

#5 如下 model

①

$$\Delta = \Delta_0 + \Delta_1 G(\omega t) \quad \boxed{\text{Real}}$$

$$H(t) = \left(\begin{array}{cc|cc} \epsilon_{k+h} & x+iy & 0 & \Delta(t) \\ x-iy & \epsilon_{k-h} & -\Delta(t) & 0 \\ \hline 0 & \Delta(t) & -(\epsilon_{k+h}) & x(-iy) \\ \Delta(t) & 0 & (x+iy) & -(\epsilon_{k-h}) \end{array} \right) = H_0 + \Delta_1 G \omega t \quad \text{B}$$

• IF $\Delta_1 = 0$ $H(t)$ is a time-independent Equation, thus

$$H(t) = H_0$$

The PH symmetry of this model is

$$\Sigma = \left(\begin{array}{c|c} 0 & I_{2 \times 2} \\ \hline I_{2 \times 2} & 0 \end{array} \right) K = \left(\begin{array}{c|c} 0 & 1 \\ \hline 1 & 0 \end{array} \right) K$$

$$= \sigma_x K$$

check

$$H_0 = \left(\begin{array}{cc} H(k) & \bar{\Delta} \\ \bar{\Delta}^\dagger & -H^*(-k) \end{array} \right)$$

$$\text{where } \bar{\Delta} = \left(\begin{array}{c|c} \Delta & \\ \hline -\Delta & \end{array} \right) \Rightarrow \boxed{\bar{\Delta}^T = -\bar{\Delta}}$$

$\Sigma H(k) \Sigma^{-1}$ well-known result in fact.

$$= \left(\begin{array}{c|c} 0 & 1 \\ \hline 1 & 0 \end{array} \right) \left(\begin{array}{cc} H^*(k) & \bar{\Delta}^* \\ \bar{\Delta}^\dagger & -H(-k) \end{array} \right) \left(\begin{array}{cc} 0 & 1 \\ \hline 1 & 0 \end{array} \right)$$

$$= \left(\begin{array}{cc} 0 & 1 \\ \hline 1 & 0 \end{array} \right) \left(\begin{array}{cc} \bar{\Delta}^* & H^*(k) \\ -H(-k) & \bar{\Delta}^\dagger \end{array} \right)$$

$$= \left(\begin{array}{cc} -H(-k) & \bar{\Delta}^\dagger \\ \hline \bar{\Delta}^* & H^*(k) \end{array} \right) = - \left(\begin{array}{cc} H(-k) & \Delta \\ \hline \Delta^\dagger & -H^*(k) \end{array} \right) = -H(-k)$$

So: $H_0(k) \Rightarrow$ ^{PH operator}
is Σ

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~~Note~~

We expect the whole model after mapping from

time-dependent $H(t)_{4N} \rightarrow$ time-independent $H_{4N \times 4N}$

The PH Operator should be.

$$\oplus \quad \Sigma_{4N \times 4N} = \begin{pmatrix} \Sigma & & & & \\ & \Sigma & & & \\ & & \Sigma & & \\ & & & \ddots & \\ & & & & \Sigma \end{pmatrix}$$

$$= (\oplus) \Sigma = I_{N \times N} \otimes \Sigma$$

Reason: We can write $H_{4N \times 4N}$ in the following way

$H_{4N \times 4N} =$ diagonal terms + off-diagonal coupling terms.

where diagonal terms = eigenvalue of $(i\partial_t - H_0)$

So we write

$$H_{4N \times 4N} = H_{\text{diag}} + H_{\text{off-diag}}$$

$$\Sigma_{4N \times 4N} H_{4N \times 4N}(k) \Sigma_{4N \times 4N}^{-1} = -H_{4N \times 4N}(k)$$

Requires that

$$\begin{aligned} \Sigma_{4N \times 4N} H_{\text{diag}}(k) \Sigma_{4N \times 4N}^{-1} &= -H_{\text{diag}}(-k) \\ \Sigma_{4N \times 4N} H_{\text{off-diag}}(k) \Sigma_{4N \times 4N}^{-1} &= -H_{\text{off-diag}}(-k) \end{aligned}$$

We prove this point in the following by using the following model.

$$H(k, t) = H = H_0 + \Delta_i G_i(\omega_i t) B$$

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where

$$B = \left(\begin{array}{c|c} 0 & 1 \\ \hline 1 & 0 \end{array} \right)$$

Obviously:

~~$$\Sigma H_0 \Sigma^{-1}$$~~

$$\Sigma H(k, t) \Sigma^{-1} = -H(k, t) \text{ for any } t$$

that is

$$\begin{aligned} \Sigma H_0(k) \Sigma^{-1} &= -H_0(-k) \\ \Sigma B \Sigma^{-1} &= -B \end{aligned}$$

this equations ensures that

$H_{4N \times 4N}$ is always PH symmetric with respect to

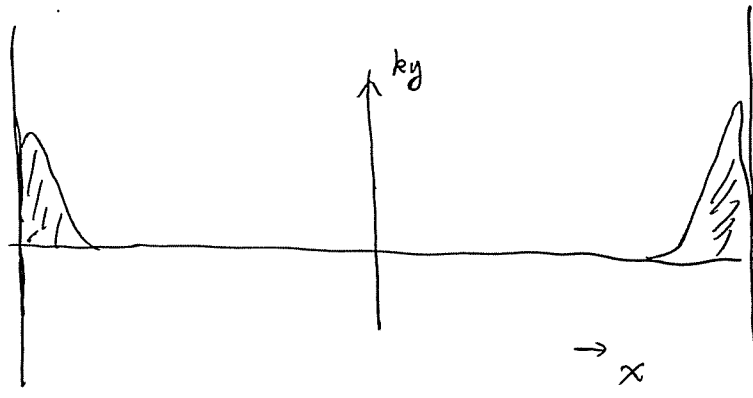
$$\Sigma = I_{N \times N} \otimes \Sigma_{4N \times 4N}$$

Lin, Can you check this PH operator in your numerical calculation?

推广到 $H(t) = H_0 + \sum_i \Delta_i G_i(\omega_i t) B_i$ 也是一样的。

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Hard wall boundary:



At $k_y = 0$, $H(\vec{k}) \rightarrow H(k_x) = H(-i\partial_x)$
is a one-dimensional model

For the state $E=0 \Rightarrow$ Majorana Fermions..

$$\boxed{\Sigma H = -H \Sigma}$$

The MFS at the same edge at the same edge have the same chirality defined as

$$\Sigma \psi = \eta \psi, \text{ where } \eta = \pm 1 \leftarrow$$

Why?

For ~~CH~~ PTH symmetry

$$H \psi = \epsilon \psi \rightarrow H \Sigma^{-1} \psi = -\epsilon \Sigma^{-1} \psi$$

Particle Hole

Now assume $\epsilon=0$ thus ψ & $\Sigma^{-1}\psi$ are degenerate

$$\Rightarrow \Phi_{\pm} = \psi \pm \Sigma^{-1}\psi$$

Then we find that $\boxed{\Sigma \Phi_{\pm} = \pm \Phi_{\pm}}$

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a lot of MFS Here, we need to show that the robustness of those edge modes requires that they have the same chirality.

For any potential V_d with PH symmetry $\Sigma V_d \Sigma^{-1} = -V_d$.

$$\langle \psi_i | V_d | \psi_j \rangle$$

$$= \eta_i \eta_j \langle \psi_i | \underbrace{\Sigma^\dagger V_d \Sigma^{-1}}_{-V_d} | \psi_j \rangle$$

$$= -\eta_i \eta_j \langle \psi_i | V_d | \psi_j \rangle$$

Thus when $\eta_i \eta_j = +1 \Rightarrow \langle \psi_i | V_d | \psi_j \rangle = 0$

The above result show that the edge state ~~should~~ have the same chirality to be robust.

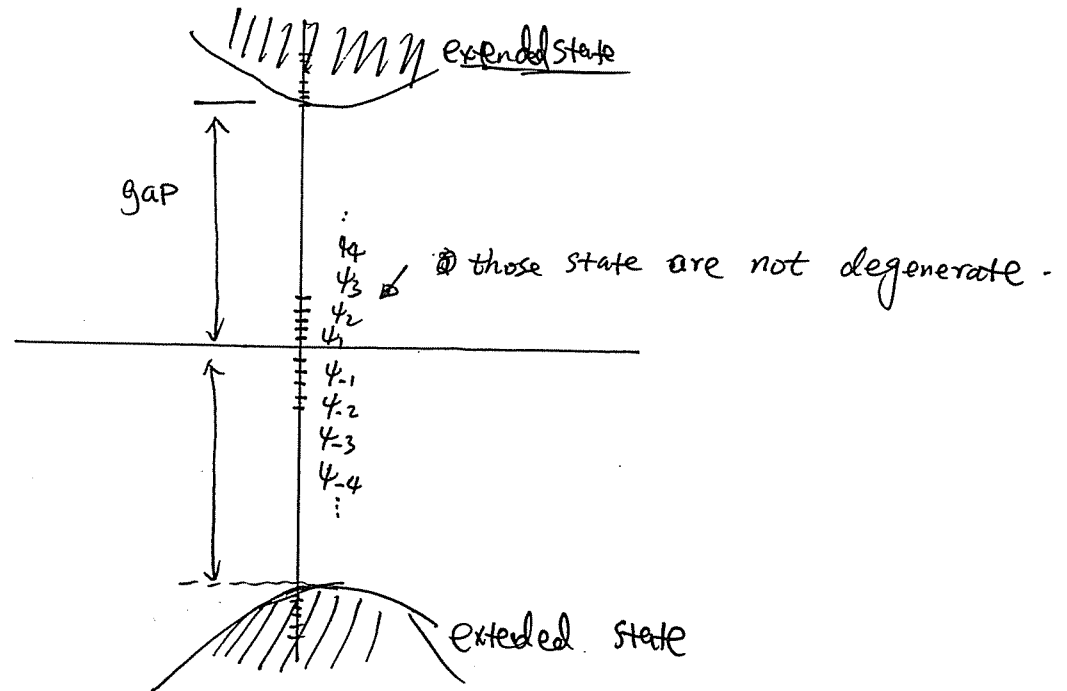
When $\eta_i \eta_j = -1$, the above result always holds for any V_d and $i \neq j \Rightarrow$ we can choose some V_d to make $\langle \psi_i | V_d | \psi_j \rangle \neq 0 \Rightarrow$ can be gapped out.

How to define it in ~~ans~~ numerical calculation?

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A: In a finite chain, L is not required to be very large.
for example $L = 100/k_F$.

Then at $k_y = 0$, we have the following spectrum



Now we have $\psi_{-1} = \Sigma^{-1} \psi_1 = \Sigma \psi_1$

(Note $\Sigma^{-1} = \Sigma$
since $\Sigma^2 = 1$)

Thus the edge state can be
well-defined using

$$\begin{cases} \phi_{\pm 1} = \psi_1 \pm \Sigma \psi_1 \\ \phi_{\pm 2} = \psi_2 \pm \Sigma \psi_2 \\ \phi_{\pm 3} = \psi_3 \pm \Sigma \psi_3 \\ \vdots \end{cases} \Rightarrow \boxed{\Sigma \phi_{\pm} = \pm \phi_{\pm}}$$

Lin, you can check that those wfs $\phi_{\pm 1}, \phi_{\pm 2}, \dots$ are
localized at the ends.

How about $L = 200/k_F$?

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In this case $\epsilon_i \sim 0$ ($|\epsilon_i| < 10^{-7}$)

of course, we can still do that if the energy levels are not degenerate. \therefore

For degenerate energy levels, any superpositions of the wavefunctions are also the eigenvector of the model and Lapack can't automatically ensure that they are ~~the~~ localized edge states.

So, I suggest to study a relative smaller system $L = 100/k_F$ to study the edge states and their chirality.

MFS:

$$\begin{cases} \gamma = \sum_i (u_i c_i + v_i c_i^\dagger) \\ \gamma^\dagger = \sum_i (u_i^* c_i^\dagger + v_i^* c_i) \end{cases}$$

$$\Rightarrow \boxed{v_i = u_i^*, \quad u_i = v_i^*}$$

check that the wavefunctions have the above feature when $|\epsilon| < 10^{-7}$.

Here

$$\boxed{c_i \Rightarrow \phi_{n\sigma} \sim \sin\left(\frac{n\pi x}{L}\right) \sigma}$$

σ is spin.

Pfaffian

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$$\sum H(k) \Sigma^{-1} = -H(-k)$$

Thus $\sum H(0) \Sigma^{-1} = -H(0)$

defines $\Sigma = \Lambda K = \textcircled{H}$

$\Rightarrow \Lambda H^* \Lambda^{-1}$

$$\Sigma H \Sigma^{-1} = \Lambda K H K^{-1} \Lambda^{-1} = \Lambda H^* \Lambda^{-1} = -H$$

$$\Rightarrow \Lambda H^* = -H \Lambda = -W$$

$$W^T = (H \Lambda)^T = \Lambda^T H^T = \Lambda^T H^* = \Lambda H^*$$

because $H^\dagger = H^*$, $\Lambda^T = \Lambda$

$\Rightarrow \boxed{W = -W^T}$

W is a skew matrix.

$$\begin{aligned} \text{Pf}(W) &= \text{Pf}(H \Lambda) = \det(H \Lambda) = \det(H) \cdot \det(\Lambda) \\ &= \det(H) \\ &\text{because } \det(\Lambda) = 1 \end{aligned}$$

$$\Rightarrow \text{Pf}(W) = \pm \sqrt{\det(H)}$$

$\mathcal{N} = -\text{sign}(\text{Pf}(W))$ is topological index, which will not change sign upon deformation when the gap is not closed, that is, ~~det~~ $\det(H) \neq 0$.

method to calculate pfaffian

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First, $\det(H) = \prod_{i=1}^N \lambda_i$ is impossible to be calculated numerically

because when N is very big ($N=200$, say)

$$\det(H) = \begin{cases} \rightarrow 0 \\ \rightarrow \pm 10 \end{cases}$$

Computer can't store this data.

The following have definite pfaffian.

$$V = \begin{pmatrix} 0 & a_1 & & & \\ & 0 & a_2 & & \\ -a_1 & 0 & 0 & & \\ & a_2 & 0 & & \\ & & & \ddots & \\ 0 & & & & 0 \end{pmatrix}$$

then $\text{pf}(w) = a_1 a_2 \dots a_N$

$$\Delta = \text{sign}(\text{pf}(w)) = \prod_i \text{sign}(a_i) \Rightarrow \text{well-defined.}$$

in numerical
calculation.

This can be done using the following Eq.

$$A = B^T B$$

$$\text{pf}(A) = \text{pf}(B^T B) = \text{pf}(B) \cdot \det(B)$$

$$= \text{pf}(B) \quad \text{if } \det(B) = 0$$

phase transition \Rightarrow Pfaffian charge sign

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