

The circular restricted three body problem (CR3BP) is a special case of the three body problem. In the CR3BP (much like in Keplerian 2-body dynamics), we neglect the mass of satellite S , while treating the larger celestial body m_1 and smaller celestial body m_2 as point masses. Crucially, these bodies must orbit one another in circular orbits. In other words, they both orbit about their inertially fixed barycenter c at *constant velocity and distance*.

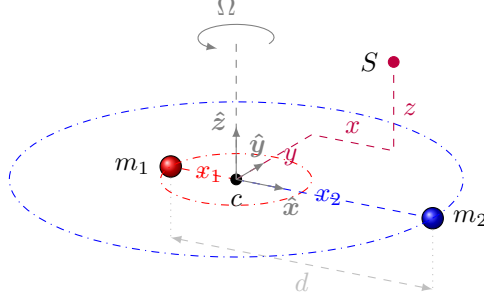


Figure 1: Geometry

The purple satellite is located by $\mathbf{r} = x\hat{x} + y\hat{y} + z\hat{z}$, and the spheres are the celestial bodies. The xyz frame is not inertially stationary. Instead, they rotate with the larger bodies at a rate Ω . The origin c is the center of mass of m_1 and m_2 (and therefore inertially fixed). The xy plane defines the plane in which m_1 and m_2 orbit, so \hat{z} is defined by the direction of their angular momenta- the constant rotation at a rate Ω is a consequence of this. Note that \mathbf{r} is not constrained to the xy plane, but m_1 and m_2 definitionally are.

Useful Relations

We can also find a relationship between x_1 , x_2 , d , m_1 , and m_2 . Because c is the barycenter,

$$x_1 m_1 = x_2 m_2$$

or

$$\frac{x_1}{m_2} = \frac{x_2}{m_1}$$

We can use this to find that

$$\begin{aligned} \frac{x_1}{x_1 + x_2} &= \frac{x_1/m_2}{x_1/m_2 + x_2/m_2} \\ &= \frac{x_2/m_1}{x_2/m_1 + x_2/m_2} \\ &= \frac{1/m_1}{1/m_1 + 1/m_2} \\ &= \frac{m_2}{m_2 + m_1} \end{aligned}$$

Defining $d = x_1 + x_2$ as the distance between the two celestial bodies, and $M = m_1 + m_2$ as their total mass, we get that

$$\boxed{\frac{x_1}{d} = \frac{m_2}{M}}$$

and similarly

$$\boxed{\frac{x_2}{d} = \frac{m_1}{M}}$$

Which can be rewritten as

$$\boxed{\frac{x_1}{m_2} = \frac{x_2}{m_1} = \frac{d}{M}}$$

Lastly, we will define two additional vectors \mathbf{r}_1 and \mathbf{r}_2 which point from the first and second body respectively to the satellite.

$$\mathbf{r}_1 = (x + x_1) \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}$$

and

$$\mathbf{r}_2 = (x - x_2) \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}$$

Now we will solve for Ω .

$$\begin{aligned} \mathbf{F}_{\text{on}2} &= -\frac{Gm_1m_2}{d^2} \hat{\mathbf{x}} \\ m_2 \mathbf{a}_2^f &= -\frac{Gm_1m_2}{d^2} \hat{\mathbf{x}} \\ \frac{d^f}{dt} \left(\frac{d^f}{dt} x_2 \hat{\mathbf{x}} \right) &= -\frac{Gm_1}{d^2} \hat{\mathbf{x}} \\ \frac{d^f}{dt} \left(\frac{d^c}{dt} x_2 \hat{\mathbf{x}} + \Omega \hat{\mathbf{z}} \times x_2 \hat{\mathbf{x}} \right) &= -\frac{Gm_1}{d^2} \hat{\mathbf{x}} \\ \frac{d^f}{dt} \Omega x_2 \hat{\mathbf{y}} &= -\frac{Gm_1}{d^2} \hat{\mathbf{x}} \\ \frac{d^c}{dt} \Omega x_2 \hat{\mathbf{y}} + \Omega \hat{\mathbf{z}} \times \Omega x_2 \hat{\mathbf{y}} &= -\frac{Gm_1}{d^2} \hat{\mathbf{x}} \\ -\Omega^2 x_2 \hat{\mathbf{x}} &= -\frac{Gm_1}{d^2} \hat{\mathbf{x}} \\ \Omega^2 x_2 &= \frac{Gm_1}{d^2} \\ \Omega &= \sqrt{\frac{G}{d^2} \frac{m_1}{x_2}} \\ \Omega &= \sqrt{\frac{GM}{d^3}} \end{aligned}$$

Defining μ conventionally as $\mu = GM$,

$$\Omega = \sqrt{\frac{\mu}{d^3}}$$

Kinematics

The transport theorem states that the inertial (fixed f frame) derivative of a vector \mathbf{u} (expressed in the rotating c frame) is

$$\frac{d^f \mathbf{u}}{dt} = \frac{d^c \mathbf{u}}{dt} + \boldsymbol{\omega}^{cf} \times \mathbf{u}$$

Where $\frac{d^f}{dt}$ denotes the derivative in the coordinates of the fixed frame f , and $\frac{d^c}{dt}$ denotes derivative in the coordinates of the rotating frame c , and $\boldsymbol{\omega}^{cf}$ denotes the angular velocity of c in f . For this case, $\boldsymbol{\omega}^{cf} = \Omega \hat{\mathbf{z}}$. We can find Ω

For the satellite's position in the CR3BP frame $\mathbf{r} = x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}$, we will find the inertial acceleration to generate equations of motion.

$$\begin{aligned} \dot{\mathbf{r}} &= \frac{d^c \mathbf{r}}{dt} + \boldsymbol{\omega}^{cf} \times \mathbf{r} \\ &= \frac{d^c}{dt} (x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}) + (\Omega \hat{\mathbf{z}} \times (x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}})) \\ &= (\dot{x} \hat{\mathbf{x}} + \dot{y} \hat{\mathbf{y}} + \dot{z} \hat{\mathbf{z}}) + (\Omega x \hat{\mathbf{y}} - \Omega y \hat{\mathbf{x}}) \\ &= (\dot{x} - \Omega y) \hat{\mathbf{x}} + (\dot{y} + \Omega x) \hat{\mathbf{y}} + \dot{z} \hat{\mathbf{z}} \end{aligned}$$

$$\begin{aligned}
\ddot{\mathbf{r}} &= \dot{\mathbf{r}} \\
&= \frac{d^c}{dt} ((\dot{x} - \Omega y) \hat{\mathbf{x}} + (\dot{y} + \Omega x) \hat{\mathbf{y}} + \dot{z} \hat{\mathbf{z}}) \\
&\quad + \Omega \hat{\mathbf{z}} \times ((\dot{x} - \Omega y) \hat{\mathbf{x}} + (\dot{y} + \Omega x) \hat{\mathbf{y}} + \dot{z} \hat{\mathbf{z}}) \\
&= (\ddot{x} - \Omega \dot{y}) \hat{\mathbf{x}} + (\ddot{y} + \Omega \dot{x}) \hat{\mathbf{y}} + \ddot{z} \hat{\mathbf{z}} \\
&\quad + ((\Omega \dot{x} - \Omega^2 y) \hat{\mathbf{y}} - (\Omega \dot{y} + \Omega^2 x) \hat{\mathbf{x}}) \\
&= (\ddot{x} - 2\Omega \dot{y} - \Omega^2 x) \hat{\mathbf{x}} + (\ddot{y} + 2\Omega \dot{x} - \Omega^2 y) \hat{\mathbf{y}} + \ddot{z} \hat{\mathbf{z}} \\
\boxed{\ddot{\mathbf{r}} &= (\ddot{x} - 2\Omega \dot{y} - \Omega^2 x) \hat{\mathbf{x}} + (\ddot{y} + 2\Omega \dot{x} - \Omega^2 y) \hat{\mathbf{y}} + \ddot{z} \hat{\mathbf{z}}}
\end{aligned}$$

Equations of Motion

We can now generate the equations of motion

$$\begin{aligned}
\sum_i \mathbf{F}_i &= m \ddot{\mathbf{r}} \\
\mathbf{F}_1 + \mathbf{F}_2 &= m \ddot{\mathbf{r}} \\
-\frac{\mu_1 m}{r_1^3} \mathbf{r}_1 - \frac{\mu_2 m}{r_2^3} \mathbf{r}_2 &= m \ddot{\mathbf{r}} \\
-\frac{\mu_1}{r_1^3} \mathbf{r}_1 - \frac{\mu_2}{r_2^3} \mathbf{r}_2 &= (\ddot{x} - 2\Omega \dot{y} - \Omega^2 x) \hat{\mathbf{x}} + (\ddot{y} + 2\Omega \dot{x} - \Omega^2 y) \hat{\mathbf{y}} + \ddot{z} \hat{\mathbf{z}}
\end{aligned}$$

We now write this as three equations, one each in x , y , and z

$$\begin{aligned}
-\frac{\mu_1}{r_1^3}(x + x_1) - \frac{\mu_2}{r_2^3}(x - x_2) &= \ddot{x} - 2\Omega \dot{y} - \Omega^2 x \\
-\frac{\mu_1}{r_1^3}y - \frac{\mu_2}{r_2^3}y &= \ddot{y} + 2\Omega \dot{x} - \Omega^2 y \\
-\frac{\mu_1}{r_1^3}z - \frac{\mu_2}{r_2^3}z &= \ddot{z}
\end{aligned}$$

Isolating the second derivatives,

$$\begin{aligned}
\ddot{x} &= -\frac{\mu_1}{r_1^3}(x + x_1) - \frac{\mu_2}{r_2^3}(x - x_2) + 2\Omega \dot{y} + \Omega^2 x \\
\ddot{y} &= -\frac{\mu_1}{r_1^3}y - \frac{\mu_2}{r_2^3}y - 2\Omega \dot{x} + \Omega^2 y \\
\ddot{z} &= -\frac{\mu_1}{r_1^3}z - \frac{\mu_2}{r_2^3}z
\end{aligned}$$

We can now substitute $\Omega = \sqrt{\frac{\mu}{d^3}} = \sqrt{\frac{\mu_1 + \mu_2}{d^3}}$

$$\begin{aligned}
\ddot{x} &= -\frac{\mu_1}{r_1^3}(x + x_1) - \frac{\mu_2}{r_2^3}(x - x_2) + 2\sqrt{\frac{\mu}{d^3}}\dot{y} + \frac{\mu}{d^3}x \\
\ddot{y} &= -\frac{\mu_1}{r_1^3}y - \frac{\mu_2}{r_2^3}y - 2\sqrt{\frac{\mu}{d^3}}\dot{x} + \frac{\mu}{d^3}y \\
\ddot{z} &= -\frac{\mu_1}{r_1^3}z - \frac{\mu_2}{r_2^3}z
\end{aligned}$$

We can also get rid of μ_2 and x_2

$$\begin{aligned}
\ddot{x} &= -\frac{\mu_1}{r_1^3}(x + x_1) - \frac{\mu_2}{r_2^3}(x - d + x_1) + 2\sqrt{\frac{\mu}{d^3}}\dot{y} + \frac{\mu}{d^3}x \\
\ddot{y} &= -\frac{\mu_1}{r_1^3}y - \frac{\mu - \mu_1}{r_2^3}y - 2\sqrt{\frac{\mu}{d^3}}\dot{x} + \frac{\mu}{d^3}y \\
\ddot{z} &= -\frac{\mu_1}{r_1^3}z - \frac{\mu - \mu_1}{r_2^3}z
\end{aligned}$$

Now, we will non-dimensionalize these to make the EOMs more general. The distance unit is $L = x_1 + x_2$. The time unit is $T = \sqrt{d/\mu}$ (the inverse of angular velocity). To non-dimensionalize, we must cancel units. The mass unit is $M = m_1 + m_2$

$$\begin{aligned}\ddot{x} \frac{T^2}{L} &= -\frac{\mu_1 \left(\frac{T^2}{L^3}\right)}{r_1^3 \left(\frac{1}{L}\right)^3} (x + x_1) - \frac{\mu_2}{r_2^3} (x - d + x_1) + 2\sqrt{\frac{\mu}{d^3}} \dot{y} + \frac{\mu}{d^3} x \\ \ddot{y} \frac{T^2}{L} &= -\frac{\mu_1 \left(\frac{T^2}{L^3}\right)}{r_1^3 \left(\frac{1}{L}\right)^3} y - \frac{\mu - \mu_1}{r_2^3} y - 2\sqrt{\frac{\mu}{d^3}} \dot{x} + \frac{\mu}{d^3} y \\ \ddot{z} \frac{T^2}{L} &= -\frac{\mu_1 \left(\frac{T^2}{L^3}\right)}{r_1^3 \left(\frac{1}{L}\right)^3} z - \frac{\mu - \mu_1}{r_2^3} z\end{aligned}$$