

Most quantum states are useless for  
measurement-based quantum  
computation

Steve Flammia

Perimeter Institute

QIP 2009, Santa Fe

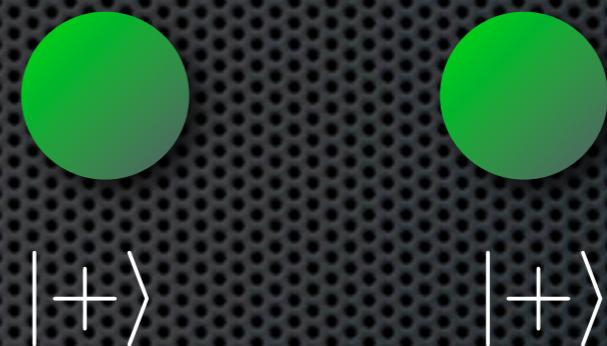
D. Gross, SF, J. Eisert 0810.4331

M. Bremner, C. Mora, A. Winter 0812.3001

# Measurement-based QC

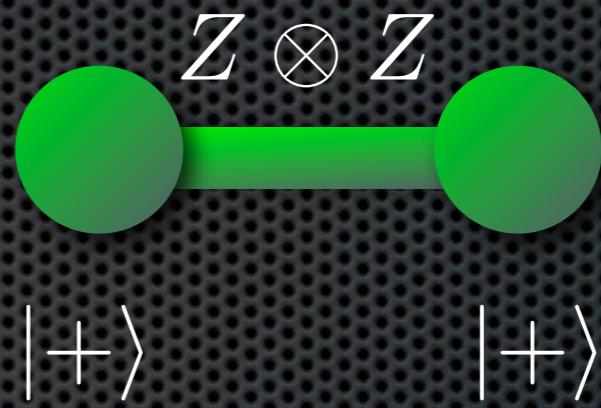
# Measurement-based QC

- prepare X eigenstates



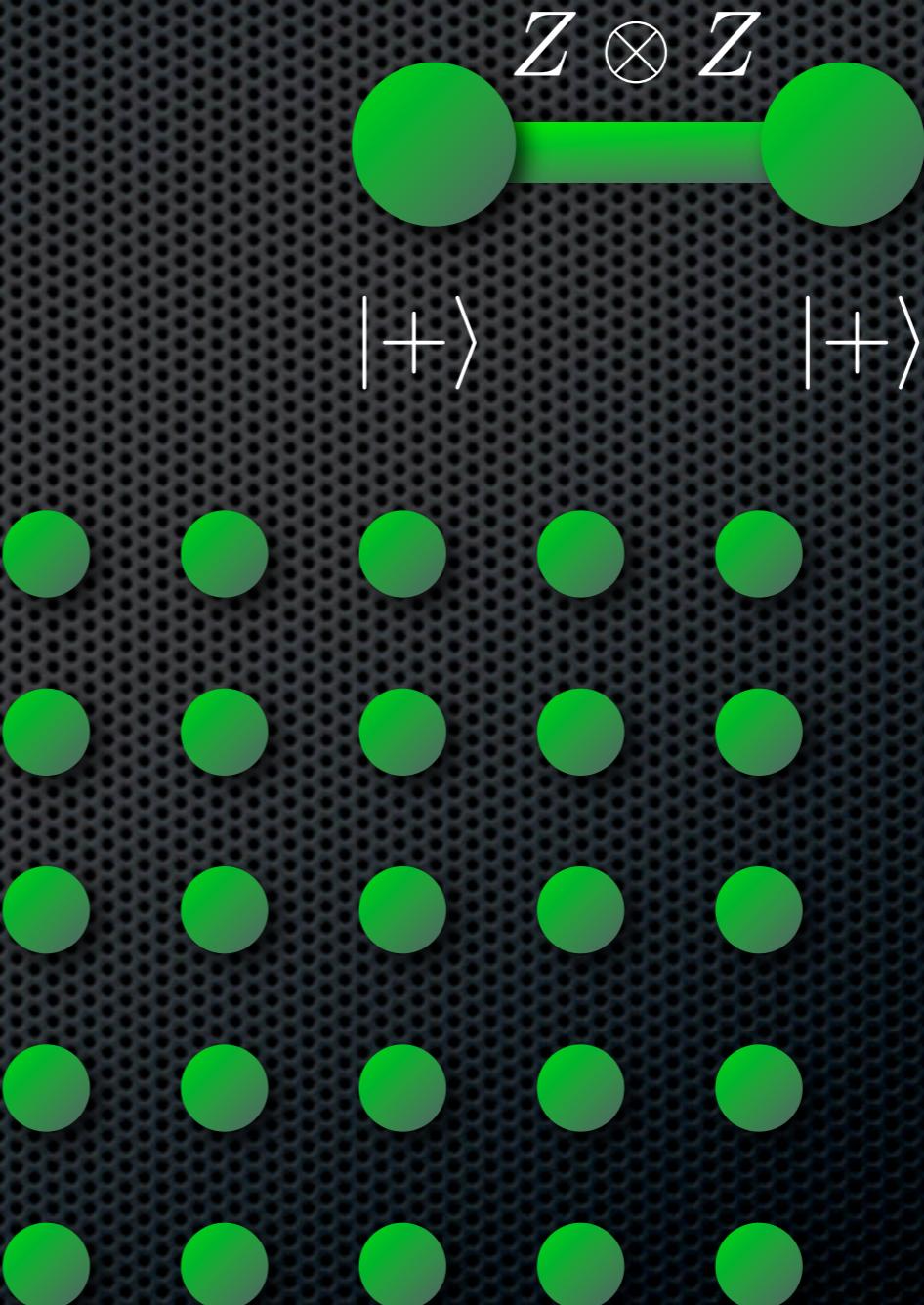
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling



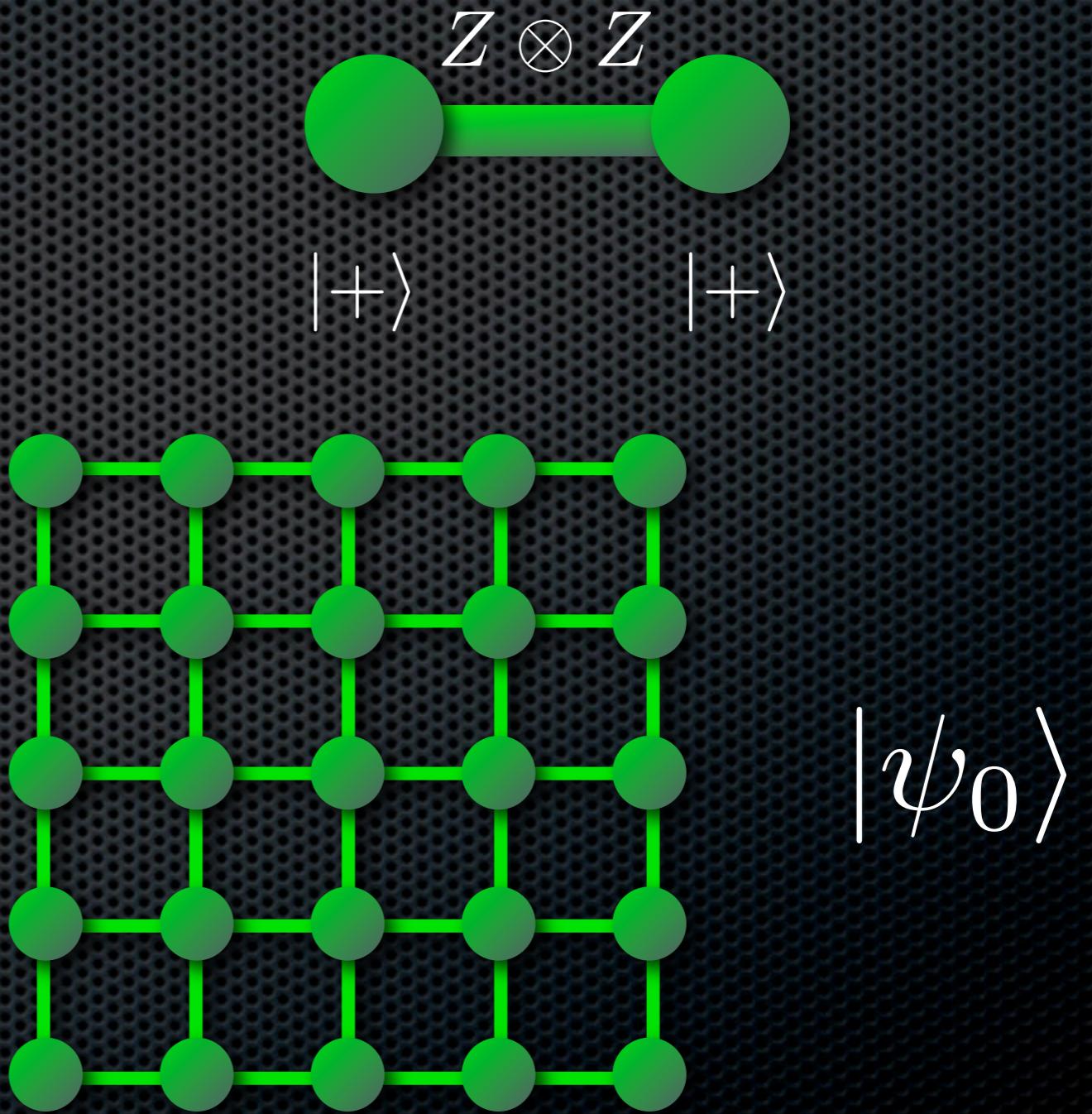
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE



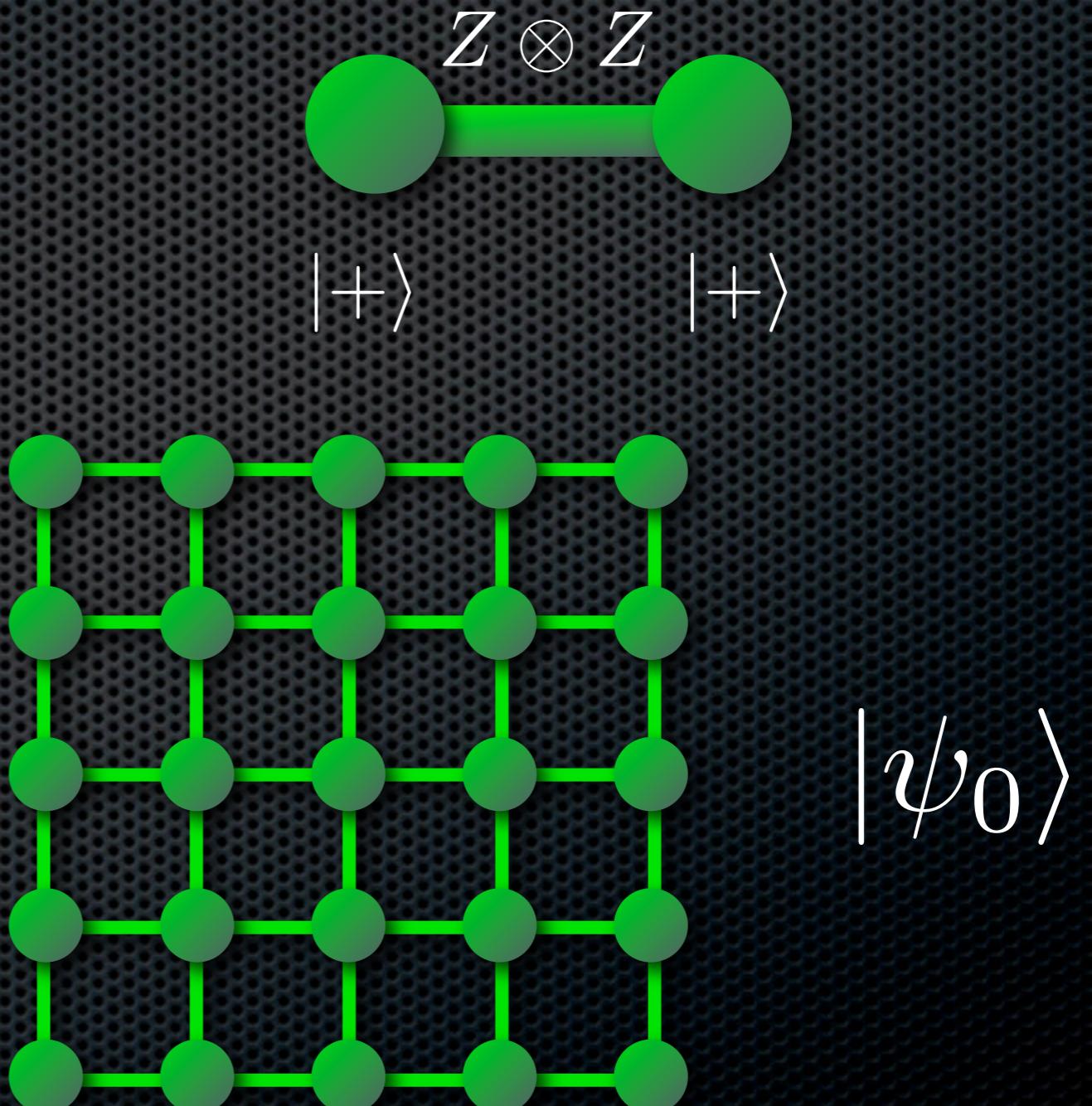
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE



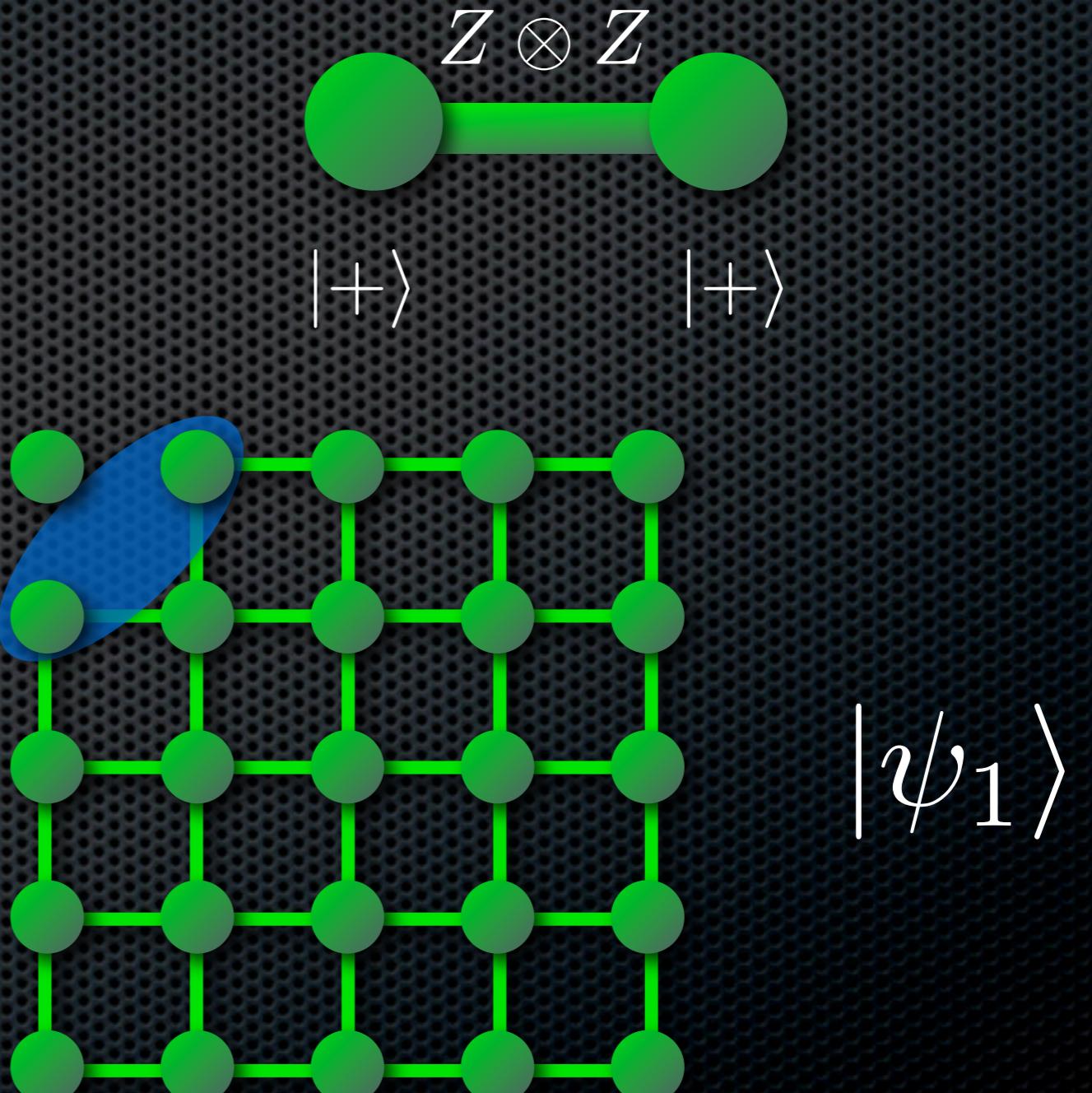
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE
- arbitrary single-qubit measurements with feedforward to compute



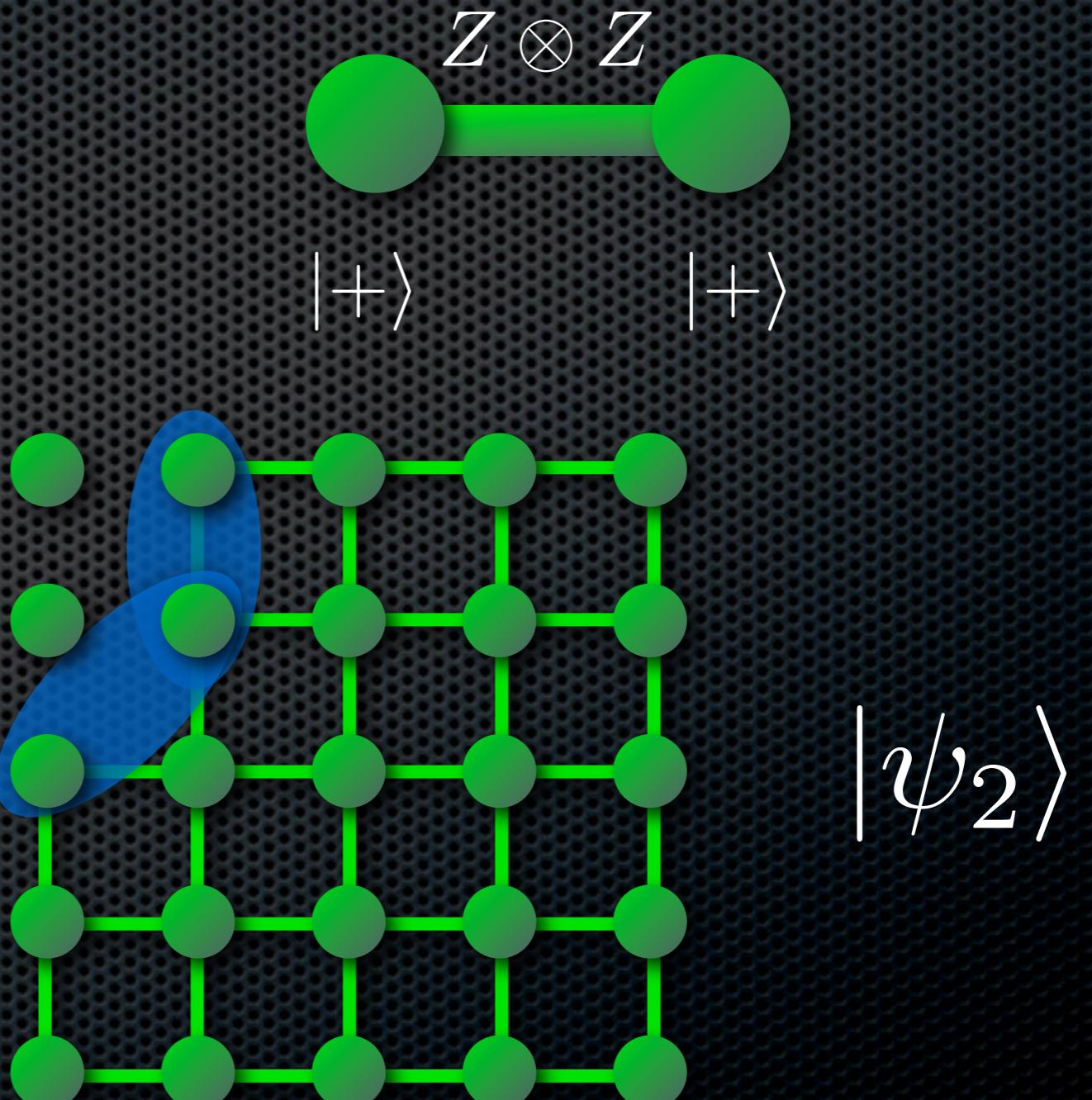
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE
- arbitrary single-qubit measurements with feedforward to compute



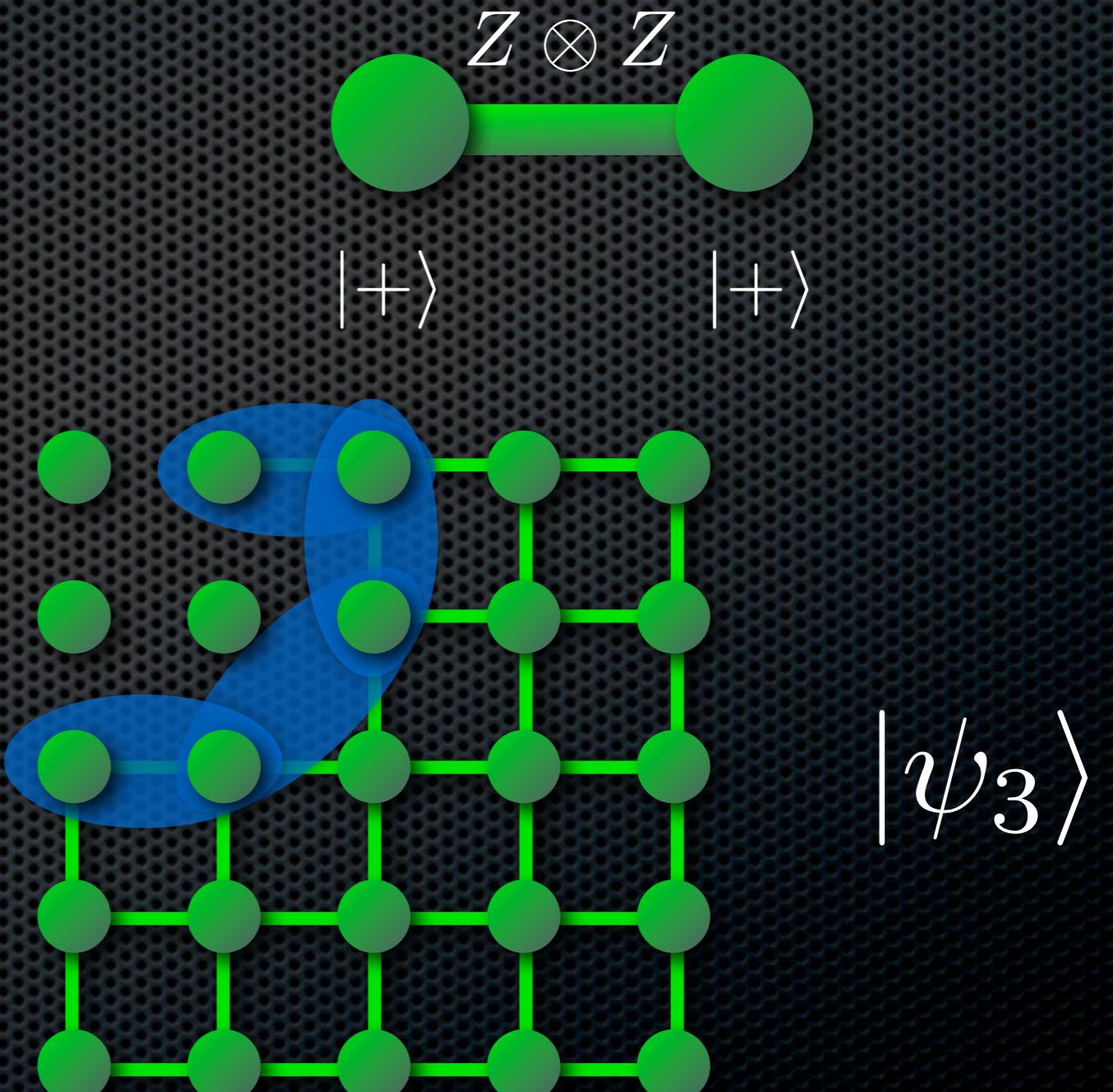
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE
- arbitrary single-qubit measurements with feedforward to compute



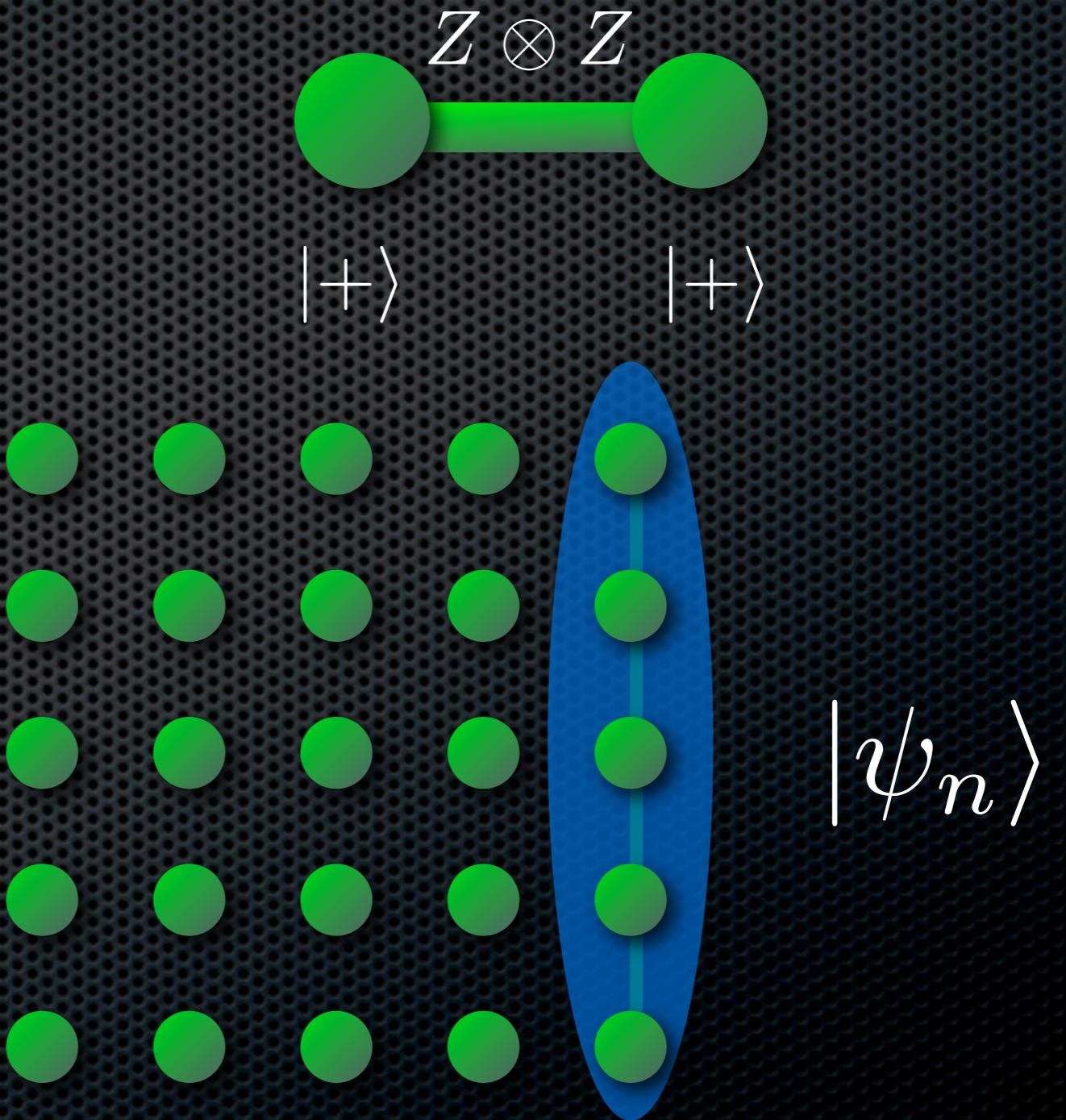
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE
- arbitrary single-qubit measurements with feedforward to compute



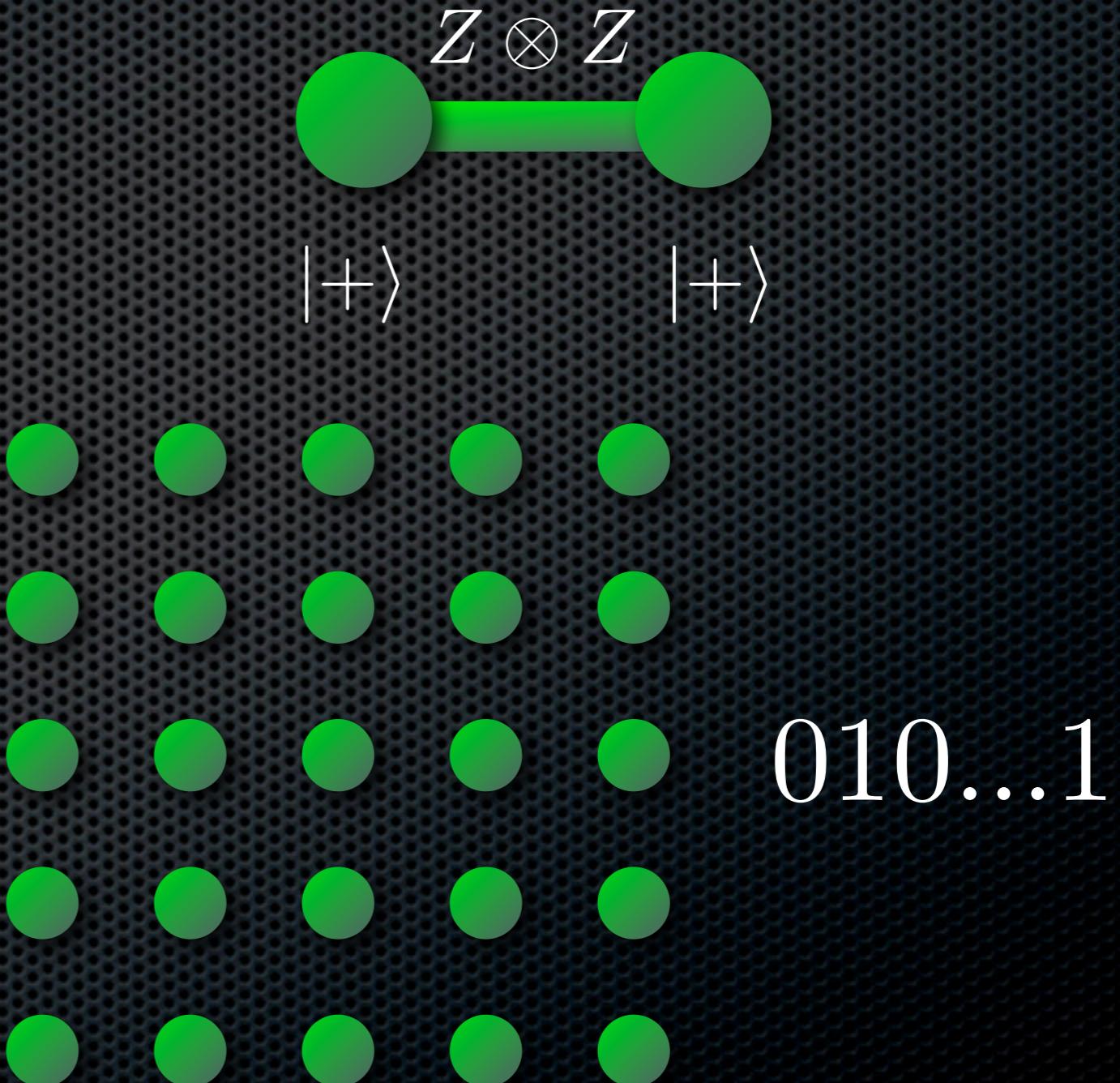
# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE
- arbitrary single-qubit measurements with feedforward to compute



# Measurement-based QC

- prepare X eigenstates
- entangle neighbors with a Z-Z coupling
- Build a large lattice for universality:  
the CLUSTER STATE
- arbitrary single-qubit measurements with feedforward to compute





- In general, MBQC requires:

- In general, MBQC requires:

- A family of  $n$  qubit quantum states

- In general, MBQC requires:

- A family of  $n$  qubit quantum states



- In general, MBQC requires:

- A family of n qubit quantum states

$|\Psi\rangle$



- In general, MBQC requires:
  - A family of n qubit quantum states
  - A classical control computer determines where to measure, the measurement basis and how to interpret the measurement outcomes

$|\Psi\rangle$



- In general, MBQC requires:
  - A family of n qubit quantum states
  - A classical control computer determines where to measure, the measurement basis and how to interpret the measurement outcomes

$|\Psi\rangle$



- In general, MBQC requires:
  - A family of n qubit quantum states
  - A classical control computer determines where to measure, the measurement basis and how to interpret the measurement outcomes

1  
 $|\Psi\rangle$



- In general, MBQC requires:
  - A family of n qubit quantum states
  - A classical control computer determines where to measure, the measurement basis and how to interpret the measurement outcomes

1  
 $|\Psi'\rangle$



- In general, MBQC requires:
  - A family of n qubit quantum states
  - A classical control computer determines where to measure, the measurement basis and how to interpret the measurement outcomes

$$|\Psi'\rangle$$


- In general, MBQC requires:
  - A family of  $n$  qubit quantum states
  - A classical control computer determines where to measure, the measurement basis and how to interpret the measurement outcomes

$$|\Psi'\rangle$$


- Without initial entanglement, it's clear you can't do better than BPP.

# Universality and entanglement

Question:

What are the necessary and sufficient conditions for a family of  $n$  qubit quantum states to be universal for MBQC?



# Universality and entanglement

Question:

What are the necessary and sufficient conditions for a family of  $n$  qubit quantum states to be universal for MBQC?



Necessary conditions:

van den Nest, Miyake, Dür, Briegel 2006  
find entanglement measures  
that must grow “quickly” with  $n$ .

# Universality and entanglement

Question:

What are the necessary and sufficient conditions for a family of  $n$  qubit quantum states to be universal for MBQC?



Sufficient conditions:

Gross, Eisert, Schuch, Pérez-García 2007  
find states with special structure  
in the many-body correlations.

Brennen & Miyake 2008, Doherty & Bartlett 2008  
find ground states with special structure.

# Bridging the divide

Quantum world

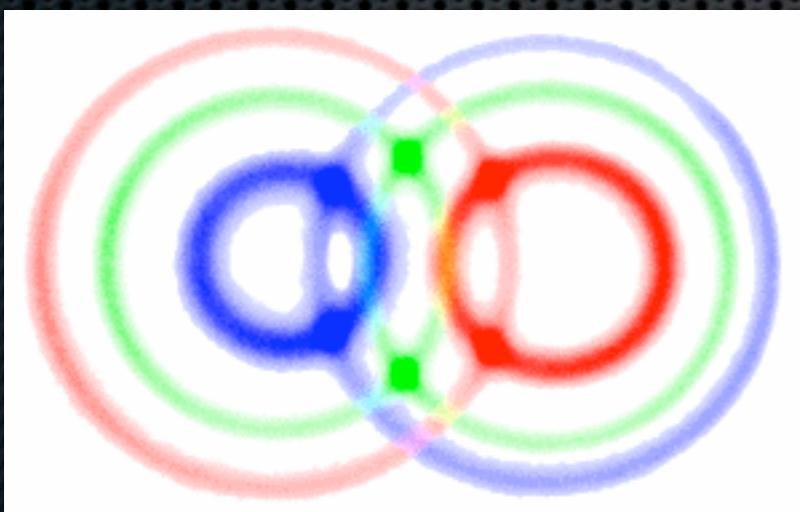
Classical world



# Bridging the divide

Quantum world

Entanglement  
and  
correlations

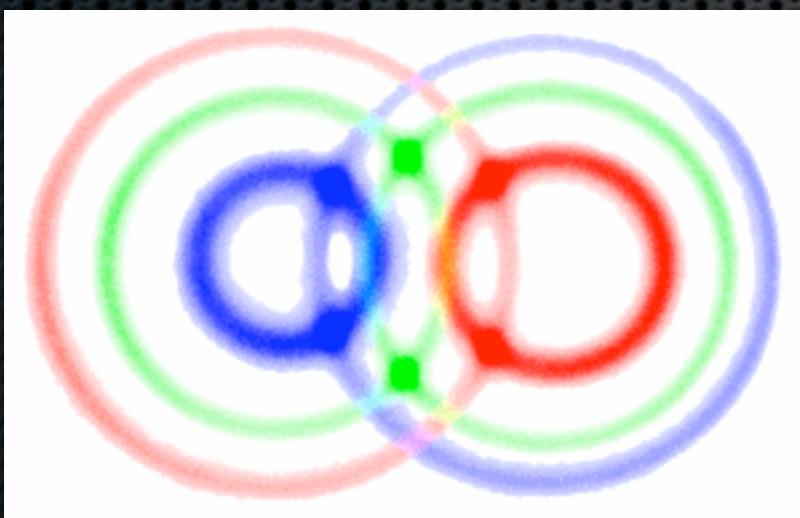


Classical world

# Bridging the divide

Quantum world

Entanglement  
and  
correlations



Classical world

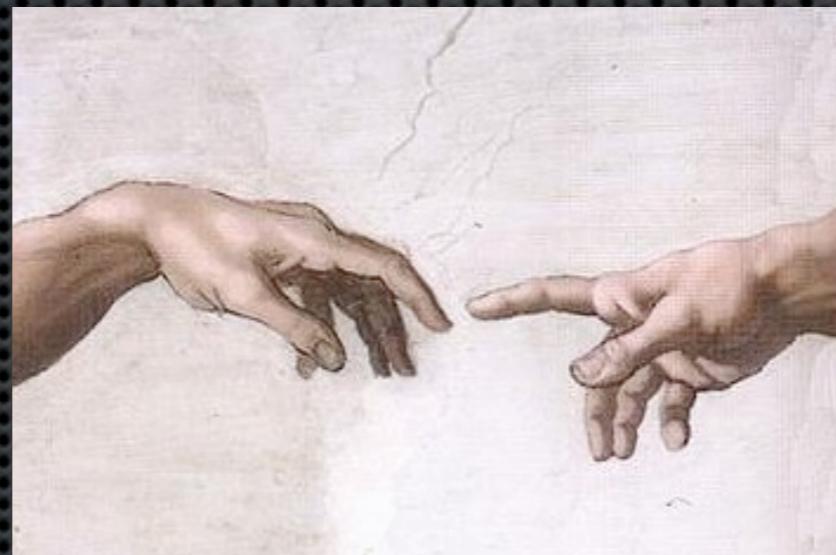
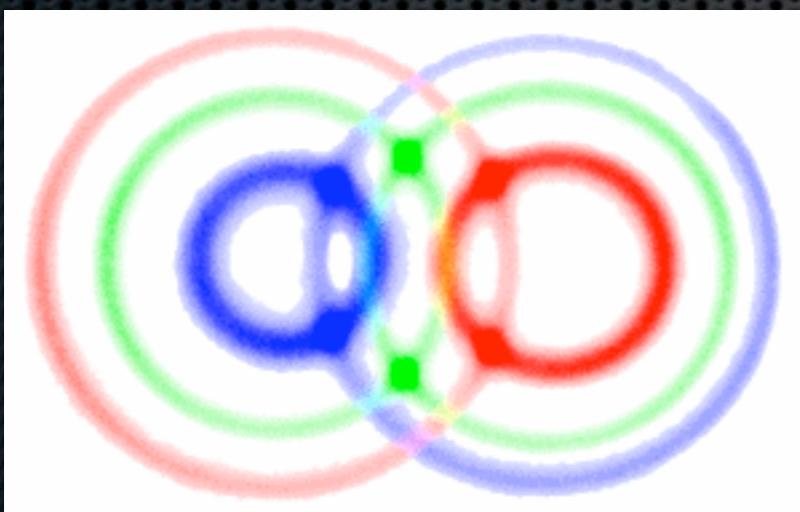
Local bases,  
Limited processing  
power.



# Bridging the divide

Quantum world

Entanglement  
and  
correlations



MBQC

Classical world

Local bases,  
Limited processing  
power.



# Local bases, geometric measure

$$E_g(\Psi) = -\log_2 \sup_{\alpha \in \mathcal{P}} |\langle \alpha | \Psi \rangle|^2$$



the set of product states

Answers the question:  
How far is the nearest  
collection of local bases  
 $\alpha_1, \alpha_2, \dots, \alpha_n$ ?

Large geometric measure  
Far from all product states

# Local bases, geometric measure

$$E_g(\Psi) = -\log_2 \sup_{\alpha \in \mathcal{P}} |\langle \alpha | \Psi \rangle|^2$$



the set of product states

Answers the question:  
How far is the nearest  
collection of local bases  
 $\alpha_1, \alpha_2, \dots, \alpha_n$ ?

Large geometric measure



Far from all product states

Theorem 1 (GFE): n qubit states with  
 $E_g > n - O(\log n)$   
are useless for MBQC.

# Local bases, geometric measure

For concreteness, a state is **useless** if it fails to provide a polynomial-time MBQC algorithm for Factoring.

Theorem 1 (GFE): n qubit states with  
 $E_g > n - O(\log n)$   
are useless for MBQC.

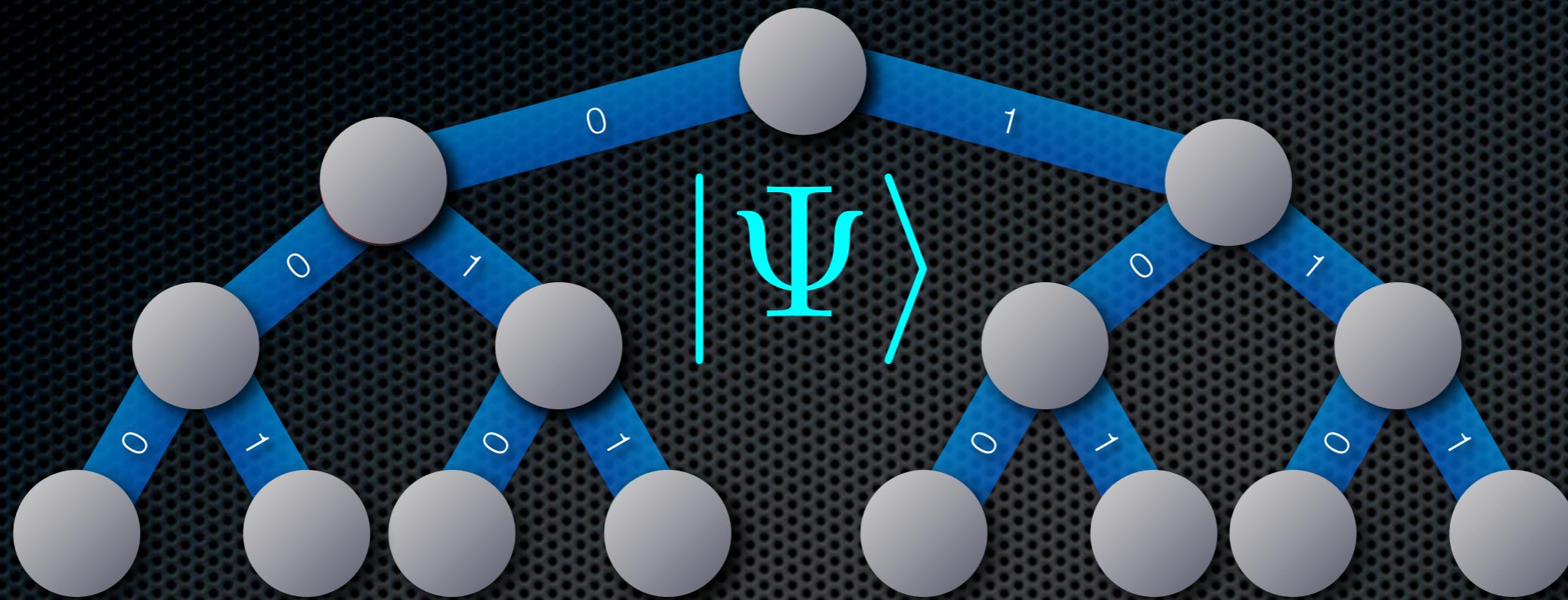
# Local bases, geometric measure

For concreteness, a state is **useless** if it fails to provide a polynomial-time MBQC algorithm for Factoring.

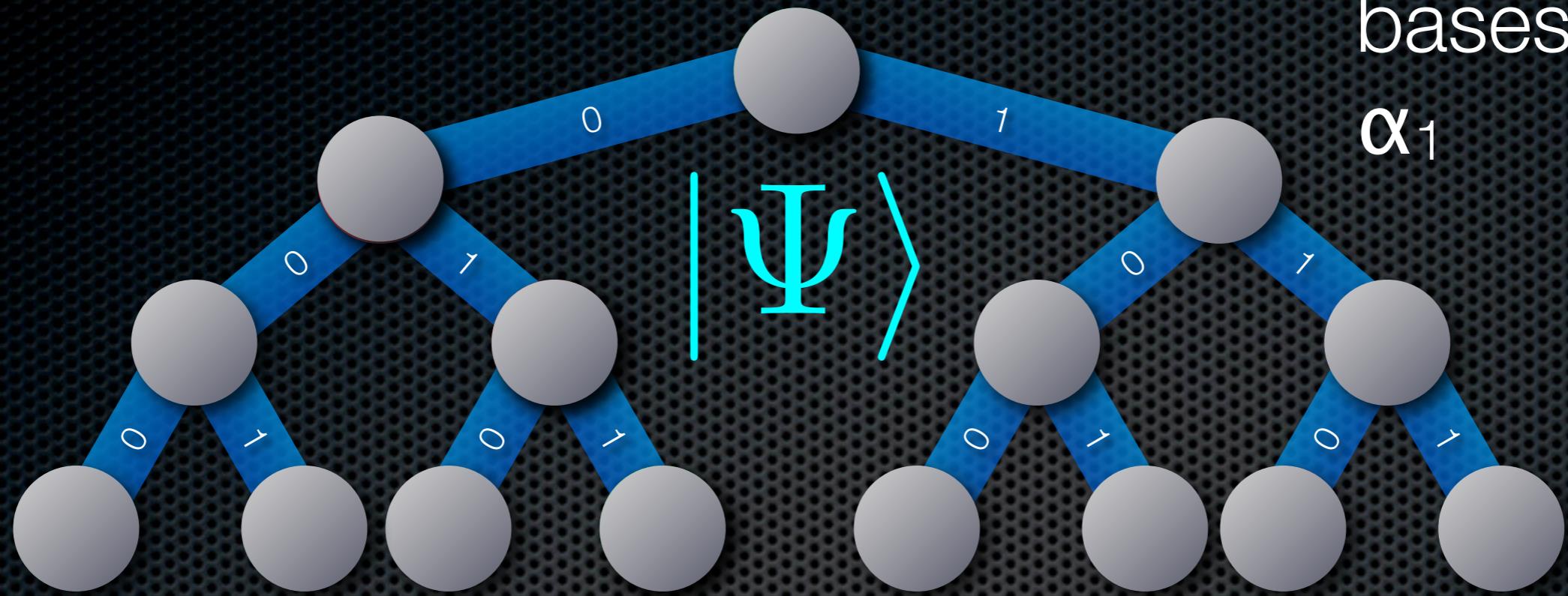
Proof strategy: replace  $\Psi$  with a classical coin and show there exists a classical algorithm that factors just as well (within poly factors).



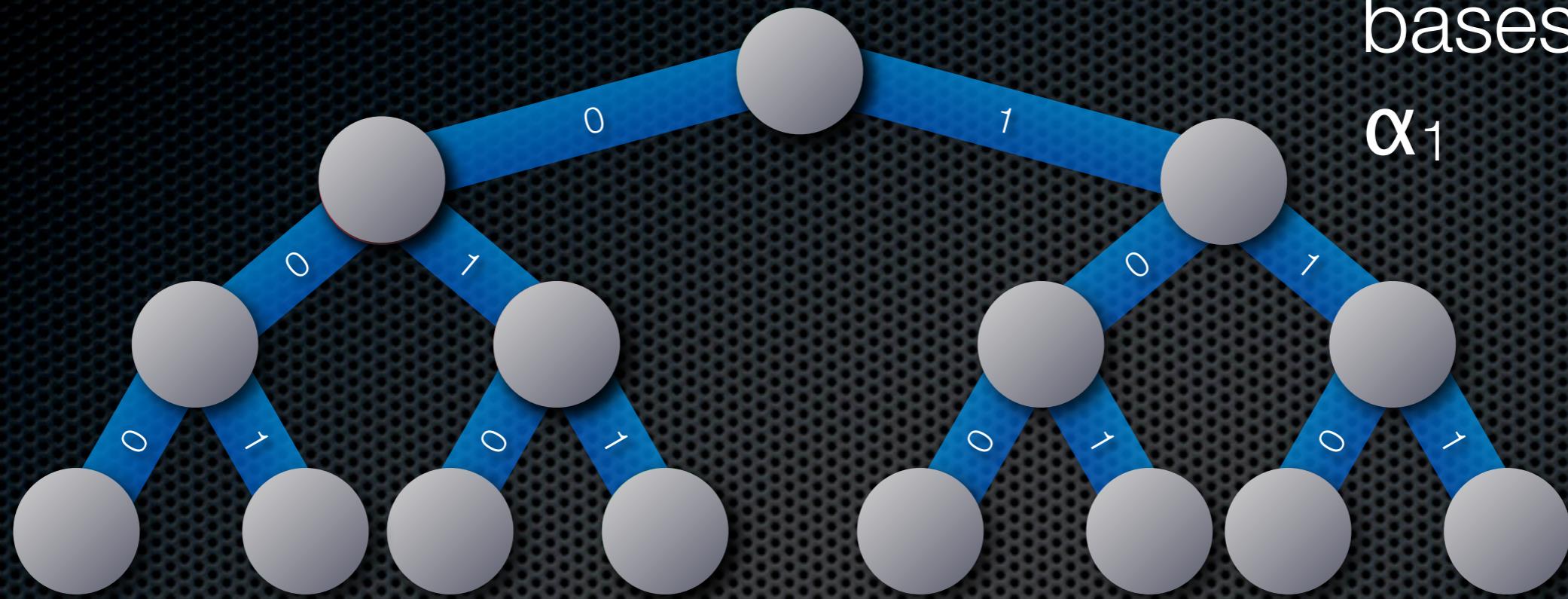
Theorem 1 (GFE): n qubit states with  
 $E_g > n - O(\log n)$   
are useless for MBQC.



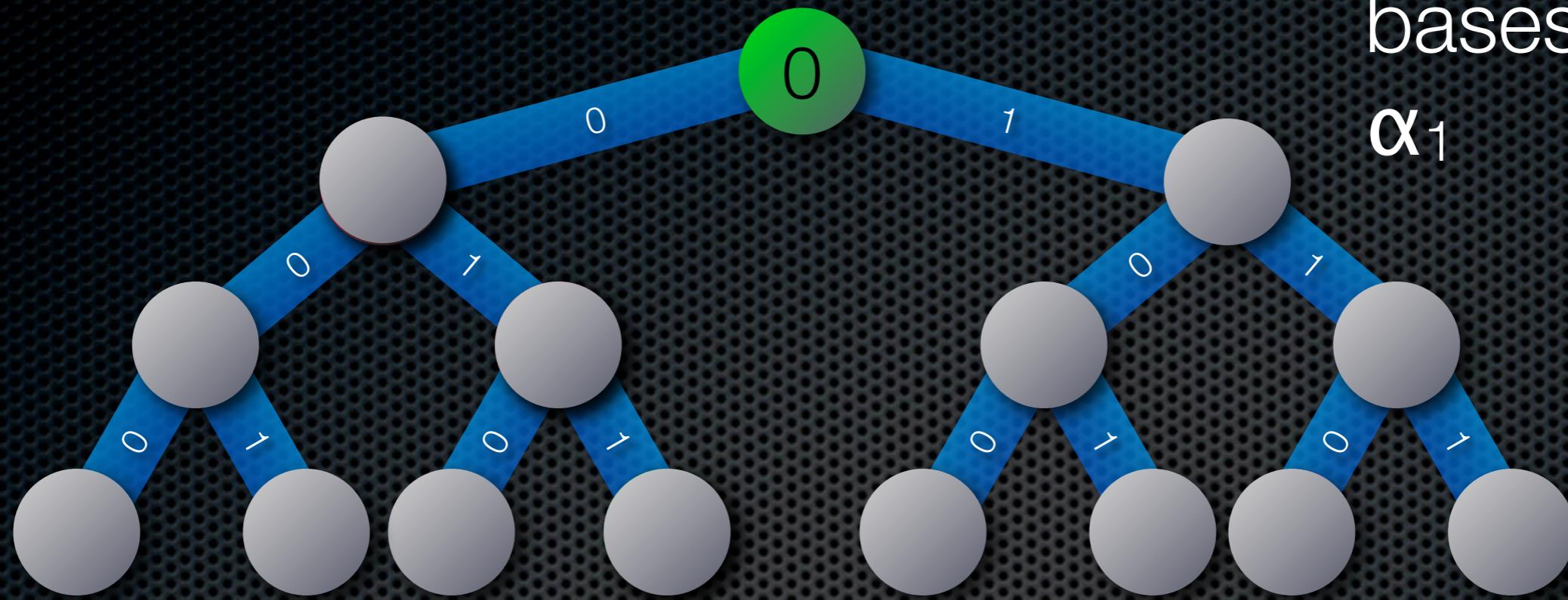
bases:  
 $\alpha_1$



bases:  
 $\alpha_1$

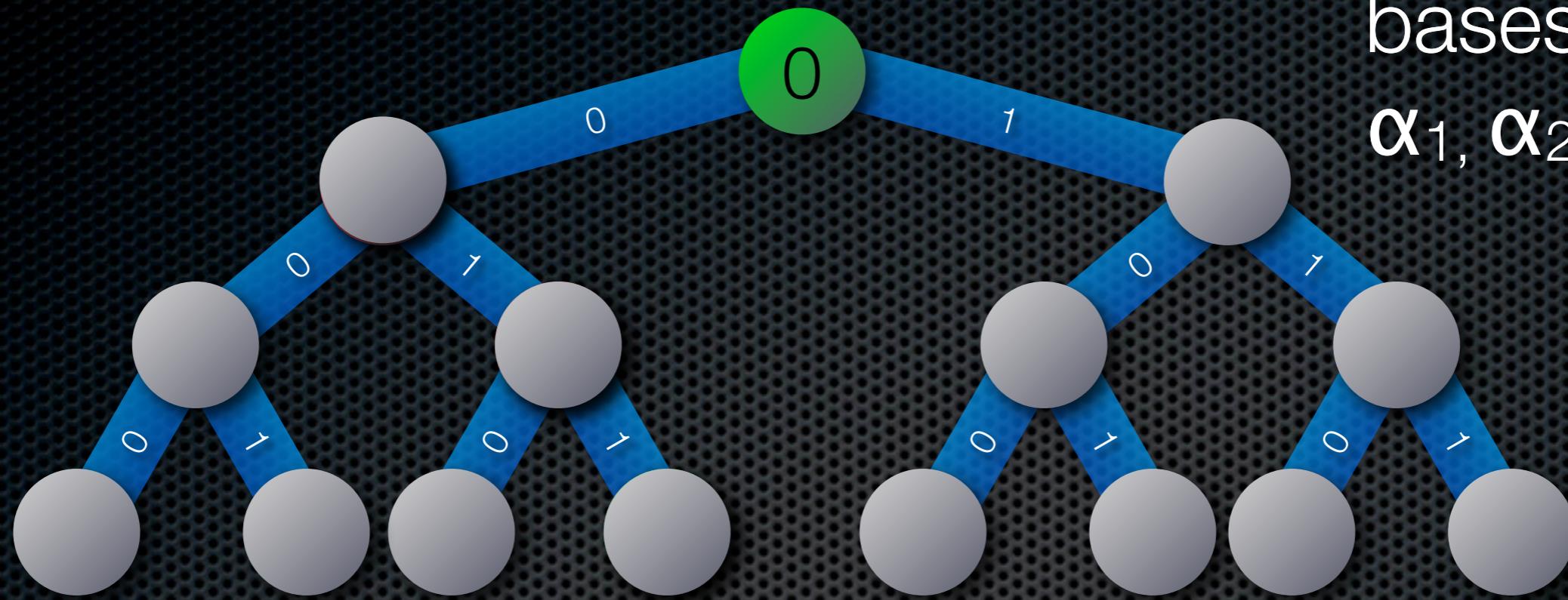


bases:  
 $\alpha_1$



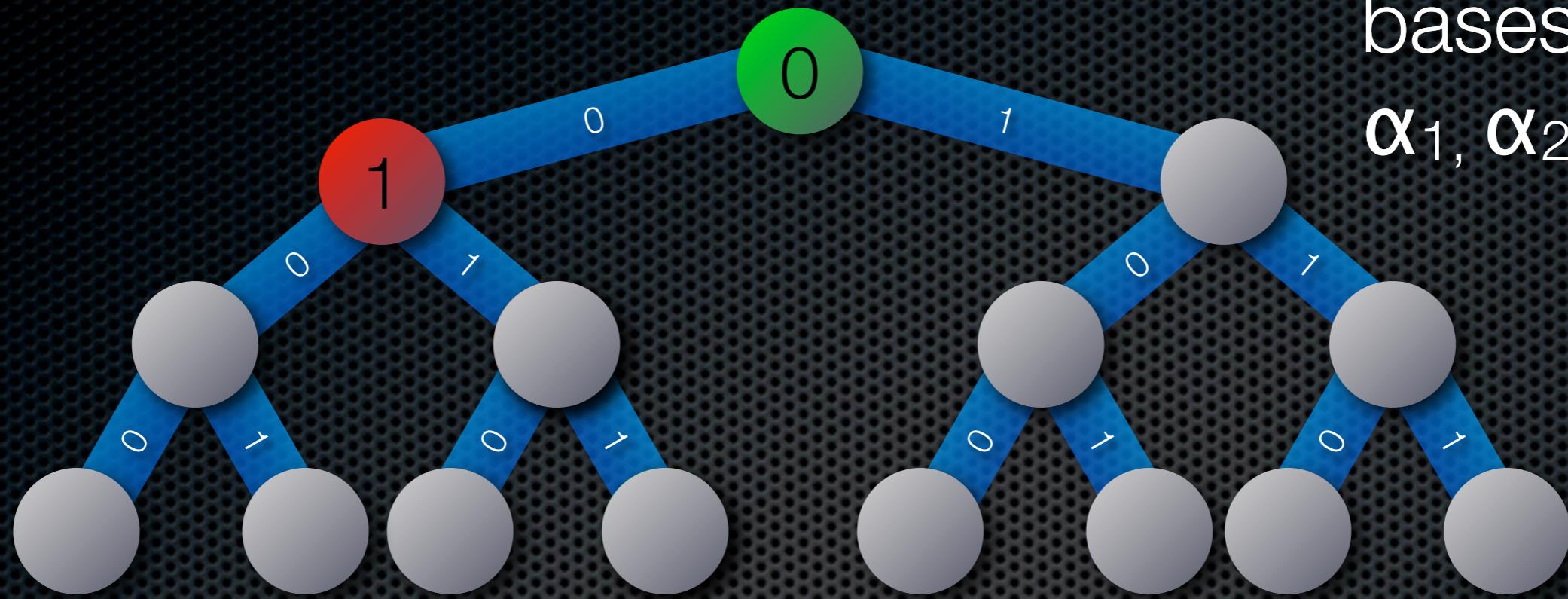
0

bases:  
 $\alpha_1, \alpha_2$



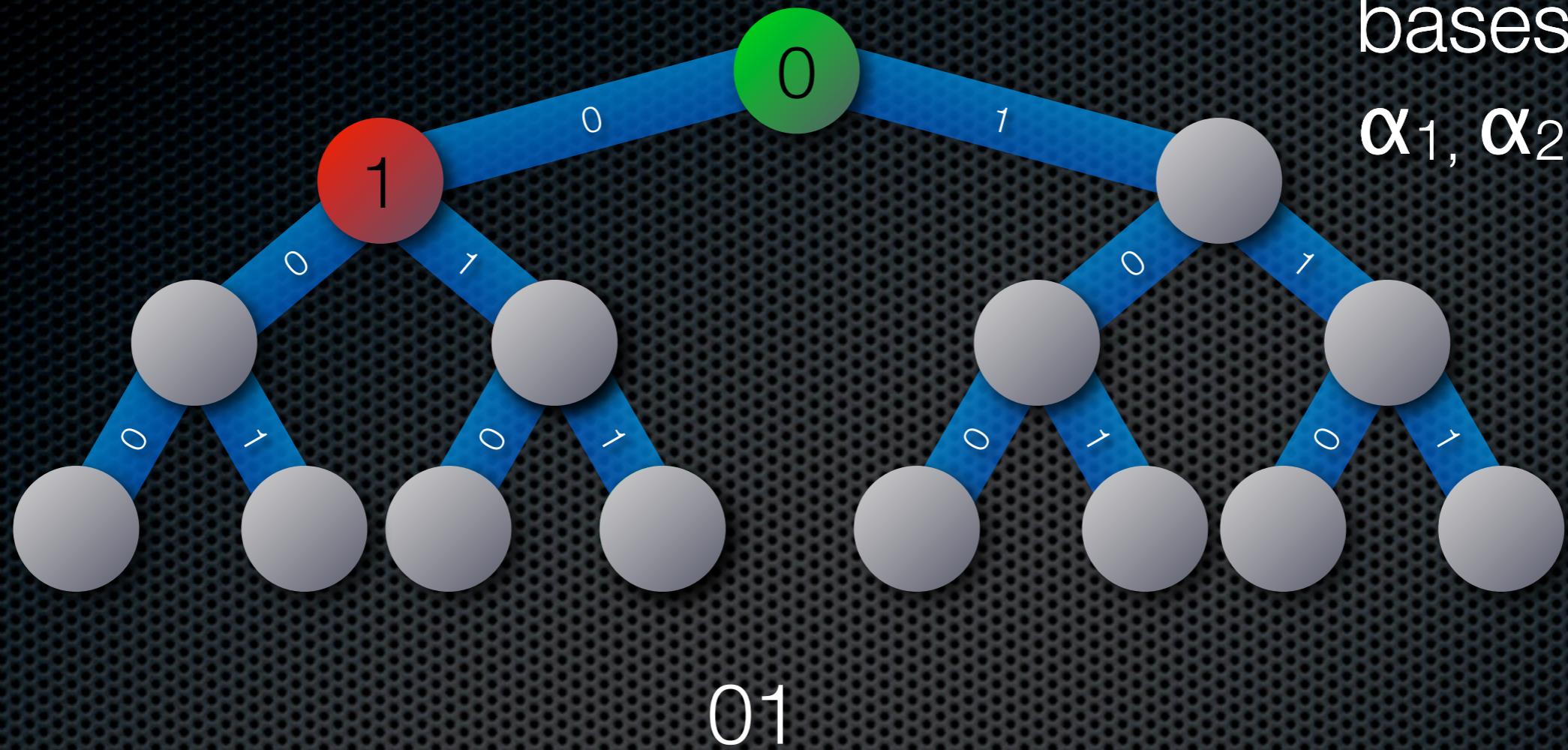
0

bases:  
 $\alpha_1, \alpha_2$

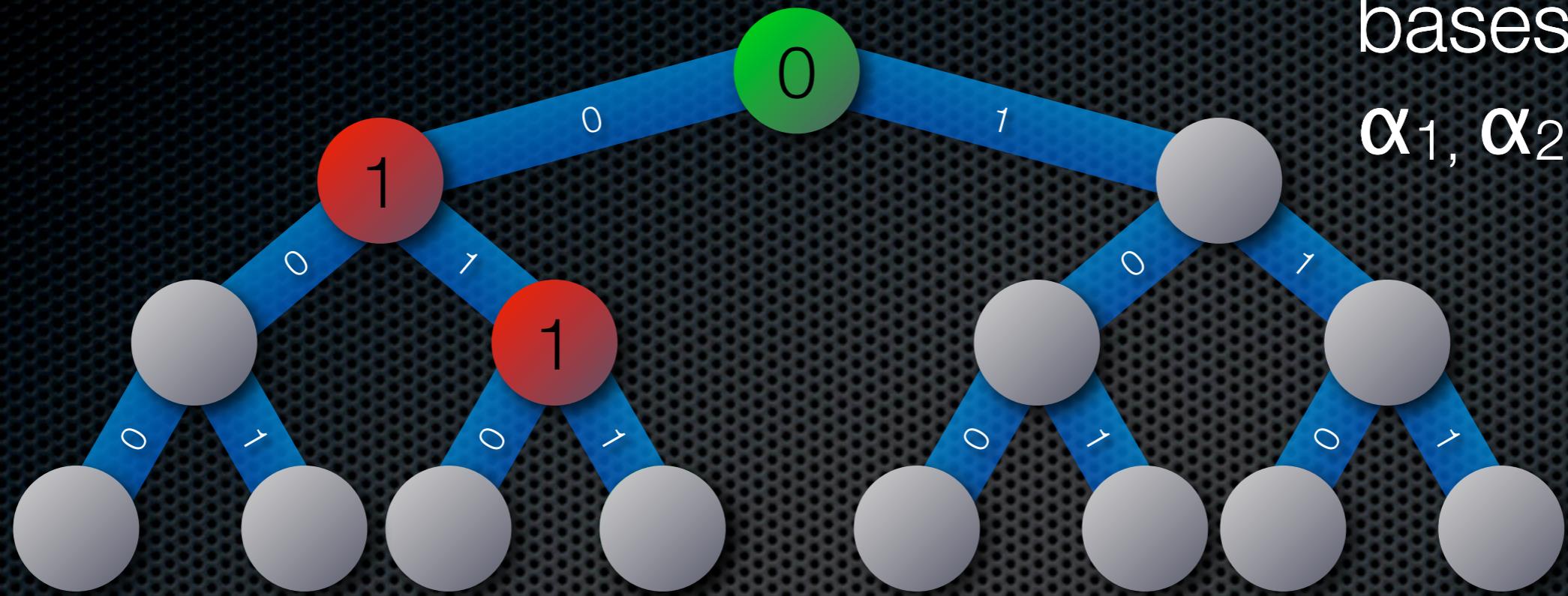


01

bases:  
 $\alpha_1, \alpha_2, \alpha_3$

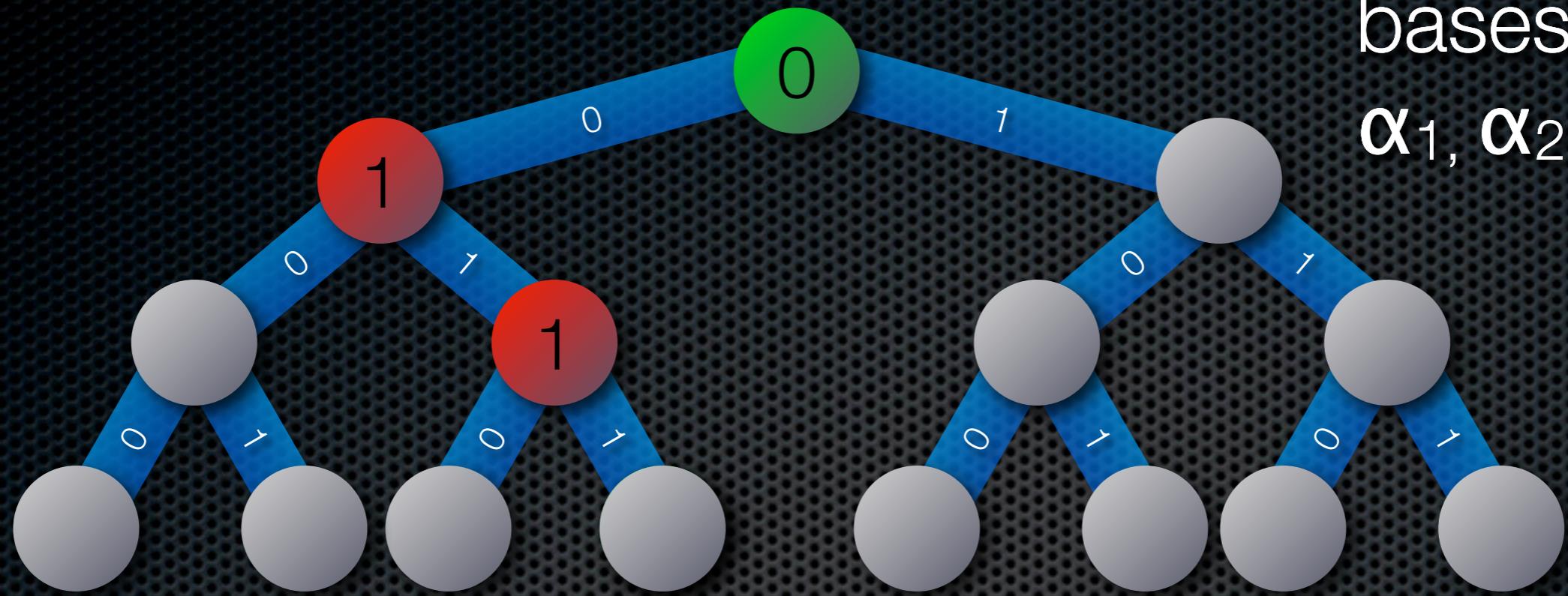


bases:  
 $\alpha_1, \alpha_2, \alpha_3$



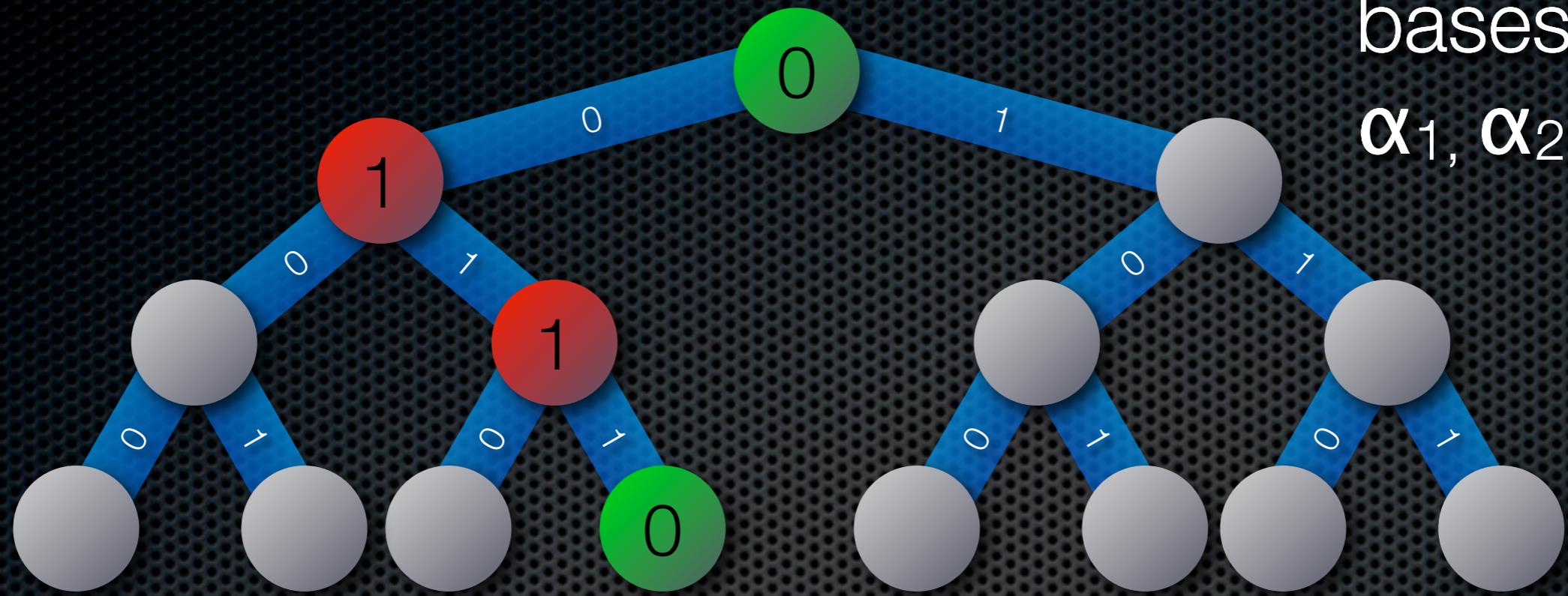
011

bases:  
 $\alpha_1, \alpha_2, \alpha_3, \alpha_4$



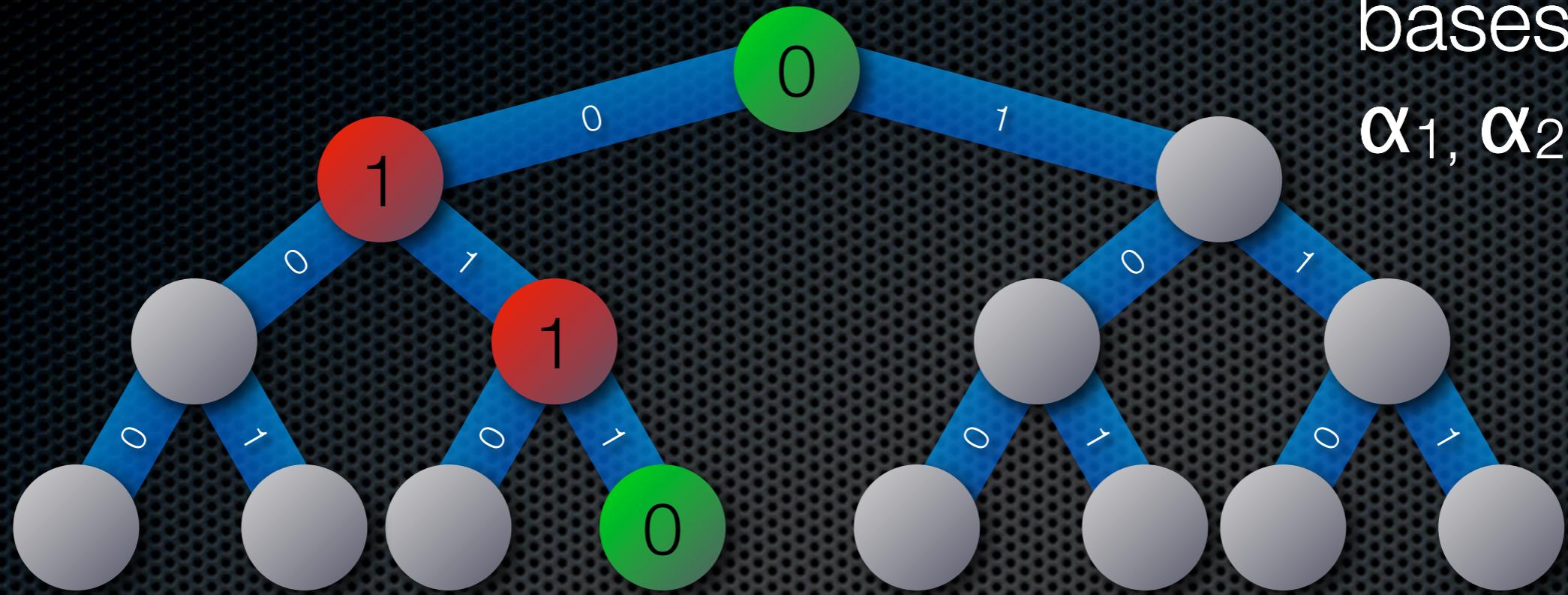
011

bases:  
 $\alpha_1, \alpha_2, \alpha_3, \alpha_4$



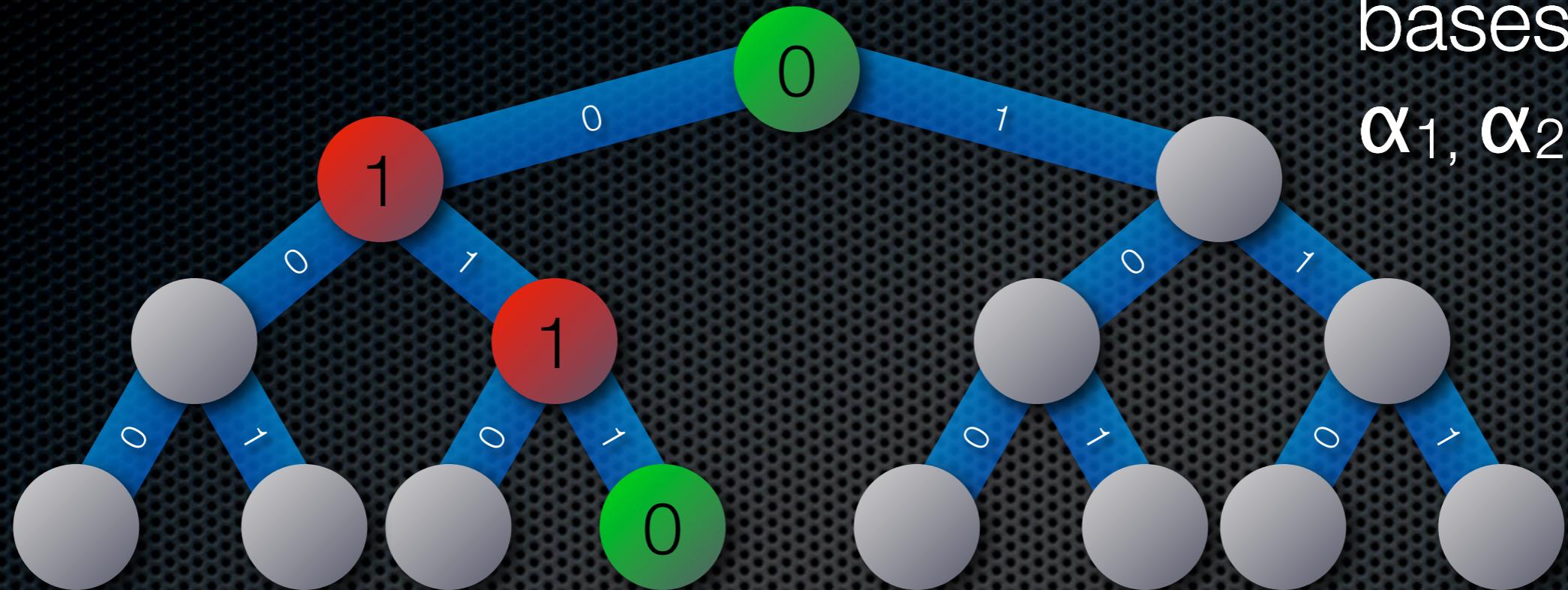
0110

bases:  
 $\alpha_1, \alpha_2, \alpha_3, \alpha_4$



The “good” outcomes G cause the classical control computer to output a valid factorization. We want this to succeed with constant probability, say  $p>.5$

bases:  
 $\alpha_1, \alpha_2, \alpha_3, \alpha_4$



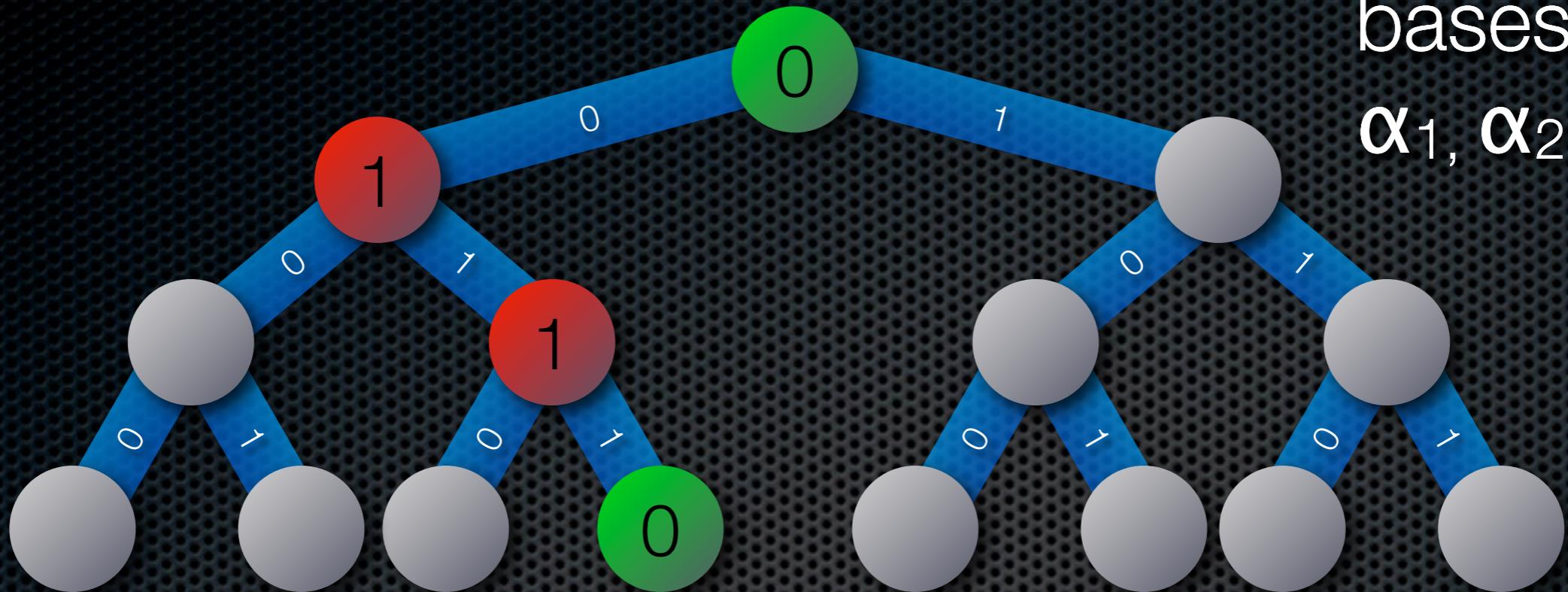
The “good” outcomes  $G$  cause the classical control computer to output a valid factorization. We want this to succeed with constant probability, say  $p > .5$

Suppose

$$E_g > n - \delta,$$

$$\delta = O(\log n)$$

bases:  
 $\alpha_1, \alpha_2, \alpha_3, \alpha_4$



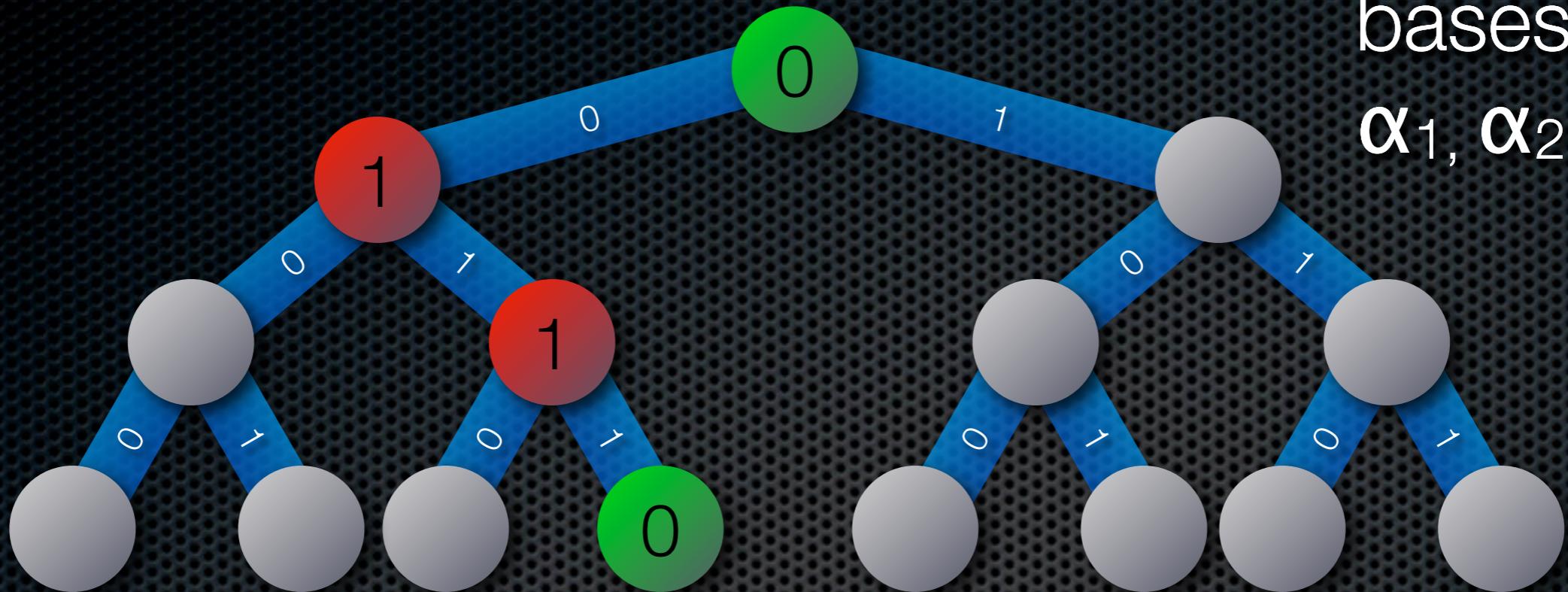
The “good” outcomes  $G$  cause the classical control computer to output a valid factorization. We want this to succeed with constant probability, say  $p > .5$

Suppose  $E_g > n - \delta$ ,  
 $\delta = O(\log n)$

$$|\langle \alpha | \Psi \rangle|^2 \leq 2^{-E_g} \leq 2^{-n+\delta}$$

$$\Rightarrow \frac{|G|}{2^n} > 2^{-\delta-1} = \text{poly}(1/n).$$

bases:  
 $\alpha_1, \alpha_2, \alpha_3, \alpha_4$



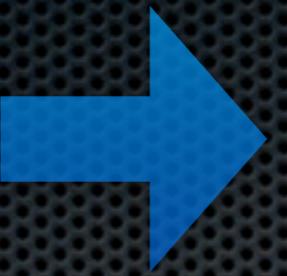
The “good” outcomes  $G$  cause the classical control computer to output a valid factorization. We want this to succeed with constant probability, say  $p > .5$

Suppose  $| \langle \alpha | \Psi \rangle |^2 \leq 2^{-E_g} \leq 2^{-n+\delta}$   
 $E_g > n - \delta$ ,  
 $\delta = O(\log n)$   $\Rightarrow \frac{|G|}{2^n} > 2^{-\delta-1} = \text{poly}(1/n).$

To simulate classically, just ignore the measurement results and use a classical coin!

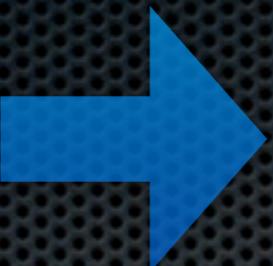


Large geometric measure



Useless for MBQC

Large geometric measure

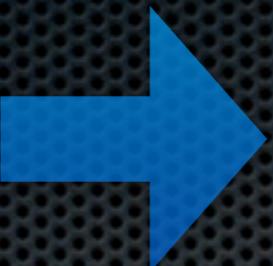


Useless for MBQC

This is vacuous unless such states exist.



Large geometric measure



Useless for MBQC

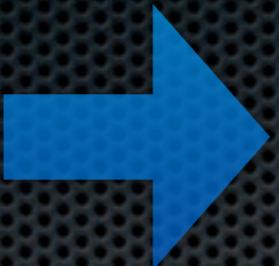
This is vacuous unless such states exist.

In fact, they are abundant.



Theorem 2 (GFE): The fraction of  $n$  qubit states with  
 $E_g < n - O(\log n)$   
is less than  $\exp(-n^2)$ .

Large geometric measure



Useless for MBQC

This is vacuous unless such states exist.

In fact, they are abundant.



Theorem 2 (GFE): The fraction of  $n$  qubit states with  
 $E_g < n - O(\log n)$   
is less than  $\exp(-n^2)$ .

The proof involves standard measure concentration arguments (via  $\epsilon$ -nets) and known results about random states

Random states  
are extravagant.



Random states  
are extravagant.



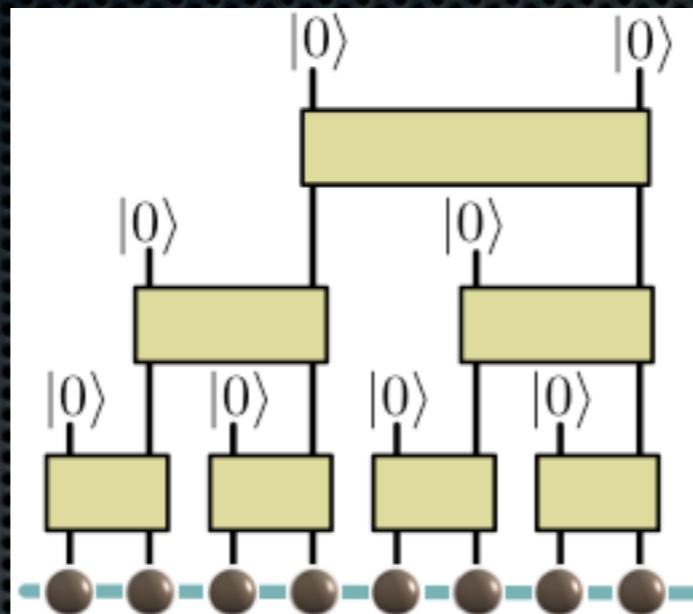
Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?



Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?

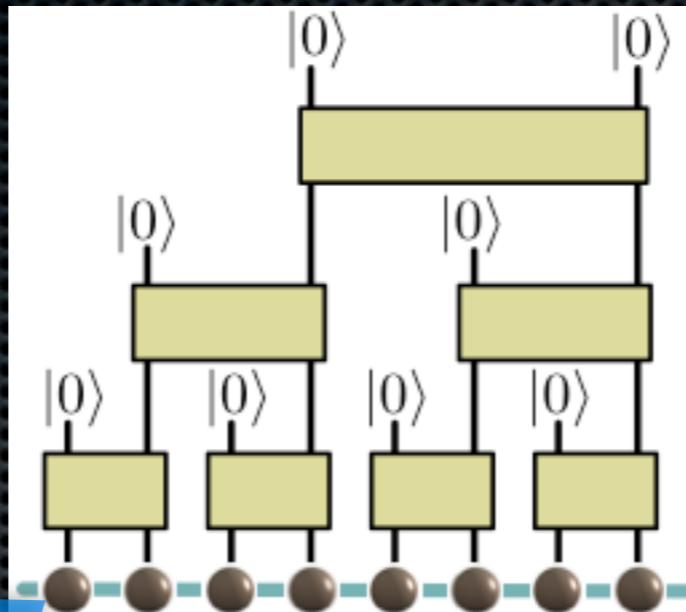


We can get to  
 $E_g > n \cdot o(n)$   
using a TTN  
construction.

Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?

d-level systems

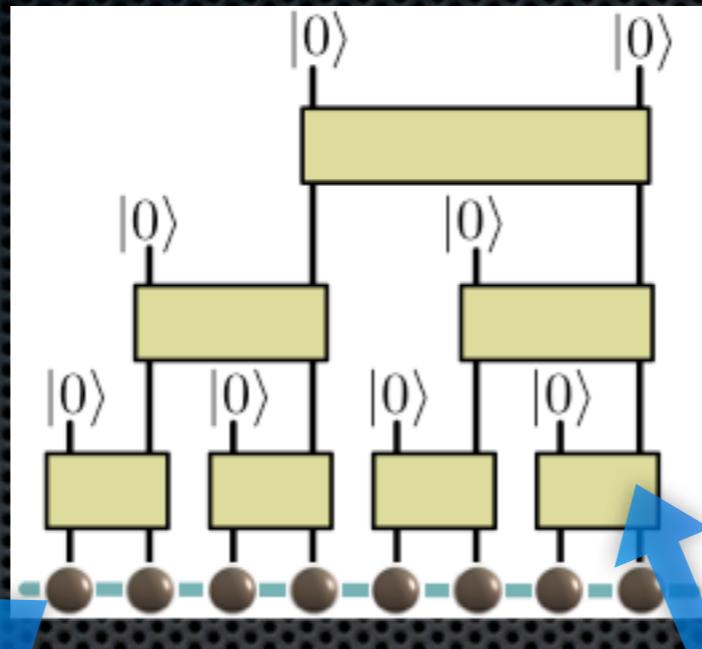


We can get to  
 $E_g > n \cdot o(n)$   
using a TTN  
construction.

Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?

d-level systems



We can get to  
 $E_g > n \cdot o(n)$   
using a TTN  
construction.

Isometry  $V = V(U)$

$$V : \mathbb{C}^d \rightarrow \mathbb{C}^d \otimes \mathbb{C}^d$$

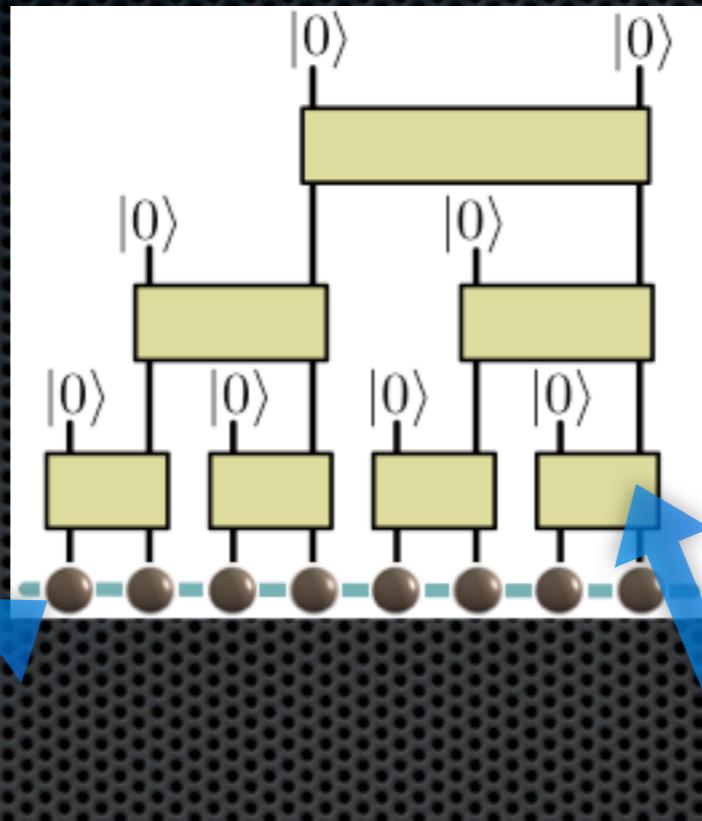
$$V|\beta\rangle = U|0\rangle \otimes |\beta\rangle$$

Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?

d-level systems

Concatenate to get the  
state of  $2^k$  qudits at level k.



We can get to  
 $E_g > n \cdot o(n)$   
using a TTN  
construction.

Isometry  $V = V(U)$

$$V : \mathbb{C}^d \rightarrow \mathbb{C}^d \otimes \mathbb{C}^d$$

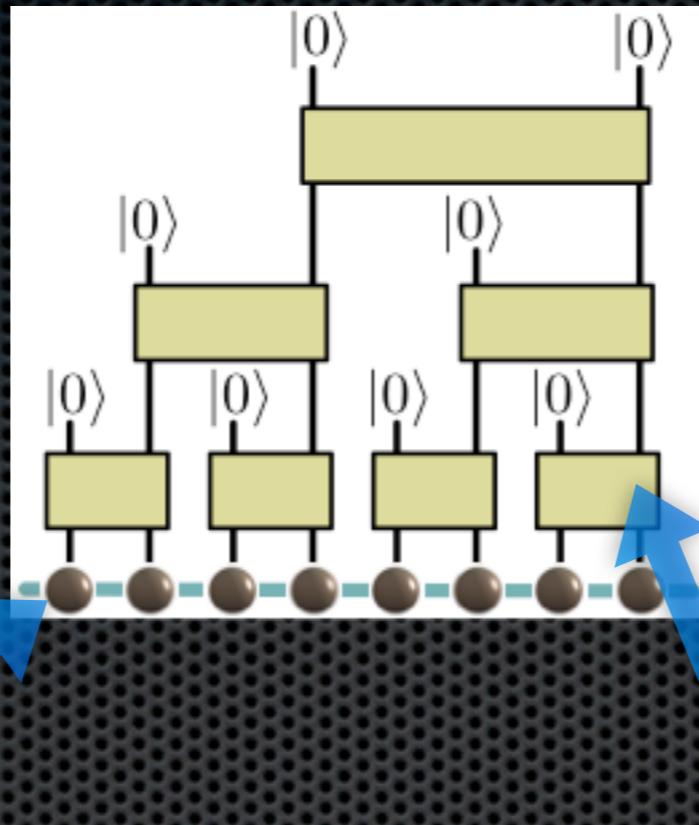
$$V|\beta\rangle = U|0\rangle \otimes |\beta\rangle$$

Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?

d-level systems

Concatenate to get the  
state of  $2^k$  qudits at level k.



We can get to  
 $E_g > n \cdot o(n)$   
using a TTN  
construction.

Isometry  $V = V(U)$

$$V : \mathbb{C}^d \rightarrow \mathbb{C}^d \otimes \mathbb{C}^d$$

$$V|\beta\rangle = U|0\rangle \otimes |\beta\rangle$$

Now choose each  $U$  randomly,  
and let  $d$  grow slowly,  $(\log n)^{1/2}$ .

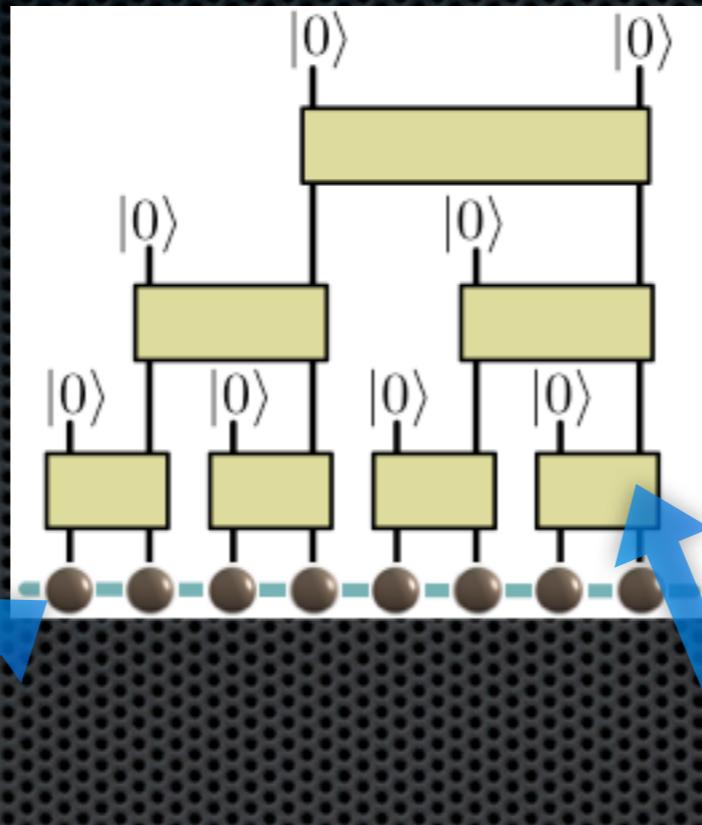
Random states  
are extravagant.

Can provably  
useless states be  
created efficiently?

d-level systems

Concatenate to get the  
state of  $2^k$  qudits at level k.

Now choose each  $U$  randomly,  
and let  $d$  grow slowly,  $(\log n)^{1/2}$ .



We can get to  
 $E_g > n - o(n)$   
using a TTN  
construction.

$$V : \mathbb{C}^d \rightarrow \mathbb{C}^d \otimes \mathbb{C}^d$$

$$V|\beta\rangle = U|0\rangle \otimes |\beta\rangle$$

$$E_g > n - o(n)$$

# Decision problems

I have a generic  $\Psi$ .  
Can I compute  
anything with it?

$|\Psi\rangle$



- For almost every state  $\Psi$ , there is no poly-bounded classical control circuit which allows a significant advantage over classical randomness.  
Only problems in BPP can be solved. (BMW '08)

$$\Pr_{\Psi} \left\{ \exists C \mid C(\Psi) - C(2^{-n} \mathbf{1}) \mid > \epsilon \right\} \leq (8^8 w)^{3v} e^{-c\epsilon^2 2^n}$$

# Randomness vs entanglement?

Random states such that  $E_g \leq \log K + O(1)$  also offer no advantage!

- Choose  $nK$  states at random from  $\mathbb{C}^2$  to construct the following (where  $K$  is superpolynomial in  $n$ ):

$$R := \sum_{j=1}^K |\psi_j^{(1)}\rangle\langle\psi_j^{(1)}| \otimes \cdots \otimes |\psi_j^{(n)}\rangle\langle\psi_j^{(n)}|$$

- Randomly pick a state from the support of  $R$  then:

$$|\Psi\rangle = \frac{1}{\sqrt{\langle\Psi_0|R|\Psi_0\rangle}} \sqrt{R} |\Psi_0\rangle$$

$$\Pr_{\Psi} \left\{ \exists C \mid C(\Psi) - C(2^{-n} \mathbf{1}) \mid > \epsilon \right\} \leq \left( 2^n + (8^8 w)^{3v} \right) e^{-c' \epsilon^2 K^{1/3}}$$



# Questions

- Can we derandomize these constructions?
- Can Hastings' techniques give improved bounds?
- Are efficiently created states subject to this effect?
- What happens with a polynomial number of copies?
- What implications does this have for the circuit model?