

1 Coulomb Scattering

2 Bhabha Scattering

$$e^-(p) + e^+(k) \rightarrow e^-(p') + e^+(k')$$

The t -channel matrix element:

$$\begin{aligned} -i\mathcal{M}_t &= [\bar{u}(p')(-ie\gamma^\mu)u(p)] \frac{-i\eta_{\mu\nu}}{(p-p')^2 + i\epsilon} [\bar{v}(k)(-ie\gamma^\nu)v(k')] \\ &= \frac{ie^2}{t} \bar{u}(p')\gamma^\mu u(p) \bar{v}(k)\gamma_\mu v(k') \end{aligned}$$

and the s -channel matrix element has a relative minus sign with the t -channel diagram:

$$\begin{aligned} i\mathcal{M}_s &= [\bar{v}(k)(-ie\gamma^\mu)u(p)] \frac{-i\eta_{\mu\nu}}{s^2 + i\epsilon} [\bar{u}(p')(-ie\gamma^\nu)v(k')] \\ &= \frac{ie^2}{s} \bar{v}(k)\gamma^\mu u(p) \bar{u}(p')\gamma_\mu v(k') \end{aligned}$$

The sum-squared matrix elements are:

$$\begin{aligned} |\mathcal{M}_t + \mathcal{M}_s|^2 &= |\mathcal{M}_t|^2 + 2\text{Re}[\mathcal{M}_t^\dagger \mathcal{M}_s] + |\mathcal{M}_s|^2 \\ \sum_{\text{spins}} |\mathcal{M}_t|^2 &= \frac{e^4}{t^2} \text{Tr} [\not{p}'\gamma^\mu \not{p}\gamma^\nu] \text{Tr} [\not{k}\gamma_\mu \not{k}'\gamma_\nu] \\ &= \frac{e^4}{t^2} 4(p'^\mu p^\nu + p'^\nu p^\mu - (p \cdot p')\eta^{\mu\nu}) 4(k_\mu k'_\nu + k_\nu k'_\mu - (k \cdot k')\eta_{\mu\nu}) \\ &= \frac{32e^4}{t^2} ((p' \cdot k')(p \cdot k) + (p' \cdot k)(p \cdot k')) \\ &= 8e^4 \frac{s^2 + u^2}{t^2} \\ \sum_{\text{spins}} |\mathcal{M}_s|^2 &= \frac{e^4}{s^2} \text{Tr} [(\not{k} - m)\gamma^\mu (\not{p} + m)\gamma^\nu] \text{Tr} [(\not{p}' + m)\gamma_\mu (\not{k}' - m)\gamma_\nu] \\ &= \frac{e^4}{s^2} 4[k^\mu p^\nu + k^\nu p^\mu - (k \cdot p)\eta^{\mu\nu}] 4[p'_\mu k'_\nu + p'_\nu k'_\mu - (p' \cdot k')\eta_{\mu\nu}] \\ &= \frac{32e^4}{s^2} [(k \cdot p')(p \cdot k') + (k \cdot k')(p \cdot p')] \\ &= 8e^4 \frac{u^2 + t^2}{s^2} \end{aligned}$$

$$\begin{aligned}
\sum_{spins} \mathcal{M}_t^\dagger \mathcal{M}_s &= -\frac{e^4}{ts} \text{Tr} [\bar{v}(k') \gamma_\mu v(k) \bar{v}(k) \gamma^\nu u(p) \bar{u}(p) \gamma^\mu u(p') \bar{u}(p') \gamma_\nu v(k')] \\
&= -\frac{e^4}{ts} \text{Tr} [k' \gamma_\mu k \gamma^\nu \not{p} \gamma^\mu \not{p}' \gamma_\nu] \\
&= -\frac{e^4}{ts} (-2) \text{Tr} [k' \not{p} \gamma^\nu k \not{p}' \gamma_\nu] \\
&= -\frac{e^4}{ts} (-8) (k \cdot p') \text{Tr} [k' \not{p}] \\
&= \frac{32e^4}{st} (k \cdot p') (k' \cdot p) \\
&= 8e^4 \frac{u^2}{st} \\
\frac{1}{4} \sum_{spins} |\mathcal{M}_s + \mathcal{M}_t|^2 &= 2e^4 \left[\frac{u^2 + t^2}{s^2} + 2 \frac{u^2}{st} + \frac{s^2 + u^2}{t^2} \right]
\end{aligned}$$

The differential cross-section is then

$$\begin{aligned}
\frac{d\sigma}{d\Omega} &= \frac{1}{(8\pi)^2} \frac{|\overline{\mathcal{M}}|^2}{s} \\
&= \frac{e^4}{32\pi^2 s} \left(u^2 \left(\frac{1}{s} + \frac{1}{t} \right)^2 + \frac{s^2}{t^2} + \frac{t^2}{s^2} \right) \\
&= \frac{\alpha^2}{2s} \left(u^2 \left(\frac{1}{s} + \frac{1}{t} \right)^2 + \frac{s^2}{t^2} + \frac{t^2}{s^2} \right)
\end{aligned}$$

3 Muon Decay

$$\mu^-(p) \rightarrow e^-(p') + \nu_\mu(k) + \bar{\nu}_e(k')$$

Approximating the propagator as $\sim -i/M_W^2$

$$\begin{aligned}
-i\mathcal{M} &= \left[\frac{g}{\sqrt{2}} \bar{u}_{\nu_\mu}(k) \gamma^\mu \frac{1}{2} (1 - \gamma^5) u_\mu(p) \right] \frac{-i}{M_W^2} \left[\frac{g}{\sqrt{2}} \bar{v}_{\nu_e}(k') \gamma_\mu \frac{1}{2} (1 - \gamma^5) u_e(p') \right] \\
&= \frac{-ig^2}{8M_W^2} \bar{u}_{\nu_\mu}(k) \gamma^\mu (1 - \gamma^5) u_\mu(p) \bar{v}_{\nu_e}(k') \gamma_\mu (1 - \gamma^5) u_e(p') \\
\sum_{spins} |\mathcal{M}|^2 &= \frac{G^2}{2} \text{Tr} [k \gamma^\mu (1 - \gamma^5) (\not{p} + m_\mu) \gamma^\nu (1 - \gamma^5)] \text{Tr} [k' \gamma_\mu (1 - \gamma^5) \not{p}' \gamma_\nu (1 - \gamma^5)]
\end{aligned}$$

The first trace is

$$\begin{aligned}
L_{muon}^{\mu\nu} &= \text{Tr} [k \gamma^\mu (1 - \gamma^5) \not{p} \gamma^\nu (1 - \gamma^5)] \\
&= \text{Tr} [k \gamma^\mu \not{p} \gamma^\nu] - \text{Tr} [k \gamma^\mu \gamma^5 \not{p} \gamma^\nu] - \text{Tr} [k \gamma^\mu \not{p} \gamma^\nu \gamma^5] + \text{Tr} [k \gamma^\mu \gamma^5 \not{p} \gamma^\nu \gamma^5] \\
&= 2\text{Tr} [k \gamma^\mu \not{p} \gamma^\nu] - 2\text{Tr} [k \gamma^\mu \gamma^5 \not{p} \gamma^\nu] \\
&= 8(k^\mu p^\nu + k^\nu p^\mu - (k \cdot p) \eta^{\mu\nu}) - 8i\epsilon^{\mu\nu\alpha\beta} k_\alpha p_\beta
\end{aligned}$$

The second is similarly

$$L_{elec}^{\mu\nu} = 8(k'^\mu p'^\nu + k'^\nu p'^\mu - (k' \cdot p') \eta^{\mu\nu} - i\epsilon^{\mu\nu\alpha\beta} k'_\alpha p'_\beta)$$

The amplitude squared is then:

$$\begin{aligned}
\sum_{spins} |\mathcal{M}|^2 &= 32G^2 (2(k \cdot k')(p \cdot p') + 2(k \cdot p')(k' \cdot p) - \epsilon^{\mu\nu\alpha\beta} \epsilon_{\mu\nu\rho\sigma} k_\alpha p_\beta k'^\rho p'^\sigma) \\
&= 64G^2 ((k \cdot k')(p \cdot p') + (k \cdot p')(k' \cdot p) - (\delta_\rho^\alpha \delta_\sigma^\beta - \delta_\sigma^\alpha \delta_\rho^\beta) k_\alpha p_\beta k'^\rho p'^\sigma) \\
&= 128G^2 (k \cdot p')(p \cdot k') \\
\frac{1}{2} \sum_{spins} |\mathcal{M}|^2 &= 64G^2 (k \cdot p')(p \cdot k')
\end{aligned}$$

The decay rate is then

$$\begin{aligned}
d\Gamma &= \frac{1}{2m_\mu} \overline{|\mathcal{M}|^2} (2\pi)^4 \delta^4(p - p' - k - k') \frac{d^3 p'}{(2\pi)^3 2E'} \frac{d^3 k}{(2\pi)^3 2\omega} \frac{d^3 k'}{(2\pi)^3 2\omega'} \\
&= \frac{1}{2m_\mu} \overline{|\mathcal{M}|^2} (2\pi)^4 \delta^4(p - p' - k - k') \frac{d^3 p'}{(2\pi)^3 2E'} \frac{d^3 k'}{(2\pi)^3 2\omega'} d^4 k \theta(\omega) \delta(k^2) \\
&= \frac{1}{2m_\mu} \overline{|\mathcal{M}|^2} \frac{1}{(2\pi)^5} \frac{d^3 p'}{2E'} \frac{d^3 k'}{2\omega'} \theta(m_\mu - E' - \omega') \delta((p - p' - k')^2)
\end{aligned}$$

We may use the 4-momentum conservation to rewrite the amplitude as

$$k \cdot p' = \frac{(k - p')^2 - k^2 - (p')^2}{2} = \frac{1}{2}(k' - p)^2 = \frac{1}{2}((\omega' - m_\mu)^2 - (\omega')^2) = \frac{1}{2}m_\mu(m_\mu - 2\omega')$$

$$p \cdot k' = \omega' m_\mu$$

$$\overline{|\mathcal{M}|^2} = 32G^2 \omega' m_\mu^2 (m_\mu - 2\omega')$$

$$\begin{aligned}
\Gamma &= \frac{m_\mu G^2}{2\pi^5} \int \omega' (m_\mu - 2\omega') \frac{d^3 p'}{2E'} \frac{d^3 k'}{2\omega'} \theta(m_\mu - E' - \omega') \delta(m_\mu^2 - 2m_\mu(E' + \omega') - 2p' \cdot k') \\
&= \frac{m_\mu G^2}{2\pi^5} \int \omega' (m_\mu - 2\omega') \frac{4\pi(E')^2 dE'}{2E'} \frac{2\pi(\omega')^2 d\omega' d\cos\theta}{2\omega'} \\
&\quad \times \theta(m_\mu - E' - \omega') \delta(m_\mu^2 - 2m_\mu(E' + \omega') + 2E'\omega'(1 - \cos\theta))
\end{aligned}$$

The integration of the δ -function of $\cos\theta$ is over the interval $[-1, 1]$, so we gain the conditions:

$$\begin{aligned}
0 &= m_\mu^2 - 2m_\mu(E' + \omega') + 2E'\omega'(1 - \cos\theta) \\
|\cos\theta| &= \left| 1 + \frac{m_\mu(m_\mu - 2E' - 2\omega')}{2E'\omega'} \right| \leq 1 \\
1 + \frac{m_\mu(m_\mu - 2E' - 2\omega')}{2E'\omega'} &\leq 1 \\
m_\mu - 2E' - 2\omega' &\leq 0 \\
\frac{1}{2}m_\mu - E' &\leq \omega' \\
1 + \frac{m_\mu(m_\mu - 2E' - 2\omega')}{2E'\omega'} &\geq -1 \\
m_\mu^2 - 2m_\mu E' - 2m_\mu \omega' &\geq -4E'\omega' \\
m_\mu(m_\mu - 2E') &\geq 2\omega'(m_\mu - 2E') \\
\omega' &\leq \frac{1}{2}m_\mu \quad \text{and} \quad 0 \leq E' \leq \frac{1}{2}m_\mu
\end{aligned}$$

Furthermore there's a $1/2E'\omega'$ factor from integrating the δ -function over $\cos \theta$:

$$\begin{aligned}
\Gamma &= \frac{m_\mu G^2}{2\pi^5} \int_0^{m_\mu/2} dE' \int_{\frac{1}{2}m_\mu - E'}^{m_\mu/2} d\omega' \frac{4\pi(E')^2}{2E'} \frac{2\pi(\omega')^2}{2\omega'} \omega' (m_\mu - 2\omega') \frac{1}{2E'\omega'} \\
&= \frac{m_\mu G^2}{2\pi^3} \int_0^{m_\mu/2} dE' \left[m_\mu \frac{\omega'^2}{2} - \frac{2\omega'^3}{3} \right]_{\frac{1}{2}m_\mu - E'}^{m_\mu/2} \\
&= \frac{m_\mu G^2}{2\pi^3} \int_0^{m_\mu/2} dE' \left[\frac{m_\mu^3}{8} - \frac{m_\mu^3}{12} - \frac{m_\mu}{2} \left(\frac{1}{2}m_\mu - E' \right)^2 + \frac{2}{3} \left(\frac{1}{2}m_\mu - E' \right)^3 \right] \\
&= \frac{m_\mu G^2}{2\pi^3} \int_0^{m_\mu/2} dE' \left[\frac{1}{2}m_\mu E'^2 - \frac{2}{3}E'^3 \right] \\
&= \frac{m_\mu G^2}{2\pi^3} \left[\frac{1}{6}m_\mu E'^3 - \frac{1}{6}E'^4 \right]_0^{m_\mu/2} \\
&= \frac{m_\mu^5 G^2}{196\pi^3} \approx 3.0019 \times 10^{-19} \text{ GeV} = \frac{\hbar}{2.187 \text{ } \mu\text{s}}
\end{aligned}$$

The actual mean lifetime is $2.2 \text{ } \mu\text{s}$, so first order is pretty close!