Instability of Flow In Magnetic Nozzle

Hunt Feng

University of Saskatchewan

June 17, 2023

Outline of Presentation

Outline of Presentation

Instability of Plasma Flow

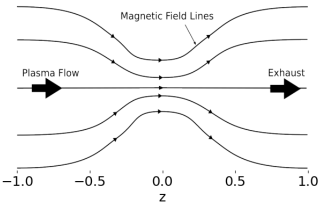
- The instability of plasma flow refers to the tendency of a plasma system to deviate from a stable, equilibrium state and exhibit perturbations or fluctuations in its behavior. [1]

To investigate instability, we assume oscillating perturbed quantities , $\tilde{n}, \tilde{v} \sim \exp(-i\omega t)$.

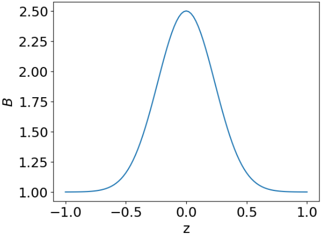
1. If $\mathrm{Im}(\omega) > 0$, then it is unstable flow since the perturbations grow exponential in time, $\exp(\mathrm{Im}(\omega)t)$.
2. If $\mathrm{Im}(\omega) <= 0$, then it is stable flow since the perturbations decay/unchanged in time.

Magnetic Nozzle

- A magnetic nozzle is a device that uses a magnetic field to shape and control the flow of charged particles in a plasma propulsion system.
- Instabilities may affect magnetic nozzle operation and the resulting thrust. [4]



(a) Simplified representation of magnetic nozzle.



(b) A simplified magnetic field of magnetic nozzle.

Figure 1: Simplified representation of magnetic nozzle. Length is normalized.

Governing Equations

The nondimensionalized governing equations for the plasma flow in magnetic nozzle are

$$\text{Cons. of Den.} \qquad \frac{\partial n}{\partial t} + n\frac{\partial v}{\partial z} + v\frac{\partial n}{\partial z} - nv\frac{\partial_z B}{B} = 0 \tag{1}$$

$$\text{Cons. of Mom.} \qquad n\frac{\partial v}{\partial t} + nv\frac{\partial v}{\partial z} = -\frac{\partial n}{\partial z} \tag{2}$$

where n, v are density and velocity, respectively.
The equilibrium quantities n_0, v_0 must satify the condition,

$$\frac{\partial}{\partial z}\left(\frac{n_0 v_0}{B}\right) = 0 \tag{3}$$

$$v_0\frac{\partial v_0}{\partial z} = -\frac{1}{n_0}\frac{\partial n_0}{\partial z} \tag{4}$$

Polynomial Eigenvalue Problem

By linearizing the governing equations, and assume oscillating perturbed quantities, $\tilde{n}, \tilde{v} \sim \exp(-i\omega t)$. We can derive the following equation,

$$
\begin{aligned}
&\omega^2 \tilde{v} \\
&+2i\omega \left(v_0 \frac{\partial}{\partial z} + \frac{\partial v_0}{\partial z} \right) \tilde{v} \\
&+ \left[(1 - v_0^2) \frac{\partial^2}{\partial z^2} - \left(3v_0 + \frac{1}{v_0} \right) \frac{\partial v_0}{\partial z} \frac{\partial}{\partial z} \right. \\
&\quad \left. - \left(1 - \frac{1}{v_0^2} \right) \left(\frac{\partial v_0}{\partial z} \right)^2 - \left(v_0 + \frac{1}{v_0} \right) \frac{\partial^2 v_0}{\partial z^2} \right] \tilde{v} = 0
\end{aligned}
\tag{5}
$$

It is a polynomial eigenvalue problem.

Outline of Presentation

Eq.(5) can be reformulate as

$$\begin{bmatrix} 0 & 1 \\ \hat{M} & \hat{N} \end{bmatrix} \begin{bmatrix} \tilde{v} \\ \omega \tilde{v} \end{bmatrix} = \omega \begin{bmatrix} \tilde{v} \\ \omega \tilde{v} \end{bmatrix} \tag{6}$$

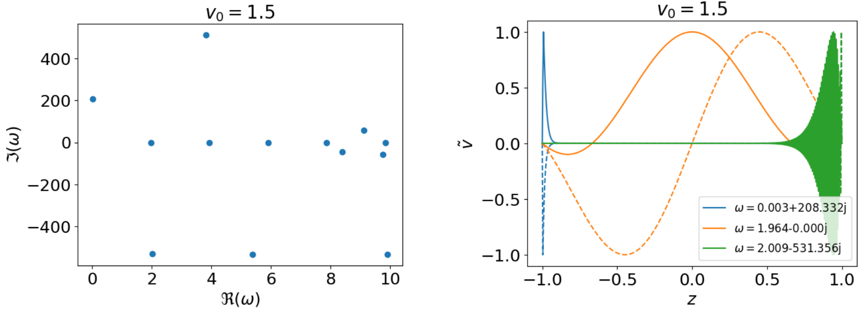
where the operators $\hat{M}$ and $\hat{N}$ are defined as

$$\begin{aligned} \hat{M} = & -\Bigg[(1 - v_0^2) \frac{\partial^2}{\partial z^2} - \left(3v_0 + \frac{1}{v_0} \right) \frac{\partial v_0}{\partial z} \frac{\partial}{\partial z} \\ & - \left(1 - \frac{1}{v_0^2} \right) \left(\frac{\partial v_0}{\partial z} \right)^2 - \left(v_0 + \frac{1}{v_0} \right) \frac{\partial^2 v_0}{\partial z^2} \Bigg] \\ \hat{N} = & -2i \left(v_0 \frac{\partial}{\partial z} + \frac{\partial v_0}{\partial z} \right) \end{aligned}$$

Then by discretizing operators $\hat{M}, \hat{N}$, this becomes an algebraic eigenvalue problem.

Spectral Pollution

- All modes of Eq.(5) with $v_0 = \text{const}$ are stable.
- Yet all discretization show unstable modes.

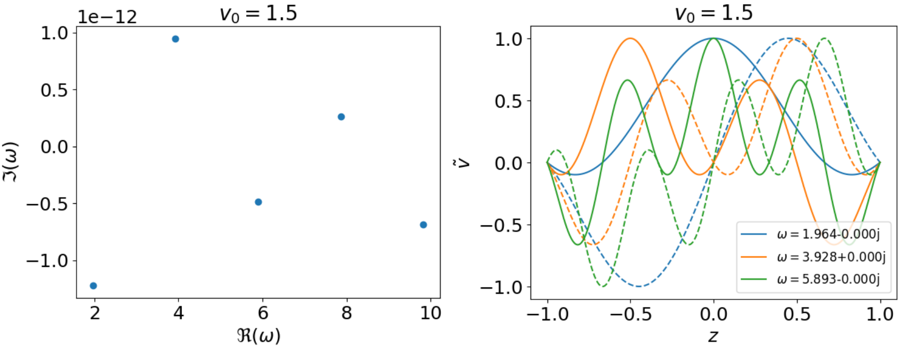


(a) Unfiltered eigenvalues.

(b) First few unfiltered eigenfunctions.

Figure 2: Finite difference discretization was used. Spurious modes occurs regardless of the resolution.

Filtering Spurious Modes



(a) Filtered eigenvalues.

(b) First few filtered eigenfunctions.

Figure 3: The spurious modes are changing under different resolution. We can filter them by convergence test.

Outline of Presentation

Existence of Singularity

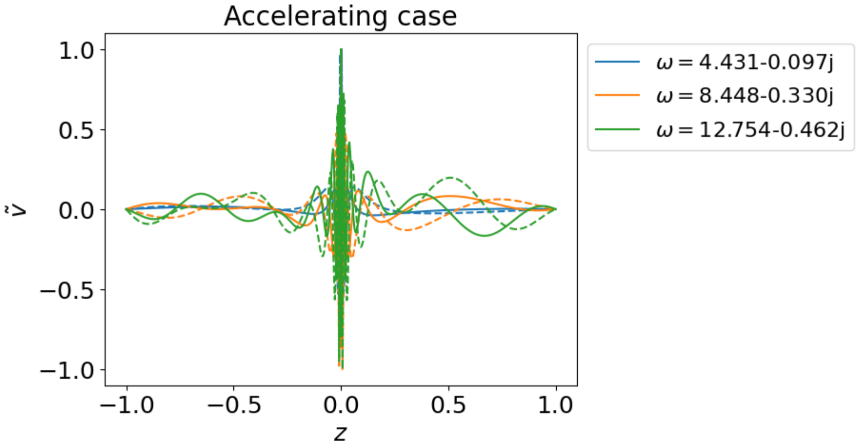


Figure 4: Dirichlet boundary conditions are set at the two ends, all eigenfunctions are squeezed to the singular point.

Interesting Connection to Black Hole

A quasi-1D fluid flow is ruled by the Euler-Lagrange equations and the continuity relation in fluid mechanics, given by [2]

$$\frac{\partial}{\partial t}(\rho A) + \frac{\partial}{\partial x}(\rho A v) = 0 \tag{7}$$

$$\frac{\partial}{\partial t}(\rho A v) + \frac{\partial}{\partial x}[(\rho v^2 + p)A] = 0 \tag{8}$$

$$\frac{\partial}{\partial t}\left(\frac{\rho v^2}{2} - \frac{p}{1-\gamma}A\right) + \frac{\partial}{\partial x}\left[\left(\frac{\rho v^2}{2} - \frac{\gamma}{1-\gamma}A\right)Av\right] = 0 \tag{9}$$

In [2, 3], the tortoise coordinate is introduced to remove singularity,

$$x^* = c_{s0}\int [c_s(x)(1 - M(x)^2)]^{-1} dx$$

where c_{s0} denotes the stagnation speed of sound, $c_s = \mathrm{d}p/\mathrm{d}\rho$ is the local speed of sound and $M(x) = v(x)/c_s(x)$ is the Mach number.

Shooting Method

- Shooting method is employed.
- Initial values are obtained by expanding solution at the singularity.

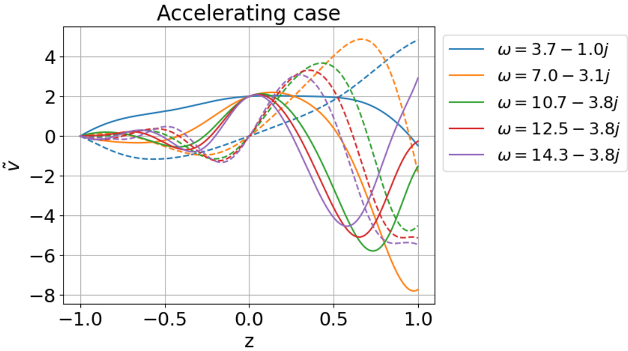


Figure 5: The solutions crosses the singular point smoothly. All modes are stable.

Outline of Presentation

Future Work

- Investigate and interpret the instability of an accelerating flow with non-zero left boundary. See Fig.6
- Compare results to analytically solvable problems with similar configuration.

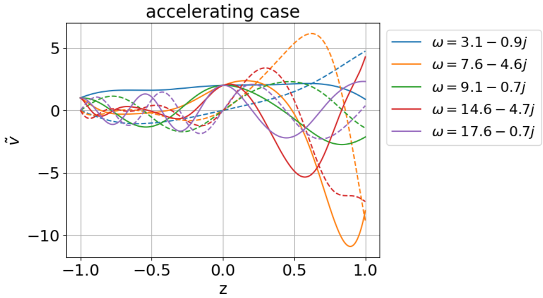


Figure 6: What is the physical interpretation of "non-zero" boundary value? How do we interpret these eigenvalues?

F. F. Chen.
Introduction to Plasma Physics and Controlled Fusion.
Springer International Publishing, Cham, 2016.

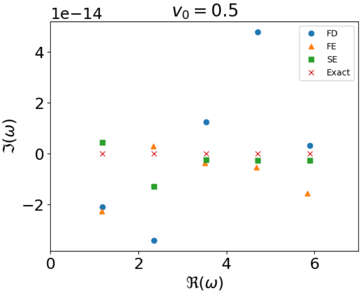
R. da Rocha.
Black hole acoustics in the minimal geometric deformation of a de Laval nozzle.
The European Physical Journal C, 77(5):355, 5 2017.
[Online; accessed 2023-03-02].

H. Furuhashi, Y. Nambu, and H. Saida.
Simulation of an acoustic black hole in a Laval nozzle.
Class. Quantum Grav., 23(17):5417–5438, Sept. 2006.

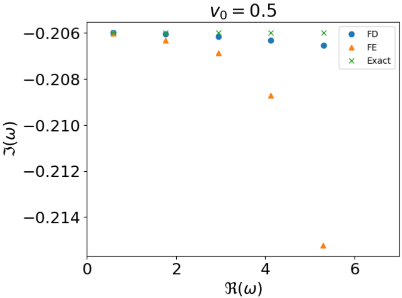
I. D. Kaganovich, A. Smolyakov, Y. Raitses, E. Ahedo, I. G. Mikellides, B. Jorns, F. Taccogna, R. Gueroult, S. Tsikata, A. Bourdon, J.-P. Boeuf, M. Keidar, A. T. Powis, M. Merino, M. Cappelli, K. Hara, J. A. Carlsson, N. J. Fisch, P. Chabert, I. Schweigert, T. Lafleur, K. Matyash, A. V. Khrabrov, R. W. Boswell, and A. Fruchtman.
Physics of E × B discharges relevant to plasma propulsion and similar technologies.
Physics of Plasmas, 27(12):120601, 12 2020.

Outline of Presentation

Constant Velocity Case - Subsonic



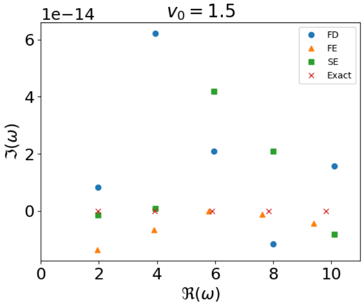
(a) Dirichlet boundary, all modes are stable.



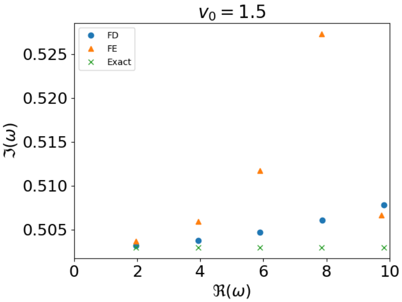
(b) Fixed-open boundary, all modes are stable.

Figure 7: Showing the first 5 eigenvalues. In the Dirichlet boundary case, all methods are close to the exact eigenvalues. Meanwhile, finite-difference method has higher accuracy than finite-element method in fixed-open case.

Constant Velocity Case - Supersonic



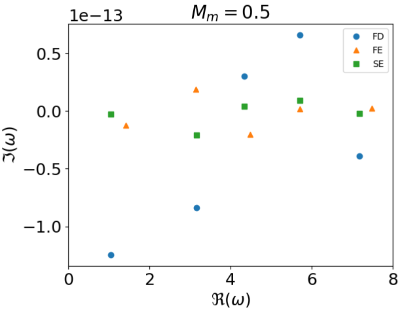
(a) Dirichlet boundary, filtered modes are stable.



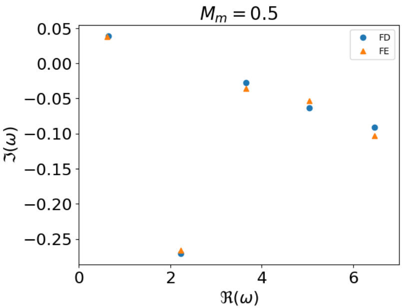
(b) Fixed-open boundary, all modes are unstable.

Figure 8: Showing the first 5 eigenvalues. In the Dirichlet boundary case, all methods are close to the exact eigenvalues. Meanwhile, finite-difference method has higher accuracy than finite-element method in fixed-open case.

Subsonic Case



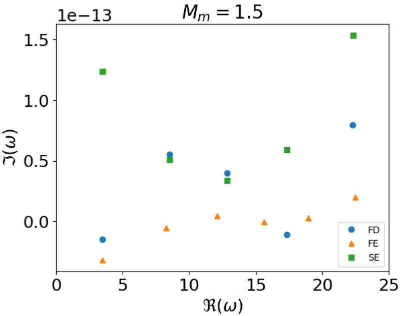
(a) Dirichlet boundary, all modes are stable.



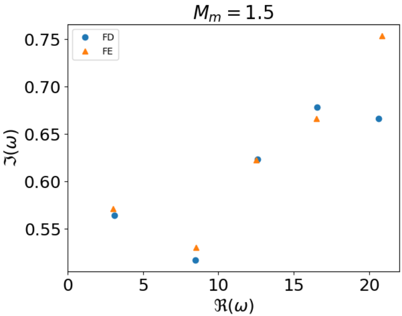
(b) The ground mode is unstable, other modes are stable.

Figure 9: Showing the first 5 modes. It suggests that the subsonic flow in magnetic nozzle is stable.

Supersonic Case



(a) Dirichlet boundary, filtered modes are stable.



(b) Fixed-open boundary, all modes are unstable.

Figure 10: This suggests that the supersonic flow is stable if the boundary is Dirichlet and unstable if the boundary is left-fixed-right-open.

Accelerating Case

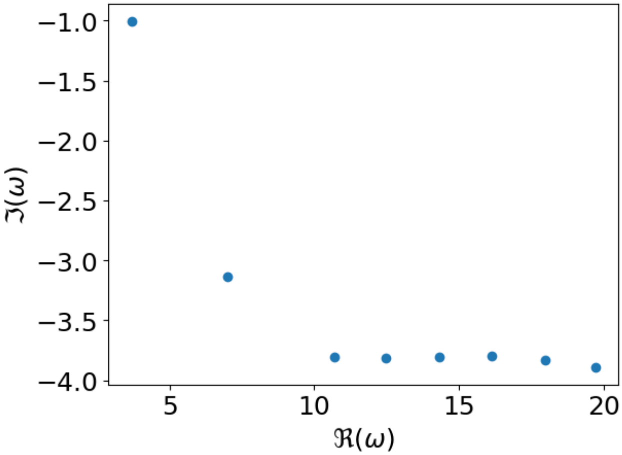


Figure 11: All modes are stable.