Indirect dark matter detection Using Gamma-ray telescopes

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- ▶ Indirect detection via γ -rays
- Annihilation to standard model particles
- $\blacktriangleright \chi \chi \to b\bar{b} \to \gamma \gamma ...$
- $\chi \chi \to \tau^+ \tau^- \to \gamma \gamma \dots$
- ▶ DM map of the Universe

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Theoretical gamma-ray flux

$$I_{\gamma}(E,\Theta) = \frac{1}{2} \frac{\langle \sigma v \rangle}{m_{DM}^2} \frac{dN_{\gamma,ann}}{dE_{\gamma}} \frac{1}{4\pi} \int dI \rho_{DM}^2(r[I,\Theta])$$

- Particle physics factor
- Astrophysics factor

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$$\frac{1}{2} \frac{\langle \sigma v \rangle}{m_{DM}^2} \frac{dN_{\gamma,ann}}{dE_{\gamma}}$$

Relic abundance \sim annihilation rate:

$$\Omega_{DM}h^2\sim rac{3 imes 10^{-27}cm^3s^{-1}}{\langle\sigma v
angle}$$

Which implies:

$$\langle \sigma v \rangle \sim 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$$

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$$\frac{1}{2} \frac{\langle \sigma v \rangle}{m_{DM}^2} \frac{dN_{\gamma,ann}}{dE_{\gamma}}$$

Differential spectrum of emitted gamma-rays:

$$\frac{dN_{\gamma,ann}}{dE_{\gamma}}$$

Can be written as:

$$\frac{dN_{\gamma,ann}}{dE_{\gamma}} = \sum_{i} B_{i} \frac{dN_{\gamma,ann}^{i}}{dE_{\gamma}}$$

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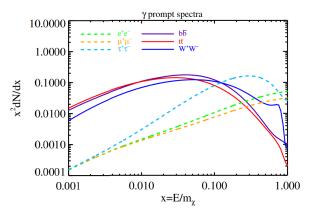


Figure: Figure from Gaskins et al. 2016: Gamma-ray spectrum from dark matter ($m_{DM} = 500\,GeV$) annihilation to six different final states, calculated using PPPC4DMID [1].

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$$\frac{1}{2} \frac{\langle \sigma v \rangle}{m_{DM}^2} \frac{dN_{\gamma,ann}}{dE_{\gamma}}$$

- ▶ Factor $\frac{1}{2}$ due to being its own anti particle.
- ▶ Factor $1/m_{DM}^2$ cancels the density squared integral.

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$$\frac{1}{4\pi}\int dl \rho_{DM}^2(r[l,\Theta])$$

J-factor definition:

$$J = rac{1}{\Delta\Omega} \int \int_{\Delta\Omega}
ho_{DM}^2(I,\Omega) dId\Omega$$

Normalization factor:

$$\int \!\! d\Omega = \int_0^{2\pi} \int_0^\pi \sin\theta d\theta d\phi = 4\pi$$

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Simple density function:

$$\rho_{NFW} = \frac{\rho_0}{\left(\frac{r}{r_{\rm s}}\right) \left[1 + \left(\frac{r}{r_{\rm s}}\right)\right]^2}$$

- R dependence
- Relate to line of sight integral

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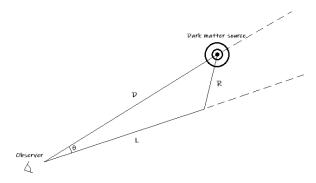


Figure: Top down view of observer looking down the line of sight with some angular separation Θ at a dark matter source.

$$R = \sqrt{L^2 + D^2 - 2L\cos\Theta}$$

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Theoretical gamma-ray flux - look back

$$I_{\gamma}(E,\Theta) = \frac{1}{2} \frac{\langle \sigma v \rangle}{m_{DM}^2} \frac{dN_{\gamma,ann}}{dE_{\gamma}} \frac{1}{4\pi} \int dI \rho_{DM}^2(r[I,\Theta])$$

- ▶ Rate of annihilation per volume: $\frac{1}{2} \frac{\langle \sigma v \rangle \rho_{DM}^2}{m_{DM}^2}$
- ▶ Differential spectrum of emitted gamma-rays: $\frac{dN_{\gamma,ann}}{dE_{\gamma}}$
- ► Volume: $\frac{1}{4\pi} \int dl$

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Sagittarius dwarf galaxy

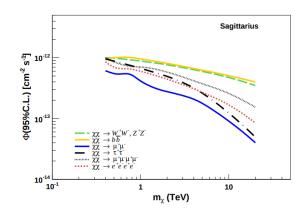


Figure: Figure from Abramowski et al. 2014: Gamma-ray flux as a function of dark matter mass (m_χ) under different annihilation channels.

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DM density profiles and observations

- ▶ What would DM-induced gamma-ray signals look like?
- Highly dependent on DM density profiles!
- Hard to constrain observationally
- Usually given as model input

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The NFW profile

- ▶ Navarro, Frenk White (1996)
- ► N-body simulation
- Showed that:

$$\rho_{NFW} = \frac{\rho_0}{\left(\frac{r}{r_s}\right) \left[1 + \left(\frac{r}{r_s}\right)\right]^2} \tag{1}$$

can be used to fit DM haloes in arbitrary cosmology.



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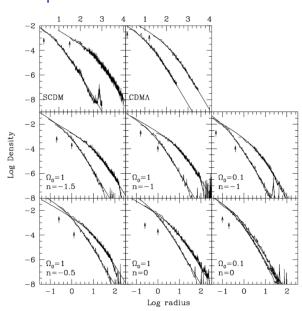
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The GNFW profile

- Generalized version of the NFW profile
- Steeper slopes close to galactic center
 - Adiabatic compression
 - Supernovae and AGN feedback
- ightharpoonup Arbitrary inner slope γ
- ▶ NFW for $\gamma = 1$

$$\rho_{GNFW} = \frac{\rho_0}{\left(\frac{r}{r_s}\right)^{\gamma} \left[1 + \left(\frac{r}{r_s}\right)\right]^{3-\gamma}} \tag{2}$$

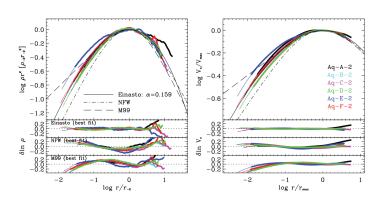
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► Radius-dependent slope

$$\rho_{Ein}(r) = \rho_0 \exp{-\frac{2}{a} \left[\left(\frac{r}{r_s} \right)^a - 1 \right]}$$
 (3)



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The Burkert profile

- Problem: spheroidal dwarf galaxies (dSphs)
 - Low density
 - Homogeneous
- ► Flattens towards the center

$$\rho_{Burk}(r) = \frac{\rho_0}{\left(1 + \frac{r}{r_s}\right) \left(1 + \left(\frac{r}{r_s}\right)^2\right)} \tag{4}$$



ESO/Digital Sky Survey 2

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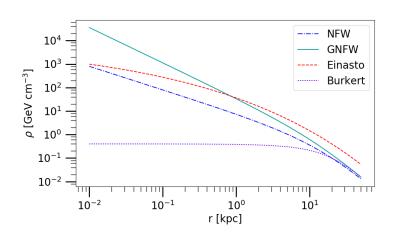
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Comparison



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Some questions

- ▶ Do observed γ -rays accurately reflect the DM density?
- Scattering processes?
- ► Background sources?

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Compton scattering

- Optical depth?
- Assuming $n_e \approx 10^0 \text{ cm}^{-3}$ (typical for ISM)
- ▶ Thomson approximation $\sigma_{K-N} \leq \sigma_{Th} \approx 10^{-24}$ cm
- Max distance d 20 kpc $\approx 10^{22}$ cm
- $\tau \approx n_e \sigma_{Th} d \approx 10^{-5}$
- ► → of little consequence

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Observational esults and

annihilation of two photons into particle-antiparticle

$$\gamma\gamma \longrightarrow e^+e^-$$
 (5)

► Suppose a DM-induced photon annihilates with a CMB photon. What energy should the photon minimally have for pair production?

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annihilation of two photons into particle-antiparticle

$$\gamma\gamma \longrightarrow e^+e^-$$
 (6)

- Suppose a DM-induced photon annihilates with a CMB photon. What energy should the photon minimally have for pair production?
- $ightharpoonup E_{\gamma} \geq 10^{14} \; \mathrm{eV}$
- Of litte consequence to Fermi-LAT
- ▶ Possibly important for CTA

Constraints from Fermi-LAT (2017)

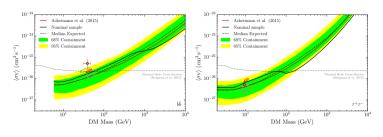


Figure: [2]

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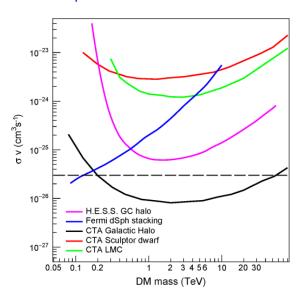
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Future telescopes - CTA



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Telescopes

Table: [4], [2], [5], [6], [7], [8], [3], [9], [10], [11], [12]

Name	Energy range	$\sigma v (cm^3/s)$
Fermi-LAT	20 MeV - 300 GeV	$\sim 4\cdot 10^{-26}$
H.E.S.S	300 GeV - 70 TeV	$4 \cdot 10^{-28}$
Veritas	85 GeV - 30 TeV	$2.85 \cdot 10^{-24}$
Magic	100 GeV - 100 TeV	$3.8 \cdot 10^{-24}$
HAWC	500 GeV - 100 TeV	$\sim 10^{-24}$
CTA	100 GeV - 10 TeV	$5 \cdot 10^{-27} - 3 \cdot 10^{-26}$
CALET	100 GeV - 1 TeV	_
GAMMA400	20 MeV - 1 TeV	_
HERD	10 GeV - 1 TeV	$ 6 \cdot 10^{-30} - 5 \cdot 10^{-29} $ $ 4 \cdot 10^{-27} - 9 \cdot 10^{-27} $

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