

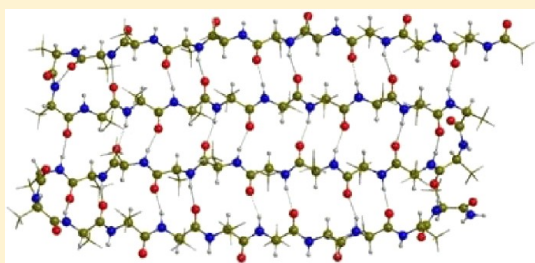
# The Folding of Acetyl(Ala)<sub>28</sub>NH<sub>2</sub> and Acetyl(Ala)<sub>40</sub>NH<sub>2</sub> Extended Strand Peptides into Antiparallel $\beta$ -Sheets. A Density Functional Theory Study of $\beta$ -Sheets with $\beta$ -Turns

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## S Supporting Information

**ABSTRACT:** We report ONIOM calculations using B3LYP/D95\*\* and AM1 on  $\beta$ -sheet formation from acetyl(Ala)<sub>N</sub>NH<sub>2</sub> ( $N = 28$  or  $40$ ). The sheets contain from one to four  $\beta$ -turns for  $N = 28$  and up to six for  $N = 40$ . We have obtained four types of geometrically optimized structures. All contain only  $\beta$ -turns. They differ from each other in the types of  $\beta$ -turns formed. The unsolvated sheets containing two turns are most stable. Aqueous solvation (using the SM5.2 and CPCM methods) reduces the stabilities of the folded structures compared to the extended strands.



## ■ INTRODUCTION

Several reports from our group<sup>1–3</sup> and others<sup>4–9</sup> have dealt with molecular orbital studies of the energetics and structures of both parallel and antiparallel  $\beta$ -sheets. However, they have generally focused upon the aggregation of two or more strands into sheet-like structures, thereby neglecting the effects upon the stability attributable to the turns that are generally present in antiparallel  $\beta$ -sheets. In this paper, we report ONIOM DFT/AM1 calculations upon two sets of isomeric conformations of Acetyl(Ala)<sub>28</sub>NH<sub>2</sub> and Acetyl(Ala)<sub>40</sub>NH<sub>2</sub> (hereafter, referred to as A28 and A40, respectively) that contain different numbers of turns and strands. We used two different length peptides to obtain some information upon the effect of the peptide length upon the energetics and structures of the various isomeric structures. We chose the peptide lengths (24 and 40 residues) since these allow for folding into sheets with equal number of residues for most structures.

The  $\beta$ -sheet structures considered contain from one to four turns for A28 and one to six turns for A40. Each sheet has one more strand than the number of turns. We chose conformations that have close to the same number of residues per strand of the sheet as they provide the most H-bonds within the sheet. We found several different conformations of the sheets that differ in the nature of the turns, but focus on the lowest enthalpy structures.

## ■ CALCULATIONAL DETAILS

We used the ONIOM<sup>10,11</sup> method as programmed in the Gaussian 09<sup>12</sup> suite of computer programs. ONIOM divides the system into up to three segments, which can be treated at different levels of calculational complexity. Thus, one can treat the essential part of the system at the high level, while the less critical parts of the system might be calculated at the medium or low level. For this study, we only used two levels (high and

medium). We treated the backbones of each of the peptides at the high level, with only the side chains (methyls) at the medium level. The high level used hybrid DFT methods at the B3LYP/D95(d,p) level. This method combines Becke's 3-parameter functional,<sup>13</sup> with the nonlocal correlation provided by the correlation functional of Lee, Yang, and Parr.<sup>14</sup> In the ONIOM method, there are unsatisfied valences in the high level at the interface between it and medium level. These valences were satisfied by using the default method of capping them with a hydrogen atom in the direction of the connecting atom in the medium level with a C–H distance of 0.723886 times the C–C distance. We used the AM1<sup>15</sup> semiempirical molecular orbital method for the ONIOM medium level.

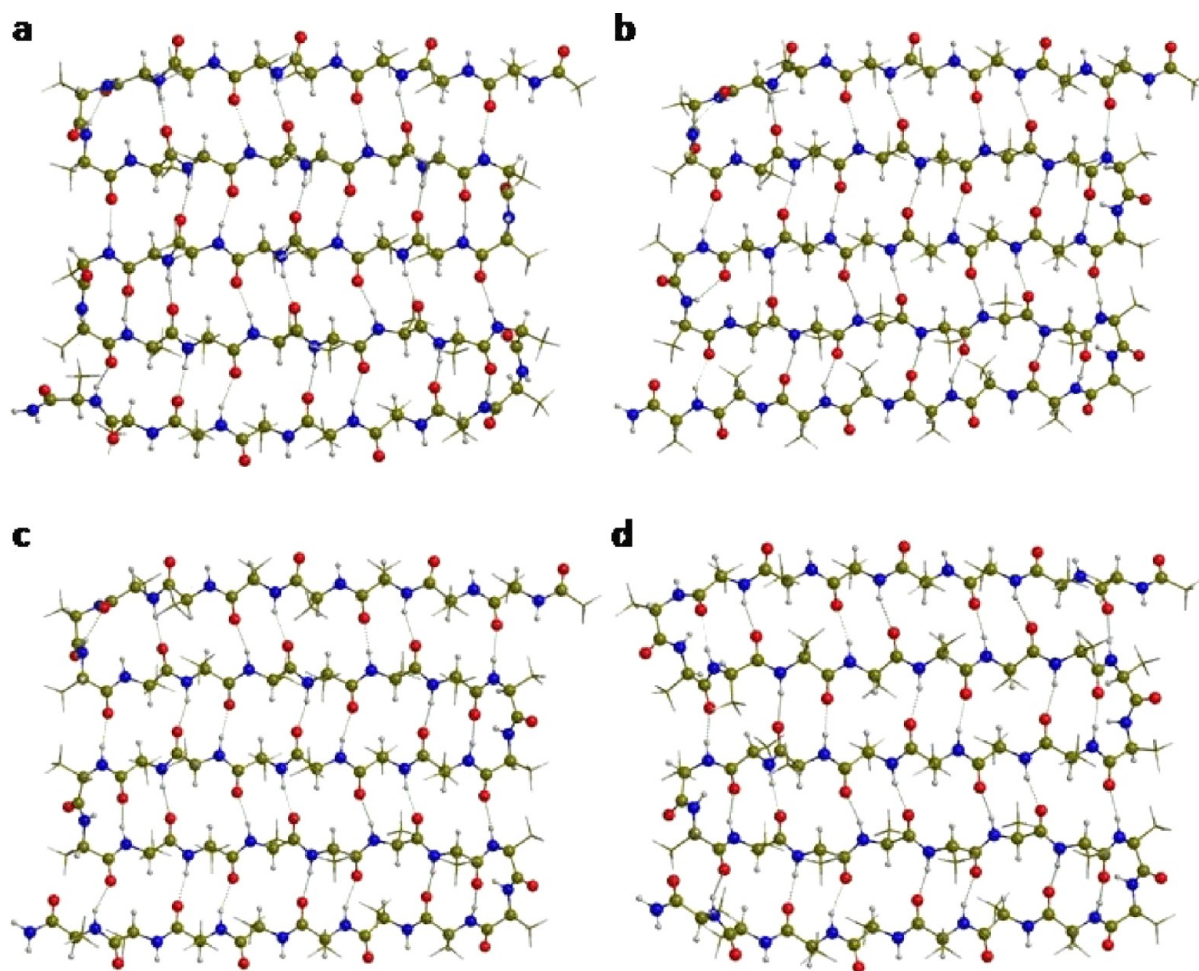
With the ONIOM method, Gaussian performs a high level (density functional theory, DFT) calculation only on those atoms designated with the valences terminated as described above.

All geometries were completely optimized in all (up to 1227) internal degrees of freedom. Vibrational calculations confirmed that the reported geometries are true minima on the potential energy surfaces (PESs), as there are no imaginary vibrational frequencies. We used these frequencies to calculate the enthalpies and free energies of the optimized species. We evaluated the counterpoise corrections (CP) for basis set superposition error (BSSE) only for the H-bonds between the strands using the single point a posteriori procedure. To do this, we cut the C–C bonds that connect strands to the turns and terminated them with H's. We then calculated the CP for each individual strand using the ghost orbitals of all the other strands and took the sum as the total CP following the

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**Figure 1.** Examples of the four kinds of structures found using A40 (presented in increasing energetic order): the lowest energy structure, **a**, contains C<sub>7</sub> H-bonds in the 1st and 4th turns and C<sub>10</sub> H-bonds in the 2nd and 3rd; structure **b** contains C<sub>7</sub> H-bonds in the 1st and 3rd turns and C<sub>10</sub> H-bonds in the 2nd and 4th; structure **c** has only one C<sub>7</sub> H-bonding turn, the 1st and C<sub>10</sub> H-bonds in the others, while the highest energy structure, **d**, contains C<sub>10</sub> H-bonds in all turns.

procedure we have previously suggested.<sup>16</sup> We did not apply the CP-opt procedure<sup>17</sup> (which optimizes on a CP corrected surface) as this would not be possible after the turns are cut away.

In a previous study of five 17-amino acid peptides,<sup>18</sup> we found little difference in relative energies between this procedure and another where the side chains (in this case, the methyls) were subsequently optimized using DFT, with the (previously optimized) peptide chain held fixed. The current procedure also gave relative energies that agreed well with complete DFT optimizations for a series of five small 3<sub>10</sub>-helical peptides.<sup>19</sup> We recently found this procedure to give reasonable results for peptides when compared to results from experimental databases, while several DFT methods that are specifically parametrized to include dispersion do not provide results in accord with experiments.<sup>20</sup> We believe this may be due to the interaction of induction and dispersion, as we have discussed elsewhere.<sup>21</sup> B3LYP/D95\*\* also gives reasonable results for H-bonding interactions.<sup>22,23</sup>

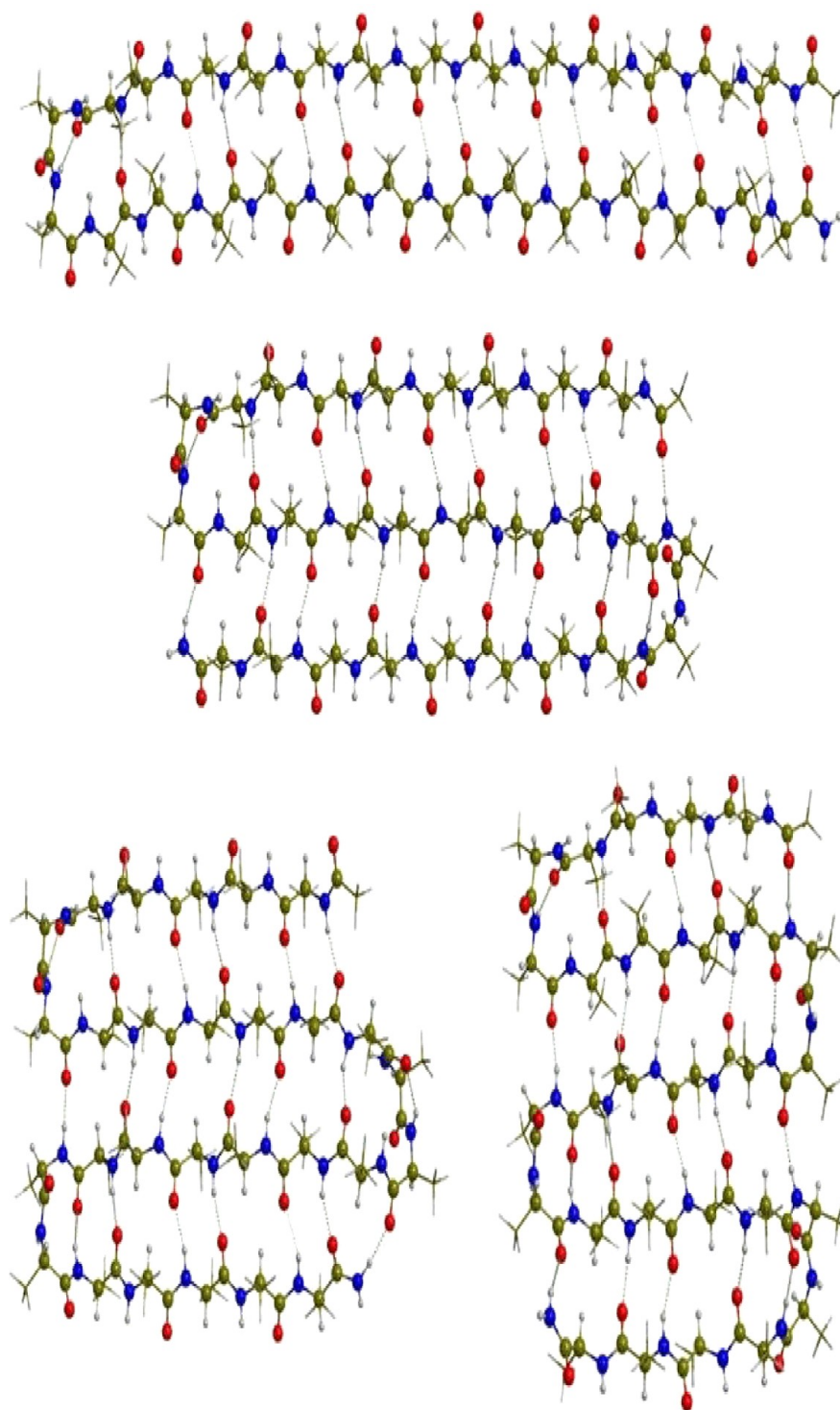
We calculated free energies using both the SM5.2<sup>24</sup> and CPCM<sup>25,26</sup> models, and geometrically optimized using CPCM with the Pauling radii.<sup>27</sup> We used the AMPAC 8.16<sup>28</sup> program to calculate single-point AM1 energies and solvation free energies of the peptides in their ONIOM optimized structures as we have done previously.<sup>29–31</sup> We used only the solvation

free energies from these calculations. Solvated enthalpies combine the ONIOM  $\Delta H$ 's with the SM5.2 solvation  $\Delta G$ 's, while solvated free energies combine the ONIOM  $\Delta G$ 's with the SM5.2 solvation  $\Delta G$ 's. We followed the (optimized) CPCM, calculations with vibrational calculations to verify the minima and to obtain thermodynamic properties. We added the CPCM solvation free energies to the  $\Delta H$ 's and  $\Delta G$ 's obtained from the vibrations at the optimized CPCM geometries to obtain the solvated  $\Delta H$ 's and  $\Delta G$ 's reported.

## RESULTS AND DISCUSSION

We found four different kinds of local minima for the  $\beta$ -sheets containing turns, which we have illustrated for A40 and labeled (in order of stability) **a**, **b**, **c**, and **d** in Figure 1. All four contain only  $\beta$ -turns using the definition that the distance between the  $\alpha$ -carbons of residues  $i$  and  $i+3$  be less than 7 Å and the structure not be helical,<sup>32,33</sup> which is still widely used.<sup>34</sup> However, they differ in the types of  $\beta$ -turns present as seen from the H-bonding topologies in the turns.

Venkatachalam<sup>35</sup> originally classified  $\beta$ -turns as types I, I', II, and II' according to their  $\phi$  and  $\psi$  dihedral angles at  $i+1$  and  $i+2$ . This classification has been widely used since.<sup>33</sup> However, this classification requires both C<sub>14</sub> and C<sub>10</sub> H-bonds between the  $i$ th and  $i+3$ rd residues in the turn in all these types and that



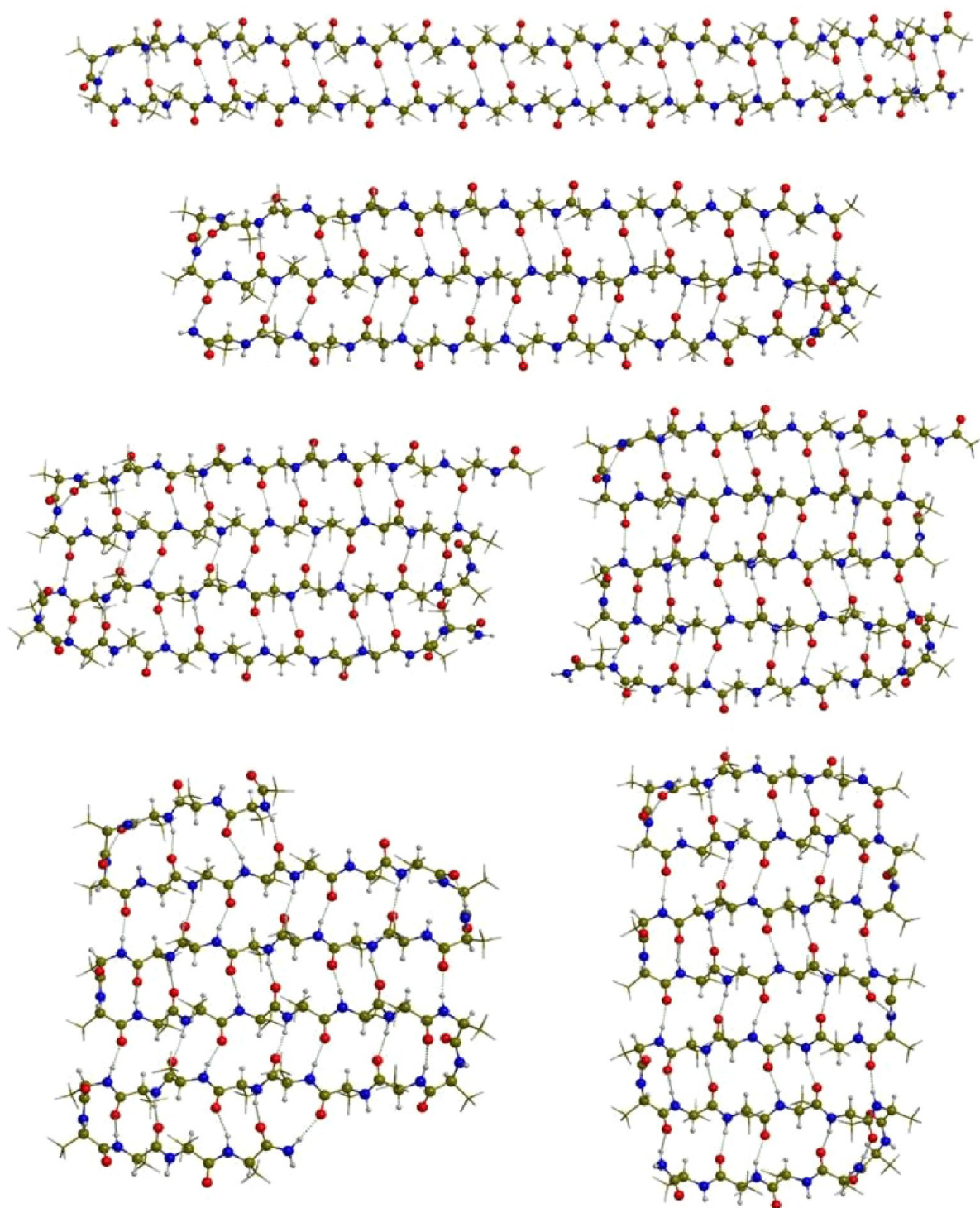
**Figure 2.** Structures of A28 containing one to four  $\beta$ -turns.

the  $i$ -first and  $i$ -second residues of type II or II' have opposite chirality (or have one glycine which is achiral). Each turn contains up to three C=O's, one of which cannot form an H-bond due to its conformation. The other two have the potential to form H-bonds. We note here that we and others have previously found C<sub>10</sub> H-bonds to be particularly weak.<sup>3,6</sup> We have recently proposed a new classification of  $\beta$ -turns, and we shall apply that classification here.<sup>36</sup> This classification divides  $\beta$ -turns into two general classes: type A (which contain a C<sub>10</sub> H-bond between the  $i$ th and  $i+3$ rd residues) and type B (which

lack this C<sub>10</sub> H-bond, but may contain one or more C<sub>7</sub> H-bonds). The C<sub>7</sub> H-bonds that we observed may be more stable than the C<sub>10</sub> (C<sub>N</sub> H-bonds refer to cyclic H-bonding systems that contain N atoms).

The most stable, **a**, structure contains only B<sub>I</sub> and A<sub>II</sub> turns, which are the most stable structures for isolated Type A and B turns, respectively.<sup>36</sup> While the A<sub>II</sub> turns have  $\phi$  and  $\psi$  dihedral angles close to those that are used to define type II' turns in the Venkatachalam classification,<sup>36</sup> the type II' turn requires that the  $i+1$ st residue to be D (or glycine).<sup>33</sup> Both the **b** and **c**





**Figure 3.** Structures of A40 containing from one to six turns.

structures contain only  $A_I$  and  $B_{II}$  turns, while **d** contains mostly  $A_I$  and  $A_{III}$  turns (there is one  $B_{II}$  turn in the A40 structure containing five turns). Thus, the  $C_{10}$  H-bonds appear in only some of the turns of the **a**, **b**, and **c** structures. We have calculated each kind of sheet that we report using each of these geometries for the turns (which are all completely optimized to local minima). In each case we have found the **a** type turn to give the lowest energy structures. For this reason, we consider only the **a** structures in the following discussion unless

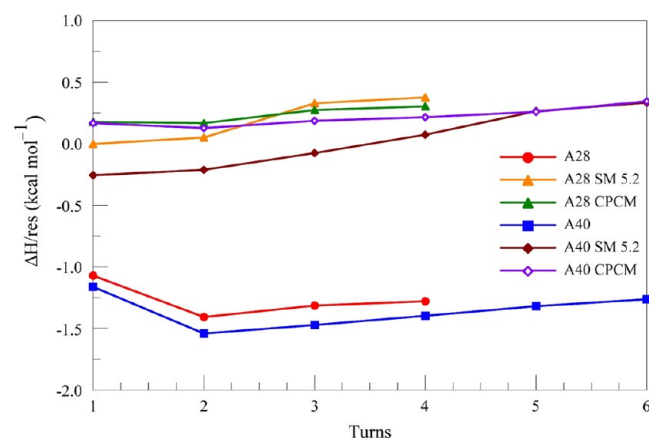
specifically noted. The data for the **b**, **c**, and **d** structures can be found in the Supporting Information.

The lowest energy structures found for A28 conformations containing from one to four turns and for A40 containing one to six turns appear in Figures 2 and 3. We tried to keep the numbers of residues in each strand as close to equal as possible to make the numbers of interstrand H-bonds follow the number of turns. We confirmed all unsolvated structures and all but three structures optimized with CPCM as true minima by calculating the vibrational frequencies. Two of the remaining

three have three low imaginary frequencies, and the other has seven. Repeated attempts at optimization of these three solvated structures failed to find true minima. Since the imaginary frequencies are all quite small ( $<15\text{ cm}^{-1}$ ), we do not expect the energies to be substantially affected. However, these imaginary frequencies do affect the enthalpies and free energies calculated from the vibrations. Since the minimum contribution to the molar enthalpy for a low frequency vibration will be  $RT$ , we added this for each imaginary frequency to the calculated enthalpies. Table 1 displays the energetic folding data for both

**Table 1. Energetics of Folding and Solvation (versus Extended Strand) Using (Single Point) SM5.2 and (Optimized) CPCM**

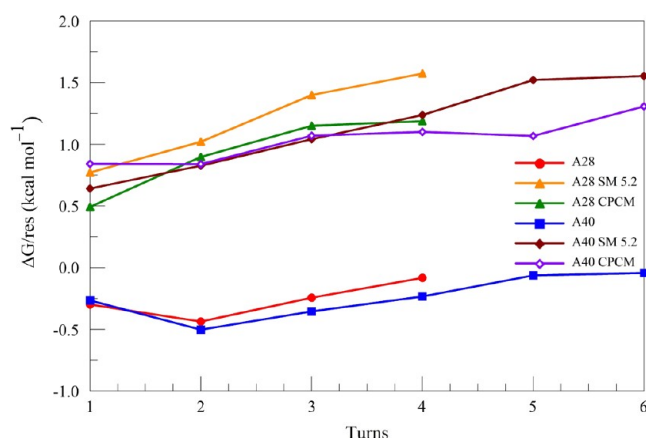
turns	H-bonds	$\Delta E$	$\Delta H$	$\Delta G$	$\Delta G_{\text{solv}}$	
					SM5.2	CPCM
A28						
1	14	−33.5	−30.0	−8.4	30.0	22.9
2	18	−45.8	−39.4	−12.2	40.8	39.5
3	20	−43.3	−36.8	−6.8	46.0	41.4
4	20	−42.2	−35.9	−2.3	46.4	38.2
A40						
1	20	−51.9	−46.5	−10.6	36.3	45.2
2	26	−70.7	−61.7	−20.1	53.2	56.2
3	27	−68.2	−58.6	−14.2	55.8	59.3
4	28	−65.3	−55.9	−9.3	58.8	56.2
5	30	−62.9	−52.8	−2.5	63.4	48.6
6	30	−60.6	−50.6	−1.7	63.9	58.9



**Figure 4.**  $\Delta H$ 's (relative to extended strands) versus number of turns with and without solvation (indicated as SM5.2 or CPCM).

polyalanines, and Figures 4 and 5 present the differences (per residue) in the  $\Delta H$ 's and  $\Delta G$ 's between the folded and extended structures (with and without solvation) for the A28 and A40. As seen from Figures 4 and 5, the most stable sheets for both A28 and A40 contain two turns (three strands). The small enthalpic preference for two versus three turns for unsolvated A28 increases for A40, which suggests that this property will remain for larger systems

For the unsolvated structures, the  $\Delta H$ 's per residue become less negative as the number of turns increases beyond two, as seen from Figure 4. Figure 5 displays a similar trend for  $\Delta G$ . The foregoing results from an increase in the number of H-bonds with the number of turns while each H-bond becomes



**Figure 5.**  $\Delta G$ 's versus extended strands with and without solvation free energies (marked SM5.2 or CPCM).

less stabilizing due to the strain induced by the turns. The increased number of H-bonds dominates for small numbers of turns, but the decrease in stability per H-bond dominates for larger numbers of turns. We have included detailed data on the H-bond lengths and average H-bond stabilities for each sheet in the Supporting Information.

Aqueous solvation tends to decrease the stability of the sheets versus the extended strands. One might expect this behavior as the extended strands have more solvent-exposed H-bond acceptors and donors. The number of exposed acceptors and donors decreases with increasing numbers of turns. As seen from Figures 4 and 5, the energetic behavior of the optimized CPCM solvated structures and those calculated using (single-point) SM5.2 differ somewhat. While both methods show the sheets to be less solvated than the extended strands, we see from the  $\Delta H$ 's of Figure 4 that, compared to CPCM, SM5.2 better solvates sheets with fewer turns and more poorly solvates those with more turns. In fact, the plots for the solvated enthalpies versus number of turns cross for both A28 and A40 in Figure 4.

Comparison of the enthalpic and free energy data (Table 1 and Figures 4 and 5) suggest that the increasing rigidity of the sheets over the strands decreases the relative entropy of the former. The unsolvated A28 with one turn has a slightly lower  $\Delta G$ /residue than the analogous structure for A40. We note that calculating the entropy (and thus the free energy) of these structures from the harmonic vibrations can introduce errors. These errors become particularly problematic for low barrier torsional modes, such as those involved in methyl group rotations. As we only use relative free energies, we expect the methyl rotational errors to cancel as the number of such modes does not change upon folding.

The solvation free energies cannot readily be calculated from first principles. When quantum mechanical methods are used, solvation generally is calculated using a continuum model of some kind that relies on empirical parameters. The other popular alternative, molecular dynamics, generally depends upon parametrized force fields and does not usually rely on quantum mechanical first principles. Such methods can be useful, but treat the problem from an entirely different approach.

The solvation  $\Delta G$  relative to the extended strand becomes more positive as the number of turns increases, presumably due to the decrease in H-bond donors (N-H's) and acceptors (C=O's) that are exposed to the solvent as more H-bonds

within the sheets form with increasing numbers of turns. Since the A40 has more exposed H-bond donors and acceptors than A28 for the same number of turns, the aqueous solvation stabilization for the former is always greater. Large  $\beta$ -sheets without polar side chains tend to become insoluble as they become larger, which is exemplified by silk and amyloids. The observed trend in  $\Delta G_{\text{solv}}$  accords with this observation.

We note that the results presented here apply to polyananines and that sheets containing other residues may be stabilized by specific interactions between residues, i.e., glutamines.<sup>37–40</sup>

## CONCLUSIONS

Unsolvated polyananines containing 28 and 40 residues prefer  $\beta$ -sheets containing two turns (three strands) enthalpically. Considering the free energy reduces the stabilities of the sheets with increasing numbers of turns. Not surprisingly, the results suggest that longer peptides can sustain  $\beta$ -sheets with more turns. Aqueous solvation reduces the stabilities of the folded sheets in comparison with the extended strand.

## ASSOCIATED CONTENT

### Supporting Information

Complete data including structures **b**, **c**, and **d** and complete ref 12. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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### Notes

The authors declare no competing financial interest.

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