

For $n + 1 \leq s \leq 2n - 1$, there is no nonzero term since there is no i for which $k_i = 0$ (see Eq. (S24)), and thus

$$G_{2n-1,s} = 0, \quad n + 1 \leq s \leq 2n - 1. \quad (\text{S30})$$

Now we can use Eqs. (S25)–(S30) to show that any pair of columns of \mathbf{G} are linearly independent. For $1 \leq s \leq n$ and $n + 1 \leq s' \leq 2n - 1$, the last two components ($k = 2n - 2$ and $k = 2n - 1$) of the s th and s' th column vectors form the two-dimensional vectors $(n - 1, 1)^T$ and $(1, 0)^T$, respectively, which are linearly independent. This implies that the full $(2n - 1)$ -dimensional vectors in the s th and s' th columns of \mathbf{G} are also linearly independent. For $1 \leq s < s' \leq n$, since the last components ($k = 2n - 1$) of the s th and s' th column vectors are both equal to one, it suffices to show that the second component ($k = 2$) is different in order to establish that they are linearly independent. From Eq. (S26), we have

$$\begin{aligned} G_{2s} - G_{2s'} &= \sum_{t \neq s} \prod_{i \neq t,s} \bar{\alpha}_i - \sum_{t \neq s'} \prod_{i \neq t,s'} \bar{\alpha}_i \\ &= \left(\sum_{t \neq s,s'} \prod_{i \neq t,s} \bar{\alpha}_i + \prod_{i \neq s',s} \bar{\alpha}_i \right) - \left(\sum_{t \neq s,s'} \prod_{i \neq t,s'} \bar{\alpha}_i + \prod_{i \neq s,s'} \bar{\alpha}_i \right) \\ &= \sum_{t \neq s,s'} \left(\prod_{i \neq t,s,s'} \bar{\alpha}_i \right) (\bar{\alpha}_{s'} - \bar{\alpha}_s) > 0, \end{aligned} \quad (\text{S31})$$

since $\bar{\alpha}_{s'} > \bar{\alpha}_s$ and $\bar{\alpha}_i > 0, \forall i$, and hence all terms in the summation are positive. Thus, we have $G_{2s} \neq G_{2s'}$, implying that the s th and s' th column vectors are linearly independent. For $n + 1 \leq s < s' \leq 2n - 1$, the argument is similar to the case of $1 \leq s < s' \leq n$; it suffices to show that the first component ($k = 1$) is different, and Eq. (S25) yields

$$\begin{aligned} G_{1s} - G_{1s'} &= \sum_{t \neq \hat{s}} \prod_{i \neq t,\hat{s}} \bar{\alpha}_i - \sum_{t \neq \hat{s}'} \prod_{i \neq t,\hat{s}'} \bar{\alpha}_i \\ &= \sum_{t \neq \hat{s},\hat{s}'} \left(\prod_{i \neq t,\hat{s},\hat{s}'} \bar{\alpha}_i \right) (\bar{\alpha}_{\hat{s}'} - \bar{\alpha}_{\hat{s}}) > 0, \end{aligned} \quad (\text{S32})$$

where we have defined $\hat{s}' \equiv s' - n + 1$ (and recall that $\hat{s} = s - n + 1$). Combining all of the above, we have that the s th and s' th column vectors of \mathbf{G} are linearly independent for all pairs of distinct s and s' , which establishes that \mathbf{G} is non-singular.

The IFT can now be applied to Eq. (S18) to conclude that the functions $c_i(\varepsilon)$ and $d_i(\varepsilon)$ are continuously differentiable. Furthermore, their derivatives satisfy the set of equations obtained by substituting these functions into Eq. (S18) and differentiating both sides with respect to ε :

$$\frac{d}{d\varepsilon} F_k(c_1(\varepsilon), \dots, c_n(\varepsilon), d_1(\varepsilon), \dots, d_n(\varepsilon), \varepsilon) = 0. \quad (\text{S33})$$

Using Eq. (S20), this can be written as

$$\sum_{\{k_i\}} \chi(\sum_i k_i = k) \cdot \sum_{\ell=1}^n \left[\left(\prod_{i \neq \ell} a_i^{(k_i)}(\varepsilon) \right) \cdot \frac{da_\ell^{(k_\ell)}(\varepsilon)}{d\varepsilon} - \sum_{\sigma} \text{sign}(\sigma) \left(\prod_{i \neq \ell} E_{i\sigma(i)}^{(k_i)}(\varepsilon) \right) \cdot \frac{dE_{\ell\sigma(\ell)}^{(k_\ell)}(\varepsilon)}{d\varepsilon} \right] = 0, \quad (\text{S34})$$

where we have now written the dependence of $a_i^{(k_i)}$ on ε explicitly. Setting $\varepsilon = 0$ and using Eq. (S21), we obtain

$$\sum_{\ell=1}^n \left(\sum_{\{k_i\}} \prod_{i \neq \ell} a_i^{(k_i)}(0) \cdot \chi(\sum_i k_i = k) \right) \cdot \left[\frac{da_\ell^{(k_\ell)}(0)}{d\varepsilon} - \frac{dE_{\ell\ell}^{(k_\ell)}(0)}{d\varepsilon} \right] = 0. \quad (\text{S35})$$

Noting that $k_\ell = 0, 1, 2$ and that the derivatives are both zero if $k_\ell = 2$ (see Eqs. (S15) and (S17)), we can rearrange the summation in Eq. (S35) to write

$$\begin{aligned} & \sum_{s=1}^n \left(\sum_{\{k_i\}} \prod_{i \neq s} a_i^{(k_i)}(0) \cdot \chi(s, k, \{k_i\}) \right) \cdot [x'_s(0) - y'_s(0)] \\ & + \sum_{s=n+1}^{2n-1} \left(\sum_{\{k_i\}} \prod_{i \neq s} a_i^{(k_i)}(0) \cdot \chi(s, k, \{k_i\}) \right) \cdot [x'_s(0) - y'_s(0)] = 0, \end{aligned} \quad (\text{S36})$$

where we recall that $x_s(\varepsilon) = c_s(\varepsilon) = a_s^{(1)}(\varepsilon)$ for $s = 1, \dots, n$ and $x_s(\varepsilon) = d_{s-n+1}(\varepsilon) = a_s^{(0)}(\varepsilon)$ for $s = n+1, \dots, 2n-1$, and we use the notations $y_s(\varepsilon) = \bar{B}_{ss}(\varepsilon)$ for $s = 1, \dots, n$ and $y_s(\varepsilon) = \bar{\alpha}_{s-n+1}$ for $s = n+1, \dots, 2n-1$. From Eq. (S23), we see that Eq. (S36) is equivalent to

$$\sum_{s=1}^{2n-1} G_{ks} [x'_s(0) - y'_s(0)] = 0, \quad (\text{S37})$$

which can be put in vector form as $\mathbf{G}(\mathbf{x} - \mathbf{y}) = \mathbf{0}$ using the notations $\mathbf{x} = (x'_1(0), \dots, x'_{2n-1}(0))^T$ and $\mathbf{y} = (y'_1(0), \dots, y'_{2n-1}(0))^T$. Thus, since \mathbf{G} is non-singular, we have $\mathbf{x} = \mathbf{y}$, and hence $x'_i(0) = y'_i(0)$. It then follows that

$$c'_i(0) = x'_i(0) = y'_i(0) = \frac{d\bar{B}_{ii}(0)}{d\varepsilon} = \frac{d}{d\varepsilon} \left(\sum_{\ell=1}^n u_{i\ell} v_{i\ell} \gamma_\ell(\varepsilon) / \beta_+ \right) \Big|_{\varepsilon=0} = \sum_{\ell=1}^n u_{i\ell} v_{i\ell} \gamma'_\ell(0) / \beta_+ \quad (\text{S38})$$

for $i = 1, \dots, n$, where we used the expression for \bar{B}_{ij} from Eq. (S12). It also follows that

$$d'_i(0) = x'_{i+n-1}(0) = y'_{i+n-1}(0) = \frac{d}{d\varepsilon} (\bar{\alpha}_i) \Big|_{\varepsilon=0} = 0 \quad (\text{S39})$$

for $i = 1, \dots, n$, since $\bar{\alpha}_i$ is a constant that does not depend on ε . In particular, we have

$$c'_2(0) = \sum_{\ell=1}^n u_{2\ell} v_{2\ell} \gamma'_\ell(0) / \beta_+ \quad \text{and} \quad d'_2(0) = 0. \quad (\text{S40})$$

Now, recall the argument at the end of Sec. 1 that $d'_2(0) = 0$ implies $c'_2(0) = 0$ if λ^{\max} decreases along the path $\gamma(\varepsilon)$ in the β -space. In view of that argument, we see from Eq. (S40) that, if the eigenvalues of \mathbf{P} are all distinct, the vector $(\gamma'_1(0), \dots, \gamma'_n(0))^T$ is parallel to the hyperplane L whenever λ^{\max} decreases along the path, where L is uniquely defined by the equation

$$\sum_{i=1}^n u_{2i} v_{2i} (\beta_i - \beta_+) = 0. \quad (\text{S41})$$

In other words, any descending path on the λ^{\max} -landscape must be tangent to the hyperplane L at β_+ .

If the path $\gamma(\varepsilon)$ is not tangent to L (and hence $c'_2(0) \neq 0$), the expansion of $c_2(\varepsilon)$ and $d_2(\varepsilon)$ around $\varepsilon = 0$, which reads

$$\begin{aligned} c_2(\varepsilon) &= 1 + c'_2(0)\varepsilon + O(\varepsilon^2), \\ d_2(\varepsilon) &= \frac{1}{4} + O(\varepsilon^2), \end{aligned} \quad (\text{S42})$$

can be substituted into Eq. (S8) to obtain the following approximation for $\lambda^{\max}(\varepsilon)$:

$$\lambda^{\max}(\varepsilon) = \begin{cases} \lambda_-^{\max} - \beta_+ c'_2(0)\varepsilon/2 + O(\varepsilon^2), & \varepsilon \leq 0, \\ \lambda_-^{\max} + \beta_+ \sqrt{c'_2(0)\varepsilon/2} + O(\varepsilon), & \varepsilon > 0. \end{cases} \quad (\text{S43})$$

(Recall we are assuming $c'_2(0) \geq 0$.) This establishes that the point β_+ is a local minimizer of λ^{\max} along any path that transversally intersects with the hyperplane L at β_+ .

4 Explicit calculation of matrix \mathbf{G} for small n

Here, we derive an expression for the matrix \mathbf{G} in Eq. (S22) for $n = 2$ and for $n = 3$. As in Sec. 1–3 above, we assume $\beta_+ > 0$, or equivalently, $\alpha_2 > 0$.

Example 1: For $n = 2$, Eq. (S11) becomes

$$\begin{aligned} \det(\bar{\mathbf{J}} - \nu \mathbf{I}) &= \det(\nu^2 \mathbf{I} + \nu \bar{\mathbf{B}} + \mathbf{D}) = \det \begin{pmatrix} \nu^2 + \bar{B}_{11}\nu & \bar{B}_{12}\nu \\ \bar{B}_{21}\nu & \nu^2 + \bar{B}_{22}\nu + 1/4 \end{pmatrix} \\ &= (\nu^2 + \bar{B}_{11}\nu)(\nu^2 + \bar{B}_{22}\nu + 1/4) - \bar{B}_{12}\bar{B}_{21}\nu^2 \\ &= \nu^4 + (\bar{B}_{11} + \bar{B}_{22})\nu^3 + (\bar{B}_{11}\bar{B}_{22} - \bar{B}_{12}\bar{B}_{21} + 1/4)\nu^2 + \bar{B}_{11}\nu/4, \\ &= \nu^4 + \frac{\beta_1 + \beta_2}{\beta_+} \nu^3 + \left(\frac{\beta_1\beta_2}{\beta_+^2} + \frac{1}{4} \right) \nu^2 + \bar{B}_{11}\nu/4, \end{aligned}$$