## **Decoupling:**

# A building block for quantum information theory



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# Information (Shannon) theory

- A practical question:
  - How to best make use of a given communications resource?
- A mathematico-epistemological question:
  - How to quantify uncertainty and information?
- Shannon:
  - Solved the first by considering the second.
  - −x mathematical theory of communication [1948]
    The

# Shannon theory provides

- Practically speaking:
  - A holy grail for error-correcting codes
- Conceptually speaking:
  - A operationally-motivated way of thinking about correlations
- What's missing (for a quantum mechanic)?
  - Features from linear structure:
     Entanglement and non-orthogonality

## Quantum Shannon Theory provides

- General theory of interconvertibility between different types of communications resources: qubits, cbits, ebits, cobits, sbits...
- Relies on a
  - Major simplifying assumption:
     Computation is free
  - Minor simplifying assumption:
     Noise and data have regular structure

## These lectures

- You will learn one, very powerful result
- Its various names:
  - State transfer
  - Fully quantum Slepian-Wolf
  - The mother of all protocols
- Part I: applications
- Part II: the proof

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VON NEUMANN ENTROPY:

H(A) := H(PA) = -trpa log pa

O & H(A) & log din A

PORE

MUTUAL INFORMATION:

I(A:B) := H(A) + H(B) - H(AB)

O & I(A:B) & 2 log din A

O & T(A:B) & 2 log din A

PRODUCT:

P MAX ENTANGLED

PAS = PA & 9B

CONDITIONAL ENTROPY ! COHERENT INFO

H(AIB) := H(AB) - H(B) =: - I(A)B)

- log din A & H(AIB) & log din A
```

BASIC RESOURCES: SINGLE SENDER &
SINGLE RECEIVER

ALICE & BOB &

[ c→ c]: ONE USE OF AN ALICE → BOB NOISELESS

CBIT CHANNEL ( id, )

[ 9 -> 9]: ONE USE OF AN ALICE -> BOB HOISELESS

[99]: ONE ALICE-BOB EBIT (107107+117117)

RESOURCE INEQUALITIES: A SHORTHAND

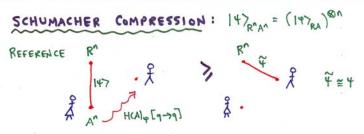
OBVIOUS: [979] > [c rc]

CLHS CAN SIMULATE RHS

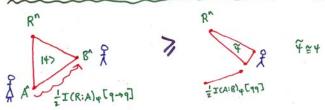
ENTANGLEMENT DISTRIBUTION: [9-79] > [99]

SUPERDENSE CODING: [9-79]+[99] > 2[c->c]

TELEPORTATION: [99]+ 2[c >c] > [9 >9]



DUR MONDAY GOAL: STATE TRANSFER: 147 RAB ON



NOTE: STATE TRANSFER GENERALIZES SCHUMACHER COMPRESSION

IF  $|\Psi\rangle_{RA}$  IS PURE, THEN  $\frac{1}{2}\text{T}(R;A)_{\psi} = \frac{1}{2}\left[H(R)_{\psi} + H(A)_{\psi} - H(RA)_{\psi}\right]$   $= \frac{1}{2}\left[H(A)_{\psi} + H(A)_{\psi} - O\right] = H(A)_{\psi}$   $\uparrow \qquad \qquad \uparrow \qquad \qquad \downarrow \qquad \qquad \downarrow \qquad \qquad \uparrow \qquad \qquad \downarrow \qquad \qquad \downarrow$ 

SAME NON-ZERO EIGENVALVES

#### RESOURCE INEQUALITY FROM STATE TRANSFER

$$\begin{array}{c} \left( \begin{array}{c} P_{AB} \right) \ + \ \frac{1}{2} \, \mathbf{I}(R;A)_{\phi} \left[ q \rightarrow q \right] \ \geqslant \ \frac{1}{2} \, \mathbf{I}(A;B)_{\phi} \left[ q q \right] \end{array} \right) \\ \left( \begin{array}{c} \left( \begin{array}{c} E_{XTRACT} \ E_{N} \ E_{BITS} \ Where \\ I_{Im} \ E_{A}/_{c} = \frac{1}{2} \, \mathbf{I}(A;B)_{\phi} \\ - ALLOW \ Q_{N} \ USES \ OF \ id_{2} \ Where \\ I_{Im} \ Q_{A}/_{c} = \frac{1}{2} \, \mathbf{I}(R;A)_{\phi} \\ - \frac{1}{2} \, \mathbf{I}(R;A)_{\phi} \end{array} \right) \\ ALLOW \ ALICE \ AND \ Bob \ MANY \ USES \ (n \rightarrow \infty) \ Of \\ MIXED \ STATE \ P_{AB} \ ic \ P_{AB} \end{array} \right)$$

SMALL IMPERFECTIONS ALLOWED BUT MUST VANISH AS A +> 00.

#### ENTANGLEMENT DISTILLATION

GOAL: 
$$\langle \Psi_{AB} \rangle + \mathbb{C} [c \to c] \geqslant \mathbb{E} [c \to c]$$

WANT TO REPLACE  $[q \to q]$  IN  $(s\tau)$  BY  $[c \to c]$ . How?

TELEPORTATION:  $[qq] + 2[c \to c] \geqslant [q \to q]$   $(\tau p)$ 

SUBSTITUTE  $(\tau p)$  INTO  $(s\tau)$ :

 $\langle \Psi_{AB} \rangle + \frac{1}{2} I(R;A)_{\phi} [[qq] + 2[c \to c]] \geqslant \frac{1}{2} I(A;B)_{\phi} [qq]$ 
 $\langle \Psi_{AB} \rangle + I(R;A)_{\phi} [c \to c] \geqslant \left[ \frac{1}{2} I(A;B)_{\phi} - \frac{1}{2} I(R;A)_{\phi} \right] [qq]$ 
 $\langle \Psi_{AB} \rangle + I(R;A)_{\phi} [c \to c] \geqslant I(A)_{\phi} [qq]$ 

HASHING INEQUALITY

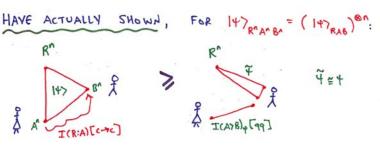
CHECK: 
$$15 \frac{1}{2} I(A:B)_{\phi} - \frac{1}{2} I(R:A)_{\phi} = I(A)_{\phi}^{2}$$
 For  $147_{RAB}^{2}$ ?

$$\frac{1}{2} I(A:B)_{\phi} - \frac{1}{2} I(R:A)_{\phi} = \frac{1}{2} \left\{ H(A)_{\phi} + H(B)_{\phi} - H(AB)_{\phi} - H(R)_{\phi} + H(R)_{\phi}^{2} \right\}$$

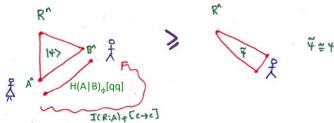
$$= \frac{1}{2} \left\{ H(B)_{\phi} - H(AB)_{\phi} - H(AB)_{\phi} + H(B)_{\phi}^{2} \right\}$$

$$= H(B)_{\phi}^{2} - H(AB)_{\phi}$$

$$= I(A)_{\phi}^{2} B)_{\phi}$$



## OR, EQUIVALENTLY,



NOTE: H(AIB) ? O ARE BOTH POSSIBLE!

>0: ENTANGLEMENT IS CONSUMED

CO: ENTANGLEMENT IS GENERATED

INTERPRETATION OF H(AIB) 4:

(QUANTUM) UNCERTAINTY ABOUT A WHEN B IS GIVEN

CO: MORE THAN CERTAIN

ANOTHER APPLICATION: NOISY SUPERDENSE CODING

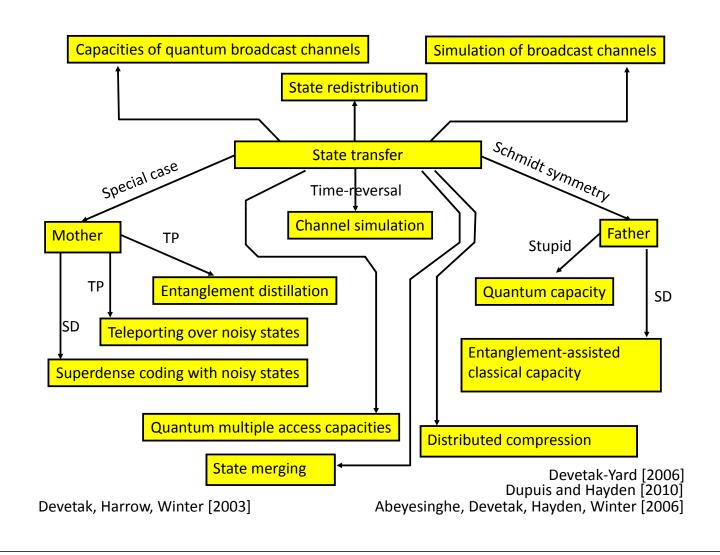
EXPLOIT MIXED PAB TO EMMANCE CLASSICAL COMMUNICATION

WANT [ C+C] OH RHS. HOW?

USE SUPERPEASE CODING:

SUBSTITUTE (SDC) INTO (ST):

CHECK: 
$$\frac{1}{2} I(R;A)_{\varphi} + \frac{1}{2} I(A;B)_{\varphi}$$
  
= $\frac{1}{2} \left\{ H(R)_{\varphi} + H(A)_{\varphi} - H(RA)_{\varphi} + H(A)_{\varphi} + H(B)_{\varphi} + H(AB)_{\varphi} \right\}$   
= $\frac{1}{2} \left\{ H(R)_{\varphi} + H(A)_{\varphi} - H(RA)_{\varphi} + H(A)_{\varphi} + H(RA)_{\varphi} - H(R)_{\varphi} \right\}$   
=  $H(A)_{\varphi}$ 



# Part II: Proof of the state transfer theorem



(The easy route to quantum information guru status)

### MIXED STATE DISTANCE MEASURES

#### 1) FIDELITY (GENERALIZES INNER PRODUCT)

$$F(p, s) = \begin{cases} 1 & \text{iff } p = s \\ 0 & \text{iff } p = s \end{cases}$$

#### 2) TRACE DISTANCE

RECALL  $\|X\|_{L^{\infty}} = \frac{1}{2} \frac{1}{2}$ 

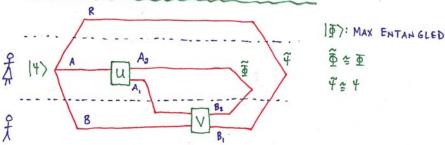
USE NORM TO MEASURE DISTANCE:
D(p, 6):= ||p-6||,

#### (1) \$ (2): ROUGH EQUIVALENCE

 $1-\sqrt{F} \leq \frac{D}{2} \leq \sqrt{1-F}$ 

1-D & F & 1- \frac{D^2}{4}

### STATE TRANSFER AS A CIRCUIT



NEED TO FIND ISOMETRIC U, V ACCOMPLISHING THIS GOAL WANT: (VoU) | 47 RAB ≅ | 47 RB, | \$\Pi\$ \rangle\_{A\_2B\_2}\$

### THE DECOUPLING ARGUMENT

SUPPOSE  $1-\epsilon < F( \psi_R \otimes \Phi_{A_2} / \omega_{RA_2} ) \leftarrow (*)$ 

DEFINITION OF FIDELITY YIELDS V AUTOMATICALLY

WILL SUFFICE TO DEMONSTRATE (\*)

STRATEGY: SHOW THAT A "TYPICAL" U WILL DO THE JOB

TH'M [DECOUPLING]: 
$$\int_{U(A)} \| \delta(u)_{RA_2} - \psi_R \otimes^{T_{A_2}} /_{J_{A_2}} \|_1^2 dU$$

$$\leq \frac{d_R d_A}{d_A^2} + \Gamma \left[ (\psi_{RA})^2 \right]$$
TO INTERPRET: RECALL SCHUMACHER COMPRESSION

LET  $\Pi_X$  BE THE TYPICAL PROJECTOR ON  $X$  FOR  $P^{\otimes n}$ 

THEN: (i) It  $\Pi_X P_X^{\otimes n} > 1 - \epsilon$ 
(ii)  $\Gamma_{A_1} P_X^{\otimes n} = \Gamma_{A_2} P_{A_2} P_{A_2}$ 

### APPLY TO DECOUPLING TH'M

SPS EVERYTHING IS TYPICAL:

147 RAB & (TTR & TTA O TTB) 147 RAB

THEN:  $\log \left\{ \frac{d_{R} d_{A}}{d_{A}^{2}} + \Gamma \left( (4_{RA})^{2} \right) \right\} = \log d_{R} + \log d_{A} + \log 4 \Gamma \left( (4_{RA})^{2} \right) - 2 \log d_{A},$   $= \log d_{R} + \log d_{A} + \log 4 \Gamma \left( (4_{R})^{2} \right) - 2 \log d_{A},$   $\sim n H(R)_{\varphi} + n H(A)_{\varphi} - n H(R)_{\varphi} - 2 \log d_{A},$   $= n H(R)_{\varphi} + n H(A)_{\varphi} - n H(RA)_{\varphi} - 2 \log d_{A},$   $= n I(R;A)_{\varphi} - 2 \log d_{A},$ 

SO, GOOD DECOUPLING PROVIDED

2 log da, >> n I (R:A) &

(EXACTLY RATE REQUIRED FOR

STATE TRANSFER!

## ASIDE : INTEGRATION ON THE UNITARY GROUP

WARM-UP: INTEGRATION ON A SPHERE

WHAT DO WE MEAN BY UNIFORM?
INVARIANCE UNDER ROTATIONS

$$\int_{S^d} f(\vec{x}) d\vec{x} = \int_{S^d} f(R\vec{x}) d\vec{x} \quad \forall \text{ Rotations } R.$$

GROUP INTEGRATION AND SYMMETRY

NATURAL GROUP SYMMETRIES:

- · LEFT AND RIGHT MULTIPLICATION
- · UNITARY GROUP HAS A UNIQUE MEASURE INVARIANT UNDER BOTH (HAAR MEASURE)

$$\int_{U(d)} f(g) dg = \int_{U(d)} f(hg) dg$$

$$= \int_{U(d)} f(gh) dg$$

$$\forall h \in U(d)$$

NORMALIZATION: FIX July 1 dg = 1.

THESE RELATIONS ARE SUFFICIENT FOR CALCULATION

CHEAT SHEET: 
$$\int f(g) dg = \int f(hg) dg = \int f(gh) dg$$
  $\forall h$ 

$$\int dg = 1$$

EXAMPLE 1: LET  $D(x) = \int U x u^{4} du$ .

THEN  $D(x) = hr(x) \frac{1}{d}$ . is  $D$  is completely depolarizing

To see this, note that:
$$V D(x) V^{4} = V \left\{ \int u x u^{4} du \right\} V^{4}$$

$$= \int (vu) Y \left( vu^{4} \right) du \quad (BY \text{ Linearity of } \int du)$$

$$= \int u x u^{4} du \quad (BY \text{ Invariance of } du)$$

$$= D(x)$$

So  $V D(x) = 0$  For all unitaries  $V D(x) = 0$ 

$$V D(x) = 0$$

= Str X du (u unitary)

= tr X (NORMALIZATION Sour 1)

• CONCLUDE 
$$b(x) = tr(x) \frac{1}{4}$$

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EXAMPLE 2: CALCULATE 2(x)= (uou) x (utout) du
   AS BEFORE, NOTE
   (\vee \otimes \vee) \chi(\chi) (\vee \otimes \vee^{+}) = [(\vee u \otimes \vee u) \chi [(\vee u)^{+} \otimes (\vee u)^{+}] du
                      = ( ( U & W ) X ( W + O U + ) d U ( INVAPIANCE OF ) du)
  = Z(X)

So [X(X), V&V]=D FOR ALL UNITARY V
  FIND Y SUCH THAT [Y, VOV] = 0 YVE ((d)
       · Y= I WOULD WORK. OTHERS?
       · CONSIDER SWAP OFFRATOR F DEFINED SO THAT
               4197, 167: F147 167 = 167 147
        THEN
        · (V&V) F 19716) = (V&V) 167147 = V1678 V147
        . F(VOV) 147 (W) = F(V1470V1W)= VIW>0V19>
        SO [F, VOV]= O YVEU(d)
     · TAKE LINEAR COMBINATIONS Y= & I + BF.
          THAT'S IT! PROOFS: · LINEAR ALGEBRA BRUTE FORCE
                                 · REPRESENTATION THEORY (BETTER)
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EXAMPLE 2 (CONT.D.):  $\chi(x):=\int (u \otimes u) \chi(u^{\dagger} \otimes u^{\dagger}) du = d(x) T + \beta(x) F$ More on F:

Note  $F^2 = I$  so  $Fic(F) = \begin{bmatrix} \pm 1 \\ \end{bmatrix}$ LET S<sub>+</sub> BE THE +1 (SYMMETRIC) EIGENSPACE

LET S<sub>-</sub> BE THE -1 (ANTISYMMETRIC) EIGENSPACE

LET  $\Pi_+$  AND  $\Pi_-$  BE ORTHOGONAL PROJECTORS ONTO S<sub>+</sub> AND S<sub>-</sub>

THEN:  $I = \Pi_+ + \Pi_-$  AND  $F = \Pi_+ - \Pi_-$ so  $\Pi_{\pm} = \frac{I \pm F}{2}$ BACK TO  $\chi(x)$ :

CAN WRITE  $\chi(x)$  IN TERMS OF  $\Pi_{\pm}$ .

SINCE  $\int (u \otimes u) \Pi_{\pm} (u^{\dagger} \otimes u^{\dagger}) du = \Pi_{\pm}$ ,  $G \in T$   $\chi^2 = \chi$ .

ALSO,  $\chi$  IS TRACE-PRESERVING.

CONCLUDE:  $\chi(x) = fr(\chi \Pi_+) \frac{\Pi_+}{r \kappa \Pi_+} + fr(\chi \Pi_-) \frac{\Pi_-}{r \kappa \Pi_-}$ 

NOTE S = span [li7]j>-lj>li7; Isicjsd}

So rk T = dim S = ( 2) = d(d-1)

rk TI = d- dim S = d(d+1)

FINDING TK TIT:

DECOUPLING THEOREM: 
$$\int \| \delta(u)_{RA_{2}} - \psi_{R} \otimes^{T_{A_{2}}} / \|_{1}^{2} du \leq \frac{d_{R} d_{A}}{d_{A_{1}}^{2}} + r \left[ (\psi_{RA})^{2} \right]$$

$$| \psi \rangle \bigwedge_{R} \frac{A_{2}}{d_{A_{1}}^{2}} = \int \psi_{R} \psi_{R}$$

WHY 2? . USES tr [ X12 (Y, & I2)] = tr X, Y, AND 6 = tR.

MUST EVALUATE [tr[(6(u) RA,)2] du

$$\begin{split} & \| Y \|_{R}^{R_{A}} \| \frac{\Delta_{A}}{\Delta_{A}} \|_{L^{2}}^{2} \| dU \\ & = \int & \{ \{ (\delta_{RA_{a}})^{2} \} dU \} \\ & = \int & \{ (\delta_{RA_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}})) \} dU \} \\ & = \int & \{ (\delta_{RA} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}})) \} dU \} \\ & = \int & \{ (\delta_{RA} \otimes \delta_{R_{a}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}})) \} dU \} \\ & = \int & \{ (\delta_{RA} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}})) ) \} dU \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ) \} dU \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ) \} dU \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ) \} dU \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) \} dU \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ) \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ) \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) \} \\ & = \int & \{ (\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ((\delta_{RR_{a}} \otimes \delta_{R_{A}}) ) ((\delta_{RR_{a}} \otimes \delta$$

ALMOST PONE!

HAVE 
$$\| \omega(u)_{RA_{2}} - \psi_{R} \otimes^{T_{A_{1}}} / d_{A_{2}} \|_{2}^{2}$$

$$= \operatorname{tr} \left[ (\omega(u)_{RA_{2}})^{2} - \frac{1}{d_{A_{1}}} \operatorname{tr} \left[ (\psi_{R})^{2} \right] \right]$$

AND  $\int \operatorname{tr} \left[ (\omega(u)_{RA_{2}})^{2} \right] du \leq \frac{1}{d_{A_{1}}} \operatorname{tr} \left[ (\psi_{R})^{2} \right] + \frac{1}{d_{A_{2}}} \operatorname{tr} \left[ (\psi_{R})^{2} \right]$ 

CANCELLATION!

$$\Rightarrow \int \| \omega(u)_{RA_{2}} - \psi_{R} \otimes^{T_{A_{2}}} / d_{A_{2}} \|_{2}^{2} du$$

$$\leq \frac{1}{d_{A_{1}}} \operatorname{tr} \left[ (\psi_{RA})^{2} \right] + \frac{1}{d_{A_{2}}} \operatorname{tr} \left[ (\psi_{R})^{2} \right] - \frac{1}{d_{A_{2}}} \operatorname{tr} \left[ (\psi_{R})^{2} \right]$$

BUT BY CAUCHY-SCHWARZ,  $\sum_{i=1}^{m} |\psi_{i}|^{2} = \sum_{i=1}^{m} |\psi_{i}|^{2} = \sum_{i=1}^{m$ 

# Further reading on arXiv (beyond Nielsen and Chuang)

- State transfer:
  - Abeyesinghe, Devetak, Hayden and Winter: The mother of all protocols
- Resource inequalities
  - Devetak, Harrow and Winter: A family of quantum protocols
  - Devetak, Harrow and Winter: A resource framework for quantum Shannon theory
- State merging: interpretation of H(A|B)
  - Horodecki, Oppenheim and Winter: Quantum information can be negative
- Further development of decoupling approach
  - Oppenheim: State redistribution as merging
  - Dupuis: The decoupling approach to quantum information theory
- State transfer and black holes:
  - Hayden and Preskill: Black holes as mirrors