

1 Free Particle Wavefunctions

1.1 Non-Relativistic Particles

Schrodinger equation

$$0 = i\partial_t\psi - H\psi = \left(i\partial_t + \frac{\nabla^2}{2m}\right)\psi$$

Solution, where $E = \mathbf{p}^2/2m$:

$$\psi_p = Ne^{-iEt+i\mathbf{p}\cdot\mathbf{x}}$$

Density and current

	ρ	\mathbf{j}
Equation	$ \psi ^2$	$\frac{1}{2mi}(\psi^*\nabla\psi - \psi\nabla\psi^*)$
Free	$ N ^2$	$\frac{\mathbf{p}}{m} N ^2$

1.2 Scalar Field Equation

Klein-Gordon Equation

$$(-\partial_t^2 + \nabla^2)\phi = m^2\phi$$

Solution, where $p^0 = E = \sqrt{\mathbf{p}^2 + m^2}$:

$$\phi = Ne^{-iEt+i\mathbf{p}\cdot\mathbf{x}} = Ne^{-ip\cdot x}$$

Density and current

	ρ	\mathbf{j}	j^μ
Equation	$i(\phi^*\partial_t\phi - \phi\partial_t\phi^*)$	$-i(\phi^*\nabla\phi - \phi\nabla\phi^*)$	$i(\phi^*\partial^\mu\phi - \phi\partial^\mu\phi^*)$
Free	$2E N ^2$	$2\mathbf{p} N ^2$	$2p^\mu N ^2$

Antiparticles (Feynman-Stuckelberg Interpretation)

$$j^\mu(\phi, -p) = 2p^\mu|N|^2 = j^\mu(\phi^*, p)$$

$$\phi^* = Ne^{-i(-p)\cdot x} = Ne^{ip\cdot x}$$

Identify antiparticles ϕ^* as having opposite charge in the current. Antiparticles have positive energy and can be used instead of equivalent negative energy solutions.

1.3 Electrodynamics

Minimal coupling prescription

$$i\partial^\mu \rightarrow i\partial^\mu + eA^\mu$$

1.4 Scalar QED

Perturbative potential

$$V = -ie(\partial_\mu A^\mu + A^\mu \partial_\mu) - e^2 A^2$$

Electromagnetic current between states $i \rightarrow f$

$$j_\mu^{fi} = -ie(\phi_f^* \partial_\mu \phi_i - \phi_i \partial_\mu \phi_f^*) = -e N_i N_f (p_i + p_f)_\mu e^{i(p_f - p_i) \cdot x}$$

Photon equation of motion and solution (from Maxwell's equations)

$$\begin{aligned} (-\partial_t^2 + \nabla^2) A^\mu &= j^\mu \\ A^\mu &= -\frac{1}{q^2} j^\mu \\ &= \frac{-ig^{\mu\nu}}{q^2} j_\nu \end{aligned}$$

Photon propagator

$$\frac{-ig^{\mu\nu}}{q^2}$$

2 Perturbation Theory

State transistion amplitude $i \rightarrow f$:

$$\begin{aligned} V_{fi}(t) &= \int d^3x \phi_f^*(\mathbf{x}, t) V(\mathbf{x}, t) \phi_i(\mathbf{x}, t) \\ T_{fi} &= -i \int d^4x \phi_f^*(x) V(x) \phi_i(x) \\ &= -i \int_{-\infty}^{\infty} V_{fi}(t) dt \end{aligned}$$

Fermi's Golden Rule (assuming V is time-independent)

$$\begin{aligned} W_{fi} &= \int \rho(E_f) dE_f \lim_{T \rightarrow \infty} \frac{|T_{fi}|^2}{T} \\ &= 2\pi |V_{fi}|^2 \int \rho(E_f) dE_f \delta(E_f - E_i) \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} dt \\ &= 2\pi |V_{fi}|^2 \rho(E_i) \end{aligned}$$

State transition amplitude for $2 \rightarrow 2$ scattering (scalar QED):

$$\begin{aligned} T_{fi} &= -i \int j_\mu^{(1)} A^\mu d^4x \\ &= -i \int j_\mu^{(1)} \frac{-ig^{\mu\nu}}{q^2} j_\nu^{(2)} d^4x \\ &= N_A N_B N_C N_D (2\pi)^4 \delta(p_A + p_B - p_C - p_D) (-i\mathcal{M}) \\ -i\mathcal{M} &= ie(p_A + p_C)^\mu \left(-i \frac{g_{\mu\nu}}{q^2} \right) ie(p_B + p_D)^\nu \end{aligned}$$

3 Cross-Section

The transition matrix element is

$$T_{n \rightarrow m} = -i \langle m | \int_{-\infty}^{\infty} H_I(t) dt | n \rangle = -i \langle m | H_{int} | n \rangle 2\pi\delta(\omega)$$

The probability is then

$$P_{n \rightarrow m} = |T_{n \rightarrow m}|^2 = |\langle m | H_{int} | n \rangle|^2 (2\pi)^2 \delta(\omega)^2$$

To compute the squared delta function, rewrite it as a Fourier transform over some finite time T and use the first delta function to set the second:

$$\begin{aligned} (2\pi)^2 \delta(\omega)^2 &= 2\pi\delta(\omega) \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} e^{-i\omega t} dt \\ &= 2\pi\delta(\omega) \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} dt \\ &= 2\pi\delta(\omega) \lim_{T \rightarrow \infty} T \end{aligned}$$

This is obviously divergent, so instead we take the probability rate

$$W_{n \rightarrow m} = \lim_{T \rightarrow \infty} \frac{P_{n \rightarrow m}}{T} = |\langle m | H_{int} | n \rangle|^2 2\pi\delta(\omega)$$