

Q&A (2.61-2.70)

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Exercise 2.61: Calculate the probability of obtaining the result +1 for a measurement $\vec{v} \cdot \vec{\sigma}$, given that the state prior to measurement is $|0\rangle$. What is the state of the system after the measurement if +1 is obtained?

Answer:

$$\begin{aligned} p(+1) &= \langle 0 | \lambda_1 \rangle \langle \lambda_1 | 0 \rangle \\ &= \frac{1}{2} \langle 0 | (I + \vec{v} \cdot \vec{\sigma}) | 0 \rangle \\ &= \frac{1}{2} \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{bmatrix} 1 + v_3 & v_1 - iv_2 \\ v_1 + iv_2 & 1 - v_3 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \frac{1}{2} (1 + v_3). \end{aligned} \tag{1}$$

The state of the quantum system immediately after the measurement is

$$\begin{aligned} \frac{|\lambda_1\rangle \langle \lambda_1 | 0 \rangle}{\sqrt{p(+1)}} &= \frac{|\lambda_1\rangle \langle \lambda_1 | 0 \rangle}{\sqrt{\frac{1}{2}(1 + v_3)}} \\ &= \frac{1}{\sqrt{\frac{1}{2}(1 + v_3)}} * \frac{1}{2} \begin{bmatrix} 1 + v_3 & v_1 - iv_2 \\ v_1 + iv_2 & 1 - v_3 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\ &= \frac{1}{\sqrt{\frac{1}{2}(1 + v_3)}} \frac{1}{2} \begin{bmatrix} 1 + v_3 \\ v_1 + iv_2 \end{bmatrix} \\ &= \sqrt{\frac{1 + v_3}{2}} \begin{bmatrix} 1 \\ \frac{v_1 + iv_2}{1 + v_3} \end{bmatrix}. \end{aligned} \tag{2}$$

Exercise 2.62: Show that any measurement where the measurement operators and the POVM elements coincide is a projective measurement.

Answer:

Exercise 2.53: What are the eigenvalues and eigenvectors of H ?

Answer:

$$\det|A - \lambda I| = \det \left| \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} - \begin{bmatrix} \lambda & 0 \\ 0 & \lambda \end{bmatrix} \right| = \begin{vmatrix} \frac{1}{\sqrt{2}} - \lambda & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} - \lambda \end{vmatrix} = 0$$

Eigenvalues are ± 1 and associated eigenvectors are $|\lambda_{\pm}\rangle = \frac{1}{\sqrt{4 \pm 2\sqrt{2}}} \begin{bmatrix} 1 \\ -1 \pm \sqrt{2} \end{bmatrix}$.

Exercise 2.54: Suppose A and B are commuting Hermitian operators. Prove that $\exp(A)\exp(B) = \exp(A+B)$. (Hint: Use the results of Section 2.1.9.)

Answer:

Since $[A, B] = 0$, A and B are simultaneously diagonalize, $A = \sum_i a_i |i\rangle \langle i|$, $B = \sum_j b_j |j\rangle \langle j|$.

$$\begin{aligned} \exp(A)\exp(B) &= \sum_i \exp(a_i) |i\rangle \langle i| B = \sum_j \exp(b_j) |j\rangle \langle j| \\ &= \sum_{i,j} \exp(a_i + b_j) |i\rangle \langle i| j\rangle \langle j| \\ &= \sum_{i,j} \exp(a_i + b_j) |i\rangle \langle j| \delta_{i,j} \\ &= \sum_i \exp(a_i + b_i) |i\rangle \langle i| \\ &= \exp(A+B). \end{aligned} \tag{3}$$

Exercise 2.55: Prove that $U(t_1, t_2)$ defined in Equation (2.91) is unitary.

Answer:

$$\begin{aligned} U(t_1, t_2)^\dagger U(t_1, t_2) &= \exp \left(\frac{iH(t_2 - t_1)}{\hbar} \right) \exp \left(\frac{-iH(t_2 - t_1)}{\hbar} \right) \\ &= \exp \left(\frac{i \sum_{E_1} E_1 |E_1\rangle \langle E_1| (t_2 - t_1)}{\hbar} \right) \exp \left(\frac{-i \sum_{E_2} E_2 |E_2\rangle \langle E_2| (t_2 - t_1)}{\hbar} \right) \\ &= \sum_{E_1, E_2} \exp \left(\frac{iE_1(t_2 - t_1)}{\hbar} \right) |E_2\rangle \langle E_2| \exp \left(\frac{-iE_2(t_2 - t_1)}{\hbar} \right) |E_1\rangle \langle E_1| \\ &= \sum_{E_1, E_2} \exp \left(\frac{i(E_1 - E_2)(t_2 - t_1)}{\hbar} \right) |E_1\rangle \langle E_1| E_2\rangle \langle E_2| \\ &= \sum_{E_1, E_2} \exp \left(\frac{i(E_1 - E_2)(t_2 - t_1)}{\hbar} \right) |E_1\rangle \langle E_2| \delta_{E_1, E_2} \\ &= \sum_{E_1} |E_1\rangle \langle E_1| \\ &= I. \end{aligned} \tag{4}$$

Thus $U(t_1, t_2)$ is unitary.

Exercise 2.56: Use the spectral decomposition to show that $K - i \log(U)$ is Hermitian for any unitary U , and thus $U = \exp(iK)$ for some Hermitian K .

Answer:

Since U is unitary, then U can perform spectral decomposition, $U = \sum_i \lambda_i |i\rangle \langle i|$

$$\begin{aligned}
 K^\dagger &= (-i \log(U))^\dagger \\
 &= (-i \log \left(\sum_i \lambda_i |i\rangle \langle i| \right))^\dagger \\
 &= (i \sum_i \log(\lambda_i) |i\rangle \langle i|).
 \end{aligned} \tag{5}$$

Exercise 2.57: (Cascaded measurements are single measurements) Suppose L_l and M_m are two sets of measurement operators. Show that a measurement defined by the measurement operators L_l followed by a measurement defined by the measurement operators M_m is physically equivalent to a single measurement defined by measurement operators N_{lm} with the representation $N_{lm} = M_m L_l$.

Answer:

If the state of the quantum system is $|\psi\rangle$ immediately before the measurement. The state of the system after the first measurement is $|\psi_L\rangle = \frac{L_l |\psi\rangle}{\sqrt{\langle \psi | L_l^\dagger L_l | \psi \rangle}}$ and the second measurement is $|\psi_M\rangle = \frac{M_m |\psi_L\rangle}{\sqrt{\langle \psi_L | M_m^\dagger M_m | \psi_L \rangle}}$.

$$\begin{aligned}
 \langle \psi_L | &= \frac{\langle \psi | L_l^\dagger}{\sqrt{\langle \psi | L_l^\dagger L_l | \psi \rangle}} \\
 |\psi_M\rangle &= \frac{M_m |\psi_L\rangle}{\sqrt{\langle \psi_L | M_m^\dagger M_m | \psi_L \rangle}} \\
 &= \frac{M_m \frac{L_l |\psi\rangle}{\sqrt{\langle \psi | L_l^\dagger L_l | \psi \rangle}}}{\sqrt{\langle \psi_L | M_m^\dagger M_m | \psi_L \rangle}} \\
 &= \frac{\frac{\langle \psi | L_l^\dagger}{\sqrt{\langle \psi | L_l^\dagger L_l | \psi \rangle}} M_m^\dagger M_m \frac{L_l |\psi\rangle}{\sqrt{\langle \psi | L_l^\dagger L_l | \psi \rangle}}}{\sqrt{\langle \psi_L | M_m^\dagger M_m | \psi_L \rangle}} \\
 &= \frac{M_m L_l |\psi\rangle}{\sqrt{\langle \psi | L_l^\dagger M_m^\dagger M_m L_l | \psi \rangle}}.
 \end{aligned} \tag{6}$$

The state of the system after the measurement operators N_{lm} ($N_{lm} = M_m L_l$) is

$$\begin{aligned}
 |\psi_N\rangle &= \frac{N_{lm} |\psi\rangle}{\sqrt{\langle \psi | N_{lm}^\dagger N_{lm} | \psi \rangle}} \\
 &= \frac{M_m L_l |\psi\rangle}{\sqrt{\langle \psi | L_l^\dagger M_m^\dagger M_m L_l | \psi \rangle}} = |\psi_M\rangle.
 \end{aligned} \tag{7}$$

Thus we proved that Cascaded measurements are single measurements.

Exercise 2.58: Suppose we prepare a quantum system in an eigenstate $|\psi\rangle$ of some observable M , with corresponding eigenvalue m . What is the average observed value of M , and the standard deviation?

Answer:

$$\begin{aligned}
\langle M \rangle &= \langle \psi | M | \psi \rangle \\
&= \langle \psi | m | \psi \rangle \\
&= m \langle \psi | \psi \rangle \\
&= m [\Delta M]^2 = \langle M^2 \rangle - \langle M \rangle^2 \\
&= \langle \psi | m^2 | \psi \rangle - m^2 \\
&= m^2 - m^2 \\
&= 0.
\end{aligned} \tag{8}$$

Exercise 2.49: Suppose we have qubit in the state $|0\rangle$, and we measure the observable X . What is the average value of X ? What is the standard deviation of X ?

Answer:

$$\begin{aligned}
\langle X \rangle &= \langle 0 | X | 0 \rangle \\
&= 0 \\
\langle X^2 \rangle &= \langle 0 | X^2 | 0 \rangle \\
&= 1 \\
[\Delta X] &= \sqrt{\langle X^2 \rangle - \langle X \rangle^2} = 1.
\end{aligned} \tag{9}$$

Exercise 2.50: Show that $\vec{v} \cdot \vec{\sigma}$ has eigenvalues ± 1 , and that the projectors onto the corresponding eigenspaces are given by $P_{\pm} = (I \pm \vec{v} \cdot \vec{\sigma})/2$.

Answer:

$$\begin{aligned}
\vec{v} \cdot \vec{\sigma} &= \sum_{i=1}^3 v_i \sigma_i \\
&= v_1 \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \\
&= v_2 \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \\
&= v_3 \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \\
&= \begin{bmatrix} v_3 & v_1 - iv_2 \\ v_1 + iv_2 & -v_3 \end{bmatrix} \\
\det(\vec{v} \cdot \vec{\sigma} - \lambda I) &= (v_3 - \lambda)(-v_3 - \lambda) - (v_1 - iv_2)(v_1 + iv_2) \\
&= \lambda^2 - (v_1^2 + v_2^2 + v_3^2) \\
&= \lambda^2 - 1.
\end{aligned} \tag{10}$$

Eigenvalues are ± 1 . 未完