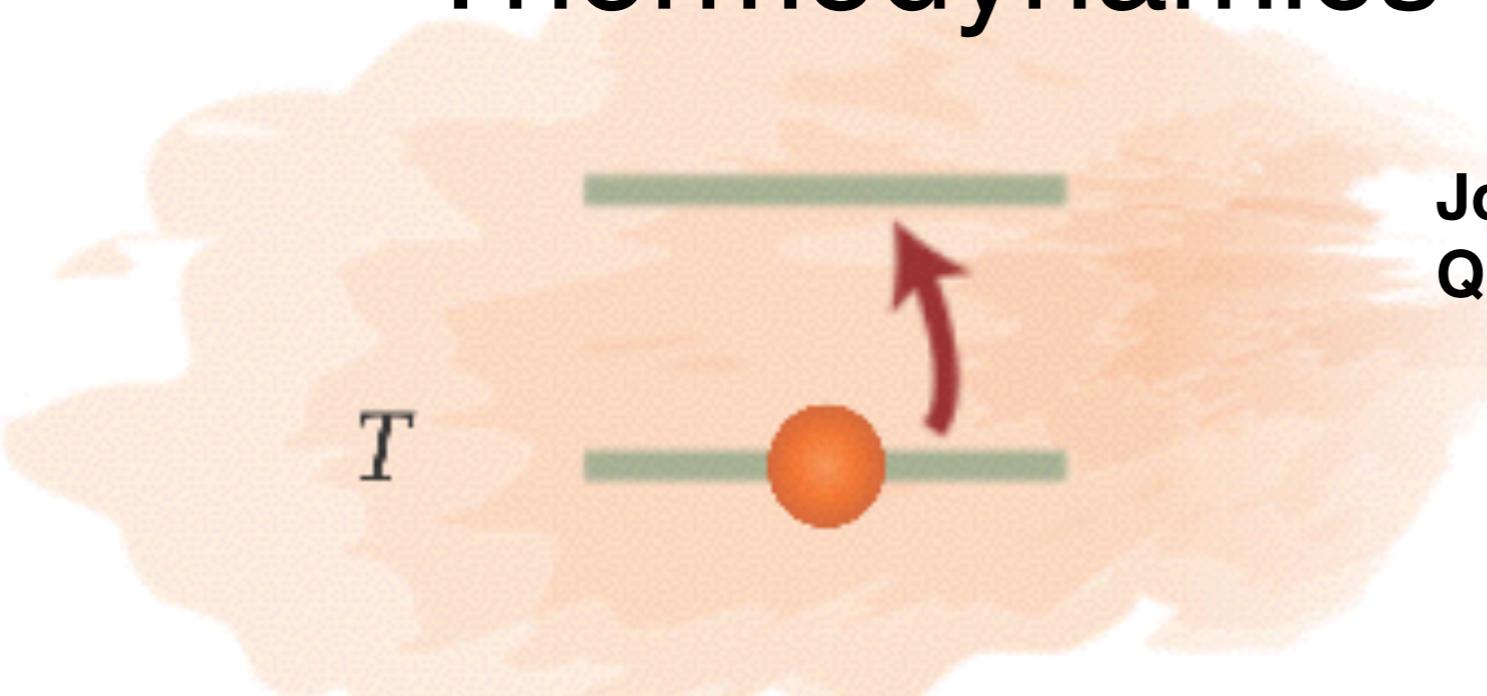


The 2nd laws of Quantum Thermodynamics



Jonathan Oppenheim (UCL)
QIP 2014

1305.5278 Brandao, Horodecki, Ng, J.O., Wehner

1111.3834 Horodecki, JO; (Nature Comms)

1111.3882 Brandao, Horodecki, JO, Renes,
Spekkens (PRL)

0212019 Horodecki², JO (PRA)

Outline

- 2nd law: allowable state transitions $(\rho, H) \rightarrow (\sigma, H')$
- Many 2nd laws: $F \longrightarrow F_a$
- Many families of 2nd laws depending on “how cyclic” the process is
- 0th, 1st law: class of operations
- Tools from QI:
 - resource theories
 - catalytic majorization
 - quantum Renyi-divergences
 - entanglement embezzling

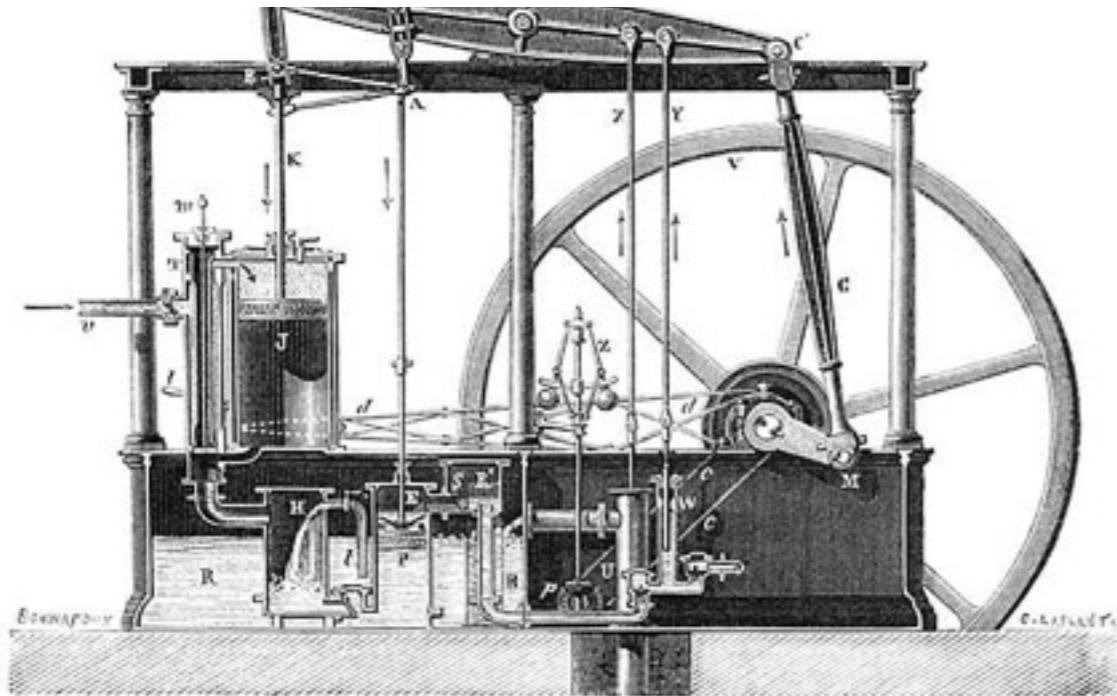
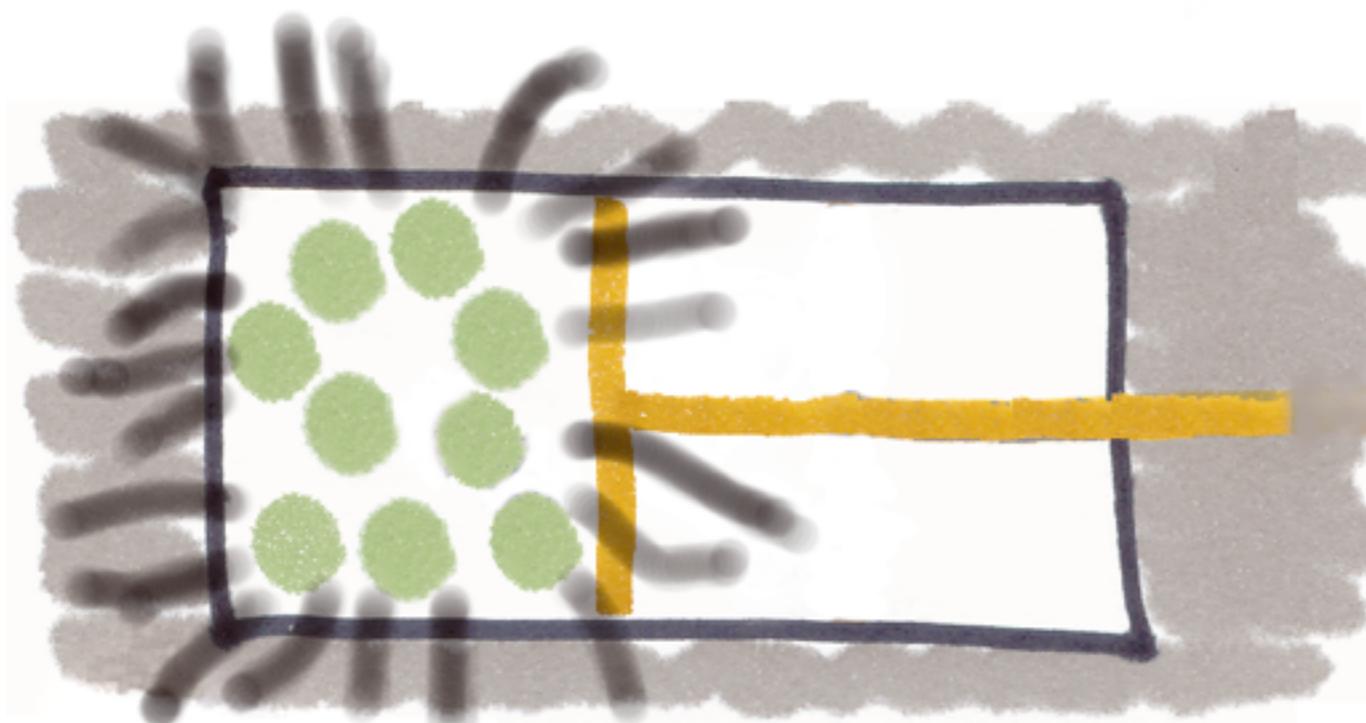


Fig. 59. — Machine à balancier de Watt.

T, tuyau de prise de vapeur; T, tiroir; J, cylindre; H, condenseur; P, pompe d'épuisement; W, pompe alimentaire de la chaudière; UX, pompe d'alimentation de la bûche R; p, Z régulateur; d, excentrique; ABCD parallélogramme; GM bielle et manivelle; V, volant.

1st wave
(phenomenology)
Carnot (1824)
Joule (1843)
Kelvin (1849)
Clausius (1854)



2nd wave
(stat mech)
Maxwell (1871)
Boltzman (1875)
Gibbs (1876)

Quantum Thermodynamics



- **Gibbs state with full information**
 - Gemmer, Michel, Mahler (2005), Popescu, Short, Winter (2006)
- **Meaning of Negative Entropy, conditional erasure**
 - Del Rio et al. (2011), Faist et al. (2013)
- **Smallest possible fridges**
 - Linden, Popescu, Skrzypczyk (2010)
- **Deterministic work extraction**
 - HHO (2003), Dahlston et al. (2010), HO (2011), Aaberg (2011), Egloff et al. (2012)
- **Average work extraction**
 - Brandao et. al. (2011), Skrzypczyk et. al. (2013)
- **Non-ideal heat baths, correlations, entanglement**
 - Reeb, Wolf (2013); Gallego et. al. (2013), Hovhannisyan et. al. (2013)
- **Thermalisation times**
- **Micro engines & machines**
 - Scovil & Schultz-Dubois (1959), Howard (1997)
 - Rousselet et al. (1994), Faucheux et al. (1995), Scully (2002)
- **Prehistory**
 - Ruch and Mead (75), Janzing et. al. (2000)

Thermodynamics of quantum systems

- No thermodynamic limit!
- Coherences in energy basis
- More precise control?
- Rigorous theory?

3 laws of classical thermodynamics

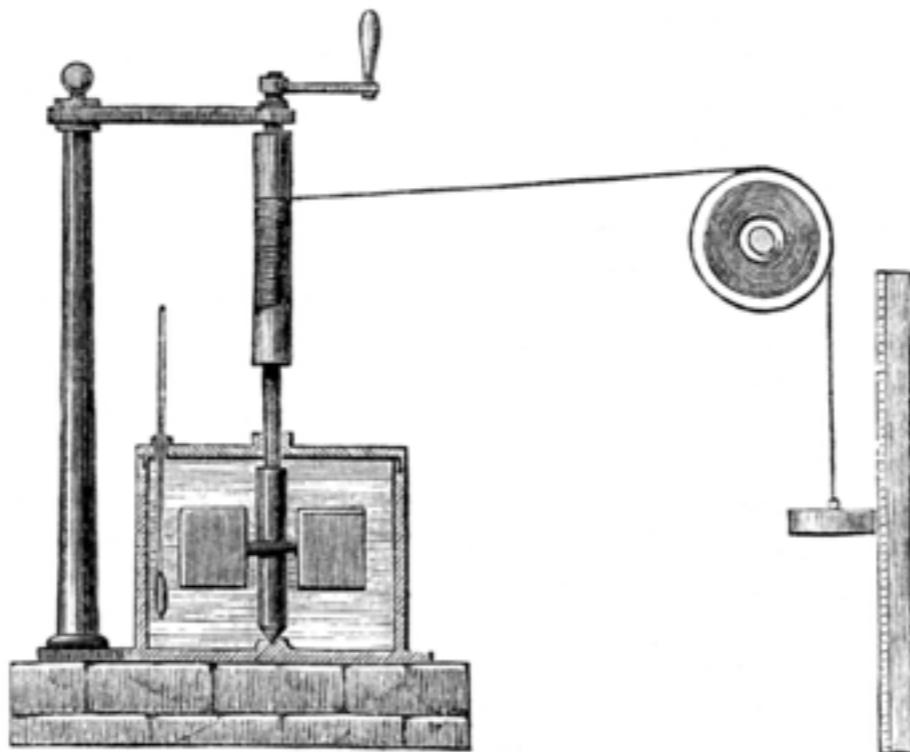
0) If R_1 is in equilibrium with R_2 and R_3 then R_2 is in equilibrium with R_3

1) $dE = dQ - dW$ (energy conservation)

2) Heat can never pass from a colder body to a warmer body without some other change occurring. – Clausius

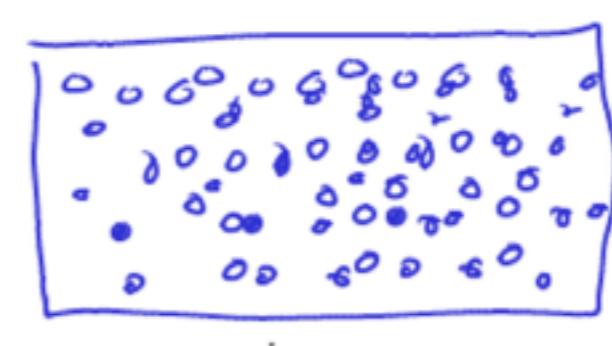
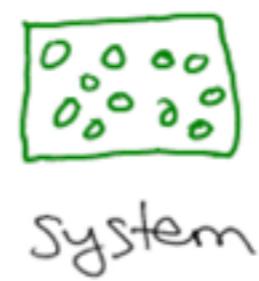
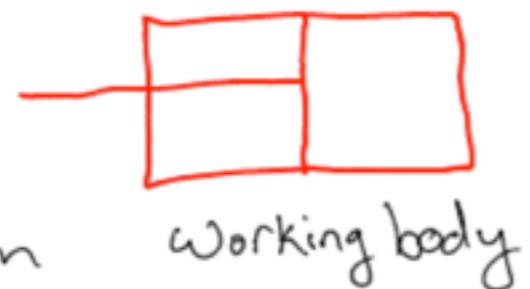
3) One can never attain $S(p) = C$ in a finite number of steps

$$1) \quad dE = dQ - dW \quad (\text{energy conservation})$$



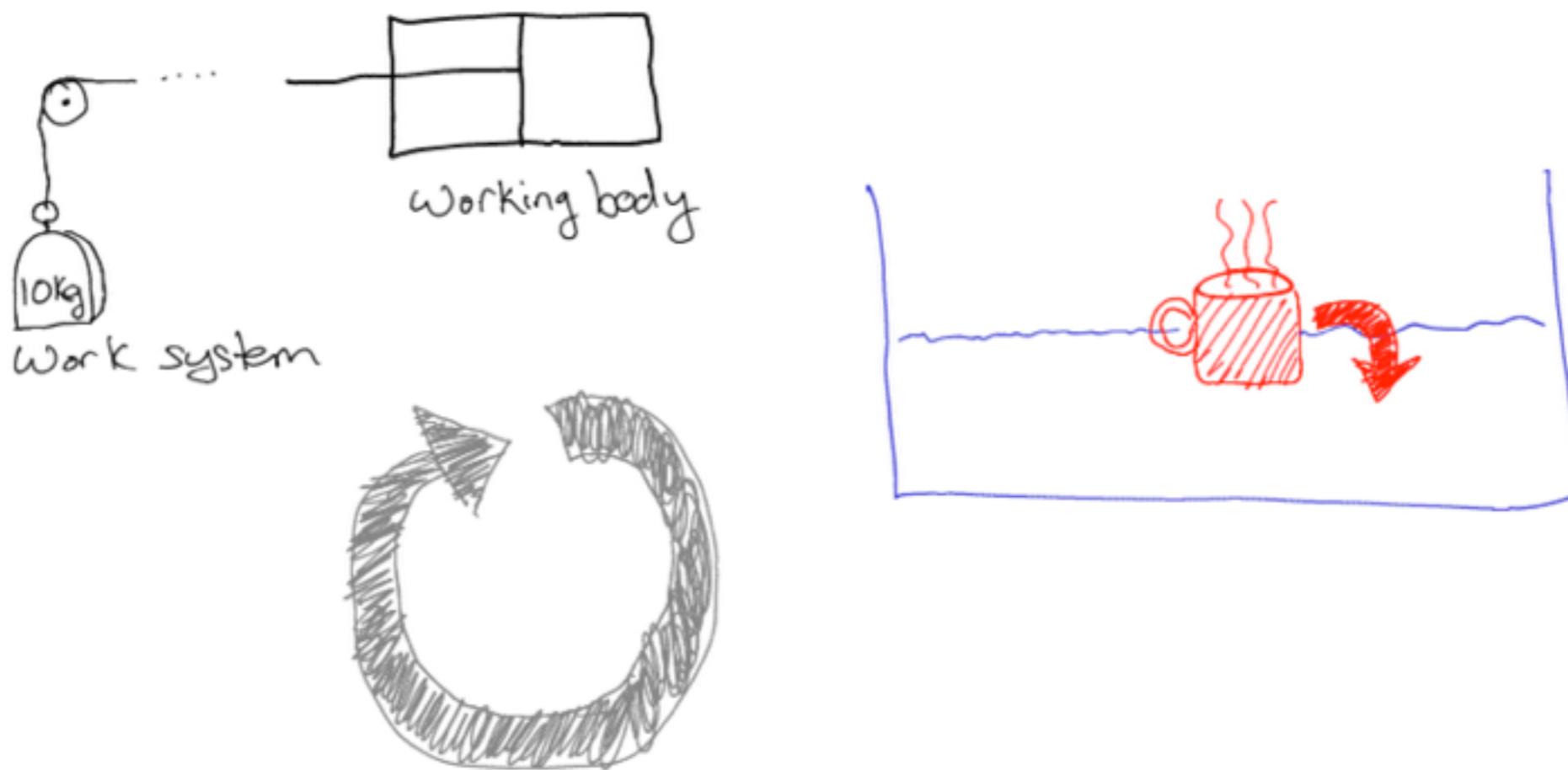
Joule (1843)

Not a consequence, but part of the class of operations



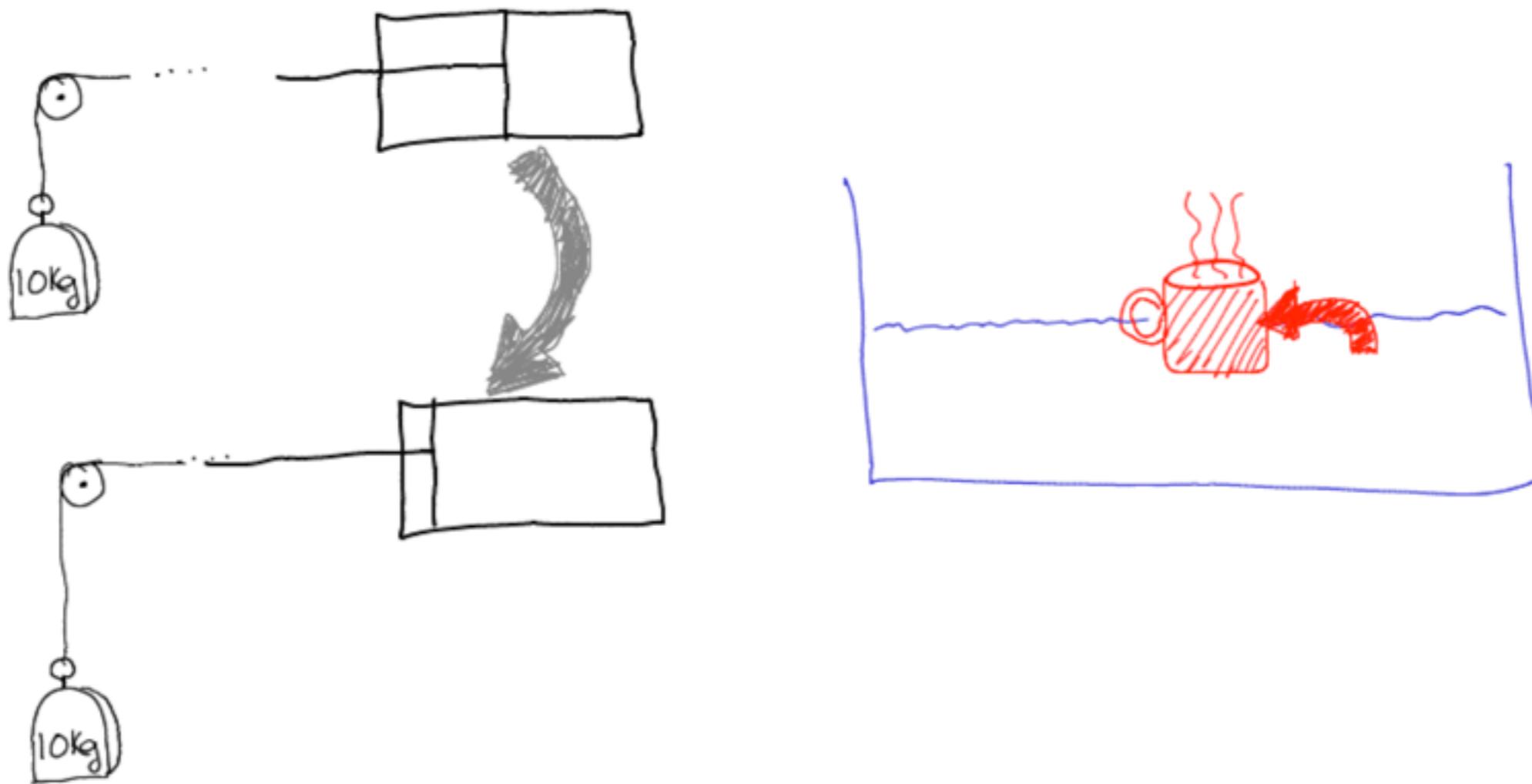
The second law

Heat can never pass from a colder body to a warmer body without some other change occurring. – Clausius



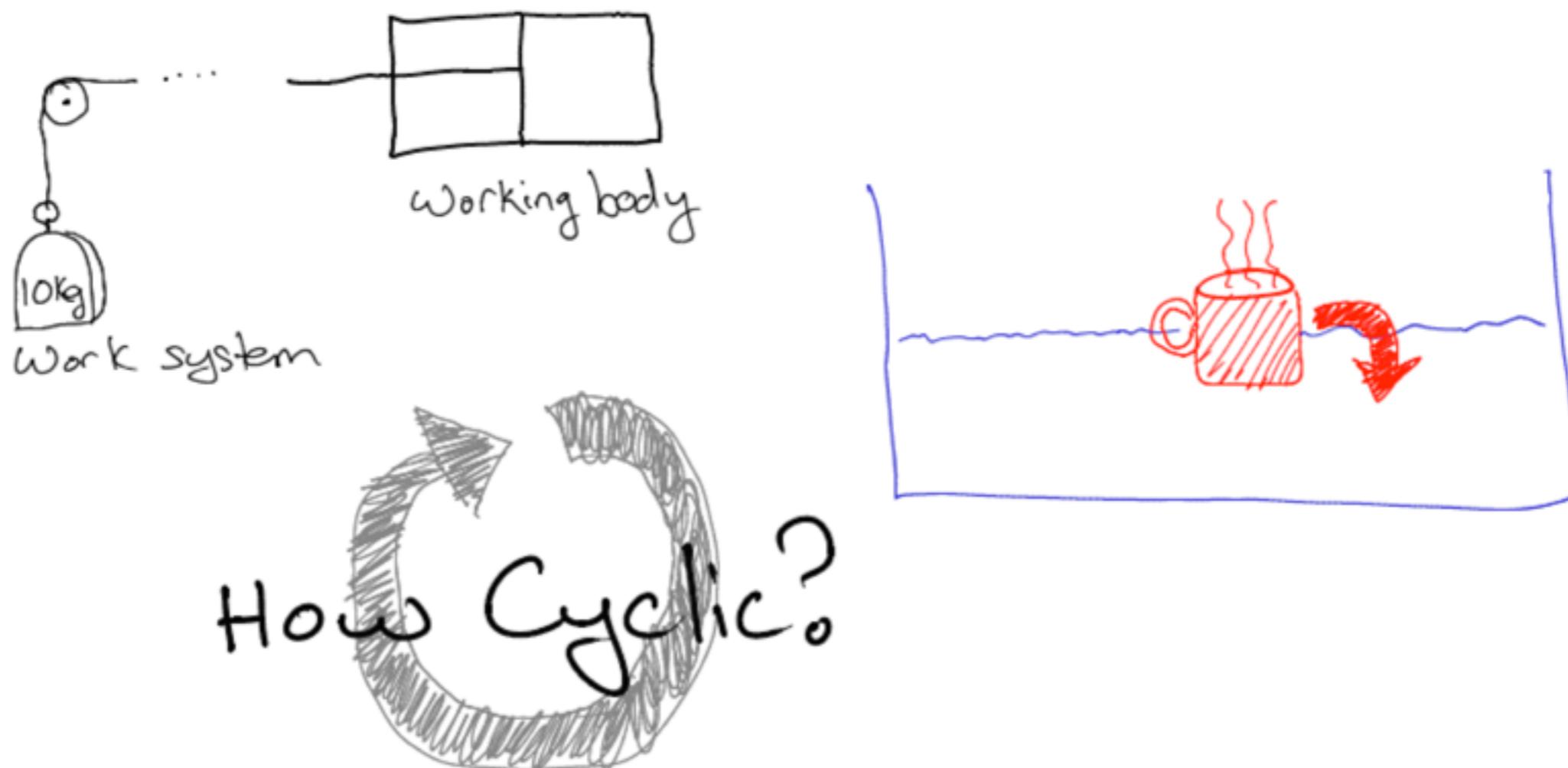
The second law

Heat can never pass from a colder body to a warmer body without some other change occurring. – Clausius



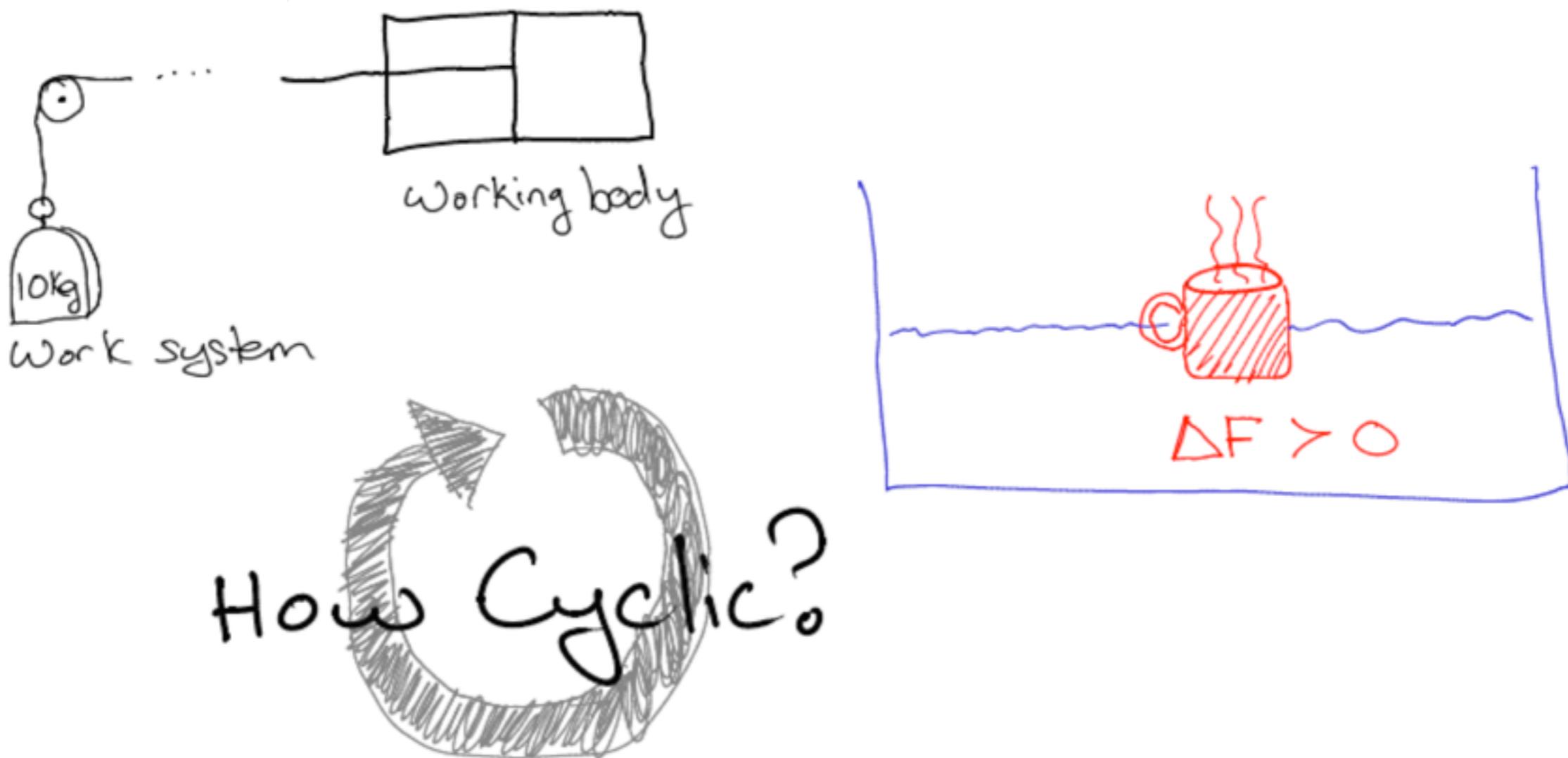
The second law

Heat can never pass from a colder body to a warmer body without some other change occurring. – Clausius

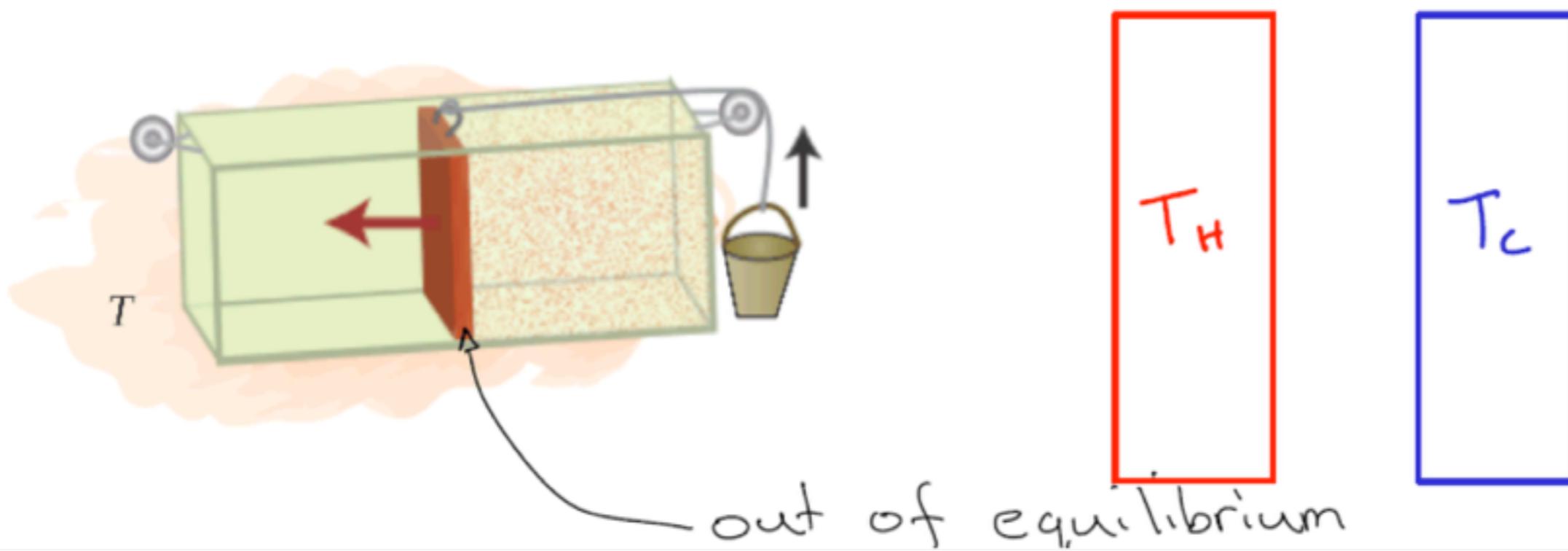


The second law (+ first law)

In any cyclic process, the free energy of a system can only decrease.



Free Energy



$$F = E - TS$$

$$W_{\text{gain}} = F(\rho_{\text{initial}}) - F(\rho_{\text{final}})$$

$$\rho_{\text{initial}} \rightarrow \rho_{\text{final}} \quad \text{only if} \quad \Delta F \geq 0$$

Catalytic Thermal Operations \wedge_{τ}

- (ρ_s, H_s)
 - ρ_s = resource
 - H_s = Hamiltonian
- adding arbitrary many copies of state τ_R
- borrowing ancillas (working body) and returning them in the “same” state σ_c
- energy conserving unitaries U
(1st law) $[U, H_s + H_R + H_c] = 0$
- tracing out (trash)

Thermal Operations



or in the micro - regime



1 law of quantum thermodynamics

Class of operations:

- 0) The only free state τ which doesn't enable arbitrary transitions is the thermal state $\tau = e^{-\beta H_R} / Z$
- 1) Energy conserving unitaires

Thermo-monotones:

- 2)* The state ρ_s must get closer to τ_s in terms of free energy type distances $F_a(\rho_s \parallel \tau_s)$ $\forall a \leq 0$

Zeroeth Law

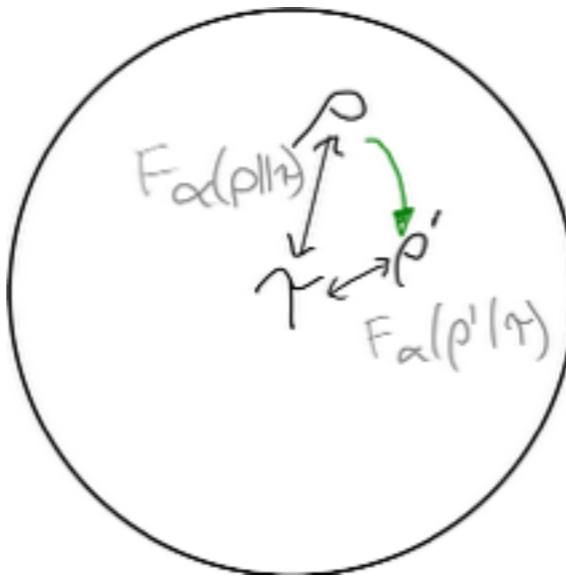
After decohering in the energy eigenbasis, one can extract work from many copies of any state which is not passive
 $(p_i \leq p_j \text{ iff } E_i \geq E_j)$

[just swap levels i and j , while raising the weight, and repeat over many blocks]

Many copies of any state except the thermal state results in a state which is not passive after decohering (Pusz and Woronowicz (78), Brandao et. al. 2011).

This gives us an equivalence class, of allowed free states (τ_β, H_R) labelled by β . Any other free state allows arbitrary transitions.

The second laws (pseudo-classical)



Thermal monotonies: $F_a(\rho||\tau) = \frac{kT}{1-a} \log \text{tr } \rho^a \tau^{1-a} - kT \log Z \quad a \geq 0$

- Ordinary 2nd law is one of many
- In macroscopic limit, weak interactions, all $F_a \simeq F$
- For ρ block diagonal, 2nd law is necessary and sufficient

Quantum Second Laws

$$\hat{F}_a(\rho||\tau) := \frac{kT}{a-1} \log \left(\text{tr} \left(\tau^{\frac{1-a}{2a}} \rho \tau^{\frac{1-a}{2a}} \right)^a \right) - kT \log Z$$

Müller–Lennert et al.
Wilde et al.
Jaksic et al. (2013)

or

$$\tilde{F}_a(\rho||\tau) := kT \frac{\text{sgn}(a)}{a-1} \log \text{tr} \rho^a \tau^{1-a} - kT \log Z \quad \text{Petz (86)}$$

\tilde{F}_a monotonic under CPT maps for $\alpha \in [0, 1]$

\hat{F}_a monotonic for $\alpha \geq \frac{1}{2}$ Frank and Lieb (2013)

$$\tilde{F}(\rho||\tau) \geq \tilde{F}(\wedge_\tau(\rho)||\wedge_\tau(\tau)) = \tilde{F}(\rho'||\tau)$$

Quantum Second Laws

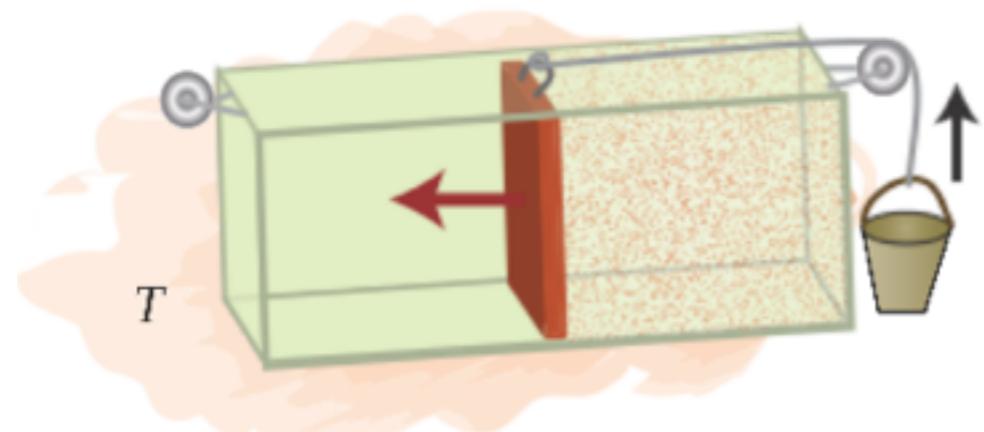
$$\rho \xrightarrow{\text{C.T.O.}} \rho' \quad \text{only if}$$

$$\hat{F}_\alpha(\rho||\tau) \geq \hat{F}_\alpha(\rho'||\tau) \quad \alpha \geq \frac{1}{2}$$

$$\hat{F}_\alpha(\tau||\rho) \geq \hat{F}_\alpha(\tau||\rho) \quad \frac{1}{2} \leq \alpha \leq 1$$

$$\hat{F}_\alpha(\rho||\tau) \geq \hat{F}_\alpha(\rho'||\tau) \quad 0 \leq \alpha \leq 2$$

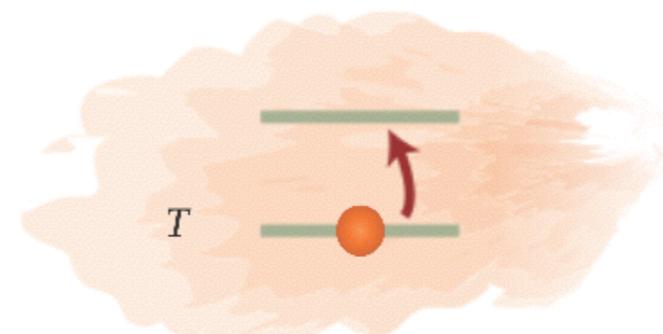
Macro



weight



Micro



ω +

$$F_\infty = kT \log \min\{\lambda : \rho \leq \lambda \tau\} - kT \ln Z$$



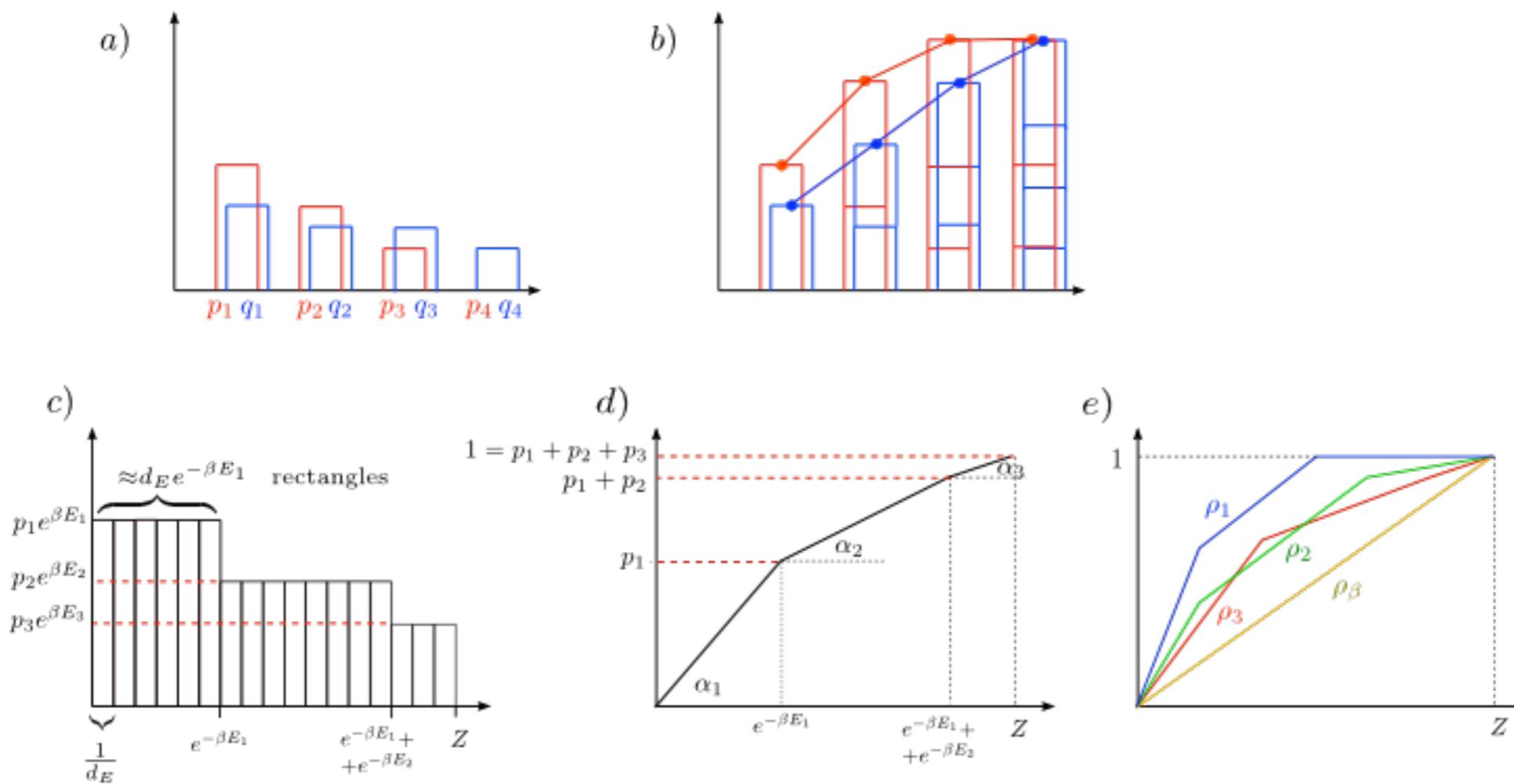
$$F_0 = kT \log \sum h(\omega, g, E_i) e^{-\beta E_i}$$

$$W_{land} = kT \log(1 + e^{-\beta E_e})$$

Quasi-classical 2nd laws

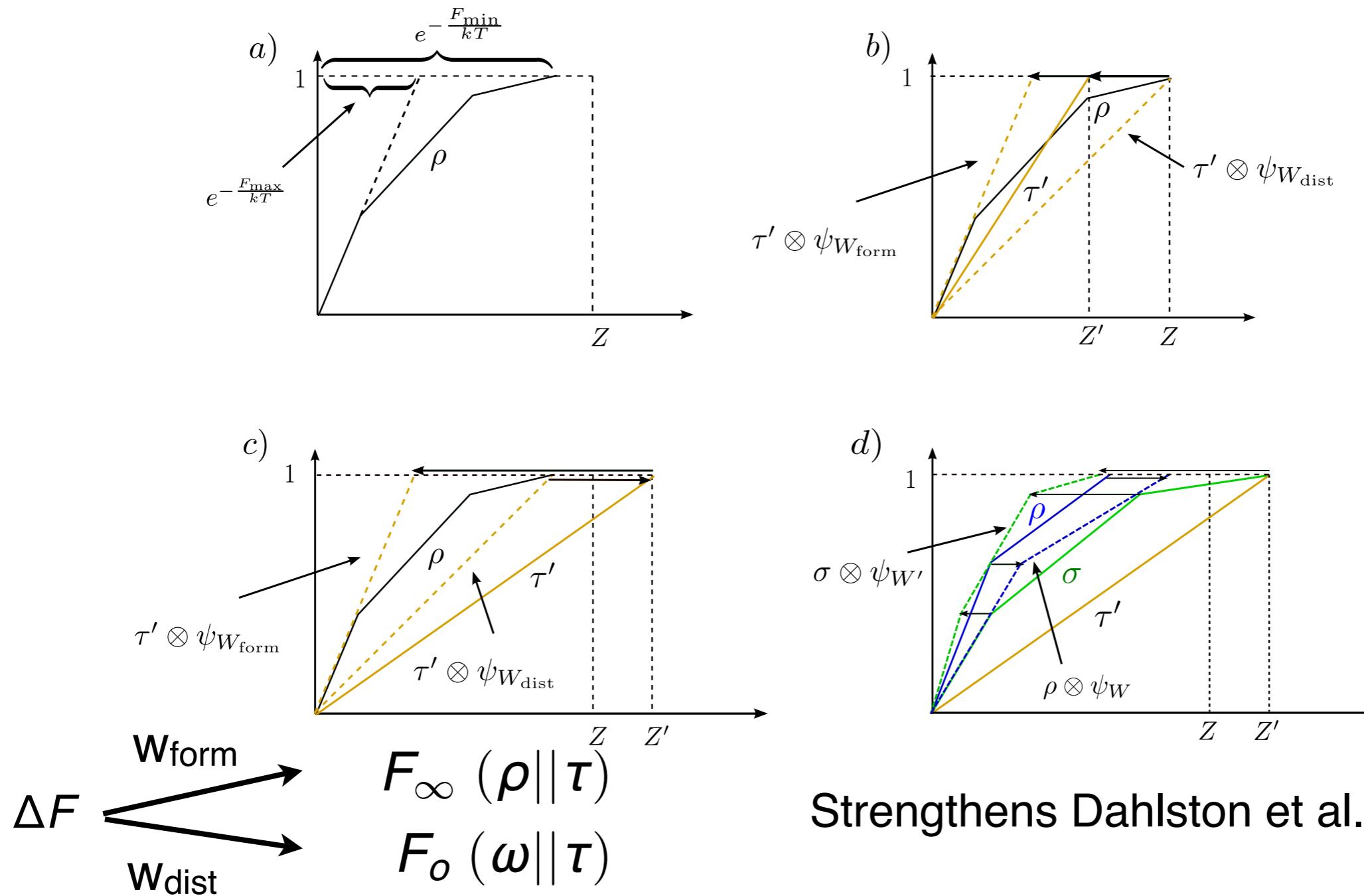
For degenerate energy levels

$$x \xrightarrow{\text{N.O.}} y \quad \text{iff } x > y \text{ Horodecki, JO (2003)}$$



$$P(E_1, g_1) e^{\beta E_1} \geq P(E_2, g_2) e^{\beta E_2} \geq \dots \beta - \text{ordered conjecture of Ruch \& Mead (75)}$$

Work of $(\rho, H) \rightarrow (\sigma, H')$



Strengthens Dahlston et al. for $H = 0$

c.f. Ahlberg (2011)

Thermo-majorization

Majorization within energy blocks

$\tau_R \otimes \rho_S$ Thermo-majorization

$E = E_R + E_S$ fixed i.e. $P_E \tau_R \otimes \rho_S P_E$

For each E_S , $g_R(E_R) = g_R(E - E_S) = e^{-\beta E_S} g_R(E)$

eigen values of $P_E \tau_R \otimes \rho_S P_E$ are $\frac{e^{\beta E_S}}{g_R(E)} P(E, g)$

with multiplicity $g_R(E) e^{-\beta E_S}$

Catalytic transformations

Working body, clock, etc.

C.N.O.

$$x \rightarrow y \quad \exists \quad z \text{ s.t. } x \otimes z > y \otimes z$$

“trumping” e.g. Klimesh - Turgut

$$x = \left(\frac{1}{2}, \frac{1}{4}, \frac{1}{4}, 0 \right) \quad y = \left(\frac{4}{10}, \frac{4}{10}, \frac{1}{10}, \frac{1}{10} \right) \quad z = \left(\frac{6}{10}, \frac{4}{10} \right)$$

$$x \succ y \quad x \otimes z > y \otimes z$$

$$D_a(x \parallel \eta) \geq D_a(y \parallel \eta) \quad \forall a \in (-\infty, \infty)$$

$$\eta = \left(\frac{1}{k}, \dots, \frac{1}{k} \right)$$

How cyclic?

$$\sigma_{\text{in}} \otimes \rho_s \rightarrow \sigma_{\text{out}} \otimes \dot{\rho}_s \quad \| \sigma_{\text{in}} - \sigma_{\text{out}} \|_F \leq \epsilon$$

Embezzlement

$$\sigma_{\text{in}} = \sum_i \frac{1}{i} |i\rangle\langle i|$$

$$\| \sigma_{\text{out}} - \sigma_{\text{in}} \|_F \leq \epsilon$$

c.f. van Dam, Hayden (2002)

$$\sigma_{\text{in}} \otimes \frac{\mathbb{I}}{2} \longrightarrow \sigma_{\text{out}} \otimes |0\rangle\langle 0|$$

2nd laws

1) All transformations are possible

$$\|\sigma_{\text{in}} - \sigma_{\text{out}}\|, \leq \epsilon$$

2) Ordinary 2nd law: $F_1(\rho \parallel \tau)$ goes down

$$\|\sigma_{\text{in}} - \sigma_{\text{out}}\|, \leq \epsilon/\log d_{\text{catalyst}}$$

3) $F_a(\rho \parallel \tau)$ must go down $a \geq 0$

Small work distance

$$D_{\text{work}} (\sigma^{\text{in}} > \sigma^{\text{out}}) := kT \inf_{a>0} [F_a(\sigma_{\text{in}} \parallel \tau) - F_a(\sigma_{\text{out}} \parallel \tau)]$$

Resource Theories

Theory	Entanglement	Purity	Asymmetry
Class of Operations	LOCC	Noisy $\frac{\mathbb{I}}{d}, u$	$T(g) \otimes \varepsilon \otimes T^t(g) = \varepsilon$
Free States	Separable σ	\mathbb{I}/d	$\sigma = T(g)\sigma T^t(g)$
Monotones	$\inf_{\sigma} R(\rho \sigma)$	$R(\rho \frac{\mathbb{I}}{d}) = N - S(\rho)$	$\inf_{\sigma} R(\rho \sigma)$

Horodecki, JO (2002, 2012)

“Quantumness”

Thermodynamics

Conclusions

Laws of thermodynamics

- Thermal Operations (U_E, τ, tr)

1st law

0th : τ must be thermal for non-trivial theory

2nd * : $F_a(\rho \parallel \tau)$ must go down

* : embezzlement, how cyclic $\forall a \leq 0$

Doesn't depend on how much control one has!

quasi-classical states: 2nd laws are necessary and sufficient

Many free energies \longrightarrow irreversibility

Limitations due to finite size, quantumness