BTP-1

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0.1 Introduction

introduction

Chapter 1

Mathematical Formulation

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The geometry considered is a thin film of Silicon in the domain $-L/2 \le X \le L/2$ and $0 \le Y \le H$. Consider a certain particle, initially located at the coordinate X. During deformation, this particle follows a path

$$\boldsymbol{x} = \boldsymbol{x}(\boldsymbol{X}, t) \tag{1.1}$$

Let u(X,t) be the displacement of a material particle located at X. Then

$$\boldsymbol{u}(\boldsymbol{X},t) = \boldsymbol{x}(\boldsymbol{X},t) - \boldsymbol{X} \tag{1.2}$$

The total deformation gradient and Green-Lagrange strain are denoted by \mathbf{F} and \mathbf{E} , respectively. Therefore,

$$\mathbf{F} = \frac{\partial \mathbf{x}}{\partial \mathbf{X}} = \nabla_{\mathbf{X}} \mathbf{u} + \mathbf{I} \tag{1.3}$$

$$\mathbf{E} = \frac{1}{2} (\mathbf{F}^\mathsf{T} \cdot \mathbf{F} - \mathbf{I}) \tag{1.4}$$

where, \mathbf{I} is the second-order isotropic tensor.

Let $\{\hat{e}_1, \hat{e}_2, \hat{e}_3\}$ be the orthonormal basis in the reference configuration. Denoting the corresponding components of X by X, Y and Z and that of u by u, v and w, and assuming plane strain deformation the components of Y are given by:

Both Inelastic and elastic deformation gradients are considered to be finite. Hence, a multiplicative decomposition of **F** into elastic and inelastic deformation is necessary. The body is first considered to reach an intermediate stress-free state and then it undergoes elastic deformation to reach the current configuration (Lee 1969).

$$\mathbf{F} = \mathbf{F}^{\text{el}} \cdot \mathbf{F}^{\text{inel}} \tag{1.6}$$

in indicial notation,
$$F_{ij} = F_{ik}^{\text{el}} F_{kj}^{inel}$$
 (1.7)

where \mathbf{F}^{el} and \mathbf{F}^{inel} are the deformation gradients due to elastic deformation and inelastic deformation respectively.

The inelastic deformation gradient tensor, \mathbf{F}^{inel} , has contribution from two sources - deformation due to concentration gradient, \mathbf{F}^{c} , and viscoplastic deformation, \mathbf{F}^{p} .

$$\mathbf{F}^{\text{inel}} = \mathbf{F}^{c} \cdot \mathbf{F}^{p} \tag{1.8}$$

1.1 Viscoplastic Deformation

A viscoplastic constitutive relation of the following form is considered.

$$\mathbf{D}^{\mathrm{P}} = \frac{\partial G(\sigma_{\mathrm{eff}})}{\partial \boldsymbol{\tau}} \tag{1.9}$$

Where \mathbf{D}^{P} is the rate dependent plastic deformation tensor, $G(\sigma_{\mathrm{eff}})$ is the flow potential, σ_{eff} is the von Mises stress and $\boldsymbol{\tau}$ is the deviatoric part of Cauchy stress tensor. citation

$$G(\sigma_{\text{eff}}) = \frac{\sigma_{\text{f}} \dot{d}_{0}}{\text{m} + 1} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)^{\text{m} + 1} \text{H} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)$$
(1.10)

$$\Rightarrow \mathbf{D}^{P} = \frac{3\tau \dot{d}_{0}}{2\sigma_{\text{eff}}} \left(\frac{\sigma_{\text{eff}}}{\sigma_{f}} - 1\right)^{m} H\left(\frac{\sigma_{\text{eff}}}{\sigma_{f}} - 1\right)$$
(1.11)

where, H is the unit step function, σ_f is the yield strength of Silicon, m is the stress exponent for plastic flow and \dot{d}_0 is the strain rate for plastic flow.

$$\mathbf{D}^{\mathrm{P}} = \mathbf{F}^{\mathrm{el}} \mathbf{F}^{\mathrm{c}} \dot{\mathbf{F}}^{\mathrm{p}} (\mathbf{F}^{\mathrm{p}})^{-1} (\mathbf{F}^{\mathrm{c}})^{-1} (\mathbf{F}^{\mathrm{el}})^{-1}$$
(1.12)

$$\Rightarrow \dot{\mathbf{F}}^{p} = (J)^{-1} \frac{3}{2} \frac{\mathbf{M_{0}^{el}} \mathbf{F}^{p}}{\sigma_{\text{eff}}} \dot{d}_{0} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)^{m} \mathbf{H} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)$$
(1.13)

where,
$$\mathbf{M_0^{el}} = J(\mathbf{F}^{el})^\mathsf{T} \boldsymbol{\tau} (\mathbf{F}^{el})^{-\mathsf{T}}$$
 (1.14)

$$J = \det(\mathbf{F}) \tag{1.15}$$

 $\mathbf{M_0^{el}}$ is the deviatoric part of Mandel stress (). The expression for Mandel stress is $\mathbf{M^{el}} = J(\mathbf{F^{el}})^\mathsf{T} \boldsymbol{\sigma}(F^{el})^\mathsf{-T}$. $\mathbf{F^p}$ is assumed to be of the following form:

$$[\mathbf{F}^{\mathbf{p}}] = \begin{bmatrix} \lambda_{11} & \lambda_{12} & 0\\ \lambda_{21} & \lambda_{22} & 0\\ 0 & 0 & \lambda_{33} \end{bmatrix}$$
 (1.16)

Since, $det(\mathbf{F}^p) = 1$

$$\lambda_{33} = 1/(\lambda_{11}\lambda_{22} - \lambda_{12}\lambda_{21}) \tag{1.17}$$

1.2 Deformation due to concentration gradient

The compound formed between Lithium and Silicon is of the form $\text{Li}_{\chi}\text{Si}$. Let the stoichiometric concentration and maximum concentration of Silicon atoms per atom of Lithium be denoted by χ_0 and χ_{max} . Defining a non-dimensional measure of the Li-ions concentration as $\tilde{c} = (\chi - \chi_0)/\chi_{\text{max}}$. since, χ_0 is the stoichiometric ratio it signifies the stress free state of the particle and hence, \tilde{c} is a measure of the deviation of the particle from undeformed state.

$$\mathbf{F}^{c} = (J^{c})^{1/3}\mathbf{I} \tag{1.18}$$

where
$$J^c = 1 + 3\eta \chi_{\text{max}} \tilde{c}$$
 (1.19)

 η is a material parameter giving rate of change in volume w.r.t. \tilde{c} . It may be noted that as \tilde{c} approaches 1, $\det(\mathbf{F}^c)$ approaches 4. Therefore the body undergoes a volumetric change of about 300% due to diffusion of Li-ions, justifying the use of large deformation analysis.

1.3 Momentum Conservation

From equations 1.6 and 1.8, \mathbf{F}^{el} can be expressed as,

$$\mathbf{F}^{\text{el}} = \mathbf{F} \cdot (\mathbf{F}^{\text{p}} \cdot \mathbf{F}^{\text{c}})^{-1} \tag{1.20}$$

The elastic Green-Lagrange strain, $\mathbf{E}^{\mathrm{el}} = \frac{1}{2} \left[(\mathbf{F}^{\mathrm{el}})^\mathsf{T} \cdot \mathbf{F}^{\mathrm{el}} - \mathbf{I} \right]$

The strain energy per unit volume in the reference configuration, $W(\mathbf{F}, \tilde{c})$, is expressed as $W(\mathbf{F}, \tilde{c}) = J^{\text{inel}}\bar{w}(\mathbf{F}, \tilde{c})$, where $\bar{w}(\mathbf{F}, \tilde{c})$ is the strain energy per unit volume in the

intermediate configuration and $J^{\text{inel}} = \det(\mathbf{F}^{\text{inel}}) = J^c$.

$$W(\mathbf{F}, \tilde{c}) = \frac{J^c}{2} \frac{E(\tilde{c})}{1+\nu} \left(\frac{\nu}{1-2\nu} (\operatorname{tr} \mathbf{E}^{el})^2 + \operatorname{tr} (\mathbf{E}^{el} \cdot \mathbf{E}^{el}) \right). \tag{1.21}$$

The elastic modulus of Silicon is concentration dependent with $E(\tilde{c}) = E_{\rm si}(1 + \eta_{\rm E}\chi_{\rm max}\tilde{c})$. The elastic second Piola-Kirchhoff stress is denoted by ${\bf S}^{\rm el}$. Differentiating W w.r.t ${\bf E}^{\rm el}$ gives,

$$\mathbf{S}^{\text{el}} = J^{c}[2\mu_{\text{si}}(\tilde{c})\mathbf{E}^{\text{el}} + \lambda_{\text{si}}(\tilde{c})\text{tr}(\mathbf{E}^{\text{el}})\mathbf{I}]$$
(1.22)

Let **P** and **S** denote the first and second Piola-Kirchhoff stress, respectively. Thus

$$\mathbf{S} = (\mathbf{F}^{c})^{-1} \cdot (\mathbf{F}^{p})^{-1} \cdot \mathbf{S}^{el} \cdot (\mathbf{F}^{p})^{-\mathsf{T}} \cdot (\mathbf{F}^{c})^{-\mathsf{T}}$$
(1.23)

$$\mathbf{P} = \mathbf{F} \cdot \mathbf{S} \tag{1.24}$$

The Cauchy stress tensor, $\boldsymbol{\sigma}$ is given by $\boldsymbol{\sigma} = (J)^{-1} \mathbf{P} \cdot \mathbf{F}^{\mathsf{T}}$. And, the deviatoric part of Cauchy is $\boldsymbol{\tau} = \boldsymbol{\sigma} - (1/3) \mathrm{tr}(\boldsymbol{\sigma}) \mathbf{I}$. The von Mises stress is $\sigma_{\text{eff}} = \sqrt{\frac{3}{2} (\tau_{11}^2 + \tau_{22}^2 + \tau_{33}^2 + 2\tau_{12}^2)}$.

Conservation of momentum leads to

$$\nabla_{\mathbf{X}} \cdot \mathbf{P} = 0. \tag{1.25}$$

1.4 Mass Conservation

Assuming flux to be negligible in the z direction, the conservation of mass is given by

$$\frac{\partial c}{\partial t} = -\nabla_{\mathbf{X}} \cdot \mathbf{j} = -\left(\frac{\partial j_X}{\partial X} + \frac{\partial j_Y}{\partial Y}\right). \tag{1.26}$$

Where, \boldsymbol{j} is the flux vector and c is a dimensional measure of Li-ions concentration, defined as $c = \chi_{\text{max}}/V_{\text{m}}^{\text{B}} \tilde{c}$.

1.5 Non-Dimensionalization

$$\tilde{j}_X, \tilde{j}_Y, \tilde{J}_0, \tilde{\boldsymbol{j}} = \frac{HV_{\rm m}^{\rm B}}{(\chi_{\rm max} D_0)} (j_X, j_y, J_0, \boldsymbol{j})$$

$$\tag{1.27}$$

$$\tilde{X}, \tilde{Y}, \tilde{u}, \tilde{v} = \frac{1}{H}(X, Y, u, v) \tag{1.28}$$

$$\tilde{t} = D_0 t / H^2 \tag{1.29}$$

$$\tilde{\mu}_{\rm si}, \tilde{\lambda}_{\rm si}, \tilde{E}_{\rm si} = \frac{1}{E_0} (\mu_{\rm si}, \lambda_{\rm si}, E_{\rm si}), \text{where } E_0 = \frac{R_g T}{V_{\rm m}^{\rm B}}$$
 (1.30)

$$\tilde{\mu}_0, \tilde{\mu}_1, \tilde{\mu}_2, \tilde{\mu}_3 = \frac{1}{R_q T} (\mu_0, \mu_1, \mu_2, \mu_3)$$
 (1.31)

$$\tilde{D} = \frac{D}{D_0} \tag{1.32}$$

$$\tilde{\dot{d}}_0 = \frac{\dot{d}_0 H^2}{D_0} \tag{1.33}$$

$$\tilde{\mathbf{S}}^{\text{el}}, \tilde{\mathbf{S}}, \tilde{\mathbf{P}}, \tilde{\boldsymbol{\sigma}}, \tilde{\boldsymbol{\tau}}, \tilde{\mathbf{M}}_{0}^{\text{el}}, \tilde{\sigma}_{\text{eff}}, \tilde{\sigma}_{\text{f}} = \frac{1}{E_{0}} (\mathbf{S}^{\text{el}}, \mathbf{S}, \mathbf{P}, \boldsymbol{\sigma}, \boldsymbol{\tau}, \mathbf{M}_{0}^{\text{el}}, \sigma_{\text{eff}}, \sigma_{\text{f}})$$
(1.34)

1.6 Definition of the state of charge

state of charge is a measure of the degree of lithiation. It can be expressed as an average concentration over the domain as follows:

$$\operatorname{soc} = \frac{\int_{-L/2}^{L/2} \int_{0}^{H} \tilde{c} dy dx}{LH}$$
 (1.35)

$$=H^{2}\frac{\int_{-L/2H}^{L/2H}\int_{0}^{1}\tilde{c}(\tilde{x},\tilde{y})\mathrm{d}\tilde{y}\mathrm{d}\tilde{x}}{LH}$$
(1.36)

$$= H \frac{\int_{-L/2H}^{L/2H} \int_{0}^{1} \tilde{c}(\tilde{x}, \tilde{y}) d\tilde{y} d\tilde{x}}{L}$$
 (1.37)

1.7 Non-Dimensional Equations in Component Form

$$[\mathbf{F}] = \begin{bmatrix} 1 + \frac{\partial \tilde{u}}{\partial \tilde{X}} & \frac{\partial \tilde{u}}{\partial \tilde{Y}} & 0\\ \frac{\partial \tilde{v}}{\partial \tilde{X}} & 1 + \frac{\partial \tilde{v}}{\partial \tilde{Y}} & 0\\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} F_{11} & F_{12} & 0\\ F_{21} & F_{22} & 0\\ 0 & 0 & F_{33} \end{bmatrix}$$
 (1.38)

$$F_{11}^{el} = \frac{F_{11} \,\lambda_{22} - F_{12} \,\lambda_{21}}{J^{c^{1/3}} \,(\lambda_{11} \,\lambda_{22} - \lambda_{12} \,\lambda_{21})} \tag{1.39}$$

$$F_{21}^{el} = \frac{F_{21} \lambda_{22} - F_{22} \lambda_{21}}{J^{c^{1/3}} \left(\lambda_{11} \lambda_{22} - \lambda_{12} \lambda_{21}\right)}$$
(1.40)

$$F_{12}^{el} = -\frac{F_{11}\,\lambda_{12} - F_{12}\,\lambda_{11}}{J^{c^{1/3}}\,(\lambda_{11}\,\lambda_{22} - \lambda_{12}\,\lambda_{21})} \tag{1.41}$$

$$F_{22}^{el} = -\frac{F_{21}\,\lambda_{12} - F_{22}\,\lambda_{11}}{J^{c^{1/3}}\,(\lambda_{11}\,\lambda_{22} - \lambda_{12}\,\lambda_{21})}\tag{1.42}$$

$$F_{33}^{el} = \frac{1}{J^{c^{1/3}} \lambda_{23}} \tag{1.43}$$

$$E_{11}^{el} = \frac{F_{11}^{el}^2}{2} + \frac{F_{21}^{el}^2}{2} - \frac{1}{2}$$
 (1.44)

$$E_{21}^{el} = \frac{F_{11}^{el} F_{12}^{el}}{2} + \frac{F_{21}^{el} F_{22}^{el}}{2}$$
 (1.45)

$$E_{12}^{el} = \frac{F_{11}^{el} F_{12}^{el}}{2} + \frac{F_{21}^{el} F_{22}^{el}}{2}$$
 (1.46)

$$E_{22}^{el} = \frac{F_{12}^{el}^2}{2} + \frac{F_{22}^{el}^2}{2} - \frac{1}{2}$$
 (1.47)

$$E_{33}^{el} = \frac{F_{33}^{el^2}}{2} - \frac{1}{2} \tag{1.48}$$

$$\tilde{S}_{11}^{\text{el}} = J^c \left(2 E_{11}^{\text{el}} \, \tilde{\mu}_{\text{si}} + \tilde{\lambda}_{\text{si}} \, \left(E_{11}^{\text{el}} + E_{22}^{\text{el}} + E_{33}^{\text{el}} \right) \right) \tag{1.49}$$

$$\tilde{S}_{21}^{\text{el}} = 2 E_{21}^{\text{el}} J^c \,\tilde{\mu}_{\text{si}} \tag{1.50}$$

$$\tilde{S}_{12}^{el} = 2 E_{12}^{el} J^c \,\tilde{\mu}_{si} \tag{1.51}$$

$$\tilde{S}_{22}^{el} = J^c \left(2 E_{22}^{el} \,\tilde{\mu}_{si} + \tilde{\lambda}_{si} \, \left(E_{11}^{el} + E_{22}^{el} + E_{33}^{el} \right) \right) \tag{1.52}$$

$$\tilde{S}_{33}^{el} = J^c \left(2 E_{33}^{el} \, \tilde{\mu}_{si} + \tilde{\lambda}_{si} \, \left(E_{11}^{el} + E_{22}^{el} + E_{33}^{el} \right) \right) \tag{1.53}$$

$$\tilde{S}_{11} = \frac{\tilde{S}_{11}^{\text{el}} \lambda_{22}^2 + \tilde{S}_{22}^{\text{el}} \lambda_{12}^2 - \tilde{S}_{12}^{\text{el}} \lambda_{12} \lambda_{22} - \tilde{S}_{21}^{\text{el}} \lambda_{12} \lambda_{22}}{J^{c^{2/3}} (\lambda_{11} \lambda_{22} - \lambda_{12} \lambda_{21})^2}$$
(1.54)

$$\tilde{S}_{21} = \frac{\tilde{S}_{12}^{\text{el}} \lambda_{12} \lambda_{21} - \tilde{S}_{22}^{\text{el}} \lambda_{11} \lambda_{12} - \tilde{S}_{11}^{\text{el}} \lambda_{21} \lambda_{22} + \tilde{S}_{21}^{\text{el}} \lambda_{11} \lambda_{22}}{J^{c^{2/3}} (\lambda_{11} \lambda_{22} - \lambda_{12} \lambda_{21})^{2}}$$
(1.55)

$$\tilde{S}_{12} = \frac{\tilde{S}_{12}^{\text{el}} \lambda_{11} \lambda_{22} - \tilde{S}_{22}^{\text{el}} \lambda_{11} \lambda_{12} - \tilde{S}_{11}^{\text{el}} \lambda_{21} \lambda_{22} + \tilde{S}_{21}^{\text{el}} \lambda_{12} \lambda_{21}}{J^{c^{2/3}} (\lambda_{11} \lambda_{22} - \lambda_{12} \lambda_{21})^{2}}$$
(1.56)

$$\tilde{S}_{22} = \frac{\tilde{S}_{11}^{\text{el}} \lambda_{21}^2 + \tilde{S}_{22}^{\text{el}} \lambda_{11}^2 - \tilde{S}_{12}^{\text{el}} \lambda_{11} \lambda_{21} - \tilde{S}_{21}^{\text{el}} \lambda_{11} \lambda_{21}}{J^{c^{2/3}} (\lambda_{11} \lambda_{22} - \lambda_{12} \lambda_{21})^2}$$
(1.57)

$$\tilde{S}_{33} = \frac{\tilde{S}_{33}^{\text{el}}}{J^{c^2/3} \lambda_{33}^2} \tag{1.58}$$

$$\tilde{P}_{11} = F_{11}\,\tilde{S}_{11} + F_{12}\,\tilde{S}_{21} \tag{1.59}$$

$$\tilde{P}_{21} = F_{21}\,\tilde{S}_{11} + F_{22}\,\tilde{S}_{21} \tag{1.60}$$

$$\tilde{P}_{12} = F_{11}\,\tilde{S}_{12} + F_{12}\,\tilde{S}_{22} \tag{1.61}$$

$$\tilde{P}_{22} = F_{21}\,\tilde{S}_{12} + F_{22}\,\tilde{S}_{22} \tag{1.62}$$

$$\tilde{P}_{33} = \tilde{S}_{33} \tag{1.63}$$

$$\tilde{\sigma}_{11} = \frac{F_{11}\,\tilde{P}_{11} + F_{12}\,\tilde{P}_{12}}{J} \tag{1.64}$$

$$\tilde{\sigma}_{21} = \frac{F_{11}\,\tilde{P}_{21} + F_{12}\,\tilde{P}_{22}}{J} \tag{1.65}$$

$$\tilde{\sigma}_{12} = \frac{F_{21}\,\tilde{P}_{11} + F_{22}\,\tilde{P}_{12}}{J} \tag{1.66}$$

$$\tilde{\sigma}_{22} = \frac{F_{21}\,\tilde{P}_{21} + F_{22}\,\tilde{P}_{22}}{J} \tag{1.67}$$

$$\tilde{\sigma}_{33} = \frac{\dot{P}_{33}}{J} \tag{1.68}$$

$$\tilde{\tau}_{11} = \frac{2\,\tilde{\sigma}_{11}}{3} - \frac{\tilde{\sigma}_{22}}{3} - \frac{\tilde{\sigma}_{33}}{3} \tag{1.69}$$

$$\tilde{\tau}_{21} = \tilde{\sigma}_{21} \tag{1.70}$$

$$\tilde{\tau}_{12} = \tilde{\sigma}_{12} \tag{1.71}$$

$$\tilde{\tau}_{22} = \frac{2\,\tilde{\sigma}_{22}}{3} - \frac{\tilde{\sigma}_{11}}{3} - \frac{\tilde{\sigma}_{33}}{3} \tag{1.72}$$

$$\tilde{\tau}_{33} = \frac{2\,\tilde{\sigma}_{33}}{3} - \frac{\tilde{\sigma}_{22}}{3} - \frac{\tilde{\sigma}_{11}}{3} \tag{1.73}$$

$$\tilde{\sigma}_{\text{eff}} = \sqrt{\frac{3}{2}(\tilde{\tau}_{11}^2 + \tilde{\tau}_{22}^2 + \tilde{\tau}_{33}^2 + 2\tilde{\tau}_{12}^2)}$$
(1.74)

$$\tilde{M}_{11}^{el} = -\frac{J\left(F_{11}^{el} F_{12}^{el} \tilde{\tau}_{12} - F_{11}^{el} F_{22}^{el} \tilde{\tau}_{11} + F_{12}^{el} F_{21}^{el} \tilde{\tau}_{22} - F_{21}^{el} F_{22}^{el} \tilde{\tau}_{21}\right)}{F_{11}^{el} F_{22}^{el} - F_{12}^{el} F_{21}^{el}}$$
(1.75)

$$\tilde{M}_{21}^{el} = -\frac{J\left(F_{12}^{el^2}\tilde{\tau}_{12} - F_{22}^{el^2}\tilde{\tau}_{21} - F_{12}^{el}F_{22}^{el}\tilde{\tau}_{11} + F_{12}^{el}F_{22}^{el}\tilde{\tau}_{22}\right)}{F_{11}^{el}F_{22}^{el} - F_{12}^{el}F_{21}^{el}}$$
(1.76)

$$\tilde{M}_{12}^{el} = \frac{J\left(F_{11}^{el^2}\tilde{\tau}_{12} - F_{21}^{el^2}\tilde{\tau}_{21} - F_{11}^{el}F_{21}^{el}\tilde{\tau}_{11} + F_{11}^{el}F_{21}^{el}\tilde{\tau}_{22}\right)}{F_{11}^{el}F_{22}^{el} - F_{12}^{el}F_{21}^{el}}$$

$$\tilde{M}_{22}^{el} = \frac{J\left(F_{11}^{el}F_{12}^{el}\tilde{\tau}_{12} - F_{12}^{el}F_{21}^{el}\tilde{\tau}_{11} + F_{11}^{el}F_{22}^{el}\tilde{\tau}_{22} - F_{21}^{el}F_{22}^{el}\tilde{\tau}_{21}\right)}{F_{11}^{el}F_{22}^{el} - F_{12}^{el}F_{21}^{el}}$$

$$(1.77)$$

$$\tilde{M}_{22}^{el} = \frac{J\left(F_{11}^{el} F_{12}^{el} \tilde{\tau}_{12} - F_{12}^{el} F_{21}^{el} \tilde{\tau}_{11} + F_{11}^{el} F_{22}^{el} \tilde{\tau}_{22} - F_{21}^{el} F_{22}^{el} \tilde{\tau}_{21}\right)}{F_{11}^{el} F_{22}^{el} - F_{12}^{el} F_{21}^{el}}$$
(1.78)

$$\tilde{M}_{33}^{el} = J\,\tilde{\tau}_{33} \tag{1.79}$$

Viscoplastic rates: These equations need editing if Fpdot.mat.tex is changed

$$\dot{\tilde{F}}_{11}^{p} = \dot{\tilde{d}}_{0} \left(\frac{\tilde{\sigma}_{\text{eff}}}{\tilde{\sigma}_{f}} - 1 \right)^{m} \left(\frac{3 \, \tilde{M}_{11}^{\text{el}} \, \lambda_{11}}{2 \, J \, \tilde{\sigma}_{\text{eff}}} + \frac{3 \, \tilde{M}_{12}^{\text{el}} \, \lambda_{21}}{2 \, J \, \tilde{\sigma}_{\text{eff}}} \right) \mathcal{H} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)$$

$$(1.80)$$

$$\dot{\tilde{F}}_{21}^{p} = \dot{\tilde{d}}_{0} \left(\frac{\tilde{\sigma}_{\text{eff}}}{\tilde{\sigma}_{f}} - 1 \right)^{m} \left(\frac{3 \,\tilde{M}_{21}^{\text{el}} \,\lambda_{11}}{2 \,J \,\tilde{\sigma}_{\text{eff}}} + \frac{3 \,\tilde{M}_{22}^{\text{el}} \,\lambda_{21}}{2 \,J \,\tilde{\sigma}_{\text{eff}}} \right) \mathcal{H} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)$$

$$(1.81)$$

$$\dot{\tilde{F}}_{12}^{p} = \dot{\tilde{d}}_{0} \left(\frac{\tilde{\sigma}_{\text{eff}}}{\tilde{\sigma}_{f}} - 1 \right)^{m} \left(\frac{3 \,\tilde{M}_{11}^{\text{el}} \,\lambda_{12}}{2 \,J \,\tilde{\sigma}_{\text{eff}}} + \frac{3 \,\tilde{M}_{12}^{\text{el}} \,\lambda_{22}}{2 \,J \,\tilde{\sigma}_{\text{eff}}} \right) \mathcal{H} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)$$

$$(1.82)$$

$$\dot{\tilde{F}}_{22}^{p} = \dot{\tilde{d}}_{0} \left(\frac{\tilde{\sigma}_{\text{eff}}}{\tilde{\sigma}_{f}} - 1 \right)^{m} \left(\frac{3 \, \tilde{M}_{21}^{\text{el}} \, \lambda_{12}}{2 \, J \, \tilde{\sigma}_{\text{eff}}} + \frac{3 \, \tilde{M}_{22}^{\text{el}} \, \lambda_{22}}{2 \, J \, \tilde{\sigma}_{\text{eff}}} \right) \mathcal{H} \left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1 \right)$$
(1.83)

$$\dot{\tilde{F}}_{33}^{p} = \frac{3 \, \tilde{M}_{33}^{el} \, \dot{\tilde{d}}_{0} \, \lambda_{33} \, \left(\frac{\tilde{\sigma}_{eff}}{\tilde{\sigma}_{f}} - 1\right)^{m}}{2 \, J \, \tilde{\sigma}_{eff}} H \left(\frac{\sigma_{eff}}{\sigma_{f}} - 1\right)$$

$$(1.84)$$

Viscoplastic rate equations in non-dimensional form:

$$\frac{d\lambda_{11}}{dt} = \frac{1}{t_{\text{ref}}} \frac{d\lambda_{11}}{d\tilde{t}} = \frac{3\dot{d}_0}{2J\tilde{\sigma}_{\text{eff}}} (\tilde{M}_{11}^{\text{el}} \lambda_{11} + \tilde{M}_{12}^{\text{el}} \lambda_{21}) \left(\frac{\tilde{\sigma}_{\text{eff}}}{\tilde{\sigma}_{\text{f}}} - 1\right)^{\text{m}} H\left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1\right)$$
(1.85)

$$\frac{d\lambda_{11}}{d\tilde{t}} = \frac{3\dot{\tilde{d}}_0}{2J\tilde{\sigma}_{\text{eff}}} (\tilde{M}_{11}^{\text{el}}\lambda_{11} + \tilde{M}_{12}^{\text{el}}\lambda_{21}) \left(\frac{\tilde{\sigma}_{\text{eff}}}{\tilde{\sigma}_{\text{f}}} - 1\right)^{\text{m}} H\left(\frac{\sigma_{\text{eff}}}{\sigma_{\text{f}}} - 1\right)$$
(1.86)

where,
$$\dot{\tilde{d}}_0 = \dot{d}_0 t_{\text{ref}}$$
 (1.87)

Momentum Conservation equations in non-dimensional form:

$$\frac{\partial \tilde{P}_{11}}{\partial \tilde{X}} + \frac{\partial \tilde{P}_{12}}{\partial \tilde{Y}} = 0 \tag{1.88}$$

and,
$$\frac{\partial \tilde{P}_{21}}{\partial \tilde{X}} + \frac{\partial \tilde{P}_{22}}{\partial \tilde{Y}} = 0.$$
 (1.89)

Mass conservation Equation in non-dimensional form:

$$\frac{\partial c}{\partial t} = -\nabla_{\mathbf{X}} \cdot \mathbf{j} = -\left(\frac{\partial j_x}{\partial X} + \frac{\partial j_y}{\partial Y}\right) \tag{1.90}$$

$$\frac{\chi_{\text{max}}}{V_{\text{m}}^{\text{B}}} \frac{D_0}{H^2} \frac{\partial \tilde{c}}{\partial \tilde{t}} = -\left(\frac{1}{H} \frac{\partial \tilde{j}_x}{\partial \tilde{X}} + \frac{1}{H} \frac{\partial \tilde{j}_y}{\partial \tilde{Y}}\right) \frac{\chi_{\text{max}}}{V_{\text{m}}^{\text{B}}} \frac{D_0}{H}$$
(1.91)

$$\Longrightarrow \frac{\partial \tilde{c}}{\partial \tilde{t}} = -\left(\frac{\partial \tilde{j}_x}{\partial \tilde{X}} + \frac{\partial \tilde{j}_y}{\partial \tilde{Y}}\right) \tag{1.92}$$

For one way coupling:

$$\boldsymbol{j} = -D_0 \boldsymbol{\nabla}_{\boldsymbol{X}} c \tag{1.93}$$

$$\tilde{j}_x = j_x H / (\frac{\chi_{\text{max}}}{V_{\text{m}}^{\text{B}}} D_0) \tag{1.94}$$

$$= -D_0 \frac{\partial c}{\partial X} H / (\frac{\chi_{\text{max}}}{V_{\text{m}}^{\text{B}}} D_0)$$
 (1.95)

$$= -\frac{\partial \tilde{c}}{\partial \tilde{X}} \tag{1.96}$$

$$\tilde{j}_y = j_y H / (\frac{\chi_{\text{max}}}{V_{\text{m}}^{\text{B}}} D_0) \tag{1.97}$$

$$= -\frac{\partial \tilde{c}}{\partial \tilde{Y}} \tag{1.98}$$

For two way coupling:

$$\boldsymbol{j} = -\frac{1}{R_o T} \frac{D \chi_{\text{max}} \tilde{c}}{V_m^b} (\mathbf{F})^{-1} (\mathbf{F})^{-\mathsf{T}} \boldsymbol{\nabla}_{\boldsymbol{X}} \mu$$
 (1.99)

$$\tilde{\boldsymbol{j}} = \frac{\boldsymbol{j}H}{\frac{\chi_{\text{max}}}{V_{\text{D}}^{\text{B}}}} D_0 = -\frac{1}{R_g T} \tilde{D}\tilde{c}(\mathbf{F})^{-1} (\mathbf{F})^{-\mathsf{T}} \nabla_{\boldsymbol{X}} \mu$$
(1.100)

$$D = D_0 \exp\left(\frac{\alpha S_h}{E_0}\right) = D_0 \exp\left(\alpha \tilde{S}_h\right) = D_0 \exp\left(\alpha \frac{\tilde{S}_{11} + \tilde{S}_{33}}{2}\right)$$
(1.101)

$$\mu = \mu_0 + \mu_s; \tilde{\mu} = \frac{\mu}{R_a T} \tag{1.102}$$

$$\mu_0 = R_g T \log(\gamma \tilde{c}); \tilde{\mu_0} = \log(\gamma \tilde{c}) \tag{1.103}$$

$$\gamma = \frac{1}{1 - \tilde{c}} \exp\left(\frac{1}{R_a T} \left[2(A_0 - 2B_0)\tilde{c} - 3(A_0 - B_0)(\tilde{c}^2)\right]\right)$$
(1.104)

$$\mu_{s} = \frac{V_{m}^{b}}{\chi_{\text{max}}} \left[-\frac{1}{3} \frac{\partial J^{c}}{\partial \tilde{c}} \tilde{F}_{im}^{\text{el}} \tilde{F}_{in}^{\text{el}} C_{mnkl} \tilde{E}_{kl}^{\text{el}} + \frac{1}{2} \left(J^{c} \frac{\partial C_{ijkl}}{\partial \tilde{c}} + \frac{\partial J^{c}}{\partial \tilde{c}} C_{ijkl} \right) \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.105)

$$\tilde{\mu}_s = \frac{1}{R_g T} \mu_s \tag{1.106}$$

$$= \frac{V_m^b}{R_g T \chi_{\text{max}}} \left[-\frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} \tilde{F}_{im}^{\text{el}} \tilde{F}_{in}^{\text{el}} C_{mnkl} \tilde{E}_{kl}^{\text{el}} + \frac{1}{2} \left(J^c \frac{\partial C_{ijkl}}{\partial \tilde{c}} + \frac{\partial J^c}{\partial \tilde{c}} C_{ijkl} \right) \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.107)

$$= \frac{1}{\chi_{\text{max}}} \left[-\frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} \tilde{F}_{im}^{\text{el}} \tilde{F}_{in}^{\text{el}} \tilde{C}_{mnkl} \tilde{E}_{kl}^{\text{el}} + \frac{1}{2} \left(J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} + \frac{\partial J^c}{\partial \tilde{c}} \tilde{C}_{ijkl} \right) \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.108)

$$\tilde{C}_{ijkl}\tilde{E}_{kl}^{\text{el}} = \tilde{P}_{ij}^{\text{el}} = \tilde{S}_{ij}^{\text{el}}/J^c \Longrightarrow \tilde{C}_{ijkl} = \tilde{\lambda}_{si}(\tilde{c})\delta_{ij}\delta_{kl} + \tilde{\mu}_{si}(\tilde{c})(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})$$
(1.109)

$$\tilde{\mu}_s = \frac{1}{\chi_{\text{max}}} \left[-\frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} \tilde{F}_{im}^{\text{el}} \tilde{F}_{in}^{\text{el}} \tilde{F}_{mn}^{\text{el}} + \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} + \frac{1}{2} \frac{\partial J^c}{\partial \tilde{c}} \tilde{C}_{ijkl} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.110)

$$= \frac{1}{\chi_{\text{max}}} \left[-\frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} \tilde{F}_{im}^{\text{el}} \tilde{F}_{in}^{\text{el}} \tilde{P}_{mn}^{\text{el}} + \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} + \frac{1}{2} \frac{\partial J^c}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{P}_{kl}^{\text{el}} \right]$$
(1.111)

$$= \frac{1}{\chi_{\text{max}}} \left[\frac{\partial J^c}{\partial \tilde{c}} \tilde{P}_{mn}^{\text{el}} \left(-\frac{1}{3} \tilde{F}_{im}^{\text{el}} \tilde{F}_{in}^{\text{el}} + \frac{1}{2} \tilde{E}_{mn}^{\text{el}} \right) + \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.112)

$$= \frac{1}{\chi_{\text{max}}} \left[\frac{\partial J^c}{\partial \tilde{c}} \tilde{P}_{mn}^{\text{el}} \left(-\frac{1}{3} (2\tilde{E}_{mn}^{\text{el}} + \delta_{mn}) + \frac{1}{2} \tilde{E}_{mn}^{\text{el}} \right) + \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.113)

$$= \frac{1}{\chi_{\text{max}}} \left[-\frac{1}{6} \frac{\partial J^c}{\partial \tilde{c}} \tilde{P}_{mn}^{\text{el}} \tilde{E}_{mn}^{\text{el}} - \frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} \tilde{P}_{mn}^{\text{el}} \delta_{mn} + \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.114)

$$= \frac{1}{\chi_{\text{max}}} \left[-\frac{1}{6} \frac{\partial J^c}{\partial \tilde{c}} \tilde{P}_{mn}^{\text{el}} \tilde{E}_{mn}^{\text{el}} - \frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} \tilde{P}_{mm}^{\text{el}} + \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{ij}^{\text{el}} \tilde{E}_{kl}^{\text{el}} \right]$$
(1.115)

$$= \frac{1}{\chi_{\text{max}}} (\tilde{\mu}_1 + \tilde{\mu}_2 + \tilde{\mu}_3)$$
 (1.116)

$$\tilde{\mu}_0 = \log(\gamma \tilde{c}) \tag{1.117}$$

$$\tilde{\mu}_{1} = -\frac{1}{6} \frac{\partial J^{c}}{\partial \tilde{c}} \left[\tilde{P}_{11}^{\text{el}} \tilde{E}_{11}^{\text{el}} + \tilde{P}_{22}^{\text{el}} \tilde{E}_{22}^{\text{el}} + 2 \tilde{P}_{12}^{\text{el}} \tilde{E}_{12}^{\text{el}} + \tilde{P}_{33}^{\text{el}} \tilde{E}_{33}^{\text{el}} \right]$$
(1.118)

$$\tilde{\mu}_2 = -\frac{1}{3} \frac{\partial J^c}{\partial \tilde{c}} [\tilde{P}_{11}^{\text{el}} + \tilde{P}_{22}^{\text{el}} + \tilde{P}_{33}^{\text{el}}] \tag{1.119}$$

$$\tilde{\mu}_3 = \frac{1}{2} J^c \frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} \tilde{E}_{kl}^{\text{el}} \tilde{E}_{ij}^{\text{el}} \tag{1.120}$$

$$\frac{\partial \tilde{C}_{ijkl}}{\partial \tilde{c}} = \frac{\partial \tilde{\lambda}_{si}(\tilde{c})}{\partial \tilde{c}} \delta_{ij} \delta_{kl} + \frac{\partial \tilde{\mu}_{si}(\tilde{c})}{\partial \tilde{c}} (\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk})$$
(1.121)

$$\tilde{\mu}_{3} = \frac{1}{2} J^{c} \frac{\partial \tilde{\lambda}_{si}(\tilde{c})}{\partial \tilde{c}} \delta_{ij} \delta_{kl} \tilde{E}_{kl}^{\text{el}} \tilde{E}_{ij}^{\text{el}} + \frac{1}{2} J^{c} \frac{\partial \tilde{\mu}_{si}(\tilde{c})}{\partial \tilde{c}} (\delta_{ik} \delta_{jl} \tilde{E}_{kl}^{\text{el}} \tilde{E}_{ij}^{\text{el}} + \delta_{il} \delta_{jk} \tilde{E}_{kl}^{\text{el}} \tilde{E}_{ij}^{\text{el}})$$
(1.122)

$$= \frac{1}{2} J^{c} \frac{\partial \tilde{\lambda}_{si}(\tilde{c})}{\partial \tilde{c}} \tilde{E}_{kk}^{\text{el}} \tilde{E}_{ii}^{\text{el}} + \frac{1}{2} J^{c} \frac{\partial \tilde{\mu}_{si}(\tilde{c})}{\partial \tilde{c}} (\tilde{E}_{ij}^{\text{el}} \tilde{E}_{ij}^{\text{el}} + \tilde{E}_{ji}^{\text{el}} \tilde{E}_{ij}^{\text{el}})$$

$$(1.123)$$

$$= \frac{1}{2} J^{c} \frac{\partial \tilde{\lambda}_{si}(\tilde{c})}{\partial \tilde{c}} (\operatorname{tr}(\mathbf{E}^{el}))^{2} + J^{c} \frac{\partial \tilde{\mu}_{si}(\tilde{c})}{\partial \tilde{c}} \tilde{E}_{ij}^{el} \tilde{E}_{ij}^{el}$$
(1.124)

$$= \frac{1}{2} J^{c} [\tilde{\lambda}'_{si}(c)(\tilde{E}^{el}_{11} + \tilde{E}^{el}_{22} + \tilde{E}^{el}_{33})^{2} + 2\tilde{\mu}'_{si}(\tilde{c})((\tilde{E}^{el}_{11})^{2} + (\tilde{E}^{el}_{22})^{2} + (\tilde{E}^{el}_{33})^{2} + 2(\tilde{E}^{el}_{12})^{2})]$$

$$(1.125)$$

$$\boldsymbol{j} = -\frac{D\chi_{\max}\tilde{c}}{V_m^b}\tilde{\mathbf{F}}^{-1}(\tilde{\mathbf{F}}^{-1})^\mathsf{T}\nabla_{\boldsymbol{X}}\tilde{\mu}$$
(1.126)

$$\tilde{\boldsymbol{j}} = \boldsymbol{j}HV_m^b/(\chi_{\text{max}}D_0) \tag{1.127}$$

$$= -\frac{D}{D_0} H \tilde{c} \tilde{\mathbf{F}}^{-1} (\tilde{\mathbf{F}}^{-1})^\mathsf{T} \nabla_{\mathbf{X}} \tilde{\mu}$$
(1.128)

$$\tilde{j}_x = \frac{-\bar{D}H\tilde{c}}{J^2} \left(\frac{\partial \tilde{\mu}}{\partial X} (\tilde{F}_{12}^2 + \tilde{F}_{22}^2) - \frac{\partial \tilde{\mu}}{\partial Y} (\tilde{F}_{11}\tilde{F}_{12} + \tilde{F}_{21}\tilde{F}_{22}) \right)$$

$$(1.129)$$

$$= \frac{-\bar{D}\tilde{c}}{J^2} \left(\frac{\partial \tilde{\mu}}{\partial \tilde{X}} (\tilde{F}_{12}^2 + \tilde{F}_{22}^2) - \frac{\partial \tilde{\mu}}{\partial \tilde{Y}} (\tilde{F}_{11}\tilde{F}_{12} + \tilde{F}_{21}\tilde{F}_{22}) \right)$$
(1.130)

$$\tilde{j}_y = \frac{-\bar{D}H\tilde{c}}{J^2} \left(\frac{\partial \tilde{\mu}}{\partial Y} (\tilde{F}_{11}^2 + \tilde{F}_{21}^2) - \frac{\partial \tilde{\mu}}{\partial X} (\tilde{F}_{11}\tilde{F}_{12} + \tilde{F}_{21}\tilde{F}_{22}) \right)$$

$$(1.131)$$

$$= \frac{-\bar{D}\tilde{c}}{J^2} \left(\frac{\partial \tilde{\mu}}{\partial \tilde{Y}} (\tilde{F}_{11}^2 + \tilde{F}_{21}^2) - \frac{\partial \tilde{\mu}}{\partial \tilde{X}} (\tilde{F}_{11}\tilde{F}_{12} + \tilde{F}_{21}\tilde{F}_{22}) \right)$$
(1.132)

1.8 Boundary and Initial Conditions

$$\tilde{c}(\tilde{X}, \tilde{Y}, 0) = 0 \tag{1.133}$$

$$\tilde{u}(\tilde{X}, \tilde{Y}, 0) = 0 \tag{1.134}$$

$$\tilde{v}(\tilde{X}, \tilde{Y}, 0) = 0 \tag{1.135}$$

$$\tilde{u}(\tilde{X}, 0, \tilde{t}) = \tilde{v}(\tilde{X}, 0, \tilde{t}) = 0 \tag{1.136}$$

$$\tilde{u}(-1/2, \tilde{Y}, \tilde{t}) = \tilde{u}(1/2, \tilde{Y}, \tilde{t}) = 0 \tag{1.137}$$

$$\tilde{j}_x(\tilde{X}, 1, \tilde{t}) = \tilde{j}_0(1 - \tilde{c}(\tilde{X}, 1, \tilde{t})) \tag{1.138}$$

Bibliography

Lee, E. H. (1969), 'Elastic-plastic deformation at finite strains'.