

Session 10 of group theory

Ladies and gentlemen, it's now time to do some physics!

First glance of symmetry in particle physics (Gell-Mann's genius.)

$m_p \approx 938 \frac{\text{MeV}}{c^2}$	$m_{\pi^+} \approx 140 \frac{\text{MeV}}{c^2}$	$m_{\Delta^{++}} \approx 1231 \text{ MeV}/c^2$
	$m_{\pi^0} \approx 135 \text{ MeV}/c^2$	$m_{\Delta^+} \approx 1232 \text{ MeV}/c^2$
$m_n \approx 939 \frac{\text{MeV}}{c^2}$	$m_{\pi^-} \approx 140 \frac{\text{MeV}}{c^2}$	$m_{\Delta^0} \approx 1233 \text{ MeV}/c^2$
		$m_{\Delta^-} \approx 1235 \text{ MeV}/c^2$

The notion of "isospin" was first introduced by Heisenberg in 1932. He was puzzled by the fact that p 's and n 's have nearly equal masses, and apart from difference in electrical charge, most of their properties are much alike!

He thought with himself : " Like electron that has two possible spin states, and the rest of their properties are alike, It has to be something like spin "ISOSPIN" which p & n are different states of it! "

That's mesmerising how he thought about this apparent similarity.

Thus, the nucleons form a doublet, just like e^- that show a doublet structure as a consequence of two-state spin orientations.

Later on, Gell-Mann found that particles with nearly equal masses can always be arranged in isospin multiplet.

$(p, n) \rightarrow$ isospin doublet

$(\pi^+, \pi^0, \pi^-) \rightarrow$ isospin triplet

$(\Delta^{++}, \Delta^+, \Delta^0, \Delta^-) \rightarrow$ isospin quadruplet

particles within a multiplet have "equal masses" but different charges.

But isospin symmetry is actually approximation! since the masses are not exactly equal!

Isospin symmetries are internal degree of freedom.

I mean, we only transform internal d.o.f.s with some rules!

E.x: $SU(3)$ on triplets $(\vec{\phi})$

$$\begin{pmatrix} \alpha_1 & \alpha_2 & \alpha_3 \\ \beta_1 & \beta_2 & \beta_3 \\ \theta_1 & \theta_2 & \theta_3 \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix} \rightarrow \sum_i \begin{pmatrix} \alpha_i \phi_i \\ \beta_i \phi_i \\ \theta_i \phi_i \end{pmatrix}$$

They are different from the so-called "spacetime symmetries", which act on $x^\mu = (t, \vec{x})$ coordinates!

These states transform according to an irrep of the symmetry group.

just like the action of rotation on 3D vector.

$SU(3)$ acts on quark field which its bases represent different quarks (u, d, s, \dots)

Concrete examples:

Consider decay of $\Delta^{++} \rightarrow p \pi^+$

under isospin matrices (analogous to Pauli-matrices)

$$I_3^{(\Delta)} |\Delta^{++}\rangle = \frac{3}{2} |\Delta^{++}\rangle$$

$I_3^{(\Delta)}$ on $|\Delta^+, \Delta^0, \Delta^-\rangle$ has $\frac{1}{2}, -\frac{1}{2}, -\frac{3}{2}$ eigenvalues respectively

$$I_3^{(\pi)} |\pi^+\rangle = |\pi^+\rangle$$

$I_3^{(\pi)}$ on $|\pi^0, \pi^-\rangle$ has 0 and (-1) eigenval.

$$I_3^{(p)} |p\rangle = \frac{1}{2} |p\rangle$$

$I_3^{(p)}$ on $|n\rangle$ has $-\frac{1}{2}$ eigenvalue.

$$\begin{aligned} \text{so that } \underbrace{I_3}_{I_3^{(p)} \otimes \mathbb{1}^{(\pi)} + \mathbb{1}^{(p)} \otimes I_3^{(\pi)}} |p\rangle \otimes |\pi^+\rangle &= I_3^{(p)} |p\rangle \otimes \mathbb{1}^{(\pi)} |\pi^+\rangle + \mathbb{1}^{(p)} |p\rangle \otimes I_3^{(\pi)} |\pi^+\rangle \\ &= \frac{1}{2} |p\rangle \otimes |\pi^+\rangle + |p\rangle \otimes |\pi^+\rangle = \frac{3}{2} |p\rangle \otimes |\pi^+\rangle \end{aligned}$$

Remark: The index 3 in I_3 stands for third generator of isospin group.

just like $\sigma_z = \sigma_3$ acts on $|m = \pm \frac{1}{2}\rangle$ gives $\pm \frac{1}{2} |m = \pm \frac{1}{2}\rangle$.

so the isospin quantum number is conserved during this reaction.

Remark: I_{\pm} raise and lower isospin quantum number, like

$$I_-^{(\Delta)} |\Delta^{++}\rangle = I_- | \frac{3}{2}, \frac{3}{2} \rangle = \sqrt{3} | \frac{3}{2}, \frac{1}{2} \rangle = \sqrt{3} |\Delta^+\rangle$$

Let's work out a mysterious physical fact, only by representation & group theory!

It's been observed that Δ predominantly decays into nucleon + pion.

By charge conservation, Δ^{++} just decays like $\Delta^{++} \rightarrow p \pi^+$

For other Δ -particles, the situation is more complicated!

Let's start the analysis by looking into six different nucleon-pion states:

$$I_3^{(p)} \mathbb{1}^{(\pi)} + \mathbb{1}^{(p)} I_3^{(\pi)} \leftarrow I_3 (|p\rangle |\pi^+\rangle) = \frac{3}{2} |p\rangle |\pi^+\rangle$$

$$I_3 (|p\rangle |\pi^0\rangle) = \frac{1}{2} |p\rangle |\pi^0\rangle$$

$$I_3 (|p\rangle |\pi^-\rangle) = -\frac{1}{2} |p\rangle |\pi^-\rangle$$

$$I_3(|n\rangle|\pi^+\rangle) = \frac{1}{2} |p\rangle|\pi^+\rangle$$

$$I_3(|n\rangle|\pi^0\rangle) = -\frac{1}{2} |p\rangle|\pi^0\rangle$$

$$I_3(|n\rangle|\pi^-\rangle) = -\frac{3}{2} |p\rangle|\pi^-\rangle$$

$\pm\frac{3}{2}$ eigenvalues occurs only once.

$\pm\frac{1}{2}$ " " twice!

Similarities with "angular momentum addition" is STRIKING.

the above repetition of eigenvalues means that in product space of $|p,n\rangle \otimes |\pi^\pm, \pi^0\rangle$ one can have $\frac{3}{2}$ and $\frac{1}{2}$ (total isospin) values!
 $1 \otimes \frac{1}{2} = \frac{1}{2} \oplus \frac{3}{2}$.

Let's see what decay of Δ^{++} teaches us, and apply it to other Δ -particle's decay.

$$\Delta^{++} \longrightarrow p \pi^+$$

By isospin invariance of particle physics (in ordinary regimes) we have

$$\begin{array}{ccc} \Delta^{++} & \longrightarrow & p \quad \pi^+ \\ \downarrow I_- & & \underbrace{}_{\downarrow I_-} \end{array}$$

$$\sqrt{3} |\Delta^+\rangle \longrightarrow |n\rangle|\pi^+\rangle + \sqrt{2} |p\rangle|\pi^0\rangle$$

\Downarrow

$$|\Delta^+\rangle \longrightarrow \frac{1}{\sqrt{3}} |n\rangle|\pi^+\rangle + \sqrt{\frac{2}{3}} |p\rangle|\pi^0\rangle$$

Thus, the decay of Δ^+ should be this!

The theoretical calculation teaches us that the ratio of occurrence of $p\pi^0$ decay to $n\pi^+$ decay must be 2.

This is in complete agreement with observational data.

Apply once more I_- ... to get:

$$| \Delta^0 \rangle \longrightarrow \frac{1}{\sqrt{3}} | p \rangle | \pi^- \rangle + \sqrt{\frac{2}{3}} | n \rangle | \pi^0 \rangle \quad \checkmark$$

So, the 2:1 ratio of particles' decay could somehow related to its internal symmetries!

Various isospin multiplet + masses!

	isospin-multiplet	spin	isospin	mass
pion	π^\pm, π^0	0	1	139.6, 135
kaon	K^0, K^+	0	$\frac{1}{2}$	497.7, 493.7
$\overline{\text{kaon}}$	K^-, \overline{K}^0	0	$\frac{1}{2}$	493.7, 497.7
eta, η^0		0	0	547.3
eta', η'		0	0	960
nucleons	n, p	$\frac{1}{2}$	$\frac{1}{2}$	938.5, 938.3
delta	$\Delta^-, \Delta^0, \Delta^+, \Delta^{++}$	$\frac{3}{2}$	$\frac{3}{2}$	≈ 1235
sigma	$\Sigma^-, \Sigma^0, \Sigma^+$	$\frac{1}{2}$	1	1192.4, 1192, 1189
sigma*	$\Sigma^{*-}, \Sigma^{*0}, \Sigma^{*+}$	$\frac{3}{2}$	1	≈ 1385
lambda	Λ^0	$\frac{1}{2}$	0	

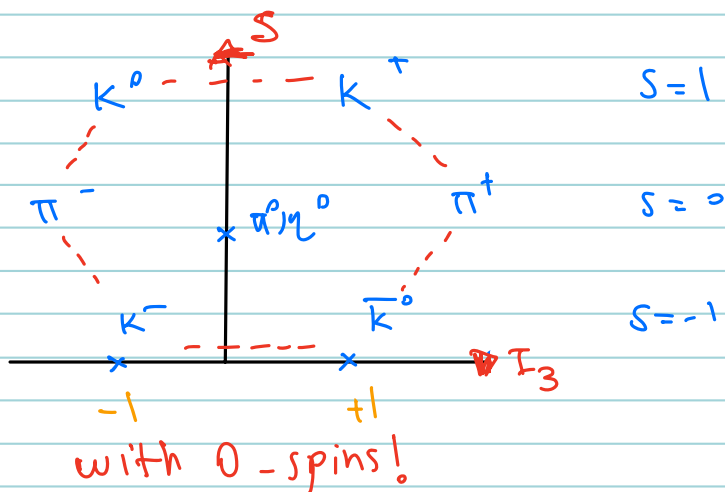
$ks_i \quad \equiv \quad \equiv^-, \quad \equiv^0$
 $ks_i^* \quad \equiv \quad \equiv^{*-}, \quad \equiv^{*0}$
 omega Ω^-

$\frac{1}{2}$
 $\frac{3}{2}$
 $\frac{3}{2}$

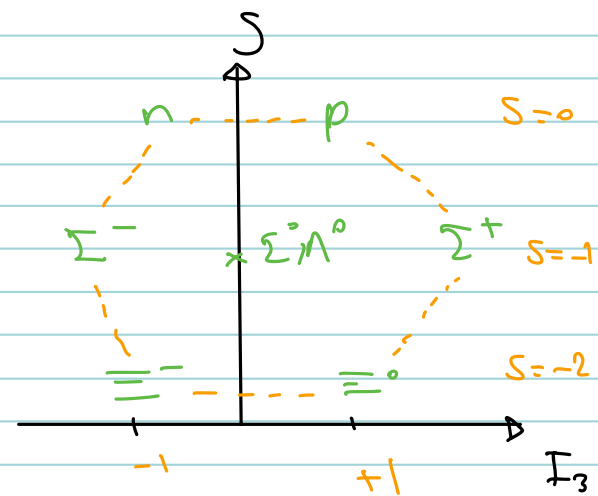
$\frac{1}{2}$
 $\frac{1}{2}$
 0

$SU(3)$ and its role in particle physics are also magical!

meson octet.



Baryons octet.



the rep. thry of $SU(3)$ gives a magical structure to all known particles! It embeds them to known reps of $SU(3)$.