

Let's zoom the picture & do a detail study; the geometry will look as follows.

$$AB = vt$$

$$BC = b$$

$$AC = r$$

$$F_{\perp} = F \sin \theta$$

$$\sin \theta = \frac{BC}{AC}$$

$$\Rightarrow \sin \theta = \frac{b}{r}$$

$$\Rightarrow \left[r = \frac{b}{\sin \theta} \right]$$

$$\tan \theta = \frac{BC}{AB} \Rightarrow \tan \theta = \frac{b}{vt} \Rightarrow t = \left(\frac{b}{v} \right) \cot \theta \Rightarrow dt = \left(\frac{b}{v} \right) \operatorname{cosec}^2 \theta d\theta$$

Now we can write the equation as

$$I = \Delta P = \int F_{\perp} dt$$

$$\Rightarrow I = \Delta P = \int F \sin \theta dt$$

$$= \int F \sin \theta \left(\frac{b}{v} \right) \operatorname{cosec}^2 \theta d\theta$$

$$= \int_0^{\pi} \left(\frac{Ze \cdot e}{4\pi\epsilon_0 r^2} \right) \sin \theta \left(\frac{b}{v} \right) \operatorname{cosec}^2 \theta d\theta$$

$$= \int_0^{\pi} \frac{Ze^2}{4\pi\epsilon_0 \left(\frac{b}{\sin \theta} \right)^2} \sin \theta \left(\frac{b}{v} \right) \operatorname{cosec}^2 \theta d\theta$$

$$= \int_0^{\pi} \frac{Ze^2}{4\pi\epsilon_0 b v} \sin \theta d\theta = \frac{Ze^2}{4\pi\epsilon_0 b v} \int_0^{\pi} \sin \theta d\theta$$

$$= \frac{Ze^2}{4\pi\epsilon_0} \frac{2}{bv}$$

$$\boxed{I = \Delta P = \frac{Ze^2}{2\pi\epsilon_0 bv}}$$

we need to calculate the kinetic energy transfer to electron mass m_e .

The v is taken outside the integral, as the change in v during the collision with a single electron is very small.

What is the total energy transfer? This can be approximated to the using Bohr's classical approach:

The energy transfer is $\Delta E = \frac{(\Delta P)^2}{2m_e}$ ~~m_e in denomi~~

m_e in denominator, which is really small than the incoming mass of the incoming particle (if it is not the electron), thus it gives a negligible contribution

$$\Rightarrow \Delta E = \frac{(\Delta P)^2}{2m_e}$$

ΔE is b dependent.

$$\Rightarrow \Delta E(b) = \left(\frac{1}{4\pi\epsilon_0} \right)^2 \frac{2Z^2 e^4}{m_e v^2 b^2}$$

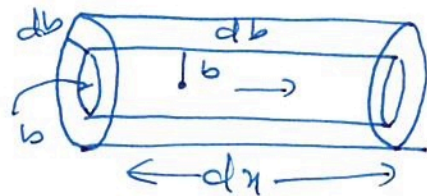
m_e is electron mass which is taking the KE from this process.

Note that this quantity depends on Z^2/m_e

$$\Rightarrow \boxed{\Delta E(b) \propto \frac{Z^2}{m_e}}$$

To find out the effect of all electrons with the given impact parameter b , we construct a cylindrical shell of electrons of radii ~~to~~ b & $b+db$ & the length dx along the path of the incoming particle.

considering the volume element $dv = 2\pi b db dx$



and the electron number density n_e , we can write the energy transfer as

$$dE(b) = \Delta E(b) n_e dv$$

This is the energy transfer to the electrons within the cylinder. This quantity.

$$\Rightarrow dE(b) = \frac{1}{(4\pi\epsilon_0)^2} \frac{4\pi z^2 e^4}{m_e v^2} n_e \frac{db}{b} d\eta \quad - \textcircled{A}$$

Thus the energy loss per unit path length of the particle due to the electrons at impact parameters between b & $b+db$ is given by above equation ~~(A)~~

$$\frac{dE(b)}{d\eta} = \frac{1}{(4\pi\epsilon_0)^2} \frac{4\pi z^2 e^4}{m_e v^2} n_e \frac{db}{b}$$

$\frac{dE}{d\eta}$ is often negative to symbolize the energy loss in material due to the passage through the material. Now the total rate of energy loss per unit path length can be calculated by integrating right hand side from b_{min} to b_{max} .

$$\Rightarrow -\frac{dE}{d\eta} = \frac{1}{(4\pi\epsilon_0)^2} \frac{4\pi z^2 e^4}{m_e v^2} n_e \int_{b_{min}}^{b_{max}} \frac{db}{b}$$

b_{min} & b_{max} are the limits for which the equation ~~for~~ ~~energy~~ ~~loss~~ is valid.

$$\Rightarrow \boxed{-\frac{dE}{d\eta} = \frac{1}{(4\pi\epsilon_0)^2} \frac{4\pi z^2 e^4}{m_e v^2} n_e \ln \frac{b_{max}}{b_{min}}} \quad - \textcircled{B}$$

$-\frac{dE}{d\eta}$ also called as stopping power

$$\Rightarrow \boxed{-\frac{dE}{d\eta} = \frac{z^2 e^4 n_e}{4\pi\epsilon_0^2 m_e v^2} \ln \left(\frac{b_{max}}{b_{min}} \right)} \quad - \textcircled{C}$$

can we put $b_{\min} = 0$ & $b_{\max} = \infty$; Think practically this will lead to a situation where the energy loss per unit path length is infinity, which is an unnatural case as incoming particle has finite energy. Thus we must avoid this situation.

Thus b_{\min} is b for which $\Delta E(b)$ has its maximum possible value, i.e. maximum possible energy transfer which occurs for a head on collision. i.e.

$$\frac{1}{2} m_e (2v)^2 = 2 m_e v^2 \quad \text{classically}$$

the relativistic approximation is $2 \gamma^2 m_e v^2$.

where $\gamma = (1 - \beta^2)^{-1/2}$ $\beta = v/c$.

$$\Delta E_{\max} = \frac{2 z^2 e^4}{m_e v^2 b^2} \left(\frac{1}{4\pi\epsilon_0} \right)^2 = 2 \gamma^2 m_e v^2$$

$$\Rightarrow b_{\min} = \frac{Ze^2}{\gamma m_e v^2} \left(\frac{1}{4\pi\epsilon_0} \right)^2$$

b_{\max} : The interaction time $\frac{b}{\sqrt{v}}$ should be less than the free mean orbital period of atom's electron $\frac{1}{\bar{\nu}}$ in order for electron to absorb energy "adiabatic invariance": $\bar{\nu}$ is the characteristic orbital frequency.

$$\Rightarrow \frac{b_{\max}}{\sqrt{v}} = \frac{1}{\bar{\nu}} \Rightarrow b_{\max} = \frac{\sqrt{v}}{\bar{\nu}}$$

This is so as electrons are not free. $\left\{ \begin{array}{l} \bar{\nu} \text{ is the} \\ \text{characteristic} \\ \text{orbital frequency} \end{array} \right.$

Further we can calculate the number of electrons n_e in terms of density of material as

$$n_e = \frac{\rho N_A Z}{A}$$

ρ = Density of the absorber material

N_A = Avogadro number to know how many atoms are there

Z, A = Atomic and mass number of atoms of material

Putting all these quantities b_{min}, b_{max}, n_e etc in equation (B) of

$$\begin{aligned} -\frac{dx}{dt} & \Rightarrow -\frac{dE}{dx} = \left(\frac{1}{4\pi\epsilon_0}\right)^2 \frac{4\pi Z^2 e^4}{m_e v^2} n_e \ln\left(\frac{b_{max}}{b_{min}}\right) \\ & = 4\pi \left(\frac{e^2}{4\pi\epsilon_0 m_e c^2}\right)^2 m_e Z^2 c^2 \left(\frac{c^2}{v^2}\right) \rho \frac{N_A Z}{A} \ln\left(\frac{b_{max}}{b_{min}}\right) \\ & = 4\pi N_A \pi e^2 m_e c^2 \frac{Z}{A} \frac{\rho}{\beta^2} Z \ln\left(\frac{rv}{\bar{v}} / \left(\frac{ze^2}{\sqrt{m_e} v^2} \frac{1}{4\pi\epsilon_0}\right)\right) \\ \Rightarrow \boxed{-\frac{dE}{dx} = 4\pi N_A \pi e^2 m_e c^2 \frac{Z^2}{A} \frac{\rho}{\beta^2} \ln\left(\frac{4\pi\epsilon_0 r^2 m_e v^2}{Ze^2 \bar{v}}\right)} \end{aligned}$$

This is the Bohr's classical formula for the energy loss for heavy charge particle.

In quantum mechanical ~~corrected~~ treatment, the inelastic collision with atomic electrons is divided into two classes hard & soft. In former, a large energy is transferred and the quantum mechanical spin and exchange effects enter into picture. In latter type of collisions, the energy transfer extends from some arbitrary energy to the minimum energy. In 1930-33 Bethe & Block deduced an approximation using quantum-mechanics treatment & Born approximation. More complete formula is:

$$-\frac{dE}{dx} = 2\pi N_A \pi e^2 m_e c^2 \rho \frac{Z}{A} \frac{Z^2}{\beta^2} \left[\ln\left(\frac{2m_e r^2 \omega_{max}}{I^2}\right) - 2\beta^2 - \delta - 2\frac{C}{Z} \right]$$
$$\text{where } \omega_{max} = \frac{2m_e (c\beta v)^2}{1 + \left(\frac{m_e}{M}\right) \sqrt{1 + (\beta v)^2 + (m_e/M)^2}} \approx 2m_e (c\beta v)^2$$

I is ionization potential.

Let's discuss about each quantities in the above equation

r_e : classical electron radius

m_e : mass of electron

N_A : Avogadro's number $6.022 \times 10^{23}/\text{mol}$

c : speed of light.

$$2\pi N_A r_e^2 m_e c^2 = 0.1535 \text{ MeV} \cdot \text{cm}^2/\text{g}.$$

Z : atomic number

A : Mass number

ρ : Density of absorbing material

z : charge of the incident particle

$$\gamma: \frac{1}{\sqrt{1-\beta^2}} : \beta = \left(\frac{v}{c}\right)$$

I : The mean excitation potential: it is determined empirically for different material from the measurement of dE/dx . This is a main parameter of the Bethe-Bloch formula and is essentially the average orbital frequency $\bar{\nu}$ from the Bohr's formula times the Planck constant, $h\bar{\nu}$. It is theoretically a logarithmic average of ν weighted by the so called oscillator strength of the atomic levels.

From the empirical estimation

$$I = h\bar{\nu} \Rightarrow b_{\max} = \frac{\sqrt{I}}{\sqrt{v}} \Rightarrow b_{\max} = \frac{\sqrt{I} h}{I}$$

$$\frac{I}{Z} = 12 + \frac{7}{Z} \text{ eV} \quad Z < 13$$

$$= 9.76 + 58.8 Z^{-1.19} \text{ eV} \quad Z \geq 13$$

δ : Density correction; the parameter which describes how transverse electric field of incident particle is screened by the charge density of the electrons in the medium. This is kind of polarization effects of atom in the medium caused by E field of the particle, more distance atom are shielded from the full E field

intensity - contribute less to energy loss.

e : it is the shell correction for the case where the velocity of the incident particle is comparable (or less) to the orbital velocity of the bound electron ($\beta \approx Z\alpha$). It has very small contribution as compared to the other effect.

W_{\max} : Maximum energy transfer in one collision

$$W_{\max} = \frac{2 m_e c^2 \beta^2}{1 + (m_e/M) \sqrt{1 + \beta^2} + (m_e/M)^2}$$

also written as

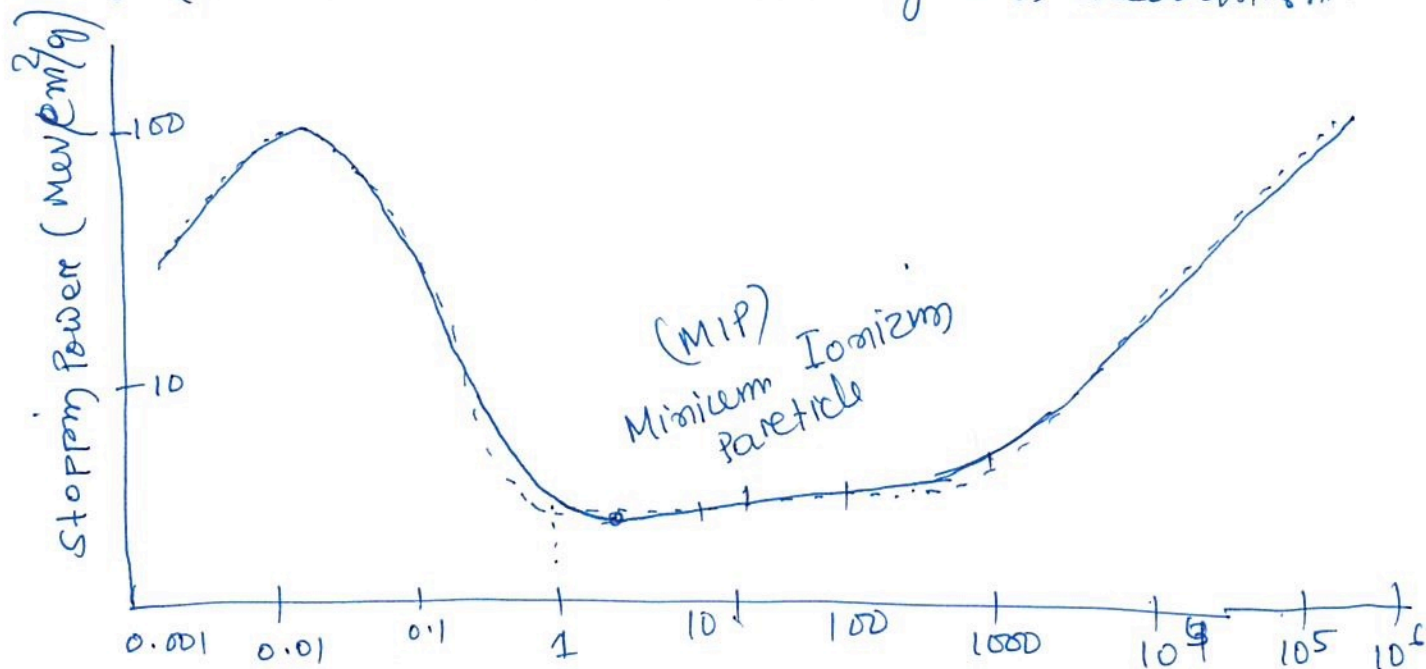
$$= \frac{2 m_e c^2 \eta^2}{1 + 2s \sqrt{1 + \eta^2 + s^2}} \quad \begin{aligned} s &= m_e/M \\ \eta &= \beta \gamma \end{aligned}$$

$$\text{if } M \gg m_e \Rightarrow W_{\max} \approx 2 m_e c^2 \eta^2$$

- The ionization energy loss does not depend on the mass of ionizing particle
- There is only a weak dependence on the medium because $\frac{Z}{A} \approx 0.5$ for most materials.
- The energy loss depends only on β and one needs the mass of incoming particle to convert it to momentum

Re B

let's look at the famous muon energy loss mechanism.



- The Bethe-Bloch equation dictates that the energy loss increases as $\sim \beta^{-2}$ as the particle speed decreases. This is shown in the muon energy loss figure as above. This means that the slow particles will be more ionizing than fast particles.
- As particle speed increases the energy loss reaches a minimum of about $1.5 - 2.0 \text{ MeV cm}^2/\text{g}$ remains at this level for almost the entire range of muon momenta one would observe in a high energy physics experiment. Hence the concept of minimum ionizing particle, which is used EHEP to refer to particles whose speed is in this regime and suffer similar energy losses ($1.5 - 2.0 \text{ MeV cm}^2/\text{g}$).
- After that it starts rising following the so called relativistic rise. However radiative effects are also important in this region.
- The relativistic rise comes from the fact that the electric field of the ionizing particle in the lab frame is important & proportional to the relativistic gamma. Hence the faster the particle is, the stronger the field becomes and therefore the particle can ionize atoms at larger distances and loses more energy. This rise of the energy loss is logarithmic as can be seen by approximating the $\frac{dE}{dx}$ equation.

$$-\frac{dE}{dx} \propto \ln\left(\frac{2me^2\beta^2\gamma^2}{I}\right) \sim \ln\frac{2me^2\omega_m}{I} + \ln(\beta^2\gamma^2) \\ \sim \ln\left(\frac{p}{m}\right)$$

Eventually the medium polarizes and cancels this effect. This prevents the energy loss from rising perpetually and eventually flattens at very high particle speeds. This is the origin of density terms, δ -correction, in the Bethe-Bloch equation.

It is a fact that the energy loss dE/dx is given in terms of $\text{MeV cm}^2/\text{g}$ in unit. It really depends very weakly on the type of medium. However to convert it to MeV/cm one has to take in to account how dense the material is as the Bethe-Bloch equation is given normalised to unit density.

There is a short notation of Bethe-Bloch equation you may encounter in some reference.

$$-\frac{dE}{dx} = K Z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \left(\frac{2 m_e \beta^2 \gamma^2 W_{\max}}{I^2} \right) - \beta^2 - \frac{\delta}{2} - \frac{C}{2Z} \right]$$

$$\frac{K}{A} = 4\pi N_A r_e^2 m_e c^2 / A \approx 0.307075 \frac{\text{MeV cm}^2}{\text{g}}$$

$$\Rightarrow -\frac{1}{f} \frac{dE}{dx} = K Z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \left(\frac{2 m_e \beta^2 \gamma^2 W_{\max}}{I^2} \right) - \beta^2 - \frac{\delta}{2} - \frac{C}{2Z} \right]$$

$\int dx = X \Rightarrow$ Radiation length. [This can be uniformly found out while working for electron i.e. X_0]

~~Energy loss of electrons:~~

The ionization energy loss is always given in terms of radiation length which is basically the $\left(\frac{dE}{dx}\right)$ per unit density thus the radiation length is given expressed in g/cm^2 unit.