# Măgurele Summer School for Computing in a Rapidly Evolving Society: Parallel Algorithms and Optimizations

One and two-qubit applications on IBM's quantum computers

by Stefan Ataman



Extreme Light Infrastructure - Nuclear Physics (ELI-NP)

July 10, 2025



### The (small) quantum part of this summer school:

### 10/07/2025:

- 9h00 10h30 The future is quantum by Radu Ionicioiu. ✓
- 11h00 12h30 Quantum information: foundations and entanglement by Radu Ionicioiu. ✓
- 14h00 15h30 One and two-qubit quantum gates, quantum circuits & measurement by S. A.  $\checkmark$
- 16h00 17h30 One and two-qubit applications on IBM's quantum computers by S. A.

### 11/07/2025:

- 9h00–10h30 Quantum information: protocols and applications by Radu Ionicioiu.
- 11h00–12h30 Quantum algorithms on IBM's quantum computers (Deutsch, Bernstein-Vazirani, Grover) by S. A.

### The IBM Quantum Platform (as of 2025)

#### Welcome to the upgraded IBM Quantum Platform

The platform provides access to quantum computers, documentation,

and learning resources.

Sign in to start running quantum workloads

Sign in to start running quantum workloads

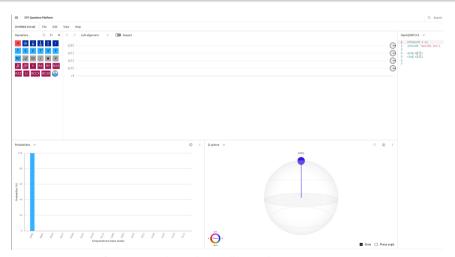
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### Click on Composer.



### The IBM Quantum Platform (as of 2025)



If you see this, then all good, we can start.

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- Single qubit gates and operations
  - Single quantum bit gates
  - IBM's quantum platform
  - The Y and Z gates
  - Measurement of a qubit
- 2 Two qubits
- 3 Entanglement. Maximally entangled states
- Quantum teleportation



### Dirac bra-ket notation - the gubit

In the "normal" z-basis,

$$|0\rangle = \begin{pmatrix} 1\\0 \end{pmatrix}$$
 and  $|1\rangle = \begin{pmatrix} 0\\1 \end{pmatrix}$ 

we have the generic qubit,

$$|\psi\rangle = c_0|0\rangle + c_1|1\rangle$$

with  $c_0, c_1 \in \mathbb{C}$  and for **normalization**:

$$|c_0|^2 + |c_1|^2 = 1$$



### More qubit examples

Let us choose  $c_1 = c_0 = \frac{1}{\sqrt{2}}$ . We have the qubit

$$|\psi\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle = \frac{1}{\sqrt{2}}\begin{pmatrix} 1\\1 \end{pmatrix} = |+\rangle$$

and we will call this "the plus state".

Set  $c_0=rac{1}{\sqrt{2}}$  and  $c_1=-rac{1}{\sqrt{2}}$ , we get

$$|\psi\rangle = \frac{1}{\sqrt{2}}|0\rangle - \frac{1}{\sqrt{2}}|1\rangle = \frac{1}{\sqrt{2}}\begin{pmatrix} 1\\ -1 \end{pmatrix} = |-\rangle$$

and we will call this "the minus state".



# The X gate

The first quantum gate we consider is Pauli's  $\sigma_x$  matrix. We call it simply the X gate:

$$X = \left(\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array}\right).$$

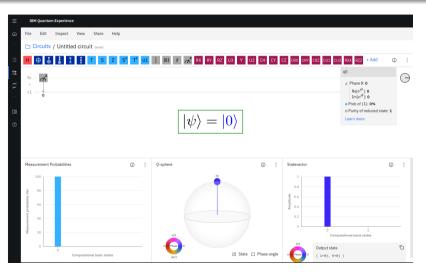
So we had the X gate:

$$\begin{cases} X|0\rangle = |1\rangle \\ X|1\rangle = |0\rangle \end{cases}$$

As a quantum circuit this is:

$$|0\rangle$$
  $X$   $|1\rangle$   $|1\rangle$   $|0\rangle$ 

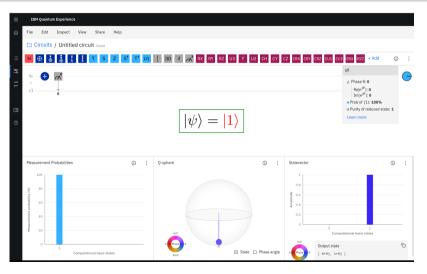
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# |1| on IBM's Quantum Platform



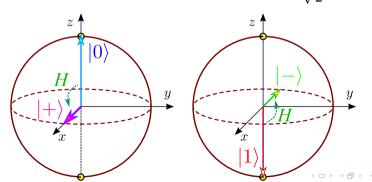
# The Hadamard gate

The Hadamard (H) gate:

$$H = \frac{1}{\sqrt{2}} \left( \begin{array}{cc} 1 & 1 \\ 1 & -1 \end{array} \right)$$

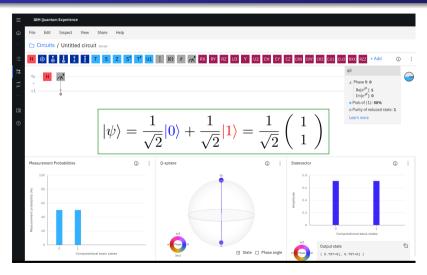
$$|0\rangle$$
  $H$   $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ 

$$|\mathbf{1}\rangle$$
  $H$   $\frac{1}{\sqrt{2}}(|0\rangle - |\mathbf{1}\rangle)$ 



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# $|0\rangle + |1\rangle$ in IBM's Quantum



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# The Y and Z gates

Pauli's  $\sigma_y$  / Y gate si defined by

$$\mathsf{Y} = \left( \begin{array}{cc} 0 & -i \\ i & 0 \end{array} \right)$$

And finally, the last Pauli gate is

$$Z = \left(\begin{array}{cc} 1 & 0 \\ 0 & -1 \end{array}\right)$$

I also recall Pauli's  $\sigma_x$  matrix i.e. the X gate,

$$X = \left(\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array}\right)$$

### The phase gate

Meet the "phase gate" called P in IBM's Quantum Experience,

$$P\left(\stackrel{\bullet}{\phi}\right) = \left(\begin{array}{cc} 1 & 0\\ 0 & e^{i\stackrel{\bullet}{\phi}} \end{array}\right)$$

It is a very handy gate when you want to create a state

$$|\psi'\rangle = c_0|0\rangle + c_1 e^{i\phi}|1\rangle$$

Indeed: start with  $|\psi\rangle = c_0|0\rangle + c_1|1\rangle$ . Apply the phase gate:

$$P\left(\begin{array}{c} \phi \right) \left| \psi \right\rangle = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\phi} \end{pmatrix} \begin{pmatrix} c_0 \\ c_1 \end{pmatrix} = \begin{pmatrix} c_0 \\ c_1 e^{i\phi} \end{pmatrix} = c_0 \left| 0 \right\rangle + c_1 e^{i\phi} \left| 1 \right\rangle$$



So we have

$$P\left(\phi\right)\left|\psi\right\rangle = c_0\left|0\right\rangle + c_1e^{i\phi}\left|1\right\rangle$$

Another way to look at the phase gate P,

$$P\left(\stackrel{\bullet}{\phi}\right) = \left(\begin{array}{cc} 1 & 0\\ 0 & e^{i\stackrel{\bullet}{\phi}} \end{array}\right)$$

is that it has the effect:

$$\begin{cases} |0\rangle \to |0\rangle \\ |1\rangle \to e^{i\phi}|1\rangle \end{cases}$$

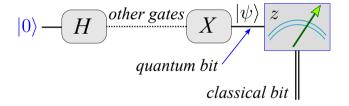
This is quite easy to remember.



# Measurement of a qubit

### Recall Max Born (1926)

Measurement is probabilistic.



#### Obvious fact:

The measurement of a qubit yields a (single) classical bit.



# Measurement of a qubit

### Recall: Max Born (1926)

(Much to Einstein's dislike.) Measurement is probabilistic.

### Specifically, for our qubit:

$$|\psi\rangle = c_0|0\rangle + c_1|1\rangle$$

The probability to find  $|0\rangle$  i. e. "0" after a measurement is  $p_0 = |c_0|^2$ .

The probability to find  $|1\rangle$  i. e. "1" after a measurement is  $p_1 = |c_1|^2$ .

#### Remark:

The total probability is the probability to find  $|0\rangle$  plus the probability to find  $|1\rangle$ . This is  $p_0 + p_1 = |c_0|^2 + |c_1|^2 = 1$ .



### Measurement of a qubit - formalized

### Specifically, for our qubit:

$$|\psi\rangle = c_0|0\rangle + c_1|1\rangle$$

The probability to find  $|0\rangle$  after a measurement is

$$p_0 = |\langle \mathbf{0} | \psi \rangle|^2 = |c_0|^2$$

The probability to find  $|1\rangle$  after a measurement is

$$p_1 = |\langle \mathbf{1} | \psi \rangle|^2 = |c_1|^2$$



# Measurement of a qubit – exercises

#### Consider the state

$$|\psi\rangle = |+\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle$$

What is the probability to find  $|0\rangle$  after a measurement? What is the probability to find  $|1\rangle$  after a measurement?

#### Consider now the state

$$|\psi\rangle = |-\rangle = \frac{1}{\sqrt{2}}|0\rangle - \frac{1}{\sqrt{2}}|1\rangle$$

Again, what is the probability to find  $|0\rangle$  after a measurement? What is the probability to find  $|1\rangle$  after a measurement?

Are these results different? How do you explain?



# Measurement of a qubit – solutions

#### Consider the state

$$|\psi\rangle = |+\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle$$

The probability to find  $|0\rangle$  after a measurement is  $p_0 = |\langle 0|\psi\rangle|^2 = |c_0|^2 = \frac{1}{2}$ . The probability to find  $|1\rangle$  after a measurement is  $p_1 = |\langle 1|\psi\rangle|^2 = |c_1|^2 = \frac{1}{2}$ .

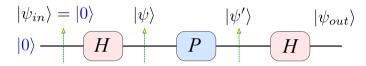
#### Consider now the state

$$|\psi\rangle = |-\rangle = \frac{1}{\sqrt{2}}|0\rangle - \frac{1}{\sqrt{2}}|1\rangle$$

The probability to find  $|0\rangle$  after a measurement is  $p_0 = |\langle 0|\psi\rangle|^2 = |c_0|^2 = \frac{1}{2}$ . The probability to find  $|1\rangle$  after a measurement is  $p_1 = |\langle 1|\psi\rangle|^2 = |c_1|^2 = \frac{1}{2}$ .



So we actually have this

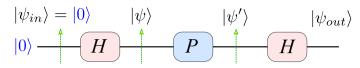


and need to compute  $|\psi_{out}\rangle$ .

The first Hadamard gate does the transformation

$$|\psi\rangle = H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$





Next, the phase gate introduces the phase shift:

$$\begin{cases} |0\rangle \to |0\rangle \\ |1\rangle \to e^{i\phi}|1\rangle \end{cases}$$

so we have the state

$$|\psi'\rangle = P(\phi) |\psi\rangle = \frac{1}{\sqrt{2}} \left( |0\rangle + e^{i\phi} |1\rangle \right)$$

Now apply a new Hadamard gate on  $|\psi'\rangle$  to get the final state

$$|\psi_{out}\rangle = H |\psi'\rangle = \frac{1}{\sqrt{2}} \left( H|0\rangle + e^{i\phi}H|1\rangle \right)$$

We're not done yet, we need  $H|0\rangle$  and  $H|1\rangle$ .

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$$|\psi_{in}\rangle = |0\rangle \qquad |\psi\rangle \qquad |\psi'\rangle \qquad |\psi_{out}\rangle$$

$$|0\rangle \qquad H \qquad P \qquad H$$

But we recall that

$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$
  $H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$ 

therefore

$$|\psi_{out}\rangle = \frac{1}{\sqrt{2}} \left( H|\mathbf{0}\rangle + e^{i\phi}H|\mathbf{1}\rangle \right) = \frac{1}{2} \left( (|\mathbf{0}\rangle + |\mathbf{1}\rangle) + e^{i\phi} \left( |\mathbf{0}\rangle - |\mathbf{1}\rangle \right) \right)$$

so in the end

$$|\psi_{out}\rangle = \frac{1 + e^{i\phi}}{2}|0\rangle + \frac{1 - e^{i\phi}}{2}|1\rangle$$

It can still be simplified. How about to force a global phase  $e^{i\frac{\phi}{2}}$ ? (It will make sense soon.)

$$|\psi_{in}\rangle = |0\rangle \qquad |\psi\rangle \qquad |\psi'\rangle \qquad |\psi_{out}\rangle$$

$$|0\rangle \qquad H \qquad P \qquad H$$

So we *force* a global phase  $e^{i\frac{\phi}{2}}$ , thus

$$|\psi_{out}\rangle = \frac{1 + e^{i\phi}}{2}|0\rangle + \frac{1 - e^{i\phi}}{2}|1\rangle$$

becomes

$$|\psi_{out}\rangle = e^{i\frac{\phi}{2}} \left( \frac{e^{-i\frac{\phi}{2}} + e^{i\frac{\phi}{2}}}{2} |0\rangle + \frac{e^{-i\frac{\phi}{2}} - e^{i\frac{\phi}{2}}}{2} |1\rangle}{2} \right) = e^{i\frac{\phi}{2}} \left( \cos\left(\frac{\phi}{2}\right) |0\rangle - i\sin\left(\frac{\phi}{2}\right) |1\rangle \right)$$

and ignoring the global phase we have

$$|\psi_{out}\rangle = \cos\left(\frac{\phi}{2}\right)|0\rangle - i\sin\left(\frac{\phi}{2}\right)|1\rangle$$

$$|\psi_{in}\rangle = |0\rangle \quad |\psi\rangle \quad |\psi'\rangle \quad |\psi_{out}\rangle \\ |0\rangle \stackrel{\uparrow}{\longrightarrow} H \stackrel{\downarrow}{\longrightarrow} H \qquad |\psi_{out}\rangle = \cos\left(\frac{\phi}{2}\right)|0\rangle - i\sin\left(\frac{\phi}{2}\right)|1\rangle$$

So we actually estimated the phase because the probability to find  $|0\rangle$  after a measurement is

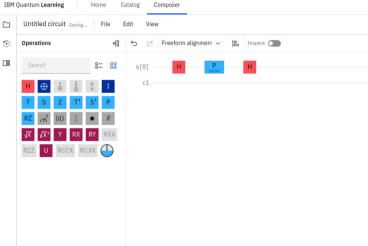
$$p_0 = \cos^2\left(\frac{\phi}{2}\right)$$

and similarly, the probability to find  $|1\rangle$  after a measurement is

$$p_1 = \sin^2\left(\frac{\phi}{2}\right)$$

# Phase measurement in IBM QE

### Open IBM Quantum and implement the following quantum circuit:



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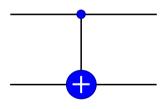
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### Two qubits. Two qubit gates

#### The CNOT gate

This is probably the most important gate we will use. It is called the controlled-NOT or simply CNOT gate. It has a **control qubit** (upper one) and a **target qubit** (the lower one):



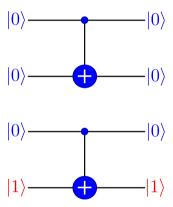
#### Remark:

The control qubit as it name says, controls. It is not affected by the CNOT gate.



# The CNOT gate

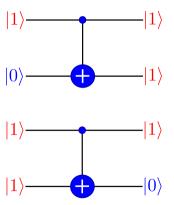
Set the control qubit to  $|0\rangle$ . Not much happens:



Boring.

# The CNOT gate

Set the control qubit to  $|1\rangle$ . Now interesting things happen:



Question: what gate has the same behaviour on the target bit?



# The CNOT gate - matrix representation

Being a two-qubit gate it must have a  $4 \times 4$  matrix representation. Here it comes:

$$\mathsf{CNOT} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 & | & 0 & 0 \\ 0 & 1 & | & 0 & 0 \\ \hline 0 & 0 & | & 0 & 1 \\ 0 & 0 & | & 1 & 0 \end{pmatrix} = \begin{pmatrix} \boxed{\mathbb{I}_2 \mid \mathbf{0}_2} \\ \boxed{\mathbf{0}_2 \mid X} \end{pmatrix}$$

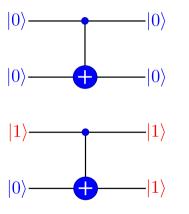
I don't suggest to use it, except in the beginning. I would go for the truth table instead:

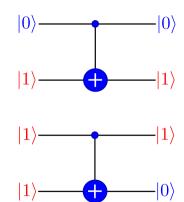
$$\begin{cases} |00\rangle \rightarrow |00\rangle \\ |01\rangle \rightarrow |01\rangle \\ |10\rangle \rightarrow |11\rangle \\ |11\rangle \rightarrow |10\rangle \end{cases}$$

Remark:

You can use the matrix representation of the gubits to check it.

# The CNOT gate in a single slide





### The control-Z gate

We saw the CNOT gate:

$$\mathsf{CNOT} = \left(\begin{array}{c|c|c} \mathbb{I}_2 & \mathbf{0}_2 \\ \hline \mathbf{0}_2 & X \end{array}\right) = \left(\begin{array}{cccc} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{array}\right) = \left(\begin{array}{ccccc} 1 & 0 & | & 0 & 0 \\ \hline 0 & 1 & | & 0 & 0 \\ \hline 0 & 0 & | & 0 & 1 \\ 0 & 0 & | & 1 & 0 \end{array}\right)$$

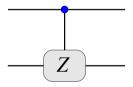
We can maybe imagine a gate

$$\mathsf{CZ} = \left(\begin{array}{c|ccc} \mathbb{I}_2 & \mathbf{0}_2 \\ \hline \mathbf{0}_2 & Z \end{array}\right) = \left(\begin{array}{cccc} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{array}\right) = \left(\begin{array}{ccccc} 1 & 0 & | & 0 & 0 \\ \hline 0 & 1 & | & 0 & 0 \\ \hline 0 & 0 & | & 1 & 0 \\ \hline 0 & 0 & | & 0 & -1 \end{array}\right)$$

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# The control-Z gate

### Meet the CZ gate:



#### Truth table:

$$\begin{cases} |00\rangle \to |00\rangle \\ |01\rangle \to |01\rangle \\ |10\rangle \to |10\rangle \\ |11\rangle \to -|11\rangle \end{cases}$$

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  - The Bell states. Generating/measuring the Bell states
  - Bell state measurement
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# The maximally entangled state - examples

The EPR (also called Bell) state:

$$|\Psi\rangle = \frac{1}{\sqrt{2}} \left( | \blacksquare_{Alice} \blacksquare_{Bob} \rangle + | \blacksquare_{Alice} \blacksquare_{Bob} \rangle \right)$$

can be made of two photons having horizontal/vertical (H/V) polarization

$$|\Psi\rangle = \frac{1}{\sqrt{2}}\left(|\textbf{\textit{H}}_{Alice}\textbf{\textit{H}}_{Bob}\rangle + |V_{Alice}V_{Bob}\rangle\right) = \frac{1}{\sqrt{2}}\left(|\textbf{\textit{H}}\textbf{\textit{H}}\rangle + |VV\rangle\right)$$

or, in quantum computing, two maximally entangles qubits

$$|\Psi\rangle = \frac{1}{\sqrt{2}}\left(|\textcolor{red}{0}_{Alice}\textcolor{blue}{0}_{Bob}\rangle + |\textcolor{blue}{1}_{Alice}\textcolor{blue}{1}_{Bob}\rangle\right) = \frac{1}{\sqrt{2}}\left(|\textcolor{blue}{00}\rangle + |\textcolor{blue}{11}\rangle\right)$$

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#### The Bell states

#### There are actually 4 maximally entangled states

$$|\Phi^{+}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$
$$|\Phi^{-}\rangle = \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$$
$$|\Psi^{+}\rangle = \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle)$$
$$|\Psi^{-}\rangle = \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$$

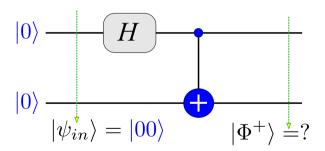
They are called the Bell states. They are all orthogonal among themselves.



## Back to the previous question

You have the quantum circuit below.

#### What is the output state?

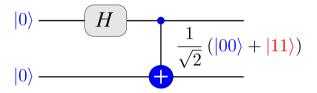


Any guesses?

### The Bell states - $|\Phi^+\rangle$

#### I claim that this scheme

generates the  $|\Phi^+\rangle$  state. Do you agree?



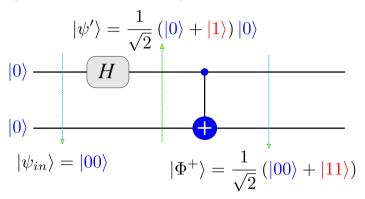
#### We know the Hadamard gate by now. Recall the CNOT gate:

$$\begin{cases} \begin{array}{c} |00\rangle \rightarrow |00\rangle \\ |01\rangle \rightarrow |01\rangle \\ |10\rangle \rightarrow |11\rangle \\ |11\rangle \rightarrow |10\rangle \end{array} \end{cases}$$

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## The Bell state generation

Let's start describing the circuit, from left to right:



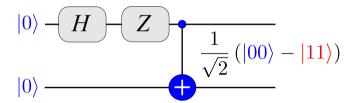
#### Everyone agrees? Did you get the same result?

Let's implement it in IBM Quantum Learning!

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#### The Bell states

I claim that this scheme generates the  $|\Phi^-\rangle$  state.



## The Bell state generation - $|\Phi^-\rangle$

Let's see:

$$|\psi'\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) |0\rangle \qquad |\Phi^{-}\rangle = \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$$

$$|0\rangle \qquad H \qquad Z$$

$$|\psi_{in}\rangle = |00\rangle \qquad |\psi''\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) |0\rangle$$

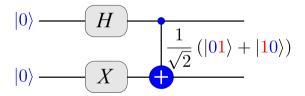
Let's implement it in IBM QE!

4D > 4A > 4B > 4B > B 990

## The Bell states - $|\Psi^+\rangle$

#### I claim that this scheme

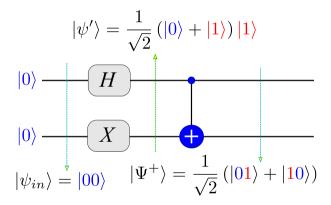
generates the  $|\Psi^+\rangle$  state. Do you agree?



Can you prove this result? Give it a try!

## The Bell state generation - $|\Psi^+ angle$

Let's see:



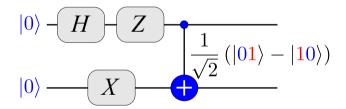
Implement it in IBM QE! Go!



## The Bell states - $|\Psi^-\rangle$

#### I claim that this scheme

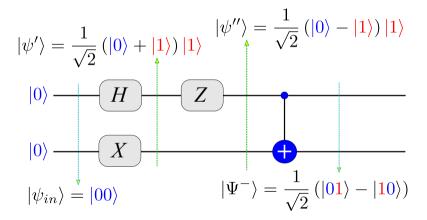
generates the  $|\Psi^-\rangle$  state. Do you agree?



Can you prove this result? Give it a try!

## The Bell state generation - $|\Psi^-\rangle$

Let's see:

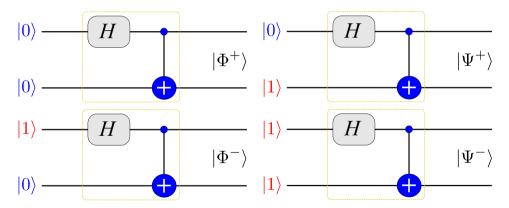


Implement it in IBM QE! Go!

4D > 4A > 4B > 4B > B 990

## The Bell state generation - homework

#### Show that:



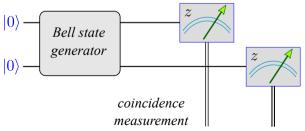
#### The Bell state measurement

#### Very subtle point

All Bell states are entangled i. e. not separable. Single measurements on one qubit will yield totally random values.

#### Measure both qubits!

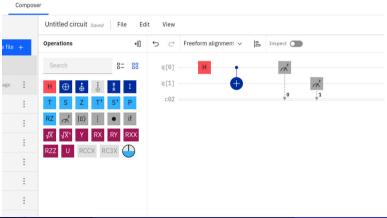
We need coincidence measurements.



### The Bell state measurement in IBM's Quantum simulator

Implementing this measurement scheme in IBM QE.

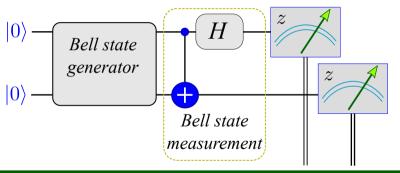
Let's test all four Bell states.



#### The Bell state measurement

#### Claim:

The circuit below distinguishes among the 4 Bell states.

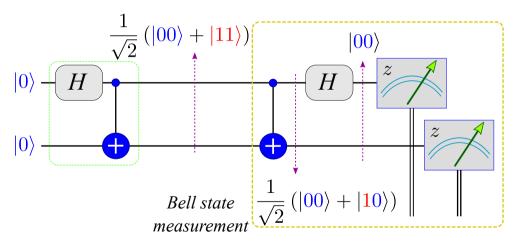


#### Can you prove it?

Start with the  $|\Phi^+\rangle$  state!

#### The Bell state measurement

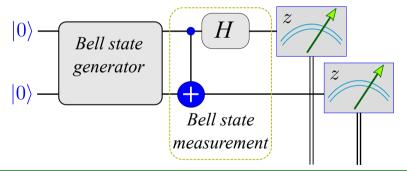
The  $|\Phi^+\rangle$  state. Here we go:



### Bell state measurement in IBM QE

#### Use IBM QE:

Implement all for Bell states and use the circuit below to measure them in IBM QE.

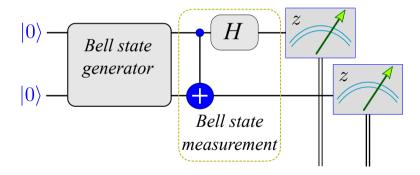


#### Can you distinguish now among the 4 Bell states?

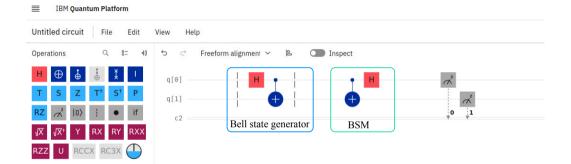
Are your results probabilistic or deterministic? Comment on your findings!

## IBM Quantum implementation

Show that the circuit below distinguishes among the  $4\ {\sf Bell}$  states.



## IBM Quantum implementation – solution



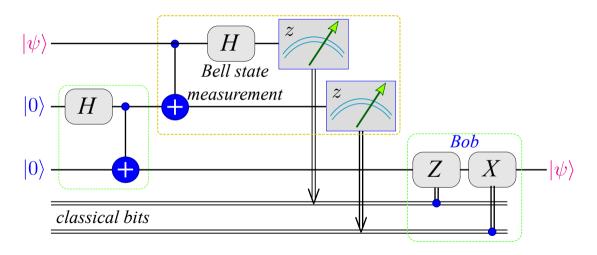


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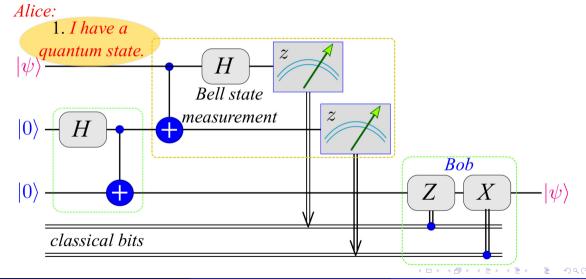
- Single qubit gates and operations
- 2 Two qubits
- 3 Entanglement. Maximally entangled states
- Quantum teleportation



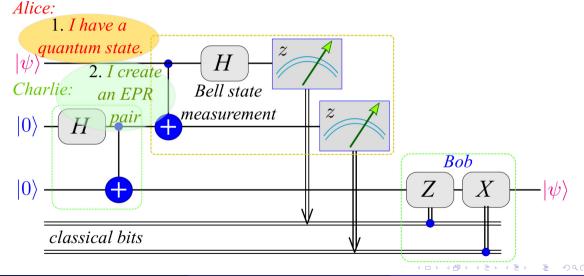
## The Quantum teleportation



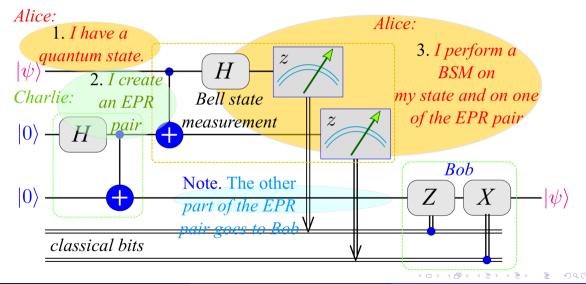
## The Quantum teleportation – 1. Alice has a $|\psi\rangle$ .



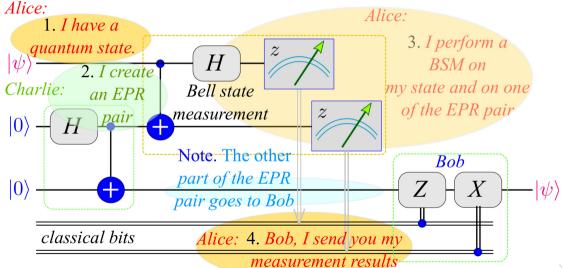
## The Quantum teleportation – 2. Charlie creates $|\Phi^+\rangle$ .



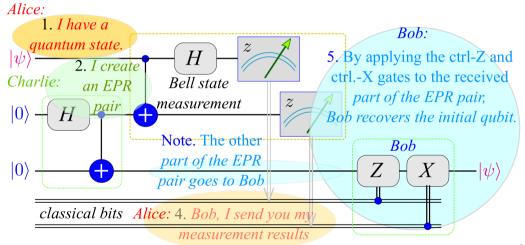
## The Quantum teleportation -3. Alice performs a BSM.



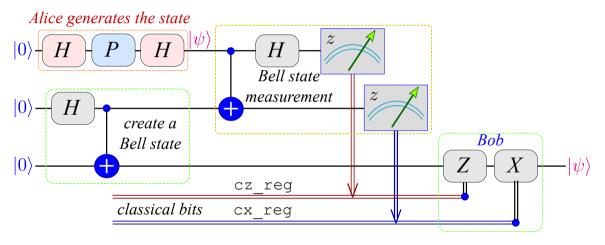
## The Quantum teleportation – 4. Alice sends her measurements to Bob.



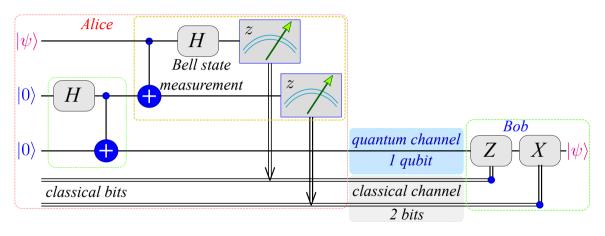
# The Quantum teleportation – 5. Bob uses the quantum plus classical channels.



## The Quantum teleportation – more practical implementation for us



## Quantum teleportation - the channels

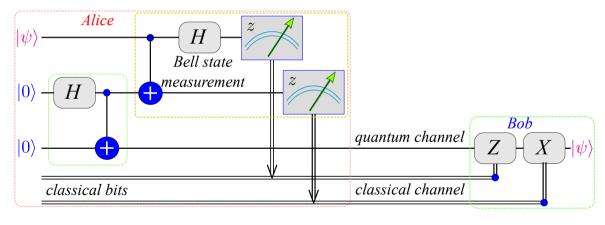


The "quantum channel" can be e. g. a fiber.

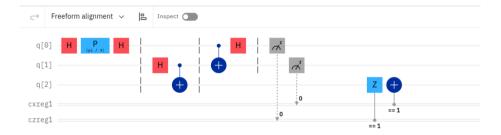
The "classical channel" can be your usual phone line.



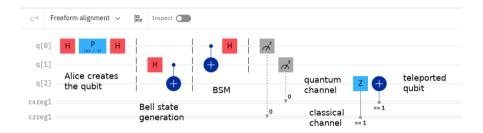
## Quantum teleportation - the channels



## Quantum teleportation in IBM Quantum



## Quantum teleportation in IBM QE

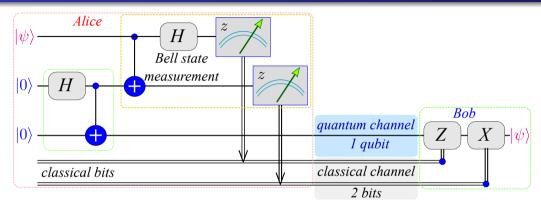


#### Now implement yourself the scheme

#### Good luck!

Hint: for control-Z and control-X (classical control) put first the quantum gate and the drag then the if from the left panel on top. Then configure the control Value to 1.

## Quantum teleportation - the big picture



#### OK, but why does it work?

We implemented it on the IBM Quantum platform. If you did it correctly, it should work. Well done!

## That's it for today.

## Thank you for your attention!

Questions are welcome.

