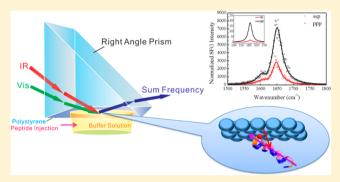


Site-Specific Orientation of an α -Helical Peptide Ovispirin-1 from **Isotope-Labeled SFG Spectroscopy**

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Supporting Information

ABSTRACT: Sum-frequency generation (SFG) vibrational spectroscopy is often used to probe the backbone structures and orientations of polypeptides at surfaces. Using the ovispirin-1 polypeptide at the solid/liquid interface of polystyrene, we demonstrate for the first time that SFG can probe the polarization response of a single-isotope-labeled residue. To interpret the spectral intensities, we simulated the spectra using an excitonic Hamiltonian approach. We show that the polarization dependence of either the label or the unlabeled amide I band alone does not provide sufficient structural constraints to obtain both the tilt and the twist of the ovispirin helix at a solid/liquid interface, but that both can



be determined from the polarization dependence of the complete spectrum. For ovispirin, the detailed analysis of the polarized SFG experimental data shows that the helix axis is tilted at roughly 138° from the surface normal, and the transition dipole of the isotope-labeled C=O group is tilted at 23° from the surface normal, with the hydrophobic region facing the polystyrene surface. We further demonstrate that the Hamiltonian approach is able to address the coupling effect and the structural disorder. For comparison, we also collected the FTIR spectrum of ovispirin under similar conditions, which reveals the enhanced sensitivity of SFG for structural studies of single monolayer peptide surfaces. Our study provides insight into how structural and environmental effects appear in SFG spectra of the amide I band and establishes that SFG of isotope-labeled peptides will be a powerful technique for elucidating secondary structures with residue-by-residue resolution.

1. INTRODUCTION

Isotope-labeling and vibrational spectroscopy provides sitespecific structural information on polypeptides and proteins. 1,2 Of the intrinsic vibrations inherent to proteins, the amide I band (mainly contributed by the backbone C=O stretching mode) is most often utilized in structural studies because its frequency and line shape are characteristic of the secondary structures and solvent environments of the backbone. Residue specific structural and environmental information can be obtained using isotope-labeling. Labeling the backbone carbonyl with 13 C, the amide I band is shifted by \sim 40 cm $^{-1}$. A shift of \sim 66 cm⁻¹ can be achieved with $^{13}C=^{18}O$ labeling. 1,3 Isotopelabeling has been used in conjunction with 1D (FTIR) and 2D IR studies to obtain residue-by-residue backbone structural information about soluble proteins, 4-6 protein/peptide folding kinetics, ^{7–9} and amyloid aggregation and structure, ^{10,11} to name only a few studies.

Isotope-labeling and vibrational spectroscopy are particularly valuable for studies of membrane proteins because they do not easily crystallize and are difficult to study with solution NMR. Solid-state NMR is a powerful technique, but it is challenging to study membrane interactions of proteins and peptides in real time. The samples for solid state NMR studies are premixed lipids and proteins/peptides and usually contain multiple layers

of lipids, not a single lipid bilayer. Vibrational spectroscopy has been used to probe the structure of membrane-bound α -helical bundles, 12 ion channels, 4 transmembrane α -helices 13–15 and helical dimers.⁵ In fact, the polypeptide that is the focus of this study, ovispirin, was previously 13C18O-labeled for 2D IR experiments on its membrane-bound structure. In that experiment, residue-by-residue structural resolution revealed the backbone orientation, tilt, and secondary structure of each residue along nearly its entire length.³ Another approach is to measure the angles of individual transition dipoles relative to the normal of the bilayer, which is done by macroscopically aligning the bilayers on an FTIR sample cell. Polarized light is then used to measure the linear dichroism of the isotope labels to back out the absolute angles. By isotope-labeling a series of amino acids, this approach was used to obtain the structure of the CD3-ζ helical bundle¹² and to study the conformation gating of the M2 ion channel from the influenza viral protein.⁴

Although orientational constraints derived from FTIR spectroscopy are very valuable, similar to solid state NMR, FTIR dichroism studies cannot be performed on single

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monolayers because the signal is too small to obtain accurate angles. For adequate signal strength, hundreds or thousands of bilayers are stacked on top of one another. Stacking works well for equilibrated structures, but precludes experiments involving kinetics, drug binding, applied potentials, or systems that cannot be stacked, such as solid interfaces. Moreover, because linear dichroism is attenuated by disorder, X-ray reflectivity must be used to independently assess the disorder of the stacks. ¹⁶ Thus, it would be quite beneficial to have a technique that is sensitive enough to measure the transition dipole angles of isotope-labeled peptides of a single monolayer.

In contrast to FTIR, sum-frequency generation (SFG) vibrational spectroscopy has the sensitivity to observe peptides at submonolayer surface coverages. SFG is a second-order nonlinear optical spectroscopy that measures the second-order nonlinear optical susceptibility, $\chi^{(2)}$, of the sample. According to the selection rule of the second-order nonlinear optical process, SFG detects signals only from media where inversion symmetry is broken, such as surfaces and interfaces. Since 2003, SFG has been successfully used to study biological molecules with various secondary structures and on different types of surfaces, including α -helices, α -hel often many approximations that go into the interpretation of the SFG spectrum of a peptide. For example, to back out the tilt of an α -helix at an interface, one often assumes that the peptide is rotationally averaged about its helical axis. 17-20 This is possible for peptides that insert perpendicularly into the cell membrane because the interior of the lipid bilayer is quite homogeneous (hydrophobic), but this is certainly unlikely for amphipathic and many other types of peptides at a hydrophobic substrate/aqueous solution interface. This assumption is required because there are not enough observables to obtain unique tilt and internal rotation angles.

Other assumptions are also common, such as that random coil regions are weak or have different peak centers in SFG spectra and that the vibrational modes follow symmetry rules. $^{17-20}$ In this paper, we show that a single residue can be resolved using 13 C isotope-labeling in the 18-residue, antibiotic ovispirin at a polystyrene/peptide solution interface. The additional observables that this label provides eliminate the need for rotational averaging. We also learn that coupling to the isotope label may need to be considered. Previously, SFG spectroscopy was used in conjunction with NMR spectroscopy to provide a structural model of the synthetic LK α 14 peptide by measuring isotope-labeled side chains. 27,28 Here, we demonstrate for the first time that SFG spectroscopy can also be used to site-selectively probe the backbone carbonyl groups themselves, thereby providing a more direct measure of secondary structure.

Ovispirin-1 is an ideal target for our initial SFG experiments because its structure has been studied extensively in solution and on model phospholipid membranes. Solution NMR experiments show that in 33% TFE/67% PBS buffer at pH 6.5 ovispirin-1 forms a slightly curved α -helix over residues 4–16, with random coil outside of this region. Solid-state NMR experiments show that on membrane bilayers, ovispirin is still predominantly α -helical, but lies primarily in the plane of the bilayer, with the helix tilted \sim 84° from the surface normal. Two-dimensional infrared spectroscopy (2DIR) and molecular dynamics (MD) simulations similarly support the α -helical structure and planar orientation and additionally show that its

hydrophobic residues face into the membrane interior, consistent with its amphipathic sequence.³

However, it is unknown what structure ovispirin will take or how it will interact with a purely hydrophobic surface, such as the polystyrene surface we utilize here. As mentioned above, in 33% TFE/67% PBS buffer at pH 6.5 ovispirin-1 forms an α helix over residues 4–16.²⁹ Similarly, our choice of solvent (40% 2,2,2-trifluoroethanol and 60% 10 mM PBS buffer of pH = 7.1) should promote helix formation similar to the solution structure (which will be confirmed later), but the tilt relative to the substrate surface may be considerably different from that in the membrane bilayer systems, and interaction with the highly hydrophobic surface may perturb the overall peptide secondary structure. In this research, we find that the peptide backbone is tilted 138° with respect to the surface normal, and the hydrophobic residues face polystyrene. Thus, the combination of isotope-labeling and SFG spectroscopy brings new light to this system and enables many new systems to be studied that were not previously possible.

2. MATERIALS AND METHODS

SFG theory and experimental details have been reported previously $^{31-35}$ and will not be repeated in detail here. SFG has been applied to study peptides and proteins at interfaces by several groups. 36-39 In our experiment, two laser beams, one visible beam at 532 nm, and an infrared beam with tunable frequency (1300-4300 cm⁻¹) are overlapped spatially and temporally at the bottom of a right-angled prism. 40 The polystyrene (PS) thin film on the prism surface was prepared by spin coating 1 wt % PS solution in toluene on the CaF₂ prism surface at 2500 rpm with a spin coater from Specialty Coating Systems. During every SFG experiment, the PS film was initially in contact with a 2 mL reservoir filled with 40% TFE and 60% 10 mM phosphate buffer (pH = 7.1) in H₂O or D_2O . Then 20 μ L ovispirin-1 peptide stock solution (1 mg/mL in the same solvent) was injected into the 2 mL reservoir. The final subphase peptide solution concentration in the reservoir was 10 μ g/mL. and the equilibration time for the peptide-PS interaction was ~ 1.5 h. To ensure that peptides in the subphase in contact with PS are homogeneously distributed, a magnetic stirrer was used during the data collection process. SFG ssp (s-SFG, s-IR, p-visible) and ppp spectra in the amide I frequency range were collected from ovispirin-1 peptide molecules at the PS/subphase peptide solution interface using a near total reflection geometry.²³ The diameter of the overlapped laser beams at the PS/peptide solution interface was ~0.5 mm. Regular ovispirin-1 (with the sequence of KNLRRIIRKIIHII-KKYG) and isotope-labeled ovispirin-1 (with the 1-carbonyl group of residue 11 Ile isotope-labeled with ¹³C=¹⁶O) were synthesized by Peptide 2.0 Inc. PS, toluene, 2,2,2-trifluoroethanol, PBS, and D2O were obtained from Sigma-Aldrich.

A Nicolet 550 spectrometer (Thermo Fisher Scientific. Inc., MA) was used to collect ATR-FTIR spectra of isotope-labeled and regular ovispirin-1 adsorbed on PS surfaces. A thin PS polymer was deposited on the ZnSe crystal surface by casting 1 wt % polymer solution in toluene and then drying it under a nitrogen gas flow. A 160 μ L portion of 1 mg/mL ovispirin-1 peptide solution in D₂O was injected into the ATR-FTIR trough (~1.6 mL); the final concentration of the ovispirin-1 peptide solution was ~100 μ g/mL (10 times larger than that used in the SFG experiment). The spectra with P and S polarizations were collected ~2 h after the injection of the peptide stock solution into the ATR-FTIR trough. The ATR-

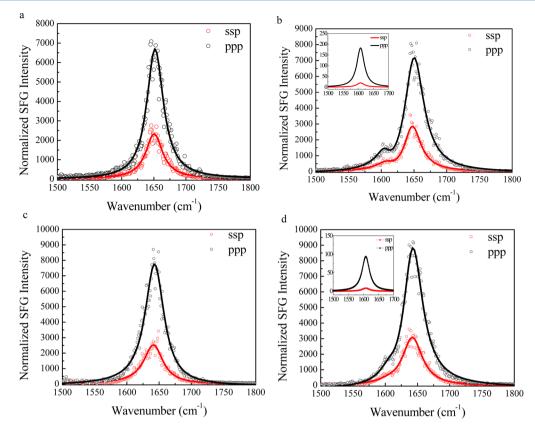


Figure 1. The ssp and ppp SFG amide I spectra of (a) regular ovispirin-1 adsorbed at the PS/peptide solution (with H_2O) interface; (b) isotope-labeled ovispirin-1 adsorbed at the PS/peptide solution (with H_2O) interface; (c) regular ovispirin-1 adsorbed at the PS/peptide solution (with D_2O) interface; (d) isotope-labeled ovispirin-1 adsorbed at the PS/peptide solution (with D_2O) interface.

FTIR sample chamber was purged with nitrogen before and throughout the measurements.

Spectra were simulated using the transition dipole coupling model and an excitonic Hamiltonian, as described in more detail in the Supporting Information. Briefly, peptide or ideal helix structures were loaded from a protein databank file. A transition dipole and Raman polarizability were assigned to each amide I local mode. Transition dipole coupling was used to calculate couplings between all pairs of local amide I modes. The resulting Hamiltonian was diagonalized to yield the eigenstates of the system, whose coefficients were used to take to calculate the normal modes by taking appropriate linear combinations of the local mode transition dipoles and Raman polarizabilities. These responses were rotated to the laboratory frame and summed to give overall intensities for the labeled and unlabeled peaks. 42

3. RESULTS AND DISCUSSION

In what follows, we begin by presenting the experimental data. The data is then interpreted with detailed orientation analyses. First, we use conventional formulas to determine the tilt of the α -helix from the polarization dependence of the unlabeled spectral region of the amide I band, assuming that the helix is isotropic (with a free rotation with the helix principal axis). Second, a similar analysis is applied to the isotope-labeled region, but as we show, these two analyses are incongruent. To understand this dilemma, our third analysis uses the ratio of the unlabeled to labeled amide I bands, which provides an additional observable that improves the structural analysis, suggesting that the peptide has a preferred orientation on the

surface without free rotation. Finally, we highlight that the Hamiltonian approach can address coupling effects and structural disorder issues.

3.1. Experimental Spectra. SFG spectra in the amide I frequency range (1500-1800 cm⁻¹) were collected from both isotope-labeled and unlabeled ovispirin-1 adsorbed at the PS/ peptide solution interface (Figure 1). A strong peak at ~1650 cm⁻¹ was detected from both systems at the interface when a peptide H₂O/TFE buffer solution was used to contact PS. With a peptide D2O/TFE buffer solution for a separate experiment, the amide I band appears at 1642 cm⁻¹, with a similar intensity. The 8 cm⁻¹ difference is typical for an amide I band detected from proteins/peptides in H₂O and in D₂O^{43,44} because the amide I vibrational mode involves the N-H stretch to a small degree. Although the SFG signal in this spectral frequency range overlaps with that contributed by the vibrational mode of random coils, we believe that here, random coils have little contribution to the SFG signal because (1) SFG signal requires the ordering of the chemical groups, and thus, the amide I signals from the random coils likely cancel each other. In other words, SFG spectroscopy is much more sensitive to α -helices than random coils. 45 (2) In both H₂O and D₂O cases, the main amide I peak has a symmetric feature with a similar bandwidth. If the 1650 cm⁻¹ peak in the H₂O case contained two different components with comparable SFG signal intensity contributed from the α -helical structure and the random coiled structure, the deuteration of the peptide solvent would lead to a different spectral feature for the 1642 cm⁻¹ peak, 44 which was not observed experimentally here. (3) We used NLOPredict,⁴⁶ developed by the Simpson group to estimate the contribution of the α -helical structure and the random coiled structure to the

SFG amide I signal, and found that the contribution of the random coiled structure is minimal for two typical orientations: namely, with the helical axis perpendicular and parallel to the polymer surface. We believe that other orientations (tilt orientations) should be intermediate to these two extremes, and thus, the random coil contribution can be ignored in the analysis (Supporting Information). For the isotope-labeling case in H_2O , a well-separated and weak peak at ~1607 cm⁻¹ was also detected, which is the Ile11 isotope-labeled amide I band (Figure 1b). Interestingly, the peak center of the isotopelabeled C=O appears at ~1607 cm⁻¹, appearing to be a shoulder to the 1642 cm⁻¹ amide I main peak when D₂O was used in the solvent (Figure 1d). Previous research on a helical dimer showed that ¹³C=O Leu in a solvent-protected region exhibited a peak at ~1606 cm⁻¹, whereas the ¹³C=O of a solvent exposed Ala absorbed at ~1585 cm⁻¹.⁴⁷ For ovispirin on polystyrene, we find that the ¹³C=O of Ile11 remains at ~1607 cm⁻¹, regardless of solvent. Thus, Ile11 is not exposed to the solvent, which is because this residue faces the polystyrene surface, as we discuss below.

We fit the SFG spectra shown in Figure 1 to extract the intensities of the isotope-labeled peak, since there exist interferences with the unlabeled amide I band. More details of the spectral fitting can be found in the Supporting Information. The experimentally deduced χ_{zzz}/χ_{xxz} ratio after Fresnel coefficients correction (see Supporting Information) is 2.03 ± 0.03 for the 1650 cm⁻¹ band in H₂O for the regular ovispirin-1. For the isotope-labeling case, we fit the spectra considering two peaks with peak centers at 1650 cm⁻¹ peak shows a χ_{zzz}/χ_{xxz} ratio of 1.95 ± 0.09 . For the 1607 cm⁻¹ peak, we obtained an experimentally measured χ_{zzz}/χ_{xxz} ratio of 3.7 ± 0.2 . The spectral fitting results obtained from peptides at the PS/peptide D₂O solution interface are similar to the H₂O cases (Table 1): 2.08 ± 0.02 (regular) and 1.97 ± 0.06 (isotope-

Table 1

	Xzzz/Xxxz	
	amide I	single residue
C12 in H ₂ O	2.03 ± 0.03	N/A
C12 in D ₂ O	2.08 ± 0.02	N/A
C13 in H ₂ O	1.95 ± 0.09	3.7 ± 0.2
$C13$ in D_2O	1.97 ± 0.06	4.1 ± 0.1

labeled) for the 1650 cm⁻¹ signal and 4.1 \pm 0.1 for the 1607 cm⁻¹ peak. To take both solvent cases into consideration, we averaged the χ_{zzz}/χ_{xxz} ratios from both solvents for the main backbone amide I peak of ovispirin-1 (2.00 \pm 0.10) and for the isotope-labeled peak (3.85 \pm 0.35). Comparing the relative helix and label peak intensities, we find that χ_{zzz} (helix)/ χ_{zzz} (label) = 6.18.

3.2. Orientational Analysis of the Helix. Having acquired the data above, the standard approach is to interpret the χ_{zzz}/χ_{xxz} ratio as a measure of the tilt of the helix at the surface. For an ideal helix, the transition dipole of each amide I local mode points 42° from the helix axis, there is a 100° rotation about this axis from one reside to the next, and each local mode has a Raman tensor associated with it. Using this ideal structure, one generates the normal modes of the helix by diagonalizing a vibrational Hamiltonian built from the local modes of the coupled amide I oscillators. The χ_{zzz}/χ_{xxz} ratio is then calculated from the transition dipoles of the normal modes

(details in the Supporting Information) as a function of the tilt (θ) and twist (ψ) angles (defined in Figure 2). Doing so for a 13-residue ideal helix produces the χ_{zzz}/χ_{xxz} ratio shown in Figure 2a when ψ is rotationally averaged around the helix axis. Rotationally averaging around ψ is necessary because there is only one experimental observable (the χ_{zzz}/χ_{xxz} ratio), but two unknowns (θ and ψ). For a perfect and infinitely long helix, one obtains the same θ -dependence whether or not ψ is averaged. For a real helix, one can in practice treat the structure as fully symmetric around the helix axis for any length of helix longer than a few turns. Here, we chose 13 residues for the helix because that is the length of the ovispirin α -helix in solution.²⁹ As discussed above, we believe that the peptide retains its α helical structure at the polymer/solution interface because the SFG spectra are dominated by a single α -helical characteristic peak at ~ 1650 cm⁻¹. This can be further confirmed by the intensity ratio of the main peak and the signal detected from the isotope-labeled unit, which will be discussed in detail in section 3.4. With these considerations in mind, the experimental ratio of $\chi_{zzz}/\chi_{xxz} = 2.00$ gives a helix tilt angle of θ =43 (or 137) \pm 5° relative to the polystyrene surface normal.

3.3. Orientational Analysis of the Isotope Label. One could apply an analysis similar to that above for just the independent isotope label. In this case, the analysis is much simpler because one does not need to construct or diagonalize a Hamiltonian. One would use the same transition dipole and Raman polarizability tensors as above to calculate the tilt-angle dependence of the χ_{zzz}/χ_{xxz} ratio for a single amide-I residue, as shown in Figure 3a. Here, we have defined θ_{label} as the angle between the transition dipole and the surface normal, and we rotationally average around the label transition dipole, ψ_{label} . Note that averaging around $\psi_{ ext{label}}$ is not equivalent to averaging around ψ for the helix or around the labeled C=O bond because the rotation axis is different. The tilt-angle dependence for both the ψ -averaged and the individual $\psi_{label} = 0^{\circ}$ and ψ_{label} = 90° cases are very similar because the Raman tensor's major axis points almost directly along the transition dipole with much smaller (though unequal) components along the perpendicular axes (see Supporting Information). The consequence is that there is a maximum of about 10° difference in tilt angles calculated from ψ -fixed or ψ -averaged curves. Thus, although ψ averaging is necessary to eliminate an unknown variable, just as for the helix analysis, it does not influence θ_{label} by more than a few degrees. Using this curve and the experimentally measured χ_{zzz}/χ_{xxz} ratio of 3.85 \pm 0.35 for the isotope-labeled peak, we find $\theta_{\rm label}$ = 26 or 154 \pm 4° from the surface normal. To do this analysis, we have to use the ψ average because otherwise, there are too many unknowns; however, that approximation is not physically reasonable because rotating around the label's axis swings the entire peptide structure, since ψ_{helix} is not equal to ψ_{label} . Taking the internal rotational average about $\psi_{
m helix}$ for the unlabeled band (as done in the last section) makes sense if the helix is long and perfect because ψ averaging is equivalent to the symmetry modes. Thus, the analysis of a single label by itself makes no physical sense, and what needs to be done is to solve for tilt and twist angles that simultaneously give the correct ratios for the both the labeled and unlabeled peaks.

Although ψ -averaging an unlabeled helix is acceptable because of the symmetry, including the isotope label destroys the rotational symmetry, and the rotational average is no longer a valid approximation. This can be shown by calculating the χ_{zzz}/χ_{xxz} ratio for the label predicted by a fully internally

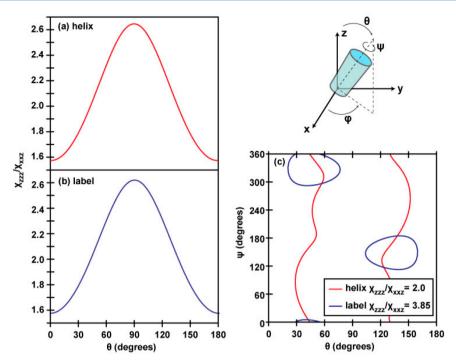


Figure 2. χ_{zzz}/χ_{xxz} ratios for (a) the unlabeled segment of an ideal α-helix and (b) the isotope-labeled peak, assuming full rotational averaging around the helix axis. (c) Contours indicating ψ/θ pairs giving the experimentally measured values for the unlabeled and labeled peak $\chi zzz/\chi xxz$ ratios. In this figure, the tilt angle θ indicates the angle between the helix axis and the surface normal.

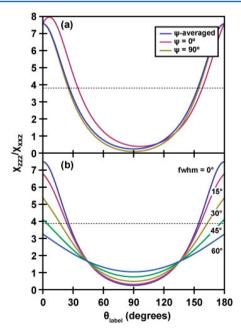


Figure 3. Tilt-angle dependence of the χ_{zzz}/χ_{xxz} ratio for a single amide-I residue, for (a) the ψ -averaged case and two fixed- ψ cases (ψ = 0 corresponds to the C(O)N bond lying in the yz plane) and (b) the ψ -averaged case for Gaussian distributions of tilt angles with different full-width-at-half-maxima (indicated). In parts a and b, the tilt angle θ is defined as the angle between the transition dipole and the surface normal.

rotationally averaged helix. When applied to the label, the orientational analysis above predicts that the χ_{zzz}/χ_{xxz} ratio of the label should match that of the unlabeled amide I band (Figure 2b) because symmetry dictates that in helical molecules, the responses of the parallel and perpendicular

normal modes (A and E1 modes) are closely related to the parallel and perpendicular components of a single local mode's response. ^{52,53} Thus, for a ψ -averaged helix tilted at 43 (or 137)° from the surface normal, one would also predict that a χ_{zzz}/χ_{xxz} ratio of 2.00 \pm 0.10 would be measured for the label, but the experimentally measured ratio is actually 3.85 \pm 0.35, which far exceeds the experimental uncertainty. In fact, a χ_{zzz}/χ_{xxz} ratio of 3.85 \pm 0.35 cannot be matched by a rotationally averaged helix at any tilt angle. That is, the ratios of the unlabeled and labeled amide I bands are incongruous for an analysis that requires rotational averaging around the helix axis.

3.4. Orientational Analysis of the Entire Peptide including Structural Disorder and Coupling. How does one reconcile the ratios of the labeled and unlabeled amide I bands? We need to consider the possibility that the helix is not isotropic about ψ_{helix} . That is to say, the helix does not have a free rotation around the principal axis. To determine the orientation of such an α -helix, we need to determine both the tilt angle and the twist angle. We retain a perfect helix with the label uncoupled from the rest of the helix and simultaneously calculate the χ_{zzz}/χ_{xxz} ratios for both the labeled and unlabeled peaks as a function of tilt and twist angles.

Using the two measured χ_{zzz}/χ_{xxz} ratios for the labeled and unlabeled peaks, we can deduce two unknowns: the tilt angle and the twist angle. As shown in Figure 2c, we find that our measured χ_{zzz}/χ_{xxz} ratios are consistent with helix tilt/twist angles of $(\theta, \psi) = (41^{\circ}, 5^{\circ})$, $(57^{\circ}, 297^{\circ})$, $(123^{\circ}, 117^{\circ})$, or $(138^{\circ}, 184^{\circ})$ (with around 4° error margins). Homodynedetected SFG cannot distinguish between the $(90 + x)^{\circ}$ and $(90 - x)^{\circ}$ tilt angles, 27 although future phase-sensitive or heterodyne-detected experiments may resolve this difficulty. However, we can narrow the possibilities using physical intuition and the experimentally observed solvent accessibility of Ile11.

Before rotation, our ideal helix is defined such that the hydrophilic region extends from -60° to $+120^{\circ}$ around the helix axis. Thus, the (41°, 5°) orientation corresponds predominantly to having the hydrophilic region of the peptide facing the polystyrene surface, whereas the (138°, 184°) orientation corresponds to having the hydrophobic region of the peptide facing the polystyrene surface. In the remaining two orientations, the hydrophobic and hydrophilic regions have roughly the same extent of interaction with the surface. Of these tilt-angle possibilities, the (138°, 184°) pair is most physically reasonable, since the solvent dependence of the peak center of the labeled unit indicates that the residue is buried at the hydrophobic interface, and the interaction of the hydrophobic side of the peptide with the hydrophobic polystyrene surface should inherently be more favorable than the interaction of the hydrophilic side of the peptide with this surface.

The following reasons further support that the $(138^{\circ}, 184^{\circ})$ orientation is the most likely orientation (Figure 4): First, this

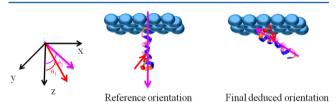


Figure 4. Schematic (right panel) showing the final deduced orientation of ovispirin (tilt angle =138°, twist angle =184°) at the polystyrene/water interface. The left panel defines the reference orientation (tilt angle = 0°, twist angle = 0°) where the hydrophilic region (blue) extends from -60° to $+120^{\circ}$ around the helix axis. θ_1 is the angle between the transition dipole of the ^{13}C =O chemical group (purple arrow) and the z axis, and θ_2 is the angle between the peptide helix axis (from N to C terminus) (red arrow) and the z axis. For the reference orientation, $\theta_1 = 138^{\circ}$ and $\theta_2 = 0^{\circ}$, and for the final deduced orientation, $\theta_1 = 23^{\circ}$ and $\theta_2 = 138^{\circ}$.

helix orientation also gives a χ_{zzz} (helix)/ χ_{zzz} (label) ratio of approximately 6.2, consistent with our experimental value of 6.18, a criterion that is not met by the $(57^{\circ}, 297^{\circ})$ and $(123^{\circ}, 117^{\circ})$ orientations (which both yield χ_{zzz} (helix)/ χ_{zzz} (label) ratios of less than 5). Second, at this orientation, the transition dipole of the isotope-labeled unit is calculated to be tilted 23° relative to the surface normal, in good agreement with the value calculated from our preceding analysis of the isolated label. To further clarify the final orientation of $(138^{\circ}, 184^{\circ})$, we also present the reference peptide orientation $(0^{\circ}, 0^{\circ})$ along with the final orientation in Figure 4.

Thus, analyzing the labeled and unlabeled peaks within the framework of an ideal helix provides insight into the overall orientation of the ovispirin peptide at the solution/polystyrene interface. However, analyzing the angles alone may not be sufficient because ovispirin is far from an ideal helix at the interface, which we learn from the observations listed below and leads to additional considerations about the interpretation of the spectra. First, comparing the intensities of the labeled and unlabeled amide I bands reveals that only a subset of the residues can contribute to the helical structure. For the (138°, 184°) tilt/twist angles calculated above, in which the angular dependence of the labeled and unlabeled peaks were considered independently, a full 18-residue helix would give a χ_{zzz} (helix)/ χ_{zzz} (label) value of 7.0, which deviates by almost 15% from the experimental value of 6.18. Reducing the length

of the helical section to anywhere between 13 to 15 residues gives a $\chi_{zzz}(\text{helix})/\chi_{zzz}(\text{label})$ value between 6.1 and 6.3. Further reducing the length of the helical section causes the $\chi_{zzz}(\text{helix})/\chi_{zzz}(\text{label})$ ratio to drop far below the experimentally measured value. Although this result is consistent with the 13-residue helix analyzed above, it reveals that the entire 18-residue ovispirin peptide cannot form a perfectly ordered helix, because if it did, then the intensity of the isotope label would be much smaller, as compared with the unlabeled amide I band (intensity is proportional to the square of the χ value).

We want to highlight here that the Hamiltonian approach we present in this paper can be used to calculate the coupling between the label and the rest of the peptide, and including these couplings in the Hamiltonian disrupts the transition dipoles of the helix, which changes the strength of many of the SFG-active normal modes. Shown in Figure 5a,b are spectra for

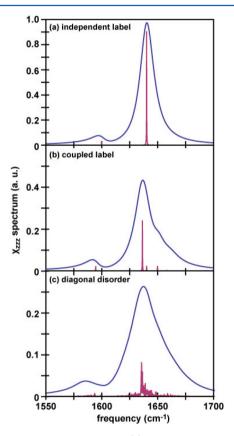


Figure 5. Simulated spectra showing (a) a 13-residue ideal helix, with the isotope label calculated independently of (or uncoupled from) the rest of the helix, (b) a 13-residue ideal helix with the isotope label incorporated into and coupled with the other residues in the helix, and (c) a 13-residue ideal helix summed over 50 spectra with 16 cm⁻¹ random disorder in the local mode frequencies. Spectra were simulated using a 15 cm⁻¹ Lorentzian for each mode. Stick spectra (purple) are also included to emphasize how the normal mode structure changes as the couplings are broken.

a 13-residue labeled helix at the (138°, 184°) tilt/twist angle. Figure 5a shows the response of an isolated label added to that of an undisrupted helix, whereas Figure 5b shows the spectrum when the isotope label is incorporated into and coupled with the rest of the helix. The stick spectrum shows significant contributions from new SFG-active modes, indicating that the label and the helix cannot be treated as independent vibrational modes. In this regard, ¹³C=¹⁸O isotope-labeling is preferable,

because the labeled amide group has better decoupling, although both $^{13}C=^{16}O$ and $^{13}C=^{18}O$ labels will affect the unlabeled spectrum about the same amount. This effect will be important when the isotope-labeled unit is incorporated into different sites of the peptide or more than one isotope-labeled unit is incorporated.

Second, the Hamiltonian approach can be used to access the effect of the structural disorder. Random disorder along the diagonal of the Hamiltonian occurs in many systems in which the individual chromophores have different local mode frequencies due to different local environments in different molecules; for example, if the local solvation environment of the amide-I modes varies from amide to amide bond. 54,55 In Figure 5c, we show the spectrum expected for a 13-residuelabeled helix with 16 cm⁻¹ standard deviation random variation in the local mode frequencies added before diagonalizing the Hamiltonian, averaged over 50 different disordered Hamiltonians; 16 cm⁻¹ is about the disorder one expects due to differences in hydrogen bonding for soluble polypeptides. We see that the SFG spectrum of the disordered helix is broader and less intense than that of the ordered helix because all different members of the ensemble now have slightly different normal modes. The overall calculated tilt angle from the labeled and unlabeled peaks does not change significantly with this added disorder, but the helix/label ratio increases. Thus, disorder that is created by the typical environment surrounding a polypeptide does not change the interpretation of the angular measurements, but does alter the relative intensities of the labeled and unlabeled amide I modes. At this stage, we cannot quantify this effect in our current measurements because it requires dynamic knowledge of the linewidths, which will be the subject of future work. Nonetheless, the point being made is that disorder breaks the symmetry of the helix, and so in real systems, one may no longer think of the normal modes purely as the fully symmetric A and E1 modes, especially for a severely disrupted helix.

Interestingly, the tilt angles deduced in this work (138°) indicate that the ovispirin is much more steeply angled away from the plane of the polystyrene surface than it is when adsorbed to a planar lipid bilayer. Using multiple label sites in future experiments will provide residue-by-residue structural information that can be used to obtain a complete backbone structure through constrained molecular dynamics simulations, as has been done with oriented polypeptides and FTIR spectroscopy. ⁵⁶

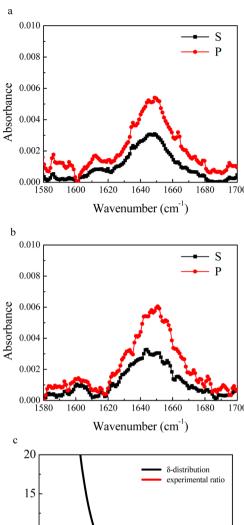
Finally, we want to mention a few other considerations that should be taken into account in future simulations. As revealed by previous NMR, 2D IR, and MD studies, ovispirin exhibits significant structural disorder, as well. For example, in the solution NMR studies, it is found that the helical portion of the peptide is somewhat curved, and the N- and C-terminal ends are disordered. The curve breaks the symmetry of the helix and adds off-diagonal disorder to the Hamiltonian. The disordered ends will have both diagonal and off-diagonal disorder. Therefore, such structures may contribute to SFG signals. In the future, we will calculate such contributions and validate the calculation method using multiple label sites in experiments. In addition, we calculated the transition dipole and Raman response of a single amide-I mode using values derived from polarized FTIR and Raman experiments.²³ However, some models use slightly different molecular responses for a single amide-I mode, \$55,56 which will yield slightly different singleresidue tilt angles. We have also assumed a δ -function

orientational distribution. If the residues or peptides are instead distributed over a wider range of tilt angles, the tilt-angle dependence flattens out, as shown in Figure 3b for a series of Gaussian tilt-angle distributions of varying widths for a single vibrationally isolated label. This flattening of the tilt-angle dependence means that the same experimental χ_{zzz}/χ_{xxz} ratio yields a center tilt angle closer to the surface normal as the distribution gets broader. Developments underway on modeling SFG spectra of peptides will help to better define these parameters. $^{57-59}$

3.5. ATR-FTIR Experiments. For comparison purposes, we also performed ATR-FTIR experiments. The concentration of ovisiprin-1 is 10 times higher because no signal was detected from the isotope-labeled group when lower peptide solution concentrations were used. The ATR-FTIR spectra collected from the PS/peptide solution interface are shown in Figure 6. We found that the PS polymer can generate a strong background ATR-FTIR signal at ~1602 cm⁻¹ from one of the benzene ring modes, ⁶⁰ appearing as either a positive or negative peak. For both isotope-labeled and regular ovispirin-1 samples, there is a broad peak at around ~1647 cm⁻¹ (Figure 6), indicating the existence of both α -helical and random coil components. For isotope-labeled sample, in addition to the negative peak at ~ 1602 cm⁻¹, an additional peak at ~ 1612 cm⁻¹ can be assigned to the isotope-labeled C=O chemical group. Data analysis shows that this C=O is ~38° vs the surface normal (Figure 6c) given that the dichroic ratio R of the isotope-labeled C=O stretching peak is ~4.13.61 Interestingly, the 1602 cm⁻¹ background signal of the PS polymer is silent in SFG spectra because this benzene ring mode is Raman-inactive and, thus, SFG-inactive. In addition, ATR-FTIR spectra have multiple contributions from the α -helix, the random coil, and the polymer background that made the orientation determination of the single C=O chemical group by ATR-FTIR subject to large errors.

Compared with ATR-FTIR, SFG studies on the C=O isotope-labeled peptides have several advantages: (1) SFG spectra are free of background contribution from the polymer substrate as well as other media in our study. For peptides and proteins, SFG spectra have minimum contributions from random coils. Unlike in IR spectra, where $\rm H_2O$ absorption band overlaps with the peptide/protein amide I band, SFG spectra have a minimum $\rm H_2O$ signal contribution, and thus, the experiment could be performed in $\rm H_2O$ and separate the signals from the isotope unit and the main amide I peak. (2) SFG is more sensitive than ATR-FTIR for detecting well-oriented α -helices; thus, it can detect signals from peptides from a single monolayer instead of bilayer stacks.

Although solid-state NMR (e.g., PISEMA) is a well-established technique in obtaining orientation restraints for membrane-bound peptides, 62 isotope-labeled SFG has some unique advantages: (1) SFG can monitor the interaction process in situ (~several minutes per spectrum), whereas NMR techniques such as PISEMA suffer from long-time data accumulation (~several hours per experiment); therefore, the isotope-labeled SFG technique can be used to monitor the dynamics of a single residue in biological processes, such as ligand titration, fibril formation, GPCR-G protein interaction, ion channel opening, and so on. (2) Only one bilayer is required, which allows very precise difference experiments, such as with ligand binding. (3) A typical solid-state NMR experiment on short peptides requires an ~100 μ g sample, whereas an SFG experiment only needs an ~10 μ g sample.



15

α 10

5

0 10 20 30 40 50 60 70 80 90

Angle(θ), degree

gure 6. (A) ATR-FTIR spectra of isotope-labeled

Figure 6. (A) ATR-FTIR spectra of isotope-labeled ovispirin-1 molecules at the PS/peptide solution interface. (B) ATR-FTIR spectra of regular ovispirin-1 molecules at the PS/peptide solution interface. (C) The relationship between the dichroic ratio, *R*, of the isotope-labeled C=O stretching peak and the tilt angle of the C=O bond direction relative to the surface normal. The dichroic ratio, *R*, of the isotope-labeled C=O stretching peak detected at the PS/peptide solution interface is shown as a horizontal line.

For ovisprin-1, we also detected SFG signals from the α -helical backbone and the isotope-labeled peptide unit when the peptide associated with a supported lipid bilayer DPPG (dipalmitoylphosphatidyglycerol)/dDPPG (deuterated dipalmitoylphosphatidyglycerol) in contact with the peptide aqueous solution. For ovispirin-1 with residue R4 mutated with $^{13}\text{C}{=}^{16}\text{O}$ incorporated glycine, the SFG ppp spectrum shows a signal contributed from the isotope-labeled unit at $\sim\!1610$ cm $^{-1}$ (Supporting Information), similar to that shown in Figure 1b. On the other hand, for the ovispirin-1 with R8 mutated with

 13 C= 16 O incorporated glycine, the SFG signal detected from the isotope-labeled unit shifted to $\sim 1600~\rm cm^{-1}$ with a broader width (Supporting Information). The different peak centers and peak widths indicate that R4 and R8 are in different microenvironments. Thus, we expect that this technique is applicable to peptides associated with lipid bilayers or cell membranes. However, since more experimental measurements and theoretical analysis are still under way and such details are beyond the scope of this research, the detailed structure and orientation analysis will be published in the future.

In this paper, we demonstrated the approach on a synthesized polypeptide that is isotope-labeled using commercially available compounds, but the method is also applicable to proteins using exciting methods developed in the past few years. One set of approaches is native chemical ligation and expressed protein ligation. 63,64 In these techniques, proteins are semisynthesized from fragments using a thiol-ene chemistry. Thus, one can isotope-label a fragment (e.g., a helix instead of a single peptide unit) in the same manner as we did here for ovispirin and then ligate it to another fragment to form the full protein. Expressed protein ligation has recently been used in conjunction with 2D IR spectroscopy.⁶⁵ It is also possible to arbitrarily isotope-label amino acids in proteins by using a cellfree expression system and a stop codon that matches an isotope-labeled tRNA.66 It is now also possible to put in nonnatural isotope labels that might be SFG chromophores, such as nitriles, by using a tRNA synthetase pair system, such as that which has been developed for phenylalanine.⁶⁷ With these new and exciting methods in hand, it is now possible to isotope-label and, thus, monitor the structures of precise locations in nearly arbitrarily sized proteins. It is also important to point out that SFG spectroscopy requires a tiny amount of sample. For example, as mentioned above, in the ovispirin-1 case, one SFG experiment requires only 20 µg sample. As a result, experiments are possible even on poorly expressing proteins or systems in which the ligation is inefficient. Furthermore, heterodyne detection can also improve the signal-to-noise ratio, $^{46-49}$ and thus, we believe that in the future, this technique will shed light on larger protein systems.

4. CONCLUSION

We have shown that isotopic labeling enables SFG spectroscopy to detect polarized amide I signals from a single peptide unit in an α -helical peptide at a buried polymer/solution interface. Such SFG signals can be used to study the microenvironment and interfacial orientation of the isotopelabeled residue. If we assume a single orientation distribution, the backbone helix of ovispirin-1 on the PS surface is 138° relative to the surface normal with a twist angle of 184°, and the transition dipole of the isotope-labeled C=O group is tilted 23° relative to the surface normal. Our result shows that using ¹³C=O isotope-labeling, SFG can be used to obtain residuespecific orientation information. In the future, using multiple label locations will allow us to measure the exact structure and orientation of surface-bound peptides without requiring assumptions about the expected secondary structure. This capability makes SFG a powerful technique for structural analysis and will bring new insight to many biophysical systems.

ASSOCIATED CONTENT

S Supporting Information

Detailed SFG data analysis and Hamiltonian approach. This information is available free of charge via the Internet at http://pubs.acs.org.

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Notes

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