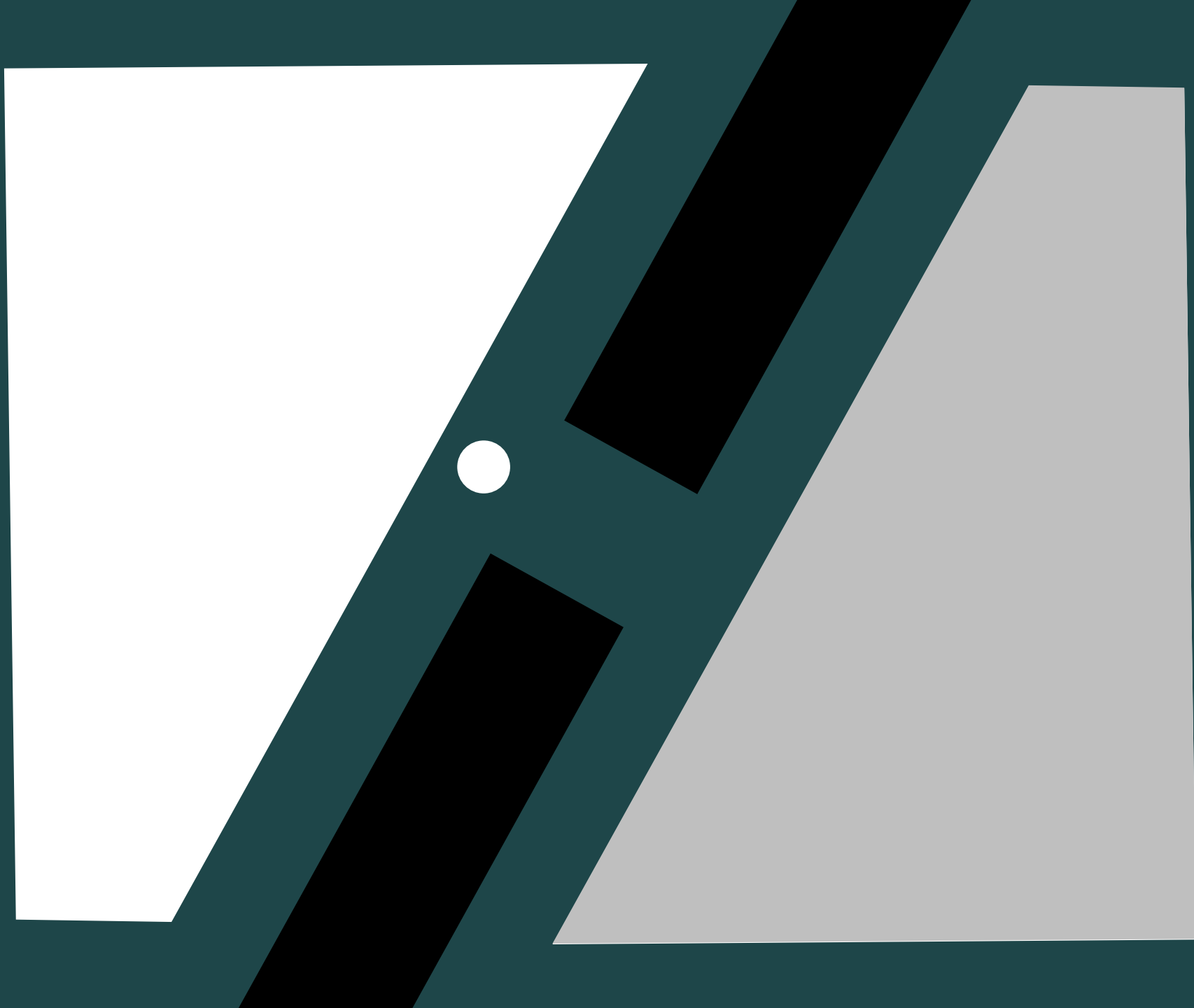


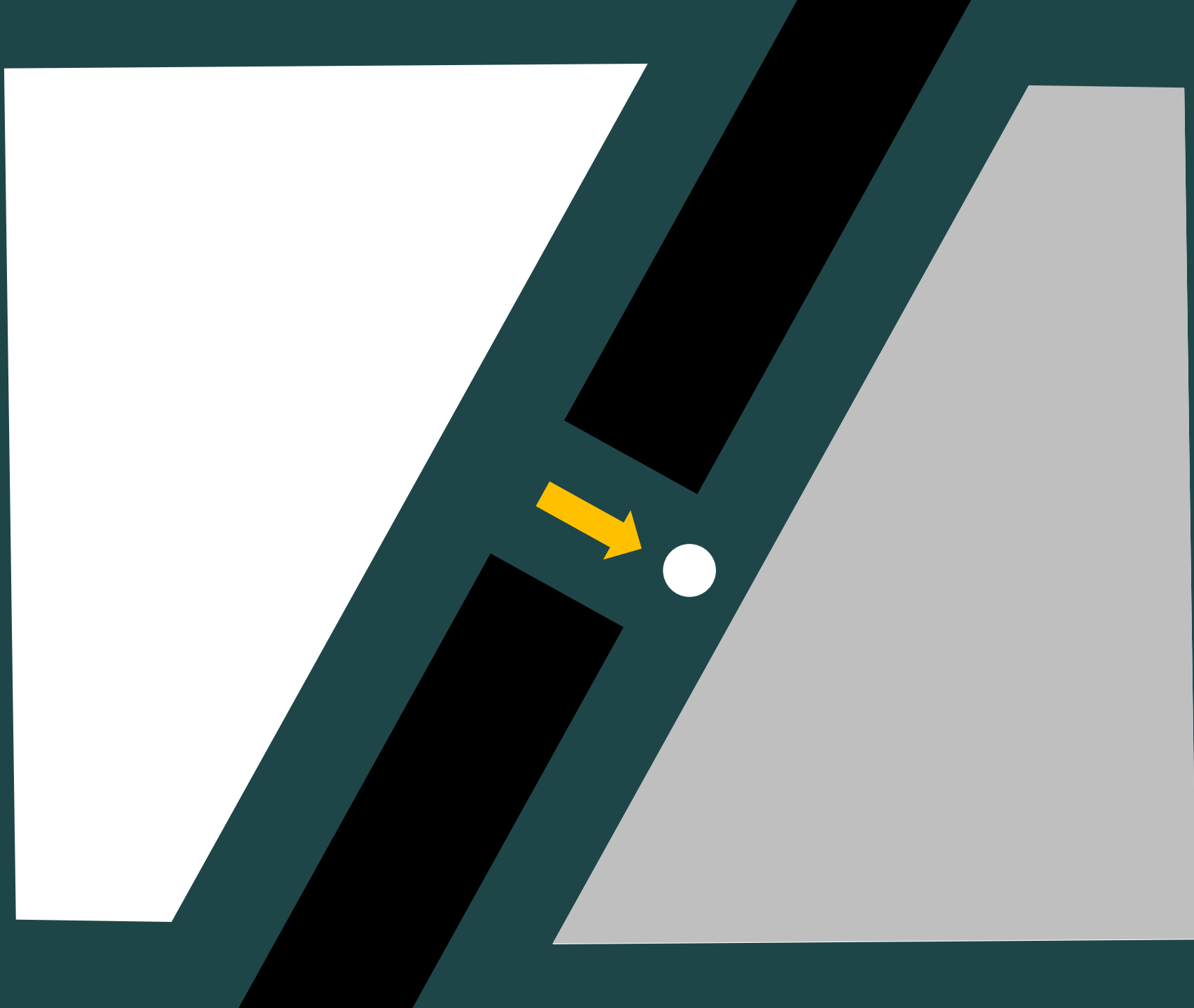
local tests of global entanglement and a counterexample to the generalized area law













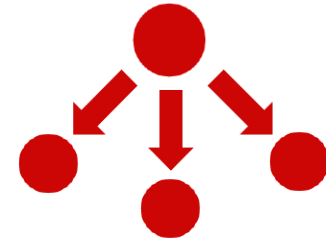






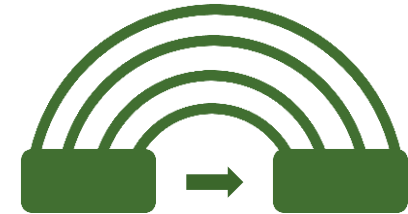
1 q. expanders

maximally entangled states



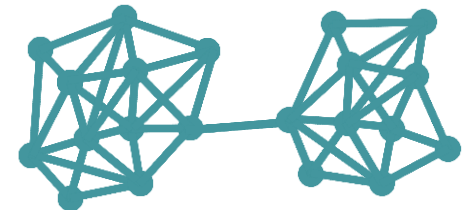
2 entanglement

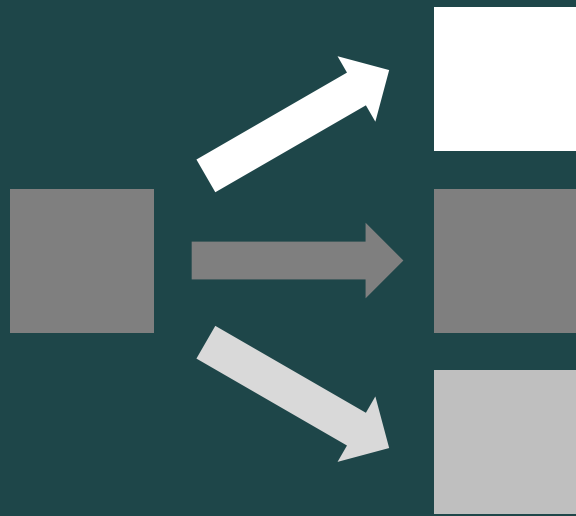
testing and communication



3 area law

gaps, connections, correlations

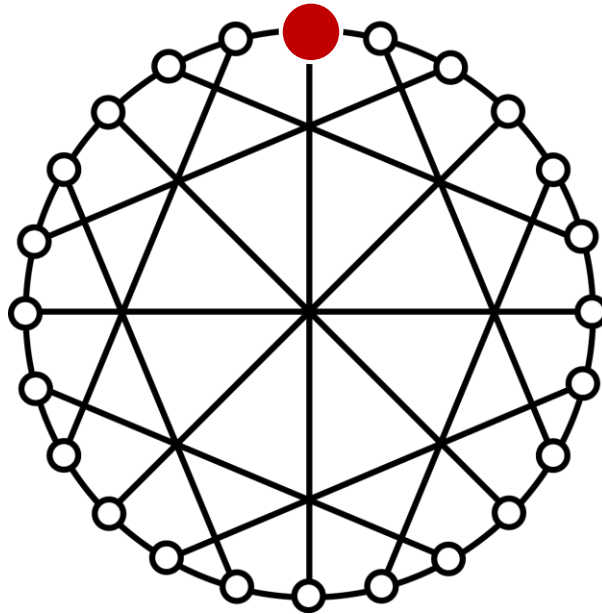




Quantum
Expanders

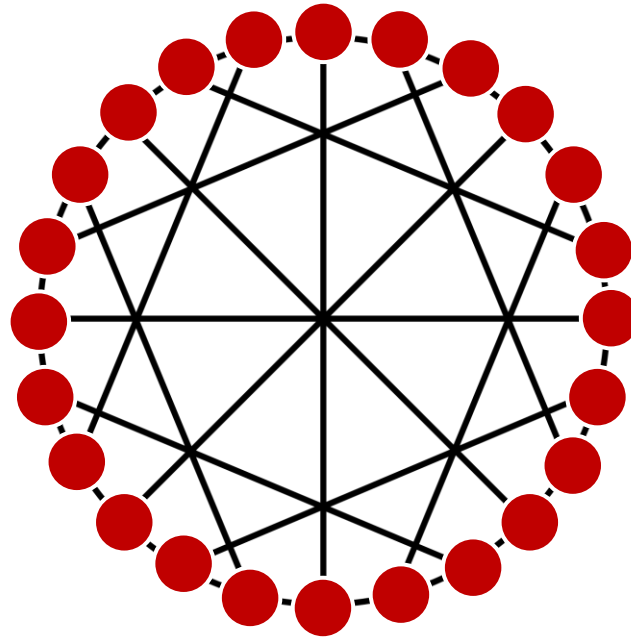
1 Classical expanders

- walk on these graphs? mix fast!



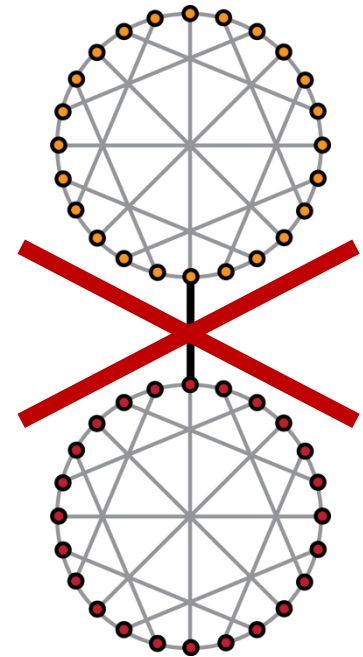
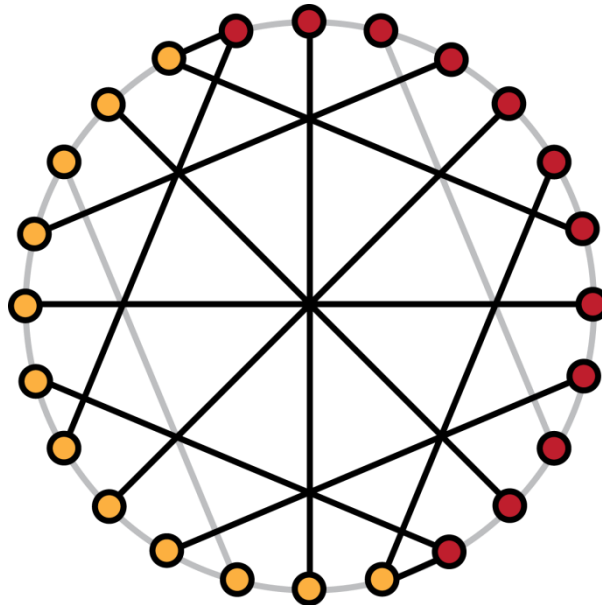
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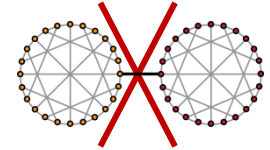
1 Classical expanders

- walk on these graphs? mix fast!
divide in two? cut a lot (fraction) of edges!

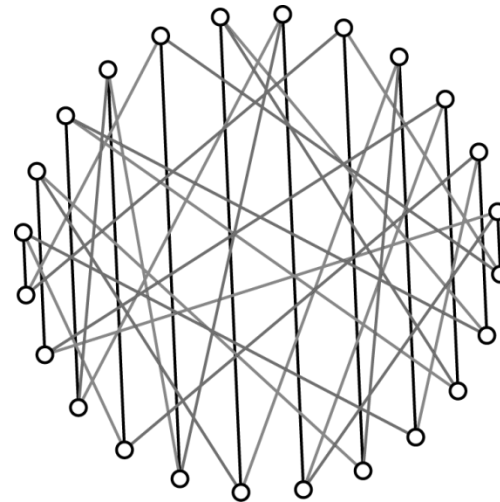
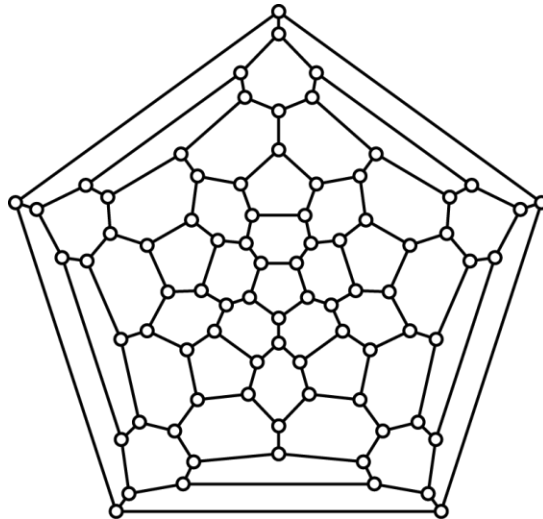


1 Classical expanders

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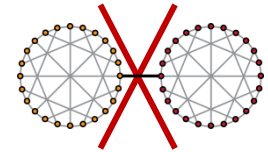


examples: Ramanujan, Cayley



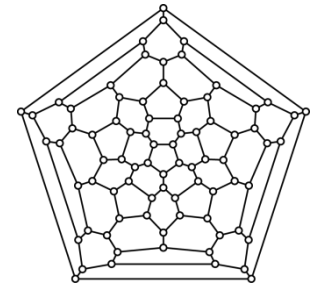
1 Classical expanders

- walk on these graphs? mix fast!
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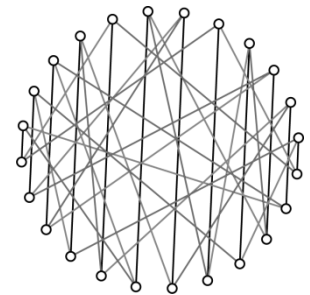


examples: Ramanujan, Cayley

- explicit, constant-degree
approximations to the full graph



normalized adjacency matrix
second largest eigenvalue $1-\lambda$



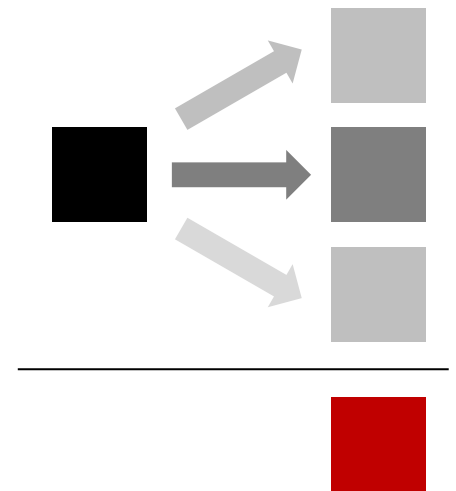
- a review [Hoory Linial Wigderson]
a talk [Harrow quantum expanders youtube]

1 Mixing up something quantum

- applying random unitaries

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

- classical expanders:
explicit, constant-degree
approximations to the full graph
fast-mixing

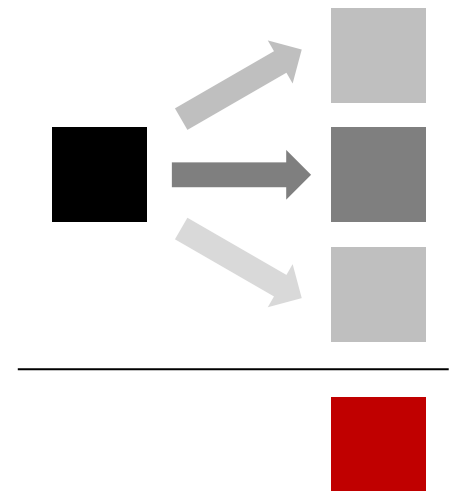


1 Quantum expanders

- applying random unitaries from a small set
a discrete approx. to Haar

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

- classical expanders:
explicit, constant-degree
approximations to the full graph
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1 Quantum expanders

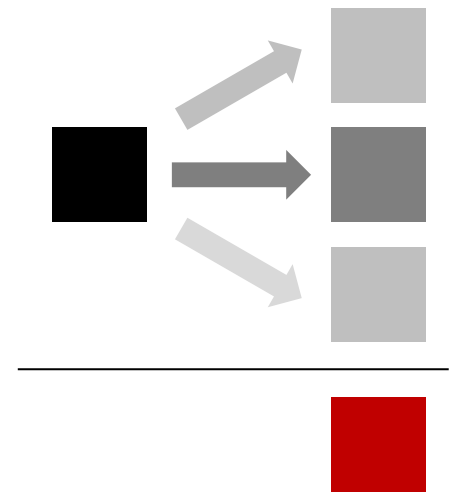
- applying random unitaries from a small set
a discrete approx. to Haar

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

- transform $N \times N$ matrices
I stays, everything else changes

$$\|\mathcal{E}(X)\|_2 \leq \lambda \|X\|_2$$

small 2nd largest sing. value



- QE constructions for fixed k (... , 8, 4, 3) $1 - \lambda \approx k^{-c}$

[Ben-Aroya+ '07, Hastings '07, Gross & Eisert '08, Hast. & Harr. '09, Gross '15]

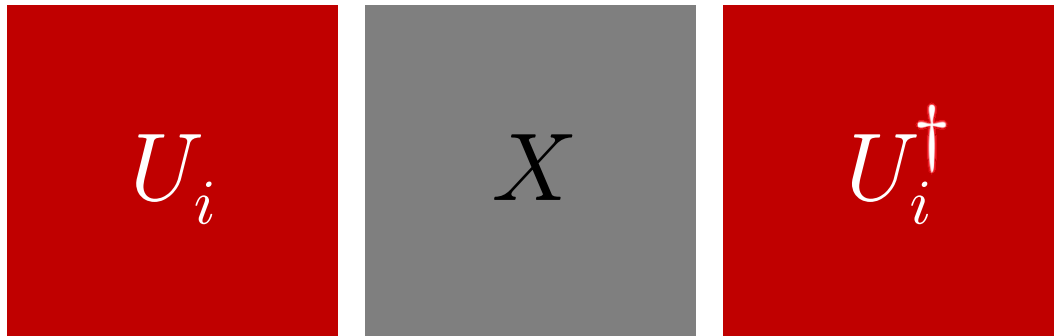
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



A diagram illustrating the commutation relation $U_i X = X U_i$. It consists of three small square boxes arranged horizontally. The first box is red and contains the symbol U_i in white. The second box is gray and contains the symbol X in black. The third box is red and contains the symbol U_i^\dagger in white.

$$U_i X = X U_i$$

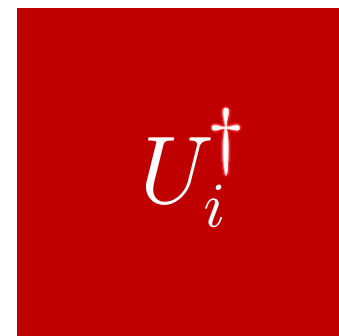
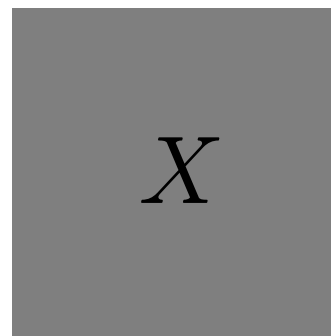
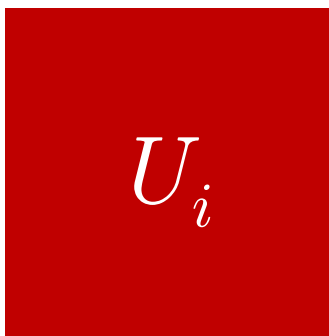
1 Quantum expanders

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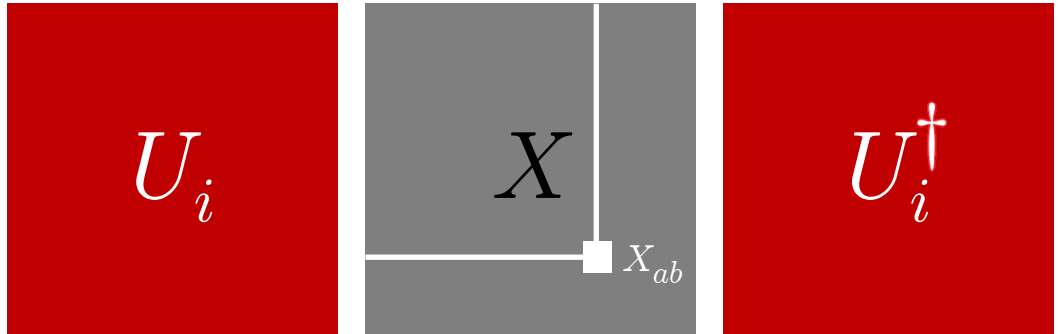
1 Quantum expanders

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$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



- interpreting matrices
as 2-register states

$$X = \sum_{a,b} X_{ab} |a\rangle \langle b|$$

density
matrix

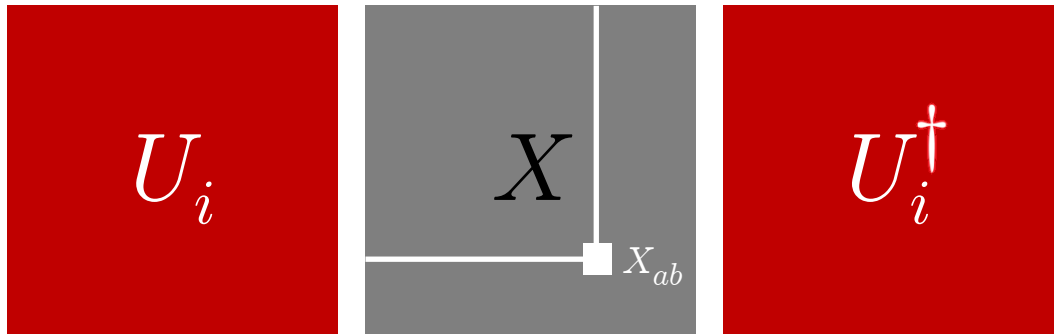
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



- interpreting matrices
as 2-register states

$$|X\rangle = \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

pure
state

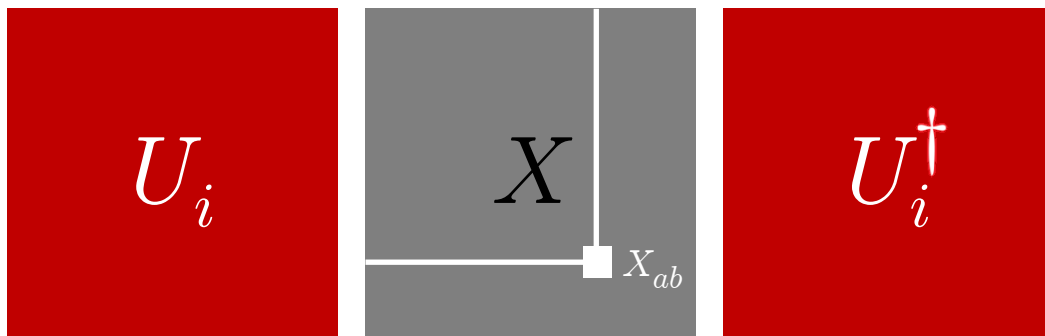
1 Quantum expanders

- transform $N \times N$ matrices

$$\mathcal{E}(X) = \frac{1}{k} \sum_{i=1}^k U_i X U_i^\dagger$$

a matrix that
doesn't change?

$$X = \mathbb{I}$$



- interpreting matrices
as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

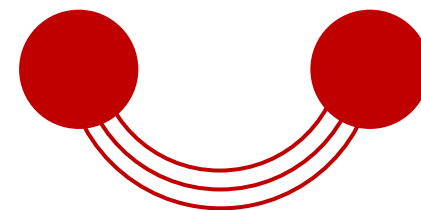
applying an expander distributively

a stationary state?

max. entangled! $\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$

1 Quantum expanders & 2 registers

- transform $N \otimes N$ states
close to the depolarizing channel



$$\|\tilde{\mathcal{E}} - |\phi_N\rangle\langle\phi_N|\| = \lambda$$

- interpreting matrices
as 2-register states

$$\frac{1}{k} \sum_{i=1}^k (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle|b\rangle$$

applying an expander distributively

a stationary state?

max. entangled! $\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$



local tests of global entanglement

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$



$O(\log\log(N) + \log(1/\lambda))$ qubits. [BDSW '96, BCGST '02]

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle|x\rangle$$



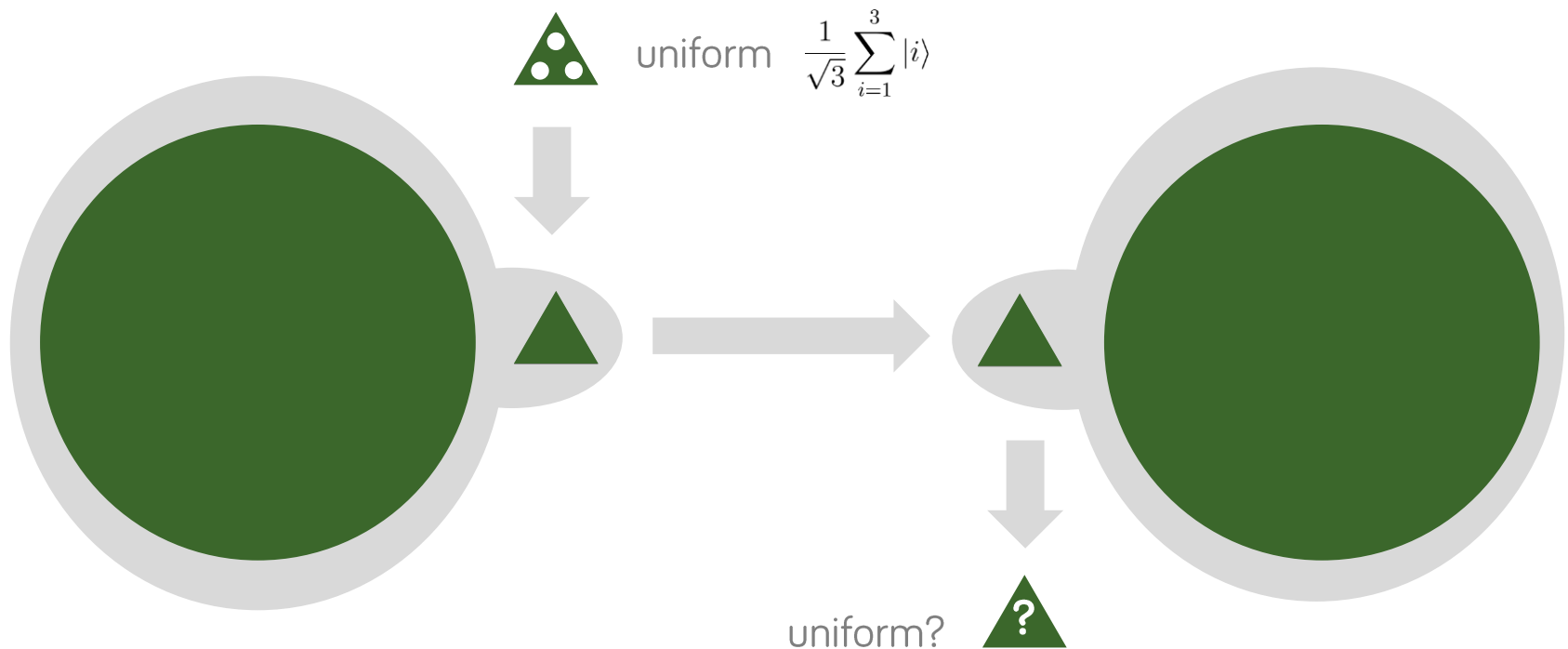
$O(\log(1/\lambda))$ qubits for error λ .

2 EPR testing

- how costly is it to certify that we share a maximally entangled state?
- apply a quantum expander distributively

$$\frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$$

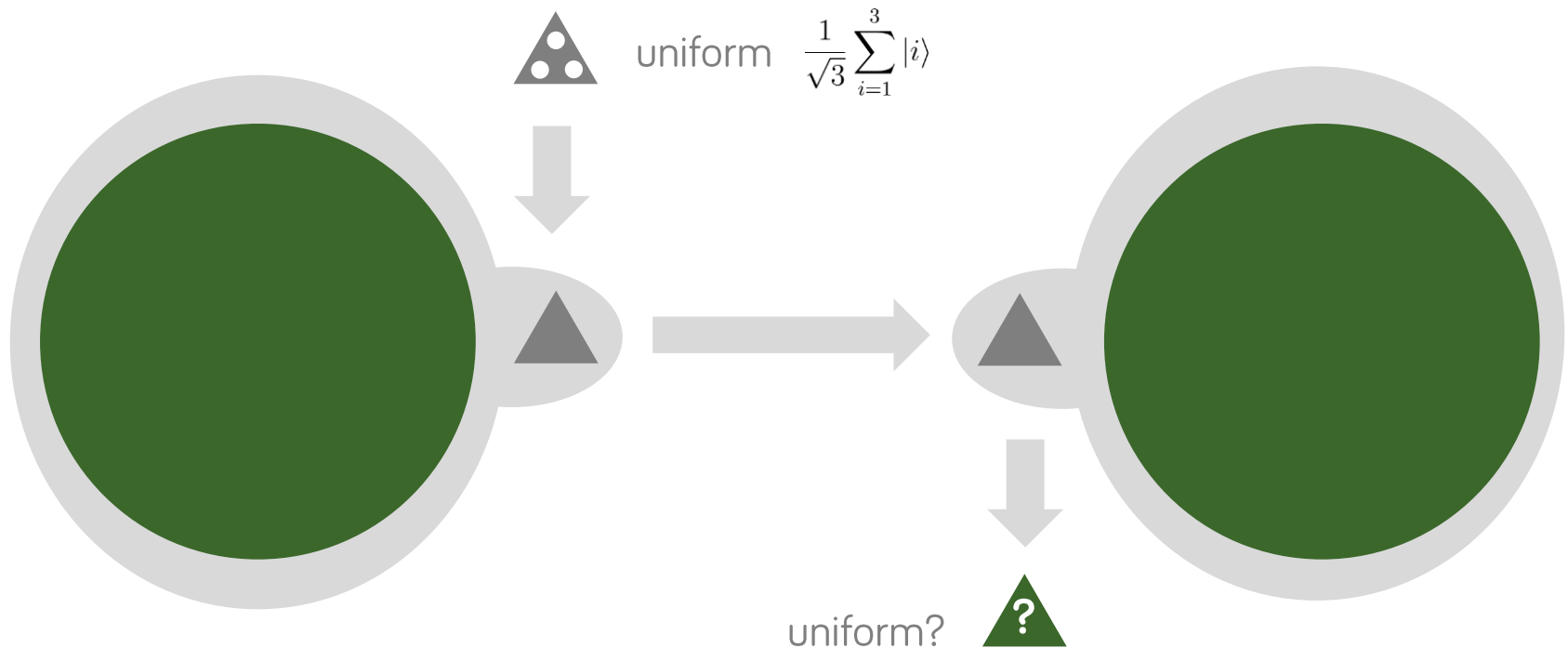
$$U_i \otimes U_i^*$$



2 EPR testing

- when does the qutrit remain uniform?

$$\frac{1}{\sqrt{3}} \sum_{i=1}^3 |i\rangle (U_i \otimes U_i^*) \sum_{a,b} X_{ab} |a\rangle |b\rangle$$

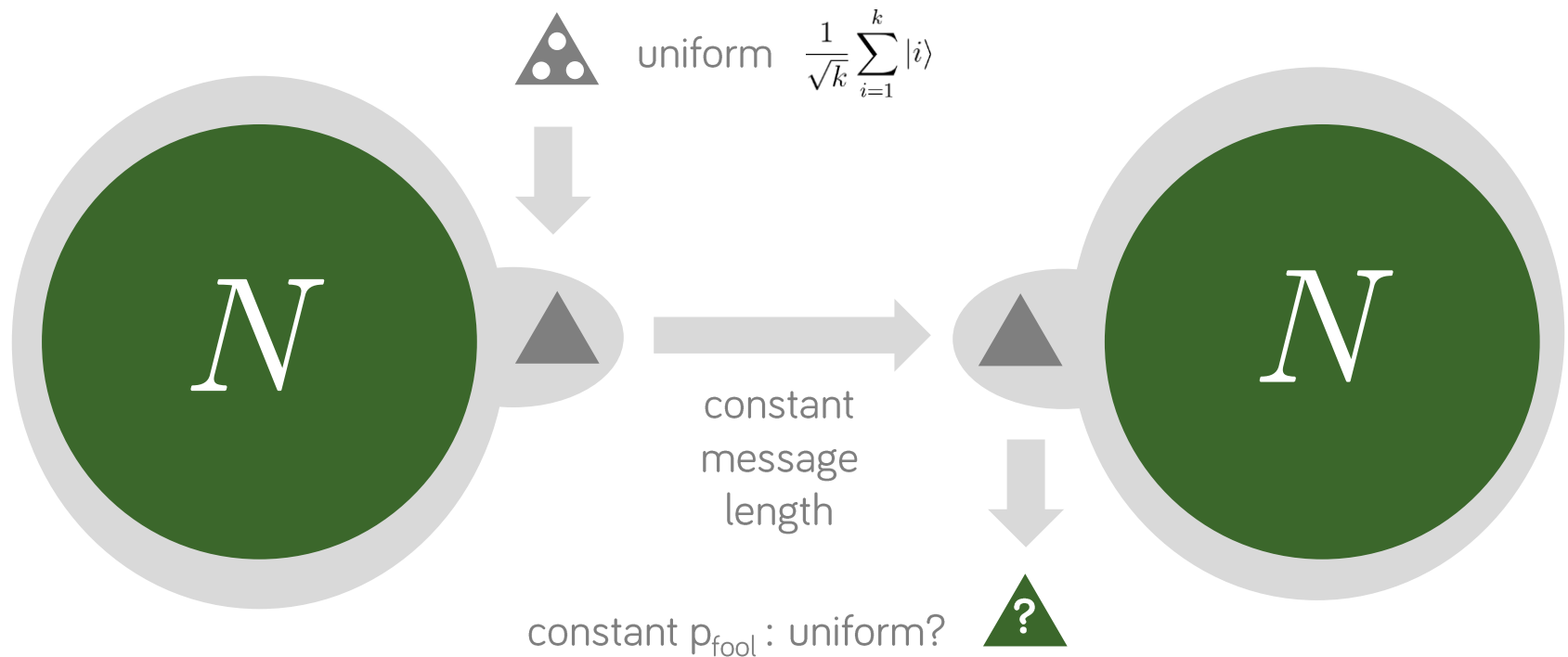


2 EPR testing

- when does the qutrit remain uniform?
for max. entangled X

$$\frac{1}{\sqrt{3}} \sum_{i=1}^3 |i\rangle (U_i \otimes U_i^*) \frac{1}{\sqrt{N}} \sum_{x=1}^N |x\rangle |x\rangle$$

- quantum expander property ... soundness

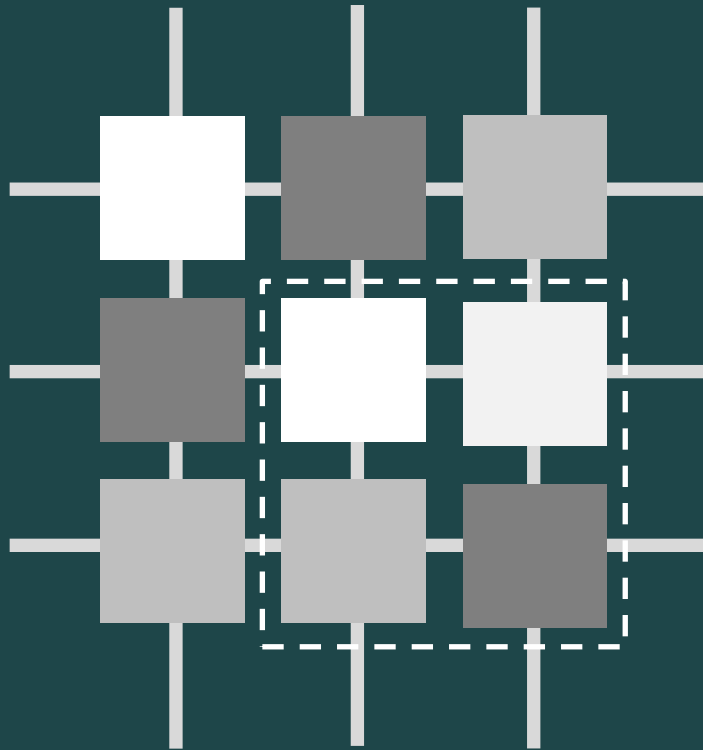


little communication

local tests

constant error

global correlations



a counterexample to the
generalized
area law

few connections

local interactions

constant gap

local correlations

few connections

local interactions

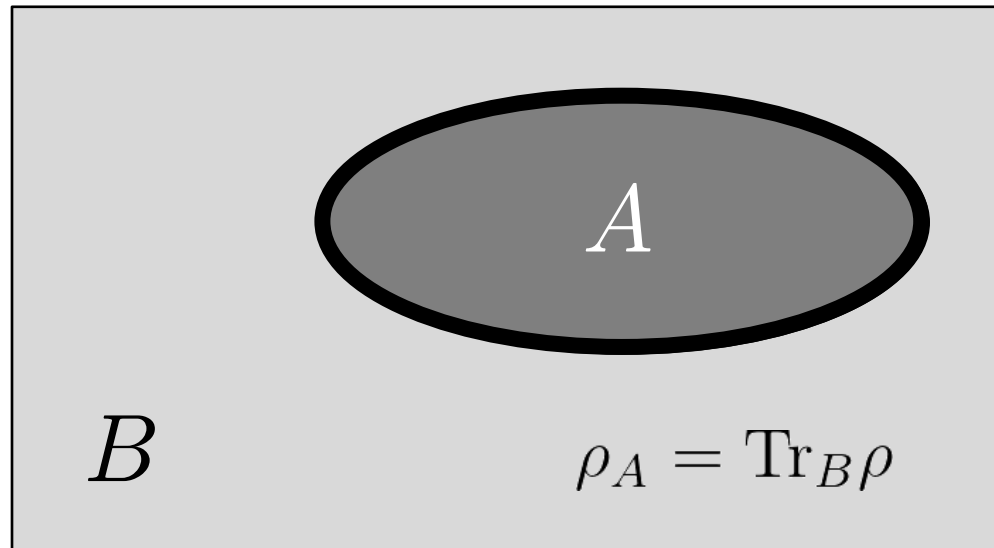
constant gap

global correlations

3 Ground states of gapped quantum spin systems

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \cancel{\text{volume}} \\ \text{surface area}$$



area law \longrightarrow a simple ground state?

3 Gapped 1D Hamiltonians

- Nothing closer than Δ to the ground state.

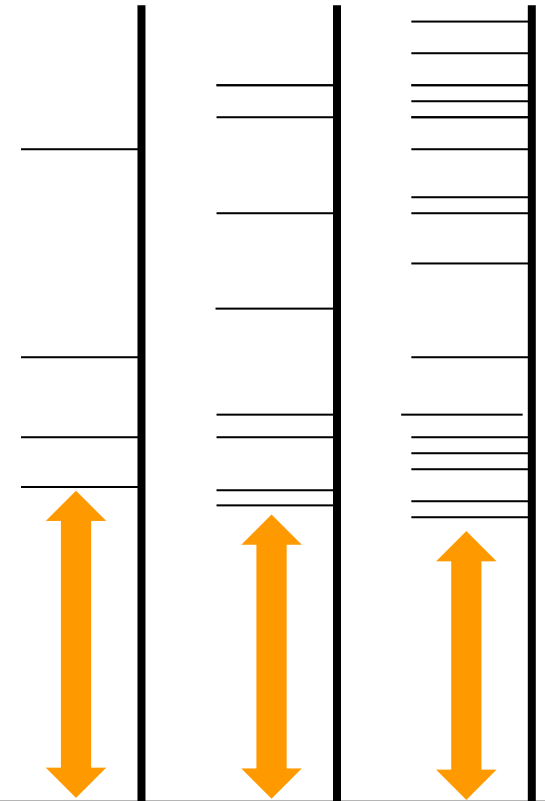
the AKLT (spin-1) chain

$$\sum_{j=1}^{N-1} \vec{S}_j \cdot \vec{S}_{j+1} + \frac{1}{3} \left(\vec{S}_j \cdot \vec{S}_{j+1} \right)^2$$

a biased walk in 1D

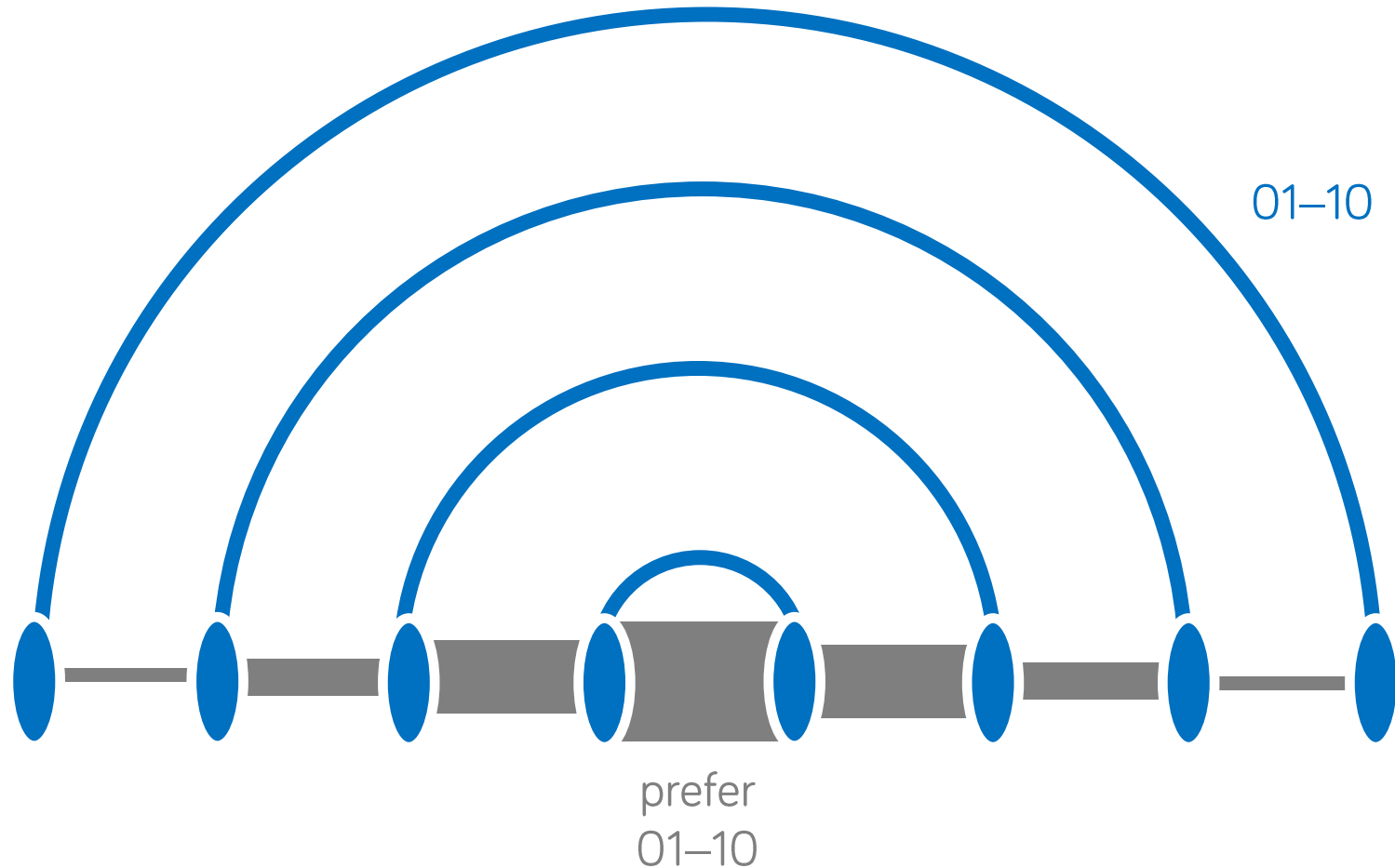
$$\sum_{j=1}^{N-1} (|j\rangle - B|j+1\rangle) (\langle j| - B\langle j+1|)$$

$N \rightarrow \infty$



exp. falloff of correlations, MPS/DMRG work well

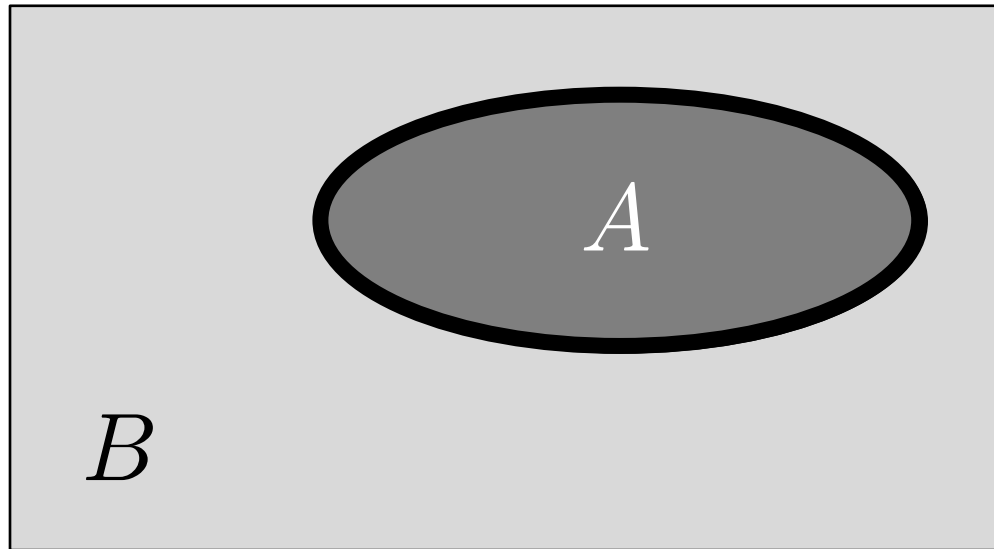
Without a gap, the entropy can be large. [Verstraete, Latorre+]



3 Ground states of gapped H's & the area law

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \cancel{\text{volume}} \\ \text{surface area}$$

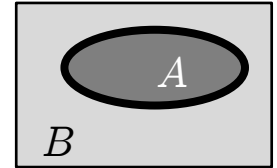


a gapped system → an area law

3 Ground states of gapped H's & the area law

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \cancel{\text{volume}} \text{ surface area}$$



1D ... theorems [Hastings 07, Arad+ 13]

algorithms [White 92, Vidal 03, Landau+ 13]

2D ... we're close

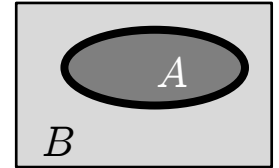
a gapped system → **an area law**

3 Ground states of gapped H's & the area law

- entanglement entropy

$$S = -\text{Tr}(\rho_A \ln \rho_A) \sim \text{volume}$$

surface area

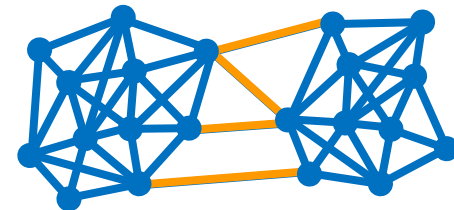


1D ... theorems [Hastings 07, Arad+ 13]

algorithms [White 92, Vidal 03, Landau+ 13]

2D ... we're close

- generalized area conjecture
entropy \sim cut size



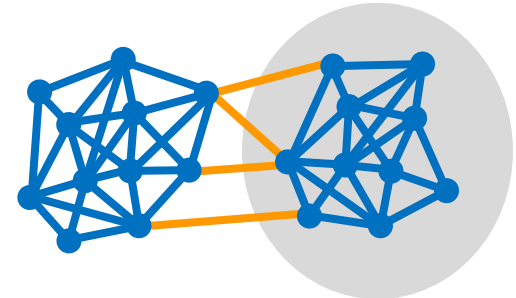
NO

a gap
a few links
 $O(1)$ terms



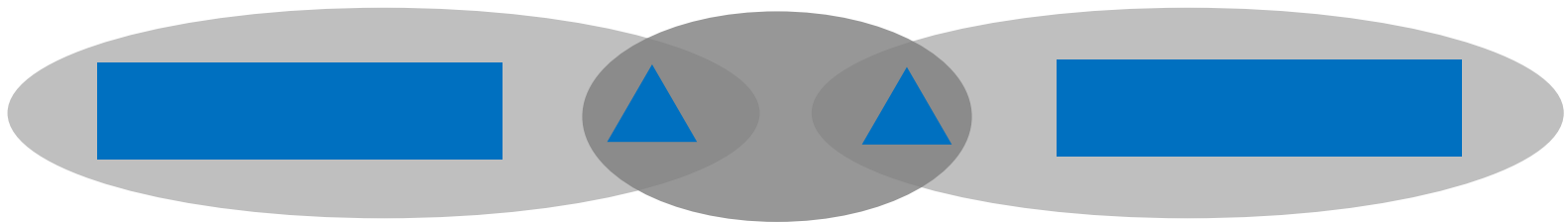
not much
entanglement
(a “simple” ground state)

- generalized area conjecture
entropy \sim cut size



3 Generalized area conjecture: the counterexample

- an $N \times 3 \times 3 \times N$ dimensional system
- a gapped, frustration-free Hamiltonian
an $O(1)$ interaction between the qutrits



3 Generalized area conjecture: the counterexample

- an $N \times 3 \times 3 \times N$ dimensional system
- a gapped, frustration-free Hamiltonian
an $O(1)$ interaction between the qutrits
- a unique, very entangled ground state
 $O(N)$ entanglement entropy across the cut

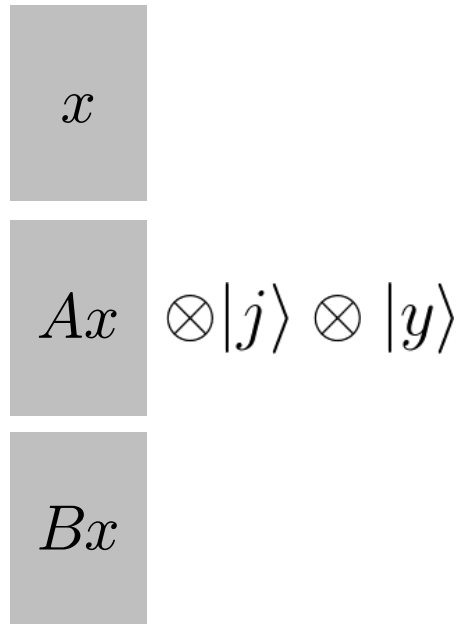


3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

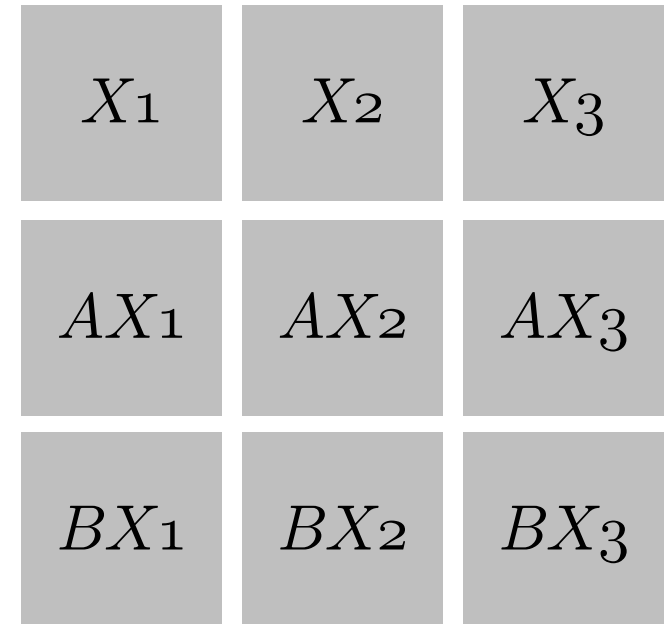
- a projector P_L with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$$

- as a vector



- as a matrix



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_R
with ground states

$$\frac{1}{\sqrt{3}} (|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle)$$

- as a vector $|i\rangle \otimes |x\rangle \otimes$

y	yA	yB
-----	------	------

as a matrix

Y_1	Y_1A	Y_1B
Y_2	Y_2A	Y_2B
Y_3	Y_3A	Y_3B



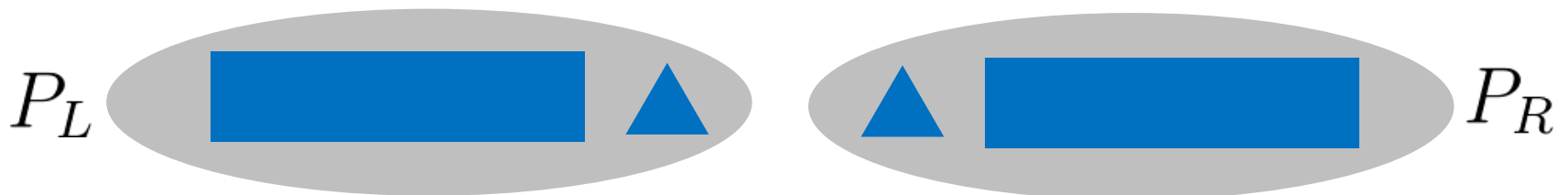
3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L
- a projector P_R

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

X_1	X_2	X_3	X	XA	XB	Y_1	Y_1A	Y_1B
AX_1	AX_2	AX_3	AX	AXA	AXB	Y_2	Y_2A	Y_2B
BX_1	BX_2	BX_3	BX	BXA	BXB	Y_3	Y_3A	Y_3B



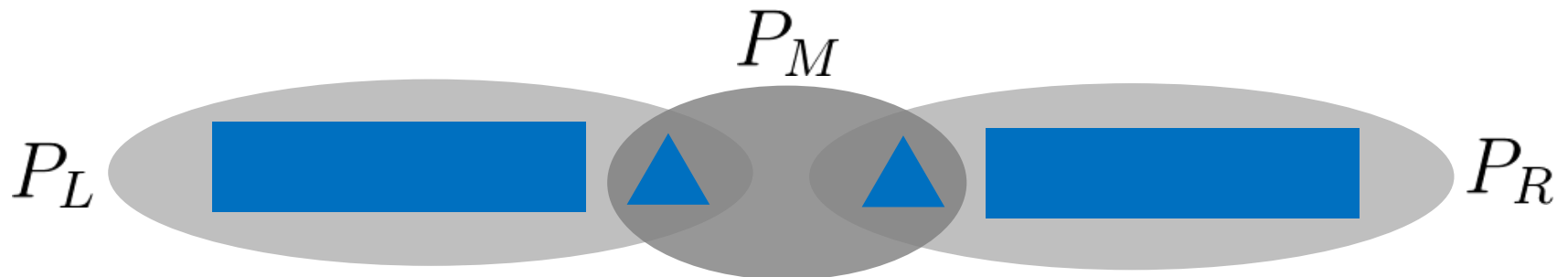
3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L
- a projector P_R
- a projector P_M

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

X	XA	XB
AX	AXA	AXB
BX	BXA	BXB



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L
- a projector P_R
- a projector P_M

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

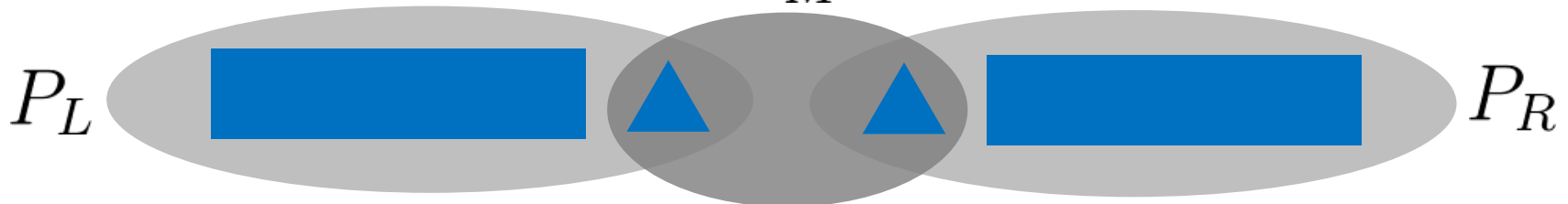
enforce symmetry: 12 & 21
13 & 31

- who commutes with A and B ?

only the identity,
as $[I, A, B]$ is
a q. expander

X	XA	XB
AX	AXA	AXB
BX	BXA	BXB

P_M



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

- a projector P_L
- a projector P_R
- a projector P_M

$$(|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle) / \sqrt{3}$$

$$(|1\rangle|y\rangle + |2\rangle A|y\rangle + |3\rangle B|y\rangle) / \sqrt{3}$$

enforce symmetry: 12 & 21
13 & 31

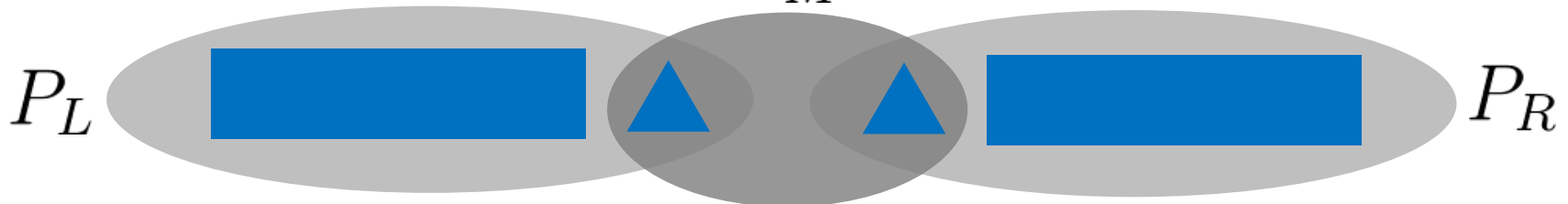
- who commutes with A and B ?

only the identity,
as $[I, A, B]$ is
a q. expander

I	A	B
A	AA	AB
B	BA	BB

$$\frac{1}{3\sqrt{N}}$$

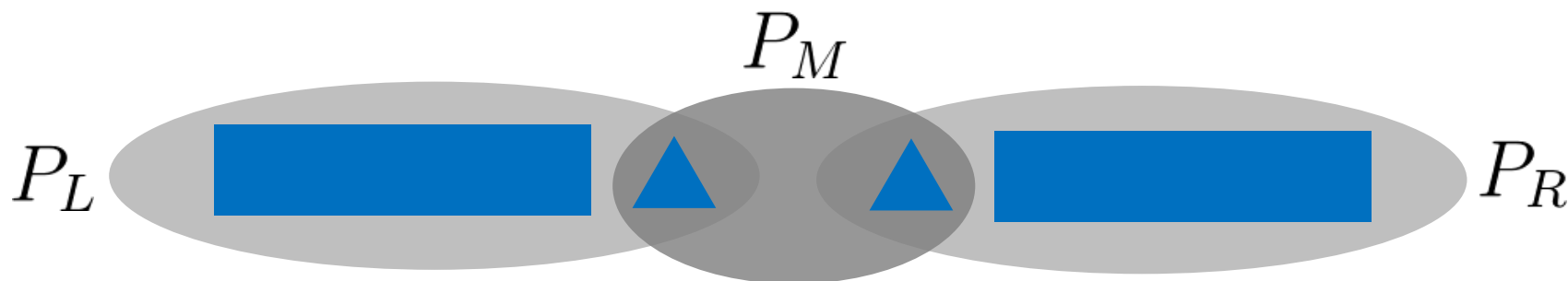
P_M



3 The 4-particle ($N \times 3 \times 3 \times N$) Hamiltonian

a unique, very entangled ground state
a gapped Hamiltonian, $O(1)$ terms, frust. free

I	A	B
A	AA	AB
B	BA	BB

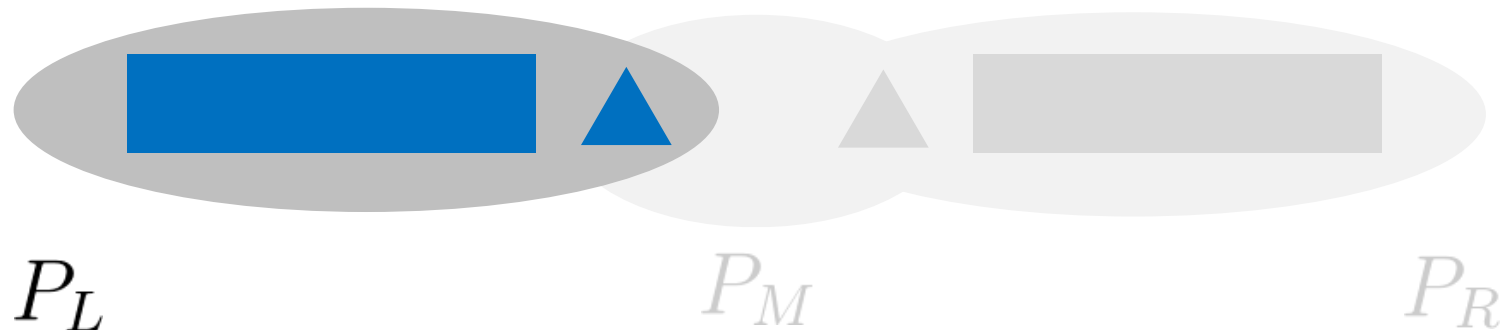
 $\frac{1}{3\sqrt{N}}$ 

3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits



prepare $\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$
from $\frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) |x\rangle$

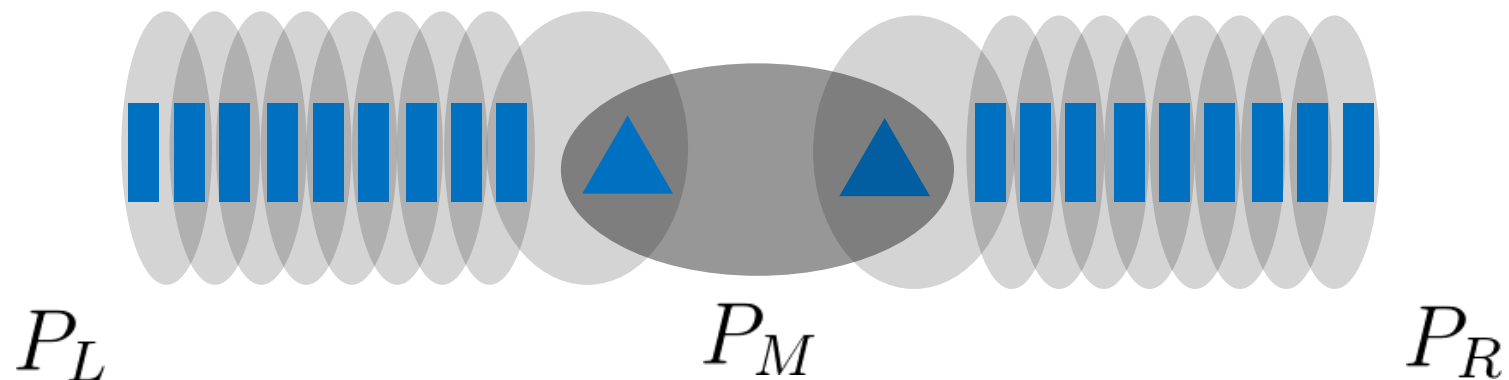


3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits

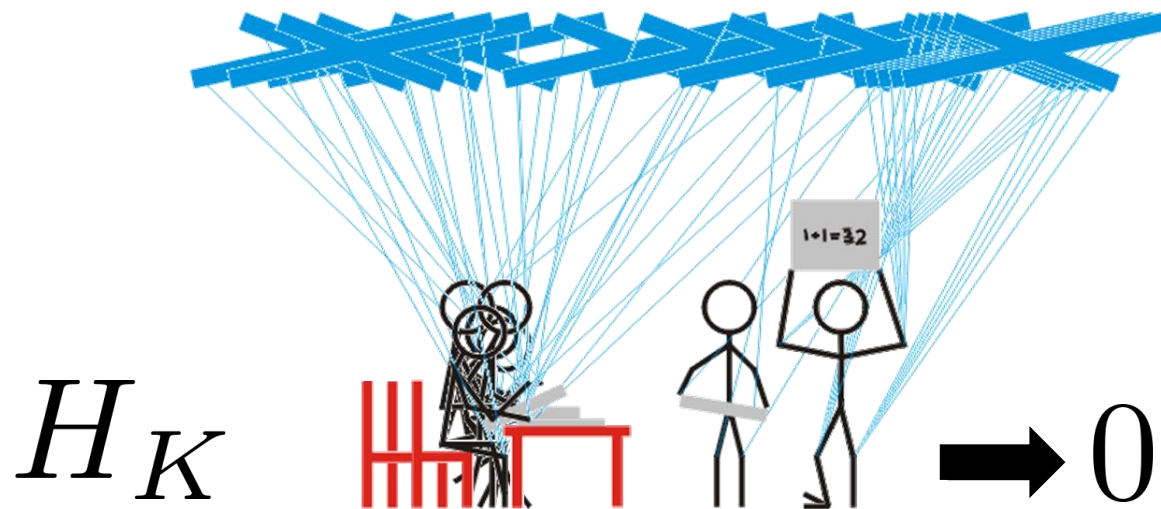


prepare $\frac{1}{\sqrt{3}} (|1\rangle|x\rangle + |2\rangle A|x\rangle + |3\rangle B|x\rangle)$
from $\frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) |x\rangle$



3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits



$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

$\underbrace{U_t \cdots U_1}_{\text{history state}} |\varphi_0\rangle$

3 The history state: a ground state of a qudit chain

2-local

c-o-n-d-i-t-i-o-n-s

d=13 [AGIK'08]

d=8 [HNN'13]

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$



3 The history state: a ground state of a qudit chain

2-local

c-o-n-d-i-t-i-o-n-s

clock encoding
state progression
initialization

$$|\cdots\rangle|0\rangle \otimes |0\rangle$$

$$|\varphi_t\rangle \otimes |t\rangle$$

$$|\varphi_{t+1}\rangle \otimes |t+1\rangle$$

$$|\psi_{hist}\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\varphi_t\rangle \otimes |t\rangle$$

idling

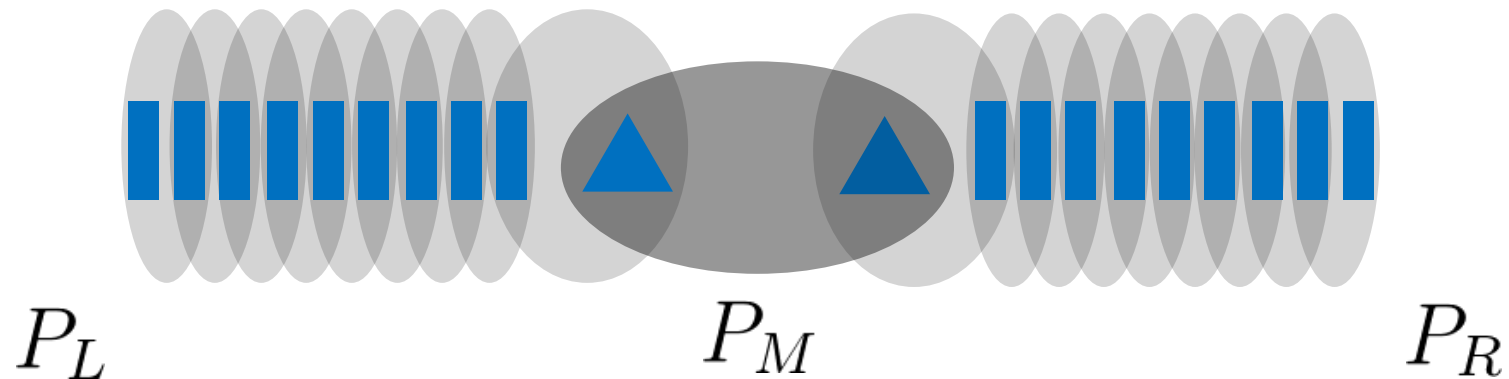
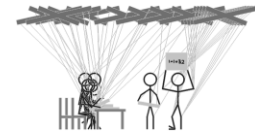


most of the state has the result

3 Making the counterexample local

- a quantum circuit, history state
Kitaev's LH, 1D, qudits

an approx. groundstate, a small $1/\text{poly}(n)$ gap



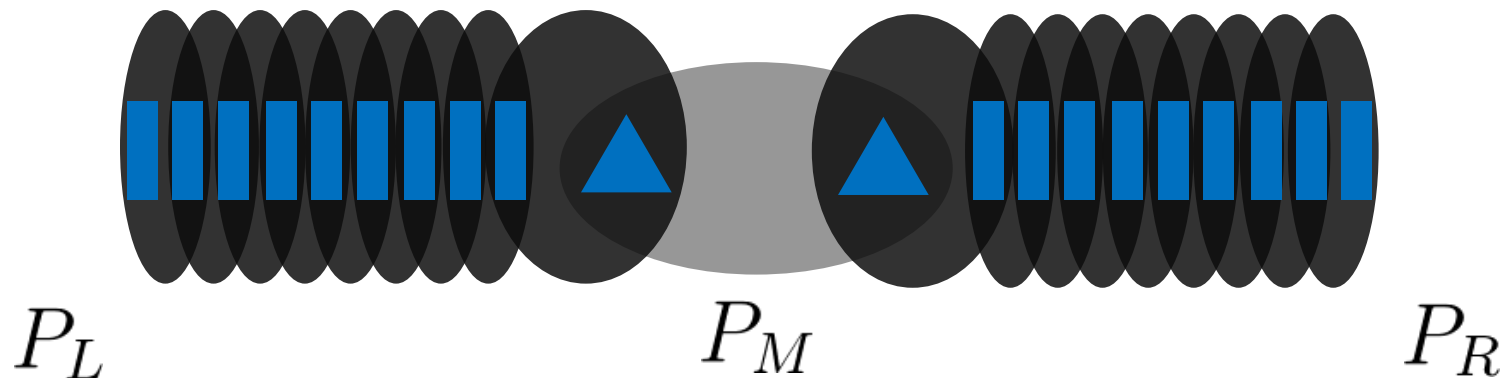
3 Making the counterexample local

- a quantum circuit, history state
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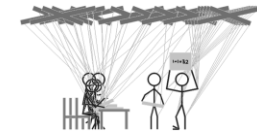
an approx. groundstate, a small $1/\text{poly}(n)$ gap

- rescale P_L, P_R (not the middle!)
a constant gap, huge couplings



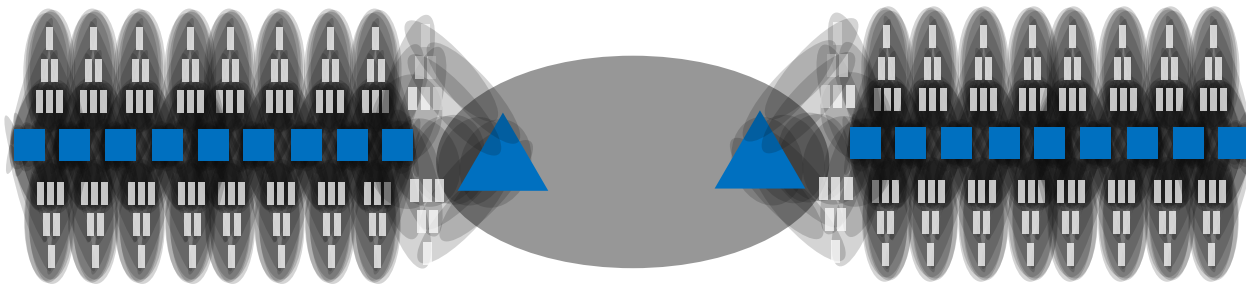
3 Making the counterexample local

- a quantum circuit, history state
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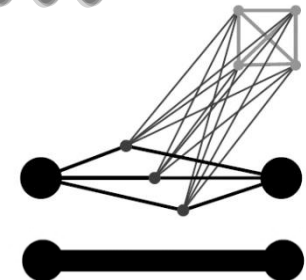


an approx. groundstate, a small $1/\text{poly}(n)$ gap

- rescale P_L, P_R (not the middle!)
a constant gap, huge couplings



- decompose using gadgets [Cao, N.]
~~huge couplings~~ many spins, high degree

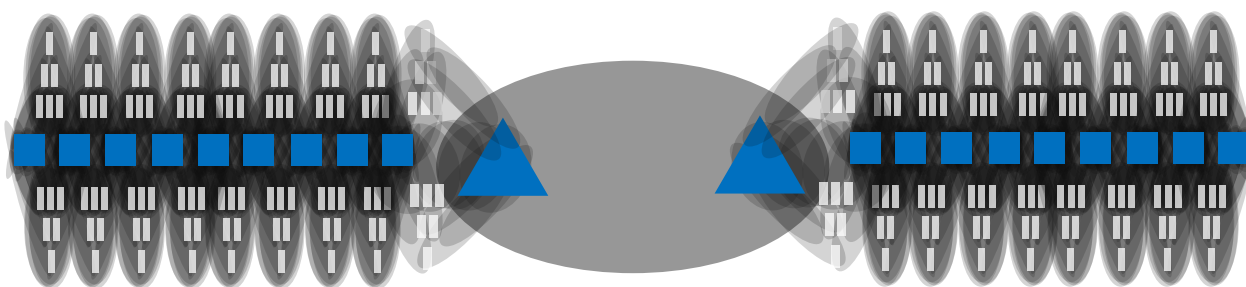


3 A local Hamiltonian, $(N \times M) \times 3 \times 3 \times (N \times M)$

- a unique and very entangled ground state

$$\approx |w\rangle \otimes$$

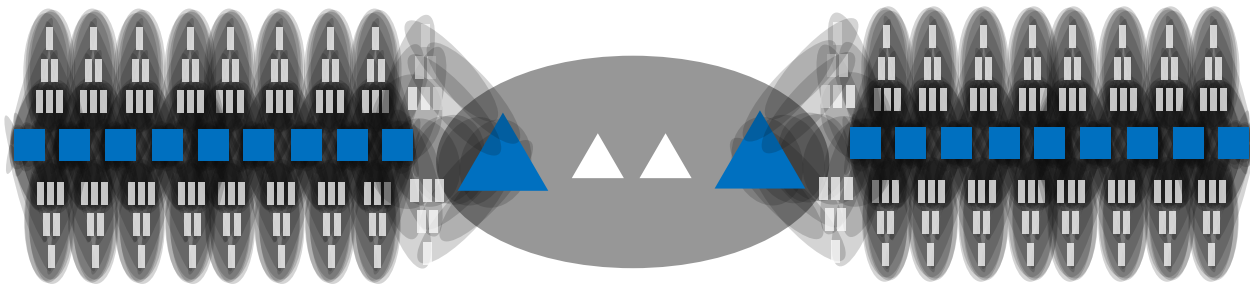
<i>I</i>	<i>A</i>	<i>B</i>
<i>A</i>	<i>AA</i>	<i>AB</i>
<i>B</i>	<i>BA</i>	<i>BB</i>



- a constant gap
 $O(1)$ or smaller norm terms, frustration

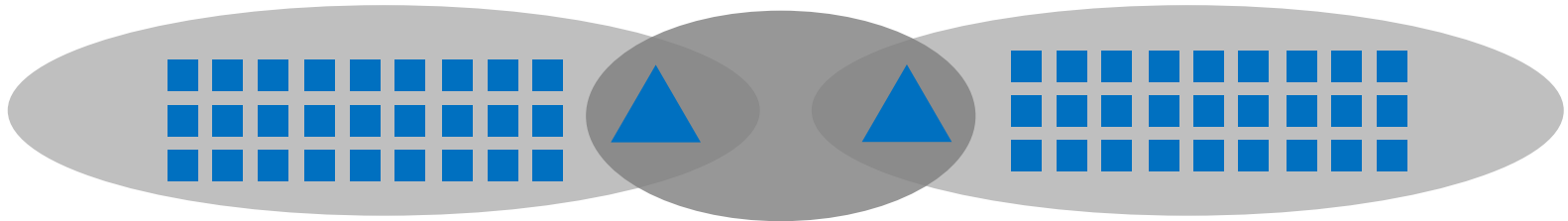
3 What next?

- a longer middle chain?



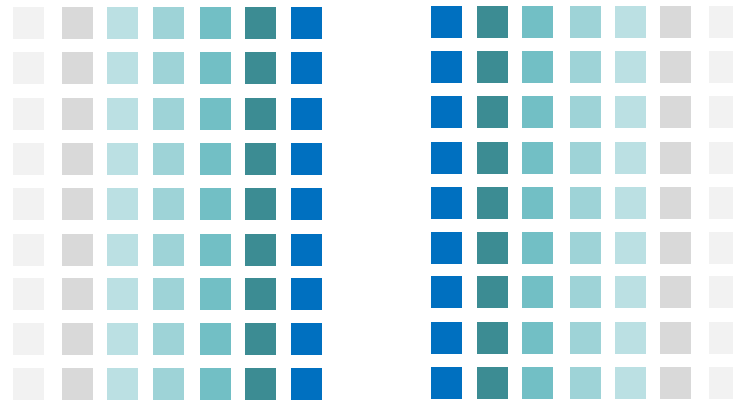
3 What next?

- a longer middle chain?
- a nice lattice on the sides?



3 What next?

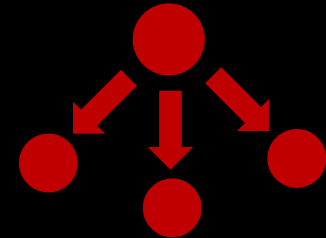
- a longer middle chain?
- a nice lattice on the sides?
- a new area conjecture: count the cut & things nearby?



1

q. expanders

maximally entangled states



2

entanglement

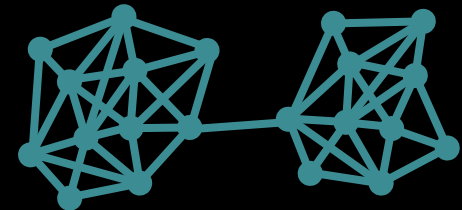
testing and communication



3

area law

gaps, connections, correlations



local tests of global entanglement and a counterexample to the generalized area law

