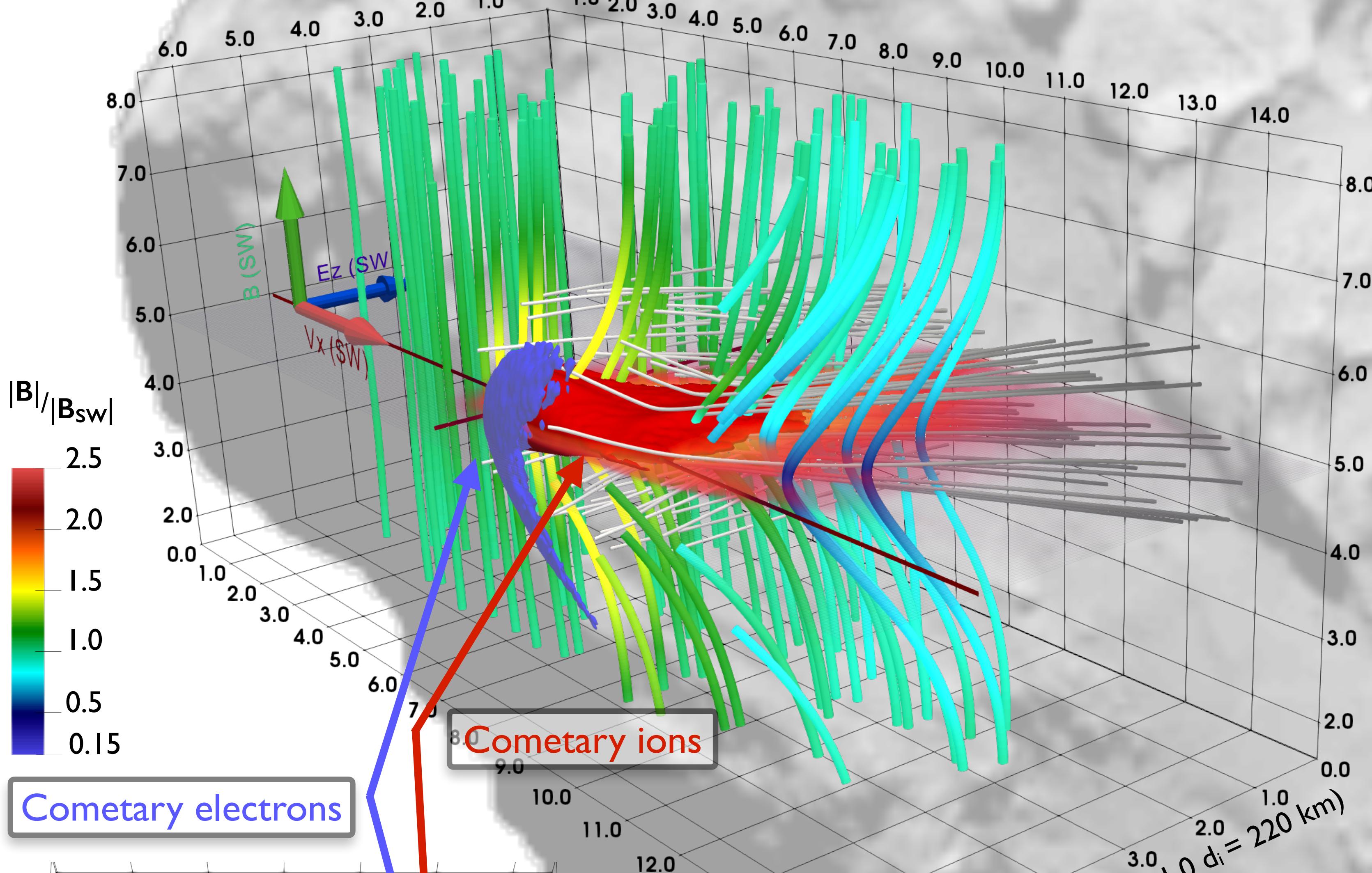


# A fully kinetic perspective of electron acceleration around a weakly outgassing comet

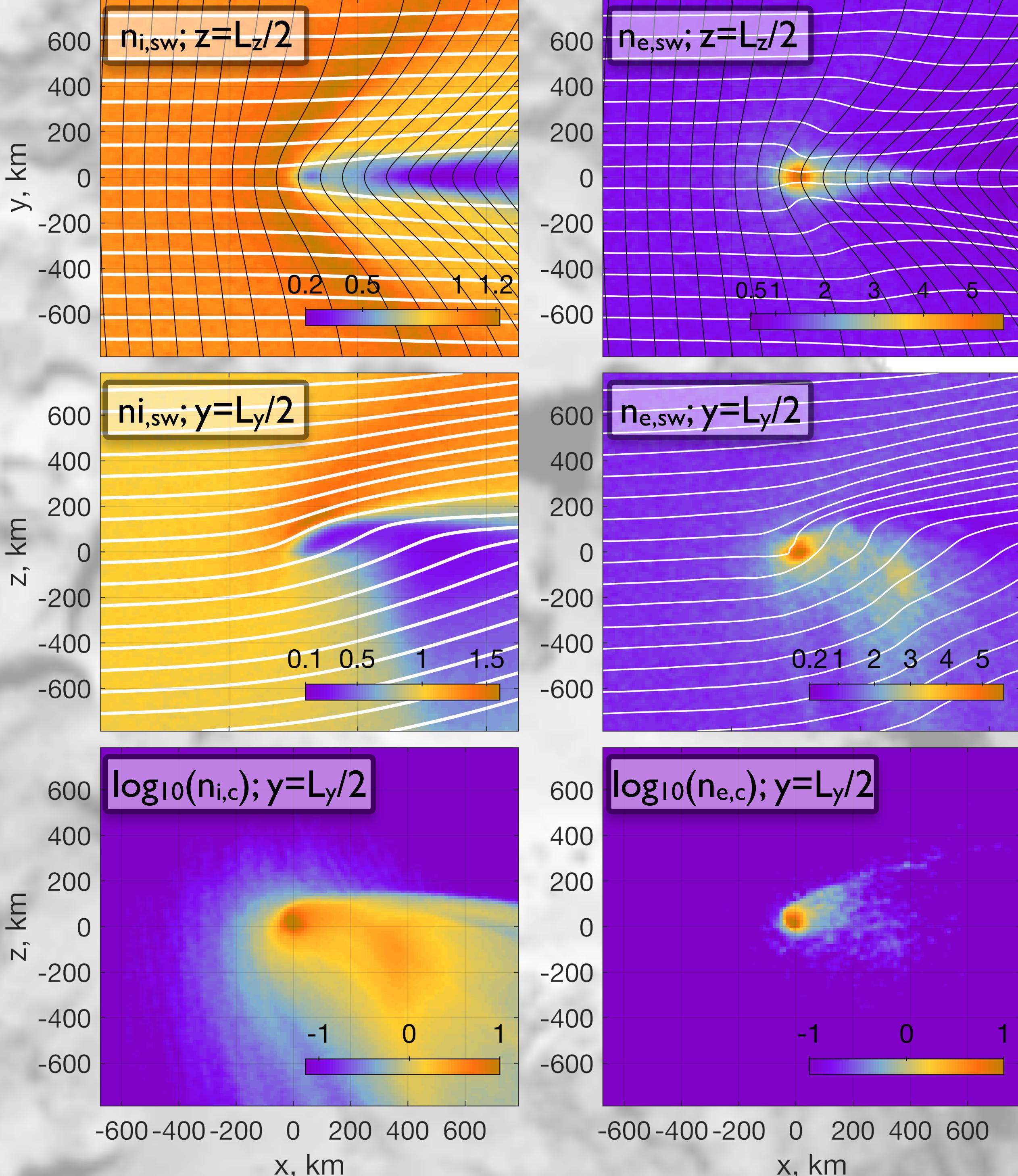
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When a comet is sufficiently close to the Sun, the sublimation of ice leads to outgassing and the formation of a coma of gas and dust. Ionisation results in mass loading of the solar wind and magnetic field draping. Here we present three-dimensional fully kinetic simulations of the solar wind interaction with comet 67P/Churyumov-Gerasimenko at a very low activity regime, before collisions have any impact on the plasma dynamics. The interaction scales are well below the relevant ion gyroradii and non-equilibrium electron distributions develop. To first order, the dynamical interaction is representative of a four-fluid coupled system [Deca et al., PRL 2017]. We focus here on the electron acceleration mechanisms around the weakly outgassing comet [Divin et al., submitted].

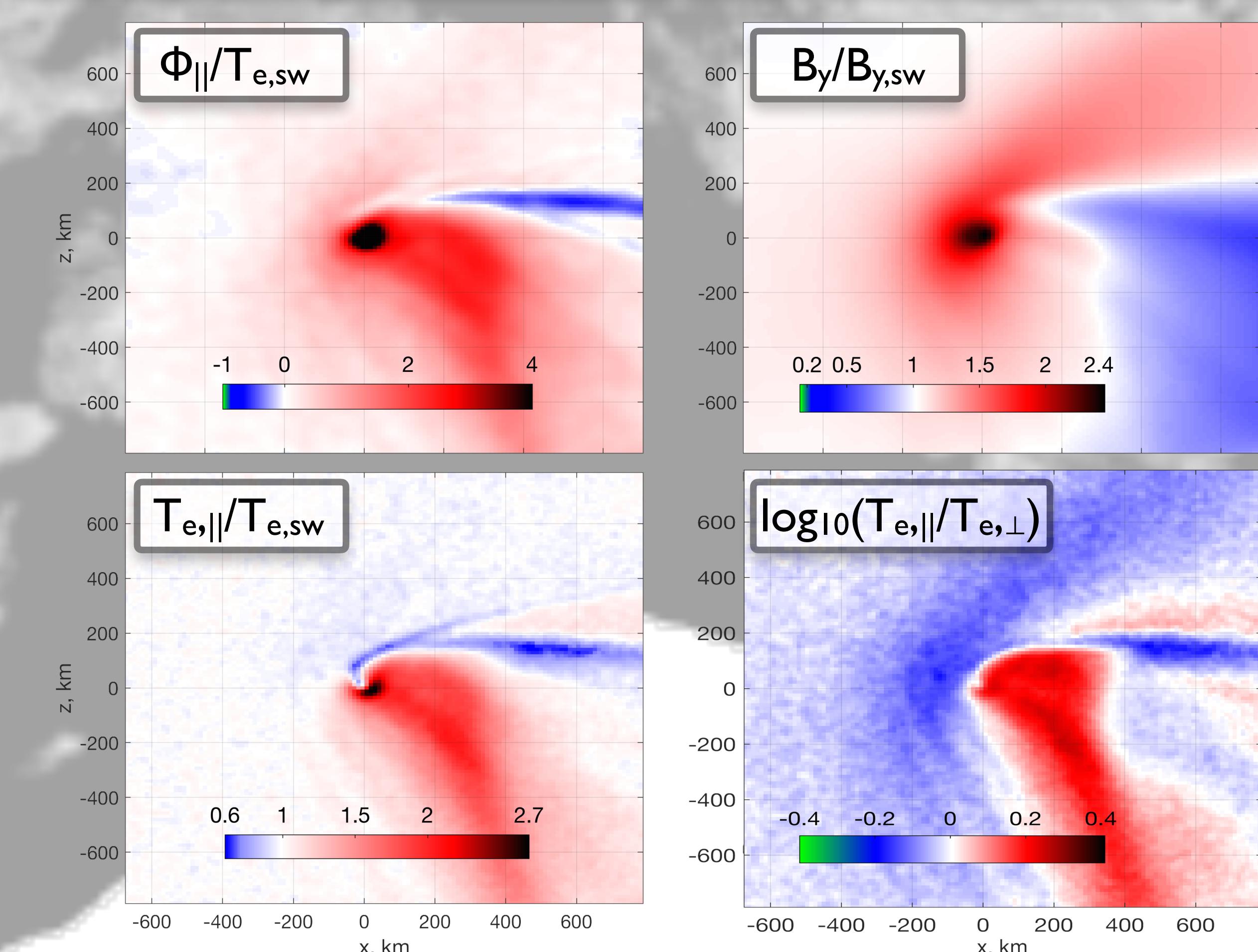


(Left and above). A 3D overview of the simulation, showing the global morphology of the interaction. In this low-outgassing regime no shock or diamagnetic cavity forms. The magnetic field drapes around the comet. We are in the collisionless regime.

(Right). The spatial distribution of the effective acceleration potential,  $\Phi_{||}$ , is in good agreement with the parallel electron temperature,  $T_{e,||}$ , indicating a strong acceleration (of solar wind/suprathermal) electrons by the parallel (ambipolar) electric field. Electrons with  $T_e < T_{e,sw}$  are picked up and create a narrow region of  $\Phi_{||} < 0$  at the edge of the wake. Areas where  $T_{e,||}/T_{e,\perp}$  is small correspond to larger magnetic field values in the pile-up region.



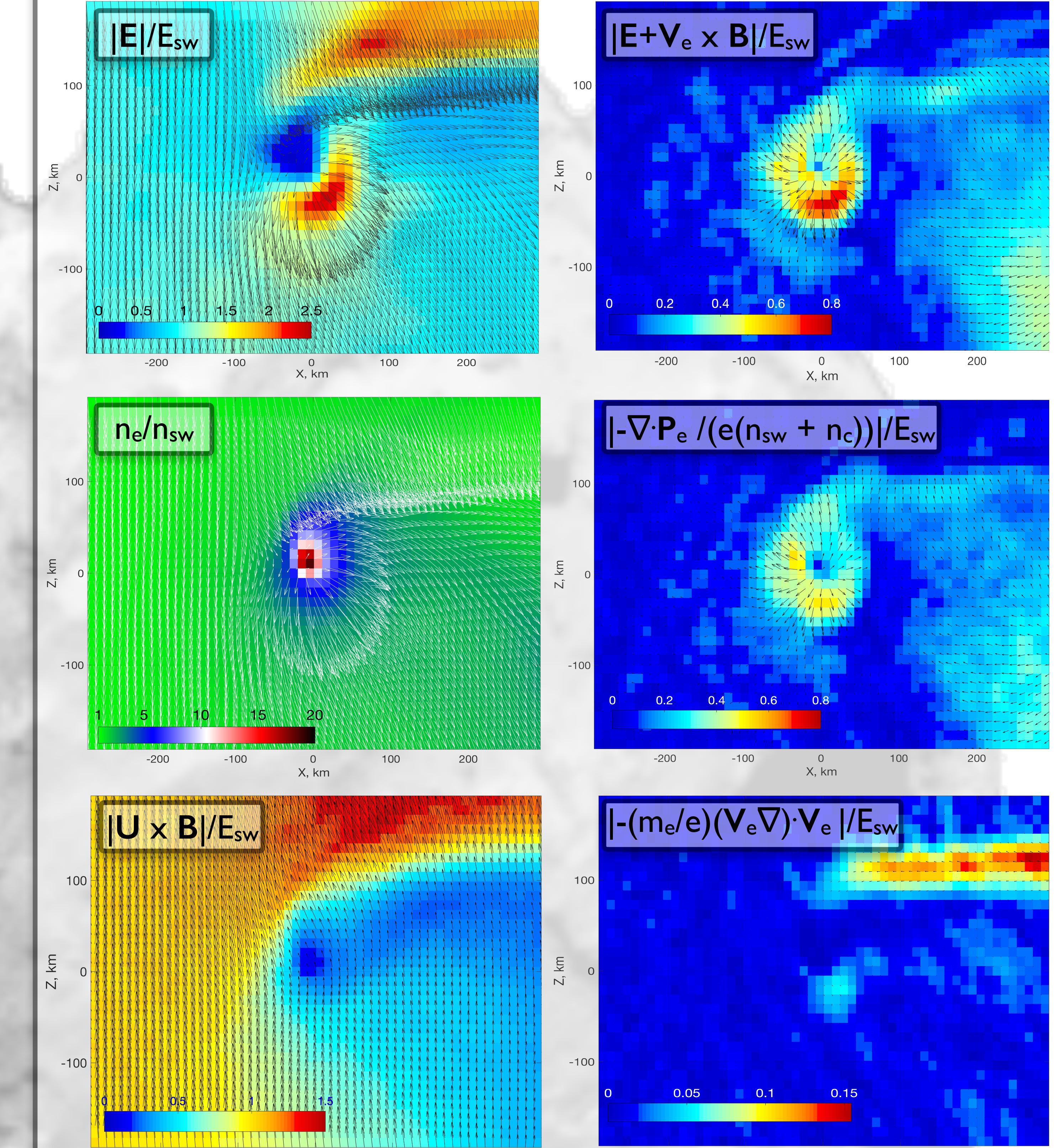
(Above). Density profiles of the 4 plasma species included in the simulation, normalised to  $n_{sw}$ . Magnetic field lines are shown in black, velocity streamlines in white. Cometary ions accelerate along the convective electric field, cometary electrons are initially accelerated in the opposite direction. Solar wind protons and electrons are decoupled close to the nucleus and in the wake.



**Ohm's law** is distilled directly from the electron dynamics in the simulation. In the vicinity of the cometary nucleus, the balance changes between the different terms in the equation. Deciphering the relative importance of each term allows to identify the driving physics in the various regions of the cometary plasma environment. It shows us what happens to all ion and electron species, and why.

$$\mathbf{V}_e \times \mathbf{B} = \mathbf{U} \times \mathbf{B} - \mathbf{j} \times \mathbf{B} / (e(n_{sw} + n_c))$$

$$\mathbf{E} + \mathbf{U} \times \mathbf{B} = \mathbf{j} \times \mathbf{B} / (e(n_{sw} + n_c)) - \nabla P_e / (e(n_{sw} + n_c)) - (m_e/e)(\mathbf{V}_e \nabla) \cdot \mathbf{V}_e$$



- Close to the outgassing nucleus the electron pressure gradient dominates.
- At sub-ion scales, the total electric field is a superposition of the solar wind convective electric field, where electrons are frozen-in, and the ambipolar electric field. The latter accelerates electrons parallel to the magnetic field and is the source of/provides feedback to the electron pressure gradient that balances the Ohm's law perpendicular to the magnetic field.
- The role of electron inertia is negligible to balance the electric field.

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