

Conformational flexibility in mammalian 15S-lipoxygenase: Reinterpretation of the crystallographic data

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ABSTRACT

Lipoxygenases (LOXs) are a family of non-heme iron dioxygenases that catalyze the regioselective and stereospecific hydroperoxidation of polyunsaturated fatty acids, and are involved in a variety of inflammatory diseases and cancers. The crystal structure of rabbit 15S-LOX1 that was reported by Gillmor *et al.* in 1997 has played key roles for understanding the properties of mammalian LOXs. In this structure, three segments, including 12 residues in the superficial $\alpha 2$ helix, are absent and have usually been described as “disordered.” By reinterpreting the original crystallographic data we were able to elucidate two different conformations of the molecule, both having well ordered $\alpha 2$ helices. Surprisingly, one molecule contained an inhibitor and the other did not, thereby adopting a closed and an open form, respectively. They differed in the conformation of the segments that were absent in the original structure, which is highlighted by a 12 Å movement of $\alpha 2$. Consequently, they showed a difference in the size and shape of the substrate-binding cavity. The new model should provide new insight into the catalytic mechanism involving induced conformational change of the binding pocket. It may also be helpful for the structure-based design of LOX inhibitors.

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Key words: crystal structure; ligand-induced conformational change; arachidonate binding mode; crystal twinning; pseudo symmetry.

INTRODUCTION

Lipoxygenases (LOXs) are nonheme iron dioxygenases that catalyze the regioselective and stereospecific hydroperoxidation of polyunsaturated fatty acids.¹ LOXs are found in a large variety of organisms, such as bacteria, plants, and mammals.^{2–4} In plants, LOXs are involved in germination and senescence and their typical substrate is linoleic acid or α -linolenic acid.⁴ In mammals, LOXs together with cyclooxygenases are key enzymes in the arachidonic acid cascade.⁵ LOXs produce hydroperoxyeicosatetraenoic acids, precursors of leukotrienes and lipoxins that have been implicated as critical signaling molecules in a variety of inflammatory diseases and cancers.^{6–9}

The crystal structures of soybean LOX-L1¹⁰ and LOX-L3,¹¹ rabbit 15S-LOX1,¹² and coral 8R-LOX¹³ have been elucidated thus far and helped understanding the properties of LOXs at the molecular level. Several structures of soybean LOX in complex with its product or inhibitor are also known.^{14–18} There are significant differences in size, sequence, and substrate preference between the plant and animal LOXs, but the overall fold and geometry of the nonheme iron-binding site are conserved.^{19,20} All LOXs are folded in a two-domain structure that is composed of a smaller N-terminal β -barrel and a larger α -helical catalytic domain. The nonheme iron essential for activity is positioned deep in a large cavity that accommodates the substrate. The regio- and stereospecificities of the various LOX isozymes are believed to be determined by the shape and depth of the cavity as well as the binding orientation of the substrate in the cavity.^{21–23}

The crystal structure of 15S-LOX1 in complex with inhibitor was reported by Gillmor *et al.*¹² Ever since, it has served key roles in understanding the specificities of mammalian LOXs as well as designing structure-based inhibitors.^{7,24,25} This structure was determined in the trigonal space group R32 at 2.4 Å resolution. In the PDB model (1lox), the atomic coordinates for many residues (177–188, 210–211, and 601–602) are not available and a superficial α -helix ($\alpha 2$) exists as disconnected even though residues 177–188 could have been modeled into it bridging the gap. According to the original article, “There is no density for surface residues 210–211 and 601–602. The density for residues 177–188 is ‘ambiguous’

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due to the nearby presence of a twofold crystallographic symmetry operator.” In subsequent articles citing this structure, this α -helix has been described as “disordered,” although it is very unlikely that a segment within a long α -helix is disordered.

Recently, Oldham *et al.* compared the structure of 8R-LOX with that of 15S-LOX1 and noted some significant discrepancies.¹³ They pointed out that the crystal packing constraints were violated in the “disordered” helical segment in 15S-LOX1 structure. We found that the two helices related by a crystallographic twofold axis on the diagonal between the *a* and *b* axes collided with each other [Fig. 1(a)], while the electron density map showed a contiguous density accommodating both ends of the helix [Fig. 1(b)]. This suggests that residues 177–187 were absent not because they were disordered but because they could not be modeled within the constraints of crystal packing. As the original model cannot be compatible with the crystal packing constraint, we have sought out alternative models. In this study, we reinterpreted the original crystallographic data using the R3 twin model and obtained the complete structure of two independent molecules that partly assumed different conformations. The new model is sound crystallographically in that there is no crystal packing problem and at the same time the “disordered” $\alpha 2$ helix is well resolved. It should provide new insight into the catalytic mechanism involving induced conformational change of the binding pocket. It may also be helpful for the structure-based design of LOX inhibitors.

METHODS

Crystallographic analysis

The electron density map calculated with the structure factors and coordinates downloaded from PDB displayed not only a contiguous density connecting the “disordered” helix but also a chunk of unaccounted density ~ 10 Å away from it [Fig. 1(b)]. The size and shape of the latter, when properly refined, might have been able to accommodate a few turns of α -helix. It appeared that we could build two kinds of conformation of the helix into these densities. In the space group R32 where the original structure had been solved, a crystallographic twofold axis passed through these densities. If two molecules of the same kind of conformation were to pack around the twofold axis, the helix would bump into each other. Thus, one conformer had to pack with the other conformer to satisfy crystal packing constraints. We inferred that this kind of packing could be achieved by employing two alternative assumptions (see Appendix for details). One (R32 disorder model) was that a protein molecule was disordered in two conformations in R32. The other (R3 twin model) was that the two protein molecules were related by pseudo-twofold symmetry

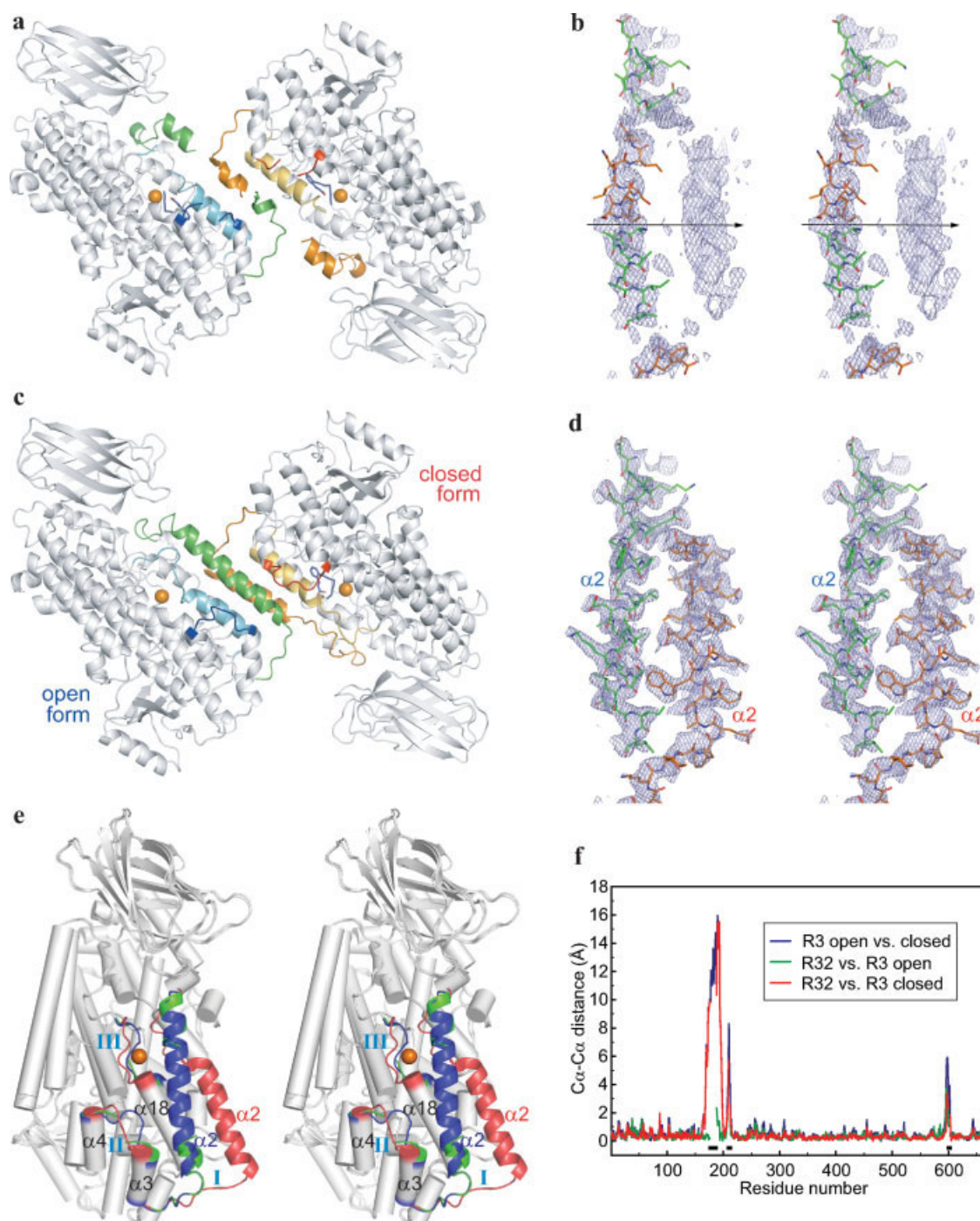
rather than crystallographic one, which implied that the space group was shifted into R3. In this case, the crystal should have been twinned according to a statistical test (see Appendix for details). Both models would be virtually indistinguishable in terms of the final refinement statistics. However, we preferred to use the R3 twin model as it was easier to perform model building and refinement than with the R32 disorder model. Refinement was carried out using the twin option in the CNS package.²⁶ Relieving the crystallographic twofold symmetry allowed us to model and refine two protein molecules that significantly differed in the $\alpha 2$ helix; the helix runs through the “disordered” region in one molecule and through the density ~ 10 Å away from it in the other molecule.

Model building and refinement

The initial R3 dimer model of 15S-LOX1 was made from the PDB coordinates: (*x*, *y*, *z*) for one molecule and ($1/3 + y$, $2/3 + x$, $2/3 - z$) obtained by twin operation for the other. In the R32 structure, the coordinates for 15 residues (177–187, 210–211, and 601–602) were absent.¹² In addition, 19 residues (Lys21, Lys45, Arg66, Lys68, Glu147, Lys157, Asp174, Ser189, Ser191, Leu192, Asn193, Lys199, Ile207, Arg212, Ser327, Gln596, Gln599, Pro600, and Ile603) were modeled as alanines; two thirds of them are near the missing residues.

The model was built using the program O.²⁷ Refinement was carried out using the twin option in CNS, typically with positional refinement, followed by simulated annealing to 3000 K, and finally, individual B-factor refinement.²⁶ Manual rebuilding was carried out between each run with σ_A -weighted, $2F_o - F_c$ maps. For the R_{free} calculation, 10% of the data were randomly selected and set aside. Even in the initial electron density map, residues in the two $\alpha 2$ helices could be easily fitted into the densities. However, further model building and refinement had to be performed meticulously and numerous iterations were necessary to trace several chains, especially, those connected to $\alpha 2$. As the model building proceeded, the symmetric relationship between the two independent molecules started to break down. Eventually, their structures became different assuming an “open” and a “closed” form, respectively; the former is more or less similar to the original R32 structure while the latter is substantially dissimilar.

The electron density for the inhibitor was found only in the closed form. The atomic positions and B-factors of the inhibitor were refined. Water molecules were picked up from the $2F_o - F_c$ and $F_o - F_c$ difference maps on the basis of the peak heights and distance criteria, and were retained only when the B-factor after refinement was below 80 Å². Noncrystallographic symmetry restraints were not applied throughout the structure refinement. The final *R* factor was 0.193 and R_{free} 0.230. The *R* factor could have been lowered but only at

**Figure 1**

Comparison of the original R32 and new R3 structures of 15S-LOX1. (a) Ribbon diagram of the original R32 structure showing a dimer related by crystallographic twofold symmetry with a viewing direction along the twofold axis. The regions near the missing residues are colored. The inhibitor is depicted as a stick model in blue and the catalytic iron as a sphere in orange. (b) Stereoview of the $2F_o - F_c$ map near $\alpha 2$ in R32, contoured at 1.0σ . The disconnected $\alpha 2$ helix and its symmetry-related mate are depicted as stick models with C—C bonds in green and orange, respectively. A crystallographic twofold axis in the diagonal between the a and b axes in R32 is depicted as an arrow in black. (c) Ribbon diagram of the new R3 structure showing two independent molecules related by pseudo-twofold symmetry. The regions that significantly differ in the open and closed forms are colored. The inhibitor is depicted as a stick model in blue and the catalytic iron as a sphere in orange. (d) Stereoview of the $2F_o - F_c$ map near the $\alpha 2$ helices in R3, contoured at 1.0σ . The $\alpha 2$ helices of the open and closed forms are depicted as stick models with C—C bonds in green and orange, respectively. (e) Superposed structures (in stereoview) of the R32 model (green) and the R3 open (blue) and closed (red) forms. Only those regions where the structures are significantly different are colored. Labels I, II, and III represent the segments that show large differences between the R3 open and closed forms. (f) Plot for the C α ...C α distance versus residue number. Short thick bars in black above the abscissa denote the residues missing in the R32 model.

Table I
Crystallographic Statistics

X-ray data	
Space group	<i>R3</i> (<i>R32</i>) ^a
Cell parameters, <i>a</i> , <i>b</i> , <i>c</i> (Å)	198.9, 198.9, 136.1
Resolution (Å)	16–2.4
Reflections, unique	76159 ^b (39132)
Completeness (%)	97.3 (97.9)
Refinement	
Number of residues	1324 (649)
Number of water molecules	219 (94)
Number of ligand	1 (1)
<i>R</i> (%) ^c	19.3 (19.8)
<i>R</i> _{free} (%)	23.0 (24.7)
<i>rms</i> bond length (Å)	0.01 (0.01)
<i>rms</i> bond angles (°)	1.45 (1.53)
Average B-factor (Å ²)	46.8 (48.6)
Main chain atoms	46.7 (47.6)
Side chain atoms	46.9 (49.6)
Ligand atoms	71.4 (73.0)
Water molecules	26.3 (39.4)
Ramachandran plot ^d	
Most favored (%)	87.2 (87.5)
Additionally allowed (%)	12.5 (12.1)
Generously allowed (%)	0.3 (0.4)
Disallowed (%)	0.0 (0.0)

^aThe values for the original *R32* structure are in parentheses.

^bThe *R3* data contains the $|F_o(hkl)|$ reflections downloaded from PDB and twin-related $|F_o(kh-l)|$'s generated by equating $|F_o(kh-l)|$ to $|F_o(hkl)|$ using the program SFTOOLS in the CCP4 package with an option to generate the symmetry related reflections.²⁸

^c $R = \sum |F_{obs}| - |F_{calc}| / \sum |F_{obs}|$. *R*_{free} is the same as *R*, but for a randomly chosen 10% of reflections that were not used in structure refinement.

^dStatistics are for the nonglycine, nonproline, and non-N- or non-C-terminal residues.

the expense of the quality of the model and *R*_{free} value. In the final density map, there was no ambiguous segment in the protein backbone. The side chains of 11 residues (out of 1324), mostly on the protein surface, were disordered or ambiguous because of weak densities, but were modeled considering the interactions with their environments. The residues in several segments of the N-terminal β-barrel and the chains connected to α2 had large B-factors; notably, average B-factors for residues 11–47, 52–56, 66–75, 169–172, 205–214, and 595–602 in the open form and residues 13–34, 40–45, 52–58, 64–75, 84–92, 126–129, 163–177, 189–214, and 595–604 in the closed form were larger than 60 Å². A summary of refinement statistics is shown in Table I. The stereochemical analysis was performed using the program PROCHECK.²⁹ The atomic coordinates have been deposited in the PDB under accession code 2P0M. As we used the twin model, the coordinates transformed by the twin operator are equally valid. Structural figures were created using the program PyMOL (Available at: <http://www.pymol.org>).

RESULTS AND DISCUSSION

In this study, we were able to determine the complete structure of 15S-LOX1 by reinterpreting the reported X-

ray data in the *R3* space group [Fig. 1(c)]. The final model contains the entire 1324 residues for two molecules (residues 2–663), an inhibitor (3-(2-octylphenyl)propanoic acid), and 219 water molecules. The geometry of the final model is good with three residues in the generously allowed regions and no residue in the disallowed regions in the Ramachandran plot.²⁹ The final $2F_o - F_c$ electron density map is well-resolved and relatively clean, as shown in Figure 1(d) for the two α2 helices. The general aspects of LOX structure as well as the structural characteristics of regulatory N-terminal β-barrel are not described here, as they are well discussed or reviewed in the previous articles.^{7,12,13,17}

Comparison of the two 15S-LOX1 structures

The crystal of 15S-LOX1 is interesting in that it has two proteins in the asymmetric unit, one with inhibitor and the other without, existing in two distinct forms. The molecule with inhibitor assumes the closed form in which the access to the binding pocket is blocked by α2. The ligand-free molecule assumes the open form in which the entrance is wide open but the pocket is shallow. These suggest that 15S-LOX1 may undergo ligand-induced conformation changes. The two molecules are related by pseudo-twofold axis that is virtually coincident with the crystallographic twofold axis in *R32*. The dimer interface formed by the two α2 helices mainly consists of hydrophobic interactions without any hydrogen bond [Fig. 1(d)].

The two *R3* structures and the *R32* structure are superposed in Figure 1(e) and the Cα...Cα distances between the equivalent residues are plotted in Figure 1(f). The original *R32* structure with inhibitor is more similar to the open form without inhibitor. It shows *rms* deviations of 0.49 Å for the 625 Cα atoms and 0.47 Å for the 641 Cα atoms with the open and closed forms, respectively. All residues missing in the *R32* model belong to the regions where the two *R3* forms are structurally different (see below). The electron densities for these regions are absent, poorly resolved or broad in the map calculated from the original *R32* model whereas they are clearly defined in the map of the current *R3* model. The *R32* structure was modeled incorrectly at several places including, for instance, Gln596 and Leu597 which are critical in defining the substrate-binding cavity.

For the two *R3* structures, the equivalent Cα atoms of the 619 residues (94% of 662) are related by pseudo-twofold symmetry with an *rms* deviation of 0.59 Å. The remaining Cα atoms deviate from the pseudo-symmetry by distances ranging from 2.4 to 16.0 Å. Major structural differences are found in the three spatially adjacent segments. These include residues 167–199 (labeled as I), 205–215 (II), and 592–604 (III). Segment I of 33 residues, including α2, shows the largest difference with the maximum Cα...Cα distance of 16.0 Å for Asn190. The

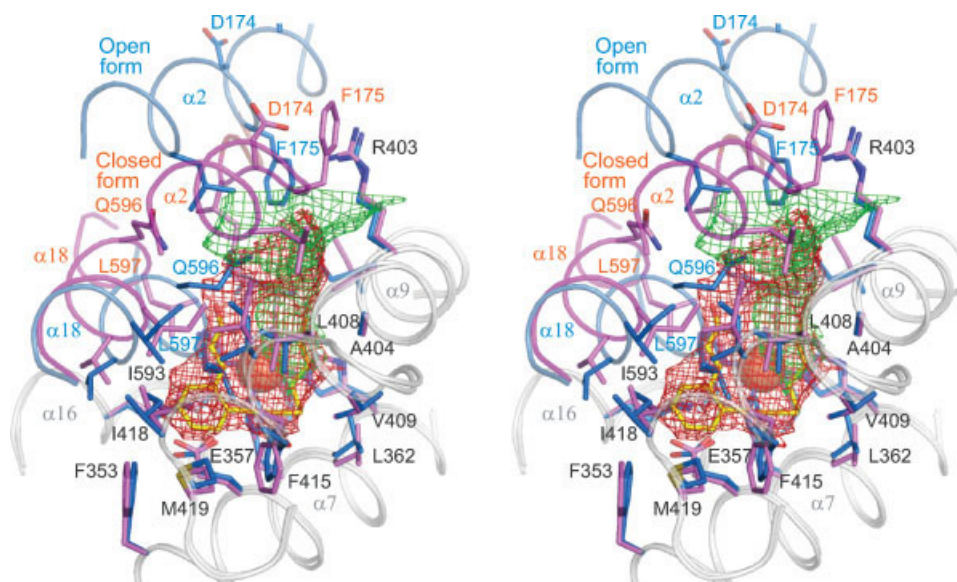


Figure 2

Comparison of the substrate-binding cavities in the open and closed forms of 15S-LOX1 (in stereoview). The secondary structural elements and the selected residues in both forms are superposed together with the cavities depicted as meshes, which are colored in green and red for the open and closed forms, respectively. The iron and the ligating ligand residues are shown behind the meshes. Note that the side chain of Leu597 in the open form partly overlays with the inhibitor (yellow stick) in the closed form.

two $\alpha 2$ helices are separated by ~ 12 Å. They differ in helical length; the open form has seven turns (Leu168-Ile194) while the closed form has five turns (Phe175-Asn193). With regard to this segment, therefore, the conformational change from the open to closed form can be envisaged to occur by unfolding residues Phe167-Asp174 and extending the chain of Asn193-Lys199 and thereby swinging $\alpha 2$ toward the open space.

The conformational differences of segments II and III are correlated with that of segment I. Segment III comprising 13 residues contains the C-terminal end of $\alpha 18$ (Ile593-Gly598), which lies beneath $\alpha 2$ and is critical in shaping the substrate-binding cavity, and the following chain of six residues (Arg599-Ile603). In the closed form, this segment is displaced by ~ 6 Å from its position in the open form because of the movement of Glu169-Phe175 in segment I toward it. This displacement, in turn, induces a change in the loop structure of segment II comprising 11 residues. Otherwise, Trp595 in segment III would collide with Cys210 in segment II.

Comparison of the substrate-binding cavities

Both R3 forms markedly differ in the shape of their substrate-binding cavities, as shown in Figure 2. In this figure, the cavity is represented by the solvent accessible surface built with a probe size of 1.8 Å to depict the space available for the hydrophobic lipid molecule. In the

open form, the cavity is funnel shaped, wide open to the protein surface, and its inside is shallow and narrow. In the closed form, the cavity is deep and wide and is isolated from the protein surface. Leu597 located at the C-terminal turn of $\alpha 18$ is a crucial determinant for controlling the shape and size of the cavity. In the open form, the side chain of Leu597 protrudes into the cavity defining one side of its bottom; it occupies the same space as the propanoic acid moiety of the inhibitor in the closed form. In the closed form, the segment containing Leu597 retreats from the cavity, opening the inner space to the inhibitor. Concurrently, the $\alpha 2$ helix closes the entry to the cavity making it inaccessible from the surface. The C α atoms of Leu597 in both forms are separated by 5.8 Å. It is worthy to note that, in soybean LOX-L3, the corresponding residue Leu773 was pinpointed as crucial for its stereospecificity.¹⁴

The binding cavity of the closed form appears like a bowed thick sheet that is concaved by the side chain of Leu408. Leu408 is on the seven residue-long loop connecting $\alpha 9$ - $\alpha 10$. The structure of this loop is conserved in all known LOXs. The cavity is formed mainly by the side chains of 23 amino acid residues from six helices and a loop. They include Ile173, Phe175, and Ser178 on $\alpha 2$; Phe353, Glu357, His361, Leu362, and His366 on $\alpha 7$; Ile400, Arg403, and Ala404 on $\alpha 9$; Leu408 and Val409 on the $\alpha 9$ - $\alpha 10$ loop; Ile414, Phe415, Ile418, and Met419 on $\alpha 10$; His545 and Gln548 on $\alpha 16$; and Ile593, Val594, Gln596, and Leu597 on $\alpha 18$. In addition, the main chain

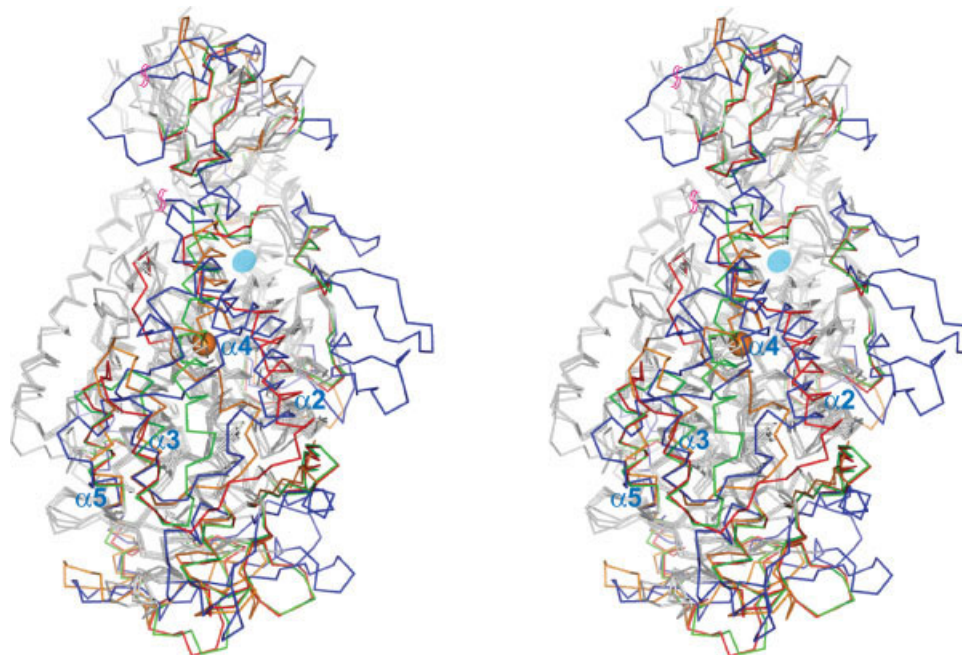


Figure 3

Comparison of the three LOX crystal structures (in stereoview). α -traces of the open and closed forms of 15S-LOX1, coral 8R-LOX, and soybean LOX-L3 are superposed and colored in green, red, orange, and blue, respectively, for the regions that markedly differ from each other. For LOX-L3, 13 residues between two wavy magenta marks in the N-terminal domain are disordered. The labels of the α helices are those for LOX-L3 and the small disk in cyan represents the plausible entry to its cavity.

atoms of Asp174, Leu179, Gly407, and Ile663 are in contact with the cavity. Among the residues in the cavity, His361, His366, His545, and Ile663 are the iron-ligating residues. The inhibitor with a planar structure fits well into the bottom half of the cavity. Its carboxylate moiety points toward the iron; however, the Fe—O distance (5.3 Å) in the present structure is longer than that (3.8 Å) in the R32 structure.

It should be emphasized that, in the R32 structure, several residues at the C-terminal region of α 18 were modeled, accommodating the inhibitor, similar to those in the closed form, while the disconnected α 2 was modeled like α 2 in the open form. If the α 2 were modeled completely, a clash could have occurred between Leu179 at α 2 and Gln596 at α 18; residue Leu179 and the side chain of Gln596 are absent in that model.

Comparison with other LOX structures

The structures of rabbit 15S-LOX1 are superposed with those of soybean LOX-L3¹¹ and coral 8R-LOX¹³ in Figure 3. 15S-LOX1 (662 residues) shares higher sequence identity with 8R-LOX (693 residues, 28.7%) than with LOX-L3 (857 residues, 19.7%), and is more similar to 8R-LOX than LOX-L3 in the overall structure. LOX-L3 has more extended surface loops than the others have. The open and closed forms of 15S-LOX1 overlay with

the structure of 8R-LOX with *rms* deviations of 1.52 Å for the 580 C α atoms and 1.60 Å for the 578 C α atoms, respectively.

By comparing the structure of 8R-LOX and the R32 structure of 15S-LOX1, Oldham *et al.* indicated that the placement of the elements of secondary structure are strikingly conserved at the “lower” half of the catalytic domain but there is deviation in the positions of the shorter helices at its entrance side.¹³ We also noticed that the long segments that form the “upper” half of the catalytic domain and may determine the site of the entry to the cavity show significant structural variations in the three proteins. They include 69 residues (Lys157–Glu225) in 15S-LOX1, 69 residues (Thr161–Glu229) in 8R-LOX, and 105 residues (Val256–Thr360) in LOX-L3. In 15S-LOX1 and 8R-LOX, these segments have the common loop- α 2-loop- α 3-loop- α 4 topology, but their structures are considerably different being observed in three distinct conformational states. For the C α atoms of these segments, 8R-LOX shows *rms* deviations of 4.28 Å and 7.02 Å with the open and closed form of 15S-LOX1, respectively. In 8R-LOX, α 2 was broken into three short helices angled to one another.¹³ These structural variations are indicative of a considerable degree of conformational flexibility present in this region of animal LOXs.

In contrast, soybean LOX-L3 seems to have a compact and rigid structure in this region, although it shows

greater flexibility than the core region indicated by large temperature factors and not well defined electron density.¹¹ LOX-L3 has an extra loop- α 4-loop structure of 36 residues (Tyr311-Pro346) intervening between α 3 and α 5, which overlay α 3 and α 4 in 15S-LOX1, respectively. Thus, unlike α 2 in animal LOXs, any helix in this segment may have no room for large movements. In support of this notion, LOX-L3 showed no major structural change upon ligand binding.^{14–18} Superposition of a free and six liganded LOX-L3 structures shows that the chain of Leu331-Gln341 juxtaposing α 4 made the largest movement but by only 3 Å. α 2 in the open form of 15S-LOX1 is situated beneath α 4 of LOX-L3, whereas that in the closed form overlays α 2 of LOX-L3. In 8R-LOX, a short helix partly overlays α 2 of LOX-L3 and the other is situated beneath α 4. As far as the movement of α 2 is concerned, therefore, it appears that both forms of 15S-LOX1 represent two extreme conformations while 8R-LOX represents an intermediate state.

In LOX-L3, the substrate-binding cavity has an elongated boot shape. The site of entry to the cavity appears to be different from that in 15S-LOX1 because of compact packing around α 2. A putative site could be found in a loop region preceding α 2 as shown in Figure 3. In soybean LOX, changes in cavity size and shape can be done within the cavity by rearranging side chains and the main chain to much lesser degree,²⁰ while much more extensive changes including movements of a long helix and reorientation of several residues in the 'hinges' are observed in 15S-LOX1 (Fig. 2). Ligand-free 8R-LOX has two well-resolved internal cavities.¹³ A constriction was formed by Leu386 and Leu628, which are equivalent to Leu362 (on α 7) and Leu597 (on α 18) of 15S-LOX1, respectively. Neither of the cavities has clear access to the surface. The cavity in 8R-LOX is deeper by \sim 4.7 Å than that in the closed form of 15S-LOX1, the depths being defined mainly by Ser377 and Phe353, respectively. Its width is narrower as Leu432 was shifted by \sim 2 Å toward the catalytic iron compared to the equivalent Leu408 of 15S-LOX1. Oldham *et al.* proposed that, when 8R-LOX binds a substrate, the movement of Trp179 (equivalent to Phe175 on α 2 of 15S-LOX1) may be required to gain access to the cavity and a rotamer change of Leu386 and a shift of Leu628 to open up the constriction.¹³ We now propose that, upon ligand binding, 8R-LOX may also undergo major conformational changes larger than local movements as in the case of 15S-LOX1.

SUMMARY

In this study, we have reinterpreted the crystal structure of 15S-LOX1 based on the data deposited in PDB as 1lox in 1997 by Gillmor *et al.*¹² We have refined the original structure as two independent molecules in the space group R3, assuming that the crystal used for their

study, reported as space group R32, was a twin. In the new model, there is no crystal packing problem and the "disordered" α 2 helix is well resolved in the two distinct conformations. It should be noted that the final molecular structures obtained from the R3 twin model could be comparably refined using the disorder model in R32 (see Appendix). However, the model building in this case would have been impractical because of the unresolved map where the densities of the disordered conformers would overlap each other. The revised structure suggests that large conformational changes may occur in mammalian LOX upon ligand binding. The regions incorrectly determined in the original structure are limited to a small portion of the whole structure, leaving the global view for LOX structure unchanged. However, these regions are functionally important as they include the residues that form the entrance side of the substrate-binding cavity. The R3 structure, especially the closed form, may provide new insight about the arachidonate binding to mammalian LOXs.

The two distinct forms of 15S-LOX1 helped us delineate the following dynamic events that may plausibly occur during the catalytic reaction. Initially, the substrate sticks into the shallow cavity through the wide-open entrance in the open form. As it penetrates into the cavity, segment III containing Leu597 is pushed toward the protein surface, enlarging the cavity. The movement of this inner segment forces nearby outer segments I and II to move outward. Segment I undergoes large conformational changes such that α 2 partially unwinds and moves to close the entrance becoming a closed form. The reverse sequence of movements may occur upon product release. The dynamic nature of α 2 may have another functional importance, because it bears Trp181, one of the residues that are implicated to be important in calcium-dependent membrane binding of 15S-LOX1.³⁰ Trp181 is positioned toward the protein surface in both forms. The flexibility in animal LOXs may be one of the features differentiating them from plant LOXs.

The elucidation of association mode between LOX and arachidonic acid is important in understanding the origin of regio- and stereospecificity of a specific LOX. As the crystal structure for the complex is not available yet, this has been done by modeling studies. There are several complex models for 15S-LOX1 made using the structure of soybean LOX or the R32 model.^{17,30–33} A complex model for 8R-LOX was also reported recently.¹³ In these models, the arachidonate has been depicted as assuming an elongated boot form bound in the methyl end-first orientation. The present study shows that an alternative complex model can be considered since the closed form of 15S-LOX1 represents a new structure for LOX. Its cavity shape suggests that the arachidonate may bind in the horseshoe-like conformation. However, it remains to be determined whether the closed form is truly a productive form. From the structures of 15S-LOX1 and 8R-LOX, we

found that, in addition to the outer $\alpha 2$, the inner $\alpha 7$, and $\alpha 9$ -loop- $\alpha 10$ as well as $\alpha 18$ also show significant deviations, albeit to a much lesser extent than $\alpha 2$, in their positions. These structures indicate that the cavity wall close to the catalytic iron may be somewhat rigid, but the other side of the cavity may be flexible enough to adjust its shape to the ligand structure, as previously noted in the comparative studies of the soybean LOX structures.²⁰ Knowledge on the conformational flexibility of mammalian LOXs may be helpful for further studies on site-directed mutagenesis and structure-based inhibitor design.

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APPENDIX

As described in the Crystallographic analysis section, the two conformers should be packed against each other to avoid bumping. This inevitably removed the crystallographic twofold symmetry, shifting the space group of the unit-cell from R32 to R3. However, Gillmor *et al.* reported rather a small value (0.073) of R_{sym} , the residual of averaging symmetry-equivalent reflection intensities for R32.¹² To resolve this apparent discrepancy we have developed several alternative hypotheses:

1. The unit-cells that had R3 symmetry locally were packed together in a disordered fashion, and thus the crystal as a whole was in R32 space group (R32 disorder model). In other words, the R3 unit-cell and the one related by twofold shared the same lattice dimensions, and thus might have been packed together to form a coherently diffracting crystal. It should be noted that there were two potential modes of disordering. In one case, only the helix might have been disordered and the rest of the molecule obeyed the R32 symmetry (Case 1). In the other case, the whole molecule might have departed from the R32 symmetry by a noticeable amount (Case 2).
2. The R3 unit-cells mentioned above were packed regularly to form a crystal in R3 space group. In this case most part of the molecules obeyed R32 crystallographic twofold symmetry except for the α_2 helices that deviated from the R32 symmetry. As their contribution to the structure factors was much less than that of the rest of the molecule, the resulting R_{sym} might have looked like that of R32. This would be similar to the R32 disorder model Case 1 in terms of various statistics and would require high resolution data to resolve. For the current data, we would not consider this model separately. Instead, we considered a model where the R3 micro-crystals just described might have been packed in such a way that formed a twinned macro-crystal (R3 twin model). In this case, homogeneous R3 micro-crystals might have grown independently but twinned because of the compatibility of their lattices. This twinning might have introduced twofold symmetry in observed diffraction pattern and resulted in such a low R_{sym} .

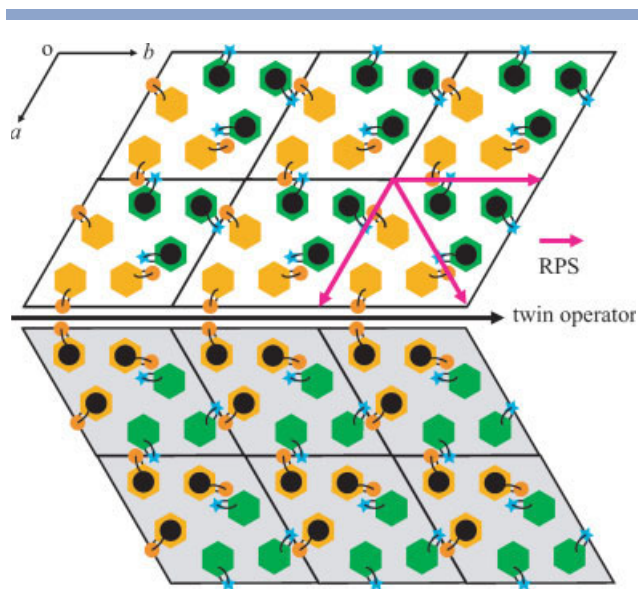


Figure 4

Schematic representation of a perfect hemihedral twin with coinciding RPS in R3 (hexagonal axes). Each protein molecule is represented by a filled hexagon. A 'green' hexagon is related to an 'orange' one by the RPS operator that is virtually parallel to a , b , or $a + b$. The features distinguishing these two molecules (the "disordered" α helices in 15S-LOX1) are signified by the matching small disk and star. The black disk within the hexagon is to denote the orientation along the c -axis: the molecules with and without the disk facing up and down, respectively, and vice versa. The molecules related by rhombohedral centering are omitted for clarity.

In the R3 twin case, the observed intensity would be the sum of contributions from the component crystallites, while in the R32 disorder case the structure factor, not the intensity, would be the vector sum of the individual contributions. As such, the refinement and model building in R3 twin case would be much more straightforward than working with R32 disorder models. It is also well known that R3 crystals are prone to be twinned.^{34,35} Recently, Lebedev *et al.* analyzed the published structures in PDB, and found that R3 crystals have been occasionally mistaken as R32 especially when twinning occurred in association with rotational pseudo-symmetry (RPS) parallel to the twinning operator (refer to Fig. 4).³⁶ In R3 space group, a , b , and $a + b$ can be twin operators. The issue of twinning potentially interfering with RPS could be very complicated; it could not be easily resolved by the usual statistical twinning tests that are typically tested prior to structure determination.^{37,38} Lebedev *et al.* suggested a set of tests that could be useful in this case.³⁶ They defined R_{twin} , the residual difference between intensities of the h and h' reflections related by the twin operator:

$$R_{\text{twin}} = \sum |I_h - I_{h'}| / \sum |I_h + I_{h'}| \quad (1)$$

where I_h and $I_{h'}$ can be based on either F_{obs} or F_{calc} . They showed that the expected value of $R_{\text{twin}}(\text{obs})$

approaches zero for perfect twin crystals, while $R_{\text{twin}}(\text{obs})$ becomes approximately equal to $R_{\text{twin}}(\text{calc})$ for untwinned normal crystals and becomes smaller than $R_{\text{twin}}(\text{calc})$ for partially twinned crystals. When RPS interferes with twinning, a correlation between observations related by twinning operators could be caused either by RPS or by both RPS and twinning. The two cases cannot be discriminated by $R_{\text{twin}}(\text{obs})$ alone, for which cases twinning detection prior to the structure determination is particularly difficult. Lebedev *et al.* showed that the expected value of $R_{\text{twin}}(\text{calc})$ approaches 0.5 when there is no RPS, but becomes smaller than 0.5 when RPS is present.³⁶ Thus, after the structure refinement is completed, the analysis of R_{twin} 's serves as a diagnostic tool to check whether the current hypothesis on the presence of twinning and pseudo-twofold symmetry is reasonable or not.

The structure factors for 15S-LOX1 were available from PDB in a form that had been merged in R32. On the basis of the previously reported R_{sym} , it was possible to estimate $R_{\text{twin}}(\text{obs})$ without the unmerged structure factor values. Whereas R_{twin} is concerned with the difference between $I(hkl)$ and $I(kh-l)$, R_{sym} considers other crystallographic symmetries as well. It can be shown that $R_{\text{twin}}(\text{obs}) \leq R_{\text{sym}}$ and becomes equivalent to R_{sym} if there is no experimental error (e.g., the reflection intensities related by threefold symmetry are exactly the same). Therefore, low R_{sym} guarantees low $R_{\text{twin}}(\text{obs})$. Having established that $R_{\text{twin}}(\text{obs})$ might have been close to zero, we then needed $R_{\text{twin}}(\text{calc})$. If $R_{\text{twin}}(\text{calc})$ also turned out to be very small, it would be likely that the crystals of 15S-LOX1 might have been untwinned in R32 (R32 disorder model Case 1). The two molecules should then be identical except in the "disordered" $\alpha 2$ helical region where each molecule would adopt its own conformation. As they had to be related by the crystallographic twofold axis, we should model them with static disorder. To test that this was indeed the case, we needed $R_{\text{twin}}(\text{calc})$ from the refined model in R3 space group. The

two molecules should be allowed to move independently throughout the refinement. Otherwise, the restriction of R32 is inevitably imposed and the resulting $R_{\text{twin}}(\text{calc})$ would be close to zero.

We carried out the refinement using CNS with the perfect hemihedral twin option in R3. In this option, CNS tries to fit $\sqrt{|F_{\text{calc}}(hkl)|^2 + |F_{\text{calc}}(kh-l)|^2}$ to $F_{\text{obs}}(hkl)$, where $F_{\text{obs}}(hkl)$ should be calculated from the reflection intensity that had been averaged over the twin operator.²⁶ As noted earlier, the F_{obs} values downloaded from PDB had been processed in R32 and were used as they were. Refinement via simulated annealing with interspersed manual model revision proceeded as usual and converged to a stage where the residuals were slightly better than the original values with comparable geometrical *rms*. $R_{\text{twin}}(\text{calc})$ based on the final model was 0.210. This value is undoubtedly larger than $R_{\text{twin}}(\text{obs})$, consistent with the presence of crystal twinning. It is also substantially lower than 0.5, indicating the presence of RPS. In our final model, the pseudo-twofold axis relating the two molecules slightly deviated from the twin axis by 0.07 Å and 0.23 Å, and the *rms* deviations between the equivalent C α atoms except for the $\alpha 2$ helices was 0.59 Å. All these results favored the perfect twin model in R3 (R3 twin model) over that of the static disorder model of the $\alpha 2$ helix in R32 (Case 1). However, it should be noted that the final molecular structures obtained from the R3 twin model could have been comparably refined in the R32 disorder model (Case 2). Although the model building in the latter case would have been impractical because of the unresolved map where the densities of the disordered conformers would overlap each other, the refinement and R-factor calculation could have been done using the final structures of the R3 twin model. The difference between these two models was in the issue whether one had to average the structure factors in complex numbers or just the intensities; it would not be trivial to differentiate them at moderate resolutions as in this case and would require high resolution data.