

1.

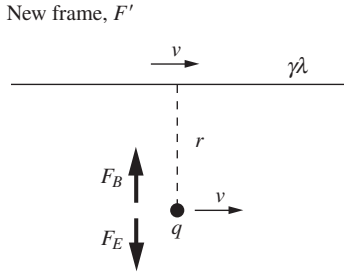


Figure 6.29.
The electric and magnetic forces in the new frame.

- (a) The force on a particle is always largest in the rest frame of the particle. It is smaller in any other frame by the γ factor associated with the speed v of the particle. The force in the particle frame (the lab frame) is $q\lambda/2\pi\epsilon_0 r$, so the force in the new frame is $q\lambda/2\gamma\pi\epsilon_0 r$.
- (b) In the new frame (call it F'), the linear charge density in the rod is increased to $\gamma\lambda$, due to length contraction. So the electric field is $E' = \gamma\lambda/2\pi\epsilon_0 r$. This field produces a repulsive electric force of $F_E = \gamma q\lambda/2\pi\epsilon_0 r$.

In F' the current produced by the rod is the density times the speed, so $I = (\gamma\lambda)v$. The magnetic field is then $B' = \mu_0 I/2\pi r = \mu_0 \gamma\lambda v/2\pi r$, directed into the page in Fig. 6.29 (assuming λ is positive). The magnetic force is therefore attractive and has magnitude (using $\mu_0 = 1/\epsilon_0 c^2$)

$$F_B = qvB' = qv \cdot \frac{\mu_0 \gamma \lambda v}{2\pi r} = \frac{\gamma q \lambda v^2}{2\pi \epsilon_0 r c^2}. \quad (6.78)$$

The net repulsive force acting on the charge q in the new frame is therefore

$$F_E - F_B = \frac{\gamma q \lambda}{2\pi \epsilon_0 r} - \frac{\gamma q \lambda v^2}{2\pi \epsilon_0 r c^2} = \frac{\gamma q \lambda}{2\pi \epsilon_0 r} \left(1 - \frac{v^2}{c^2}\right) = \frac{q \lambda}{2\gamma \pi \epsilon_0 r}, \quad (6.79)$$

where we have used $1 - v^2/c^2 \equiv 1/\gamma^2$. This net force agrees with the result in part (a).

- (c) In the lab frame, the charges in the rod aren't moving, so \mathbf{E}_\perp is the only nonzero field in the Lorentz transformations in Eq. (6.76). It is directed away from the rod with magnitude $\lambda/2\pi\epsilon_0 r$. Equation (6.76) immediately gives the electric field in the new frame as $\mathbf{E}'_\perp = \gamma \mathbf{E}_\perp$. So \mathbf{E}'_\perp has magnitude $E'_\perp = \gamma\lambda/2\pi\epsilon_0 r$ and is directed away from the rod, in agreement with the electric field we found in part (b).

Equation (6.76) gives the magnetic field in the new frame as $\mathbf{B}'_\perp = -\gamma(\mathbf{v}/c^2) \times \mathbf{E}_\perp$. The velocity \mathbf{v} of F' with respect to the lab frame F points to the *left* with magnitude v . We therefore find that \mathbf{B}'_\perp points into the page with magnitude $B'_\perp = \gamma(v/c^2)(\lambda/2\pi\epsilon_0 r)$. In terms of $\mu_0 = 1/\epsilon_0 c^2$, this can be written as $B'_\perp = \mu_0 \gamma \lambda v/2\pi r$, in agreement with the magnetic field we found in part (b). We therefore arrive at the same net force, $F_E - F_B$, as in part (b).

2. $P = 2 \text{ W}$ (P: Power).

$$\therefore I = \frac{P}{V} = \frac{10^7}{5 \times 10^4} \text{ A} = 200 \text{ A}.$$

Field due to one wire, $B = \frac{\mu_0 I}{2\pi r}$.

$$\therefore B = \frac{4\pi \times 10^{-7} \times 200}{2\pi \times 1} = 4 \times 10^{-5} \text{ T}.$$

Other wire causes equal field in same direction.

$$\therefore \text{Total } B = 8 \times 10^{-5} \text{ T}.$$

3. $B_x = 0, B_y = 0, B_z = B_0.$

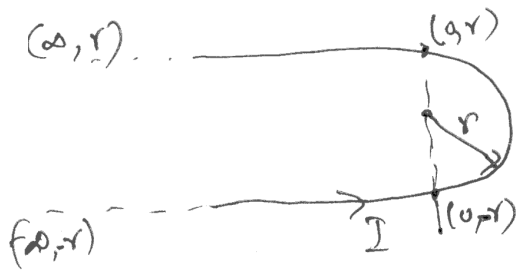
$$\vec{B} = \vec{\nabla} \times \vec{A} = \hat{x} \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) + \hat{y} \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) + \hat{z} \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right)$$

$$\therefore \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} = 0; \quad \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} = 0; \quad \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} = B_0.$$

One possible choice: $A_x = A_z = 0; A_y = B_0 x.$

What are other choices?

4.



Break wire into 3 segments:
 2 line segments (infinite)
 1 semicircular segment

1 segment goes from $(-\infty, -r)$ to $(0, -r)$

$$B_1 = \frac{\mu_0 I}{4\pi r} (\cos\theta_1 - \cos\theta_2) = \frac{\mu_0 I}{4\pi r} (\cos 0^\circ - \cos \pi) \\ = \frac{\mu_0 I}{4\pi r} (1 - (-1)) = \frac{\mu_0 I}{4\pi r} = B_3 \text{ (contribution from other line segment)}$$

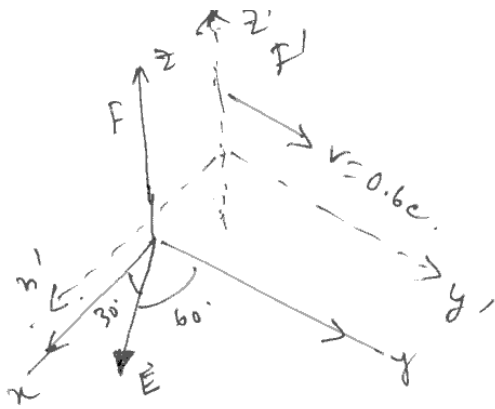
Contribution from semicircular arc,

$$B_2 = \frac{\mu_0 I}{4\pi} \int \frac{d\vec{l} \times \vec{r}}{r^2} = \frac{\mu_0 I}{4\pi} \int_0^\pi \frac{r d\theta}{r^2} = \frac{\mu_0 I}{4\pi r}$$

Note that direction of all 3 magnetic fields are out of page. If that is the $+\hat{z}$ direction, then,

$$\vec{B} = \vec{B}_1 + \vec{B}_2 + \vec{B}_3 = \left(2 \cdot \frac{\mu_0 I}{4\pi r} + \frac{\mu_0 I}{4\pi r} \right) \hat{z} \\ = \frac{\mu_0 I}{4\pi r} (2 + 1) \hat{z}$$

5.



Frame F : $\vec{E} = \begin{cases} E_x = 100 \cos 30^\circ \\ E_y = 100 \sin 30^\circ \\ E_z = 0 \end{cases}$

$\vec{B} = 0$.

Transformation eqns: $\vec{E}'_{\parallel} = \vec{E}_{\parallel}$; $\vec{E}'_{\perp} = \gamma(\vec{E}_{\perp} + \vec{v} \times \vec{B}_{\perp})$
 $\vec{B}'_{\parallel} = \vec{B}_{\parallel}$; $\vec{B}'_{\perp} = \gamma(\vec{B}_{\perp} - \frac{\vec{v}}{c^2} \times \vec{E}_{\perp})$

Hence, $\vec{v} = 0.6c \hat{y}$.

The component \parallel to direction of relative motion is the y component.

$\therefore E'_y = E_y$ & since, $\vec{E}'_{\perp} = \gamma \vec{E}_{\perp}$ ($\because \vec{B} = 0$).

$\therefore E'_x = \gamma E_x$; $E'_z = \gamma E_z = 0$ ($\because E_z = 0$)

$\gamma = \frac{1}{\sqrt{1 - \beta^2}} = \frac{1}{\sqrt{1 - (0.6)^2}} = 1.25$.

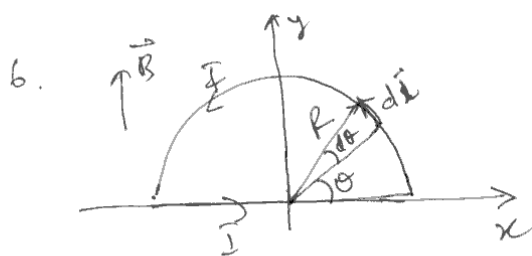
$\therefore E'_x = 1.25 E_x$; $E'_y = E_y$.

$\therefore \tan \phi' = \frac{E'_y}{E'_x} = \frac{1}{1.25} \frac{E_y}{E_x} \Rightarrow \phi' = \tan^{-1} \left(\frac{E_y}{1.25 E_x} \right)$

$|\vec{E}'| = \sqrt{E_x'^2 + E_y'^2 + E_z'^2} = \sqrt{E_x'^2 + E_y'^2}$.

Now, since, $\vec{B} = 0$ in F, $\therefore B'_y = B_y = 0$.

$\vec{B}'_{\perp} = -\frac{\gamma \vec{v} \times \vec{E}_{\perp}}{c^2} = -\frac{\gamma}{c^2} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0.6c & 0 \\ E_x & E_y & 0 \end{vmatrix} = -\frac{\gamma}{c^2} \hat{z} (0 - 0.6 E_x)$
 $= + \frac{\gamma}{c^2} 0.6 E_x \hat{z} = \frac{1.25}{c^2} E_x \hat{z}$.



$$\vec{B} = B \hat{y}$$

Length of straight segment
 $= 2R$

\therefore Magnetic force on straight segment,

$$\begin{aligned}\vec{F}_1 &= I(2R\hat{x}) \times B\hat{y} \quad (\text{recall, } \vec{F} = I \int d\vec{l} \times \vec{B}) \\ &= 2IRB\hat{z}. \quad (\text{out of page}).\end{aligned}$$

For the semicircular segment, take an infinitesimal length $d\vec{l}$ on the segment (fig.)

Then, $d\vec{l} = -R\sin\theta d\theta \hat{x} + R\cos\theta d\theta \hat{y}$.

$$\therefore d\vec{F}_2 = I(d\vec{l} \times \vec{B}) = I(-R\sin\theta d\theta \hat{x} + R\cos\theta d\theta \hat{y}) \times B\hat{y}.$$

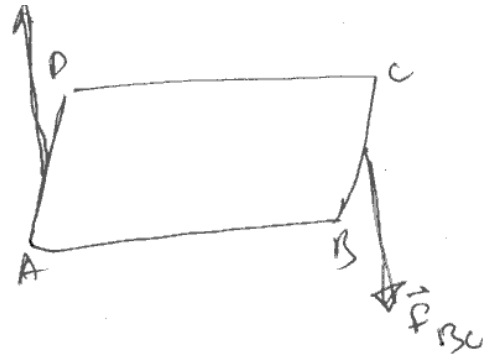
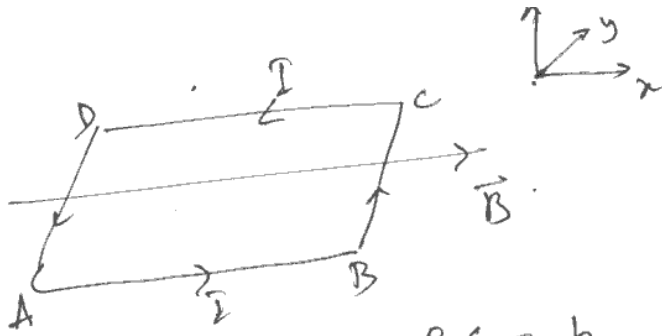
$$= -IRB\sin\theta d\theta \hat{z}$$

Integrating,
$$\begin{aligned}\vec{F}_2 &= -IRB \int_0^\pi \sin\theta d\theta \hat{z} \\ &= -2IRB\hat{z}. \quad (\text{into the page})\end{aligned}$$

\therefore Net force on the wire, $\vec{F} = \vec{F}_1 + \vec{F}_2 = 0$.

\rightarrow consistent with the fact that net magnetic force acting on a closed current carrying loop must be zero.

7.



Let, $AB = a$, $BC = b$.

Magnetic forces acting on line segments AB & CD are zero since they are \parallel & antiparallel to \vec{B} .

\therefore Cross products $= 0$.

On segment BC : ~~$\vec{F}_{BC} = I \int_a^b \vec{y} \times \vec{B} \hat{x}$~~

$$\vec{F}_{BC} = I (+b \hat{y}) \times B \hat{x} = -IbB \hat{z}.$$

On segment DA :

$$\vec{F}_{DA} = I (-b \hat{y}) \times B \hat{x} = +IbB \hat{z} = -\vec{F}_{BC}$$

Net force on loop:

$$\vec{F} = \vec{F}_{AB} + \vec{F}_{BC} + \vec{F}_{CD} + \vec{F}_{DA} = 0 - IbB \hat{z} + 0 + IbB \hat{z} = 0.$$

Although net force is zero, \vec{F}_{BC} & \vec{F}_{DA} produce a torque leading to a rotation of loop about the y axis. \therefore Torque w.r.t. centre of loop:

$$\begin{aligned}\vec{\tau} &= \left(-\frac{a}{2}\right)\hat{x} \times \vec{F}_{BC} + \left(-\frac{a}{2}\hat{x}\right) \times \vec{F}_{DA} \\ &= \frac{a}{2}\hat{x} \times (-IbB\hat{z}) + \frac{a}{2}\hat{x} \times (IbB\hat{z}) \\ &= \frac{IabB}{2}\hat{y} + \frac{IabB}{2}\hat{y} = IabB\hat{y} \\ &= IAB\hat{y} \quad \text{where, } A = ab \text{ is the area of loop.}\end{aligned}$$

\rightarrow rotation is clockwise.

$$\therefore \vec{\tau} = I (\vec{A} \times \vec{B}) \quad \text{where, } \vec{A} = A\hat{n} \quad (\hat{n} : \text{unit vector normal to plane of loop}).$$

In the general case when loop makes an angle θ with the field \vec{B} :

The lever arm,

$$\vec{r} = \frac{a}{2} (-\sin\theta \hat{n} + \cos\theta \hat{z}) = -\vec{r}'$$

$$\therefore \vec{\tau} = \vec{r} \times \vec{F}_{DA} + \vec{r}' \times \vec{F}_{BC} = 2\vec{r} \times \vec{F}_{DA}$$

$$= 2\frac{a}{2} (-\sin\theta \hat{n} + \cos\theta \hat{z}) \times IabB\hat{z}$$

$$= IabB \sin\theta \hat{y} = I\vec{A} \times \vec{B}$$

For a loop of N turns, $|\vec{\tau}| = \underbrace{N I A B \sin\theta}_{\text{magnetic moment}(\mu)}$

$$\vec{\mu} = N I \vec{A}$$

$$\therefore \boxed{\vec{\tau} = \vec{\mu} \times \vec{B}}$$

