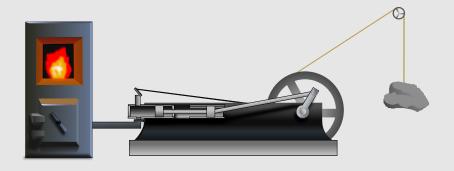
Universal operations in resource theories and local thermodynamics

Henrik Wilming, Rodrigo Gallego, Jens Eisert



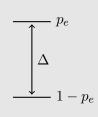
January 16th, 2015

Thermodynamics

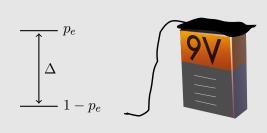


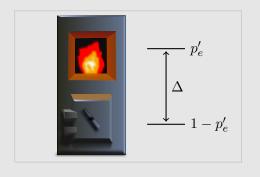




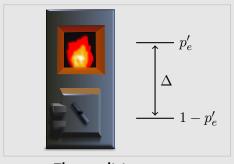








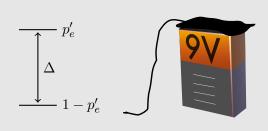




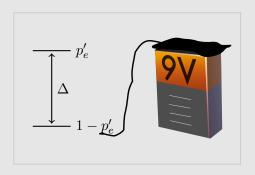


Thermalising map

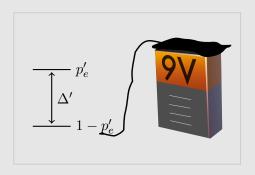




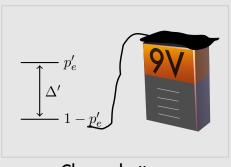












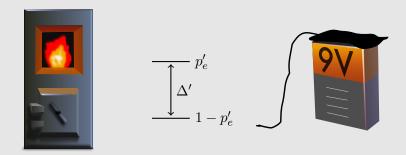
Charge battery

Unitary dynamics of the form

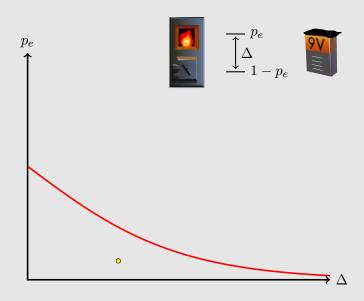
$$(\rho_0, H_0) \mapsto (U_t \rho_0 U_t^{\dagger}, H_t)$$

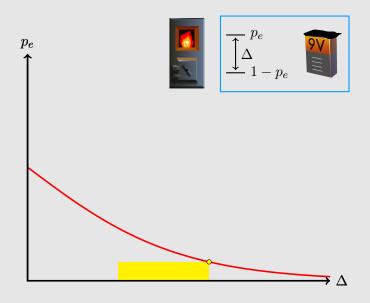
charges the battery by an amount of work

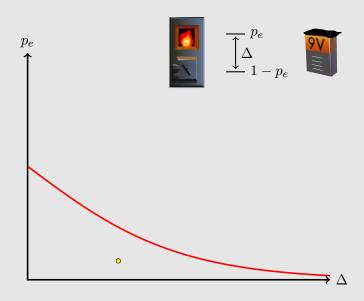
$$\langle W \rangle = \operatorname{Tr} \left(\rho_0 H_0 - U_t \rho_0 U_t^{\dagger} H_t \right).$$

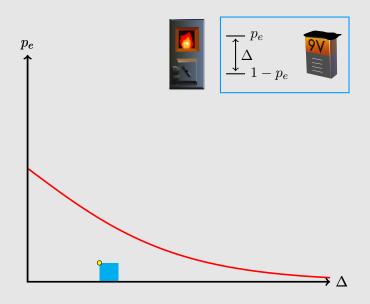


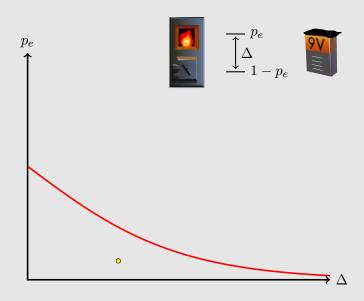
Goal: Charge the battery as much as possible











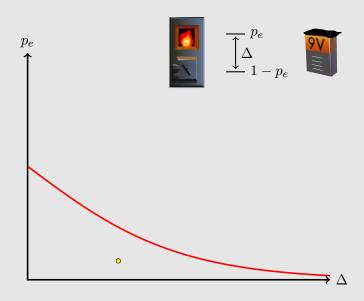


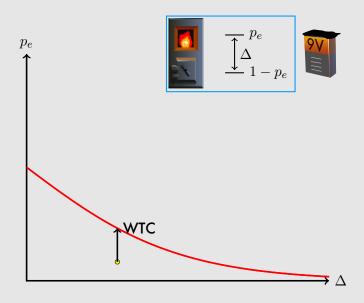
Def.: Weak thermal contact (WTC)

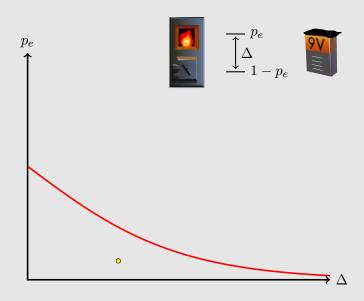
WTC puts the system into thermal equilibrium:

$$(\rho, H) \mapsto (\omega_H, H), \quad \omega_H := \frac{\mathbf{e}^{-\beta H}}{Z_H}.$$











Def.: Thermal Operations (TO) [7]

Let ${\cal H}_B$ be a Hamiltonian (on a bath). A thermal operation is of the form

$$(\rho, H) \mapsto \left(\mathsf{Tr}_B \left(U \rho \otimes \omega_{H_B} U^{\dagger} \right), H \right),$$

where U is any unitary such that $[U, H \otimes \mathbf{1} + \mathbf{1} \otimes H_B] = 0$.



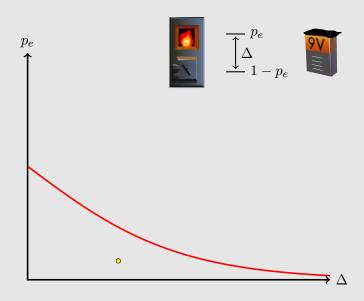
Def.: Gibbs-perserving map (GP-map)

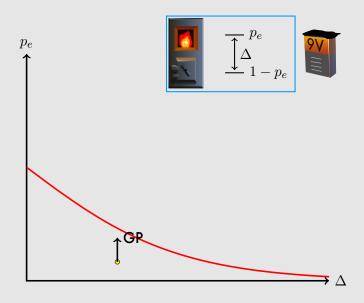
A Gibbs-preserving map is a transformation on pairs that cannot bring the system out of thermal equilibrium:

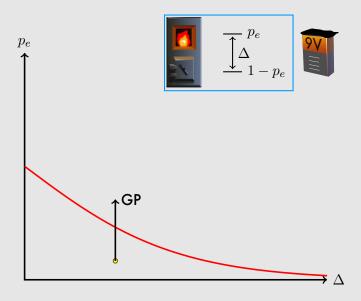
$$(\rho, H) \mapsto (\mathcal{G}_H(\rho), H), \quad \mathcal{G}_H(\omega_H) = \omega_H$$

with a quantum channel \mathcal{G}_H .

(







Thermalising maps



Weak thermal contact (WTC):

$$(\rho, H) \mapsto (\omega_H, H).$$

Thermal operations (TO):

$$(\rho, H) \mapsto \left(\mathsf{Tr}_B \left(U \rho \otimes \omega_{H_B} U^{\dagger} \right), H \right).$$

Gibbs-Preserving maps (GP):

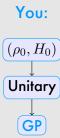
$$(\rho, H) \mapsto ((\mathcal{G}_H(\rho), H), \quad \mathcal{G}_H(\omega_H) = \omega_H.$$

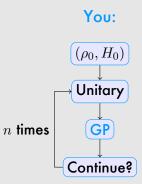
You:

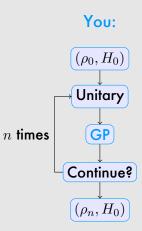
 (ρ_0, H_0)

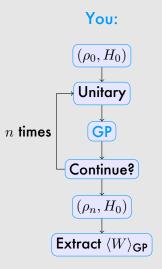


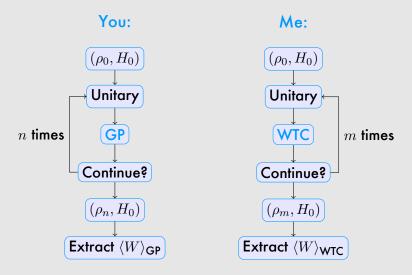




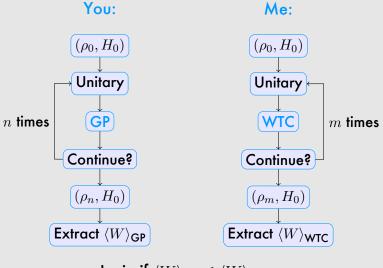






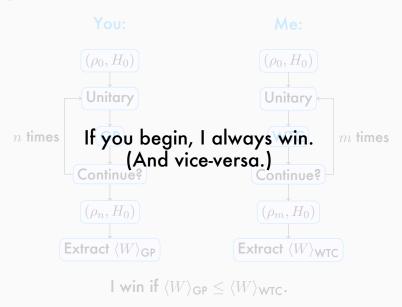


The game of work-extraction



I win if $\langle W \rangle_{\mathsf{GP}} \leq \langle W \rangle_{\mathsf{WTC}}$.

The game of work-extraction



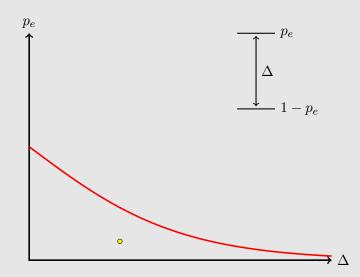
Universality of WTC

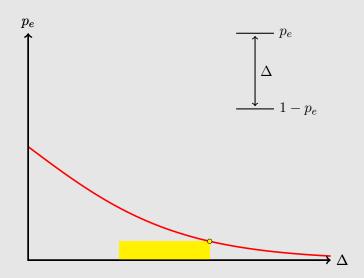
Theorem [2, 4]

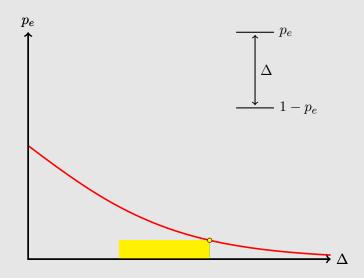
The work yield of a cyclic Hamiltonian process using GP-maps as thermalising maps is bounded as

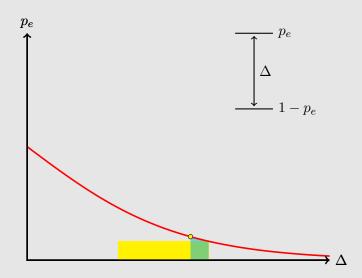
$$\langle W \rangle_{\mathsf{GP}}(\rho_0, H_0) \leq \frac{1}{\beta} S(\rho_0 || \omega_{H_0}).$$

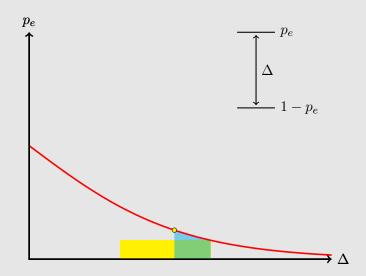
The bound can be saturated arbitrarily well already with WTC as thermalising maps.

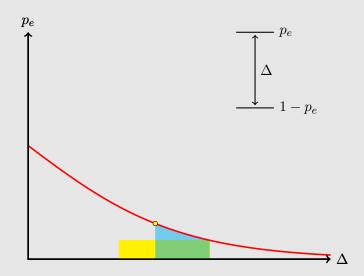


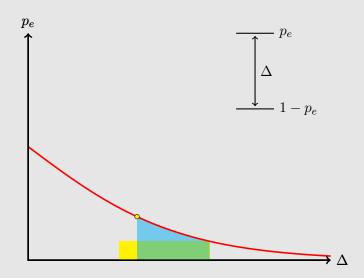


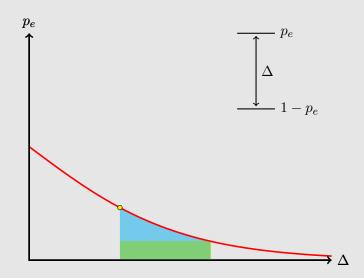


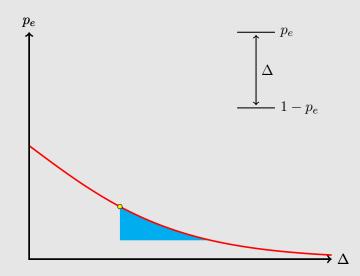


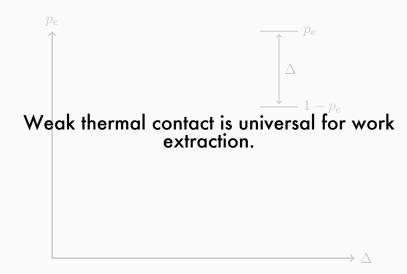


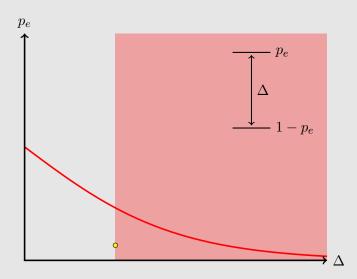


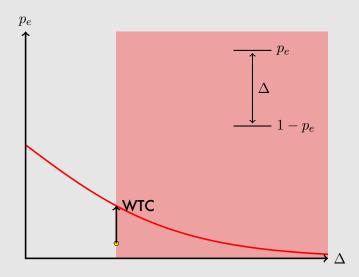


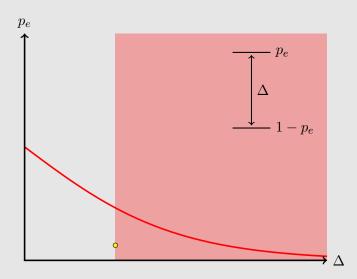


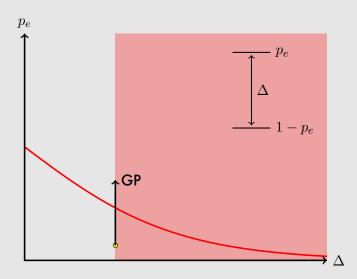


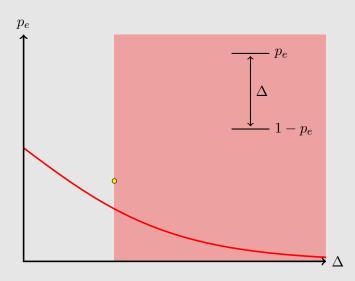


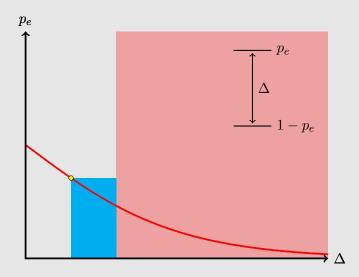


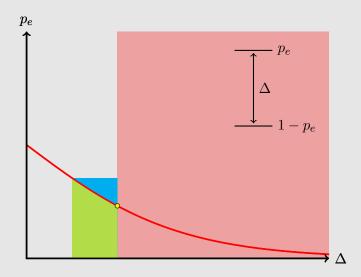


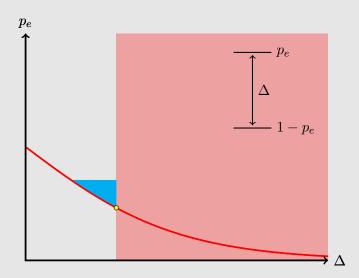


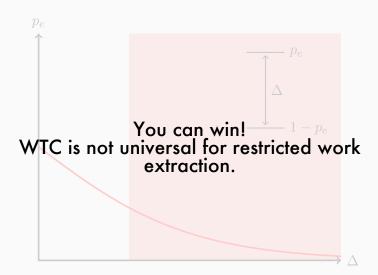












More general restrictions

Given an initial state (ρ_0,H_0) we restrict the possible Hamiltonians to a set $\mathcal{H}(H_0)$.

WTC under restrictions

Theorem (General work bound)

The work that can be extracted from a pair (ρ_0, H_0) using WTC is bounded as

$$\langle W \rangle_{\mathrm{WTC}}^{\mathcal{H}}(\rho_0, H_0) \leq \frac{1}{\beta} S(\rho_0 || \omega_{H_0}) \underset{U \in \mathcal{U}[H_0]}{-\inf} \frac{1}{\beta} S(U \rho_0 U^\dagger || \omega_{H_t}),$$

with $\mathcal{U}[H_0]$ being the unitary group generated by $\mathcal{H}(H_0)$. The bound can be saturated arbitrarily well.

Restriction on locality

$$\mathcal{H}_{\mathsf{loc}}(H_0) := \left\{ H_0 + \sum_i H_i \mid H_i \; \mathsf{local}
ight\}$$

Models the situation of local control over an interacting Hamiltonian.

Restriction on locality

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ight\}$$

Models the situation of local control over an interacting Hamiltonian.

Theorem (Non-universality of WTC)

Under the restriction to \mathcal{H}_{loc} , there exist initial states (ρ_0,H_0) such that no work can be extracted using WTC but work can be extracted using GP-maps.

Proof (Example)

■ Consider two qubits with maximally mixed initial state Ω , Hamiltonian $Z \otimes Z$ and $\beta = 1$. W.l.o.g. assume all Hamiltonians are traceless.

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- General bound implies

$$\langle W \rangle(\Omega, Z \otimes Z) \leq \sup_{\substack{H_t \in \mathcal{H}_{loc}(H_0) \\ U \in \mathcal{U}[H_0]}} \left(S(\Omega||\omega_{Z \otimes Z}) - S(U\Omega U^\dagger||\omega_{H_t}) \right),$$

Proof (Example)

- Consider two qubits with maximally mixed initial state Ω , Hamiltonian $Z \otimes Z$ and $\beta = 1$. W.l.o.g. assume all Hamiltonians are traceless.
- General bound implies

$$\begin{split} \langle W \rangle(\Omega, Z \otimes Z) &\leq \sup_{\substack{H_t \in \mathcal{H}_{loc}(H_0) \\ U \in \mathcal{U}[H_0]}} \left(S(\Omega || \omega_{Z \otimes Z}) - S(U\Omega U^\dagger || \omega_{H_t}) \right), \\ &= \sup_{\substack{H_t \in \mathcal{H}_{loc}(H_0)}} \left(S(\Omega \, || \omega_{Z \otimes Z}) - S(\Omega \, || \omega_{H_t}) \right), \end{split}$$

Proof (Exc $\frac{1}{\beta}S(\rho||\omega_H) = F(\rho,H) - F(\omega_H,H)$ state Ω , all with the free energy $\mathbf{F}(\rho,H) = \mathbf{Tr}(\rho H) - \frac{1}{\beta}S(\rho).$ $\langle W \rangle ($

$$\sup_{H_{t} \in \mathcal{H}_{\text{loc}}(H_{0})} \left(S(\Omega \, || \omega_{Z \otimes Z}) - S(\Omega \, || \omega_{H_{t}}) \right),$$

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$$\begin{split} &= \sup_{H_t \in \mathcal{H}_{\text{loc}}(H_0)} \left(S(\Omega \, || \omega_{Z \otimes Z}) - S(\Omega \, || \omega_{H_t}) \right), \\ &= \sup_{H_t \in \mathcal{H}_{\text{loc}}(H_0)} \left(F(\omega_{H_t}, H_t) - F(\omega_{Z \otimes Z}, Z \otimes Z) \right). \end{split}$$

Proof (Exd Peierls-Bogoliubov inequality: Cons state Ω . $F(\omega_{A+B}, A+B) \leq F(\omega_A, A) + \operatorname{tr}(\omega_A B)$ Hami Ham Gene Any traceless local Hamiltonian is $2U^{\dagger}||\omega_{H_t})\Big)\,,$ orthogonal to $\omega_{Z\otimes Z}$. Hence $F(\omega_{H_{\star}}, H_{t}) \leq F(\omega_{Z \otimes Z}, Z \otimes Z).$ $(F(\omega_{H_t}, H_t) - F(\omega_{Z \otimes Z}, Z \otimes Z))$.

 $H_t \in \mathcal{H}_{loc}(H_0)$

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 $< (F(\omega_{Z \otimes Z}, Z \otimes Z) - F(\omega_{Z \otimes Z}, Z \otimes Z)) \le 0$

Gene

Cons Thermomajorization [5, 6, 7] implies Ham that there exists a thermal operation ${\cal G}$ such that

$$\mathcal{G}(\Omega) = \omega_{Z \otimes Z + tZ \otimes \mathbf{1}},$$

for |t| < 0.46.

$$2U^{\dagger}||\omega_{H_t})\Big),$$

$$=\sup_{Z}\left(S(\Omega\,||\omega_{Z\otimes Z})-S(\Omega\,||\omega_{H_t})\right),$$
 Since $Z\otimes Z+tZ\otimes \mathbf{1}\in\mathcal{H}_{\mathrm{loc}}(Z\otimes Z)$, we can extract

 $\langle W \rangle = S\left(\omega_{Z \otimes Z + tZ \otimes 1} || \omega_{Z \otimes Z}\right) > 0.$

^[5] E. Ruch, R. Schranner, and T. H. Seligman, J. Chem. Phys. 69, 386 (1978).

^[6] H. D. Janzing, P. Wocjan, R. Zeier, R. Geiss, and T. Beth, Int. J. Th. Phys. 39, 2717 (2000).

^[7] M. Horodecki and J. Oppenheim, Nature Comm. 4, 2059 (2013).

Summarizing the effect of restrictions

Restrictions on the Hamiltonians

Passive states for WTC

GP-maps can activate passive states

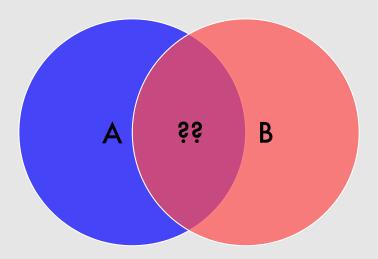
WTC ceases to be universal

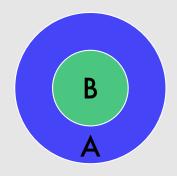
Outlook

Resource theory (roughly)

A class of operations that is closed under composition and contains the identity together with a distinction between free and costly states and operations [9].

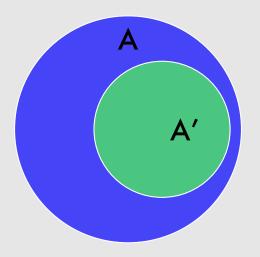
What happens if we "combine" two resource theories?



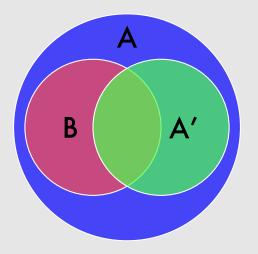


We call a sub-theory $B \subset A$ universal for a specific task if the task can already be achieved optimally only with operations from B.

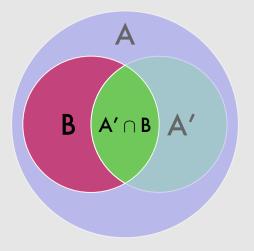
Suppose $A^{\prime}\subset A$ is universal for some task.



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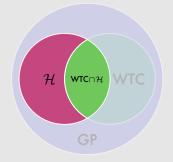
Suppose $A' \subset A$ is universal for some task.



Is $A' \cap B$ still universal in B?



WTC is universal for unrestricted work extraction



WTC is not universal for restricted work extraction



WTC is universal for unrestricted work extraction

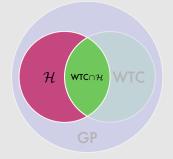


WTC is not universal for restricted work extraction

Open questions



WTC is universal for unrestricted work extraction



WTC is not universal for restricted work extraction

Open questions

Can we find similar situations outside of thermodynamics?



WTC is universal for unrestricted work extraction



WTC is not universal for restricted work extraction

Open questions

- Can we find similar situations outside of thermodynamics?
- Other operationally meaningful types of restrictions?



WTC is universal for unrestricted work extraction



WTC is not universal for restricted work extraction

Open questions

- Can we find similar situations outside of thermodynamics?
- Other operationally meaningful types of restrictions?
- What about single-shot statistical mechanics / different notions of work?

Thanks for listening! For details see arXiv:1411.3754.

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References:

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- [2] J. Aberg, Phys. Rev. Lett. 113, 150402 (2014).
- [3] R. Gallego, A. Riera, and J. Eisert, arXiv:1310.8349 (2013).
- [4] H. Wilming, R. Gallego, and J. Eisert, arXiv:1411.3754 (2014).
- [5] E. Ruch, R. Schranner, and T. H. Seligman, J. Chem. Phys. 69, 386 (1978).
- [6] H. D. Janzing, P. Wocjan, R. Zeier, R. Geiss, and T. Beth, Int. J. Th. Phys. 39, 2717 (2000).
- [7] M. Horodecki and J. Oppenheim, Nature Comm. 4, 2059 (2013).
- [8] P. Faist, J. Oppenheim, and R. Renner, arXiv:1406.3618 (2014).
- [9] B. Coecke, T. Fritz, R. W. Spekkens, arXiv:1409.5531 (2014).