

# Pouring a Fizzy Drink

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## 1 Pouring a Fizzy Drink

We will work entirely in terms of height, assuming the volume is made up of circular disks of infinitesimal height stacked upon each other. Let's suppose we have a cylindrical glass of height  $H_g \in \mathbb{R}_{>0}$ . Let the height of the drink be denoted by  $h : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ . Let the foam that is formed by the pouring process be denoted by  $b : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ .

Now suppose we wished to fill this glass up from empty to  $H_0 \in (0, H_g)$  without the foam spilling out of the top of the glass, *assuming the foam is always on top of the drink*. Then we would require for all times  $t \in \mathbb{R}_{\geq 0}$  that the height of the foam and drink are less than the height of the glass,

$$0 \leq h(t) + b(t) \leq H_g \quad (1.1)$$

along with the initial conditions

$$h(0) = b(0) = 0. \quad (1.2)$$

Whenever the foam would spill over the glass, we reduce the foam down to  $b = H_g - h$  to simulate overflowing. A diagram of this is shown in Figure 1. We can choose to pour the drink into the glass however we wish, given by some pouring function  $p : \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$  in the form

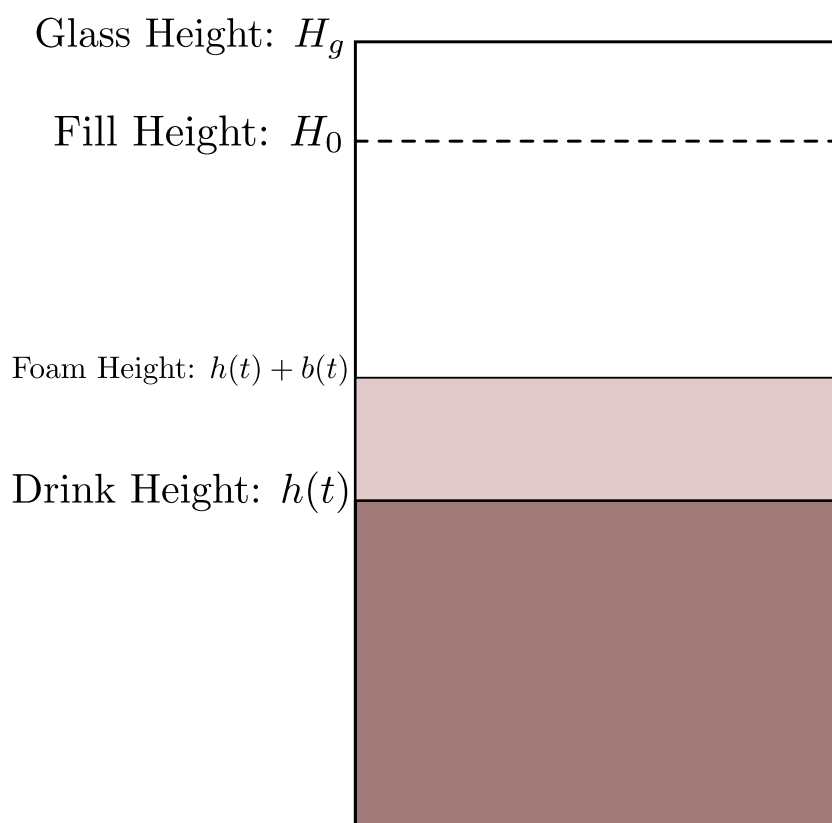
$$\frac{dh}{dt} = p(t). \quad (1.3)$$

The foam generated when the drink is poured is proportional to the amount being poured in, yielding a term

$$\text{Foam in : } \mu \frac{dh}{dt} = \mu p(t) \quad (1.4)$$

where  $\mu \in \mathbb{R}^+$  is a dimensionless proportionality constant. The foam "evaporates", leaving the glass, according to some function  $q : \mathbb{R}_{\geq 0}^2 \rightarrow \mathbb{R}_{\geq 0}$ , so that  $b$  satisfies the differential equation

$$\frac{db}{dt} = \mu p(t) - q(b). \quad (1.5)$$



**Figure 1:** Diagram of the glass, partially filled with drink and foam.

### 1.1 Another example

We could try

$$p(t) = |\cos(\ln(1 + t/3))|, \quad (1.6)$$

$$q(t, b) = \sqrt{b} + b^{10}. \quad (1.7)$$

The resulting solution is shown in Figure 2.

Again, the  $p, q$  graphs are shown in Figures 3a and 3b

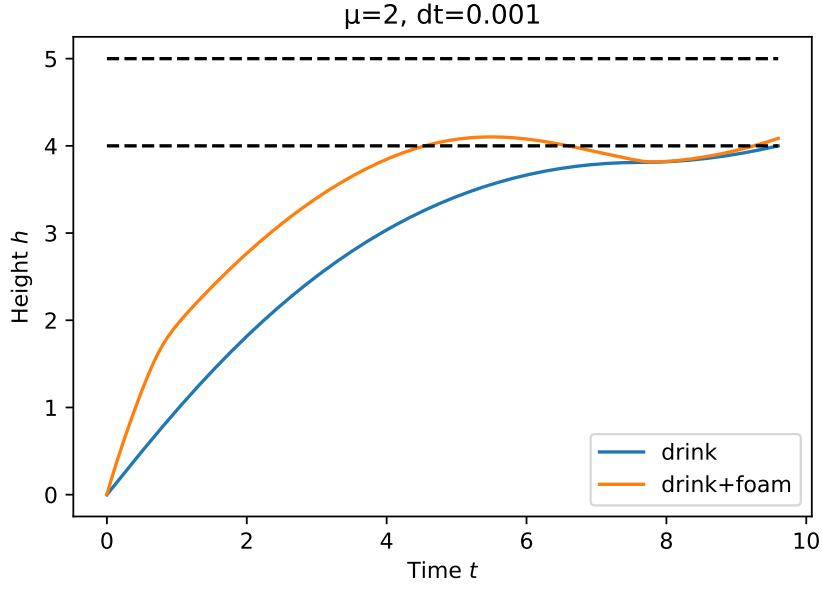
### 1.2 An exact form

Let's suppose that

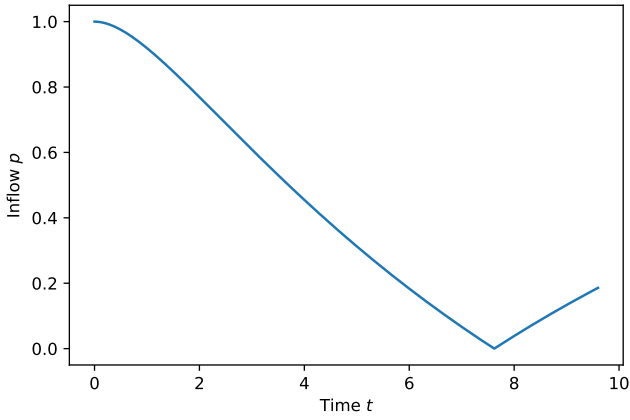
$$q(t, b) = kb(t) \quad (1.8)$$

for some constant  $k > 0$ .

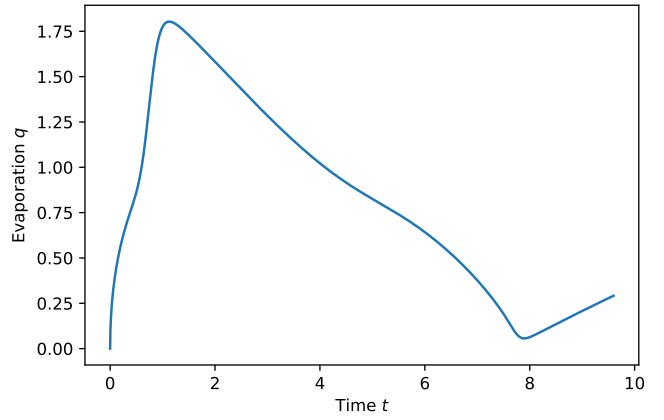
Then we have



**Figure 2:** Drink and foam height as a function of time. The orange line is the foam height, on top of the drink shown in blue. The time taken to fill the drink to the required height is  $\approx 9.599$  units.



(a) Inflow function  $p(t)$ .



(b) Evaporation function  $q(t, b(t))$ .

**Figure 3:** Derivative functions.

$$\frac{db}{dt} = \mu p(t) - kb(t). \quad (1.9)$$

Multiplying by an integrating factor  $e^{kt}$ , we rearrange to

$$\frac{d}{dt} \left( e^{kt} b(t) \right) = \mu e^{kt} p(t). \quad (1.10)$$

Integrating from 0 to  $t$  yields

$$e^{kt}b(t) - \underbrace{b(0)}_{=0} = \mu \int_0^t e^{ky} p(y) dy \quad (1.11)$$

so that  $b$  is given by

$$b(t) = \mu e^{-kt} \int_0^t e^{ky} p(y) dy. \quad (1.12)$$

The spill condition (1.1) is then

$$0 \leq \int_0^t p(y) dy + \mu e^{-kt} \int_0^t e^{ky} p(y) dy \leq H_g. \quad (1.13)$$

Now, at the final time  $t_f$  we have

$$-H_0 \leq \mu e^{-kt_f} \int_0^{t_f} e^{ky} p(y) dy \leq H_g - H_0. \quad (1.14)$$

The integrand is positive, so we can replace the lower bound with 0,

$$0 \leq \mu e^{-kt_f} \int_0^{t_f} e^{ky} p(y) dy \leq H_g - H_0. \quad (1.15)$$

### 1.2.1 Specific example

Let  $H_g = 1$  so that  $H_0 \in (0, 1)$ . We let

$$p(t) = |\cos(t)| \quad (1.16)$$

so that

$$h(t) = \sin(t) \quad (1.17)$$

omitting the "sgn(cos(t)) factor", which means that the final time occurs at

$$t_f = \sin^{-1}(H_0). \quad (1.18)$$

Then, using equation (1.12),  $b$  is given by

$$b(t) = \frac{\mu(k \cos(t) + \sin(t))}{1 + k^2} - \frac{\mu k e^{-kt}}{1 + k^2}. \quad (1.19)$$

At the final time  $t_f$  we get

$$b(t_f) = \frac{\mu(k \sqrt{1 - H_0^2} + H_0)}{1 + k^2} - \frac{\mu k e^{-k \sin^{-1}(H_0)}}{1 + k^2}. \quad (1.20)$$

For this system to abide by the spill condition (due to  $b$  and  $h$  monotonically increasing on  $(0, t_f)$ ), we require  $\mu, k, H_0$  to satisfy

$$0 \leq H_0 (1 + \mu + k^2) + \mu k \sqrt{1 - H_0^2} - \mu k e^{-k \sin^{-1}(H_0)} \leq H_g (1 + k^2) \quad (1.21)$$

## 2 Minimisation

For example, let  $\mu = 2, k = 25, H_g = 1, H_0 = 0.8$ .

We will once again use the foam function

$$q(b) = kb. \quad (2.1)$$

### 2.1 Constant input

Suppose we tried the pouring function

$$p(t) = c \quad (2.2)$$

for some  $c > 0$ . Then

$$h(t) = ct \quad (2.3)$$

and

$$b(t) = \mu e^{-kt} \int_0^t e^{ky} c \, dy = \frac{\mu c}{k} (1 - e^{-kt}). \quad (2.4)$$

Now

$$(h + b)(t) = c \left[ t + \frac{\mu}{k} (1 - e^{-kt}) \right]. \quad (2.5)$$

The finishing time is given by

$$t_f = H_0/c. \quad (2.6)$$

The derivative of  $h + b$  is given by

$$\frac{d(h + b)}{dt} = c (1 + \mu e^{-kt}) \geq 0 \quad (2.7)$$

so the only maximum that can occur is at time  $t_f$ . That means that the drink doesn't spill if and only if

$$(h + b)(t_f) = c \left[ t_f + \frac{\mu}{k} (1 - e^{-kt_f}) \right] \leq H_g. \quad (2.8)$$

That is

$$(h + b)(t_f) = \left[ H_0 + \frac{\mu c}{k} (1 - e^{-kH_0/c}) \right] \leq H_g \quad (2.9)$$

which rearranges to

$$c (1 - e^{-kH_0/c}) \leq \frac{k}{\mu} (H_g - H_0). \quad (2.10)$$

This is an increasing function, and since we want to minimise  $t_f = H_0/c$  we need to choose the largest allowed  $c$  value. We set

$$c (1 - e^{-kH_0/c}) = \frac{k}{\mu} (H_g - H_0). \quad (2.11)$$

Now, the above has an upper bound as  $c \rightarrow \infty$ , given by  $kH_0$ .

## 2.2 Harder

For an example of trying to find a minimum time, we will use a trial function

$$p(t) = \alpha t \quad (2.12)$$

for some constant  $\alpha \in \mathbb{R}^+$ .

Now,

$$h(t) = \int_0^t p(y) dy = \int_0^t \alpha y dt = \frac{1}{2} \alpha t^2. \quad (2.13)$$

Using (1.12) we have

$$b(t) = \mu e^{-kt} \int_0^t e^{ky} (\alpha y) dy \quad (2.14)$$

$$= \alpha \gamma \left( \frac{e^{-kt}}{k} + t - \frac{1}{k} \right) \quad (2.15)$$

where  $\gamma = \mu/k$ .

This means that the height of the foam at time  $t$  is given by

$$(h + b)(t) = \alpha \left[ \frac{1}{2} t^2 + \gamma \left( \frac{e^{-kt}}{k} + t - \frac{1}{k} \right) \right]. \quad (2.16)$$

This means we need to find  $\alpha \in \mathbb{R}_0^+$  to minimise  $t_f$  where

$$h(t_f) = \frac{1}{2} \alpha t_f^2 = H_0 \quad (2.17)$$

subject to

$$(h + b)(t) = \alpha \left[ \frac{1}{2} t^2 + \gamma \left( \frac{e^{-kt}}{k} + t - \frac{1}{k} \right) \right] \leq H_g \quad (2.18)$$

for all time  $t \in (0, t_f)$ . Furthermore,

$$t_f = \sqrt{\frac{2H_0}{\alpha}}. \quad (2.19)$$

Now, seeking an upper bound on  $(h + b)$ , the derivative of  $(h + b)$  is given by

$$\frac{d(h + b)}{dt} = \alpha t + \alpha(1 - e^{-kt}) = 0, \quad (2.20)$$

which reduces to

$$(t + 1)e^{kt} = 1. \quad (2.21)$$

The solution to this is given by

$$\hat{t} = \frac{W(e^k k) - k}{k} \quad (2.22)$$

where  $W$  is the Lambert  $W$  function. By the extreme value theorem,  $h+b$  takes its maximum, and we deduce that this at either  $\hat{t}$  or  $t_f$ . For the foam to not spill, we need

$$\max \left( (h+b)(\hat{t}), (h+b)(t_f) \right) \leq H_g. \quad (2.23)$$

### 2.2.1 Case 1

Suppose that

$$t_f = \sqrt{\frac{2H_0}{\alpha}} \leq \frac{W(e^k k) - k}{k} = \hat{t}. \quad (2.24)$$