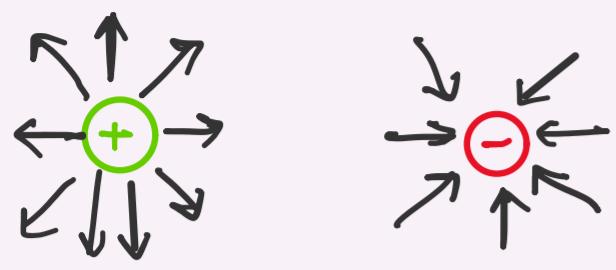


# Electric and Magnetic fields

## Electric Fields

→ Visualized by Electric Flux lines



→ Charges can be isolated

$$\rightarrow \text{Potential } V_n = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q}{n}$$

$$\rightarrow \text{Electric Field } E_n = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q}{n^2}$$

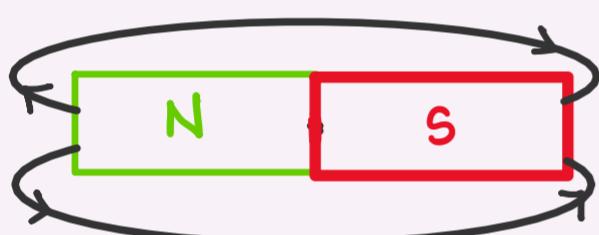
$$\text{So, } E_n = -\frac{dV_n}{dn}$$

## Magnetic Fields

→ No monopoles

→ Can be expressed in terms of flux lines

→ Flux lines are continuous



## Maxwells Equations

→ Gauss' Law

$$\nabla \vec{E} = \frac{P}{\epsilon_0}$$

→ Gauss' Law for magnetism

$$\nabla \vec{B} = 0$$

→ Faraday's Law

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

→ Ampere - Maxwell Law

$$\nabla \times \vec{B} = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\text{Note: } \frac{\partial \vec{B}}{\partial t} = \frac{\partial \vec{B} \cdot \text{Area}}{\partial t} = \frac{\partial \vec{B}}{\partial t}$$

→ 0 in free space

## Nabla / Del Operator

$$\rightarrow \vec{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k} \quad (\text{Del operator})$$

$$\rightarrow \vec{\nabla} \vec{\nabla} = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad (\text{Laplacian operator})$$

$$\text{So, } \nabla V = \frac{\partial V_x}{\partial x} \hat{i} + \frac{\partial V_y}{\partial y} \hat{j} + \frac{\partial V_z}{\partial z} \hat{k} \quad (\text{grad } V)$$

$$\vec{\nabla} \cdot \vec{V} = \frac{\partial V_x}{\partial x} + \frac{\partial V_y}{\partial y} + \frac{\partial V_z}{\partial z} \quad (\text{Divergence } V)$$

$$\vec{\nabla} \times \vec{A} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix} \quad (\text{curl } A)$$

$$\rightarrow \vec{\nabla} \times (\vec{\nabla} \times \vec{A}) = \vec{\nabla} (\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A} \quad (\text{identity})$$

$$\rightarrow \nabla^2 \vec{A} = \left( \frac{1}{v^2} \cdot \frac{\partial^2 \vec{A}}{\partial t^2} \right) \quad (\text{General wave eq.})$$

## ★ Electric waves in free space

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (\text{Faraday's Law})$$

$$\Rightarrow \vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = \vec{\nabla} \times \left( -\frac{\partial \vec{B}}{\partial t} \right)$$

$$\vec{\nabla} (\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} = \vec{\nabla} \times \left( -\frac{\partial \vec{B}}{\partial t} \right) \quad (\text{Using Identity})$$

$$= -\left( \frac{\partial \vec{\nabla} \times \vec{B}}{\partial t} \right)$$

For free space,  $\vec{\nabla} \cdot \vec{E} = 0$

$$\text{So, } -\nabla^2 \vec{E} = -\left( \frac{\partial \vec{\nabla} \times \vec{B}}{\partial t} \right)$$

$$\nabla^2 \vec{E} = \left( \frac{\partial \vec{\nabla} \times \vec{B}}{\partial t} \right)$$

$$\nabla^2 \vec{E} = \mu_0 \epsilon_0 \cdot \frac{\partial^2 \vec{E}}{\partial t^2} \quad (\text{Ampere - Maxwell law})$$

$$\nabla^2 \vec{E} = \frac{1}{c^2} \cdot \frac{\partial^2 \vec{E}}{\partial t^2} \quad (c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}) \quad (\text{wave eq. for Electric wave})$$

## ★ Magnetic waves

$$\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad (\text{Ampere - Maxwell Law})$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \vec{\nabla} \times \left( \mu_0 \epsilon_0 \cdot \frac{\partial \vec{E}}{\partial t} \right)$$

$$\vec{\nabla} (\vec{\nabla} \cdot \vec{B}) - \nabla^2 \vec{B} = \mu_0 \epsilon_0 \cdot \left( \frac{\partial \vec{\nabla} \times \vec{E}}{\partial t} \right)$$

$$-\nabla^2 \vec{B} = \mu_0 \epsilon_0 \left( \frac{\partial \vec{\nabla} \times \vec{E}}{\partial t} \right)$$

$$-\nabla^2 \vec{B} = \mu_0 \epsilon_0 \left( -\frac{\partial^2 \vec{B}}{\partial t^2} \right) \quad (\text{Faraday's Law})$$

$$\nabla^2 \vec{B} = \frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2} \quad (c = \frac{1}{\sqrt{\mu_0 \epsilon_0}})$$

## Mutually perpendicularity between B and E

Consider equation  $E_x = E_0 \cos(\omega t + kz) \hat{i}$

$$\nabla \times \vec{E} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix} = \hat{i}(0) - \hat{j}\left(-\frac{\partial E_x}{\partial z}\right) + \hat{k}\left(-\frac{\partial E_x}{\partial y}\right)$$

$$= \frac{\partial(E_0 \cos(\omega t + kz))}{\partial z} \hat{j} - \frac{\partial(E_0 \cos(\omega t + kz))}{\partial y} \hat{k}$$

$$\nabla \times \vec{E} = -kE_0 \sin(\omega t + kz) \hat{j}$$

$$\Rightarrow -\frac{\partial \vec{B}}{\partial t} = -kE_0 \sin(\omega t + kz) \hat{j} \quad (\text{Faraday's Law})$$

$$\int \partial B = kE_0 \int \sin(\omega t + kz) dt \hat{j}$$

$$B = \frac{kE}{\omega} \cos(\omega t + kz) \hat{j}$$

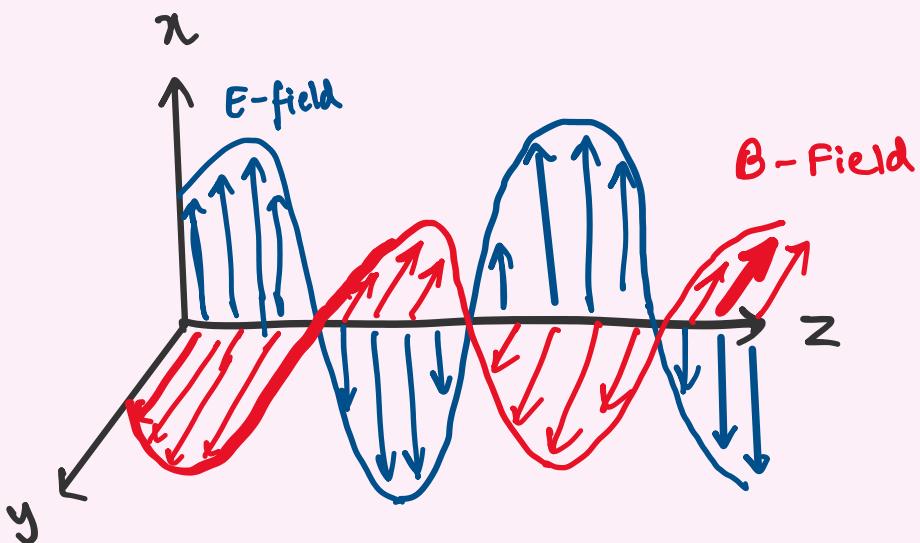
$$B = \frac{E_0}{c} \cos(\omega t + kz) \hat{j} \quad (c = \frac{\omega}{k})$$

$$B = B_0 \cos(\omega t + kz) \hat{j} \quad (B_0 = \frac{E_0}{c})$$

So, B is in y component

& E is in x component

Hence, B & E are mutually perpendicular



### 3. Energy Transportation by EMW



3. Energy  
Transporta...

#### Energy Transportation by EMW

For Electromagnetic waves

$$\rightarrow v \text{ (velocity)} = c = 3 \times 10^8$$

$$\rightarrow \lambda \text{ (wavelength)} = \frac{2\pi}{k}$$

$$\rightarrow f \text{ (frequency)} = \frac{v}{\lambda}$$

$$\rightarrow T \text{ (Period)} = \frac{1}{f}$$

$$\rightarrow B(x, t) = \frac{E(x, t)}{c}$$

$$\rightarrow A \text{ (Amplitude)} = \frac{B_0}{E_0}$$

Energy content of electric component

$$E_E = \frac{1}{2} \epsilon_0 E_n^2$$

Energy content of magnetic component

$$E_M = \frac{1}{2} \frac{B_n^2}{\mu_0}$$

Total energy content of EM Wave

$$\begin{aligned} E_T &= E_E + E_M = \frac{1}{2} \epsilon_0 E_n^2 + \frac{1}{2} \frac{B_n^2}{\mu_0} \\ &= \frac{1}{2} \epsilon_0 E_n^2 + \frac{1}{2} \frac{E_n^2}{c^2 \mu_0} \quad (B_0 = \frac{E_0}{c}) \\ &= \frac{1}{2} \epsilon_0 E_n^2 + \frac{1}{2} \epsilon_0 E_n^2 \quad (\frac{1}{c^2} = \mu_0 \epsilon_0) \\ E_T &= \epsilon_0 E_n^2 \quad (\text{in } z \text{ direction}) \end{aligned}$$

Power dissipated  
= Energy Flux  $\times$  Area  
=  $E_0 E^2 c \times A$

Average energy of EMW

$\rightarrow$  Avg. Energy of EMW = Energy transported in 1 cycle

$$\begin{aligned} \text{Total energy cycle} &= \langle \text{Energy} \rangle = \frac{c \epsilon_0}{T} \int_0^T E_n^2 dt \\ &= \frac{c \epsilon_0}{T} \int_0^T E_{0n}^2 \cos^2(\omega t + kx) dt \\ &= \frac{1}{2} \frac{c \epsilon_0}{T} \cdot E_{0n}^2 \cdot T = \frac{1}{2} c \cdot \frac{B_{0y}}{\mu_0} \\ &= \frac{1}{2} \frac{E_{0n} B_{0y}}{\mu_0} \end{aligned}$$

Intensity  $\propto |\text{Amplitude}|^2$

$$? E(x, t) = 10^3 \cos(\omega t - 3\pi \times 10^6 z)$$

Find  $v, \lambda, f, T, B(x, t), A$

$$A. v = 3 \times 10^8 \text{ m s}^{-1}$$

$$\lambda = \frac{2\pi}{k} = \frac{2\pi}{3\pi \times 10^6} = 6.67 \times 10^{-7} \text{ m}$$

$$f = \frac{v}{\lambda} = 4.5 \times 10^{14} \text{ Hz}$$

$$T = 2.22 \times 10^{-16} \text{ s}$$

$$B(x, t) = \frac{E(x, t)}{c} = \frac{10^3}{3 \times 10^8} = 3.33 \times 10^{-6}$$

Poynting Vector ( $\vec{S}$ )

$\rightarrow \vec{S}$  describes EM energy transported per unit time per unit volume

$$\vec{S} = \frac{1}{\mu_0} E \times B = c^2 \epsilon_0 E \times B = c \epsilon_0 E^2$$

$\rightarrow$  Direction of  $\vec{S}$  is direction of propagation of waves

$\rightarrow \vec{S}$  is time dependant

$$j_{\text{flux}} = \frac{E_T \times CAF}{AT}$$

$$= c \times \epsilon_0 E^2$$

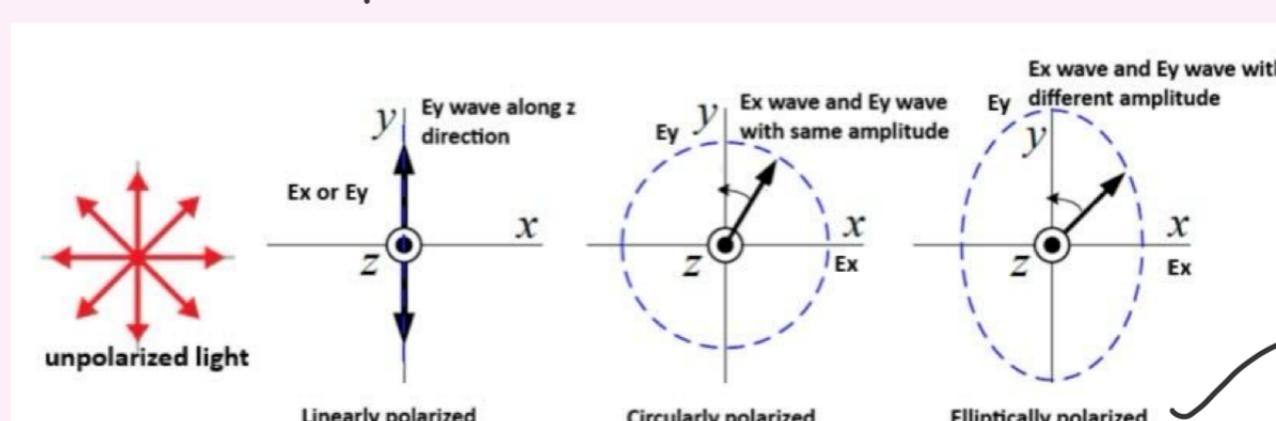
$$= c \times \epsilon_0 E (BC)$$

$$= c^2 \times \epsilon_0 E \times B$$

$$\vec{S} = \frac{1}{\mu_0} E \times B$$

Polarization of EMW

$\rightarrow$  Natural light is unpolarized  
All planes of propagation being equally probable



- Plane wave is called linearly polarized
- Addition of horizontally & vertically linearly polarized waves of same amp in same phase results in linearly polarized wave at 45°

- Light is composed of 2 plane waves of equal amplitude but differ in phase by 90° are circularly polarized

2 plane waves of differing amplitude are related in phase by 90° (or) relative phase other than 90° are elliptically polarized

## 4. Failure of classical EMW theory and blackbody radiation



4. Failure of classical...

### 4. Failure of classical EMW theory and blackbody radiation

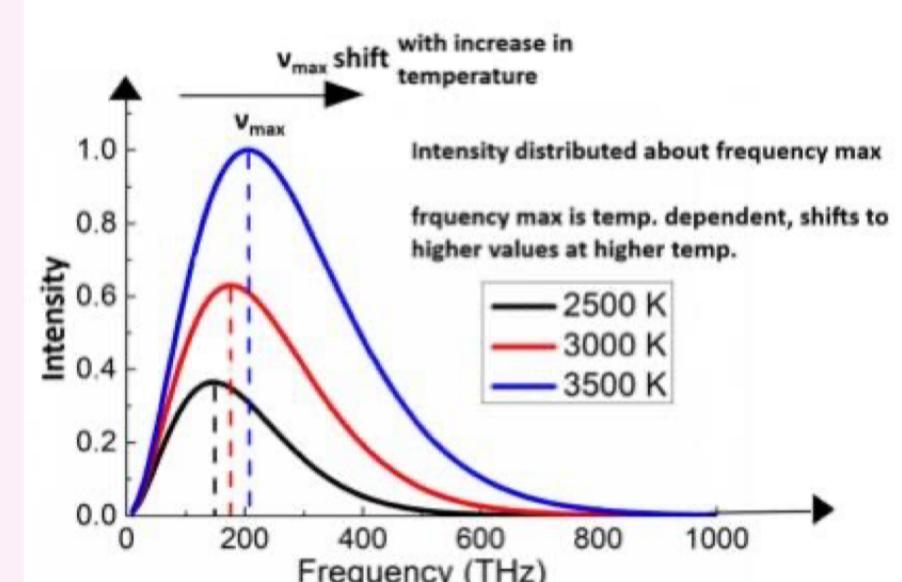
- EM Radiation  $\Rightarrow$  Mutual perpendicularity b/w E & B waves
- Classical wave theory  $\Rightarrow$  Energy content of wave  $\propto$  (Amplitude of waves)<sup>2</sup>
- Wave theory  $\Rightarrow$  Explains reflection, refraction, interference, diffraction & polarization of light

Couldn't explain -

- Photo-electric effect
- Spectrum of Hydrogen emissions (Atomic Spectra)
- Black-body Radiation Spectrum
- Compton Scattering

### Blackbody Radiation (Hypothetical Body)

- Blackbody  $\Rightarrow$  Materials which absorb all incident rays. On heating, it will emit all wavelengths of radiation that was absorbed
- ex: sun, incandescent bulb etc,  $\Rightarrow$  Emissivity = 1



$\Rightarrow$  Radiation only depends on temperature. Not material of object

In Intensity vs Wavelength graph, The maximum intensity of prominent wavelength is shifting towards shorter wavelength as temperature increases

There are distinct curves for different temperatures

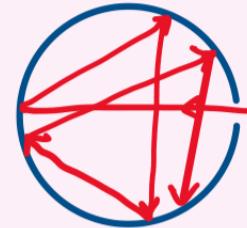
Ex: sun, incandescent bulb etc,

### Rayleigh - Jeans Law

- Assumed black body as cavity oscillators
- No of oscillators b/w  $\nu$  &  $\nu + \delta\nu$  is  $dN = \frac{8\pi\nu^2 d\nu}{c^3}$  (No of modes  $\propto \nu^2$ )
- $\langle E \rangle = k_B T$  (Maxwell - Boltzmann distribution law)
- $\rho(\nu)d\nu = \langle E \rangle dN$  (Energy density, according to  $= k_B T \times \frac{8\pi\nu^2 d\nu}{c^3}$  Rayleigh - Jeans Law)

$\hookrightarrow$  But it contradicts with experimental observations

Model that doesn't allow incident radiation to escape because of inner multiple reflections This cavity when heated emit radiations of every possible frequency & rate of emission  $\uparrow$  as Temperature  $\uparrow$



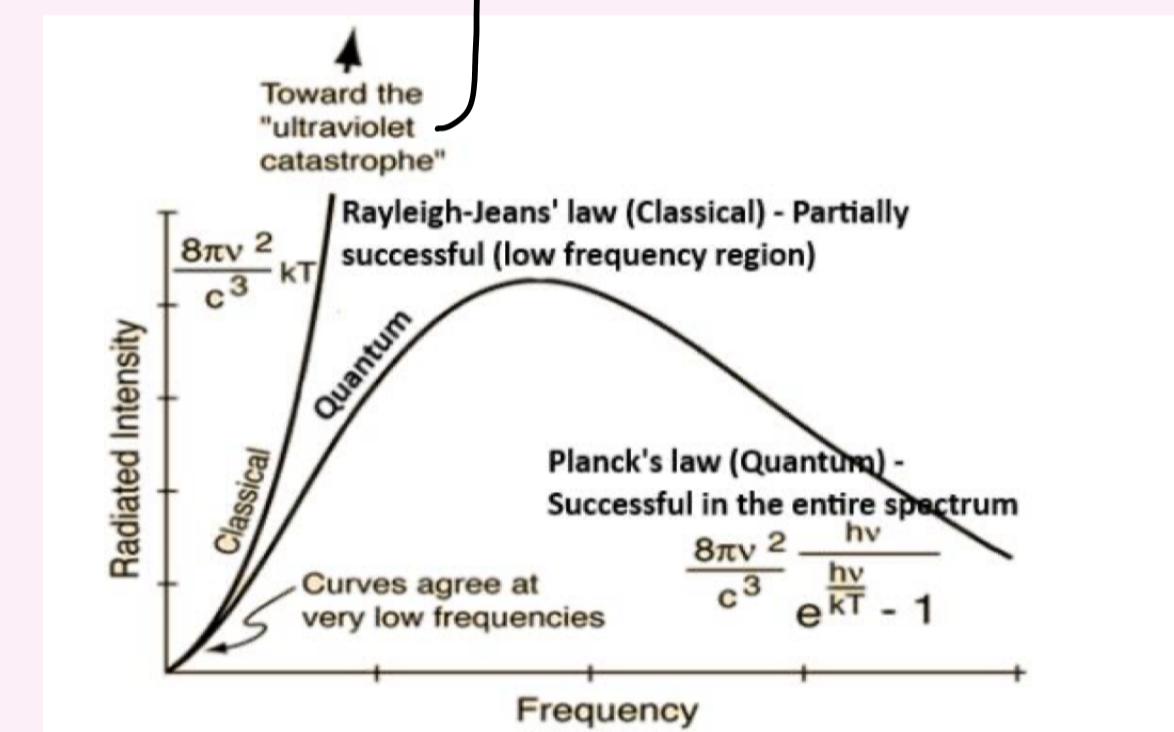
$\hookrightarrow$  Breakdown of Rayleigh Jeans at higher frequencies

### Max - Planck hypothesis

- Solution for R-JL / Ultraviolet Catastrophe
- Energy of oscillator model of black body must be restricted to finite multiples of fundamental natural frequency times constant ( $E = nh\nu$ )
- $\langle E \rangle = \frac{h\nu}{e^{\frac{h\nu}{kT}} - 1}$

and,  $\rho(\nu)d\nu = \langle E \rangle dN$

$$= \frac{h\nu}{e^{\frac{h\nu}{kT}} - 1} \cdot \frac{8\pi\nu^2 d\nu}{c^3} = \frac{8\pi h\nu^3}{c^3} \cdot \frac{1}{e^{\frac{h\nu}{kT}} - 1} \cdot d\nu$$



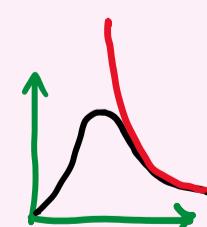
? Frequency of harmonic oscillator at 50°C is  $6.2 \times 10^{12} \text{ s}^{-1}$   
Estimate average energy of oscillator as per planck's idea of cavity oscillator, also compare the same with classical energy & average energy by R-J Law

A. As per planck's idea,  $\langle E \rangle = \frac{h\nu}{e^{\frac{h\nu}{kT}} - 1} = \frac{6.63 \times 10^{-34} \times 6.2 \times 10^{12}}{e^{\frac{6.63 \times 6.2 \times 10^{-34}}{1.38 \times 10^{-23} \times 323}} - 1} = 2.713 \times 10^{-21} \text{ J}$

As per classical analysis,  $\langle E \rangle = k_B T = 1.38 \times 10^{-23} \times 323 = 4.457 \times 10^{-21} \text{ J}$

Difference =  $1.744 \times 10^{-21} \text{ J}$

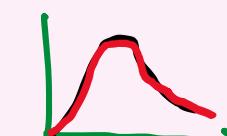
Rayleigh Jeans  $\Rightarrow u(\lambda) d\lambda = \frac{8\pi kT}{\lambda^4} d\lambda$



Wiens  $\Rightarrow u(\lambda) d\lambda = \frac{c_1}{\lambda^5} e^{-\frac{c_2}{\lambda}} d\lambda$  (Exponential decay)



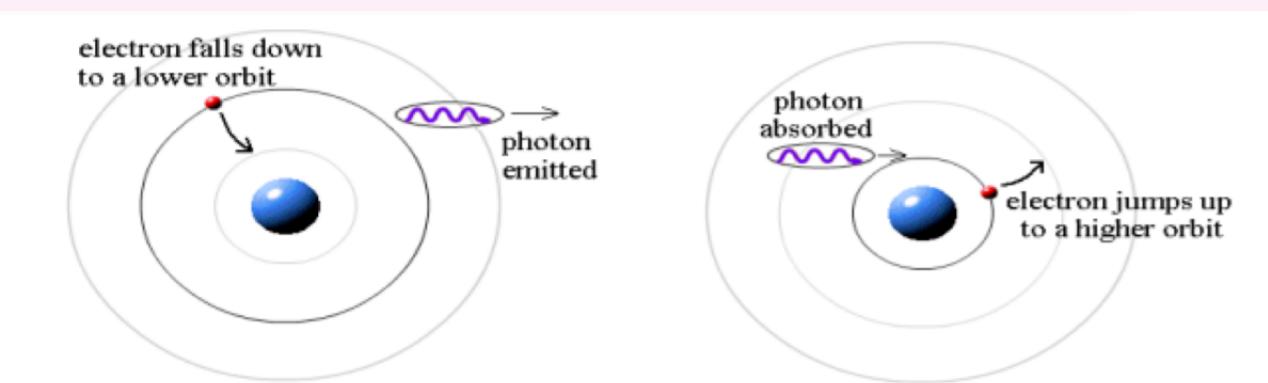
Plancks  $\Rightarrow u(\lambda) d\lambda = \frac{8\pi hc}{\lambda^5} \left[ \frac{1}{e^{\frac{hc}{kT}} - 1} \right] d\lambda$



## 5. Atomic Spectra, Photoelectric Effect and Compton Shift

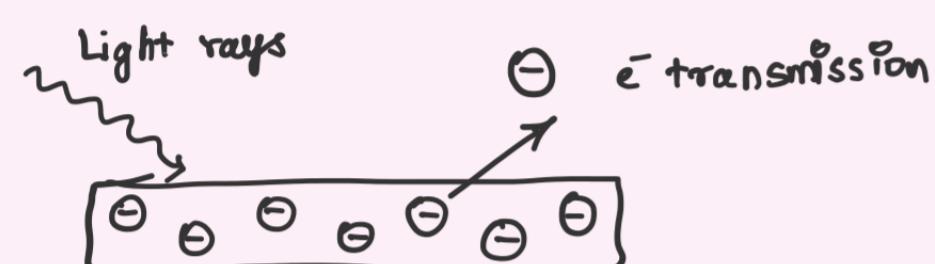
## Atomic Spectra

- Different elements have different spectra, so atomic spectroscopy allows identification of a sample's elemental composition
- Atomic absorption lines in solar spectrum are called Fraunhofer lines
  - Discrete lines  $\Rightarrow$  Emission Spectra
  - No discrete lines  $\Rightarrow$  Absorption Spectra
- Based on Max Planck, Energy comes in quanta & absorbed/emitted in terms of quanta



## Photoelectric Effect

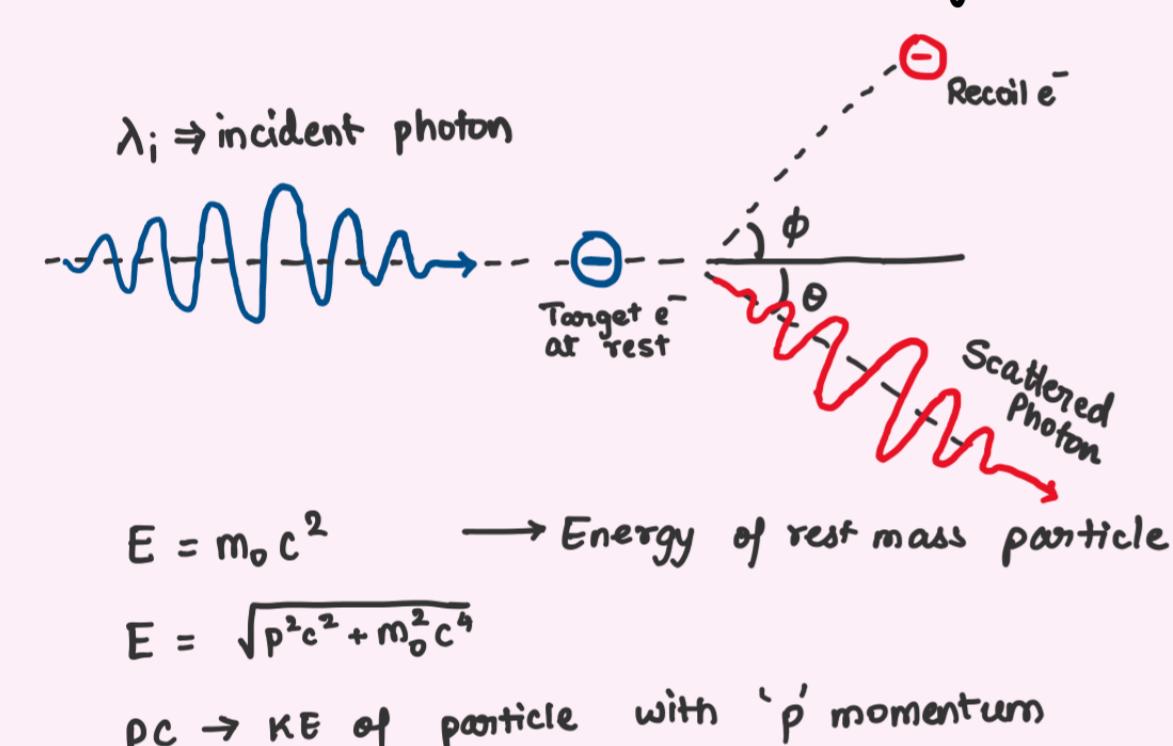
- $e^-$  emission from metals under irradiation



- Instantaneous emission of  $e^-$  with KE dependant on Wavelength of radiation
- Energy of  $e^-$  independent of intensity of radiation
- $h\nu = W + KE$
- Waves can be of dual nature

## Compton Shift

$\Rightarrow$  Process of scattering of radiation from loosely bound  $e^-$ , which is released from atom & remainder of energy is rereleased as electromagnetic radiation



$$\Delta\lambda = \lambda_f - \lambda_i = \frac{h}{m_e c} (1 - \cos\theta)$$

$$\left[ \frac{h}{m_e c} = 2.45 \text{ pm} \right]$$

$$\frac{h}{m_e c} = \lambda_c \text{ (Compton wavelength)}$$

$\Delta\lambda$  is independent of incident wavelength  
 $\Delta\lambda$  dependant on scattering angle

For  $e^-$ ,  $\lambda_c = 2.42 \times 10^{-12} \text{ m}$

If  $\theta = 180^\circ$ , max shift

$$P_i = \frac{E_i}{c} = \frac{h\nu_i}{c} = \frac{h}{\lambda_i}$$

Q. X-rays of wavelength  $0.112 \text{ nm}$  is scattered from a carbon target. Calculate wavelength of x-rays scattered at an angle  $90^\circ$  wrt original direction. What's the energy lost by x-ray photons? What's the energy gained by  $e^-$ ? If incident x-ray retraces back what will be the shift?

$$\Delta\lambda = \frac{h}{m_e c} (1 - \cos\theta)$$

$$= 2.42 \times 10^{-12} (1 - 0)$$

$$= 0.0242 \text{ Å}$$

Wavelength of scattered x-rays,  $\lambda' = \lambda + \Delta\lambda$

$$= 0.112 + 0.0242$$

$$= 1.1442 \text{ Å}$$

Energy lost by x-ray photons =  $h\nu - h\nu'$

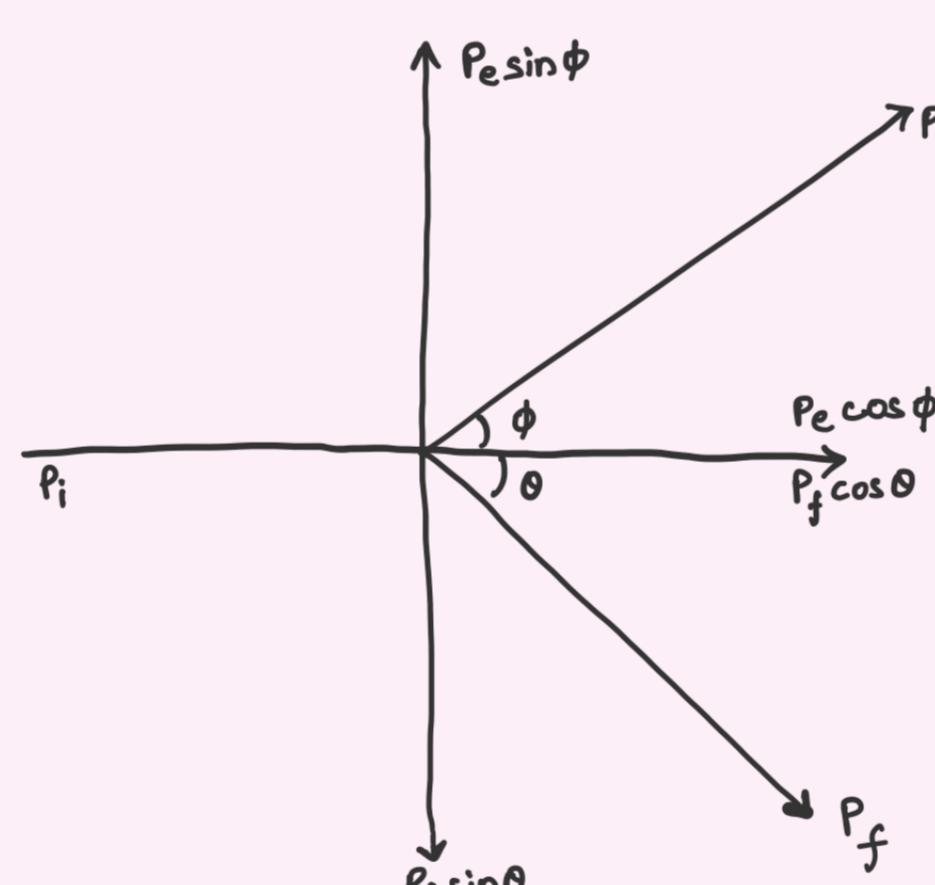
$$= \frac{hc}{\lambda} - \frac{hc}{\lambda'}$$

$$= 1.77 \times 10^{-15} - 1.74 \times 10^{-15}$$

$$= 0.03 \times 10^{-15}$$

$$= 3 \times 10^{-17} \text{ J}$$

$$= 187.2 \text{ eV} \quad (1 \text{ J} = 6.242 \times 10^{18} \text{ eV})$$



## Compton Shift Derivation

Along Incident direction  $\Rightarrow P_i + 0 = P_f \cos\theta + P_e \cos\phi$

Along perpendicular direction  $\Rightarrow 0 = P_f \sin\theta - P_e \sin\phi$

Before & After collision conservation,

$$P_e^2 = P_i^2 + P_f^2 - 2P_i P_f \cos\theta \quad \text{---} ①$$

Before & After energy conservation

$$P_i c + m_e c^2 = P_f c + \sqrt{P_e^2 c^2 + m_e^2 c^4}$$

$\Rightarrow$  Square on Both sides & divide by  $c^2$

$$P_e^2 = P_i^2 + P_f^2 - 2P_i P_f - 2m_e c(P_i - P_f) \quad \text{---} ②$$

From ① & ②

$$P_i^2 + P_f^2 - 2P_i P_f \cos\theta = P_i^2 + P_f^2 - 2P_i P_f + 2m_e c(P_i - P_f)$$

Also, substitute  $P_i = \frac{h}{\lambda_i}$  &  $P_f = \frac{h}{\lambda_f}$

$$\cancel{\lambda_i \lambda_f (1 - \cos\theta)} = \cancel{m_e c (P_i - P_f)}$$

$$(1 - \cos\theta) = m_e c \frac{(P_i - P_f)}{P_i P_f}$$

$$1 - \cos\theta = m_e c \left( \frac{1}{P_f} - \frac{1}{P_i} \right)$$

$$1 - \cos\theta = m_e c \left( \frac{\lambda_f - \lambda_i}{h} \right)$$

$$\lambda_f - \lambda_i = \frac{h}{m_e c} (1 - \cos\theta) \Rightarrow \text{Proved particle nature of EM radiation}$$

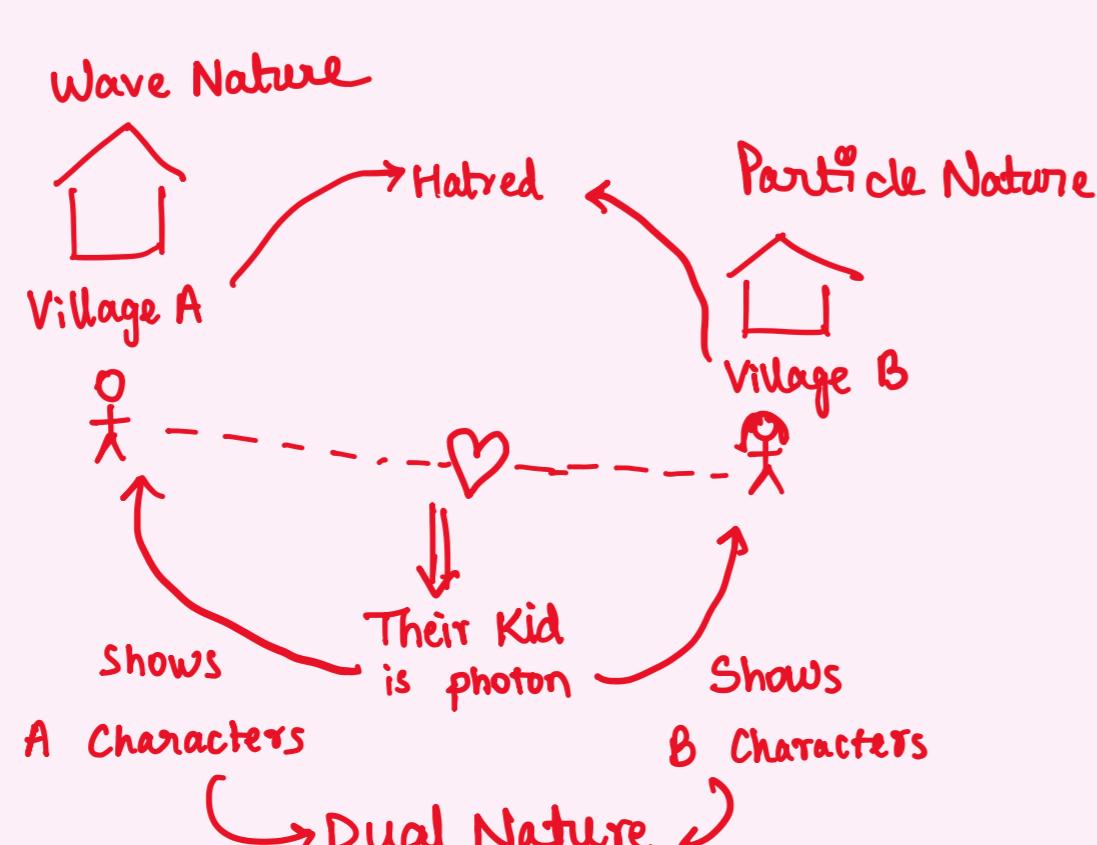
$$E = pc$$

$$E = \frac{hc}{\lambda}$$

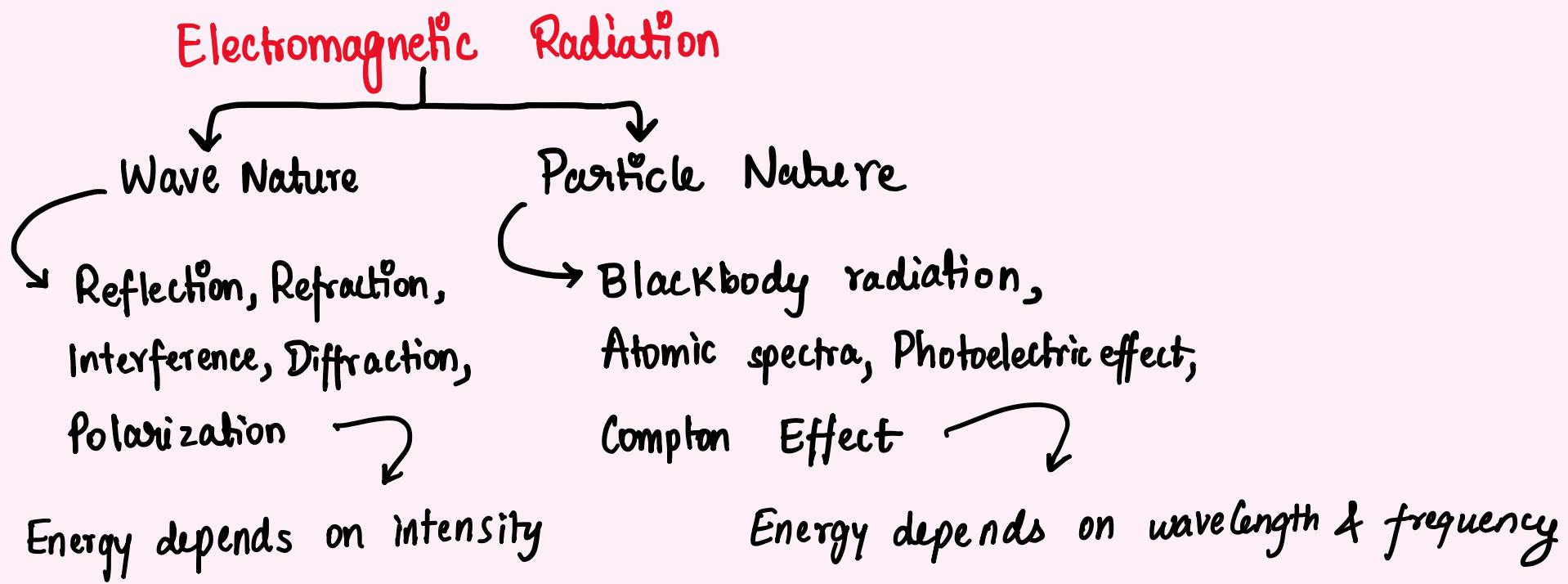
$$pc = \frac{hc}{\lambda}$$

$$\lambda = \frac{h}{p} \quad \begin{matrix} \text{wave} \\ \text{property} \end{matrix} \quad \begin{matrix} \text{particle property} \end{matrix}$$

$$\lambda = \frac{h}{mv} \quad \begin{matrix} \text{particle's speed} \\ (\text{must be moving}) \end{matrix}$$



## 6. Dual Nature of Radiation

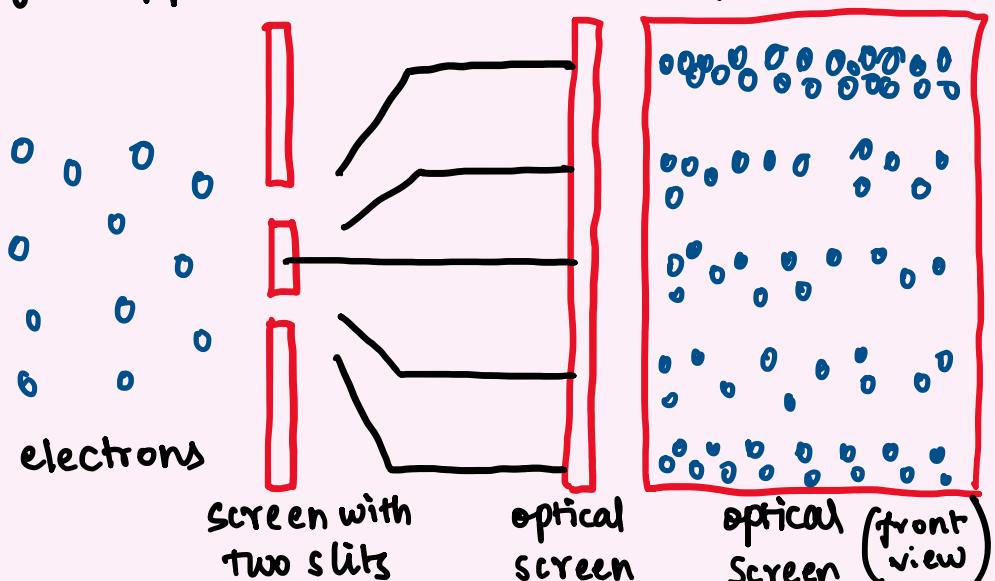


### De Broglie Hypothesis

- Moving matter should exhibit wave characteristics
- $\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{h}{\sqrt{2mE}} = \frac{h}{\sqrt{2meV}}$
- Wavelength of macro particles are too small to be measured  
Wavelength of sub-atomic particles are measurable ( $\sim 10^{-10}$  m)

### Double Slit Experiment

- Diffraction is a characteristic wave phenomenon
  - When single photon is sent through double-slit, it behaves like a wave by creating an interference pattern on screen
- However, when photons are detected individually on the screen, they appear as discrete particles hitting specific points on the screen



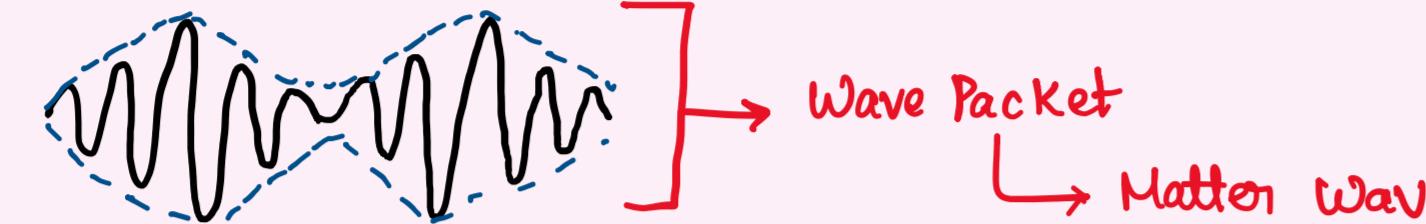
## 7. Matter Waves

→ Wave packets describe matter waves  
(Defined wavelength & amplitude maximum)

$$\rightarrow \text{Wavy line} \Rightarrow y_1 = A \sin(\omega t - kx)$$

$$\rightarrow \text{Wavy line} \Rightarrow y_2 = A \sin((\omega + \Delta\omega)t - (k + \Delta k)x)$$

$$\text{Wave packet} = y_1 + y_2$$



$$\text{If we use } \sin A + \sin B = 2 \sin\left(\frac{A+B}{2}\right) \cos\left(\frac{A-B}{2}\right)$$

$$\begin{aligned} y_1 + y_2 &= 2A \sin\left(\frac{(2\omega + \Delta\omega)t - (2k + \Delta k)x}{2}\right) \cos\left(\frac{\Delta\omega t - \Delta kx}{2}\right) \\ &\approx 2A \sin(\omega t - kx) \cos\left(\frac{\Delta\omega t - \Delta kx}{2}\right) \end{aligned}$$

⇒ We start with constant Amplitude & End with time varying Amplitude

→ Spread around central maximum is approximate position of particle

### Phase Velocity ( $v_p$ )

→ It is the velocity with which a constant phase point marked on the superposed wave moves

$$v_p = \frac{dx}{dt} = \frac{\omega}{k}$$

$$\left| \begin{array}{l} kx - \omega t = C \\ \frac{dkx}{dt} - \omega = 0 \Rightarrow \frac{dx}{dt} = \frac{\omega}{k} \end{array} \right.$$

### Group Velocity

→ It is the velocity of common velocity of the superposed wave

$$\left| \begin{array}{l} \frac{\Delta kx - \Delta \omega t}{2} = C \\ \frac{\Delta k}{dt} - \Delta \omega = 0 \Rightarrow \frac{dx}{dt} = \lim_{\Delta k \rightarrow 0} \frac{\Delta \omega}{\Delta k} = \frac{d\omega}{dk} \end{array} \right.$$

### Relation b/w $v_g$ & $v_p$

$$\rightarrow v_g = \frac{d\omega}{dk} = \frac{d(v_p k)}{dk} = v_p + k \cdot \frac{dv_p}{dk}$$

$$\text{Also, } \frac{dv_p}{dk} = \frac{dV_p}{d\lambda} \cdot \frac{d\lambda}{dk} = \frac{dV_p}{d\lambda} \cdot \frac{d\left(\frac{2\pi}{k}\right)}{dk} = \frac{dV_p}{d\lambda} \left(-\frac{2\pi}{k^2}\right)$$

$$v_g = v_p - k \cdot \frac{dv_p}{dk} \left(\frac{2\pi}{k^2}\right)$$

$$= v_p - \frac{dV_p}{d\lambda} \left(\frac{2\pi}{k}\right)$$

$$= v_p - \lambda \frac{dV_p}{d\lambda}$$

$$\text{Case 1 : } v_g = v_p$$

→ Only when  $\frac{dV_p}{d\lambda} = 0$  (Light is not dispersive)

$\lambda$  can't be 0

$$\text{Case 2 : } v_g < v_p$$

$$\rightarrow v_g = v_p/2$$

$$\Rightarrow v_g = v_p - \lambda \frac{dV_p}{d\lambda}$$

$$\frac{v_p}{2} = v_p - \lambda \frac{dV_p}{d\lambda}$$

$$\frac{v_p}{2} = \lambda \frac{dV_p}{d\lambda} \Rightarrow \int \frac{d\lambda}{\lambda} = 2 \int \frac{dV_p}{v_p}$$

$$\ln \lambda \propto 2 \ln v_p$$

$$v_p \propto \sqrt{\lambda}$$

$$\text{Case 3 : } v_g > v_p$$

$$\rightarrow v_g = 2v_p$$

$$\Rightarrow v_g = v_p - \lambda \frac{dV_p}{d\lambda}$$

$$v_p = -\lambda \frac{dV_p}{d\lambda}$$

$$\int \frac{dV_p}{v_p} = - \int \frac{d\lambda}{\lambda}$$

$$\ln v_p \propto -\ln \lambda$$

$$v_p \propto \lambda^{-1}$$

$$\rightarrow v_g = \frac{d\omega}{dk}$$

$$\omega = \frac{E}{\hbar}$$

$$\left( \hbar = \frac{h}{2\pi} \right)$$

$$d\omega = \frac{dE}{\hbar}$$

$$w = 2\pi\nu$$

$$= \frac{2\pi}{h} \times E$$

$$(E=h\nu)$$

$$k = \frac{p}{\hbar}$$

$$\lambda = \frac{h}{p}$$

$$= \frac{2\pi}{h} \times p$$

$$v_g = \frac{d\omega}{dk} = \frac{dE}{dk} \frac{K}{dp} = \frac{d(\frac{p^2}{2m})}{dp}$$

$$= \frac{1}{2m} \cdot 2p = \frac{p}{m} = v_{\text{particle}}$$

## 8. Heisenberg's Analysis

- Wave packets have uncertainties
- $\Delta x$ : Spread in estimation of position
- $\Delta K$ : propagation constant
- $\Rightarrow \Delta x \cdot \Delta K \geq \frac{1}{2}$
- $\Delta x \cdot \Delta p \geq \frac{\hbar}{4\pi}$

### Heisenberg's Uncertainty Principle

#### 1) Position Momentum Uncertainty

→ Position & momentum can't be determined simultaneously with unlimited precision

$$\Delta x \cdot \Delta p \geq \frac{\hbar}{2} \geq \frac{\hbar}{4\pi}$$

#### 2) Energy Time Uncertainty

→ Energy & time can't be determined simultaneously with unlimited precision

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{2} \geq \frac{\hbar}{4\pi}$$

#### 3) Uncertainty Relation for Circular motion

→ Angular position & Angular Momentum can't be determined simultaneously with unlimited precision

$$\Delta \theta \cdot \Delta L \geq \frac{\hbar}{2} \geq \frac{\hbar}{4\pi}$$

### Applications of Uncertainty Principle

#### 1) Non-existence of $e^-$ inside nuclei

→  $e^-$  don't exist inside nuclei but are emitted from nucleus during  $\beta$  decay with energies of order of 8 MeV

→ Assuming  $e^-$  to be inside nucleus, then  $\Delta x \approx 10^{-14} m \approx$  nuclear diameter

$$So, \Delta x \cdot \Delta p = \frac{\hbar}{2}$$

$$\text{Then, } \Delta p = \frac{6.63 \times 10^{-34}}{4\pi \times 10^{-14}} \approx 5.28 \times 10^{-21} \text{ kg m s}^{-1}$$

Hence,  $p \approx \Delta p$

$$KE = \frac{p^2}{2m} = \frac{(\Delta p)^2}{2m} = \frac{(5.28 \times 10^{-21})^2}{2 \times 9.1 \times 10^{-31} \times 1.6 \times 10^{-30}} = 96 \text{ MeV}$$

(→  $e^-$  energy should be quite high to be inside nuclei)

→ Thus,  $e^-$  doesn't reside inside nuclei

#### 2) Gamma ray microscope (a thought experiment)

→ Experiment to measure position of electron

→ To observe the electron, wavelength of radiation must be of a comparable size to  $e^-$

So, wavelength of  $\gamma$  rays =  $10^{-12} m$

Resolution of microscope  $\Rightarrow \Delta x = \frac{\lambda}{\sin \theta}$

(→ Min. distance at which 2 distinct points of a specimen can still be seen)

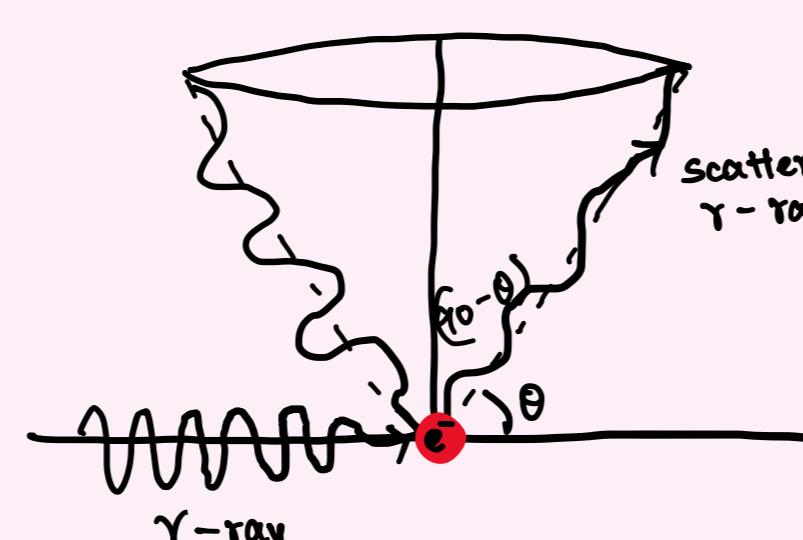
→ High energy  $\gamma$  rays impart momentum to  $e^- \Rightarrow$  Compton Effect

$$p_n \approx \pm \frac{h}{\lambda} \sin \theta$$

$$\Delta p_n \approx \frac{2h \sin \theta}{\lambda} (< p_n)$$

$$So, \Delta x \cdot \Delta p_n = \frac{\lambda}{\sin \theta} \cdot \frac{2h \sin \theta}{\lambda} \approx 2h > \frac{\hbar}{4\pi}$$

Thus, position & momentum can't be found simultaneously



### Other forms of Uncertainty Relations

#### 1) Position - Wavelength uncertainty relation

$$\Delta x \cdot \Delta p = \Delta \left( \frac{h}{\lambda} \right) = h \cdot \Delta \left( \frac{1}{\lambda} \right) = h \cdot \left( -\frac{1}{\lambda^2} \Delta \lambda \right)$$

$$\Delta x \cdot h \left( -\frac{1}{\lambda^2} \Delta \lambda \right) \geq \frac{\hbar}{4\pi}$$

$$\Delta x \cdot \Delta \lambda \geq \frac{\hbar^2}{4\pi} \Rightarrow \Delta x \cdot \Delta \lambda \geq \left| \frac{\lambda^2}{4\pi} \right|$$

#### 2) Uncertainty in terms of position & velocity

$$\Delta p = \Delta(mv)$$

$$\Delta x \cdot \Delta v \geq \frac{\hbar}{4\pi m}$$

#### 3) Position-propagation vector uncertainty relation

$$\Delta x \cdot \Delta K \geq \frac{1}{2}$$

4) Min. uncertainty in 1 parameter corresponds to Max. uncertainty in other

$$\Delta x_{\min} \cdot \Delta p_{\max} \geq \frac{\hbar}{4\pi}$$

#### 5) Accuracy & Uncertainty

Uncertainty in velocity  $\Delta v = \text{velocity} \times \text{accuracy}$   
 $= \text{velocity} \times \% \text{ error}$

#### 6) Energy-time uncertainty relation ( $\lambda$ & $t$ terms)

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{4\pi}$$

$$E = \frac{hc}{\lambda} \Rightarrow \Delta E = \Delta \left( \frac{hc}{\lambda} \right) = hc \cdot \Delta \left( \frac{1}{\lambda} \right) = hc \left( -\frac{1}{\lambda^2} \Delta \lambda \right)$$

$$\Delta t \cdot hc \left( -\frac{1}{\lambda^2} \Delta \lambda \right) \geq \frac{\hbar}{4\pi}$$

$$\Delta t \cdot \Delta \lambda \geq \frac{\hbar^2}{4\pi c} \geq \left| \frac{\lambda^2}{4\pi c} \right|$$

#### 7) Energy-time uncertainty relation ( $v$ & $t$ terms)

$$E = hv$$

$$\Delta E = \Delta hv = h \cdot \Delta v$$

$$\Delta E \cdot \Delta v \geq \frac{1}{4\pi}$$

Q. In the Balmer series of the hydrogen spectra, the longest wavelength spectral line with a wavelength  $656 \text{ nm}$  is emitted when  $e^-$  makes a transition from  $n=3$  to  $n=2$  state. This wavelength, even though pure is seen to have a small spread around  $\lambda$  of the order of  $10^{-5} \text{ nm}$ . Estimate the min. time spent by the electron before deexcitation

A.  $\Delta E \cdot \Delta t \geq \frac{\hbar}{2}$

$$\Delta E = \frac{hc}{\lambda^2} \cdot \Delta \lambda$$

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{4\pi}$$

$$\Delta t \geq \frac{\hbar}{4\pi} \cdot \frac{\lambda^2}{hc \cdot \Delta \lambda}$$

$$= \frac{(656 \times 10^{-9})^2}{4\pi \times 3 \times 10^8 \times 10^{-5} \times 10^{-9}}$$

$$= 1.54 \times 10^{-8} \text{ s}$$

$$= 11.4 \text{ ns}$$

$$\Delta \lambda = \frac{36}{1.097 \times 10^7 \times 5} = 0.000065 \text{ nm}$$

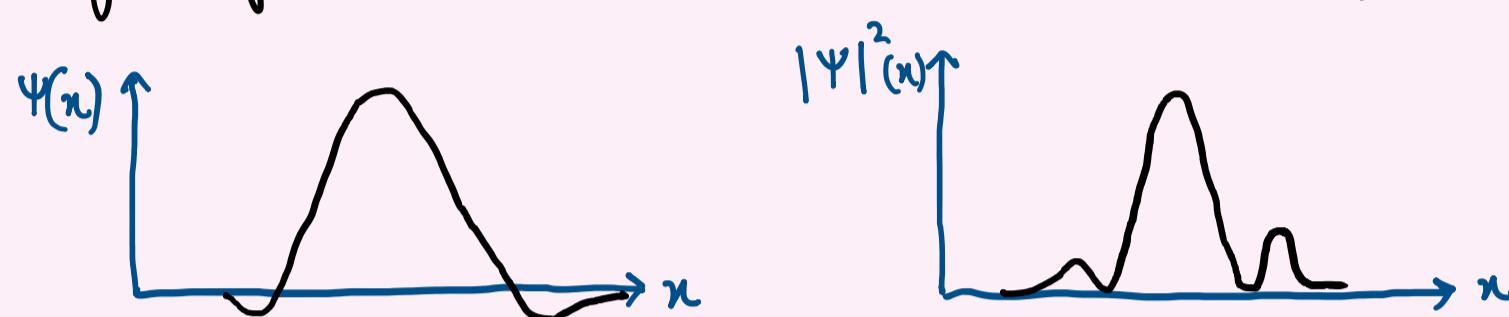
$$= 6.5 \times 10^{-8} \text{ nm}$$

## 9. Wave Functions

- Matter waves of moving bodies can be represented by wave function  $\Psi$  depends on position & time
- $\Psi(x, y, z, t) = \Psi(x) \cdot \phi(y) \cdot \chi(z) \cdot \psi(t)$
- where  $\Psi(x), \phi(y), \chi(z)$  are orthogonal functions

### Wave Function & Probability Density

- Wave Function can be +, -, complex & can change with time
- $|\Psi|^2$  is called probability density which represents probability of finding the particle in unit volume of space



### Characteristics of accepted wave functions

- Finite, Continuous, Single Valued (FCS)
- Derivatives: Finite, Continuous, Single Valued (dFCS)
- Normalisable (Normalisable)

### Normalization of wave functions

- Total probability in the range where function is defined has to be unity
- $\int \Psi \times \Psi \, dx = 1$
- Normalisation condition is  $\int_{-\infty}^{\infty} \Psi \times \Psi \, dx = 1$

### State Function

- A well-behaved function satisfying FCS, dFCS & Normalisable are state functions
- $\Psi = A e^{i(\kappa x - \omega t)}$
- $K = \frac{P}{\hbar}, \lambda, \omega = \frac{E}{\hbar}$
- $\Psi = A e^{\frac{i}{\hbar}(px - Et)}$

### Double Slit Experiment

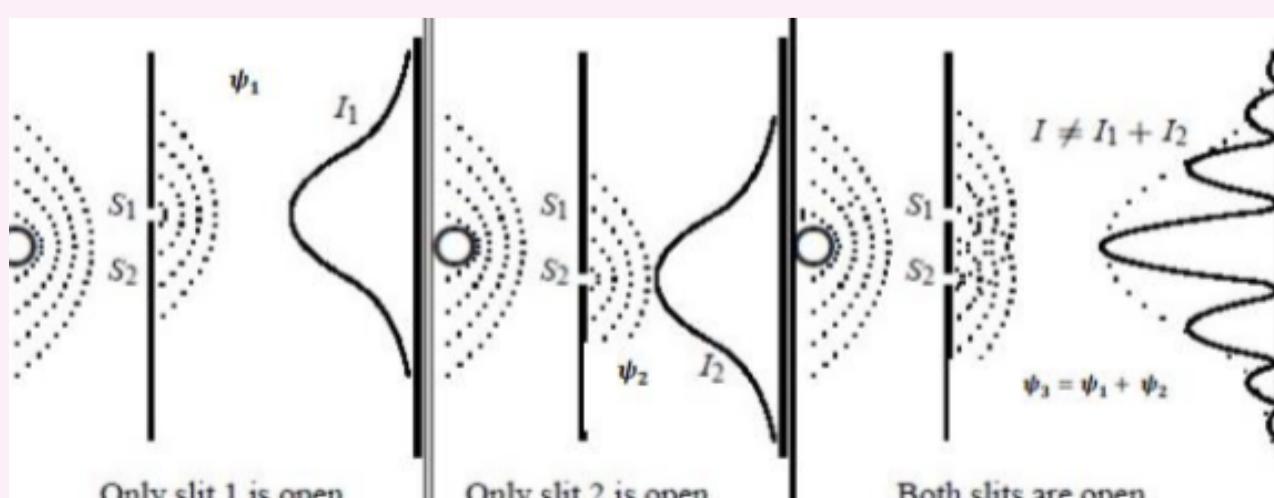
Let,  $\Psi_1$  &  $I_1 = |\Psi_1|^2$  be wave function & probability of photons in slit 1  
 $\Psi_2$  &  $I_2 = |\Psi_2|^2$  be wave function & probability of photons in slit 2

$\Psi_3 = \Psi_1 + \Psi_2$  (Superposed wave function for photons from both slits)

But,  $I_3 = |\Psi_3|^2 \neq I_1 + I_2$

$$I_3 = |\Psi_1|^2 + |\Psi_2|^2 + \Psi_1^* \Psi_2 + \Psi_1 \Psi_2^*$$

So,  $|\Psi_1 + \Psi_2|^2 = (\Psi_1^* + \Psi_2^*)(\Psi_1 + \Psi_2)$



## 10. Operators and Eigen value

### Eigen Value

→ A mathematical operator  $\hat{G}$  operating on wave function can result in eigen value  $G$  of the observable

$$\rightarrow \hat{G}\Psi = g\Psi$$

→ Operators arise because in quantum mechanics, a system is described with waves & not discrete particles

ex:  $e^{4x}$  is eigen function of the operator  $\frac{d^2}{dx^2}$ ,  
then, Eigen value equation:  $\hat{G}\Psi = G\Psi$   
 $\frac{d^2}{dx^2}(e^{4x}) = 16e^{4x}$

$$\text{Eigen value} : 16$$

$$\text{ex: } \Psi(x) = A\sin kx \quad \& \quad \hat{q} = i\hbar \frac{\partial}{\partial x}$$

$$\text{Then, } \hat{q}\Psi = q\Psi$$

$$i\hbar \frac{\partial}{\partial x}(A\sin kx) = i\hbar k A\cos kx \Rightarrow \text{NOT Eigen Function}$$

### Momentum Operator

→ Partial derivative of  $\Psi$  wrt position  
 $\frac{\partial \Psi}{\partial x} = \frac{\partial (Ae^{\frac{i}{\hbar}(px-Et)})}{\partial x} = \frac{iP}{\hbar} (Ae^{\frac{i}{\hbar}(px-Et)}) = \left(\frac{iP}{\hbar}\right)\Psi$

$$\Rightarrow \frac{i\hbar \partial \Psi}{\partial x} = P\Psi \quad \Rightarrow \left\{ -i\hbar \frac{\partial}{\partial x} \right\} \Psi = P\Psi$$

$$\rightarrow \text{Momentum operator} = -i\hbar \frac{\partial}{\partial x}$$

### Kinetic Energy Operator

→ 2nd derivative of  $\Psi$  wrt position

$$\frac{\partial^2 \Psi}{\partial x^2} = \frac{\partial^2}{\partial x^2} (Ae^{\frac{i}{\hbar}(px-Et)}) = \left(\frac{iP}{\hbar}\right)^2 \Psi$$

$$\Rightarrow \frac{-\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} = \frac{P^2}{2m} \Psi = KE\Psi$$

$$\rightarrow \text{KE operator} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$$

### Total Energy operator

→ Derivative of  $\Psi$  wrt time

$$\frac{\partial \Psi}{\partial t} = \frac{\partial (Ae^{\frac{i}{\hbar}(px-Et)})}{\partial t} = -\frac{iE}{\hbar} \Psi$$

$$\Rightarrow i\hbar \frac{\partial \Psi}{\partial t} = E\Psi$$

$$\rightarrow \text{TE operator} = i\hbar \frac{\partial}{\partial t} \Rightarrow \text{Also called Hamiltonian operator} \hat{H}$$

### Position Operator

→ Corresponds to position observable of a particle

$$\hat{x}\Psi = x\Psi$$

### Potential Energy Operator

→ It isn't explicitly described.

Eigen value can be expressed difference of TE & KE

$$\hat{V}\Psi = V\Psi$$

### Expectation values of observables

→ Expectation values  $\equiv$  avg. of repeated measurements of system

→ In general,

operator  $\hat{G}$  of observable  $g$

gives expectation value of observable  $\langle g \rangle = \frac{\int \Psi_x G\Psi dx}{\int \Psi_x \Psi dx}$

$$\Rightarrow \int \Psi_x \hat{G}\Psi dx = \int \Psi_x g\Psi dx = \langle g \rangle \int \Psi_x \Psi dx$$

$$\rightarrow \text{In 3D space, } \langle g \rangle = \frac{\int \Psi_x G\Psi dV}{\int \Psi_x \Psi dV}$$

### Time Independent Schrodinger Equation

$$\rightarrow \left[ \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + U(x) \right] \Psi(x, t) = i\hbar \frac{\partial}{\partial t} \Psi(x, t)$$

$$\begin{aligned} \Psi(x, t) &= A e^{\frac{i}{\hbar}(Px-Et)} \\ &= A e^{\frac{iPx}{\hbar}} \cdot e^{-\frac{iEt}{\hbar}} \\ &= \phi(x) \cdot e^{-\frac{iEt}{\hbar}} \end{aligned}$$

$$\text{Now, } \left[ \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + U(x) \right] \phi(x) \cdot e^{-\frac{iEt}{\hbar}} = i\hbar \frac{\partial}{\partial x} \phi(x) \cdot e^{-\frac{iEt}{\hbar}} \\ = \phi(x) \left\{ i\hbar \left( \frac{\partial}{\partial x} \right) E \cdot e^{-\frac{iEt}{\hbar}} \right\}$$

$$\text{So, } \left[ \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + U(x) \right] \phi(x) = E \phi(x)$$

$$\hat{H}\phi(x) = E\phi(x)$$

# U1 problems

- ? The electric field associated with an EM radiation (light) is given by,  
 $E(x, t) = 10^3 \cos(\omega t - \pi x \cdot 10^6 z)$   
 Evaluate  
 1. Speed of the Electric vector  
 2. Wavelength  
 3. Frequency  
 4. Period of the wave  
 5. Magnetic field associated with the wave  
 6. Direction of propagation of the magnetic transverse wave  
 7. Amplitude of the electric field vector  
 8. Amplitude and direction of the transverse magnetic wave

A. 1.  $V = 3 \times 10^8 \text{ m s}^{-1}$   
 2.  $\lambda = \frac{2\pi}{k} = \frac{2\pi}{3\pi \times 10^6} = \frac{2}{3} \times 10^{-6} \text{ m}$   
 3.  $\nu = \frac{c}{\lambda} = \frac{3 \times 10^8}{\frac{2}{3} \times 10^{-6}} = 4.5 \times 10^{14} \text{ Hz}$   
 4.  $T = \frac{1}{\nu} = 2.22 \times 10^{-15} \text{ s}$   
 5.  $B = \frac{E_0}{c} = \frac{10^3}{3 \times 10^8} = 0.33 \times 10^{-5}$   
 6.  $\perp \text{ to } E$  (in  $y$ -direction)  
 7.  $E_0 = 10^3$   
 8.  $B_0 = 3.33 \times 10^{-6}$  in  $y$ -direction

? The frequency of harmonic oscillator at 50C is  $6.2 \times 10^{12}$  per sec.  
 Estimate the average energy of the oscillator as per Planck's idea of cavity oscillator, also compare the same with classical average energy and average energy by R-J law.

A.  $T = 273 + 50 = 323 \text{ K}$   
 $\nu = 6.2 \times 10^{12} \text{ Hz}$

According to planck,  $E = \frac{h\nu}{e^{\frac{h\nu}{k_B T}} - 1} = 2.71 \times 10^{-21} \text{ J}$

According to R-J,  $E = k_B T = 4.46 \times 10^{-12} \text{ J}$

Difference =  $1.744 \times 10^{-21} \text{ J}$

? X-rays of wavelength 0.112 nm is scattered from a carbon target. Calculate the wavelength of X-rays scattered at an angle 90° with respect to the original direction. What is the energy lost by the X-ray photons? What is the energy gained by the electrons? If the incident x-ray retraces back what will be the shift?

$\lambda = 0.112 \times 10^{-9} \text{ m}$     $\lambda' = ?$    If  $\theta = 60^\circ$

$\Delta\lambda = \frac{h}{m_e c} (1 - \cos\theta) \Rightarrow \lambda' = \lambda + \frac{h}{m_e c} (1 - \frac{1}{2})$   
 $= 0.112 + 0.0012$   
 $= 1.1321 \text{ Å}$

Energy lost =  $\frac{hc}{\lambda} - \frac{hc}{\lambda'} = hc \left( \frac{1}{\lambda} - \frac{1}{\lambda'} \right)$   
 $= 1.89 \times 10^{-17} \text{ J}$   
 $= 118.31 \text{ eV}$

? Find the de Broglie wavelength of electrons moving with a speed of  $10^8 \text{ m/s}$

$$\lambda = \frac{h}{p} = \frac{h}{m_e v} = 7.27 \times 10^{-11} \text{ m}$$

? An alpha particle is accelerated through a potential difference of 1 kV. Find its de Broglie wavelength.

$$V = 1000 \text{ V} \quad \therefore \lambda = ?$$

$$\lambda = \frac{h}{p} = \frac{h}{\sqrt{2mKE}} = \frac{h}{\sqrt{2m_q V}}$$

$$m_\alpha = 4m_p \quad \& \quad q = 2e \quad V = 1000 \text{ V}$$

$$\lambda = \frac{h}{\sqrt{2 \times 4 \times 2m_p e \times 10^3}} = 4.525 \times 10^{-13} \text{ m}$$

? Compare the momenta and energy of an electron and photon whose de Broglie wavelength is 650nm

$$\lambda = 650 \text{ nm}$$

$$P = h/\lambda \quad E = \frac{p^2}{2m}$$

$$P_e = P_p = \frac{h}{650 \times 10^{-9}} = 1.02 \times 10^{-27} \text{ kg m s}^{-1}$$

$$K_e = \frac{p^2}{2m_e} = 5.71 \times 10^{-25} \text{ kg m s}^{-2}$$

$$K_p = \frac{hc}{\lambda p} = 3.056 \times 10^{-14} \text{ kg m s}^{-2}$$

$$\frac{K_e}{K_p} = 1.868 \times 10^{-6}$$

? Calculate the de Broglie wavelength of electrons and protons if their kinetic energies are i) 1% and ii) 5% of their rest mass energies.

A. rest mass energy of  $e^- = m_e c^2 = 8.2 \times 10^{-14} \text{ J}$   
 rest mass energy of photon =  $m_p c^2 = 1.5 \times 10^{-10} \text{ J}$

i)  $KE_e = 1\% \text{ of } E = \frac{1}{100} \times E = 8.2 \times 10^{-16} \text{ J}$   
 $KE_p = 1\% \text{ of } E = \frac{1}{100} \times E = 1.5 \times 10^{-12} \text{ J}$

$$\lambda_e = \frac{h}{\sqrt{2mKE_e}} = 1.72 \times 10^{-11} \text{ m} ; \lambda_p = \frac{h}{\sqrt{2mKE_p}} = 9.36 \times 10^{-15} \text{ m}$$

ii)  $KE_e = 5\% \text{ of } E = \frac{5}{100} \times E = 4.1 \times 10^{-15} \text{ J}$   
 $KE_p = 5\% \text{ of } E = \frac{5}{100} \times E = 7.5 \times 10^{-12} \text{ J}$

$$\lambda_e = \frac{h}{\sqrt{2mKE_e}} = 7.68 \times 10^{-12} \text{ m} ; \lambda_p = \frac{h}{\sqrt{2mKE_p}} = 4.18 \times 10^{-15} \text{ m}$$

? An electron and a photon have a wavelength of 2A. Calculate their momenta and total energies.

$$P_e = \frac{h}{\lambda} = \frac{h}{2 \times 10^{-10}} = 3.315 \times 10^{-24} \text{ kg m s}^{-1} = P_p$$

$$TE_e = m_e c^2 + \frac{p^2}{2m_e} \xrightarrow{\text{KE}} = 8.2 \times 10^{-14} \text{ J}$$

$$TE_p = \frac{hc}{\lambda} = 9.945 \times 10^{-16} \text{ kg m s}^{-2}$$

? The spectral line of Hg green is 546.1 nm has a width of 10-5 nm. Evaluate the minimum time spent by the electrons in the upper state before de excitation to the lower state.

$$\Delta t \cdot \Delta \lambda \geq \left| \frac{\lambda^2}{4\pi c} \right|$$

$$\Delta t \geq \left| \frac{\lambda^2}{4\pi c \Delta \lambda} \right| = \frac{(546.1 \times 10^{-9})^2}{4\pi c \times 10^{-5} \times 10^{-9}} = \frac{546.1 \times 546.1 \times 10^{-9}}{4\pi c \times 10^{-5}} = 7.91 \times 10^{-9} \text{ s}$$

? The uncertainty in the location of a particle is equal to its de Broglie wavelength. Show that the corresponding uncertainty in its velocity is approx one tenth of its velocity.

$$\Delta x = \lambda = \frac{h}{p}$$

$$p = mv$$

$$\Delta p = m \Delta v$$

$$\Delta x \cdot \Delta p \geq \frac{h}{4\pi}$$

$$\frac{K}{p} \cdot \Delta p \geq \frac{K}{4\pi} \Rightarrow \frac{m \Delta v}{p} \geq \frac{1}{4\pi} \Rightarrow \Delta v = \frac{v}{12.56} \approx \frac{v}{10}$$

? A wave packet is represented as,  $y = 10 \sin(30t - 40x) \cdot \cos(0.3t - 0.5x)$ . Find the phase and group velocities

$$v_p = \frac{\omega}{k} = \frac{30}{40} = 0.75$$

$$v_g = \frac{dw}{dk} = \frac{0.3}{0.5} = 0.6$$

? The speed of an electron is measured to be 1 km/s with an accuracy of 0.005%. Estimate the uncertainty in the position of the particle.

$$\Delta x \cdot \Delta p \geq \frac{h}{4\pi}$$

$$\Delta x \cdot m \Delta v \geq \frac{h}{4\pi}$$

$$\Delta x \cdot \Delta v \geq \frac{h}{4\pi m}$$

$$\Delta v = v \times \text{accuracy} = 1000 \text{ m/s} \times \frac{0.005}{100} = 0.05 \text{ m/s}$$

$$\Delta x = \frac{h}{4\pi m \Delta v} = \frac{6.63 \times 10^{-34} \text{ J s}}{4\pi \times 9.11 \times 10^{-31} \text{ kg} \times 0.05} = 1.159 \times 10^{-3} \text{ m}$$