Empirical Dust Models Changhoon $\operatorname{Hahn}^{1,2,*}$ and IQ Collaboratory

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ABSTRACT

We construct a dust empirical model (DEM) framework for applying dust attenuation to simulated galaxies based on handful of sensible assumptions that come from our understanding of dust attenuation. The DEM framework is essentially based on state-of-the-art attenuation curves with a flexible parameterization that allows us to statistically sample them. Applying DEMs to three different hydrodynamic simulations, SIMBA, Illustris TNG, and EAGLE, we are able to produce UV and optical photometry consistent with SDSS observations $((G - R) - M_r)$ color-magnitude and $(FUV - NUV) - M_r$ relations). This suggests that there's enough freedom in our current understanding (or lack) of dust for all simulations to reproduce observations. Meanwhile, the DEM provides some insights into dust as well as the subgrid physics that goes into the hydro simulations.

Keywords: keyword1 – keyword2 – keyword3

1. INTRODUCTION

dust is important because....

assumptions on the attenuation curve can dramatically impact the physical properties inferred from SED fitting (e.g. Kriek & Conroy 2013; ?; ?; Salim & Narayanan 2020).

motivation for an empirical dust attenuation model

attenuation vs extinction. While extinction curves have been derived from observations and theoretically, it's not easy to map this to attenuation curves. Attenuation curves are a product of complicated empirical processes since it accounts for light that gets scattered and star light that is not obscured

This makes modeling them in a complete physically motivated method expensive. People have done it Narayanan et al. (2018); Trayford et al. (2020). some detail about the radiative transfer method and such. But besides being expensive they have to make a number of assumptions anyway. e.g. Narayanan et al. (2018) assumes a fixed extinction curve.

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Moreover, because the radiative transfer method is expensive it's hard to compare many different simulations. Not only that, observables generated from simulations that take into radative transfer dust models complicates simulation to simulation comparisons. Because you're simultaneously comparing the galaxy formation prescription and all the dust prescription.

Instead, we present a framework using flexible dust empirical models that paints attenuation curves onto galaxies. describe at a high level how we are parameterizing DEMs

talk about the advantages: extremely flexible so it can encompass the wide variety of attenuation curves found in radiative transfer, easy to correlate the attenuation curve with galaxy properties.

Also DEMs make it possible to statistically apply attenuation curves for large galaxy population. Putting this ontop of simulations, we can use them to generate observables and compare them to observations to constrain the DEM. This framework allows us to learn about attenuation curves given a model for galaxy formation.

The other way around also works. If you don't care about dust at all, DEM provides a framework to easily marginalize over dust attenuation and treat dust as a nuisance parameter.

In this paper, we do above for multiple simulations.

Starkenburg et la. in prep will use this framework to marginalize over dust and compare galaxy populations predicted by multiple simulations .

2. DATA

2.1. SDSS DR7 Central Galaxies

Throughout the paper we compare the simulations and models described below to the observed SDSS central galaxy sample from the Tinker et al. (2011) group catalog. The group catalog, first, selects volume-limited sample of galaxies at $z \approx 0.04$ with $M_r < -18$ and complete above $M_* > 10^{9.4}h^{-2}M_{\odot}$ from the NYU Value-Added Galaxy Catalog (VAGC; Blanton et al. 2005) of SDSS DR7 (Abazajian et al. 2009). The stellar masses are estimated using the kcorrect code (Blanton & Roweis 2007) assuming a Chabrier (2003) initial mass function.

Central galaxies are then identified using a halo-based group finder that uses the abundance matching ansatz to iteratively assign halo masses to groups. Every group contains one central galaxy, which by definition is the most massive, and a group can contain ≥ 0 satellites. As with any group finder, galaxies are misassigned due to projection effects and redshift space distortions; however, the central galaxy sample has a purity of $\sim 90\%$ and completeness of $\sim 95\%$ (Tinker et al. 2018).

2.2. Illustris TNG

describe what galaxy properties (SFH, ZH, etc) are available

2.3. SIMBA

describe what galaxy properties (SFH, ZH, etc) are available

In Figure 1, we compare the stellar mass function (SMF) of our SDSS central galaxy sample along with central galaxy SMFs of the SIMBA (orange) and TNG (blue) simulations. The uncertainties for the SDSS SMF are derived from jackknife resampling. Although we present the SMFs for reference, we do not use stellar masses throughout the paper since they are inconsistently defined among simulations

TODO

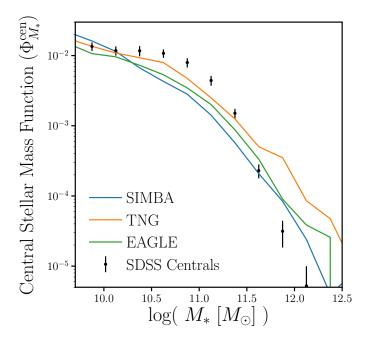


Figure 1. The stellar mass functions of central galaxies, $\Phi_{M_*}^{\text{cen}}$, of the SIMBA (orange) and TNG (blue) simulations compared to the SDSS $\Phi_{M_*}^{\text{cen}}$. The uncertainties for the SDSS $\Phi_{M_*}^{\text{cen}}$ are derived using jackknife resampling and SDSS centrals are identified using a halo-based group finder (Section 2.1). For SIMBA and TNG galaxies, we calculate M_* as the total stellar mass within the host halo, excluding contributions from any subhalos; centrals are classified based on their individual definition (Section 2.2 and 2.3). should we include total SMF...? The simulations and observations have loosely consistent $\Phi_{M_*}^{\text{cen}}$.

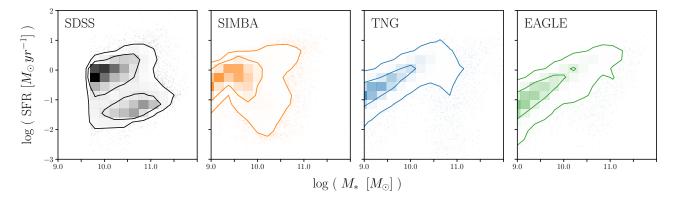


Figure 2. The M_* – SFR relation of central galaxies in SIMBA (orange), TNG (blue), and EAGLE (green) simulations and SDSS observations. CH: describe where the M* and SFRs come from for everythign. The differences among the M_* – SFR relations demonstrate that the hydro simulations predict galaxy populations with significantly different physical properties.

and observations. Instead, we compare between the simulations and SDSS using luminosity, M_r , which we consistently forward model and measure in the simulations. In these comparisons, we restrict ourselves to galaxies brighter than $M_r < -20$, where our SDSS central galaxy sample is complete.

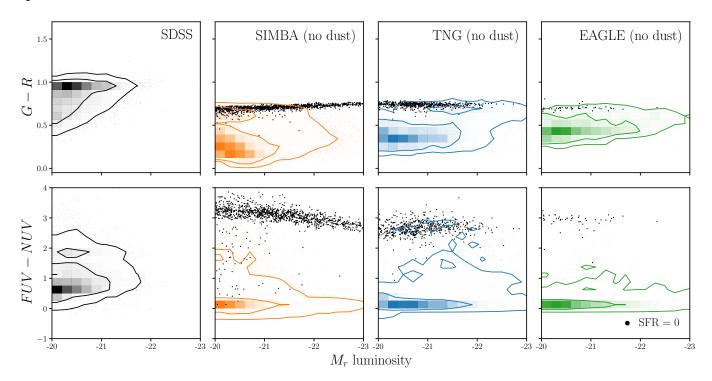


Figure 3. We present the distributions of the main observables used throughout the paper for SDSS (left), SIMBA (center), and TNG (right) centrals. We do not include any prescription for dust for the simulated galaxies. The top panels present the G-R versus M_r color magnitude relations while the bottom panels present the FUV-NUV versus M_r relations. The observables for the SIMBA and TNG simulated galaxies are derived using forward modeling and are therefore consistent with SDSS measurements (Section 2.5). The contours for SIMBA and TNG do not include galaxies with SFR= 0, which we mark separately in black. Despite similar SMFs, the simulations without any dust prescription show stark differences with observations in the color-magnitude observable-space.

instantaneous SFR=0 for $\sim 11\%$ of SIMBA galaxies, $\sim 13\%$ for TNG, $\sim 2\%$ for EAGLE

2.4. Spectral Energy Distributions

describe how the SED is generated using the SFH and ZHs

2.5. Forward Modeling SDSS Photometry and Spectra

3. DUST EMPIRICAL MODELING

3.1. Fiducial DEM

motivation for the DEM model

We begin by defining the dust attenuation curve $A(\lambda)$ as

$$F_o(\lambda) = F_i(\lambda) 10^{-0.4A(\lambda)} \tag{1}$$

where F_o is the observed flux and F_i is the intrinsic flux. We normalize the attenuation at the V band,

$$A(\lambda) = A_V \frac{k(\lambda)}{k_V}. (2)$$

TODO

For the normalization of the attenuation curve, A_V , we use the slab model from Somerville & Primack (1999); Somerville et al. (2012). In the slab model the amplitude of attuenuation depends on the inclination angle, i, and the optical depth, τ_V :

$$A_V = -2.5 \log \left[\frac{1 - e^{-\tau_V \sec i}}{\tau_V \sec i} \right] \tag{3}$$

justification of why this is enough. We sample i uniformly.

R. We TODO

Recently, Salim & Narayanan (2020) find significant dependence in A_V on both M_* and SFR. We include this dependence through τ_V , which we flexibly parameterize as

$$\tau_V(M_*, SFR) = m_{\tau,1} \log \left(\frac{M_*}{10^{10} M_{\odot}} \right) + m_{\tau,2} \log SFR + c_{\tau}.$$
 (4)

Next, for the wavelength dependence of the attenuation curve, we use $k(\lambda)$ from Noll et al. (2009):

$$k(\lambda) = (k_{\text{Cal}}(\lambda) + D(\lambda)) \left(\frac{\lambda}{\lambda_V}\right)^{\delta}.$$
 (5)

Here $k_{\text{Cal}}(\lambda)$ is the Calzetti (2001) curve:

$$k_{\text{Cal}}(\lambda) = \begin{cases} 2.659(-1.857 + 1.040/\lambda) + R_V, & 6300\text{Å} \le \lambda \le 22000\text{Å} \\ 2.659(-2.156 + 1.509/\lambda - 0.198/\lambda^2 + 0.011/\lambda^3) + R_V & 1200\text{Å} \le \lambda \le 6300\text{Å} \end{cases}$$

where λ_V is the V band wavelength. δ , the slope of the attenuation curve. also correlates with galaxy properties. So we parameterize δ and

$$\delta(M_*, SFR) = m_{\delta,2} \log \left(\frac{M_*}{10^{10} M_{\odot}} \right) + m_{\delta,2} \log SFR + c_{\delta}$$
 (6)

 $D(\lambda)$ is the UV dust bump, which we parameter using the standard Lorentzian-like Drude profile:

$$D(\lambda) = \frac{E_b(\lambda \Delta \lambda)^2}{(\lambda^2 - \lambda_0^2)^2 + (\lambda \Delta \lambda)^2}$$
 (7)

where λ_0 , $\Delta\lambda$, and E_b are the central wavelength, FWHM, and strength of the bump, respectively. In our DEM, we assume fixed $\lambda_0 = 2175 \mathring{A}$ and $\Delta\lambda = 350 \mathring{A}$.

Kriek & Conroy (2013) and Tress et al. (2018) found evidence that E_b correlates with the slope of the attenuation curve for star-forming galaxies $z \sim 2$. This was dependence was confirmed with simulations in ?. E_b :

$$E_b = m_E \ \delta + c_E \tag{8}$$

we fixed this and find our results do not change significantly.

TODO

We also split the attenuation on the star light and nebular emission

$$F_o(\lambda) = F_i^{\text{star}}(\lambda) 10^{-0.4A(\lambda)} + F_i^{\text{neb}}(\lambda) 10^{-0.4A_{\text{neb}}(\lambda)}$$
(9)

where we parameterize

$$A_{\text{neb}}(\lambda) = f_{\text{neb}}A(\lambda) \tag{10}$$

mention of how we treat SFR = 0 galaxies

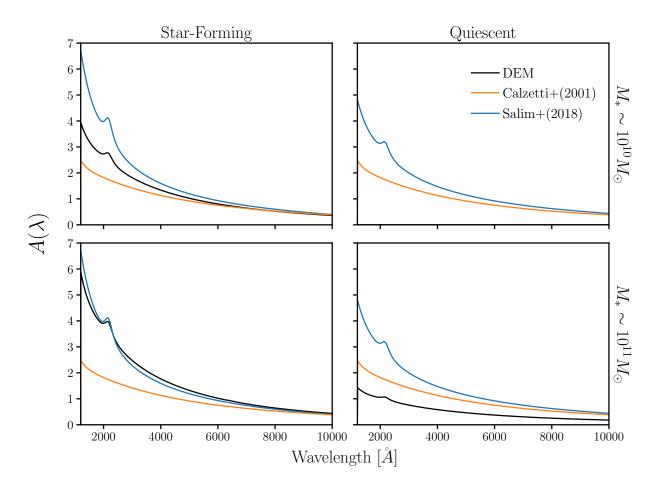


Figure 4. Comparison of the attenuation curve for our fiducial dust empirical model (DEM; black) to common attenuation curves in the literature (Calzetti 2001, orange; Salim et al. 2018, blue). We compare the curves for typical star-forming galaxies with SFR = $10^{0.5} M_{\odot}/yr$ in the left panel and for quiescent galaxies with SFR = $10^{-2} M_{\odot}/yr$ in the right. For our fiducial DEM, we use parameter values at the center of the priors listed in Table 1 and include attenuation curves for $M_* = 10^{9.5}$ (dashed) and $10^{11} M_{\odot}$ (solid). Our fiducial DEM is flexibly parameterized to incorporate both M_* and SFR dependence in the attenuation curve (Section 3.1).

3.2. Likelihood-Free Inference

Approximate Bayesian Computation with Population Monte Carlo Hahn et al. (2017),

To compare the outputs of our DEMs to observations, we first measure the color-magnitude observables $(G-R, FUV-NUV, \text{ and } M_r)$ as described in Section 2.5 consistent with SDSS measurements. Afterwards, we compare the forward modeled observables to SDSS using a L2 norm distance metric:

$$\bar{\rho}(\theta) = \sum_{i=1}^{n} \left[X_i^{\text{SDSS}} - X_i^{\text{model}}(\theta) \right]^2. \tag{11}$$

 X^{SDSS} and $X^{\text{model}}(\theta)$ are *n*-dimensional data vectors of the SDSS and model observables. In our case, we use a 3-dimensional histogram along G-R, FUV-NUV,

Table 1. Parameters of the	Dust Em	pirical M	lodels
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Parameter	Definition	
$m_{ au,M_*}$	Slope of the log M_* dependence of optical depth, τ_V	flat $[-5., 5.]$
$m_{ au, { m SFR}}$	Slope of the log SFR dependence of optical depth, τ_V	flat $[-5., 5.]$
$c_{ au}$	amplitude of the optical depth, τ_V	flat $[0., 6.]$
$m_{\delta,1}$	Slope of the $\log M_*$ dependence of the attenuation curve slope offset, δ	flat $[-4., 4.]$
$m_{\delta,2}$	Slope of the log SFR dependence of the attenuation curve slope offset, δ	flat $[-4., 4.]$
c_{δ}	amplitude of the attenuation curve slope offset, δ	flat $[-4., 4.]$
$f_{ m neb}$	fraction of nebular attenuation curve	flat $[1., 4.]$

Ishida et al. (2015)

Paragraph on the priors we choose.

- flat priors on everything. We might want to note that this doesn't give flat priors on τ_V or μ_{A_V} and σ_{A_V} for NT but we're ultimately interested in the dependence.
- priors were chosen to be uniformative and encompass constraints in the literature:
- m_E and c_E were chosen to be include Kriek & Conroy (2013); Narayanan et al. (2018); Tress et al. (2018)

4. RESULTS

We present the posterior distributions of the DEM parameters for the SIMBA (orange), TNG (blue), and EAGLE (green) hydro simulations in Figure 5. The DEM parameters include m_{τ,M_*} , $m_{\tau,\rm SFR}$, and c_{τ} , which parameterize the M_* dependence, SFR dependence, and amplitude of τ_V , the V-band optical depth. τ_V dictates the overall strength of the dust attenuation. They also include m_{δ,M_*} , $m_{\delta,\rm SFR}$, and c_{δ} , which parameterize the M_* dependence, SFR dependence, and amplitude of δ , the slope offset of the attenuation curve (Section 3.1 and Table 1). The posteriors are derived using ABC (Section 3.2) and the contours mark the 68% and 95% confidence intervals.

In addition, we present the observables predicted by the DEM with median of the posteriors for the SIMBA (orange), TNG (blue), and EAGLE (green) simulations in Figure 6. We include the observables for SDSS in the left-most panel for comparison (Section 2.1). The top panels present the $(G-R)-M_r$ color-magnitude relations while the bottom panels present the $(FUV-NUV)-M_r$ relations. In Figure 3, where we compare of the same observables but for simulations without any dust prescription, we find dramatic differences in observable-space between SDSS and the simulations. In contrast, using DEMs we produce $(G-R)-M_r$ and $(FUV-NUV)-M_r$ relations consistent with SDSS for all of the simulations.

The agreement in the observables is inspite of the significant differences in the SMFs and M_* -SFR relations (Figures ?? and 2). In other words, the DEM has the flexibility to reproduce observations even for simulations that predict galaxy populations with significantly different physical properties.

We emphasize that the DEM is based on the standard prescriptions for dust attenuation and, thus, serve as a flexible parameterization within the bounds of our current understanding of dust in galaxies.

CH: More detailed look into the DEM observables.

• how does it compare to EAGLE+SKIRT?

Figure 6 highlights two key points. First, any comparison of simulations must account for dust. Dust entirely changes the predictions of simulations in observables-space. Fortunately, the DEM provides a simple framework for including dust without the need for expensive ray-tracing methods. Second, the current limitations in our understanding of dust in galaxies significant impedes our ability to understand galaxy formation from simulations. To robustly interpret any comparison of simulations, we would need to marginalize over dust (e.g. DEM parameters). Since DEMs can produce consistent observables for a range of simulations, marginalizing over dust would leave little constraining power on the subgrid prescriptions (i.e. galaxy physics) of the simulations.

CH: What do the differences in DEM parameters say about the differences among the hydro sims? The whole SIMBA discrepancy CH: combining our understanding of dust maybe we can glean something about galaxy evolution

While DEMs demonstrate that dust is a major bottleneck for interpreting galaxy simulations, they also provide some insights into dust. For instance, the posteriors of DEM parameters in Figure 5 reveal consistent trends among the simulations. In all three simulations, we find significant positive M_* dependence of τ_V : $m_{\tau,M_*} \sim 2$. Galaxies with higher M_* have overall higher dust attenuation. CH: how does this compare to the literature? Salim et al. (2018) finds the same thing with the GALEX-SDSS-WISE Legacy Catalog 2 (GSWLC2) sample for SF galaxies.

Something about attenuation curves for quiescent galaxies since they are difficult to get from observations (CH: why?)

We also find overall little M_* and SFR dependence in δ . In fact, the amplitude of δ is roughly consistent with 0. CH: what does this mean? CH: how does this compare to the literature? This is consistent with Salim & Narayanan (2020), where they measured the attenuation curve slopes of 23,000 galaxies from GALEX-SDSS-WISE Legacy Catalog 2 (CH: cite). comparison to ? (only 129 nearby galaxies though) ? (one galaxy, jesus) Kriek & Conroy (2013) using stacked SEDs of medium-and broadband photometry of galaxies at 0.5 ; z ; 2 find an average slope of delta=-0.2.

paragraph on the variation of attenuation curves based on CH: what does this mean? While an empirical prescription like DEM doesn't allow explicit modeling of the complex dust-star geometry, it does a good job at mimicking it. CH: how does this compare to the literature? comparison to Narayanan et al. (2018) paper comparison to Salim et al. (2018) for SF population comparison to ? (only 129 nearby galaxies though)

paragraph on restating how we can learn about dust through DEMs based on trends we see across all simulations. summarize main findings again.

What are the limitations of DEM and how can it be improved?

- too many luminous galaxies
- color distribution isnt' perfect.

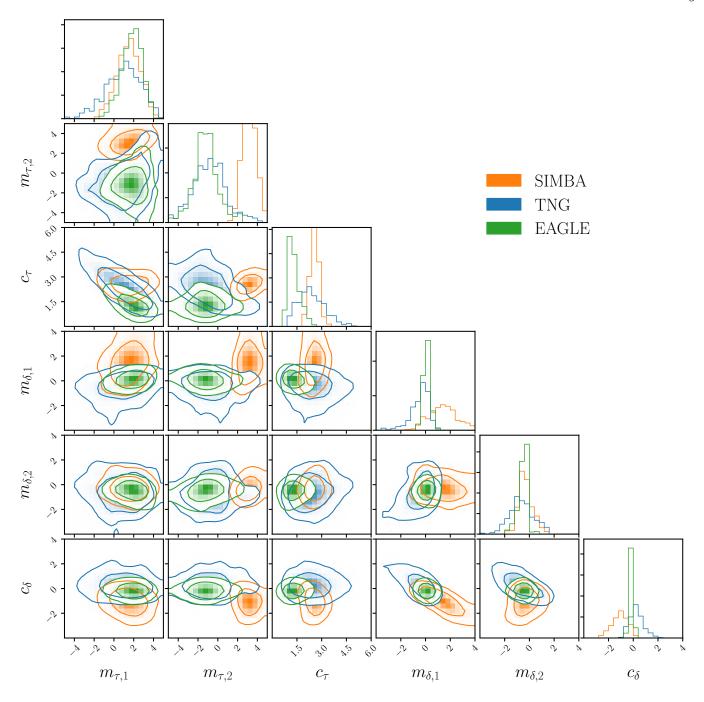


Figure 5. Posterior distributions of our DEM parameters for the SIMBA (orange), TNG (blue), and EAGLE (green) hydro simulations. The contours mark the 68% and 95% confidence intervals. We derive these posteriors using Approximate Bayesian Computation (ABC, Section 3.2) and describe the parameters in Section 3.1 and Table 1. In all simulations, dust attenuation increases for higher M_* galaxies ($m_{\tau,M_*} \sim 2$). The simulations also have consistent optical depth amplitudes (c_{τ}). However, the SFR dependence of τ_V is different among the simulations. For TNG and EAGLE, star-forming galaxies have lower τ_V ; for SIMBA quiescent galaxies have lower τ_V . Meanwhile, for the slope offset of the attenuation curve, δ , we find little M_* and SFR dependence in the simulations and that the amplitude (c_{τ}) is consistent with 0.

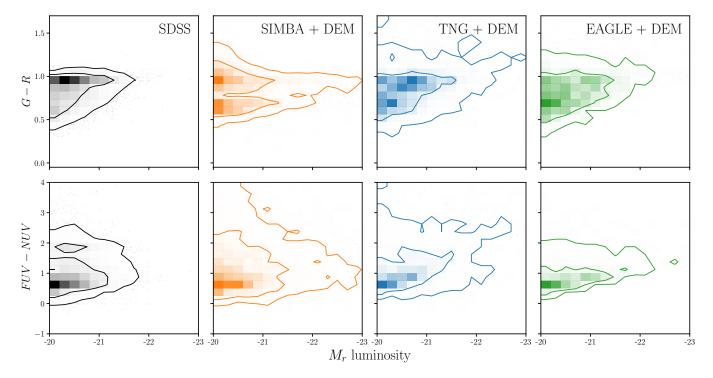


Figure 6. $(G-R)-M_r$ color-magnitude (top panels) and $(FUV-NUV)-M_r$ (bottom panels) relations predicted by the median DEM posteriors for the SIMBA (orange), TNG (blue), and EAGLE (green) hydro simulations. For comparison, we include the observables for SDSS in the left-most panel (Section 2.1). The median posterior DEMs produce dramatically different observables than when we do not include any dust prescription (Figure 3). Hence, dust must be account for when interpreting and comparing simulations. Moreover, with the DEMs, all three simulations produce observables consistent with SDSS. Since different simulations can produce reproduce observations by varying dust, dust significantly limits our ability to constrain the physical processes that go into galaxy simulations.

• There isn't a whole lot of flexibility for SFR=0 galaxies predicted by simulations and they do not agree well with observations A.

CH: How robust are our results? We fix the UV bump to reduce the number of parameters. But when run our analysis without fixing the UV bump, we find it does not impact our results. We also get no constraints on the UV bump parameters.

We rely on the slab model. But nothing changes when we use a more flexible truncated normal distribution in Appendix B. tnorm DEM model allows us to also vary the scatter of the attenuation curve

CH: How about our prior choice?

CH: If we marginalize over dust, can we learn anything about galaxy evolution from the simulation?

Are there observables that hydro sims + DEMs cannot reproduce? What does that say about the hydro sims?

What observables are unaffected by DEMs? We should chase those observables.

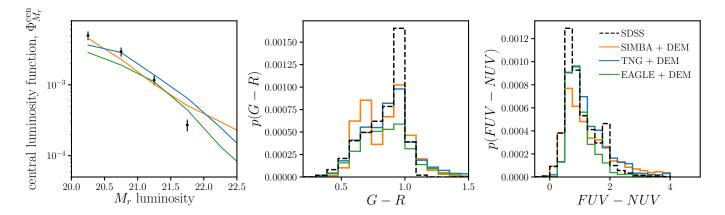


Figure 7. Comparison of the observables predicted by the simulations with the posterior DEM.

We clearly have to becareful with overinterpreting hydro sims because modifying dust allows us to reproduce whatever we want.

• Should we bother calibrating our empirical and semi-analytic models to hydrodynamic simulations when the hydro sims also require marginalizing over dust parameters? Does this mean that if our goal is to make realistic mocks, we can be relatively careless about

CH: What are some applications for DEMs? Realistic mock catalogs that reproduce observations in observable-space rather than physical parameter space.

5. SUMMARY

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APPENDIX

A. RESOLUTION EFFECTS

Figure demonstrating imprint SFR=0 leave on the observable space and how we deal with them so we can ignore them...

B. BEYOND THE SLAB DEM

A major assumption of our fiducial DEM is that we sample the amplitude of attenuation from the slab model. The slab model makes the simplifying assumption that dust in galaxies are in a slab-like geometry and illuminated by the stellar radiation source (Somerville & Primack 1999). Then, for a given τ_V , the attenuation depends solely on the orientation of the galaxy. This simplification, ignores any complexities in the star-to-dust geometry that impact the shape of the attenuation curve (Witt

Gordon 1996, 2000, Seon Drain 2016).

Besides its simplifications, the slab model predicts A_V distribution with significant differences than the A_V distributions measured from observations. In Figure 8, we compare the A_V distribution predicted by the slab model (black) to the A_V distribution of star-forming galaxies in our SDSS sample (blue). The A_V values are derived using SED fitting from the ? MPA-JHU catalog and how are the SF galaxies classified. The slab model A_V values are derived using Eq. 3 and 4 with M_* s and SFRs from the same SDSS sample and the inclinations, i, are uniformly sampled over the range $[0, \pi/2]$. With $\{m_{\tau,1}, m_{\tau,2}, c_{\tau}\}$ chosen to reproduce the observed A_V distribution, the slab model can reproduce the overall shape. However, it predicts an extended high A_V tail not found in observations.

Given these shortcomings of the slab model, we want to ensure that our results do not hinge on the slab model. Modeling the star-to-dust geometries with increased complexities, however, would involve expensive hydrodynamic simulations and dust radiative transfer calculations (e.g. Narayanan et al. 2018)jonsson2006, rocha2008, natale2015,hayward smith2015,hou2017,trayford2020. We instead take an empirical approach and implement a flexible model for sampling A_V based on a truncated normal distribution:

 $A_V \sim \mathcal{N}_T(\mu_{A_V}, \sigma_{A_V}) = \frac{\mathcal{N}(\mu_{A_V}, \sigma_{A_V})}{1 - \Phi\left(-\frac{\mu_{A_V}}{\sigma_{A_V}}\right)}.$ (B1)

Here, \mathcal{N} is the standard normal distribution and $\Phi(x) = \frac{1}{2} \left(1 + \operatorname{erf}(x/\sqrt{2}) \right)$ is the cumulative distribution function of \mathcal{N} . μ_{A_V} and σ_{A_V} are the mean and variance of the truncated normal distribution. Similar to Eq. 4, we allow μ_{A_V} and σ_{A_V} to depend on the physical properties of galaxies:

$$\mu_{A_V} = m_{\mu,1} (\log M_* - 10.) + m_{\mu,2} \log SFR + c_\mu$$
 (B2)

$$\sigma_{A_V} = m_{\sigma,1}(\log M_* - 10.) + m_{\sigma,2}\log SFR + c_{\sigma}.$$
 (B3)

The A_V distribution from our truncated normal (orange dashed) closely reproduces the observed SDSS A_V distribution (Figure 6). N_T is able to reproduce the overall skewness but unlike the slab model, it does not have a long high A_V tail. With more free parameters and a functional form that closely resembles the observed A_V distribution, the truncated normal model provides a flexible alternative to the slab model and we include it in our analysis.

REFERENCES

Abazajian K. N., et al., 2009, The Astrophysical Journal Supplement Series, 182, 543

Blanton M. R., Roweis S., 2007, The Astronomical Journal, 133, 734

Blanton M. R., et al., 2005, The Astronomical Journal, 129, 2562

Calzetti D., 2001, New Astronomy Reviews, 45, 601

Chabrier G., 2003, Publications of the Astronomical Society of the Pacific, 115, 763 Hahn C., Vakili M., Walsh K., Hearin A. P., Hogg
D. W., Campbell D., 2017, Monthly Notices of the Royal Astronomical Society, 469, 2791

Ishida E. E. O., et al., 2015, Astronomy and Computing, 13, 1

Kriek M., Conroy C., 2013, The Astrophysical Journal Letters, 775, L16

Narayanan D., Conroy C., Davé R., JohnsonB. D., Popping G., 2018, The Astrophysical Journal, 869, 70 TODO

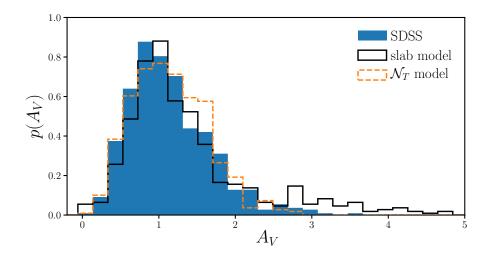


Figure 8. Comparison of A_V distribution of SDSS star-forming galaxies (blue) to predictions from the slab model (Eq. 3; black). detail on how SDSS SF galaxies are classified. The slab model assumes that there's a slab of dust in front of a galaxy. We use $\tau_V = 2$ for the slab model above. Regardless of τ_V , however, the slab model predicts a significantly more asymmetric and peaked A_V distribution than observations. Given this disagreement, we include in our analysis a DEM with an empirical prescription for A_V based on a truncated normal distribution, which better reproduce the observed A_V distribution (Section B).

Noll S., Burgarella D., Giovannoli E., Buat V., Marcillac D., Muñoz-Mateos J. C., 2009, Astronomy and Astrophysics, 507, 1793

Salim S., Narayanan D., 2020, arXiv:2001.03181 [astro-ph]

Salim S., Boquien M., Lee J. C., 2018, The Astrophysical Journal, 859, 11

Somerville R. S., Primack J. R., 1999, Monthly Notices of the Royal Astronomical Society, 310, 1087 Somerville R. S., Gilmore R. C., Primack J. R., Domínguez A., 2012, Monthly Notices of the Royal Astronomical Society, 423, 1992

Tinker J., Wetzel A., Conroy C., 2011, preprint, 1107, arXiv:1107.5046

Tinker J. L., Hahn C., Mao Y.-Y., Wetzel A. R., Conroy C., 2018, Monthly Notices of the Royal Astronomical Society, 477, 935

Trayford J. W., Lagos C. d. P., Robotham A.
S. G., Obreschkow D., 2020, Monthly Notices of the Royal Astronomical Society, 491, 3937
Tress M., et al., 2018, Monthly Notices of the

Royal Astronomical Society, 475, 2363