangle \approx |0\rangle + \frac{\Omega\lambda}{2\omega} |1\rangle + \frac{\Omega^2\lambda^2}{4\omega^2} |0\rangle$$

and

$$|\psi_2\rangle \approx |1\rangle - \frac{\Omega\lambda}{2\omega} |0\rangle + \frac{\Omega^2\lambda^2}{4\omega^2} |1\rangle$$

(f)

The perturbation is $H' = \lambda\Omega\sigma_x$: Therefore we have the first order corrections are

$$E_1^1 = \langle 0 | H' | 0 \rangle = \lambda\Omega \langle 0 | 1 \rangle = 0$$

and

$$E_2^1 = \langle 1 | H' | 1 \rangle = \lambda\Omega \langle 1 | 0 \rangle = 0$$

(g)

This make sense when comparing to the results in (e), since both methods get that the first order correction to the energy eigenvalues is 0.

(h)

We have that first order correction for the energy eigenstates are:

$$|\psi_1^1\rangle = \frac{\langle 1 | H' | 0 \rangle}{E_1 - E_0} |1\rangle$$

$$|\psi_1^1\rangle = \frac{\lambda\Omega}{2\omega} |1\rangle$$

And for the other state:

$$|\psi_2^1\rangle = \frac{\langle 0|H'|1\rangle}{E_0 - E_1} |0\rangle$$

$$|\psi_2^1\rangle = -\frac{\lambda\Omega}{2\omega} |0\rangle$$

(i)

This result makes sense compared the the results in (e) since both methods get that the shift for the energy eigenstates $\frac{\lambda\Omega}{2\omega} |1\rangle$ for $|\psi_1\rangle$ and $-\frac{\lambda\Omega}{2\omega} |0\rangle$ for $|\psi_2\rangle$.

(j)

$$E_0^2 = \frac{|\langle 1|H'|0\rangle|^2}{E_0 - E_1}$$

$$E_0^2 = \frac{\lambda^2\Omega^2}{4\omega^2}$$

and

$$E_1^2 = \frac{|\langle 0|H'|1\rangle|^2}{E_1 - E_0}$$

$$E_1^2 = \frac{\lambda^2\Omega^2}{4\omega^2}$$

As we can see this is the same as the result in (e) for the second order correction to the energy eigenvalues.

Problem 3

(a)

We have that the eigenvalues are given by the characteritic equation

$$\begin{aligned} (x - E_a^0 - \lambda V_{aa})(x - E_b^0 - \lambda V_{bb}) - \lambda^2 V_{ab} V_{ba} &= 0 \\ x^2 - (E_a^0 + \lambda V_{aa} + E_b^0 + \lambda V_{bb})x + (E_a^0 + \lambda V_{aa})(E_b^0 + \lambda V_{bb}) - \lambda^2 V_{ab} V_{ba} &= 0 \end{aligned}$$

Solving we get:

$$x = \frac{1}{2}((E_a^0 + \lambda V_{aa} + E_b^0 + \lambda V_{bb}) \pm \sqrt{(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb})^2 - 4(\lambda^2 V_{ab} V_{ba})})$$

(b)

We have that

$$\begin{aligned} \left. \frac{d}{d\lambda} \sqrt{(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb})^2 - 4(\lambda^2 V_{ab} V_{ba})} \right|_{\lambda=0} &= \frac{1}{2(E_a^0 - E_b^0)} (2(E_a^0 - E_b^0)(V_{aa} - V_{bb})) \\ &= (V_{aa} - V_{bb}) \end{aligned}$$

and

$$\left. \frac{d^2}{d\lambda^2} \sqrt{(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb})^2 - 4(\lambda^2 V_{ab} V_{ba})} \right|_{\lambda=0} = \frac{4V_{ab} V_{ba}}{E_b - E_a}$$

Therefore we get the second order taylor series expansion of the eigenvalues as:

$$x = \frac{1}{2} \left((E_a^0 + \lambda V_{aa} + E_b^0 + \lambda V_{bb}) \pm \left(E_a^0 - E_b^0 + \lambda(V_{aa} - V_{bb}) + \frac{4\lambda^2 V_{ab} V_{ba}}{E_b - E_a} \right) \right)$$

(c)

We can rewrite

$$\sqrt{(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb})^2 - 4(\lambda^2 V_{ab} V_{ba})}$$

As

$$(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb}) \sqrt{1 - \frac{4(\lambda^2 V_{ab} V_{ba})}{(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb})^2}}$$

As $\lambda \rightarrow 0$ we have that becomes

$$(E_a^0 + \lambda V_{aa} - E_b^0 - \lambda V_{bb}) \sqrt{1 - \frac{4(\lambda^2 V_{ab} V_{ba})}{(E_a^0 - E_b^0)^2}}$$

We have that the taylor series of $\sqrt{1+x}$ around $|x| \leq \frac{1}{2}$ therefore we have that the taylor series for $\sqrt{1 - \frac{4(\lambda^2 V_{ab} V_{ba})}{(E_a^0 - E_b^0)^2}}$ converges only if

$$\left| \frac{4(V_{ab} V_{ba})}{(E_a^0 - E_b^0)^2} \right|^2 \leq \frac{1}{2}$$

$$\left| \frac{4(V_{ab})}{(E_a^0 - E_b^0)^2} \right| \leq \frac{1}{4}$$