## Chapter 2

#### **Mobility**

**2.1** (a) The mean free time between collisions using Equation (2.2.4b) is

$$\mu_n = \frac{q \tau_{mn}}{m_n} \rightarrow \tau_{mn} = \frac{\mu_n m_n}{q} = 2.85 \times 10^{-13} \text{ sec}$$

where  $\mu_n$  is given to be 500 cm<sup>2</sup>/Vsec (= 0.05 m<sup>2</sup>/Vsec), and  $m_n$  is assumed to be  $m_0$ .

(b) We need to find the drift velocity first:

$$v_d = \mu_n \mathcal{E} = 50000 \, cm / \sec$$
.

The distance traveled by drift between collisions is

$$d = v_d \tau_{mn} = 0.14 \, nm$$
.

**2.2** From the thermal velocity example, we know that the approximate thermal velocity of an electron in silicon is

$$v_{th} = \sqrt{\frac{3kT}{m}} = 2.29 \times 10^7 cm / \text{sec}.$$

Consequently, the drift velocity  $(v_d)$  is  $(1/10)v_{th} = 2.29 \times 10^6$  cm/sec, and the time it takes for an electron to traverse a region of 1  $\mu$ m in width is

$$t = \frac{10^{-4} cm}{2.29 \times 10^{6} cm/\text{sec}} = 4.37 \times 10^{-11} \text{sec}.$$

Next, we need to find the mean free time between collisions using Equation (2.2.4b):

$$\mu_n = \frac{q \, \tau_{mn}}{m_n} \rightarrow \tau_{mn} = \frac{\mu_n m_n}{q} = 2.10 \times 10^{-13} \,\text{sec}$$

where  $\mu_n$  is 1400 cm<sup>2</sup>/Vsec (=0.14 m<sup>2</sup>/Vsec, for lightly doped silicon, given in Table 2-1), and  $m_n$  is 0.26m<sub>0</sub> (given in Table 1-3). So, the average number of collision is

$$\frac{t}{\tau_{mn}} = 207.7 \ collision \implies 207 \ collisions$$
.

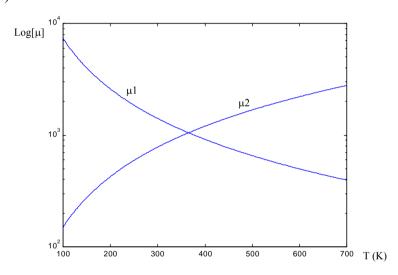
In order to find the voltage applied across the region, we need to calculate the electric field using Equation (2.2.3b):

$$v_d = -\mu_n \mathcal{E} \rightarrow \mathcal{E} = \frac{v_d}{\mu_n} = \frac{2.29 \times 10^6 \, cm/\sec}{1400 \, cm^2 / V \sec} = 1635.71 \, Vcm^{-1}.$$

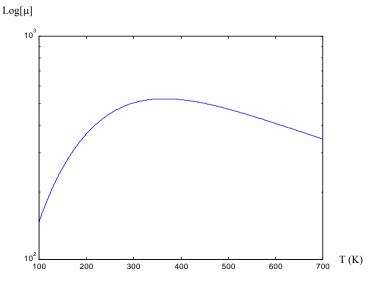
Then, the voltage across the region is

$$V = \mathbf{E} \times width = 1635.71 \, Vcm^{-1} \times 10^{-4} \, cm = 0.16 V$$
.

### **2.3** (a)



# (b) If we combine $\mu_1$ and $\mu_2$ ,



The total mobility at 300 K is

$$\mu_{TOTAL}(300 \, K) = \left(\frac{1}{\mu_1(300 \, K)} + \frac{1}{\mu_2(300 \, K)}\right)^{-1} = 502.55 \, cm^2 \, / V \, \text{sec} \, .$$

(c) The applied electric field is

$$\mathcal{E} = \frac{V}{l} = \frac{1V}{1mm} = 10V/cm.$$

The current density is

$$J_{ndrift} = q\mu_n n \mathcal{E} = q\mu_n N_d \mathcal{E} = 80.41 A / cm^2.$$

### Drift

- **2.4** (a) From Figure 2-8 on page 45, we find the resistivity of the N-type sample doped with  $1\times10^{16}$  cm<sup>-3</sup> of phosphorous is 0.5  $\Omega$ -cm.
  - (b) The acceptor density (boron) exceeds the donor density (P). Hence, the resulting conductivity is P-type, and the net dopant concentration is  $N_{net} = |N_d N_a| = p = 9 \times 10^{16} \text{cm}^{-3}$  of holes. However, the mobilities of electrons and holes depend on the total dopant concentration,  $N_T = 1.1 \times 10^{17} \text{cm}^{-3}$ . So, we have to use Equation (2.2.14) to calculate the resistivity. From Figure 2-5,  $\mu_p(N_T = 1.1 \times 10^{17} \text{cm}^{-3})$  is 250 cm<sup>2</sup>/Vsec. The resistivity is

$$\rho = \frac{1}{\sigma} = \frac{1}{qN_{net}\mu_p} = \frac{1}{q \times 9 \times 10^{16} cm^{-3} \times (250 cm^2 / V \text{ sec})} = 0.28 \ \Omega cm \ .$$

(c) For the sample in part (a),

$$E_c - E_f = kT \ln \left( \frac{N_c}{N_d} \right) = 0.026V \ln \left( \frac{2.8 \times 10^{19} cm^{-3}}{10^{16} cm^{-3}} \right) = 0.21 eV$$

$$\frac{10^{16} cm^{-3}}{10^{16} cm^{-3}} = 0.21 eV$$

For the sample in part (b),

$$E_{f} - E_{v} = kT \ln \left( \frac{N_{v}}{N_{net}} \right) = 0.026V \ln \left( \frac{1.04 \times 10^{19} \, cm^{-3}}{9 \times 10^{16} \, cm^{-3}} \right) = 0.12 \, eV$$

$$E_{c}$$

$$E_{i}$$

$$E_{f}$$

$$E_{f}$$

$$E_{f}$$

$$E_{v}$$

- **2.5** (a) Sample 1: N-type  $\square$  Holes are minority carriers.  $p = n_i^2/N_d = (10^{10} \text{cm}^{-3})^2/10^{17} \text{cm}^{-3} = 10^2 \text{ cm}^{-3}$ 
  - Sample 2: P-type  $\square$  Electrons are minority carriers.  $n = n_i^2/N_a = (10^{10} \text{cm}^{-3})^2/10^{15} \text{cm}^{-3} = 10^5 \text{ cm}^{-3}$
  - Sample 3: N-type  $\Box$  Holes are minority carriers.  $p = n_i^2/N_{net} = (10^{10} cm^{-3})^2/(9.9 \times 10^{17} cm^{-3}) \approx 10^2 cm^{-3}$
  - (b) Sample 1:  $N_d = 10^{17} cm^{-3}$   $\mu_n(N_d = 10^{17} cm^{-3}) = 750 \ cm^2/Vsec$  (from Figure 2-4)  $\sigma = q N_d \mu_n = 12 \ \Omega^{-1} cm^{-1}$ 
    - Sample 2:  $N_a = 10^{15} cm^{-3}$   $\mu_p(N_a = 10^{15} cm^{-3}) = 480 \ cm^2/Vsec$  (from Figure 2-4)  $\sigma = qN_a\mu_p = 12 \ \Omega^{-1}cm^{-1}$
    - Sample 3:  $N_T = N_d + N_a = 1.01 \times 10^{17} \text{cm}^{-3}$   $\mu_n (N_T = 1.01 \times 10^{17} \text{cm}^{-3}) = 750 \text{ cm}^2/\text{Vsec (from Figure 2-4)}$   $N_{net} = N_d - N_a = 0.99 \times 10^{17} \text{cm}^{-3}$  $\sigma = q N_{net} \mu_n = 11.88 \ \Omega^{-1} \text{cm}^{-1}$
  - (c) For Sample 1,

$$E_c - E_f = kT \ln \left( \frac{N_c}{N_d} \right) = 0.026V \ln \left( \frac{2.8 \times 10^{19} cm^{-3}}{10^{17} cm^{-3}} \right) = 0.15 eV$$
.

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For Sample 2,

$$E_{f} - E_{v} = kT \ln \left( \frac{N_{v}}{N_{a}} \right) = 0.026V \ln \left( \frac{1.04 \times 10^{19} cm^{-3}}{10^{15} cm^{-3}} \right) = 0.24 eV.$$

$$E_{c}$$

$$E_{i}$$

$$0.24 eV$$

$$E_{f}$$

$$E_{v}$$

For Sample 3,

$$E_{c} - E_{f} = kT \ln \left( \frac{N_{c}}{N_{net}} = N_{d} - N_{a} \right) = 0.026V \ln \left( \frac{2.8 \times 10^{19} cm^{-3}}{9.9 \times 10^{16} cm^{-3}} \right) = 0.15 eV .$$

$$E_{c}$$

$$E_{f}$$

$$E_{i}$$

$$E_{v}$$

**2.6** (a) From Figure 2-5,  $\mu_n(N_d = 10^{16} \text{cm}^{-3} \text{ of As})$  is 1250 cm<sup>2</sup>/Vs. Using Equation (2.2.14), we find

$$\rho = \frac{1}{\sigma} = \frac{1}{qn\mu_n} = 0.5 \,\Omega cm.$$

(b) The mobility of electrons in the sample depends not on the net dopant concentration but on the total dopant concentration  $N_T$ :

$$N_T = N_d + N_a = 2 \times 10^{16} cm^{-3}$$
.

From Figure 2-5,

$$\mu_n(N_T) = 1140 \, cm^2 / Vs$$
 and  $\mu_p(N_T) = 390 \, cm^2 / Vs$ .

 $N_{net} = N_d - N_a = 0$ . Hence, we can assume that there are only intrinsic carriers in the sample. Using Equation (2.2.14),

$$\rho = \frac{1}{\sigma} = \frac{1}{q n_i \mu_n + q p_i \mu_p} = \frac{1}{q n_i (\mu_n + \mu_p)}$$
$$= \frac{1}{q \times 1 \times 10^{10} cm^{-3} \times (1140 + 390)(cm^2 / V \text{ sec})}.$$

The resistivity is  $4.08 \times 10^5 \Omega$ -cm.

(c) Now, the total dopant concentration (N<sub>T</sub>) is 0. Using the electron and hole mobilities for lightly doped semiconductors (from Table 2.1), we have

$$\mu_n = 1400 \, cm^2 / V \sec \quad and \quad \mu_p = 470 \, cm^2 / V \sec$$
.

Using Equation (2.2.14),

$$\rho = \frac{1}{\sigma} = \frac{1}{q n_i \mu_n + q p_i \mu_p} = \frac{1}{q n_i (\mu_n + \mu_p)}$$
$$= \frac{1}{q \times 1 \times 10^{10} \text{ cm}^{-3} \times (1400 + 470)(\text{cm}^2 / V \text{ sec})}.$$

The resistivity is  $3.34 \times 10^5 \ \Omega$ -cm. The resistivity of the doped sample in part (b) is higher due to ionized impurity scattering.

- **2.7** It is given that the sample is *n*-type, and the applied electric field  $\varepsilon$  is 1000V/cm. The hole velocity  $v_{dp}$  is  $2 \times 10^5$  cm/s.
  - (a) From the velocity and the applied electric field, we can calculate the mobility of holes:

$$v_{dp} = \mu_p \mathcal{E}, \ \mu_p = v_{dp}/\mathcal{E} = 2 \times 10^5 / 1000 = 200 \text{cm}^2 / \text{V} \cdot \text{s}.$$

From Figure 2-5, we find  $N_d$  is equal to  $4.5 \times 10^{17}$ /cm<sup>3</sup>. Hence,

$$n = N_d = 4.5 \times 10^{17} / \text{cm}^3$$
, and  $p = n_i^2 / n = n_i^2 / N_d = 10^{20} / 4.5 \times 10^{17} = 222 / \text{cm}^3$ .

Clearly, the minority carriers are the holes.

(b) The Fermi level with respect to E<sub>c</sub> is

$$E_f = E_c - kT ln(N_d/N_c) = E_c - 0.107 \text{ eV}.$$

(c)  $R = \rho L/A$ . Using Equation (2.2.14), we first calculate the resistivity of the sample:

$$\sigma$$
 = q(μ<sub>n</sub> n + μ<sub>p</sub> p) ≈ qμ<sub>n</sub> n = 1.6×10<sup>-19</sup> × 400 × 4.5×10<sup>17</sup> = 28.8/Ω-cm, and ρ=  $\sigma$  <sup>-1</sup> = 0.035 Ω-cm.

Therefore,  $R = (0.035) \times 20 \mu \text{m} / (10 \mu \text{m} \times 1.5 \mu \text{m}) = 467 \Omega$ .

# Diffusion

**2.8** (a) Using Equation (2.3.2),

$$J = qn \upsilon = qD(dn/dx).$$

Therefore,

$$\upsilon = D(1/n)(dn/dx) = -D/\lambda$$
. (constant)

(b)  $J = q\mu_n n\mathcal{E} = qn\upsilon$  and  $\upsilon = \mu_n\mathcal{E}$ .

Therefore,  $\varepsilon = -D/\mu_n \lambda = -(kT/q)/\lambda$ .

(c)  $\varepsilon = -1000 \text{V/cm} = -0.026/\lambda$ . Solving for  $\lambda$  yields  $0.25 \mu\text{m}$ .

**2.9** (a) 
$$\mathcal{E} = -\frac{dV}{dx} = \frac{1}{q} \frac{dE_v}{dx} = \frac{1}{q} \frac{\Delta}{L} = \frac{\Delta}{qL}$$
.

(b)  $E_c$  is parallel to  $E_v$ . Hence, we can calculate the electron concentration in terms of  $E_c$ .

$$n(x) = n_0 e^{-(E_c(x) - E_c(0))/kT}$$
 where  $E_c(x) - E_c(0) = (\Delta/L)x$ .

Therefore,  $n(x) = n_0 e^{-x\Delta/LkT}$ .

(c) 
$$J_n q n \mu_n \mathcal{E} + q D_n \frac{dn}{dx} = 0$$
 
$$q n_i e^{-\Delta x / LkT} \mu_n \frac{\Delta}{qL} + q D_n n_i e^{-\Delta x / LkT} \left( -\frac{\Delta}{LkT} \right) = 0$$

Therefore,

$$\frac{\mu_n}{q} = \frac{D_n}{kT} \Rightarrow D_n = \frac{kT}{q} \mu_n.$$