

Physics 115C HW 5

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Problem 1

(a)

We have that

$$\Delta = \frac{\Omega}{2}$$

Therefore we must have that

$$\frac{\lambda^2}{\Delta^2} = 4 \frac{\lambda^2}{\Omega^2} = \frac{\gamma^2}{\hbar^2 \Omega^2}$$

Therefore we have that

$$\lambda = \frac{\gamma}{2\hbar}$$

Since

$$\Delta^2 = \lambda^2 + \omega^2$$

We have that

$$\omega^2 = \frac{\gamma^2}{4\hbar^2} - \frac{\Omega^2}{4}$$

Since $\Omega^2 = \delta\omega^2 + \frac{\gamma^2}{\hbar^2}$ we have:

$$\omega^2 = -\frac{\delta\omega^2}{4}$$

Therefore we have that

$$\omega = i\frac{\delta\omega}{2}$$

(b)

Therefore we have that

$$H_0 = \begin{pmatrix} 2\hbar i \frac{\delta\omega}{2} & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$H_1 = \begin{pmatrix} 0 & \frac{\gamma}{2} \\ \frac{\gamma}{2} & 0 \end{pmatrix}$$

Therefore we have that

$$H_0 + H_1 = \begin{pmatrix} 2\hbar i \frac{\delta\omega}{2} & \frac{\gamma}{2} \\ \frac{\gamma}{2} & 0 \end{pmatrix}$$

We know that the eigenvalues in terms of λ , ω and Δ are $\hbar(\omega \pm \Delta)$, therefore we have that the energy eigenvalues are

$$E_{\pm} = \hbar \left(i \frac{\delta\omega}{2} \pm \frac{\Omega}{2} \right)$$

And we have that the eigenvectors $v_+ = d \begin{pmatrix} \omega + \Delta \\ \lambda \end{pmatrix}$ and $v_- = d \begin{pmatrix} -\lambda \\ \omega + \Delta \end{pmatrix}$ where d is a normalization constant. Therefore we have that

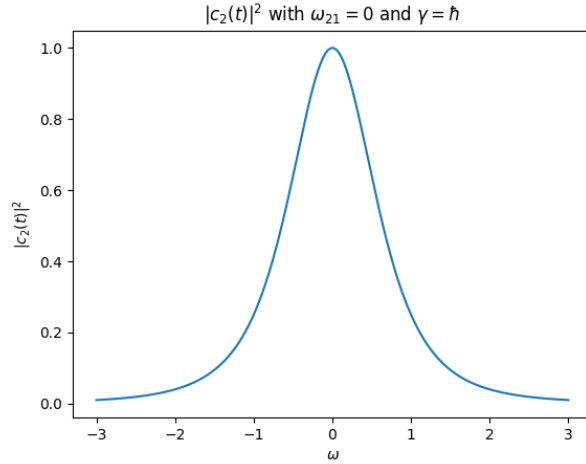
$$v_+ = d \begin{pmatrix} i \frac{\delta\omega}{2} + \frac{\Omega}{2} \\ \frac{\gamma}{2\hbar} \end{pmatrix}$$

And

$$v_- = d \begin{pmatrix} -\frac{\gamma}{2\hbar} \\ i \frac{\delta\omega}{2} + \frac{\Omega}{2} \end{pmatrix}$$

We have that $d^{-2} = \Omega$.

(c)



This is a lorentzian curve.

Problem 2

We have that the probability of a transition is given by

$$P_{if} = \left| \frac{i}{\hbar} \int_0^t e^{iE_0 t} e^{-\gamma t} \langle f | z | i \rangle dt \right|^2$$

We have that

$$\langle 211 | z | 100 \rangle = 0$$

$$\langle 21 - 1 | z | 210 \rangle = 0$$

From 6.3b. Also from applying what we have from 6.3e and changing L_z to L_y and y to z we get:

$$\langle 200 | z | 100 \rangle = 0$$

And thus we just need to solve for

$$\begin{aligned}
\langle 210 | z | 100 \rangle &= \int_0^\infty \int_0^{2\pi} \int_0^\pi r^3 \sin(\theta) \cos(\theta) \psi_{210}^*(r, \theta, \phi) \psi_{210}(r, \theta, \phi) d\theta d\phi dr \\
&= 2\pi \int_0^\infty Y_{210} Y_{100} r^3 \int_0^\pi \sin(\theta) \cos^2(\theta) d\theta dr \\
&= \frac{4\pi}{3} \int_0^\infty Y_{210} Y_{100} r^3 dr \\
&= \frac{4\pi}{3} \frac{1}{4\sqrt{2\pi}} \frac{1}{\sqrt{\pi}a_0^4} \int_0^\infty r^4 e^{-3r/2a_0} dr \\
&= \frac{2^{\frac{15}{2}} a_0}{3^5}
\end{aligned}$$

Thereforer we have that the probability of transistion is given by

$$P_{if} = \left| \frac{i}{\hbar} \frac{2^{\frac{15}{2}} a_0}{3^5} \int_0^t e|E_0| e^{-\gamma t} e^{\frac{it}{\hbar} E_f - E_i} dt \right|^2$$

Thus the limit as $t \rightarrow \infty$ is:

$$\begin{aligned}
P_{if} &= \left| \frac{i}{\hbar} \frac{2^{\frac{15}{2}} a_0}{3^5} \frac{e|E_0|}{-\gamma + \frac{i(E_f - E_i)}{\hbar}} \right|^2 \\
P_{if} &= \frac{2^{15} a_0^2}{3^{10} \hbar^2} \frac{e^2 |E_0|^2}{\gamma^2 + \frac{(E_f - E_i)^2}{\hbar^2}}
\end{aligned}$$

We have that $E_2 - E_1 = \frac{3e^2}{8a_0}$ thus we get:

$$P_{if} = \boxed{\frac{2^{15} a_0^2}{3^{10} \hbar^2} \frac{e^2 |E_0|^2}{\gamma^2 + \frac{9e^4}{64a_0^2 \hbar^2}}}$$

Problem 3

(a)

We have that

$$\begin{aligned}[L_z, x] &= [xp_y - yp_x, x] \\ &= [xp_y, x] - [yp_x, x] \\ &= x[p_y, x] + [x, p_y]x - [y, x]p_x - y[p_x, x] \\ &= -y[p_x, x] \\ &= yi\hbar\end{aligned}$$

Likewise

$$\begin{aligned}[L_z, y] &= [xp_y - yp_x, y] \\ &= [xp_y, y] - [yp_x, y] \\ &= x[p_y, y] + [x, p_y]y - [y, x]p_x - y[p_x, y] \\ &= x[p_y, y] - [y, x]p_x \\ &= -xi\hbar\end{aligned}$$

and

$$\begin{aligned}[L_z, z] &= [xp_y - yp_x, z] \\ &= [xp_y, z] - [yp_x, z] \\ &= x[p_y, z] + [x, p_y]z - [y, x]p_x - y[p_x, z] \\ &= x[p_y, z] - [y, x]p_x \\ &= 0\end{aligned}$$

(b)

We have that

$$\begin{aligned}\langle n'l'm' | [L_z, z] | nlm \rangle &= \langle n'l'm' | 0 | nlm \rangle \\ \langle n'l'm' | L_z z - z L_z | nlm \rangle &= 0 \\ \langle n'l'm' | (m' - m)z | nlm \rangle &= 0 \\ (m' - m) \langle n'l'm' | z | nlm \rangle &= 0\end{aligned}$$

Thus we can see that unless $m' = m$ the matrix element must be zero, thus if light is polarized in the z direction, it cannot change the m quantum number. Thus we cannot go from $|100\rangle \rightarrow |211\rangle$ with light polarized in the z direction.

(c)

We have that

$$\begin{aligned}\langle n'l'm' | [L_z, x] | nlm \rangle &= \langle n'l'm' | i\hbar y | nlm \rangle \\ \langle n'l'm' | L_z x - x L_z | nlm \rangle &= i\hbar \langle n'l'm' | y | nlm \rangle \\ \hbar \langle n'l'm' | (m' - m)x | nlm \rangle &= i\hbar \langle n'l'm' | y | nlm \rangle \\ (m' - m) \langle n'l'm' | x | nlm \rangle &= i \langle n'l'm' | y | nlm \rangle\end{aligned}$$

(d)

We have that:

$$\begin{aligned}[L_z, x \pm iy] &= [L_z, x] \pm i[L_z, y] \\ &= i\hbar y \pm \hbar x\end{aligned}$$

$$\begin{aligned}\langle n'l'm' | [L_z, x \pm iy] | nlm \rangle &= \langle n'l'm' | i\hbar y \pm \hbar x | nlm \rangle \\ (m' - m) \langle n'l'm' | x \pm iy | nlm \rangle &= \pm \langle n'l'm' | x \pm iy | nlm \rangle \\ (m' - m \mp 1) \langle n'l'm' | x \pm iy | nlm \rangle &= 0\end{aligned}$$

Thus we can see that the only possible nonzero matrix elements are when $m' = m \pm 1$. This reminds us of the raising and lower operators.

(e)

We have that

$$\begin{aligned}\langle n'l'm' | [L_z, y] | nlm \rangle &= -\langle n'l'm' | i\hbar x | nlm \rangle \\ (m' - m) \langle n'l'm' | y | nlm \rangle &= -i \langle n'l'm' | x | nlm \rangle \\ (m' - m)^2 \langle n'l'm' | x | nlm \rangle &= \langle n'l'm' | x | nlm \rangle \\ ((m' - m)^2 - 1) \langle n'l'm' | x | nlm \rangle &= 0\end{aligned}$$

Thus we can see that $\langle n'l'm' | x | nlm \rangle = \langle n'l'm' | y | nlm \rangle = 0$ unless $(m' - m)^2 = 1$. This means that for both the $\langle n'l'm' | x | nlm \rangle$ and $\langle n'l'm' | y | nlm \rangle$ matrices will have 0 values beyond the diagonal above and below the main diagonal.