Imperial College London

Computational Physics: Fourier Transforms

Blackett Laboratory

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Today: Outline

- Introduction
- Review of Fourier Transforms (very briefly)
- Discrete Fourier Transforms
- Sampling and Aliasing
- Fast Fourier Transforms
- An Introduction to Pseudocode

Introduction

Analytically:

- As you know: represent "any" function as a sum of others
- Sinusoidal functions which form a complete basis set
- Transform into the "frequency domain"
- The coefficients of the sum carry the information
- In the limit of infinitely "close" coefficients, transform into functions
- Appears directly in numerous places in physics

Introduction

In a computer:

- Sample the original function at regular, discrete points
- Perhaps the result of solving a differential equation
- Or it could be data from an experiment

- Perform a Discrete Fourier Transform (DFT)
- DFTs can also be used to as a computational tool, e.g., to solve differential equations
- Many applications:

- •
- The "Fast Fourier Transform" is a particularly efficient implementation

Recap and notation

• Continuous Fourier Transform of f(t):

$$ilde{f}(\omega) = \int_{-\infty}^{+\infty} e^{i\omega t} f(t) \, \mathrm{d}t = \mathcal{F}(f(t)),$$

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-i\omega t} \tilde{f}(\omega) \, \mathrm{d}\omega = \mathcal{F}^{-1}(\tilde{f}(\omega)),$$

for time t and angular frequency $\omega = 2\pi\nu$ (where ν is frequency)

- Time t and ω are reciprocal 'coordinates' or variables
- $\tilde{f}(\omega)$ is the (angular) frequency spectrum of the function f(t)
- The "~" has nothing to do with the same symbol used elsewhere in the course!
- An expansion over plane waves $\exp(-i\omega t)$ which constitute an orthonormal basis
- $\tilde{f}(\omega)$ are the "coefficients" of each component wave of angular frequency ω

Recap: Derivatives

FTs can be very useful in manipulating differential equations

 Spatial or time derivatives become algebraic operations in Fourier space

$$\frac{\mathsf{d}}{\mathsf{d}x}u(x)=v(x)$$

• Taking the FT of this equation yields

$$-ik\tilde{u}(k) = \tilde{v}(k)$$

• The proof is not complicated:

$$\frac{\mathrm{d}}{\mathrm{d}x} \left(\frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ikx} \tilde{u}(k) \, \mathrm{d}k \right) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ikx} \tilde{v}(k) \, \mathrm{d}k$$

and so on...

This can be generalised to higher derivatives and higher dimensions

Introduction

A Continuous Periodic Function

The Fourier Series

The Complex Fourier Series

The Fourier Transform

A Discrete Periodic Function

Introduction

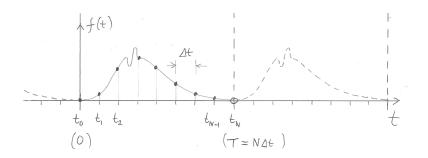
DFTs are the equivalent of complex Fourier series, which are defined on a finite length domain, but for a **discretely sampled function** rather than a continuous function

- Using time and angular frequency here
- Time domain—We assume the function is sampled by N equally—
- A time domain of length $T = N\Delta t$

$$f_n \equiv f(t_n)$$
 with $n = 0, 1, 2, ..., N - 1$ and $t_n = n \Delta t$

- The function f(t) is **assumed to be periodically extended** beyond the domain $0 \le t \le T$ so that f(t + mT) = f(t) for integer m
- For the sampled function, periodicity means $f_N = f_0$; hence we choose to discard f_N in our sequence of samples

Introduction



Introduction

- Frequency domain—discrete spectrum
- Longest wave that can fit exactly into the time domain has an angular frequency ω_{\min} =
- Shortest wave that can be *recognised* by the grid and fits periodically into the time domain has duration so that $\omega_{max} =$
- This maximum frequency is known as the Nyquist frequency

$$\omega_{max} = \frac{\pi}{\Delta t} = \Delta \omega \,\, \frac{N}{2}$$

• These frequencies define the discrete, finite, angular frequency grid on which \tilde{f} is sampled

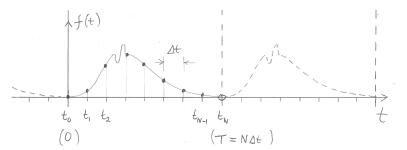
$$\tilde{f}_p pprox rac{1}{\Delta t} \tilde{f}(\omega_p)$$
 with $p = -rac{N}{2}, \dots, 0, \dots, rac{N}{2}$ and $\omega_p = p \Delta \omega = p rac{2\pi}{N \Delta t}$

Introduction

- Each discrete frequency corresponds to a wave that fits exactly an integer number of times "p" into the finite domain
- *N* degrees of freedom ($\tilde{f}_{-N/2} \equiv \tilde{f}_{N/2}$)
- Why " \approx " instead of "=" in " $\tilde{f}_p \approx \frac{1}{\Lambda t} \tilde{f}(\omega_p)$ "?

- "Aliasing" can cause the DFT to be distorted
 - will come back to this
- Will see in a minute where the Δt comes from

Introduction



Definition

The DFT—the analogue of a continuous FT, but with the continuous integral $\int \dots dt$ replaced by a finite sum $\sum \dots \Delta t$:

$$\tilde{f}_p = \sum_{n=0}^{N-1} f_n e^{i\omega_p t_n} = \sum_{n=0}^{N-1} f_n e^{i2\pi p n/N},$$

where $\omega_p t_n = \left(\frac{2\pi p}{N\Delta t}\right) (n\Delta t)$ has been used for the second form

• The backwards transform is similarly defined as:

$$f_n = rac{1}{N} \sum_{p=-N/2+1}^{N/2} \tilde{f}_p e^{-i2\pi p n/N}$$

- The different normalisation factor of 1/N accounts for the discrete sampling (from $dt \to \Delta t$ and $d\omega \to \Delta \omega = 2\pi/(N\Delta t)$)
- When going from the FT to the DFT, a factor of Δt is absorbed into \tilde{f}_p , so $\tilde{f}_p \approx \tilde{f}(\omega_p)/\Delta t$ Limit of continuously-defined points $(N \to \infty)$, and limit of infinite range

• The backwards transform

$$f_n = \frac{1}{N} \sum_{p=-N/2+1}^{N/2} \tilde{f}_p e^{-i2\pi pn/N}$$

is often written as

$$f_n = rac{1}{N} \sum_{p=0}^{N-1} \tilde{f}_p e^{-i 2\pi p n/N}$$

- Aliasing also explains how this is possible
- This is a lossless transform (up to rounding errors etc.)

Code Snippets

```
# Set up basic step sizes
N = 100
deltat = 1.0/N # increment from 0 to 1
# the DFT function of p
def ftilde(p):
   sum = 0.0j
   for n in range(N):
       sum += f(n*deltat)*np.exp(1j*2.0*pi*p*n/N)
   return sum
# the inverse DFT
def finverse(n):
   sum = 0.0j
   for p in range(\Theta, N):
       sum += ftilde(p)*np.exp(-1j*2.0*pi*p*n/N)
   return sum/N
# "Simple" function
def f(x):
   val = -100*x**2+100*x**4 # simple
   if(x>0.4 and x<0.6):
       return val*10*abs(x-0.5) # minor blip in the middle
       #return val
   return val
# plot f(t)
tarray = np.zeros(N)
farray = np.zeros(N)
for i in range(0, N):
    tarrav[i] = i*deltat
   farrav[i] = f(i*deltat)
plt.clf()
plt.plot(tarray.farray.'.-',label="Discrete f(t)",alpha=0.2)
plt.legend()
```

```
# plot f(x) at high resolution
largeN = N *20
finedeltat = deltat / 20.0
tarray2 = np.zeros(largeN)
farray2 = np.zeros(largeN)
for i in range(0, largeN):
    tarray2[i] = i*finedeltat
    farray2[i] = f(i*finedeltat)
plt.plot(tarray2, farray2, '.--', label="Original f(t)",
         alpha=0.2)
plt.legend()
# plot ftilde (both Real and Imaginary parts)
parray = np.arange(0,N)
ftildearrayR = np.zeros(parray.size)
for p in range(0, parray.size):
    ftildearrayR[p] = np.real(ftilde(p))
ftildearrayI = np.zeros(parray.size)
for p in range(0, parray.size):
    ftildearrayI[p] = np.imag(ftilde(p))
plt.clf()
plt.plot(parray, ftildearrayR,'.-',label="ftilde Real",
         alpha=0.2)
plt.legend()
plt.plot(parray, ftildearrayI,',--',label="ftilde Imag",
         alpha=0.6)
plt.legend()
```

Relation to Complex Fourier Series

Complex Fourier series represent the discrete spectrum of a **continuous function** f(t), also on a domain of finite length

$$c_k = rac{1}{T} \int_0^T f(t) \, e^{i \, \omega_k t} \, \mathrm{d}t \,, \quad f(t) = \sum_{k=-\infty}^{+\infty} c_k \, e^{-i \, \omega_k t}$$

- The allowed angular frequencies have the same spacing $\Delta\omega$, but are now infinitely many
- Intimate relationship between Discrete Fourier Transforms and Complex Fourier Series:

$$f_n = \frac{1}{N} \sum_{p=0}^{N-1} \tilde{f}_p e^{-i2\pi pn/N}$$

In Multiple Dimensions

For a function $f_{n,m} = f(x_n, y_m)$ defined on a 2D grid $x_n = m\Delta x$ for $0 \le n \le N - 1$ and $y_m = m\Delta y$ for $0 \le n \le M - 1$:

• The DFT is

$$\tilde{f}_{p,q} = \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{n,m} e^{i \, 2\pi p n/N} e^{i \, 2\pi q m/M}$$

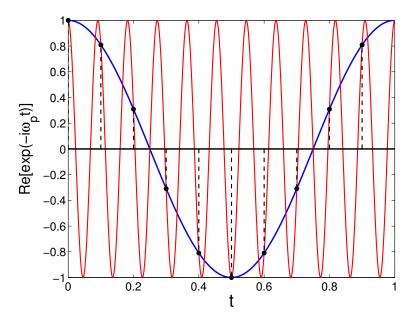
And the backward DFT is

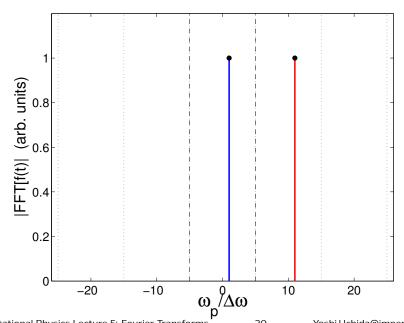
$$f_{n,m} = rac{1}{N\,M} \sum_{p=0}^{N-1} \sum_{q=0}^{M-1} ilde{f}_{p,q} \, e^{-i\,2\pi p n/N} \, e^{-i\,2\pi q m/M}$$

• This can easily be extended to three and more dimensions

- For any wave with a frequency ω_p lying in the frequency domain captured by the DFT, there are higher frequency waves with $\omega_{p'} = \omega_p + m\Omega$ (where $m \in \mathbb{Z}$ and $\Omega = 2\omega_{max}$ is the width of the frequency domain) that look exactly the same when
- This is clear from

$$\exp\left[i(\boldsymbol{\omega_p} + \boldsymbol{m}\Omega)t_n\right] = \exp\left[i\left(\frac{2\boldsymbol{\pi}\,\boldsymbol{p}\boldsymbol{n}}{\boldsymbol{N}} + \right)\right]$$





Computational Physics Lecture 5: Fourier Transforms

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The discrete sampling of the function to be transformed introduces some sampling effects and care has to be taken in allowing for them

- There is a minimum sampling rate that needs to be used so that we don't lose information
- If the function to be sampled fits into the finite length time (or spatial) domain of length *T*, and is **bandwidth limited** to below the Nyquist frequency then
- If the function has $\tilde{f}(\omega) \neq 0$ for $|\omega| > \omega_{max}$ then the frequency content for $|\omega| > \omega_{max}$ will be aliased down into the range $|\omega| < \omega_{max}$ and distort the DFT compared to the
- If the original function is bandwidth limited but is longer than T, then it will be *clipped* which will introduce a sharp cutoff at t = 0 and/or T
- This will effectively increase the bandwidth of the clipped function and aliasing will occur again during sampling of it

- This explains:
 - Why $\tilde{f}_{-N/2} \equiv \tilde{f}$
 - Why the equation for the backwards DFT works:
 - The "negative" frequency part of the spectrum $p = \left\{-\frac{N}{2}+1,\ldots,-1\right\}$ maps to the positive spectrum at $p' = p + N = \left\{\frac{N}{2}+1,\frac{N}{2}+2,\ldots,N-1\right\}$ lying beyond the Nyquist frequency
 - In other words, although the $\sum_{p=N/2+1}^{N-1}$ parts are above the Nyquist frequency, they happen to capture the desired negative frequencies within the Nyquist range
 - Aliasing:

$$ilde{f_p} = \sum_{m} \mathcal{F}(f)|_{\omega_p + m\Omega}$$

i.e., the DFT (the LHS) folds in the Fourier components from the exact FT of the continuous function from the indistinguishable set of waves $\omega_p' = p + m\Omega$.

Fast Fourier Transforms

- ullet Directly coding the DFT results in an $\mathcal{O}(\emph{N}^2)$ algorithm
- But this can be reduced to $\mathcal{O}(N \log_2 N)$ with the Fast Fourier Transform, for $N = 2^m$
- Actually invented in the early 1800s by Gauss (before Fourier), but discovered and made popular by Cooley and Tukey in 1965
- The principles are:
 - Can split an *N*-sample DFT into two *N*/2-sample ones

$$ilde{f_p} = \sum_{n=0}^{N-1} f_n e^{i2\pi pn/N} \ = \sum_{n=0,2,...}^{N-2} f_n e^{i2\pi pn/N} + \sum_{n=1,3,...}^{N-1} f_n e^{i2\pi pn/N} \ = ilde{f_p}^{\text{even}} + e^{i2\pi p/N} imes ilde{f_p}^{\text{odd}}$$

This can be done

Fast Fourier Transforms

Outline of Algorithm

• For $0 \le p < N$:

$$egin{array}{lcl} ilde{f_p} & = & \displaystyle\sum_{n=0}^{N-1} f_n e^{i2\pi p n/N} \ & = & ilde{f_p^{ ext{even}}} + e^{i2\pi p/N} imes ilde{f_p^{ ext{odd}}} \end{array}$$

- Here, $\tilde{f}_p^{\text{even,odd}}$ are N/2-sample DFTs of the even and odd samples of the original f_n
- These are each periodic such that $\tilde{f}_{p+N/2}^{\text{even,odd}} = \tilde{f}_{p}^{\text{even,odd}}$
- This results in, for $0 \le p < N/2$:

$$egin{array}{lcl} ilde{f}_p &=& ilde{f}_p^{ ext{even}} + e^{i2\pi p/ ext{N}} imes ilde{f}_p^{ ext{odd}} \ ilde{f}_{p+ ext{N}/2} &=& ilde{f}_p^{ ext{even}} - e^{i2\pi p/ ext{N}} imes ilde{f}_p^{ ext{odd}} \end{array}$$

• This takes about $\mathcal{O}(N^2/4 + N)$ operations

Pseudocode

- We often need to explain an algorithm to each other, in a way that clearly shows how you would code it up (in any language)
- In this case, it helps not to be burdened by the (language-specific) overhead that real code brings with it (as well as a need to be syntactically perfect)
 - So omit things like

- We call this "Pseudocode"; it is ultimately intended to be human-friendly, but with a code-like structure
 - No formal syntax; it is up to the author to make things clear
 - Can use language-specific elements if it is not confusing (e.g., logic statements from Python or C etc.)

Pseudocode

Examples

 Examples of pseudocode of all sorts of styles are readily available online, but here are a couple of examples:

```
N = 100 \# divisions
                                             def incircle(x, y):
deltat = 1.0/N # goes from 0 to 1
                                               return True if (x,y) is
                                                         in a circle of radius 0.5,
function f(x)
                                                       False if not
  f = -100 \times ^2 + 100 \times ^4
  if x is between 0.4 and 0.5
                                             total = 0
    return f*10*abs(x-0.5)
                                             accepted = 0
  else
                                             loop 10 times
    return f
                                               assign random numbers between
                                                         0 and 1 to x and y
function ftilde(p)
                                               print x, y
                                               if incircle(x, y):
  sum = 0
  for n from 0 to N
                                                 accepted += 1
  sum += f(n*deltat)
                                                 total += 1
         * exp(i*2*pi*p*n/N)
  return sum
                                             ratio = accepted / total
                                             print 4 * ratio, real value of pi
```

Fast Fourier Transforms

Representation in Pseudocode

```
function FFT(f): # f is an array
        N = size of arrav f
        if N == 1:
             return f[0] # for a sample size of 1, the FT is the value itself
        if N is not a power of 2, exit
        # recursive calls
        farray even = FFT(even entries of f) # size N/2
        farray odd = FFT(odd entries of f) # size N/2
        return an N-element array made up of
                   farray_even[p] + exp(i 2pi p/N) * farray_odd[p]
                     (for p = 0..N/2-1)
            and of farray_even[p] - exp(i 2pi p/N) * farray_odd[p]
                     (for p = N/2..N-1)
```

The expressions in the last few lines was corrected after the lecture

Fast Fourier Transforms

- Recursive calls, while logically clear, can require substantial overheads when the code is executed
 - but "Premature optimization is the root of all evil"!
- Further shortcuts exist which can speed up the code by sorting the elements of the calculations in a clever way
- The use of reversing the binary representations of the indices of the discretised function is given in the lecture notes
- The FFT algorithm leads naturally to the Fast Convolution of functions, Filtering, and finding the Correlation between functions etc.

Using Fourier Transforms to Solve Problems

Spectral Methods

• Solving PDEs using FTs, e.g.: the Poisson Equation

where ρ is a source density and u is some potential

- appears in electrostatic and gravitational potentials, etc.
- With periodic boundary conditions, FFT methods are very fast
- Using FTs, the analytical solution for the above is

$$\tilde{u}(\vec{k}) = -\frac{\tilde{
ho}(\vec{k})}{|\vec{k}|^2}$$

• The inverse transform gives:

$$oldsymbol{u}(ec{oldsymbol{x}}) = \mathcal{F}^{-1} \left[-rac{\widetilde{oldsymbol{
ho}}(ec{oldsymbol{k}})}{|ec{oldsymbol{k}}|^2}
ight]$$

 Similar methods can be applied to the calculations of convolutions etc.

Today: Summary

- Introduction
- Review of Fourier Transforms (very briefly)
- Discrete Fourier Transforms
- Sampling and Aliasing
- Fast Fourier Transforms
- An Introduction to Pseudocode
- Spectral Methods

Tomorrow we will look at Random Numbers and Monte Carlo Methods