

## Coupled Photosynthesis–Stomatal Conductance Model for Leaves of $C_4$ Plants\*

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### Abstract

Leaf based models of net photosynthesis ( $A_n$ ) and stomatal conductance ( $g$ ) are often components of whole plant, canopy and regional models of net primary productivity and surface energy balance. Since  $C_4$  metabolism shows unique responses to environmental conditions and  $C_4$  species are important agriculturally and ecologically, a realistic and accurate leaf model specific to  $C_4$  plants is needed. In this paper we develop a simple model for predicting  $A_n$  and  $g$  from leaves of  $C_4$  plants that is easily parameterised and that predicts many of the important environmental responses.

We derive the leaf model from a simple biochemical–intercellular transport model of  $C_4$  photosynthesis that includes inorganic carbon fixation by PEP carboxylase, light dependent generation of PEP and RuBP, rubisco reaction kinetics, and the diffusion of inorganic carbon and  $O_2$  between the bundle sheath and mesophyll. We argue that under most conditions these processes can be described simply as three potentially limiting steps. The leaf photosynthesis model treats  $A_n$  as first order with respect to either light,  $CO_2$  or the amount of rubisco present and produces a continuous transition between limitations. The independent variables of the leaf photosynthesis model are leaf temperature ( $T_l$ ), intercellular  $CO_2$  levels and the absorbed quantum flux.

A simple linear model of  $g$  in terms of  $A_n$  and leaf surface  $CO_2$  level ( $p_s$ ) and relative humidity ( $h_s$ ) is combined with the photosynthesis model to give leaf photosynthesis as a function of absorbed quantum flux,  $T_l$  and  $p_s$  and  $h_s$  levels.

Gas exchange measurements from corn leaves exposed to varied light,  $CO_2$  and temperature levels are used to parameterise and test the models. Model parameters are determined by fitting the models to a set of 21 measurements. The behaviour of the models is compared with an independent set of 71 measurements, and the predictions are shown to be highly correlated with the data.

Under most conditions the leaf model can be parameterised simply by determining the level of rubisco in the leaves. The effects of light environment, nutritional status and water stress levels on  $A_n$  and  $g$  can be accounted for by appropriate adjustment of the capacity for rubisco to fix  $CO_2$ . We estimate rubisco capacity from  $CO_2$  and light saturated photosynthesis although leaf nitrogen content or rubisco assays from leaf extracts could also be used for this purpose.

### Introduction

Photosynthesis and stomatal movement are physiological processes that occur within leaves but whose influence on  $CO_2$ , water vapour and sensible heat fluxes extends to canopy, regional and global scales. Interest in predicting net primary productivity, hydrology and energy balance of vegetated surfaces in meteorological and climate models has lead to the use of simple leaf models for describing canopy processes (Lindroth and Halldin 1986; Sellers *et al.* 1986; Sellers 1987; Running and Coughlan 1988; Stewart 1988). Recently, more mechanistic models of photosynthesis and stomatal conductance at the leaf level have been proposed (Ball 1988; Leuning 1990; Tenhunen *et al.* 1990; Collatz *et al.* 1991; Sellers *et al.*

1992) that derive from the  $C_3$  photosynthesis model of Farquhar *et al.* (1980) and the empirical stomatal conductance model of Ball *et al.* (1987). These latter models are complete enough to include important responses yet simple enough to be easily parameterised and evaluated numerically.

Though the  $C_3$  pathway of photosynthesis dominates most terrestrial ecosystems, another pathway,  $C_4$ , is important in certain agricultural and natural systems. The  $C_4$  pathway is common among species native to tropical and subtropical grasslands, and important crop species such as corn, sorghum, sugar cane and pasture grasses possess  $C_4$  photosynthesis. Accurate and simple models of photosynthesis and stomatal conductance of  $C_4$  plants would be useful for predicting surface fluxes and energy balance in these systems.

$C_4$  differs from  $C_3$  photosynthesis in several important biochemical and physiological properties. In both types, rubisco fixes  $CO_2$  into the photosynthetic carbon reduction pathway common to all aerobic photosynthetic organisms, but the rubisco reaction is compartmented differently. In  $C_3$  photosynthesis,  $CO_2$  fixed by rubisco is obtained directly from the intercellular spaces of the leaf by diffusion, whereas in  $C_4$  plants  $CO_2$  is delivered to rubisco, which is localised in the bundle sheath chloroplasts, by a metabolic pump that concentrates  $CO_2$ . The elevated  $CO_2$  concentrations maintained in the bundle sheath cells at the cost of additional ATP, has the benefit of inhibiting photorespiration. Consequently,  $C_4$  plants lack several features of  $C_3$  plants that are associated with photorespiration: (a) an  $O_2$  dependent elevated  $CO_2$  compensation point; (b) an inhibition of photosynthesis by  $O_2$ ; and (c) a dependence of the quantum yield for photosynthesis on  $O_2$ ,  $CO_2$  and temperature. The kinetics of the photosynthetic  $CO_2$  response also differ, because for  $C_4$  metabolism the initial  $CO_2$  fixation is via a more efficient catalyst, PEP carboxylase. These differences in function and compartmentation result in differing sensitivities of net photosynthesis to environmental conditions such as temperature,  $CO_2$  and  $O_2$  concentrations, light intensity and nitrogen availability. Generally these differences tend to favour  $C_4$  plants over  $C_3$  plants at high temperatures where photorespiration is stimulated. In addition  $C_4$  plants may be favoured when nitrogen is limiting (Brown 1978). Finally,  $C_4$  and  $C_3$  plants exhibit differences in stomatal conductance under most conditions.  $C_4$  plants in general have lower stomatal conductances than  $C_3$  plants. Lower conductances coupled with higher photosynthetic capacity in leaves of  $C_4$  plants result in higher water use efficiencies in comparison to leaves of  $C_3$  plants (Pearcy and Ehleringer 1984).

These differences in physiological responses together with the ecological and agricultural importance of  $C_4$  plants emphasise the need for a specific model describing  $C_4$  photosynthesis and stomatal conductance. Grant (1989) developed a leaf model for  $C_3$  and  $C_4$  photosynthesis which he coupled with a radiation penetration model to predict canopy photosynthesis. His models are based on kinetic expressions for  $C_3$  photosynthesis described in the models of Laing *et al.* (1974) and Farquhar *et al.* (1980). He modifies his  $C_3$  model to make it behave like  $C_4$  photosynthesis by arbitrarily increasing the chloroplastic  $CO_2$  concentration 200-fold, thus simulating the  $CO_2$  concentrating mechanism. These models, however, give unrealistic responses because of the way that the  $C_3$  model is structured (see Collatz *et al.* 1990). In addition, his model does not accurately describe the kinetics of the  $CO_2$  limited rate of  $C_4$  photosynthesis.

We present here a  $C_4$  photosynthesis-stomatal conductance model that is based on an updated version of the simple intercellular transport (ICT) model of  $C_4$  photosynthesis developed by Berry and Farquhar (1978) and is linked with the stomatal model proposed by Ball *et al.* (1987). It is similar in form to the  $C_3$  model of Collatz *et al.* (1991). The model predicts photosynthesis and stomatal conductance as a function of leaf temperature ( $T_l$ ), photosynthetically active quantum flux density ( $Q_p$ ), and  $CO_2$  partial pressure and relative humidity at the leaf surface ( $p_s$  and  $h_s$ , respectively). The important adjustable parameters are the capacities of rubisco and PEP carboxylase to fix  $CO_2$ , which can be estimated from leaf photosynthetic responses to light and  $CO_2$ . We parameterise our model with a subset of leaf gas exchange data and then show a test of the model's capacity to predict

independent measurements of photosynthetic rate and stomatal conductance under a variety of conditions. The leaf model can be included within the framework of canopy and regional models to predict gross photosynthesis and transpiration from C<sub>4</sub> systems.

## Theory

### Stomatal Model

Stomatal conductance is sensitive to a number of environmental conditions such as light, humidity and CO<sub>2</sub> concentration. Efforts to predict stomatal conductance have focused on empirical analysis of stomatal response to variations in single factors taken one at a time and then the single factor responses are combined to give stomatal responses under conditions where all factors considered vary at the same time. The most commonly used model of this type is that of Jarvis (1976). Ball *et al.* (1987) proposed a model for stomatal conductance that derives primarily from two observations. First, that the ratio of the CO<sub>2</sub> concentration in the intercellular spaces and the leaf surface tends to be constant for leaves of the same photosynthetic pathway (i.e. C<sub>3</sub> or C<sub>4</sub>) provided that atmospheric humidity remains constant (Wong *et al.* 1979; Ball and Berry 1982; Ramos and Hall 1982). Second, the data of Ball and others (Jarvis 1980; Leuning 1990; Mott and Parkhurst 1991) show that the sensitivity of stomata to a given vapour pressure deficit decreases as leaf temperature increases. Ball *et al.* (1987) formulated a simple linear model for stomatal conductance to water vapour consistent with these observations that takes the form

$$g = m \frac{A_n h_s P}{p_s} + b, \quad (1)$$

where  $h_s$  is the leaf surface relative humidity,  $P$  is the atmospheric pressure,  $p_s$  is the leaf surface CO<sub>2</sub> partial pressure, and  $m$  and  $b$  are the slope and intercept of a linear regression (see below). Ball (1988), Collatz *et al.* (1991), Leuning (1990) and Norman and Polley (1989) found that Eqn 1 accounts for much of the variation in stomatal responses observed in C<sub>3</sub> and C<sub>4</sub> leaves with different photosynthetic capacities and exposed to different temperature, light, humidity and CO<sub>2</sub> regimes.

### Photosynthesis Model

Several models of C<sub>4</sub> photosynthesis at the cellular level have been proposed, all of which in essence link C<sub>3</sub> photosynthesis in the bundle sheath chloroplasts with a carbon pump

Table 1. Parameters used in the intercellular transport (ICT) model

Name	Symbol	Equation	Value	Source
Michaelis constant, CO <sub>2</sub>	$K_c$	2A	140 Pa	Furbank <i>et al.</i> (1989)
O <sub>2</sub> inhibition constant	$K_o$	2A	34 kPa	Woodrow and Berry (1988)
CO <sub>2</sub> /O <sub>2</sub> -specificity ratio	$\tau$	2A	2600	Collatz <i>et al.</i> (1991)
RuBP quantum requirement	$\alpha_r$	3A	0.11 mol mol <sup>-1</sup>	Berry and Farquhar (1978)
PEP quantum requirement	$\alpha_p$	5A	0.167 mol mol <sup>-1</sup>	Berry and Farquhar (1978)
Intercellular resistance to CO <sub>2</sub> diffusion	$r_c$	6A	500 m s mol <sup>-1</sup>	Furbank <i>et al.</i> (1989)
Diffusivity of O <sub>2</sub> relative to CO <sub>2</sub>	$D$	7A	40	Farquhar (1983)
Fraction of PSII in bundle sheath	$\lambda$	7A	0.1	Berry and Farquhar (1978)
Fractional RuBP quantum requirement	$f$	3	0.6	Berry and Farquhar (1978)
PEPcase rate constant for CO <sub>2</sub>	$k_p$	4	0.7 mol m <sup>-2</sup> s <sup>-1</sup>	Data, Figs 1 and 2
Maximum rubisco capacity	$V_{\max}$	5	39 μmol m <sup>-2</sup> s <sup>-1</sup>	Data, Figs 1 and 2
Leaf quantum absorptance	$a$	3	0.8	Norman and Polley (1989)
Intercellular partial pressure of O <sub>2</sub>	$O_i$	7A	21 kPa	

driven by the activity of PEP carboxylase in the mesophyll leaf cells (Berry and Farquhar 1978; Peisker 1979; Farquhar 1983; Furbank and Hatch 1987). Carbon derived from intercellular  $\text{CO}_2$  is fixed into  $\text{C}_4$  acids in the mesophyll, transported to the bundle sheath cells and released as  $\text{CO}_2$ . Leakage of inorganic carbon from the bundle sheath cells to the intercellular spaces occurs because there is a large gradient in  $\text{CO}_2$  concentration created by the pump. Following Berry and Farquhar (1978) the steady state balance of these transport processes can be simply expressed as

$$A = W_p - L, \quad (2)$$

where photosynthesis ( $A$ ) is equal to the rate of  $\text{C}_3$  photosynthesis in the bundle sheath chloroplasts.  $W_p$  is the velocity of the PEP carboxylase reaction and  $L$  is the flux of  $\text{CO}_2$  leakage from the bundles sheath to the intercellular spaces of the mesophyll. Both  $A$  and  $L$  are themselves affected by the steady-state  $\text{CO}_2$  concentration in the bundle sheath. We have developed analytical solutions to this intercellular transport (ICT) model that are further described in Appendix A. A simulation of photosynthetic response to intercellular partial pressure of  $\text{CO}_2$  ( $p_i$ ) is shown as the dashed curve in Fig. 2. Parameters used in the simulation are given in Table 1 and discussed in Appendix A.

Although the ICT model gives reasonable simulations, it is complex and unwieldy for routine simulations of leaf responses to environmental conditions. Analysis of the model behaviour shows that under most conditions, the  $\text{CO}_2$  concentrating mechanism appears to work efficiently. Assuming that it is present without formally simulating its operation makes it possible to simplify the ICT model without sacrificing important response characteristics. The resulting model has a structure similar to the  $\text{C}_3$  model of Collatz *et al.* (1991). The logic used to develop the simplified  $\text{C}_4$  model from the more detailed kinetic expressions for the ICT model is discussed below.

At rate limiting light intensities the efficiency of  $\text{CO}_2$  fixation with respect to absorbed light (quantum yield) determines the rate of photosynthesis, and empirical measurements indicate that the quantum yield is constant for  $\text{C}_4$  plants over a wide range of conditions (Ehleringer and Björkman 1977; Ehleringer and Pearcy 1983). In terms of the ICT model the light dependent rate ( $J_i$ ) as given in Eqn 3A (see Appendix A), reduces to

$$J_i = a\alpha_r f Q_p, \quad (3)$$

under conditions where the partial pressure of  $\text{CO}_2$  in the bundle sheath ( $p_{bs}$ ) is sufficiently high to suppress photorespiration. Other terms are:  $a$ , leaf absorptance,  $\alpha_r$ , the intrinsic quantum yield of  $\text{C}_3$  photosynthesis,  $f$ , the fraction of absorbed photons used by the  $\text{C}_3$  reactions, and  $Q_p$ , the incident quantum flux density. In the simplified model the terms  $\alpha_r f$  can be combined into a single constant. The theoretical value for this constant is  $0.067 \text{ mol/mol}$  which approximates the response of the ICT model at  $Q_p \approx 50 \mu\text{mol m}^{-2} \text{ s}^{-1}$ .

At low  $\text{CO}_2$  concentrations, empirical studies show that  $A$  increases linearly from the compensation point (near zero Pa) to rate saturation which occurs at an intercellular  $\text{CO}_2$  partial pressure ( $p_i$ ) of about 10 Pa. The slope is largely independent of  $\text{O}_2$ . In terms of the ICT model the absence of a strong  $\text{O}_2$  dependence arises because the rate is controlled by the initial  $\text{CO}_2$  fixation reaction ( $W_p$ , Eqn 4A) less the leak rate ( $L$ , Eqn 6A) both of which are independent of  $\text{O}_2$  partial pressure. The  $\text{CO}_2$  limited flux ( $J_c$ ) may be expressed as

$$J_c = p_i \left( k_p - \frac{L}{p_i} \right) / P, \quad (4)$$

where  $k_p$  is a pseudo-first-order rate constant for PEP carboxylase with respect to  $p_i$ . The ICT model predicts that  $L/p_i$  is small when  $p_i$  is low and increases non-linearly as  $A$  approaches the limit imposed by the  $V_{\text{max}}$  of rubisco ( $J_c$ , see below).

In terms of the simplified model,  $(k_p - L/p_i)$  is taken as a constant,  $k$ , which may be empirically evaluated from leaf gas exchange measurements at rate limiting  $p_i$  levels. The deviation of  $dA/dp_i$  from  $k$  as  $p_i$  increases is empirically taken into account through the curvature or co-limitation parameter,  $\beta$  (Eqn 3B, Table 2). The models presented here do not explicitly treat the fact that PEP carboxylase uses  $\text{HCO}_3^-$  rather than  $\text{CO}_2$  as substrate or the observations that this enzyme's activity is regulated by activation processes and metabolic pool sizes (Hatch 1987; Leegood and von Caemmerer 1989).

Empirical observations show that, when  $J_i$  and  $J_c$  are not limiting, the rate of assimilation approaches a rate,  $J_e$ , that is largely independent of  $\text{CO}_2$  and light. In the ICT model the rate under these conditions is controlled by the capacity for  $\text{CO}_2$  fixation by rubisco (Eqn 2A). The high  $\text{CO}_2$  concentration in the bundle sheath chloroplasts is close to saturating for rubisco, and the rate under these conditions approaches

$$J_e = V_{\max} \quad (5)$$

Other reactions downstream from rubisco could become limiting under some circumstances (as occurs in C<sub>3</sub> photosynthesis: see Collatz *et al.* 1991) but we currently lack an experimental basis for distinguishing this type of limitation for one imposed by  $V_{\max}$  alone. The ICT model predicts fairly smooth transitions between the  $\text{CO}_2$  limited and  $V_{\max}$  limited states with relatively little co-limitation (Fig. 2). The ICT model described here does not include effects of temperature but these can be added by including the temperature dependences of the kinetic parameters in the model.

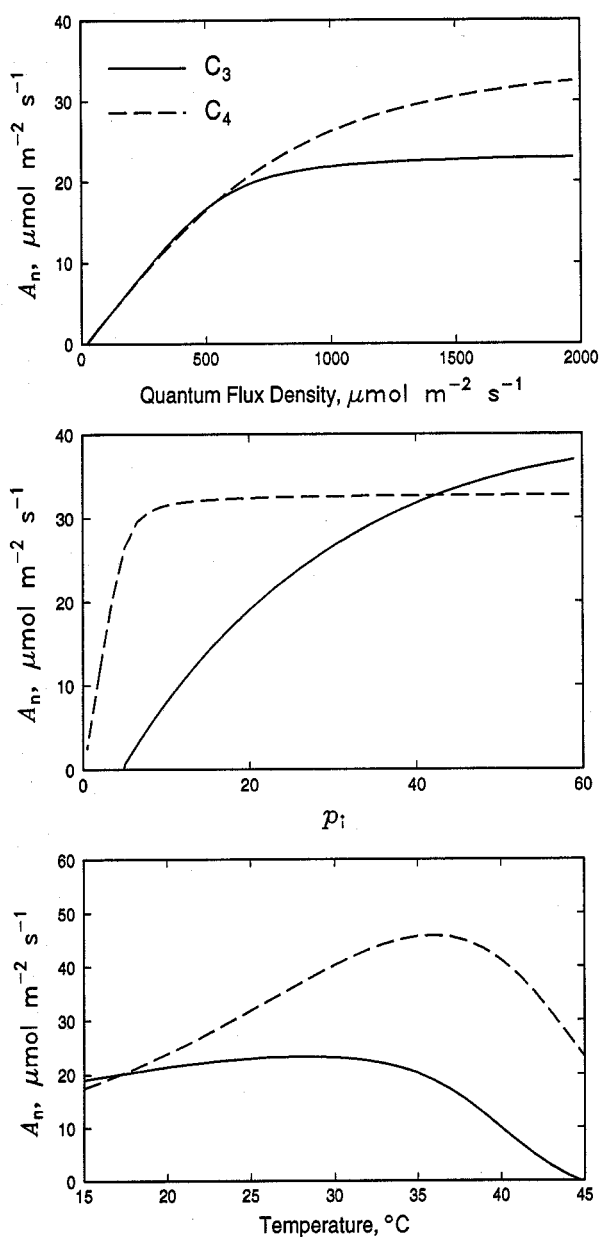
Thus, we can simplify the ICT model to describe just these three limiting states. The relationships between limiting steps and the rate of net photosynthesis ( $A_n$ ) for the simplified model can be expressed as

$$A_n \approx \min \left\{ \begin{array}{l} J_i, f(\alpha, Q_p) \\ J_c, g(p_i, k, T_l) \\ J_e, h(V_{\max}, T_l) \end{array} \right\} - R_d, i(T_l), \quad (6)$$

where  $f$ - $i$  denote separate functions of,  $T_l$  is leaf temperature,  $\alpha$  is the product of leaf absorptance and the intrinsic quantum utilisation efficiency and  $R_d$  is the rate of  $\text{CO}_2$  release or respiration. Equation 6 is the same general form proposed by Farquhar *et al.* (1980) for a model of C<sub>3</sub> photosynthesis and used by others (Kirschbaum and Farquhar 1984; Collatz *et al.* 1991). Following Collatz *et al.* (1991) we used nested quadratic formulations (Eqns 2B and 3B, Appendix B) to pick the subprocess most limiting and to assign a gradual transition from one limitation to another with some specific curvature (see Collatz *et al.* 1990).

Table 2. Parameters used in the simplified C<sub>4</sub> photosynthesis model

Name	Symbol	Equation	Value ( $Q_{10}$ )
Initial slope of photosynthetic $\text{CO}_2$ response	$k$	6	$0.7 \text{ mol m}^{-2} \text{ s}^{-1}$ (2)
Maximum rubisco capacity	$V_{\max}$	6	$39 \text{ } \mu\text{mol m}^{-2} \text{ s}^{-1}$ (2)
Initial slope of photosynthetic light response	$\alpha$	6	$0.04 \text{ mol m}^{-1}$
Leaf respiration	$R_d$	6	$0.8 \text{ } \mu\text{mol m}^{-2} \text{ s}^{-1}$ (2)
Stomatal slope factor	$m$	1	3.0
Stomatal intercept factor	$b$	1	0.08
Atmospheric pressure	$P$	1	$10^5 \text{ Pa}$
Curvature parameter	$\theta$	2B	0.83
Curvature parameter	$\beta$	3B	0.93



**Fig. 1.** Comparison of the simulated responses of  $C_3$  (solid line) and  $C_4$  (dashed line) photosynthesis. Response of net photosynthesis (a) to quantum flux, at 25 $^{\circ}\text{C}$ , and intercellular  $\text{CO}_2$  partial pressure ( $p_i$ ) of 25 and 15 Pa for  $C_3$  and  $C_4$  respectively; (b) to intercellular  $\text{CO}_2$  partial pressure at 25 $^{\circ}\text{C}$  and quantum flux of 1500  $\mu\text{mol m}^{-2} \text{s}^{-1}$  and (c) to leaf temperature at  $p_i$  of 25 and 15 Pa for  $C_3$  and  $C_4$  respectively and quantum flux of 1500  $\mu\text{mol m}^{-2} \text{s}^{-1}$ . Parameters used in the simulations of  $C_3$  photosynthesis are given by Collatz *et al.* (1991) except that the substrate saturated rubisco capacity is set to 90  $\mu\text{mol m}^{-2} \text{s}^{-1}$ .

The simplified model has the advantages of having an analytical solution and a reduced number of parameters (5 v. 13).

The general response characteristics of the simplified C<sub>4</sub> model are shown in Fig. 1 together with those of a C<sub>3</sub> model of Collatz *et al.* (1991) for comparison. The parameters of the C<sub>4</sub> model are given in Table 2 and their estimation is described in detail later. To summarise briefly, Fig. 1a shows the greater photosynthetic capacity and slower approach to light saturation that are characteristic of C<sub>4</sub> photosynthesis (Osmond *et al.* 1980). At 25°C the slope of the light response curve at low quantum fluxes (the quantum yield) is similar for the two types. However, the quantum yield of C<sub>3</sub> photosynthesis is dependent on temperature, O<sub>2</sub> and CO<sub>2</sub> partial pressures (not shown) which is not the case for C<sub>4</sub> plants. This results in higher quantum yields for C<sub>3</sub> plants at lower temperatures while the reverse is true at higher temperatures (Ehleringer and Björkman 1977).

Figure 1b illustrates responses of A<sub>n</sub> to the partial pressure of CO<sub>2</sub> in the intercellular spaces (*p<sub>i</sub>*). The greater initial slope and the rapid saturation with respect to *p<sub>i</sub>* is typical of C<sub>4</sub> photosynthesis (Osmond *et al.* 1980). At *p<sub>i</sub>* levels usually observed in non-stressed C<sub>4</sub> leaves (> 10 Pa) A<sub>n</sub> approaches the V<sub>max</sub> capacity for rubisco. C<sub>3</sub> leaves, on the other hand, may reach similar values of A<sub>n</sub> at *p<sub>i</sub>* levels twice normal, but A<sub>n</sub> remains a fraction of V<sub>max</sub> (90 μmol m<sup>-2</sup> s<sup>-1</sup>) for rubisco. The absence of significant rates of photorespiration in C<sub>4</sub> plants allows A<sub>n</sub> to remain positive at low *p<sub>i</sub>* (< 1 Pa) whereas C<sub>3</sub> photosynthesis is characterised by negative rates below the CO<sub>2</sub> compensation point (c. 4 Pa).

Figure 1c shows the differences between C<sub>3</sub> and C<sub>4</sub> pathways in response to temperature at high quantum flux and constant *p<sub>i</sub>*. The low photorespiration rates and ability to photosynthesise at close to the full rubisco capacity results in higher A<sub>n</sub> at high temperatures in C<sub>4</sub> plants. These differences disappear or are reversed at lower temperatures. The high and low temperature inhibition functions are given in Appendix B. The parameters in these functions can be adjusted to give different temperature response patterns (see below).

Leaves of C<sub>4</sub> plants may show variable photosynthetic responses depending upon the light intensity for growth (Louwerse and Zweerde 1977; Robichaux and Pearcy 1980; Wong *et al.* 1985a), nitrogen nutrition (Wong 1979; Wong *et al.* 1985a; Sage and Pearcy 1987) or drought stress (Ehleringer 1983; Wong *et al.* 1985b). These variations can generally be accommodated by adjusting the parameters V<sub>max</sub>, *k* and *R<sub>d</sub>* of our model as described below in the experimental methods. Accounting for adaptation and acclimation of C<sub>4</sub> photosynthesis to growth temperature regimes (Björkman and Pearcy 1971; Pearcy 1977) may also require adjusting the high and low temperature stability parameters of the model (see Appendix B).

### Coupling Photosynthesis to Stomatal Conductance

*p<sub>i</sub>* appears as an independent variable in the photosynthesis model (Eqns 6 and 3B) but it is a state variable at the leaf level determined by the uptake of CO<sub>2</sub> by photosynthesis and the supply of CO<sub>2</sub> by diffusion through the stomatal pores. Knowing A<sub>n</sub> (Eqn 3B) and *g* (Eqn 1), we can obtain *p<sub>i</sub>* from the one-dimensional model for diffusion of CO<sub>2</sub>,

$$p_i = p_s - A_n / 1.6P/g, \quad (7)$$

where the proportionality factor 1.6 accounts for the ratio of the diffusivities of CO<sub>2</sub> and H<sub>2</sub>O vapour in the stomatal pore (Cowan and Troughton 1971) and *P* is atmospheric pressure.

Equation 7 together with Eqns 1 and 6 can be solved simultaneously to give *g* and A<sub>n</sub> in terms of Q<sub>p</sub>, *T<sub>l</sub>*, *h<sub>s</sub>* and *p<sub>s</sub>*. The exact solution is a cubic equation described in Appendix C. A numerical approach similar to that of Collatz *et al.* (1991) can be used to couple this model with leaf energy balance and boundary layer diffusion.

## Materials and Methods

The fluxes of  $\text{CO}_2$  and  $\text{H}_2\text{O}$  vapour were measured with an open gas exchange system similar to that described and illustrated in fig. 11.4 of Field *et al.* (1989). Portions of corn leaves were enclosed in a water jacketed chamber in which gas fluxes from the top and bottom of leaves were measured separately. Boundary layer conductance was estimated at  $1.2 \mu\text{mol m}^{-2} \text{s}^{-1}$  for each side based on measured temperature, area and evaporation rates from moist filter paper. Variables derived from the flux measurements ( $A_n$ ,  $g$ ,  $p_s$ ,  $h_s$  and  $p_i$ ) were calculated according to von Caemmerer and Farquhar (1981) and Ball (1987) for each leaf surface.  $A_n$  and  $g$  are reported as the sums of measurements from both sides of the leaves.  $h_s$ ,  $p_s$  and  $p_i$  were estimated from the flux-weighted averages of estimates from each side. Corn plants were well fertilised and grown in pots ( $\approx 4 \text{ L}$ ) in a glasshouse at Stanford during the summer. A set of measurements in which quantum flux and the partial pressure of  $\text{CO}_2$  in the intercellular spaces were varied while  $h_s$  and  $T_l$  remained relatively constant ( $h_s = 0.7\text{--}0.8$ ,  $T_l = 25 \pm 1^\circ\text{C}$ ,  $n = 21$ ) were used to set the adjustable parameters of the model.

Equation 6 was solved by selecting the smaller roots of nested quadratic equations (Eqns 2B and 3B). The adjustable parameters  $V_{\text{max}}$ ,  $k$ ,  $\alpha$ ,  $\theta$ , and  $\beta$  were estimated by non-linear regression using a subset of gas exchange measurements.  $R_d$  was estimated by measuring  $\text{CO}_2$  flux from darkened leaves.  $Q_{10}$  values for  $V_{\text{max}}$ ,  $k$  and  $R_d$  were specified as given in Table 2. The parameters  $m$  and  $b$  of the stomatal model (Eqn 1) were determined by linear regression using measured values of  $A_n$ ,  $h_s$  and  $p_s$ .

An independent set of measurements ( $n = 71$ ) over a broad range of  $\text{CO}_2$  partial pressure, quantum flux, temperature and humidity (photon flux density, dark to  $1960 \mu\text{mol m}^{-2} \text{s}^{-1}$ ;  $\text{CO}_2$  partial pressure, 3–34 Pa; leaf temperature,  $18.5\text{--}32.5^\circ\text{C}$ , surface relative humidity, 43–84%) were used to test the accuracy of the parameterised models.

## Results

The predictions of  $A_n$  and  $g$  using fitted parameters were highly correlated ( $r^2 = 0.99$  and  $0.93$  respectively) with the data used to parameterise the models. Fig. 2 shows the

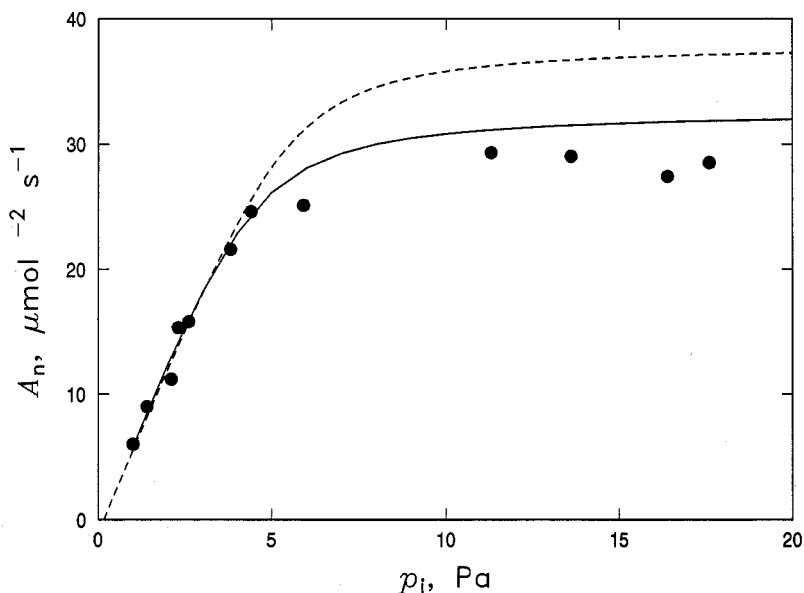
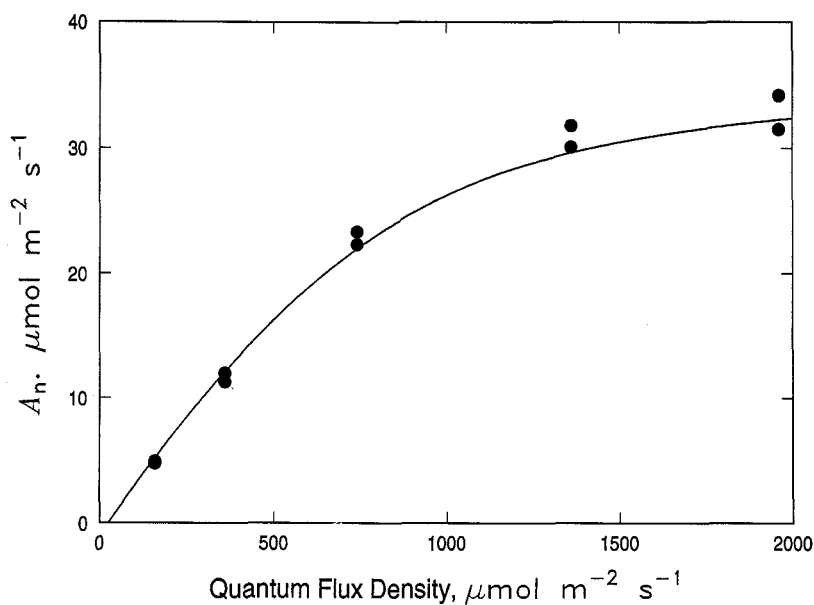


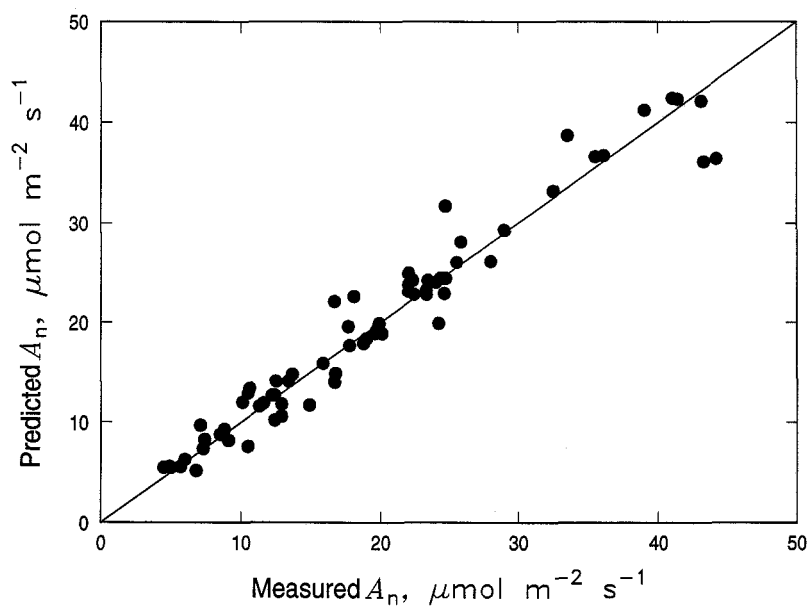
Fig. 2. Measured and predicted responses of photosynthesis ( $A$ ) to the  $\text{CO}_2$  partial pressure in the intercellular spaces of the leaf. Symbols refer to measurements. The solid line is the predicted response based on the leaf photosynthesis model using parameter values giving the best fit to measurements shown here and in Fig. 3. The dashed line shows the response predicted by the intercellular transport model described in Appendix A. Leaf temperature is  $25^\circ\text{C}$  and incident quantum flux is  $1400 \mu\text{mol m}^{-2} \text{s}^{-1}$ .



response of  $A_n$  to  $p_i$  obtained from measurements (symbols) and the model (solid line). We measured CO<sub>2</sub> efflux in the dark at 25°C to be 0.8  $\mu\text{mol m}^{-2} \text{s}^{-1}$  and this value is added to the measured  $A_n$  to give  $A$ . As is generally observed in C<sub>4</sub> plants  $A$  saturates abruptly as  $p_i$  reaches CO<sub>2</sub> partial pressures over 10 Pa. This is reflected in the value of



**Fig. 3.** Response of photosynthesis to incident quantum flux density. Symbols represent measurements and solid line gives response of the model using fitted parameters. Leaf temperature is 25°C and intercellular CO<sub>2</sub> partial pressure is 15 Pa.



**Fig. 4.** Photosynthesis predicted by the simplified C<sub>4</sub> model plotted against measured rates. The line represents the 1 to 1 relationship between predicted and measured rates.

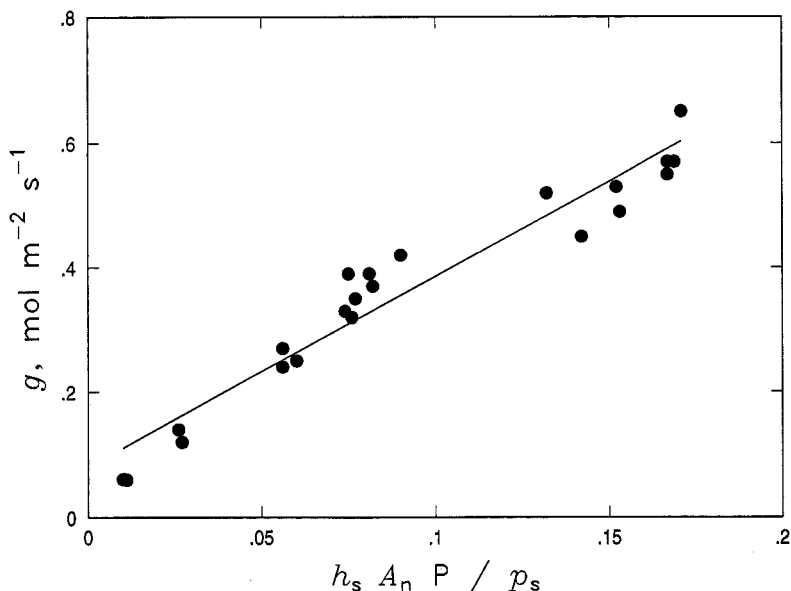
the fitted curvature parameter,  $\beta$ , which is close to 1 ( $\beta = 0.93$ ). On the other hand,  $A$  saturates more gradually with respect to quantum flux as shown in Fig. 3 (curvature parameter  $\theta = 0.83$ ) causing  $A$  to continue increasing beyond full sun quantum fluxes.

Figure 4 shows a comparison of measured  $A_n$  to  $A_n$  predicted by the model using a set of data ( $n = 71$ ) independent of that used to parameterise the photosynthesis model (see Figs 2 and 3; Table 2). Linear regression of the data for a line passing through the origin shows predictions are highly correlated with measurements ( $r^2 = 0.989$ ), and the slope of the regression line is not significantly different from 1 ( $P > 0.999$ ).

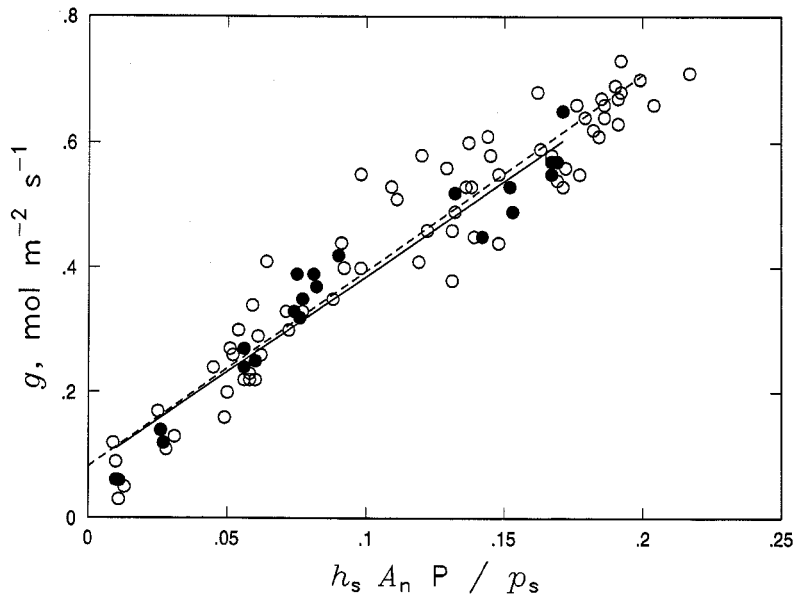
A plot of measured  $g$  versus the product ( $h_s \times A_n \times P/p_s$ ) is shown in Fig. 5 using the same data set used to parameterise the photosynthesis model (Figs 2 and 3). The value for the slope factor,  $m$ , of 3.06 (Table 2) is similar to that reported by Norman and Polley (1989) and Ball (1988) for a number of  $C_4$  species, and is less than half of the values reported for  $C_3$  species (Ball 1988; Leuning 1990). The value for the intercept,  $b$ , of 0.08 is somewhat higher than that found by Norman and Polley (1988) and Ball (1988).

In Fig. 6, the 71 independent measurements of  $g$  are plotted against ( $h_s \times A_n \times P/p_s$ ) as open symbols and the regression line for this data is dashed. For comparison the parameterisation data and regression line from Fig. 4 are re-plotted here as solid symbols and line respectively. The slope factor  $m$  for the validation data set is 3.13 and the intercept  $b$  is 0.08 ( $r^2 = 0.91$ ). The similar parameter values obtained for the parameterisation and validation data sets and high correlation coefficient show that the stomatal model can be accurately parameterised from a relatively small number of measurements. Fitting the data to a second degree polynomial resulted in a small increase in the correlation coefficient ( $r^2 = 0.92$ ).

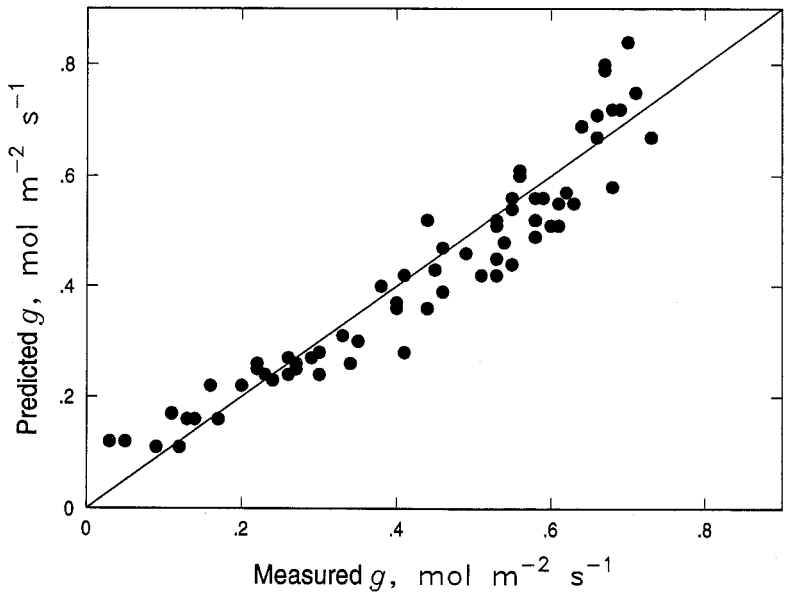
A test of the stomatal conductance and photosynthesis models in combination is provided by using predicted  $A_n$  rather than the measured  $A_n$  in the stomatal model (Eqn 1). A plot of  $g$  predicted from the photosynthesis model and the fitted parameters  $m$  and  $b$  in relation to measured  $g$  is shown in Fig. 7. The regression line is constrained to pass through the



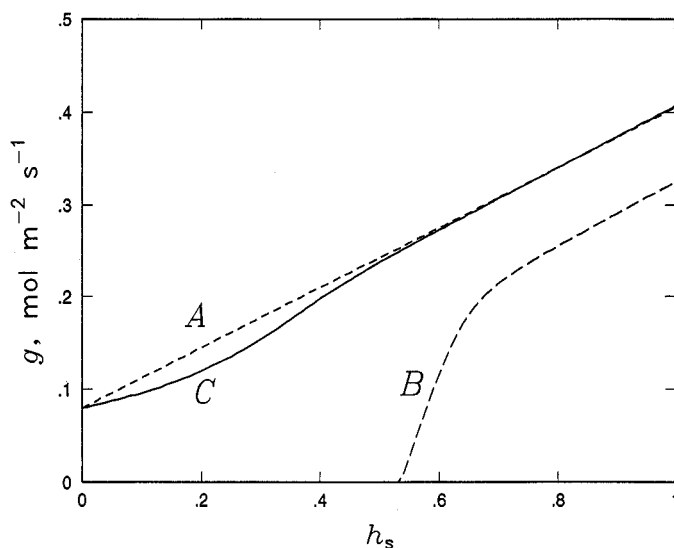
**Fig. 5.** Stomatal conductance as a function of the product of surface humidity, net photosynthesis and the reciprocal of surface  $CO_2$  partial pressure. The regression line is given as  $y = 3.06x + 0.08$  ( $r^2 = 0.93$ ).



**Fig. 6.** Predicted stomatal conductance as a function of net photosynthesis, surface relative humidity and CO<sub>2</sub> partial pressure from an independent data set (open circles). Data and regression line from Fig. 4 (closed symbols and solid line respectively) are included for comparison. The independent data set yields the linear regression equation  $y = 3.13x + 0.08$  ( $r^2 = 0.91$ ).



**Fig. 7.** Stomatal conductance predicted from measured net photosynthesis in comparison to measured stomatal conductance. The line represents the 1 to 1 relationship between predicted and measured rates.



**Fig. 8.** Simulated response of stomatal conductance to leaf surface relative humidity using the coupled photosynthesis–stomatal conductance model. Curve *A* is the response in the absence of photosynthesis feedback on *g*. Curve *B* includes feedback interaction between net photosynthesis and stomatal conductance when the intercept *b* of the stomatal model is set to zero. Curve *C* is the same as *B* except the parameter *b* is set to 0.08 (Table 2). For these simulations quantum flux is  $1500 \mu\text{mol m}^{-2} \text{s}^{-1}$ , leaf temperature is  $25^\circ\text{C}$  and surface  $\text{CO}_2$  partial pressure is 30 Pa.

origin. The correlation between measured and predicted *g* remains high for the linear regression ( $r^2 = 0.984$ ) and the slope of the regression line is not significantly different from 1 ( $P > 0.999$ ).

We did not determine whether our measurements of *g* reflected a spatially homogeneous distribution of conductance across the measured leaf surfaces. It is possible that our calculations of average  $p_i$  levels could be in error due to patchy stomatal closure (Laisk 1983; Daley *et al.* 1989; Farquhar 1989). However, because of the linear responses of  $A_n$  to  $p_i$  at both high and low  $\text{CO}_2$  partial pressures (Figs 1*b* and 2), patchy stomatal closure should be a problem only when photosynthesis of some parts of the leaf is saturated with respect to  $\text{CO}_2$  while in others parts photosynthesis is  $\text{CO}_2$  limited. Under these conditions the calculated  $p_i$  and thus  $A_n$  would be overestimated. The effects of non-homogeneous distribution of *g* are not large enough to cause serious discrepancies between predictions of our model and our leaf measurements.

## Discussion

The ICT model discussed here includes consideration of photorespiration and  $\text{CO}_2$  and  $\text{O}_2$  leakage as component processes of  $\text{C}_4$  photosynthesis. Analysis of the behaviour of the ICT model shows that photorespiration is negligible under most conditions if the  $\text{CO}_2$  concentrating mechanism is given plausible parameters. The simplified model does not explicitly treat photorespiration or the  $\text{CO}_2$  concentrating mechanism but these processes are implicitly contained in the structure of the model. The advantages gained from this simplification are a more tractable model that can be solved analytically and a relatively small number of adjustable parameters. These characteristics make this model useful as a subcomponent of more complex models designed to predict photosynthesis at larger scales such as the leaf and canopy.

Of the five main parameters in the simplified photosynthesis model (Eqn 6),  $V_{\max}$  and  $k$  are likely to be the most variable among leaves grown under different conditions. Additionally,  $R_d$  may vary in a systematic way with growth conditions. Several studies (Wong *et al.* 1979; Robichaux and Pearcy 1980; Wong *et al.* 1985a) show that  $A_n$  measured at high light both under the CO<sub>2</sub> saturation ( $V_{\max}$  limited) and CO<sub>2</sub> limitation ( $k$  limited) varies with the nitrogen nutrition and light levels that C<sub>4</sub> plants are exposed to during growth. Robichaux and Pearcy (1980) and Pearcy (1977) showed that  $R_d$  co-varies with CO<sub>2</sub>-saturated  $A_n$  for differing growth conditions. For unstressed leaves of corn we obtained  $V_{\max} = 39 \mu\text{mol m}^{-2} \text{s}^{-1}$ ,  $k = 0.7 \text{ mol m}^{-2} \text{s}^{-1}$  and  $R_d = 0.8 \mu\text{mol m}^{-2} \text{s}^{-1}$ , or  $k = 18 \times 10^3 V_{\max}$  and  $R_d = 0.021 V_{\max}$ . Assuming  $V_{\max}$ ,  $R_d$  and  $k$  co-vary with leaf nitrogen and remain approximately proportional to one another as conditions change,  $V_{\max}$  (or its surrogate, leaf nitrogen) becomes the most important parameter to estimate in our C<sub>4</sub> model. This is also the case for the C<sub>3</sub> model of Collatz *et al.* (1991).  $V_{\max}$  can be estimated from leaf photosynthesis measurements as was done here or possibly from leaf nitrogen levels.

We expect that the other parameters of the simplified model will vary over smaller ranges of values than  $V_{\max}$  and  $k$ . For example, the kinetic parameters of rubisco, PEP carboxylase and respiration are biochemical properties that should not vary much from leaf to leaf. Similarly, the  $Q_{10}$  values for the temperature responses of  $V_{\max}$ ,  $k$  and  $R_d$  under non-stressed conditions should not vary with the photosynthetic capacity of the leaf. Studies of rubisco (Badger and Collatz 1977),  $R_d$  (Percy 1977; Polley *et al.* 1992) and  $k$  (Polley *et al.* 1992 and our measurements) were used to select the  $Q_{10}$  values used here. Furthermore,  $\alpha$  may also be conservative since leaf absorbance and the intrinsic quantum yield of C<sub>4</sub> photosynthesis vary over narrow ranges and are fairly independent of environmental conditions. Finally, the fitted values for the parameters  $\theta$  and  $\beta$  which reflect the degree of curvature in the light and CO<sub>2</sub> responses of  $A$ , are also likely to be similar for most C<sub>4</sub> plants.

The stomatal model contains two parameters, the slope  $m$  and intercept  $b$  of Eqn 1. The fitted values for these parameters are similar to those reported by others for C<sub>4</sub> plants (Ball 1988; Norman and Polley 1989) and may not need to be routinely adjusted. The smaller value for  $m$  found in C<sub>4</sub> plants relative to C<sub>3</sub> plants is consistent with the observation that  $g$  is usually smaller in the former for a given set of conditions.

The intercept,  $b$ , of the linear equation Eqn 1 gives a minimum value for  $g$  as  $h_s A_n P/p_s \rightarrow 0$ . The existence of a positive intercept in this equation has important implications for the behaviour of the coupled photosynthesis-stomatal conductance model. This is shown graphically in Fig. 8 which gives simulations of  $g$  as a function of  $h_s$ . Setting  $b$  to zero causes  $g$  to go to zero at moderate humidities, because humidity-induced stomatal closure causes  $A_n$  to decrease as  $p_i$  is reduced, which in turn closes stomata further (a positive feedback loop). This behaviour can be illustrated algebraically by combining Eqns 1 and 7 and setting  $b = 0$  to give

$$p_i = p_s(1 - 1.6/mh_s). \quad (8)$$

It is clear from Eqn 8 that as  $h_s$  approaches  $1.6/m$  or  $0.53$ ,  $p_i$  approaches zero, as must  $A$  and  $g$ . If we use the fitted value of  $b$  (curve C in Fig. 8), the modelled response of  $g$  to  $h_s$  is more realistic. The minimum value of  $g$  that we measured in the dark, however, is much smaller than the fitted value of  $b$  at high light levels ( $0.01$  v.  $0.08$ ) implying that minimum conductance is dependent to some extent on light. If  $b$  is held constant, transpiration could be overestimated at night, especially from canopies with high leaf area index. One possible approach to this problem is to make  $b$  dependent upon  $Q_p$ . For instance, a day-night switch or a light-dependent exponential function similar to that used to describe temperature inhibition (Eqn 5B) can be used to obtain a low minimum value of  $g$  at night while avoiding the unrealistic behaviour at low  $h_s$  (Fig. 8) and high  $Q_p$ . It seems possible that other conditions such as severe drought stress may affect the minimum value of  $g$  as well, though we currently have no evidence to support this. Further studies on the response

of  $g$  to stress and low light conditions are needed in order to resolve the behaviour of the minimum  $g$  and to improve the model's treatment of this phenomenon.

The combination of Eqns 1, 6 and 7 takes into account important environmental and physiological factors that affect photosynthesis and stomatal conductance in leaves of  $C_4$  plants with minimum complexity. Such a model should be useful for predicting  $C_4$  canopy photosynthesis provided leaf surface conditions can be specified. Simple canopy models which assume single bulk values for surface concentrations through the whole canopy may be adequate for predicting fluxes of  $CO_2$  and  $H_2O$  (Goudriaan 1989; Sellers *et al.* 1992) from canopies, when used in conjunction with a radiation penetration model and the leaf photosynthesis model reported here.

Simple yet realistic models of photosynthesis and stomatal conductance of canopies of  $C_4$  plants should be useful for the study of the impacts of large-scale land use change in the tropics. Since rainforests which are largely composed of  $C_3$  species are increasingly being replaced by tropical pasture grasses which are  $C_4$  type, characteristics of  $C_4$  metabolism such as high temperature tolerance (see Fig. 1c) and lower stomatal conductance could alter regional hydrology and surface energy balance (Shukla *et al.* 1990). A realistic model of leaf gas exchange in  $C_4$  plants such as that reported here can be incorporated as an interactive part of an atmospheric climate model to predict the effects of land use change in the tropics.

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## Appendix A.

### Intercellular Transport Model of C<sub>4</sub> Photosynthesis

Berry and Farquhar (1978) proposed the following expression to describe C<sub>4</sub> photosynthesis:

$$A = W_p - L, \quad (1A)$$

where  $W_p$  is the activity of PEP carboxylase and  $L$  is the leak of inorganic carbon between the bundle sheath cells and the intercellular spaces of the leaf mesophyll.  $A$  is assumed to be equal to the rate of net CO<sub>2</sub> fixation in the bundle sheath chloroplasts which is a function



of the CO<sub>2</sub> and O<sub>2</sub> partial pressures in the bundle sheath cells. Using the expression for rubisco-limited C<sub>3</sub> photosynthesis which takes into account photorespiration (Farquhar *et al.* 1980),  $A$  can also be expressed as

$$A = \frac{V_{\max} \left( p_{bs} - \frac{0.5 O_{bs}}{\tau} \right)}{p_{bs} + K_c \left( 1 + \frac{O_{bs}}{K_o} \right)}. \quad (2A)$$

Likewise the light-limited net CO<sub>2</sub> fixation in the bundle sheath is given as

$$A = \alpha \alpha_r f Q_p \frac{p_{bs} - \frac{0.5 O_{bs}}{\tau}}{p_{bs} + \frac{1.17 O_{bs}}{\tau}}, \quad (3A)$$

where  $V_{\max}$  is the capacity for CO<sub>2</sub> fixation by rubisco at saturating substrate concentrations,  $K_c$  is the Michaelis Constant for CO<sub>2</sub>,  $K_o$  is the competitive inhibition constant,  $p_{bs}$  and  $O_{bs}$  are the bundle sheath CO<sub>2</sub> and O<sub>2</sub> partial pressures respectively,  $\tau$  is the specificity of rubisco for CO<sub>2</sub> relative to O<sub>2</sub>,  $\alpha_r$  is the intrinsic quantum efficiency for CO<sub>2</sub> fixation (or the reciprocal of the quantum requirement for the production of RuBP),  $Q_p$  is the quantum flux density incident on the leaf and  $f$  is the fraction of absorbed  $Q_p$  that is used for production of RuBP. The constant 1.17 reflects the extra light-generated ATP required by photorespiration (see Berry and Farquhar 1978).

We used somewhat simpler expressions than Berry and Farquhar (1978) for the dependence of  $W_p$  on CO<sub>2</sub> and  $Q_p$ . CO<sub>2</sub>-limited mesophyll photosynthesis is defined as

$$W_p = k_p p_i / P, \quad (4A)$$

where  $p_i$  is the CO<sub>2</sub> concentration in the intercellular spaces of the mesophyll. The apparent first order rate constant for CO<sub>2</sub>,  $k_p$ , is in theory the ratio of  $V_{\max}/K_m$  of PEP carboxylase at the pH of the mesophyll cells. Empirically it may be derived from the initial slope of CO<sub>2</sub> response curves of leaves (Fig. 2 and Table 1). Note that  $k$  derived in this manner would also include any non-stomatal resistance to CO<sub>2</sub> transport from the intercellular spaces to the sites of fixation.

The light-limited  $W_p$  is defined as

$$W_p = \alpha \alpha_p (1 - f) Q_p, \quad (5A)$$

where  $\alpha_p$  is the reciprocal of the quantum requirement for the production of PEP. The term  $1 - f$  refers to the proportion of absorbed  $Q_p$  linked to the production of PEP.

The leak of CO<sub>2</sub> between the bundle sheath and the mesophyll is simply expressed by

$$L = (p_{bs} - p_i) / Pr_c, \quad (6A)$$

where  $r_c$  is the resistance to CO<sub>2</sub> diffusion through the bundle sheath cell walls and membranes. Large gradients in both HCO<sub>3</sub><sup>-</sup> and CO<sub>2</sub> concentrations between the bundle sheath and the mesophyll may occur but the resistance to HCO<sub>3</sub><sup>-</sup> diffusion is estimated to be ten times larger than to CO<sub>2</sub> diffusion (Furbank *et al.* 1989) so we neglect the former.

The rate of O<sub>2</sub> production and its diffusion between the bundle sheath and mesophyll is represented as

$$\lambda A = (O_{bs} - O_i) / P D r_c, \quad (7A)$$

where  $\lambda$  is the proportion of total net O<sub>2</sub> production occurring in the bundle sheath (corn bundle sheath chloroplasts show little photosystem II activity so  $\lambda$  is assumed to be small).

$O_i$  is the partial pressure of  $O_2$  in the intercellular spaces and  $D$  is a constant that takes into account the different solubilities and diffusivities of  $CO_2$  and  $O_2$  assuming that both gases diffuse along the same pathway.

Solving Eqns 6A and 7A for  $p_{bs}$  and  $O_{bs}$  respectively, and substituting these solutions into either the  $CO_2$  limited (Eqns 2A and 4A) or the light limited (Eqns 3A and 5A) equations and combining with Eqn 1A yields a quadratic solution for either the  $CO_2$  or light limited  $A_n$  in terms of the variables  $p_i$ ,  $O_i$ , and  $Q_p$ :

$$aA^2 + bA + c = 0, \quad (8A)$$

where the coefficients are defined for the  $CO_2$  and light limited conditions respectively as

$$a = \begin{cases} P \left( \frac{\lambda r_c D K_c}{K_o} - r_c \right) \\ P \left( \frac{\lambda D 1.17}{\tau} - r_c \right) \end{cases};$$

$$b = \begin{cases} Pr_c p_i k_p + p_i + K_c + \frac{O_i K_c}{K_o} + \frac{0.5 \lambda Pr_c D V_{max}}{\tau} + V_{max} Pr_c \\ a\alpha_p(1-f)Pr_c Q_p + p_i + \frac{1.17 O_i}{\tau} + \frac{0.5 a\alpha_r f Q_p \lambda Pr_c D}{\tau} + a\alpha_r f Q_p Pr_c \end{cases};$$

$$c = \begin{cases} V_{max} \left( \frac{0.5 O_i}{\tau} - Pr_c k_p p_i - p_i \right) \\ a\alpha_r f Q_p \left( \frac{0.5 O_i}{\tau} - a\alpha_p(1-f)Q_p Pr_c - p_i \right) \end{cases}.$$

The predicted response of the cellular model to  $CO_2$  using parameters listed in Table 1 in comparison to those predicted by the leaf model and measured from a corn leaf at  $25^\circ C$  is shown in Fig. 2. The  $CO_2$  limited region of the curve expresses the limitation imposed by PEP carboxylase activity while the  $CO_2$  saturated region expresses the limitation by the rubisco capacity ( $V_{max}$ ). The curvature in the transition between limitations is primarily a function of  $r_c$ ,  $K_c$ ,  $K_o$  and the ratio of  $k_p$  to  $V_{max}$ . Increasing  $r_c$ ,  $K_o$ ,  $k_p/V_{max}$  or decreasing  $K_c$  causes more abrupt saturation with respect to  $CO_2$ . Other parameters ( $O_i$ ,  $\lambda$ ,  $\tau$ ,  $D$ ) have relatively small influence on curvature in the region of transition between  $CO_2$  limited and  $CO_2$  saturated rates. The cellular model overestimates  $A$  at high  $CO_2$  levels because in this simulation  $A$  is assumed to be independent of the quantum flux while measured rates and those predicted by the simplified model are not completely light saturated.

Preliminary analysis of Eqn 8A has shown some interesting behaviours of the ICT model. First, there exists an optimum value for  $f$  which is dependent in particular on  $r_c$  and  $Q_p$ . The optimum value for  $f$  (the value which maximises  $A$ ) is close to 0.6. The quantum requirement for RuBP re-generation in the absence of photorespiration relative to the total quantum requirement for PEP and RuBP re-generation is about 0.6. The optimum value of  $f$  increases as  $Q_p$  goes to zero.

Another interesting model response occurs at very low quantum fluxes. When  $A$  approaches zero it is obvious that  $p_{bs}$  will equal  $p_i$  (in the absence of dark respiration,  $R_d$ ). As  $Q_p$  increases,  $p_{bs}$  increases, which causes the quantum yield to increase as photorespiration is inhibited. When  $Q_p$  is above about  $50 \mu\text{mol m}^{-2} \text{s}^{-1}$  photorespiration is largely inhibited and the quantum yield becomes constant with increasing  $Q_p$ . The large changes in the degree to which rubisco is  $CO_2$ -limited over the range of quantum fluxes from  $0$ – $50 \mu\text{mol m}^{-2} \text{s}^{-1}$  implies that discrimination by rubisco against  $^{13}\text{C}$  relative to  $^{12}\text{C}$  will decrease with increasing light over this range.

## Appendix B.

### Equations Describing the Coupled Photosynthesis–Stomatal Conductance Model

The C<sub>4</sub> leaf model is derived from the following equations that define the three unknowns, stomatal conductance ( $g$ ), net photosynthesis ( $A_n$ ) and the partial pressure of CO<sub>2</sub> in the intercellular spaces of the leaf ( $p_i$ ).

Following Ball *et al.* (1987) stomatal conductance is given by

$$g = m \frac{h_s A_n P}{p_s} + b, \quad (1B)$$

where  $h_s$  is the leaf surface relative humidity,  $P$  is atmospheric pressure,  $p_s$  is leaf surface partial pressure of CO<sub>2</sub> and  $m$  and  $b$  are the slope and intercept of a linear regression with respect to the other variables in the equation.

Gross photosynthesis ( $A$ ) is given as a function of incident quantum flux density ( $Q_p$ ) and the intercellular partial pressure of CO<sub>2</sub> ( $p_i$ ) and leaf temperature ( $T_l$ ) in the form of a pair of nested quadratic equations. The first quadratic is expressed as

$$\theta M^2 - M(V_T + \alpha Q_p) + V_T \alpha Q_p = 0, \quad (2B)$$

where  $V_T$  is the temperature-dependent, substrate-saturated rubisco capacity,  $\alpha$  is the quantum efficiency and  $M$  is the flux determined by the rubisco and light limited capacities. The curvature parameter  $\theta$  gives a gradual transition between the light limited and  $V_T$  limited flux. The limitation on the overall rate by  $M$  and the CO<sub>2</sub> limited flux, ( $k_T p_i / P$ ), where  $k_T$  is the temperature-dependent pseudo-first order rate constant with respect to  $p_i$ , is likewise expressed as a quadratic

$$\beta A^2 - A(M + k_T p_i / P) + M k_T p_i / P = 0. \quad (3B)$$

$\beta$  is analogous to  $\theta$  and specifies the degree of co-limitation between  $M$  and the CO<sub>2</sub> limited flux. The smaller roots are the appropriate solutions for both quadratics.  $A_n$  is defined as

$$A_n = A - R_T, \quad (4B)$$

where  $R_T$  is the temperature-dependent rate of leaf respiration. The temperature dependence of the substrate saturated rubisco capacity ( $V_{\max}$ ), the pseudo-first order rate constant with respect to CO<sub>2</sub> ( $k$ ) and the leaf respiration ( $R_T$ ) are given as  $Q_{10}$  functions.  $V_{\max}$  and  $R_d$  are also adjusted with inhibition functions that place upper and lower temperature limits on  $A_n$ . These limits were adjusted to approximate responses of C<sub>4</sub> plants to extreme temperatures (see Berry and Björkman 1980). The temperature dependencies are given as follows.

$$\begin{aligned} V_T &= \frac{V_{\max} Q_{10}^{\frac{T_l - 25}{10}}}{(1 + e^{0.3(13 - T_l)})(1 + e^{0.3(T_l - 36)})}; \\ R_T &= \frac{R_d Q_{10}^{\frac{T_l - 25}{10}}}{1 + e^{1.3(T_l - 55)}}; \\ k_T &= k Q_{10}^{\frac{T_l - 25}{10}}, \end{aligned} \quad (5B)$$

where  $Q_{10}$  is the proportional increase in a parameter value for a 10°C increase in leaf temperature (Table 2).

Finally, the partial pressure of CO<sub>2</sub> in the intercellular spaces is defined as

$$p_i = p_s - \frac{1 \cdot 6 A_n P}{g} \quad (6B)$$

The solution to these simultaneous equations can be obtained by substitution to yield a cubic equation which is described further in Appendix C. Alternatively, a numerical solution may be used which could include consideration of diffusion through the leaf boundary layer and of the leaf energy balance, as was done for the C<sub>3</sub> model of Collatz *et al.* (1991).

### Appendix C.

#### Analytical Solution to the Coupled Photosynthesis-Stomatal Conductance Model

Eqns 1B through 6B can be combined to eliminate  $p_i$  and  $g$  to yield a cubic equation,

$$a'A_n^3 + b'A_n^2 + c'A_n + d' = 0, \quad (1C)$$

where the coefficients are defined as

$$\begin{aligned} a' &= \beta m h_s \frac{P}{p_s}; \\ b' &= \beta b - \beta m h_s R_T \frac{P}{p_s} + 1 \cdot 6 k_T - M m h_s \frac{P}{p_s} - k_T m h_s; \\ c' &= R_T M m h_s \frac{P}{p_s} - M b + R_T k_T m h_s - k_T b \frac{P}{P} - 1 \cdot 6 k_T R_T + M k_T m h_s - 1 \cdot 6 k_T M; \\ d' &= M k_T b \frac{P}{P} + 1 \cdot 6 M k_T R_T - M k_T m h_s R_T. \end{aligned}$$

An explicit solution to Eqn 1C can be obtained and the same root was found to be applicable for widely different CO<sub>2</sub>, light and temperature conditions. The solution is given by the following root

$$A_n = -2\sqrt{Q} \cos\left(\frac{S + 4\pi}{3}\right) - \frac{b'}{3a'}, \quad (2C)$$

where  $Q$ ,  $S$  and  $R$  are defined as

$$\begin{aligned} Q &= \frac{\left(\frac{b'}{a'}\right)^2 - 3\frac{c'}{a'}}{9}; \\ R &= \frac{2\left(\frac{b'}{a'}\right)^2 - 9\frac{b'c'}{a'^2} + 27\frac{d'}{a'}}{54}; \\ S &= \arccos\left(\frac{R}{\sqrt{Q^3}}\right) \end{aligned}$$

(see Press *et al.* 1989).

The solution for  $A_n$  can be substituted into Eqn 1B of Appendix B to give  $g$ . Fig. 8B and C show the responses of  $g$  calculated in this manner to surface relative humidity.