

Riemannian geometry

- x Riemannian manifolds
- x Connections & Covariant derivative
- x Geodesics
- x Curvature & Ricci tensors.
- x Isometries
- x Conformal symmetry & CKVs.

x Homogenous Riemannian manifolds.

x Constant curvature spaces

x metrics on Lie manifolds

x operators ($\vec{r}, \nabla^\perp, \nabla^2 = \Delta$, \star -operators.)

x Affine, Killing - Conformal-projective, Jacobi-Hamilton vector fields.

x First & second fundamental forms.

Most of the interesting stuff has not been covered in the classroom.

So the burden of penetrating through all these topics is on your shoulders.

My Aim is to just solve problems, and provide some definitions on new topics ...

Chern-Forms and Chern-classes + RR - GRR - RRH - As

themselves are a possible direction here!

(Problem 1) (on Riemannian Manifolds.)

Let (M, g) be a Riemannian n -manifold. Prove that

i) Given $\alpha, \beta \in T_p^* M$, and orthonormal bases $\{e_i\}$, $i=1,\dots,n$ of $T_p M$, one has $g^{-1}(\alpha, \beta) = \sum_i \alpha(e_i) \beta(e_i)$

(when g^{-1} denotes the contravariant metric associated g , in physics jargon, its g^μ rather than g_μ)

ii) For $X \in T_p M$, one has $g^{-1}(\alpha, X^\flat) = \alpha(X) = g(\alpha^\sharp, X)$
where :

$$\begin{array}{lll} \text{soft, } & b: T_p M \longrightarrow T_p^* M & X^\flat = g(X, \cdot) \\ \text{flat, b'nat} & & \\ \text{sharp, diese} & \# : T_p^* M \longrightarrow T_p M & \alpha^\sharp = g^{-1}(\alpha, \cdot) \\ & \cancel{\#} & \end{array}$$

are musical isomorphisms.

i) In general, $g_{ij}(p)$ is a matrix in $\{e_i\}$ bases, then $g^{ij}(p) = (g_{ij}(p))^{-1}$ is a matrix w.r.t. dual $\{e^j\}$ bases, base of $T_p^* M$.

Starting with generic symmetric metric g_{ij} , its bases might not constitute an orthonormal frame for $T_p M$, e.g. let $n=3$, then g is given by $g = (g_{ij})$ {Not diagonal!}, then

$$e_1^r g_{rr} e_2^s = (1 \ 0 \ 0) \begin{pmatrix} g_{11} & g_{12} & g_{13} \\ g_{21} & g_{22} & g_{23} \\ g_{31} & g_{32} & g_{33} \end{pmatrix} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = g_{12} \neq 0$$

This ailment could be cured by going into bases of eigenvectors of $g_{ij} \rightarrow$ which necessarily fixes $g_{rr} = \delta_{rr} = 1$. Hence,

$$g^{-1}(\alpha, \beta) = (g^{ij} e_i \otimes e_j)(\alpha_r e^r, \beta_s e^s)$$

$$= g^{ij} \alpha_i \beta_j = \delta^{ij} \alpha_i \beta_j = \sum_i \alpha_i \beta_i$$

$$\stackrel{\text{by def}}{=} \sum_i \alpha(e_i) \beta(e_i)$$

$$\text{(ii)} \quad g^{-1}(\alpha, X^b) = g^{ij}(p) \alpha_i (X^b)_j$$

$$= \underbrace{g^{ij}(p)}_{\text{They contract}} \alpha_i (\underbrace{g_{kj}(p)}_{}) X^k$$

They contract

$$= \delta_{ik}^i \alpha_i X^k = \alpha_i X^i = \alpha(X)$$

$$\text{and } g(\alpha^\#, X) = g_{ij}(p) (\alpha^\#)^i X^j$$

$$= g_{ij}(p) g^{ki}(p) \alpha_k X^j$$

$$= \delta_{jk}^k \alpha_k X^j = \alpha_j X^j = \alpha(X) \quad \square$$

(problem 2) (On Connections & Covariant derivative)

$$\forall p: J: T_p M \rightarrow T_p M : J^2 = -1$$

Let g be a "Hermitian metric" on an almost Complex manifold (M, J)

$$\text{i.e. } g(JX, JY) = g(X, Y) \quad \forall X, Y \in X(M)$$

Define $(1,2)$ -tensor F :

$$F(X, Y) = g(X, JY) \quad \forall X, Y \in X(M)$$

prove the following:

(i) F is skew-symmetric.

(ii) F is J -invariant: $F(JX, JY) = F(X, Y)$

Suppose further that ∇ is any metric compatible connection.

$$\text{(iii)} \quad (\nabla_X F)(Y, Z) = g(Y, (\nabla_X J)Z)$$

$$(iv) g((\nabla_X J)Y, Z) + g(Y, (\nabla_X J)Z) = 0$$

i) Easy!

$$\begin{aligned} F(Y, X) &= g(Y, JX) = g(JY, J^2X) \\ &= -g(JY, X) = -g(X, JY) \\ &= -F(X, Y) \end{aligned}$$

ii)

$$\begin{aligned} F(JX, JY) &= g(JX, J^2Y) = -g(JX, Y) \\ &= -g(J^2X, JY) \\ &= g(X, JY) \\ &= F(X, Y) \end{aligned}$$

iii) First notice

$$\nabla_X \{F(Y, Z)\} = (\nabla_X F)(Y, Z) + F(\nabla_X Y, Z) + F(Y, \nabla_X Z)$$

properties of action of ∇ on scalars.

equally ↓ sometimes called generalized pseudo-Leibnitz rule!

by compatibility

$$\text{And } \nabla_X g(Y, JZ) = (\nabla_X g)(Y, JZ) + g(\nabla_X Y, JZ) + g(Y, \nabla_X JZ)$$

$$\begin{aligned} &= \underbrace{g(\nabla_X Y, JZ)}_{F(\nabla_X Y, Z)} + \underbrace{g(Y, (\nabla_X J)Z)}_{g(Y, (\nabla_X J)Z)} + \underbrace{g(Y, J\nabla_X Z)}_{F(Y, \nabla_X Z)} \end{aligned}$$

$$\text{Hence } (\nabla_X F)(Y, Z) = g(Y, (\nabla_X J)Z)$$

(iv) since $F(x, Y)$ is skew-symmetric by (i), so is $\nabla_X F(Y, Z)$

or a direct computation gives its anti-symmetric property.

$$\begin{aligned}
 \text{by (iii)} : g((\nabla_X J)Y, Z) &\stackrel{\text{sym}}{=} g(Z, (\nabla_X J)Y) \\
 &\stackrel{(iii)}{=} \nabla_X F(Z, Y) \\
 &= -\nabla_Y F(X, Z) \\
 &= -g(Y, (\nabla_X J)Z)
 \end{aligned}$$

Hence, we're done. $g((\nabla_X J)Y, Z) + g(Y, (\nabla_X J)Z) = 0$

(problem 3) (on connection and geodesics)

Consider \mathbb{R}^3 with metric $ds^2 = (1+x^2)dx^2 + dy^2 + e^z dz^2$

(i) Find all Γ -symbols.

(ii) Solve geodesic equation.

(iii) Consider $y(t) = (x=t, y=t, z=t)$ curve

obtain parallel transport of the vector (a, b, c) at origin along y .

(iv) Is y , itself, a geodesic?

(v) Find two parallel vector fields $X(t), Y(t)$ on y , such that $g(X(t), Y(t))$ is constant.

$$i) \quad g = \begin{pmatrix} 1+x^2 & & \\ & 1 & \\ & & e^z \end{pmatrix}, \quad g^{-1} = \begin{pmatrix} \frac{1}{1+x^2} & & \\ & 1 & \\ & & e^{-z} \end{pmatrix}$$

since the metric is diagonal, according to the prescription in your lecture notes ...

They involve
derivative
w.r.t
other
Components!

$$\left\{ \begin{array}{l} \Gamma_{rr}^{\lambda} = 0 \\ \Gamma_{rr}^{\lambda} = -\frac{1}{2} g^{\lambda\lambda} \partial_{\lambda} g_{rr} \\ \Gamma_{r\lambda}^{\lambda} = \partial_r (\ln \sqrt{|g_{\lambda\lambda}|}) \end{array} \right.$$

Repeated indices are
not
summed over!

$$\Gamma_{\lambda\lambda}^{\lambda} = \partial_{\lambda} (\ln \sqrt{|g_{\lambda\lambda}|})$$

Hence, only nonvanishing Γ s are:

$$\left\{ \begin{array}{l} \Gamma_{11}^1 = \partial_1 (\ln (1+x^2)^{\frac{1}{2}}) = \frac{x}{1+x^2} \\ \Gamma_{13}^3 = \partial_3 \ln \sqrt{e^1} = \frac{1}{2} \end{array} \right.$$

(ii) According to geodesic equation:

$$\frac{d^2x}{dt^2} + \frac{x}{1+x^2} \left(\frac{dx}{dt} \right)^2 = 0$$

$$\frac{d^2y}{dt^2} = 0$$

$$\frac{d^2z}{dt^2} + \frac{1}{2} \left(\frac{dz}{dt} \right)^2 = 0$$

And, solutions are:

$$\frac{x''}{x'} + \frac{xx'}{1+x^2} = 0 \quad \int \rightarrow \ln x' + \frac{1}{2} \ln (1+x^2) = \ln A.$$

$$\Rightarrow x' = \frac{A}{\sqrt{1+x^2}} \quad \rightarrow \text{integrate } dx \sqrt{1+x^2} = A dt$$

$$\rightarrow \frac{1}{2} \left\{ x \sqrt{1+x^2} + \log(x + \sqrt{1+x^2}) \right\} = At + B$$

$$\text{and } y(t) = Ct + D$$

$$\text{For } z'' + \frac{1}{2}(z')^2 = 0 \rightarrow \text{let } p = z' \Rightarrow \frac{dp}{dt} + \frac{1}{2}p^2 = 0$$

$$\frac{dp}{p^2} + \frac{dt}{2} = 0 \quad \int \quad \frac{1}{p} = \frac{t}{2} + E/2$$

$$\text{so we find : } \frac{2}{t+E} = z' \rightarrow z = 2\log(t+E) + 2\log F$$

$$z = \log(\underbrace{Ft+G^2}_{FE})^2$$

(iii) $\nabla_y X = 0$ is equation of parallel transporting X along y .

$$\frac{dv^i}{dt} + \Gamma_{jk}^i \frac{dx^j}{dt} v^k = 0$$

$$\left\{ \begin{array}{l} \frac{dv^1}{dt} + \frac{x}{1+x^2} \frac{dx}{dt} v^1 = 0 \\ \frac{dv^2}{dt} = 0 \\ \frac{dv^3}{dt} + \frac{1}{2} \frac{d^2}{dt^2} v^3 = 0 \end{array} \right.$$

when $x=t$
 $y=t$
 $z=t$

reduces to

$$\left\{ \begin{array}{l} \frac{dv^1}{dt} + \frac{t}{1+t^2} v_1 = 0 \\ \frac{dv^2}{dt} = 0 \\ \frac{dv^3}{dt} + \frac{v^3}{2} = 0 \end{array} \right.$$

we integrate above equations:

$$v^1 \rightarrow \log v^1 = \log A - \frac{1}{2} \log(1+t^2) \quad \text{with } v^1(0) = a \text{ initial condition} \rightarrow v^1 = \frac{a}{\sqrt{1+t^2}}$$

$$v^2 \rightarrow v^2 = C \quad \text{with } v^2(0) = b \text{ initial const} \rightarrow v^2 = b$$

$$v^3 \rightarrow v^3 = D e^{-t\gamma_2} \quad \text{with } v^3(0) = c \text{ initial const} \rightarrow v^3 = c e^{-t\gamma_2}$$

(iv) Does y satisfy geodesic eq? for example:

$$\frac{d^2x}{dt^2} + \frac{x}{1+x^2} \left(\frac{dx}{dt} \right)^2 ? = 0$$

zero $x=t$

Does $\frac{t}{1+t^2} = 0$ for all t ? \times

so this curve is not geodesic.

(v) In (iii) we found parallel transported VFs.

$$(a, b, c) \xrightarrow{\text{parallel transp}} v(\tau) = \left(\frac{a}{\sqrt{1+\tau^2}}, b, c e^{-\tau\gamma_2} \right)$$

$$\text{take } X(0) = (1, 0, 0) \rightarrow X(\tau) = \left(\frac{1}{\sqrt{1+\tau^2}}, 0, 0 \right)$$

$$Y(0) = (0, 1, 0)$$

$$Y(t) = (0, 1, 0)$$

And $g(x(t), y(t)) = 0$ ⊗

(problem 4) (pedagogical on geodesic)

Consider \mathbb{R}^2 with usual flat metric, $ds^2 = dx^2 + dy^2$.

Does $y(t) = (x=t^3, y=t^3)$, a geodesic?

Since $y(t) = (t^3, t^3)$ like $y=x, \dots$ and straight lines are geodesics of the plane, ... it might be true!

$$\frac{dy}{dt} = 3t^2(1, 1)$$

→ geodesic equation $\frac{D}{dt} \frac{dy}{dt} = 0$

BUT, $\frac{D}{dt} \frac{dy}{dt} = \frac{d}{dt} (3t^2 \partial_x + 3t^2 \partial_y)$

$$= 6t \partial_x + 6t \partial_y \neq 0 !$$

Hence $y(t)$ is not a geodesic

observation: geodesics have constant velocity $|y'(t)|$.

but $|y'(t)| = 3\sqrt{2}t^2 \dots$ it's not constant!

The fact that a curve is a geodesic depends both on its shape and its parametrization!

(problem 5) (on isometries)

Let (\mathbb{H}, g) be UHP, define $SL(2, \mathbb{R})$ action on it by

$$z = x + iy, M = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{R})$$

$$\omega = M z = \frac{az+b}{cz+d}$$

prove that $SL(2, \mathbb{R})$ is group of isometries of (\mathbb{H}, g)

Given $z \in UHP$, $z = x + iy \rightarrow \operatorname{Im}(z) = y > 0$

$$\text{we have } \omega = \frac{az+b}{cz+d} = \frac{a(x+iy) + b}{c(x+iy) + d} \times \frac{(cx+d) - icy}{(cx+d) - icy}$$

$\underbrace{}_{c.c.}$

$$= \frac{(ax+b) + iay}{(cx+d)^2 + (cy)^2} =$$

$$\frac{(ax+b)(cx+d) + acy^2 + i \cdot \{ ay(cx+d) - cy(ax+b) \}}{(cx+d)^2 + (cy)^2}$$

$$(cx+d)^2 + (cy)^2$$

$$+ y(\cancel{ad-bc}) \quad \text{since } \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, \mathbb{R})$$

$$\operatorname{Im} \omega = \frac{\cancel{aycx} + \cancel{ayd} - \cancel{ayc} - \cancel{cyb}}{(cx+d)^2 + (cy)^2} = \frac{y}{(cx+d)^2 + (cy)^2}$$

$$\rightarrow \operatorname{Im} \omega = \frac{\operatorname{Im} z}{|cz+d|^2} \quad (\text{takes real line} \mapsto \text{half-plane})$$

Thus we use UHP under $SL(2, \mathbb{R})$ linear fractional transformation.

Let us check the metric g

$$g = \frac{1}{y^2} (dx^2 + dy^2) = \frac{dz \wedge d\bar{z}}{\text{Im}^2(z)} = \frac{1}{2} \frac{dz \otimes d\bar{z} + d\bar{z} \otimes dz}{\text{Im}^2(z)}$$

on the other hand:

$$dw = \frac{1}{(cz+d)^2} dz$$

$$d\bar{w} = \frac{1}{(c\bar{z}+d)^2} d\bar{z}$$

so under $z \rightarrow w = Mz$

$$g \mapsto \frac{1}{2} \frac{(cz+d)^2 dw \wedge (c\bar{z}+d)^2 d\bar{w}}{|cz+d|^4 \text{Im}^2(w)} = \frac{dw \wedge d\bar{w}}{\text{Im}^2(w)}$$

$$\frac{(cz+d)^2 (c\bar{z}+d)^2}{(c\bar{z}+d)^2}$$

so it leaves the metric invariant $\Rightarrow SL(2, \mathbb{R})$ is indeed the isometry group of H with Poincaré metric.

(problem 6) (Constant curvature spaces)

Related topic: the Riemannian manifold which possesses the maximum number of Killing vectors, is called maximally symmetric space.

By a simple Country argument :

- R is constant.
- $R_{rr} = \frac{1}{n} R g_{rr}$
- $R_{r\lambda r\mu} = \frac{R}{n(n-1)} (g_{r\lambda} g_{\mu\nu} - g_{r\nu} g_{\mu\lambda})$

In Lorentzian signature :

$$\text{MSS} \begin{cases} R > 0 \rightarrow \text{dS} \\ R = 0 \rightarrow \text{mink.} \\ R < 0 \rightarrow \text{AdS} \end{cases}$$

In Euclidean signature:

$$\text{MSS} \begin{cases} R > 0 \rightarrow \text{sphere} \\ R = 0 \rightarrow \text{Euclidean } \mathbb{R}^n \\ R < 0 \rightarrow \text{saddle or hyperbolic manifolds.} \end{cases}$$

Let (M, g) be a Riemannian manifold of Constant Curvature.

We define metric \tilde{g} on $M \times M$ by : $(x_1, x_2, y_1, y_2 \in X(M))$

$$\tilde{g}((x_1, y_1), (x_2, y_2)) = g(x_1, x_2) \circ p_{r_1} + g(y_1, y_2) \circ p_{r_2}$$

We say that metric decomposes to $\tilde{g}_{n \times n} = g \otimes g = \begin{pmatrix} g_m & \\ & g_m \end{pmatrix}$

Is $(M \times M)$ a Constant Curvature Space?!

Let $(p, q) \in M \times M$. Then are charts along with coordinates in M

such that $\begin{cases} p \in U \text{ with coordinates } x^1, \dots, x^n \end{cases}$

$$g_{\alpha\beta} \quad " \quad " \quad x^{n+1}, \dots, x^{2n}$$

Hence $U \times V$ is a chart on $M \times M$ with (x^1, \dots, x^n) coordinates!

Metric \tilde{g} is : $\tilde{g}_{AB}^{(x,y)} = \begin{pmatrix} g_{ij}(x) \\ & g_{i+n, j+n}(y) \end{pmatrix}_{n \times n}$

$A, B \in \{1, \dots, 2n\}$

christoffel symbols :

$$\tilde{\Gamma}_{BC}^A = \frac{1}{2} \tilde{g}^{AD} (\partial_C \tilde{g}_{DB} + \partial_B \tilde{g}_{DC} - \partial_D \tilde{g}_{BC})$$

since \tilde{g} is block diagonal \rightarrow only $\tilde{\Gamma}_{jk}^i(x, y) = \tilde{\Gamma}_{jk}^i(x)$

$$\tilde{\Gamma}_{j+n, k+n}^{i+n}(x, y) = \tilde{\Gamma}_{jk}^i(y)$$

with $i, j, k \in \{1, 2, \dots, n\}$ survives!

Similarly $\tilde{R}_{jkl}^i(x, y) = R_{jkl}^i(x)$ $\tilde{R}_{j+n, k+n, l+n}^{i+n}(x, y) = R_{jkl}^i(y)$

Hence \tilde{R} is not zero.

But if $M \times M$ is going to be a maximally symmetric spaces

then $\tilde{R}_{jkl}^i = \tilde{R}(g_{ik}^i g_{jl} - g_{il}^i g_{jk})$
 $= \tilde{R}(\delta_{ik}^i g_{jl} - \delta_{il}^i g_{jk})$

Note : $g^{\alpha\beta} g_{\beta}^{\gamma} = g^{\alpha\gamma} \rightarrow \left[\quad \right]_{n \times n} (?) = \left(\quad \right)_{\text{itself}}$

In particular $\tilde{R}_{j, k+n, j}^{k+n} = \tilde{K}(g_{jj}) \equiv 0$ by these rules!

$$\Rightarrow \tilde{K} = 0 !$$

so the product manifold can't be nothing except flat
(if possible!)

(problem 7) (on metrics on Lie manifolds!)

Let H be the Heisenberg group $\rightarrow \begin{pmatrix} 1 & x & y \\ 0 & 1 & z \\ 0 & 0 & 1 \end{pmatrix}, (x, y, z) \in \mathbb{R}^3$

(i) Find left-invariant metric g on H , with bases dual to left-invariant VFs!

(ii) Find ∇ -connection of g .

(iii) Is (H, g) space of constant curvature?

(i) The basis of left-invariant 1-forms on H , gives us the metric! just like (dx, dy, dz) in Euclidean space $ds^2 = dx^2 + dy^2 + dz^2$. why it's true? It's hard-core Lie-algebra! and it only works for simple-semi-simple lie algebras! (those without proper ideals!)

We've had left-invariant VFs $\rightarrow B = \left\{ \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, x \frac{\partial}{\partial y} + z \frac{\partial}{\partial z} \right\}$

By a simple linear combination of $\beta_i = A_i dx + B_i dy + C_i dz$

and applying β_i on Left-inv VFs \rightarrow

$$\beta_1 = dx \quad \rightarrow \quad \beta_1(\partial_x) = 1 \quad \text{and on } \partial_y, \partial_z \quad \partial_y, \partial_z = 0$$

$$\beta_2 = dy - x dz \quad " \quad " \quad "$$

$$\beta_3 = dz \quad " \quad " \quad "$$

$$\Rightarrow ds^2 = dx^2 + (dy - x dz)^2 + dz^2$$

(ii) Only by hand, the ∇_x could be found or Koszul identity (not metrics needed) ...

$$2g(\nabla_x X, Z) = g([X_1 X], Z) - g([X_2 Z], X) + g([Z X], X)$$

We only need to know $\nabla_{X_i} X_j$.

The only non-zero brackets are $[X_1, X_3] = -[X_3, X_1] = X_2$

$$\Rightarrow g(\nabla_{X_1} X_1, X_1) = 0 \quad \xrightarrow{\substack{\text{metric's} \\ \text{no syante}}} \quad \nabla_{X_1} X_1 = 0$$

$$\text{Similarly} \rightarrow \nabla_{X_1} X_2 = -\frac{X_3}{2} \quad \text{and} \quad \nabla_{X_2} X_1 = +\frac{X_3}{2}$$

by torsionless connection!

$$\nabla_{X_1} X_3 = -\nabla_{X_3} X_1 = \frac{1}{2} X_2$$

$$\nabla_{X_2} X_3 = -\nabla_{X_3} X_2 = -\frac{1}{2} X_1$$

(iii) Direct computation of curvature gives it's not constant!

It's sectional curvature also isn't constant along

any cross-section.

sectional curvature on $K = \langle x_1|_p, x_2|_p \rangle$ plane.

$$K(p) = R(x_1, x_2, x_1, x_2)(p)$$

$$= g(\nabla_{x_1} \nabla_{x_2} x_2 - \underbrace{\nabla_{x_2} \nabla_{x_1} x_2}_{-\frac{1}{2}x_3} - \nabla_{[x_1, x_2]} x_2, x_1)(p)$$
$$= \frac{1}{2}(\frac{1}{2})x_1 = \frac{1}{4}x_1$$

$$= g(\frac{1}{4}x^1, x^1)(p) = \frac{1}{4}$$

take $K' = \langle x_1|_p, x_3|_p \rangle$

$$K(p') = \dots = -\frac{1}{4} \quad \textcircled{*}$$

(problem 8) (on Operators on Riemannian manifolds)

Remember musical isomorphisms:

$$b : T_p M \longrightarrow T_p^* M$$
$$x \longmapsto x^b \qquad x^b(\gamma) = g(x, \gamma)$$

$$\# : T_p^* M \longrightarrow T_p M$$
$$\omega \longmapsto \omega^\# \qquad \omega^\#(\xi) = g^{-1}(\omega, \xi)$$

Define grad of f as : $\text{grad } f = (\text{d}f)^\#$

Prove :

$$(i) g(\text{grad } f, X) = X(f)$$

$$(ii) (\frac{\partial}{\partial x^i})^b$$

$$(iii) (dx^i)^*$$

(iv) write grad f in local Coordinates.

(v) Recover the usual grad in \mathbb{R}^3 ,

$$b : TM \longrightarrow T^* M$$

$$X \longmapsto g^{ij} x^i e^j = x^j e^j$$

$$\# : T^* M \longrightarrow TM$$

$$\omega \longmapsto g^{ij} \omega_i e_j = \omega^j e_j$$

$$\text{so that } (x^b)^\# : X \xrightarrow{\#} (x^b)^j e_j$$

$$\text{But } (x^b)^j b_j = x^j$$

$$\Rightarrow (x^b)^\# = x^j e_j = X !$$

$\#$, b are inverse of each other!

$$(i) g(\text{grad } f, X) = g((df)^\#, X) = \underbrace{(df)^\#}_\text{by def of } b(X)$$

$$= df(X) = Xf$$

✓

$$(ii) \text{Def: } (\frac{\partial}{\partial x^i})^b (\frac{\partial}{\partial x^j}) = g\left(\frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j}\right) = g_{ij}$$

$$\rightarrow \left(\frac{\partial}{\partial x^i} \right)^b = g_{ik} dx^k$$

$$= g_{ik} dx^k \frac{\partial}{\partial x^i}$$

(iii) Take $\#$ from (ii) $\rightarrow \frac{\partial}{\partial x^i} = g_{ik} (dx^k)^\#$

$$\Rightarrow (dx^i)^\# = g^{ji} \frac{\partial}{\partial x^i}$$

(iv) $\text{grad } f = (\nabla f)^\# = \left(\frac{\partial f}{\partial x^i} dx^i \right)^\#$

$$= \frac{\partial f}{\partial x^i} g^{ji} \frac{\partial}{\partial x^j}$$

$$= g^{ji} \frac{\partial f}{\partial x^i} \frac{\partial}{\partial x^j}$$



Q1 In \mathbb{R}^3 $g^{ij} = 1 \rightarrow \text{grad } f = \sum_i \frac{\partial f}{\partial x^i} \frac{\partial}{\partial x^i}$ ✓

(problem 9) (Operator gymnastics)

Consider \mathbb{R}^n with Euclidean metric. and $w = dx^1 \wedge \dots \wedge dx^n$

(i) prove that given $\Omega_k \in \Lambda^k(\mathbb{R}^n)$ there's only one
 * Ω_k of $\text{deg} n-k$:

$$*\Omega_k(x_1, \dots, x_{n-k}) w = \Omega_k \wedge x_1^b \wedge \dots \wedge x_{n-k}^b$$

(ii) Hodge star operator is defined by previous formula

$$*: \Lambda^k \mathbb{R}^n \mapsto \Lambda^{n-k} \mathbb{R}^n$$

$$\text{prove: } \star^2 = (-1)^{k(n-k)}$$

$$\star^{-1} = (-1)^{k(n-k)} \star$$

$$\Omega_k \wedge (\star \oplus_k) = \bigoplus_k \wedge (\star \Omega_k)$$

(iii) The Codifferential operator $\delta: \Lambda^k \mathbb{R}^n \rightarrow \Lambda^{k-1} \mathbb{R}^n$ is defined:

$$\delta = (-1)^{n(k+1)+1} \star d \star$$

prove that $\delta^2 = 0 \rightarrow$ it establishes a cohomological complex.

(iv) $\Delta: \Lambda^k \mathbb{R}^n \rightarrow \Lambda^k \mathbb{R}^n$ is defined to be:

$$\Delta = (d + \delta)^2 = d\delta + \delta d$$

$$\text{prove } \Delta f = - \sum_{i=1}^n \frac{\partial^2 f}{\partial (x^i)^2}.$$

We've only to prove them to a basis of ext. algebra.

Let $\{x^1, \dots, x^n\}$ be an orthonormal bases of VFs on \mathbb{R}^n and $\{\theta^1, \dots, \theta^n\}$ the dual bases.

And also the $i_1 < \dots < i_{n-k}, j_1 < \dots < j_{n-k}$ ordered indices!

$$(i) \star (\theta^{i_1} \wedge \dots \wedge \theta^{i_{n-k}}) (x_{j_1}, \dots, x_{j_{n-k}}) w = \text{by def}$$

$$(x_{j_1})^\flat = \theta^{j_1}$$

$$= \theta^{i_1} \wedge \dots \wedge \theta^{i_{n-k}} \wedge \theta^{j_1} \wedge \dots \wedge \theta^{j_{n-k}}$$

First of all, this vanishes if i, j indices are not complement.

so the above equation gives:

$$\star(\theta^{i_1} \wedge \dots \wedge \theta^{i_k}) = \text{sgn}(i_1, \dots, i_k, j_1, \dots, j_{n-k}) \theta^{j_1} \wedge \dots \wedge \theta^{j_{n-k}}$$

so we deduce that

$$\{\star(\theta^{i_1} \wedge \dots \wedge \theta^{i_k})\}_{(x_{j_1}, \dots, x_{j_{n-k}})} w =$$

$$\text{sgn}(i_1, \dots, i_k, j_1, \dots, j_{n-k}) w$$

Since action of bases specify a form uniquely,
therefore the correspondence is 1-1. \square

(ii) From $\boxed{\quad}$ above we take another \star .

$$\star\{\star(\theta^{i_1} \wedge \dots \wedge \theta^{i_k})\} = \text{sgn}(i_1, \dots, i_k, j_1, \dots, j_{n-k}) \times \\ \star\{\theta^{j_1} \wedge \dots \wedge \theta^{j_{n-k}}\}$$

$$= \underbrace{\text{sgn}(i_1, \dots, i_k, j_1, \dots, j_{n-k})}_{\text{work this out when both } n, k \text{ are even/odd / } e0/o0e} \underbrace{\text{sgn}(j_1, j_2, \dots, j_{n-k}, i_1, \dots, i_k)}_{(-1)^{k(n-k)}} \theta^{i_1} \wedge \dots \wedge \theta^{i_k}$$

$$(-1)^{k(n-k)}$$

$$\Rightarrow \star^2 = (-1)^{k(n-k)}$$

From \star results ($\star \star^{-1}$) $\rightarrow \star = (-1)^{k(n-k)} \star^{-1}$

$$\rightarrow \star^{-1} = (-1)^{k(n-k)} \star.$$

$$\text{Consider } \Omega_k = \theta^{i_1} \wedge \dots \wedge \theta^{i_k}$$

$$\Theta = \theta^{j_1} \wedge \dots \wedge \theta^{j_{n-k}}$$

$$\rightarrow \Omega_k \wedge (\star \Theta) = \begin{cases} 0 & \text{if } \{i_1, \dots, i_k\} \neq \{j_1, \dots, j_k\} \\ \omega & \text{if } " = " . \end{cases}$$

$$\text{similarly } \Theta_k \wedge (\star \Omega_k) = \begin{cases} 0 & \text{if } \{i_1, \dots, i_k\} \neq \{j_1, \dots, j_k\} \\ \omega & \text{if } " = " . \end{cases}$$

$$\rightarrow \Omega_k \wedge (\star \Theta_k) = \Theta_k \wedge \Omega_k$$

$$(ii)$$
 $\delta^2 = \star \underline{\star} \star \star \underline{\star} \star \star$ \star
 $\star^2 \quad \leftarrow \Omega_k \quad \rightarrow \star^2 = (-1)^{(n-k+1)(k-1)}$
 $(n-k)$
 $(n-k+1)$

$$= (-1)^{(n-k+1)(k-1)} \star \underline{\star^2} \star = 0 \quad \text{Q.E.D.}$$

(iv) Note that $\delta f = 0$ since $\star f \propto \omega \rightarrow$

$$\star \delta f = \star \omega = 0$$

$$\Rightarrow \Delta f = (\star \delta + \delta \star) = \delta \star f$$

$$= - \star \star \star \star f$$

$$= - \star \star \star \frac{\partial f}{\partial x_i} dx^i$$

$$= - \star \star \left\{ (-1)^{i-1} \frac{\partial f}{\partial x_i} dx^1 \wedge \dots \wedge \overset{i}{\cancel{dx^i}} \wedge \dots \wedge dx^n \right\}$$

$$= - \star (-1)^{i-1} \frac{\partial^2 f}{\partial x_i \partial x_j} dx^j \wedge dx^1 \wedge \dots \wedge \overset{j}{\cancel{dx^i}} \wedge \dots \wedge dx^n$$

$$= - \star \left(\sum_{j=1}^n \frac{\partial^2 f}{(\partial x_j)^2} \times dx^1 \wedge \dots \wedge dx^n \right)$$

$$= - \sum_{j=1}^n \frac{\partial^2 f}{(\partial x_j)^2} \star \omega = 1$$

$$= - \sum_{j=1}^n \frac{\partial^2 f}{(\partial x_j)^2}$$

(problem 10) (Def of Killing KVF)

If $X \in X(M)$ is KVF iff $L_X g = 0$

our starting def is that X is KVF if g is invariant under pull back along its integral curve!

so if ϕ is integral curve of $X \rightarrow \phi_t^* g = g$

$$\text{By def } L_X g = \lim_{t \rightarrow 0} \frac{\phi_t^* g - g|_{t=0}}{t} = 0$$

 (look at previous notes ...)

since $\mathcal{L}_X g = 0$, for any tensor field K :

$$\phi_s^*(\mathcal{L}_X K) = - \left(\frac{d}{dt} (\phi_t^* K) \right)_{t=s}$$

Hence, by hypothesis $\rightarrow 0 = \phi_s^* (\underbrace{\mathcal{L}_X g}_{\text{zero}}) = - \left(\frac{d}{dt} (\phi_t^* g) \right) |_{t=s}$

$$\rightarrow \phi_t^* g \text{ doesn't depend on } t \rightarrow \phi_0^* g = g$$

$$= \phi_0^* g = g$$

□

(problem 11) (\mathbb{R}^3 's KVF's)

The real Lie algebra generated by

$$\langle \partial_x, \partial_y, \partial_z, x\partial_y - y\partial_x, -y\partial_z + z\partial_y, z\partial_x - x\partial_z \rangle$$

are KVF's of \mathbb{R}^3 .

Let $X = \lambda^i \frac{\partial}{\partial x^i}$ be a KVF $\rightarrow \mathcal{L}_X g = 0 \rightarrow$

$$\mathcal{L}_X g = \sum_{i,j} \left(\frac{\partial \lambda^j}{\partial x^i} + \frac{\partial \lambda^i}{\partial x^j} \right) dx^i \otimes dx^j$$

we deduce \rightarrow

$$\begin{cases} \frac{\partial \lambda^1}{\partial x^1} = \frac{\partial \lambda^2}{\partial x^2} = \frac{\partial \lambda^3}{\partial x^3} = 0 \\ \frac{\partial \lambda^1}{\partial x^2} + \frac{\partial \lambda^2}{\partial x^1} = \frac{\partial \lambda^1}{\partial x^3} + \frac{\partial \lambda^3}{\partial x^1} = \frac{\partial \lambda^2}{\partial x^3} + \frac{\partial \lambda^3}{\partial x^2} = 0 \end{cases}$$

so $\lambda_1(y, z)$, $\lambda_2(x, z)$, $\lambda_3(x, y)$ from first eqs.

The last line ∂_2 (first) & ∂_3 (second)

$$\frac{\partial^2 \lambda^1}{\partial (x^2)^2} = \frac{\partial^2 \lambda^1}{\partial (x^3)^2} = 0$$

$$\Rightarrow \lambda^1 = A_1 yz + B_1 y + C_1 z + D_1$$

similarly $\rightarrow \lambda^2 = A_2 xz + B_2 z + C_2 x + D_2$

$$\lambda^3 = A_3 xy + B_3 x + C_3 y + D_3$$

play into last equations \rightarrow

$$\lambda^1 = -C_2 y + C_1 z + D_1$$

$$\lambda^2 = -C_3 z + C_2 x + D_2$$

$$\lambda^3 = -C_1 x + C_3 y + D_3$$

This gives the generators in the statement.

By using $L_{(x,y)} = [L_x, L_y]$ one can check that

All these linear combinations actually form

a Lie algebra.

There are a lot of VFs on manifolds.

X is a projective VF if

$$(\mathcal{L}_X \nabla)(Y, Z) = [X, \nabla_Y Z] - \nabla_{[X, Y]} Z - \nabla_Y [X, Z]$$

is of the form $\theta(X)Z + \theta(Y)Y$

for some $\theta \in \Lambda^1(M)$!

CKVs are easier $\rightarrow X$ is CKV if $\mathcal{L}_X g = f g$

They preserve the conformal class of metric on the man.

X is harmonic VF if X^\flat is ker of Laplace
 $\Delta X^\flat = 0$

but how Δ acts on forms? $\Delta = \cancel{\nabla} \times \cancel{\nabla}$

such questions are in the realm of harmonic analysis.

Hamiltonian VFs are defined via symplectic structure on M .

$$dH(Y) = \underbrace{\omega(X_H, Y)}$$

Asymm. non-deg form w

like cot. bundle.

Hamiltonian VF of X .

(problem 12) (on Stokes Theorem)

Let M be a compact Riemannian manifold.

- (i) prove that $\delta \alpha \star d \star$ is adjoint of the "d"-operator w.r.t. inner product of integration.

$$\langle d\alpha, \beta \rangle = \langle \alpha, \delta \beta \rangle$$

$\alpha, \beta \in \Lambda^r(M)$

and vice-versa

- (ii) $\Delta = d\delta + \delta d$ is self-adjoint w.r.t inner product of integration

$$\langle \Delta \alpha, \beta \rangle = \langle \alpha, \Delta \beta \rangle$$

- (i) Def of inner product of integration

$$\langle \alpha, \beta \rangle = \int_M (\alpha \wedge * \beta) \quad \{ \text{on Compact } M \}$$

why on compact manifold M .

$$0 = \int_{\partial M} (\alpha \wedge * \beta) = \int_M d(\alpha \wedge * \beta)$$

$$= \int_M dd^* \alpha \wedge * \beta + (-1)^r \alpha \wedge \underbrace{d}_{\star^{-1}} * \beta$$

$$d = (-1)^{n(n-r+1)} \star \circ \star$$

$$\star^{-1} d = (-1)^{n(n-r+1)} d \star$$

$$= \int_M d\alpha \wedge * \beta - \int_M \alpha \wedge \underbrace{\delta}_{\star^{-1} d} \beta$$

$$(-1)^{n(n-r)} \star \delta = (-1)^{\#} d \star$$

$$\Rightarrow \langle d\alpha, \beta \rangle = \langle \alpha, \delta\beta \rangle$$

(ii) $\langle d\alpha, \beta \rangle = \langle d\delta\alpha + \delta d\alpha, \beta \rangle$

$$= \langle d\delta\alpha, \beta \rangle + \langle \delta d\alpha, \beta \rangle$$

problem (i) ↓

$$= \langle \delta d\alpha, \beta \rangle + \langle d\alpha, d\beta \rangle$$

$$= \langle \alpha, d\delta\beta \rangle + \langle \alpha, \delta d\beta \rangle$$

$$= \langle \alpha, (d\delta + \delta d)\beta \rangle = \langle \alpha, \Delta\beta \rangle$$

✓

