

Neutrons and X-rays

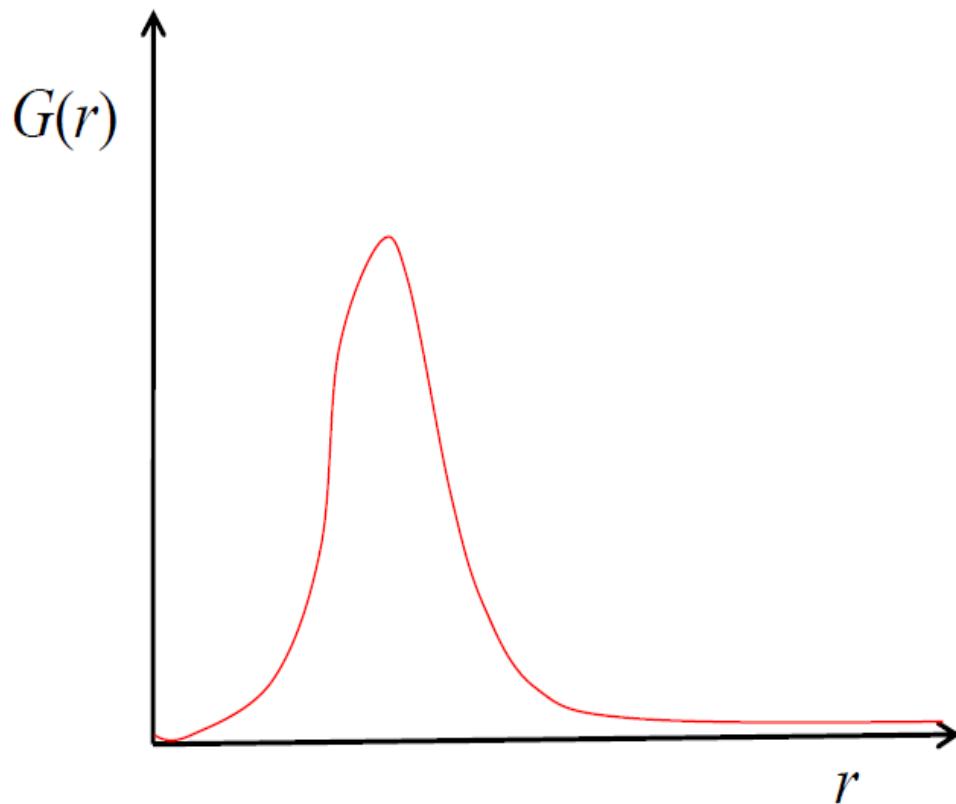


Andrew Harrison, Diamond Light Source
14th Oxford School on Neutron Scattering - 2015

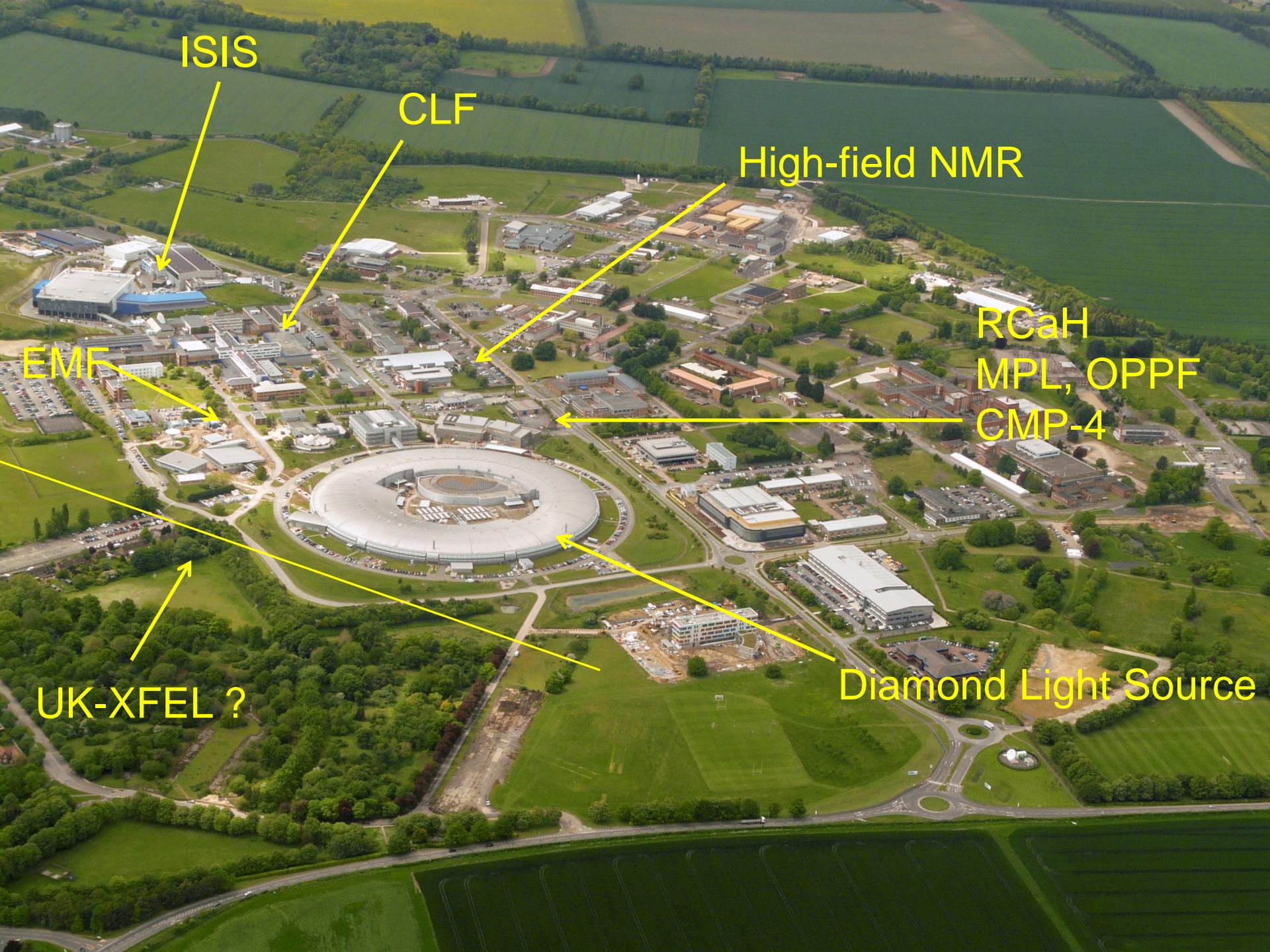




Neutron source – synchrotron distribution



Acknowledgement: Jon Goff



ISIS

CLF

High-field NMR

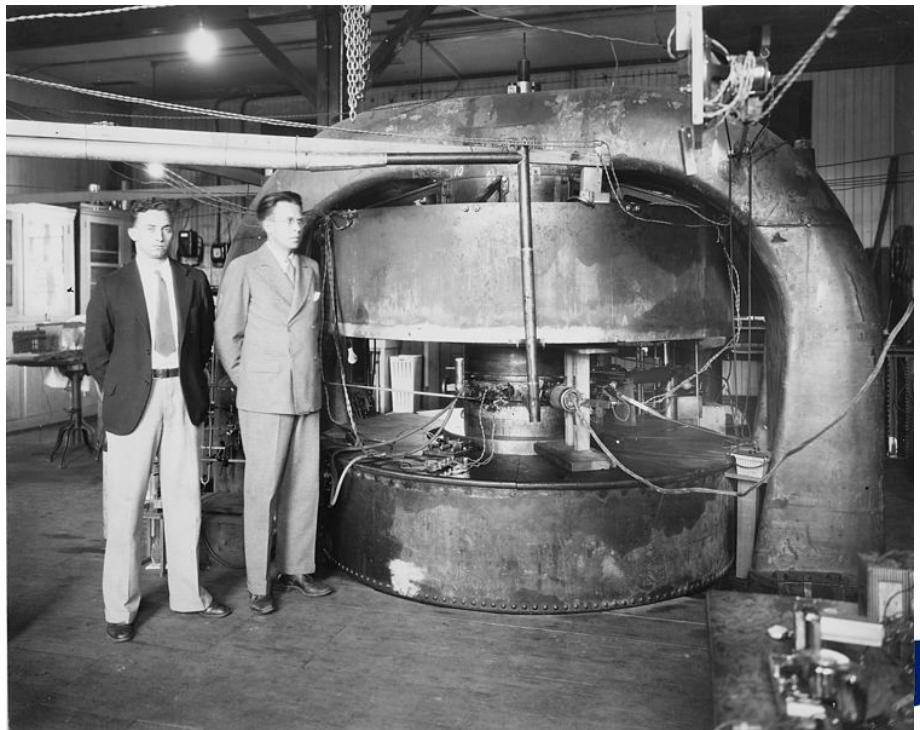
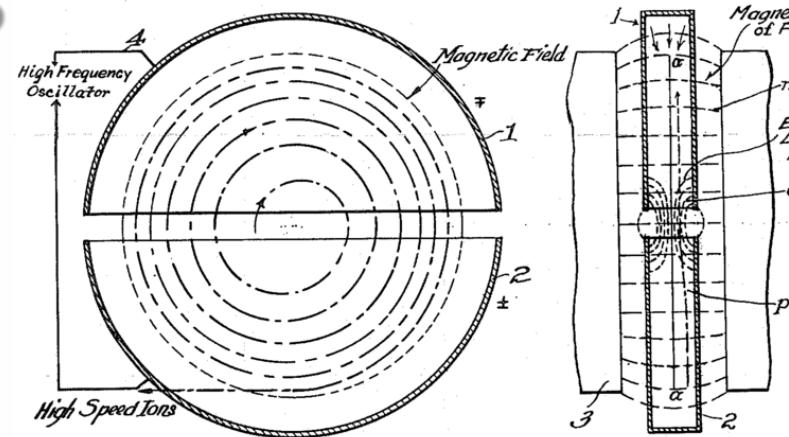
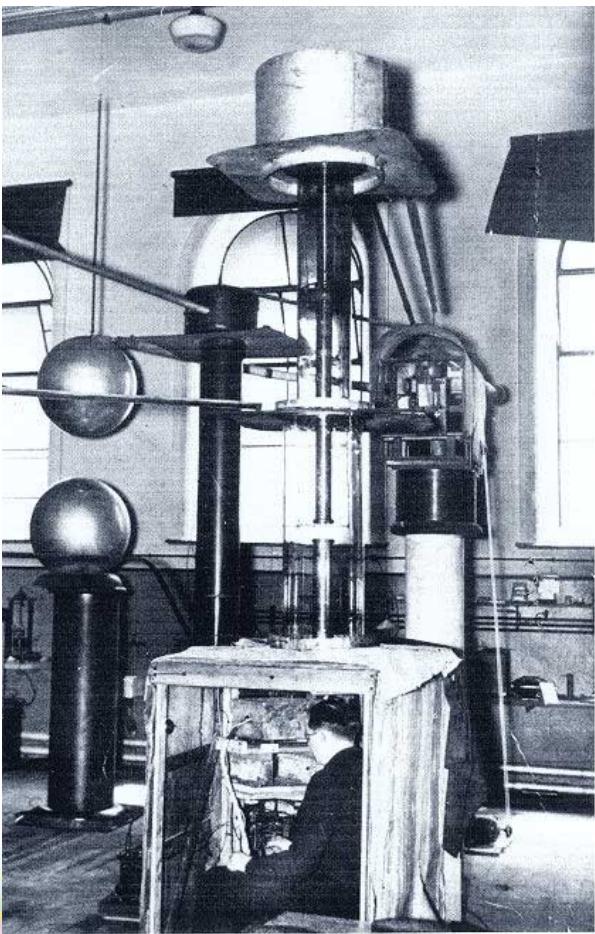
EMF

RCaH
MPL, OPPF
CMP-4

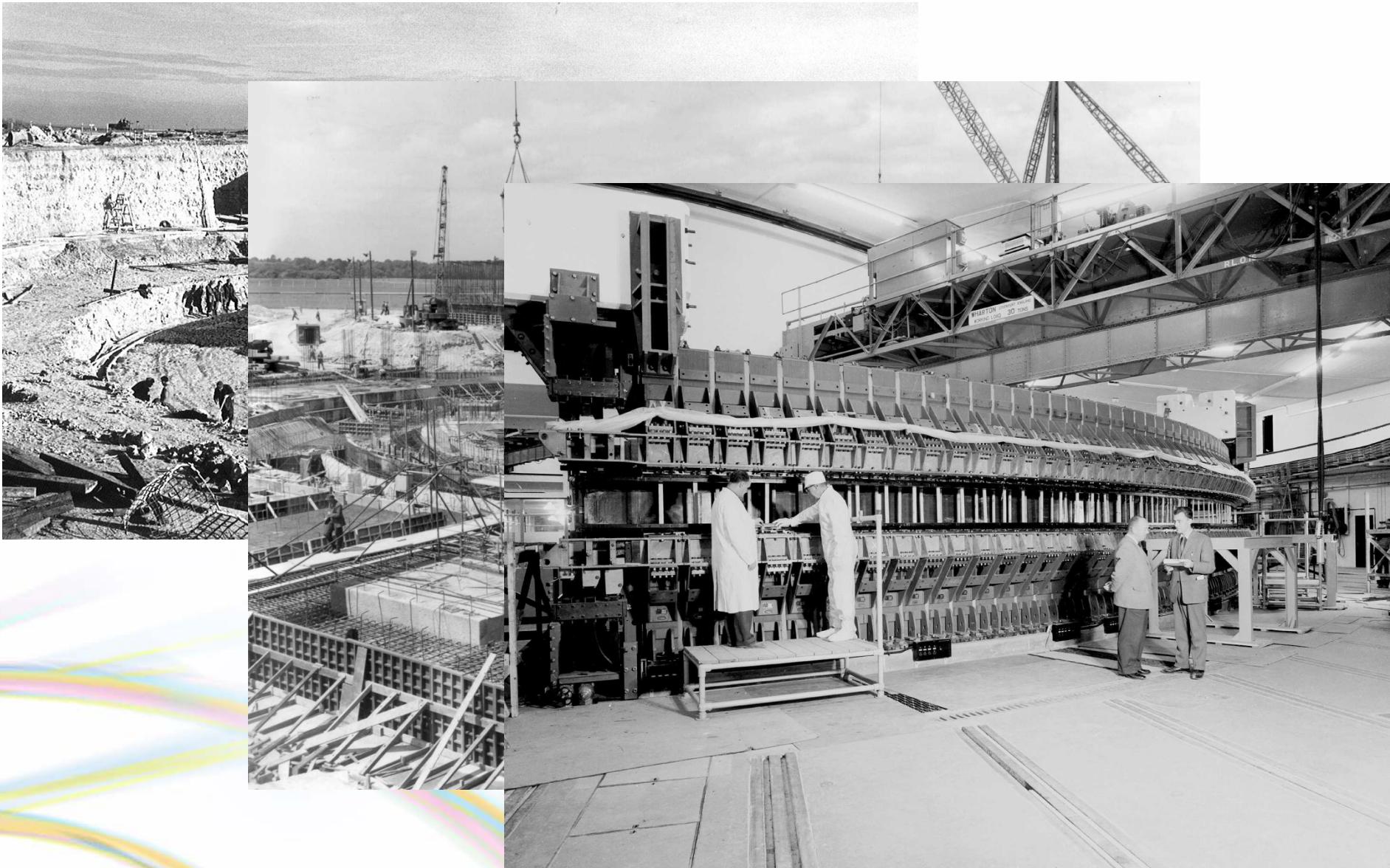
UK-XFEL ?

Diamond Light Source

Origins



Ever increasing circles

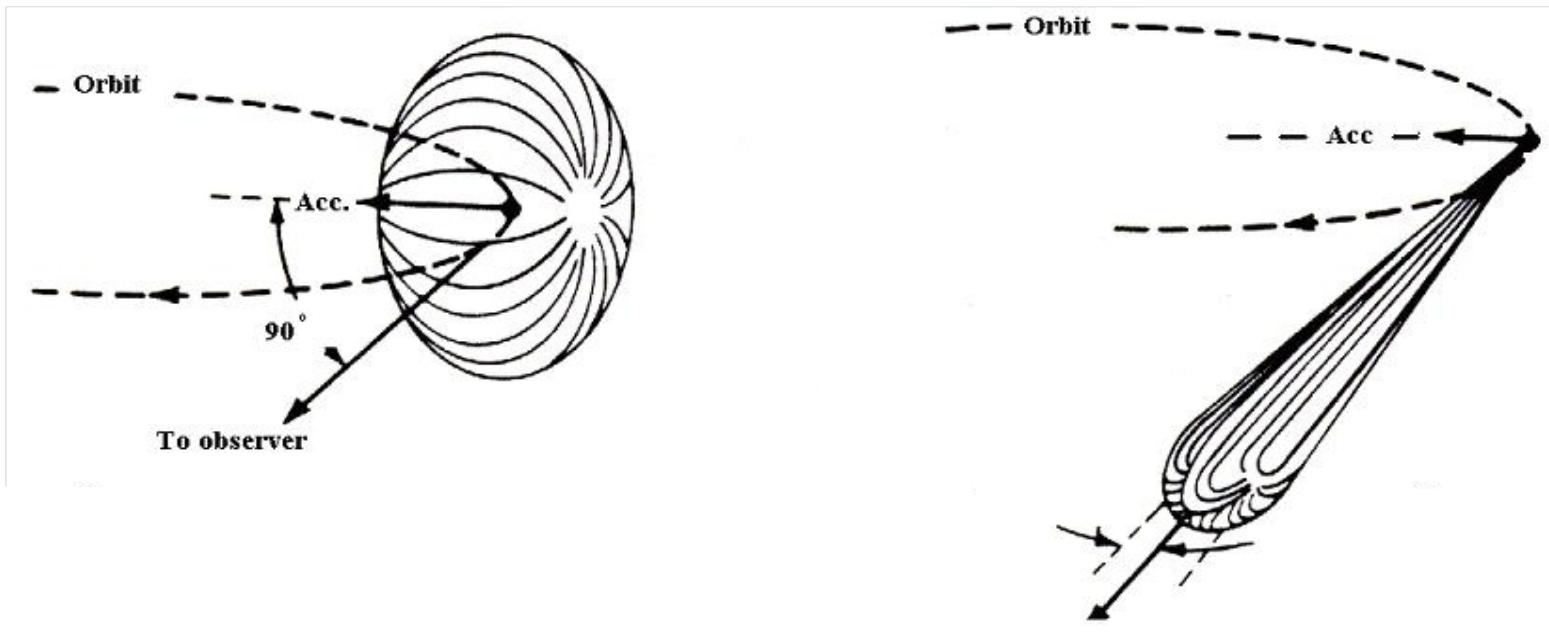


Ever increasing circles



ond

From Maxwell to synchrotrons



$$V \sim C$$

First observation ?

Radiation from Electrons in a Synchrotron

F. R. ELDER, A. M. GUREWITSCH, R. V. LANGMUIR,

AND H. C. POLLACK

Research Laboratory, General Electric Company,

Schenectady, New York

May 7, 1947

HIGH energy electrons which are subjected to large accelerations normal to their velocity should radiate electromagnetic energy.¹⁻⁴ The radiation from electrons in a betatron or synchrotron should be emitted in a narrow cone tangent to the electron orbit, and its spectrum should extend into the visible region. This radiation has now been observed visually in the General Electric 70-Mev synchrotron.⁵ This machine has an electron orbit radius of 29.3 cm and a peak magnetic field of 8100 gauss. The radiation is seen as a small spot of brilliant white light by an observer looking into the vacuum tube tangent to the orbit and toward the approaching electrons. The light is quite bright when the x-ray output of the machine at 70 Mev is 50 roentgens per minute at one meter from the target and can still be observed in daylight at outputs as low as 0.1 roentgen.

The synchrotron x-ray beam is obtained by turning off the r-f accelerating resonator and permitting subsequent changes in the field of the magnet to change the electron orbit radius so as to contract or expand the beam to suitable targets. If the electrons are contracted to a target at successively higher energies, the intensity of the light radiation is observed to increase rapidly with electron energy. If, however, the electrons are kept in the beam past the

peak of the magnetic field and then expanded to a target, the intensity of the radiated light appears to be independent of the energy at which the electrons are removed from the beam. This is to be expected, for in a given machine the radiation is proportional to the fourth power of the electron energy. The light radiation is not observed if the beam is contracted before its energy is about 30 Mev. When the electron beam has been accelerated to the peak of the magnetic field and then decelerated to low energy, a rough measurement of the phase angle over which the light was visible gave a value of 90-100 degrees. The light was viewed through a slotted disk rotating at synchronous speed.

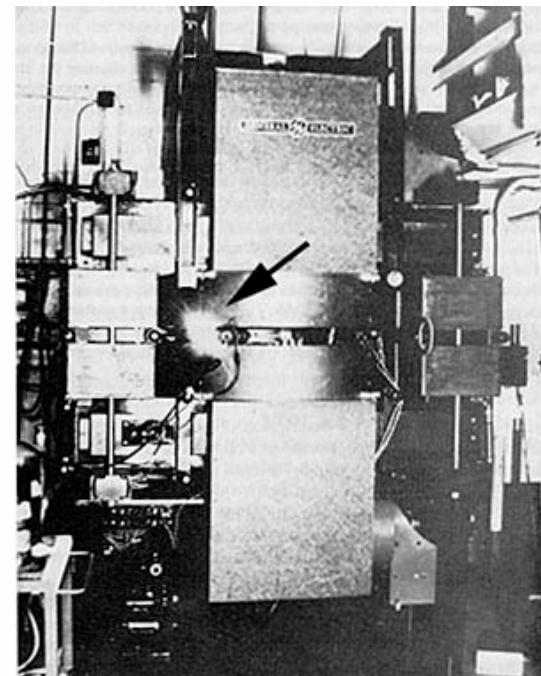
If the r-f resonator is turned off a short time before the peak of the magnetic field, the electron beam slowly contracts to a radius just larger than that of the interior

target and then expands as the magnetic field decreases. In this case, the observer no longer sees a single point of light but a short line with extension in the plane of the orbit.

The light emitted from the beam is polarized with the electric vector parallel to the plane of the electron orbit. It disappears as the observer rotates a piece of Polaroid before the eye through ninety degrees. An investigation of the spectral distribution of the energy is in progress and will be reported.

This work has been supported by the Office of Naval Research under contract N5ori-178.

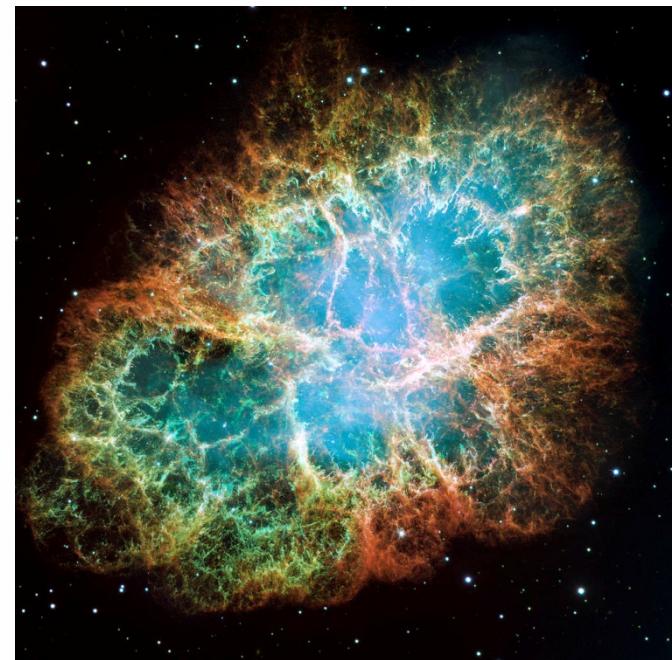
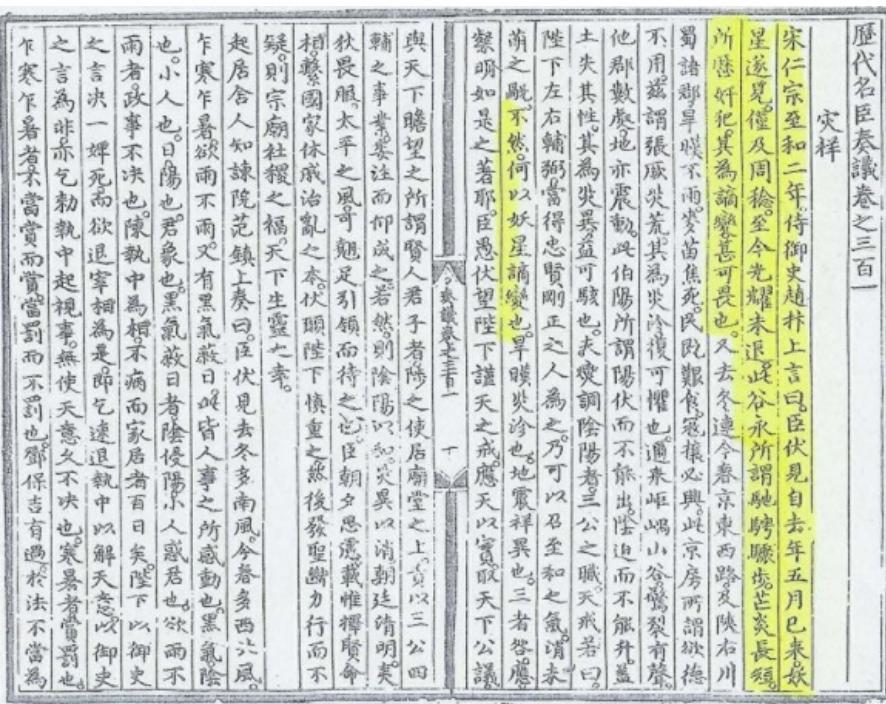
¹ D. Iwanenko and I. Pomeranchuk, Phys. Rev. 65, 343 (1944).
² J. P. Blewett, Phys. Rev. 69, 87 (1946).
³ L. I. Schiff, Rev. Sci. Inst. 17, 6 (1946).
⁴ J. S. Schwinger, Phys. Rev. 70, 798 (1946).
⁵ H. C. Pollack *et al.*, Phys. Rev. 70, 798 (1946).



1946

 diamond

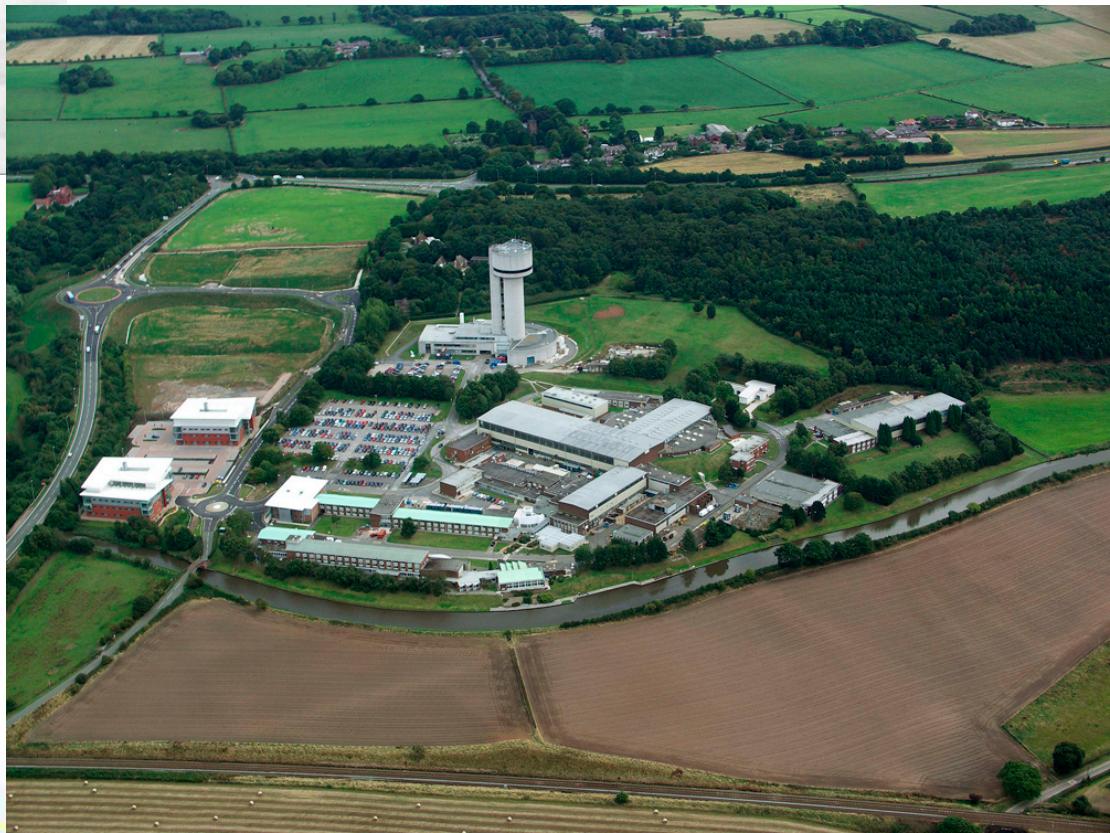
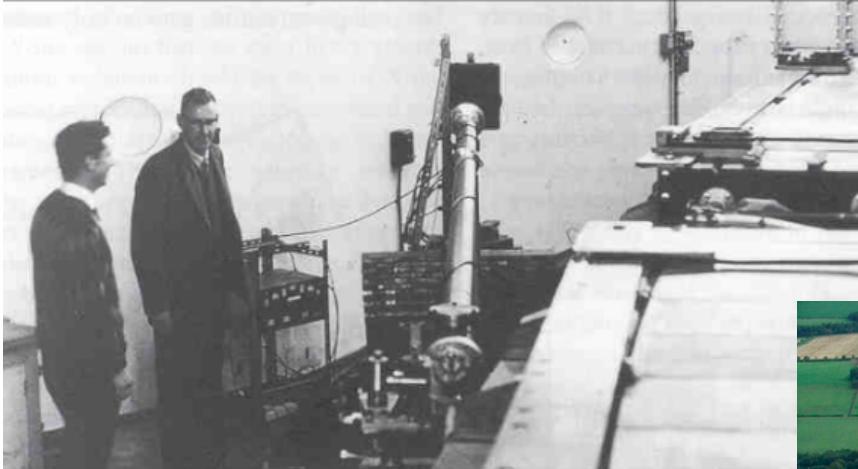
First observation ?



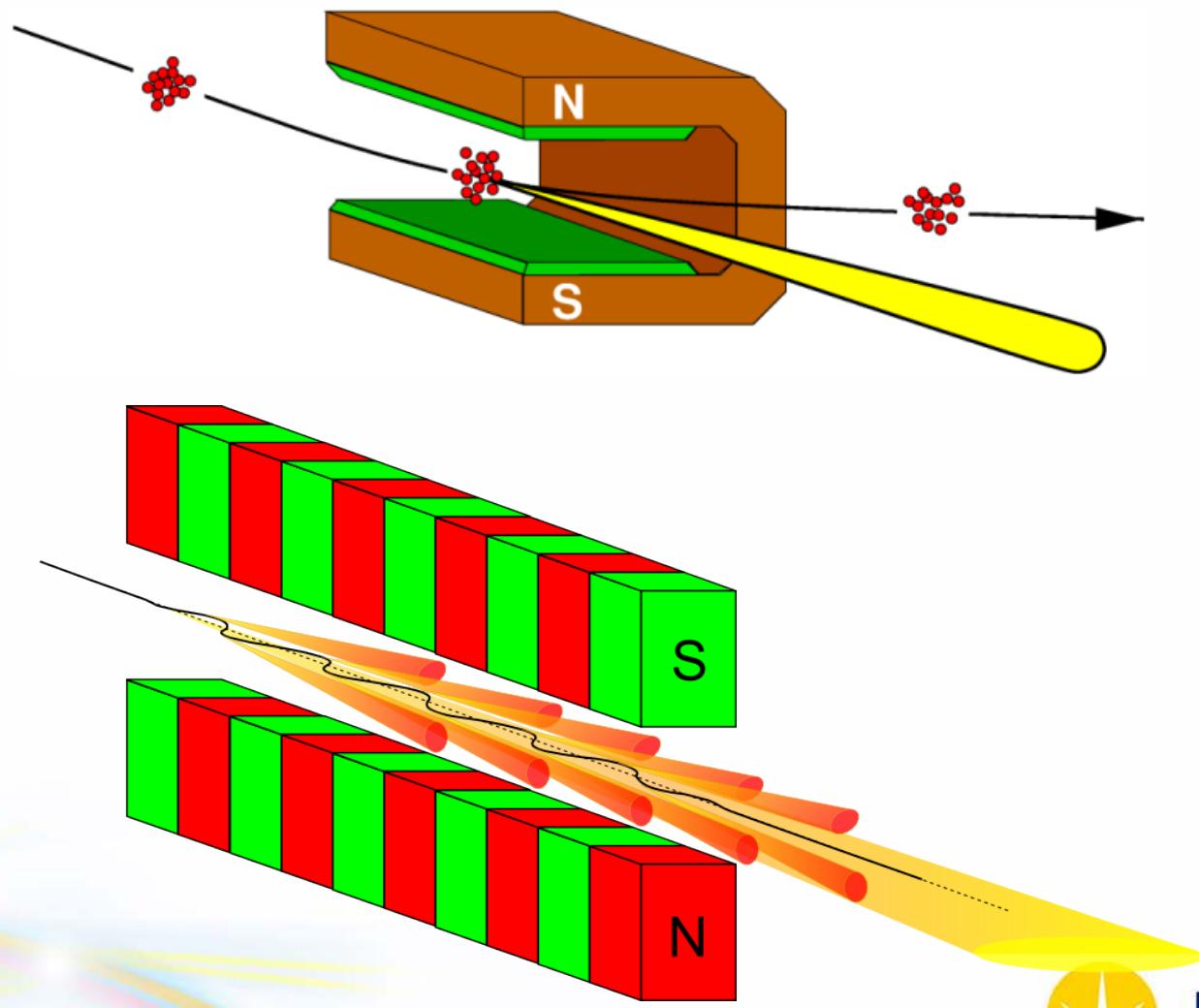
1054

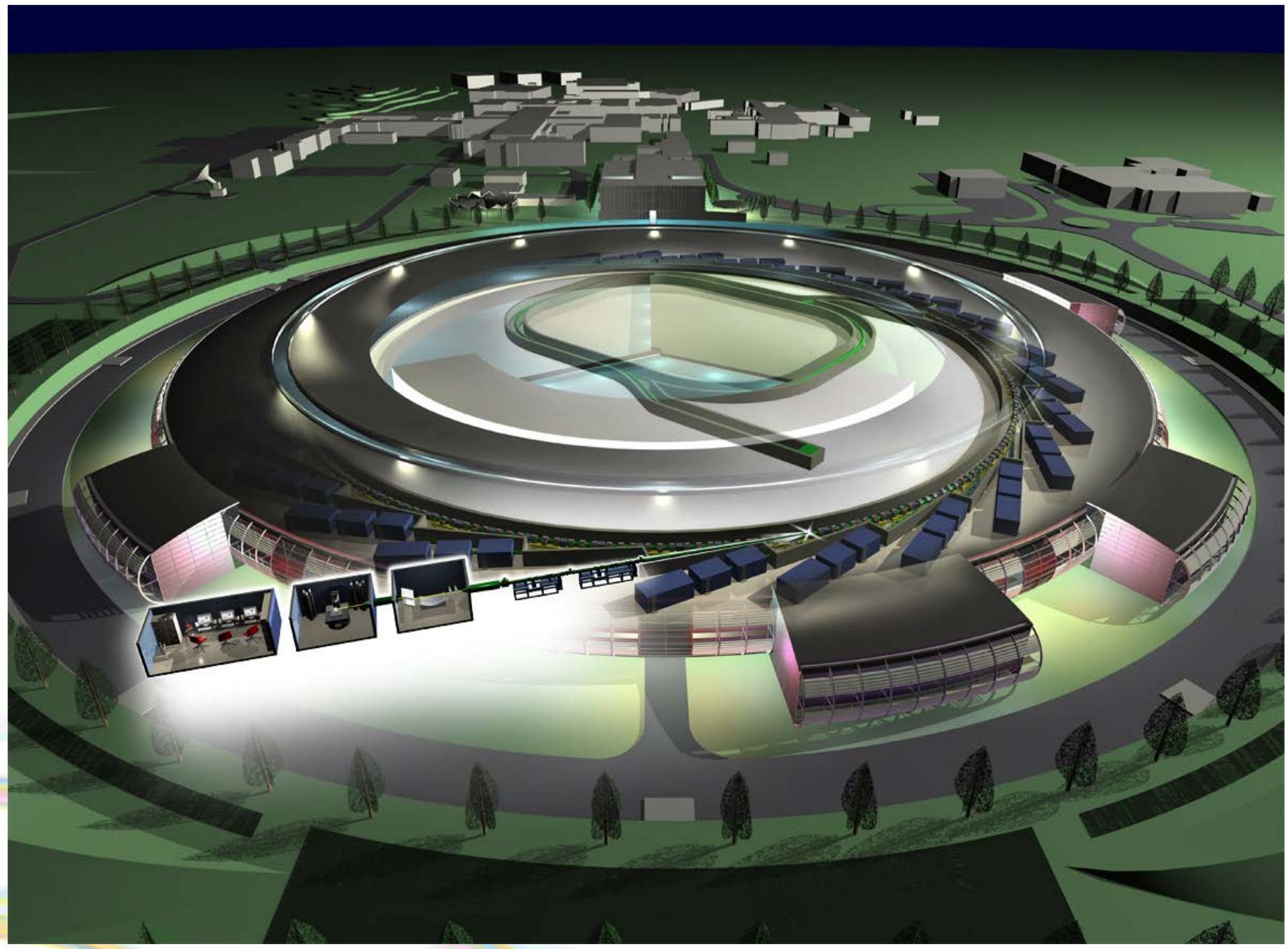


Jumping a generation (2nd incarnation)

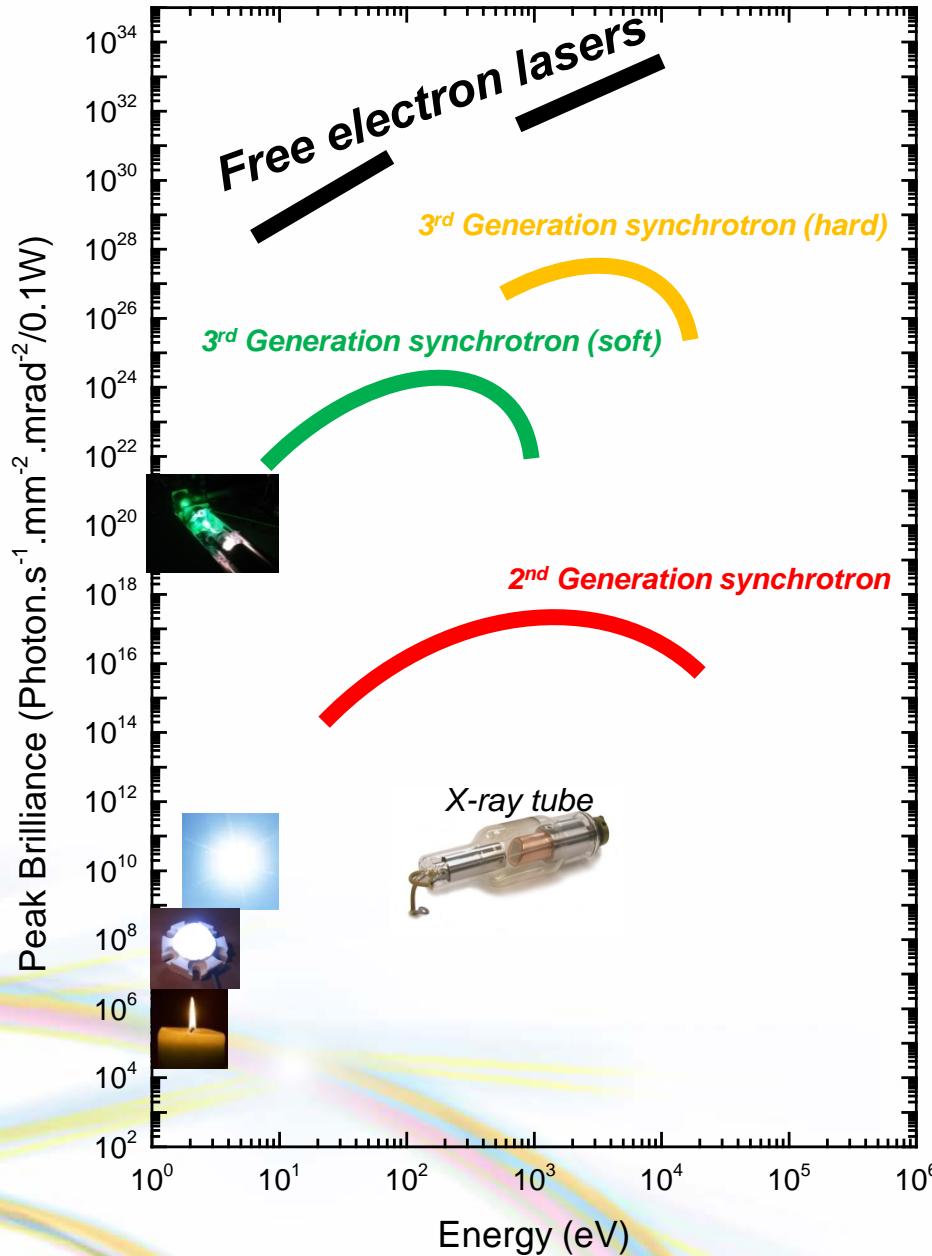


2nd and 3rd generation sources





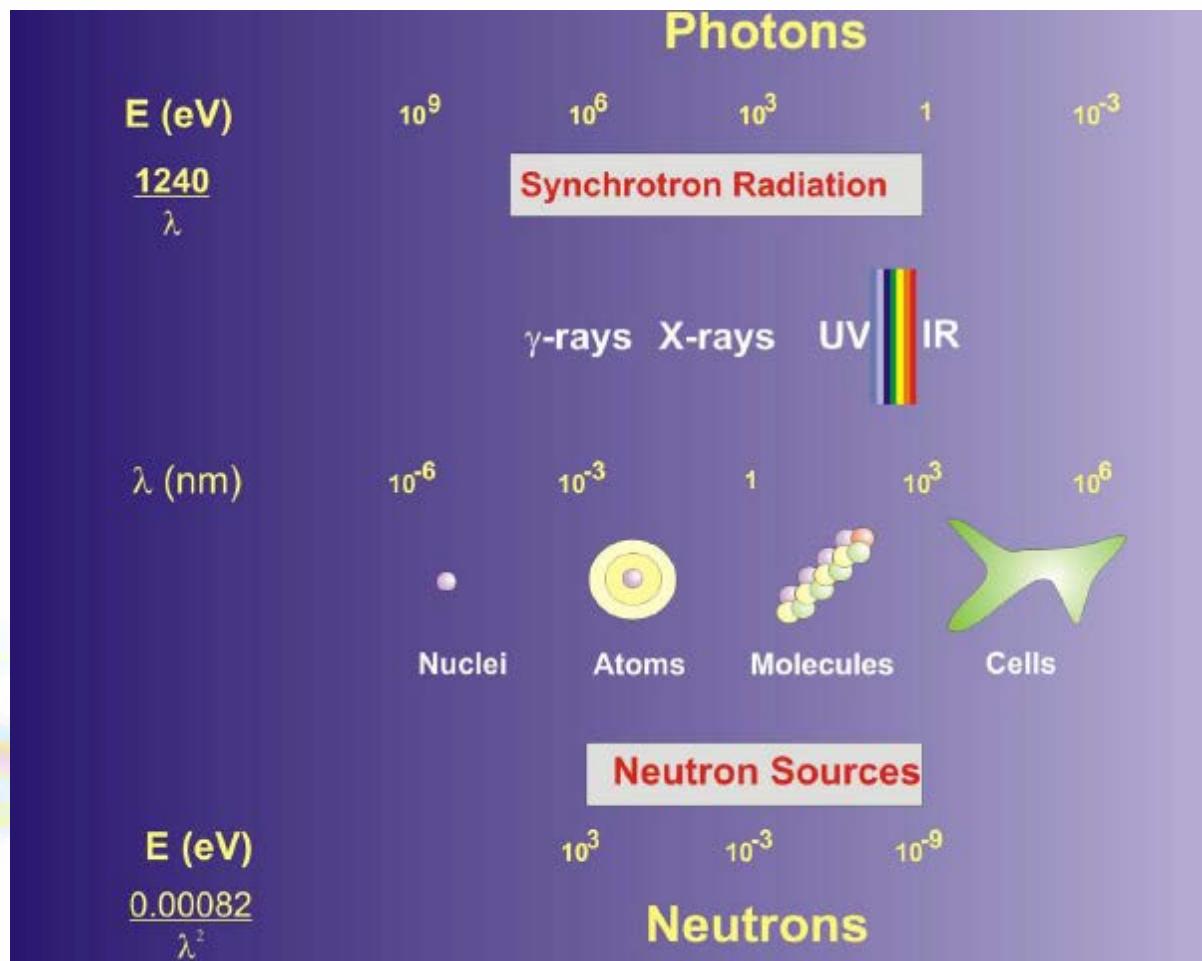
Not just ordinary light...



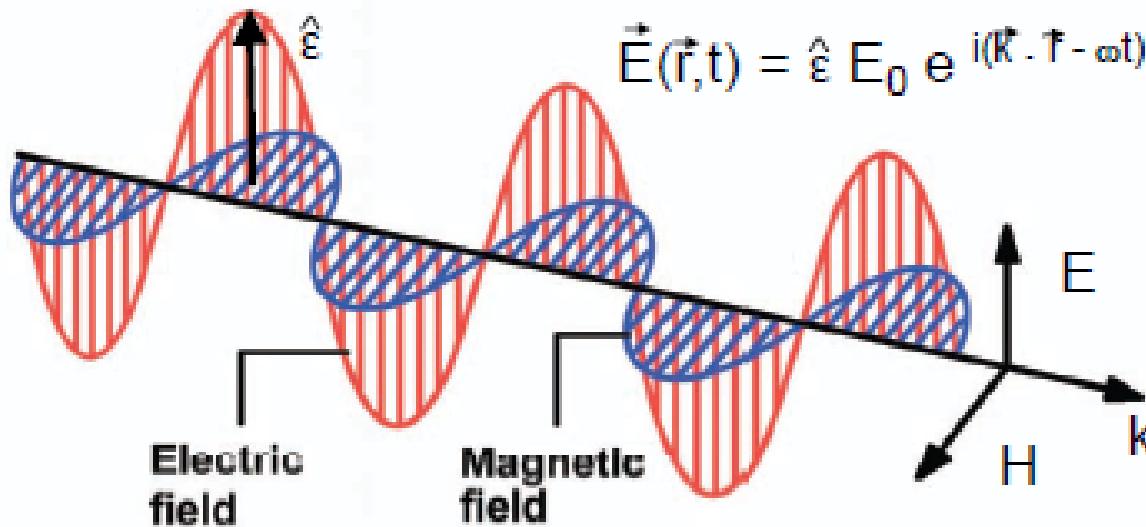
- Brilliant
- Tuneable energy – from IR to X-ray
- Focussing without loss of intensity – 100 mm to 10 nm
- High degree of coherence

Properties of neutrons and X-rays

	Neutron	X-ray
Charge	0	0
Mass/kg	1.675×10^{-27}	0
Spin	$\frac{1}{2}$	1
Magnetic moment	$-1.913\mu_N$	0 (Nuclear magneton $\mu_N \approx 5.051 \times 10^{-27} J/T$)



EM wave characteristics



Wavelength

λ

Wavenumber

$k=2\pi/\lambda$

Frequency

$\nu=\omega/2\pi$

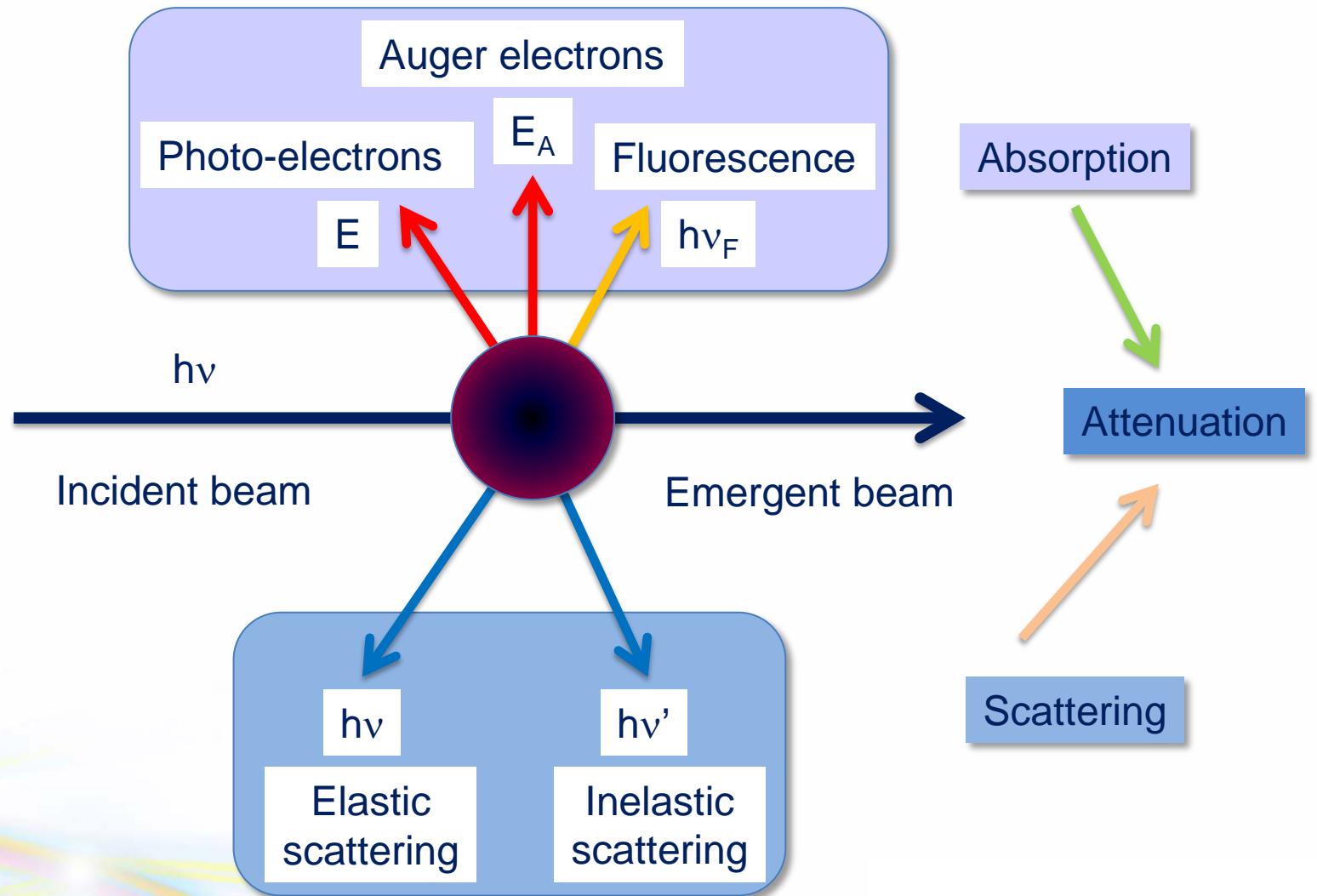
Polarisation

$\hat{\epsilon}$

Energy

$h\nu = \hbar\omega = 12.398 \text{ [keV]}/\lambda[\text{\AA}]$

X-ray – matter interactions



X-ray absorption/emission processes

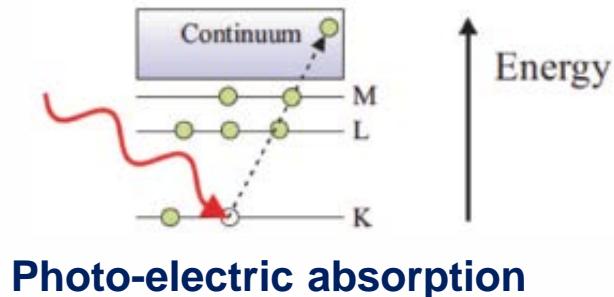
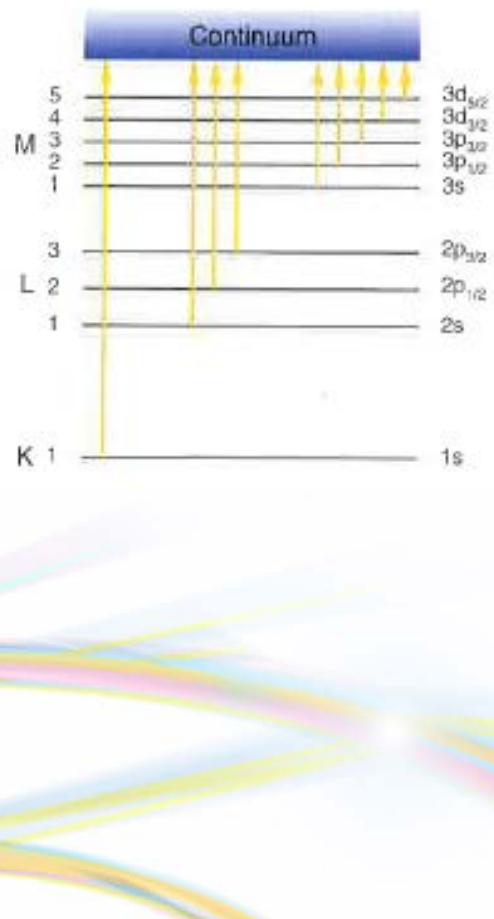
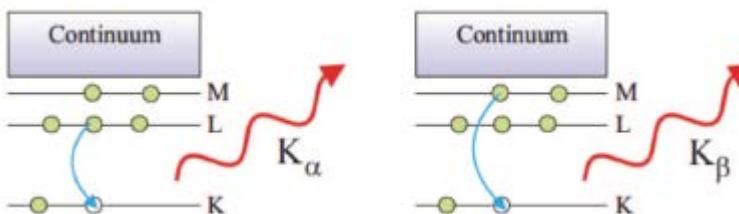
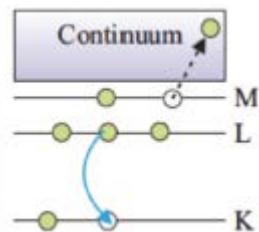


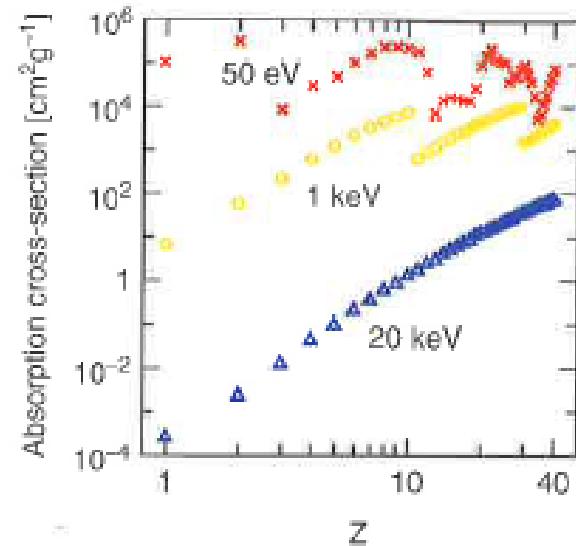
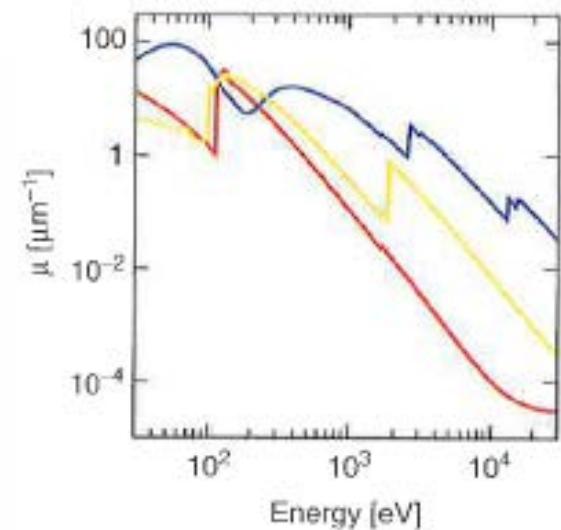
Photo-electric absorption



Fluorescent emission

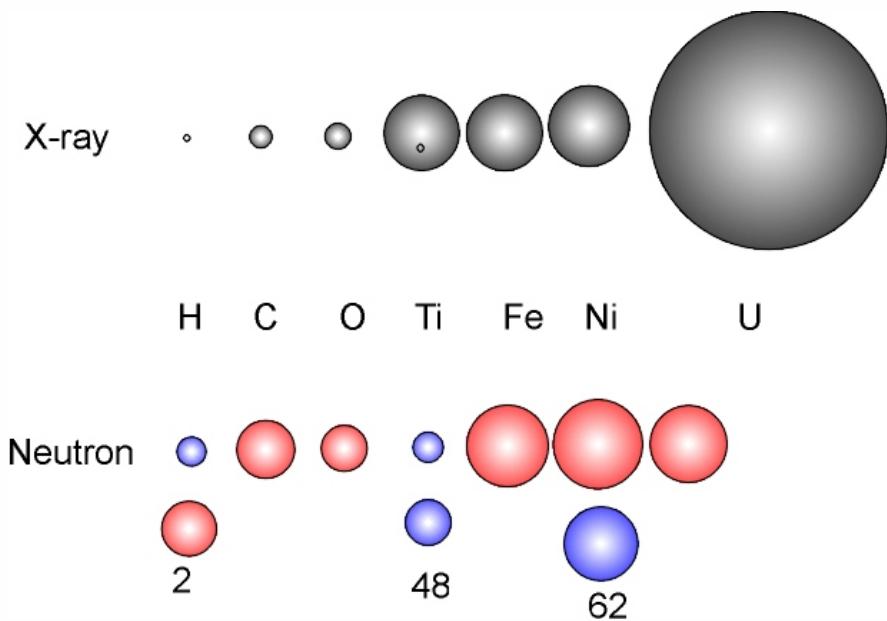
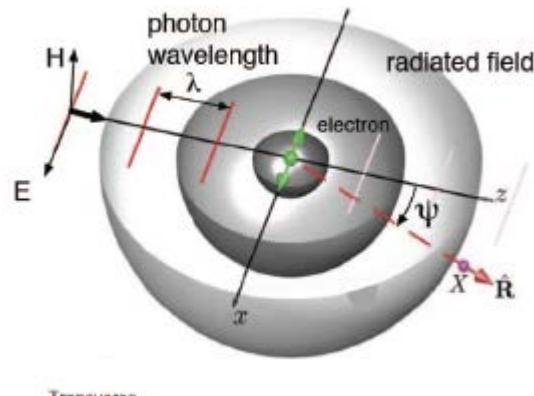


Auger electron emission



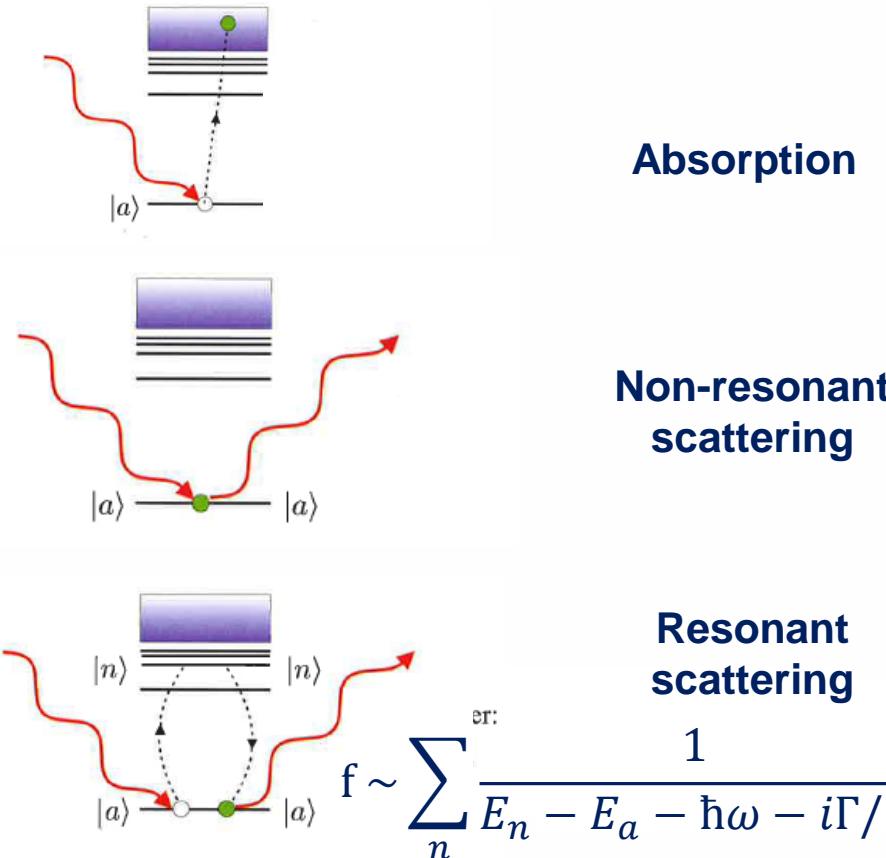
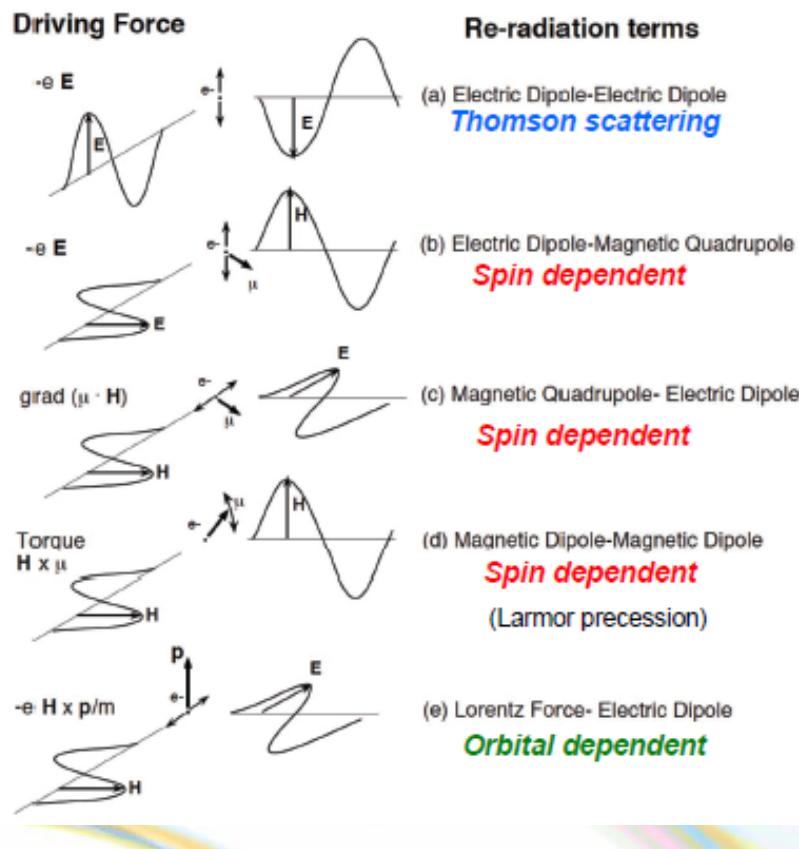
X-ray scattering - elastic

- Leading term is the (Thomson) scattering by electrons
 - The electric field of the incident EM wave drives motion of the electrons and produces spherical waves of the same energy which may interfere to produce a diffraction pattern
 - Elastic scattering intensity is proportional to Z^2



X-ray scattering - elastic

- Secondary terms are weaker
 - Various (non-resonant) magnetic terms 10^5 - 10^7 x weaker
 - If the energy of the incident photons is tuned to resonate with (virtual) transitions to intermediate states very strong amplification of the photon scattering process may occur (see later for RIXS too)



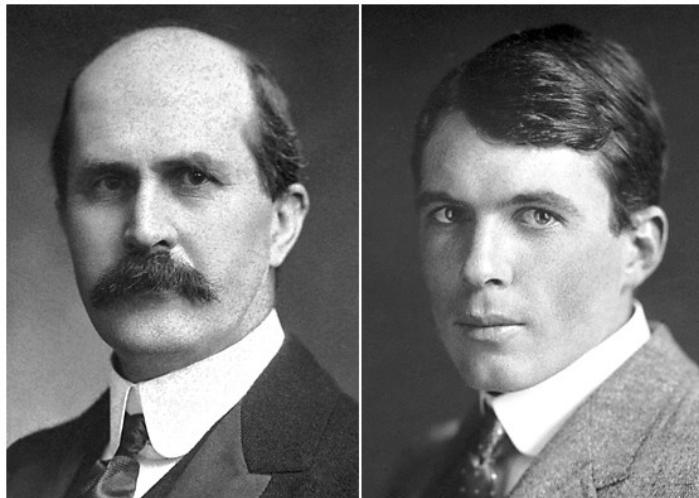
Diversity - imaging, diffraction, spectroscopy



Wilhelm Röntgen



Max von Laue



William and Lawrence Bragg



Henry Moseley

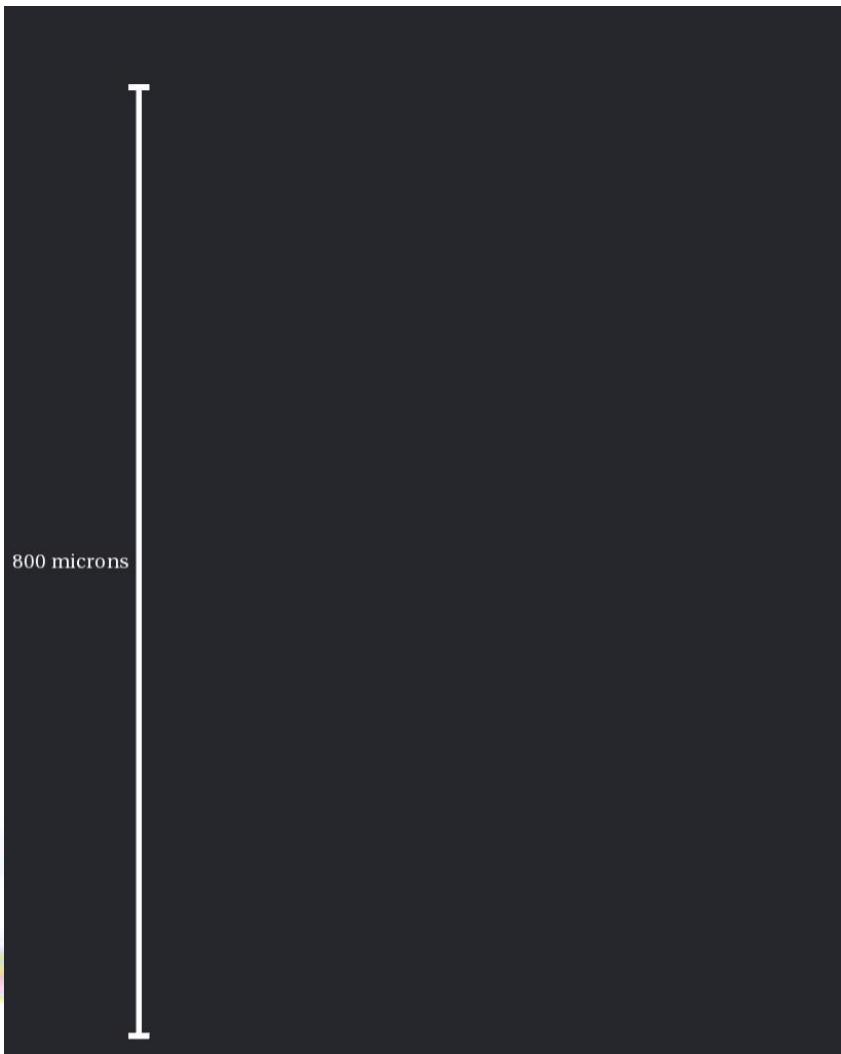


Origins of imaging



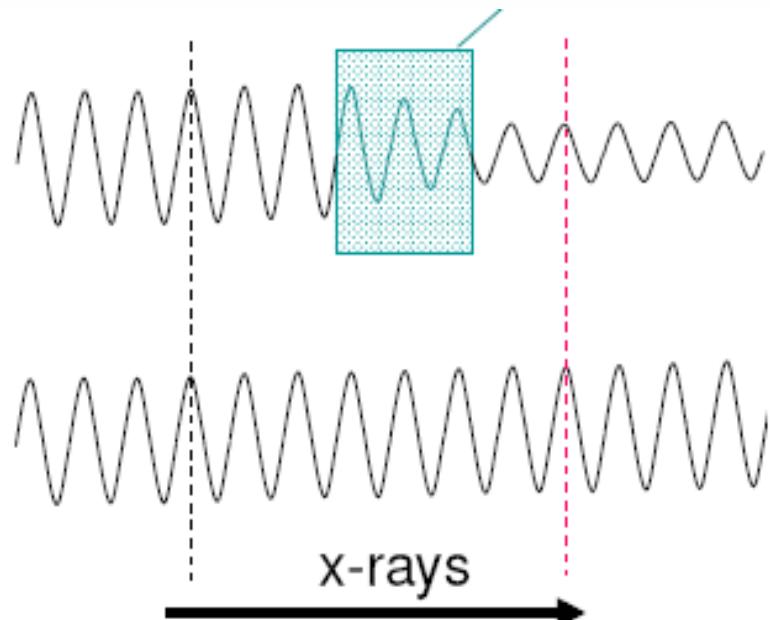
I have seen my death!
(Anna Roentgen, 1900)

Finer, faster, three dimensional ...



...more subtle

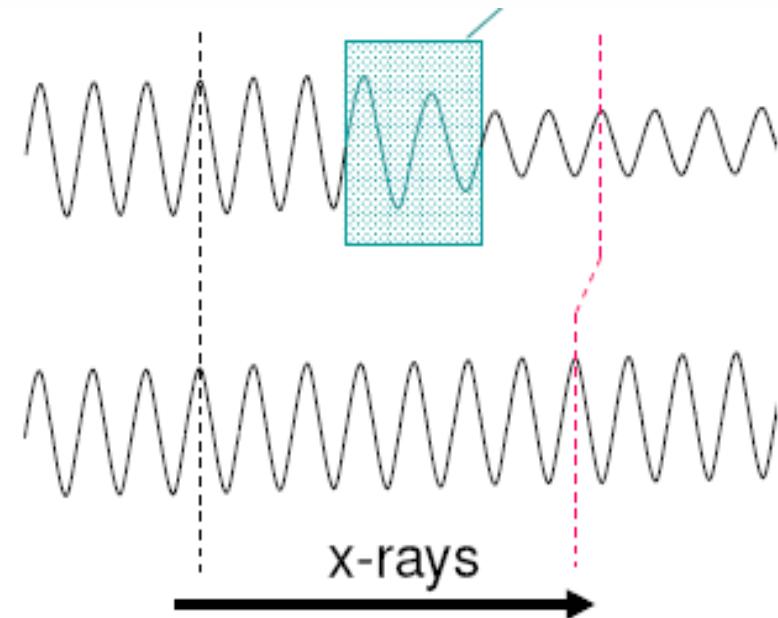
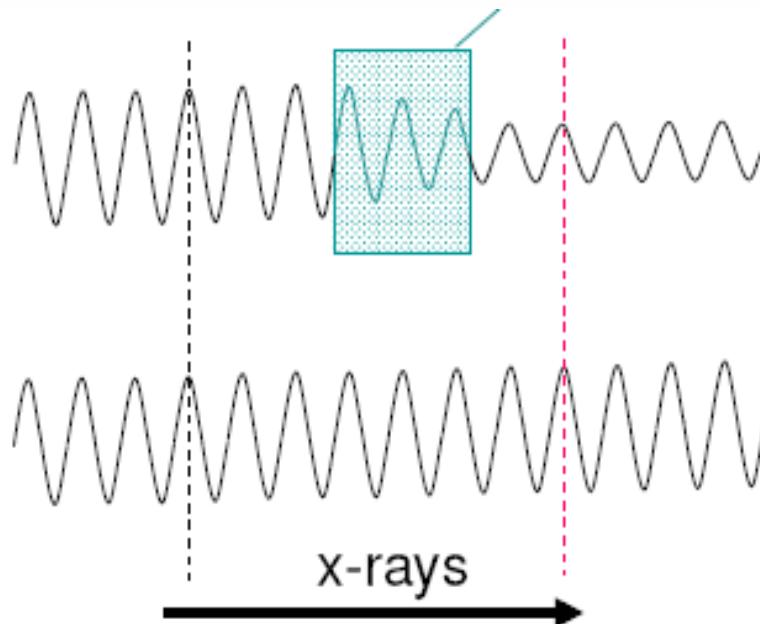
Contrast very poor
if absorption
contrast low



...more subtle

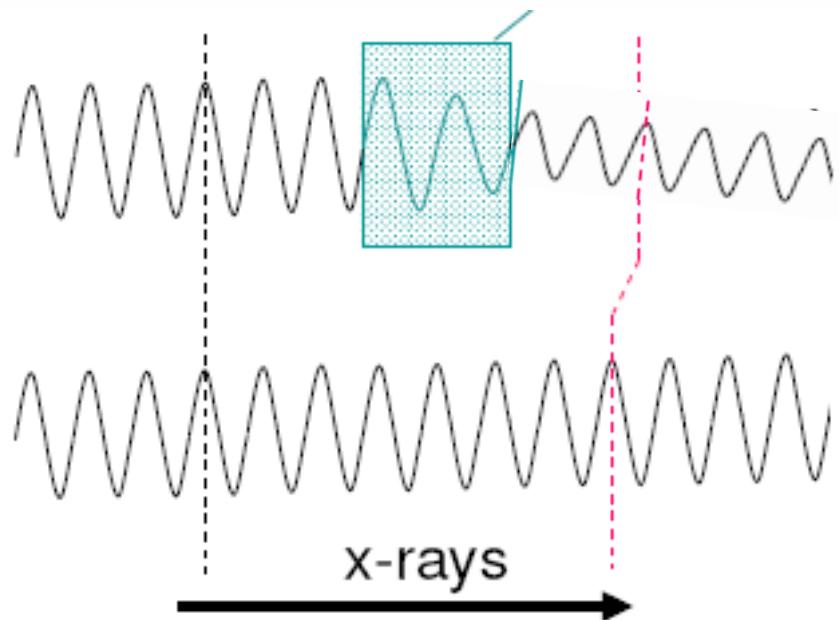
Contrast very poor
if absorption
contrast low

If refractive index
different, phase
shift occurs

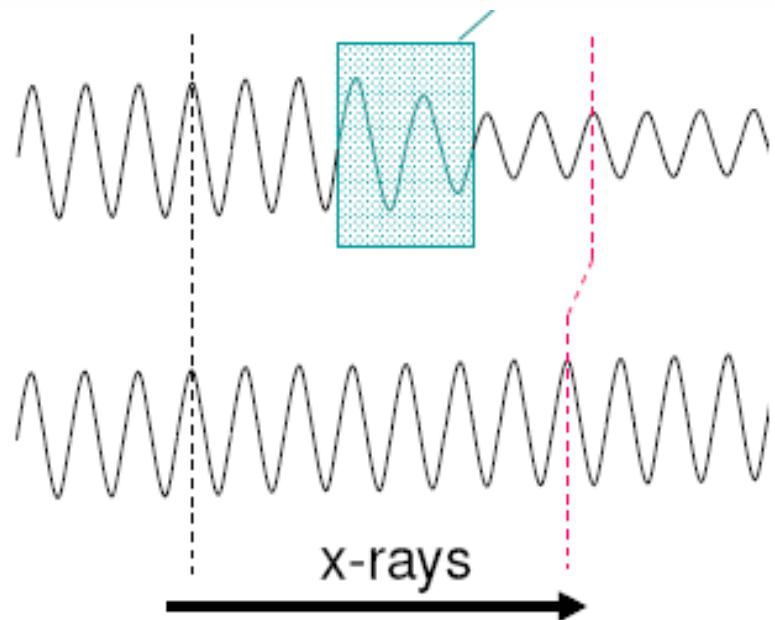


...more subtle

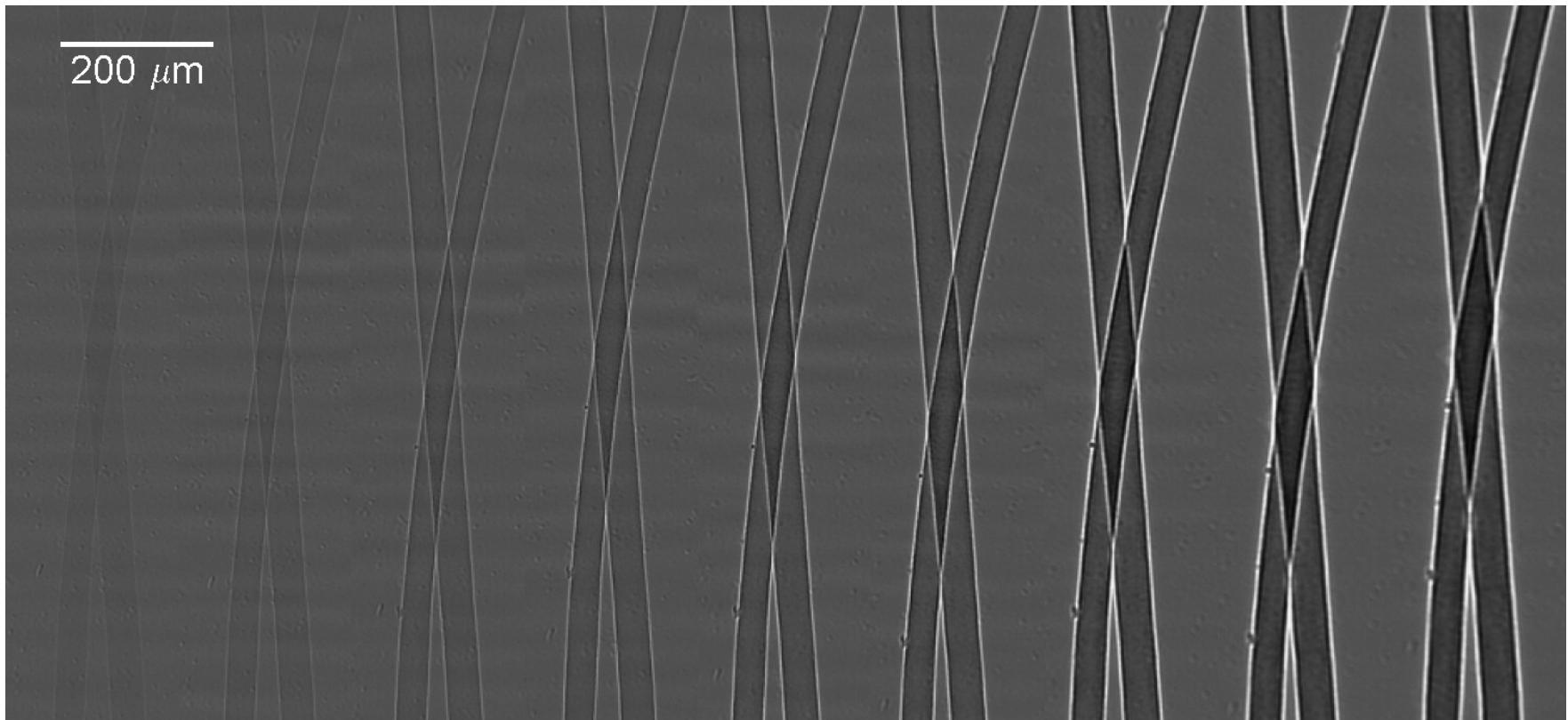
Some diffraction
may also occur



If refractive index
different, phase
shift occurs



...more subtle



20mm



1140mm

Distance

$E = 20\text{keV}$

From: A. Rack et al. NIM A, 586 (2), 327-344, (2008).

Low Z, soft tissue



*Reconstructed slice
of mouse knee*

*Results courtesy
K. Maadi*

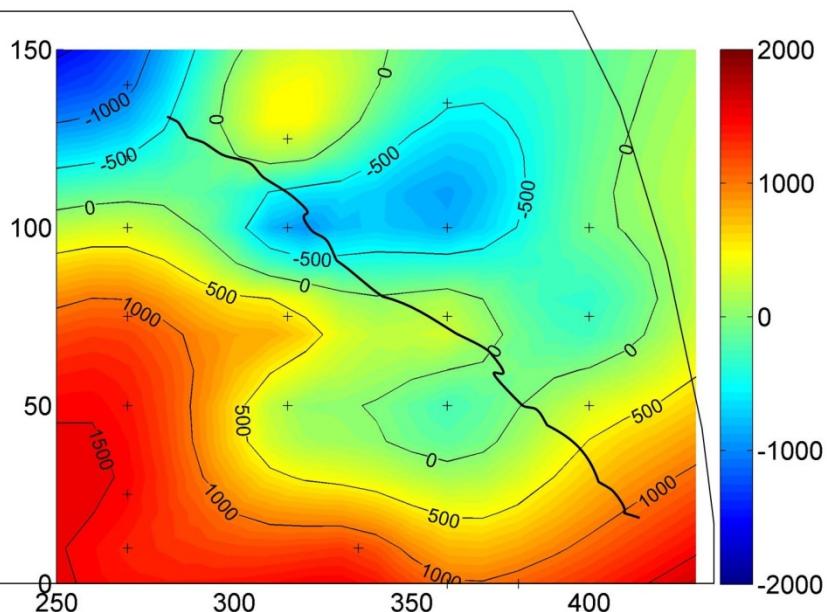
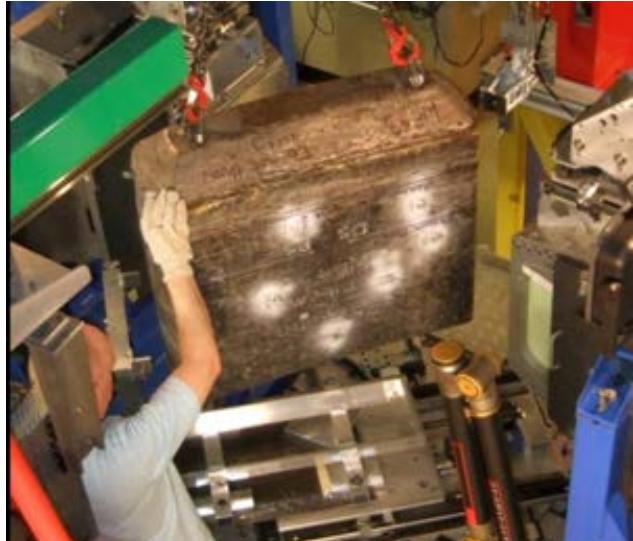
*Image processed with
support of A. Bodey*

Complementarity of neutron imaging

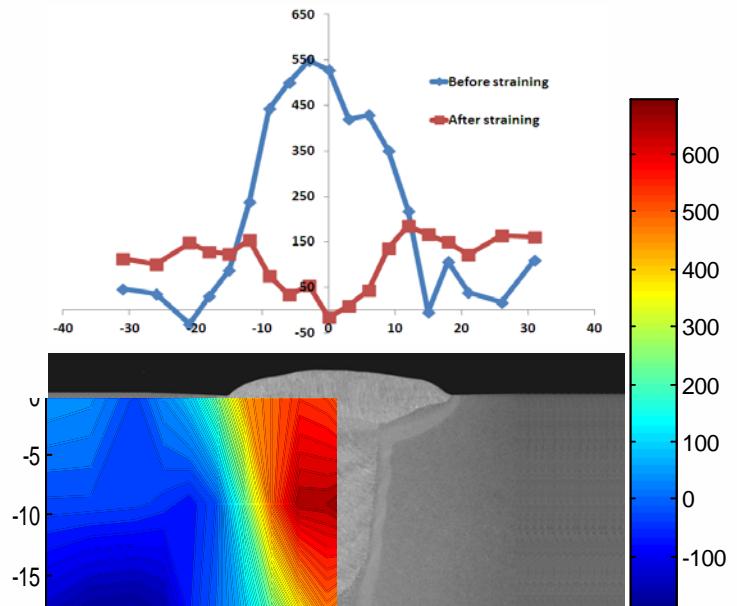
	Neutrons	Synchrotron X-rays
Penetration •Steels	40-60 mm	20-25 mm
Spatial resolution	0.5 mm ³ (cubic gauge volume)	0.1 mm lateral to incident beam; 1 mm parallel to beam (elongated gauge volume)
Stress complete tensor	Yes	With problems
Coarse grain problem	No	Yes
Counting time	10-30 minutes	Seconds-minutes



Complementarity of neutron imaging

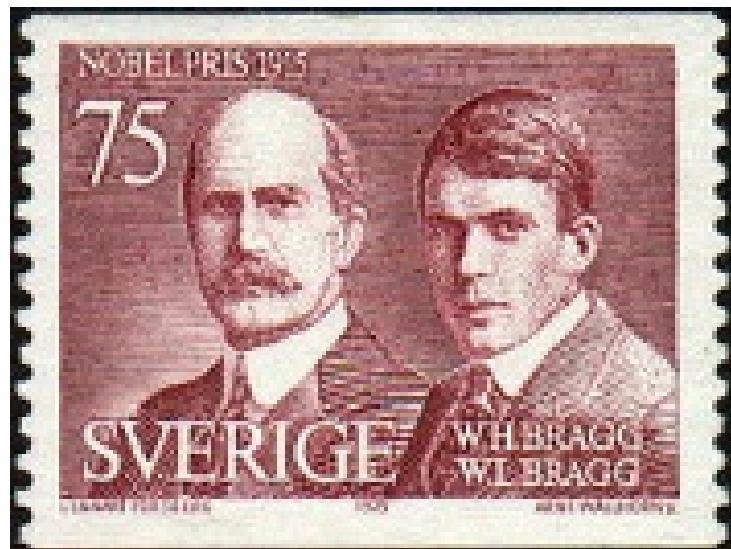


Residual strain profile within the Direct Chill cast Mg alloy slab, with estimated position of the crack in the measurement plane.
M Turski, et al., Metall Mater Trans A 43A (5) 1547-1557 (2012)

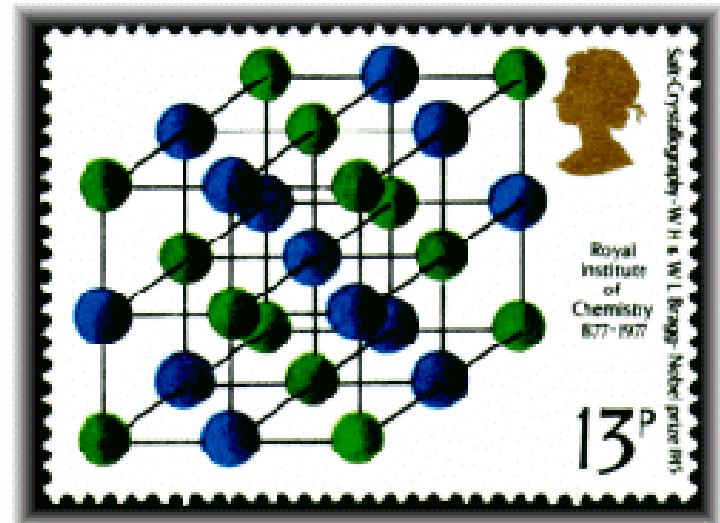
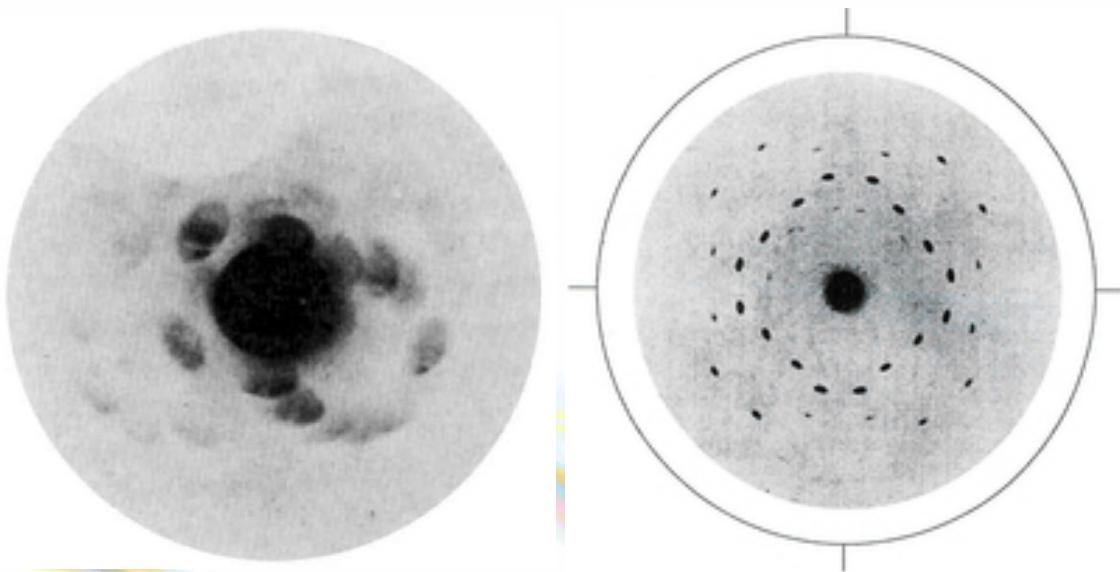


Measure welding residual stresses in pipeline for on-offshore transport of oil and gas
With thanks to Shu Yan Zhang, ISIS

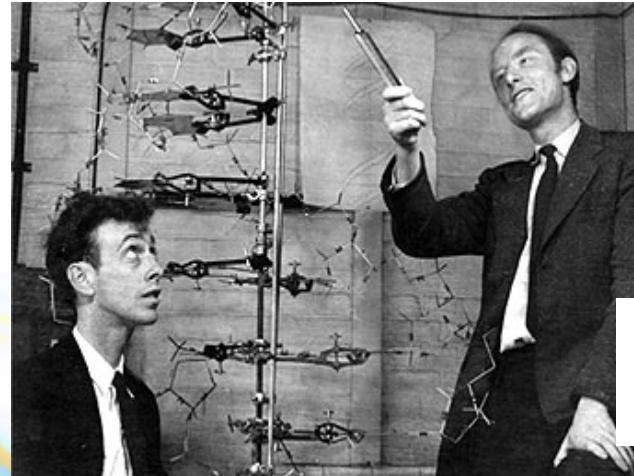
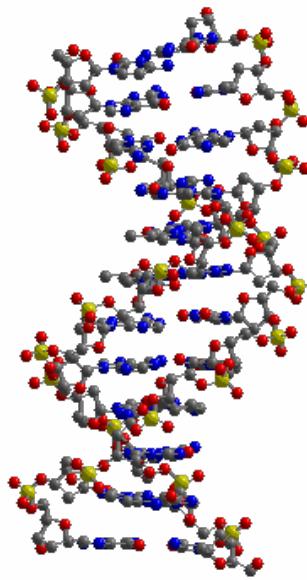
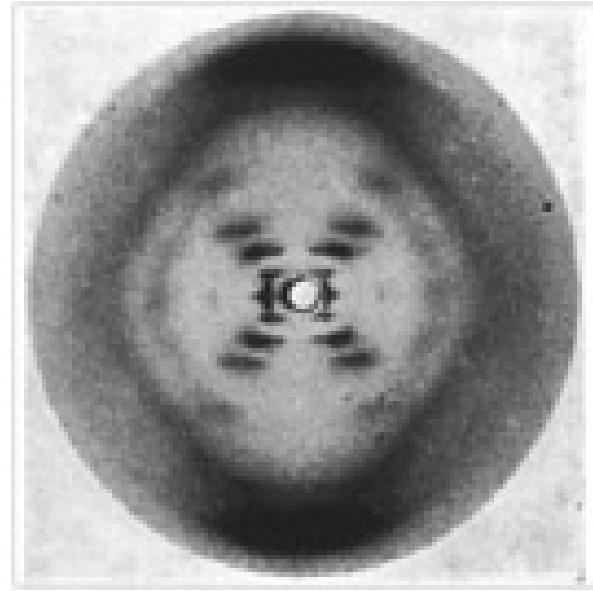
Origins of diffraction at atomic scales



$$n\lambda = 2d \cdot \sin\theta$$

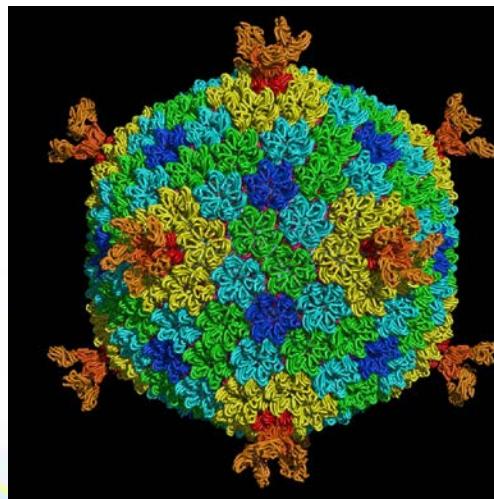
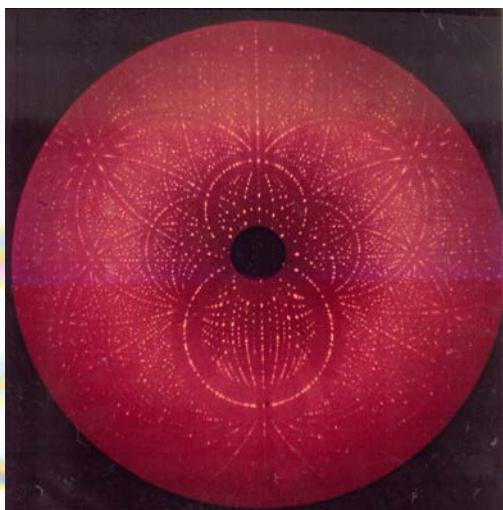
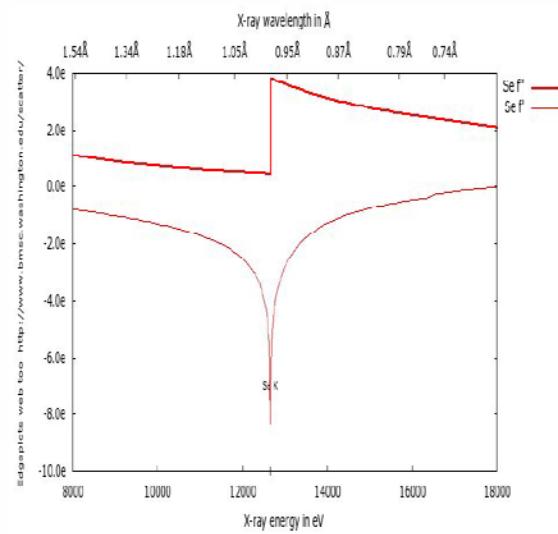
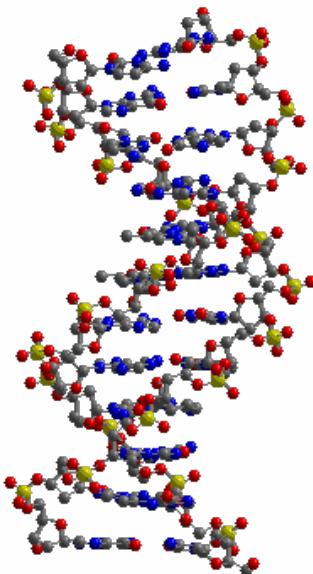
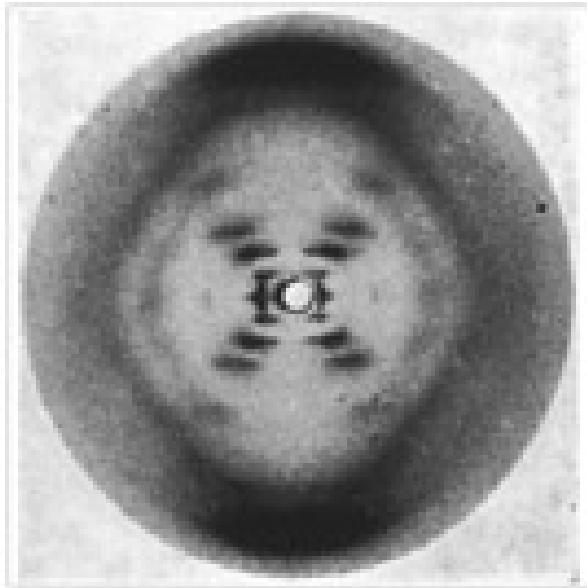


40 years on...



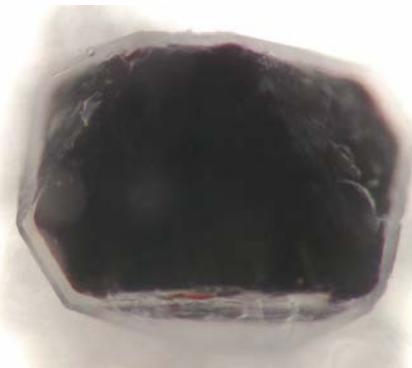
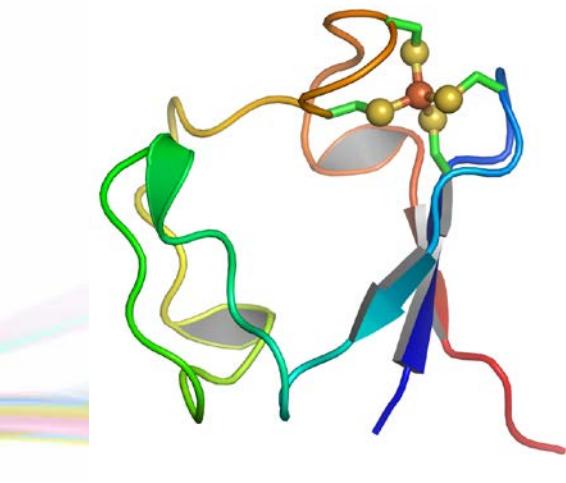
Macromolecular crystallography today

- Brilliance and tunability of photon energy, high-performance computing – $> 10^{10}$ times more information/second



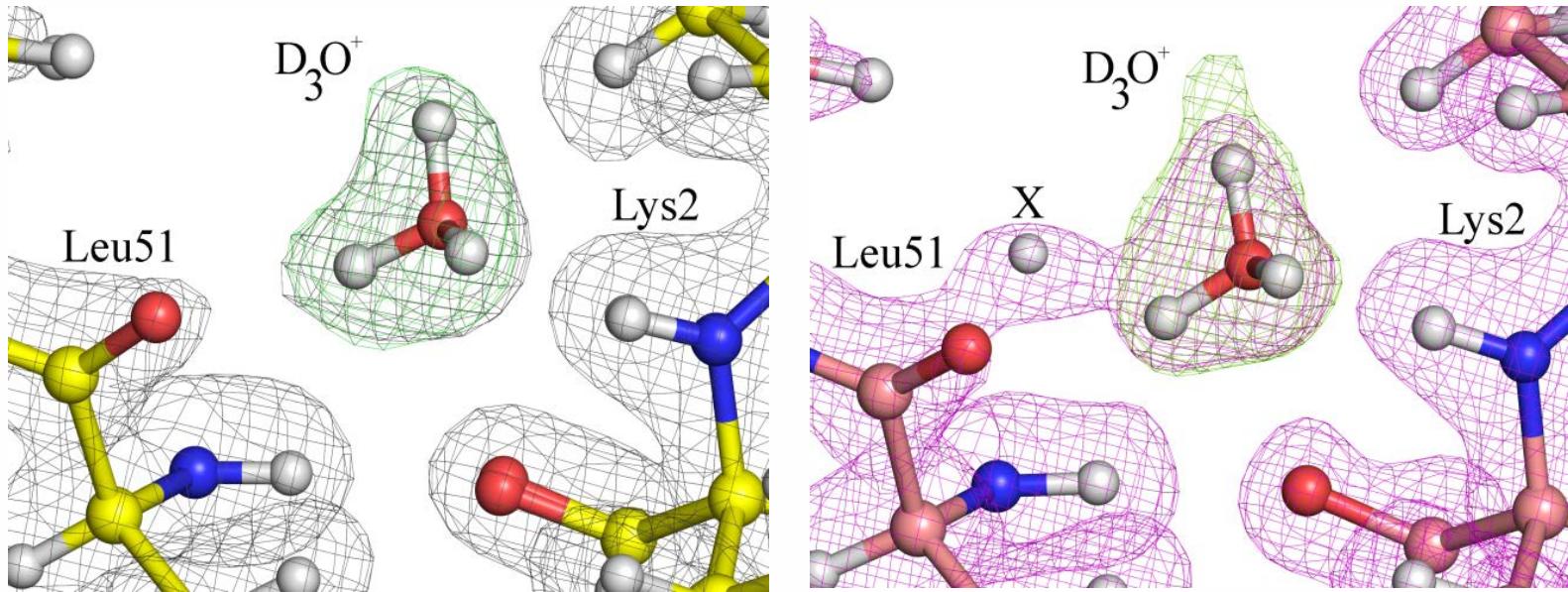
Complementarity: neutrons

- Location of H atoms when X-rays/chemical ‘rules’ don’t help
- Study of samples susceptible to radiation damage
- Study of Rubredoxin structure illustrates both
 - Small (~6kD) iron-sulphur containing redox protein – important model system to understand electron transfer processes using redox systems – here Fe^{3+} - Fe^{2+}
 - Fe^{3+} form very easily reduced in the X-ray beam

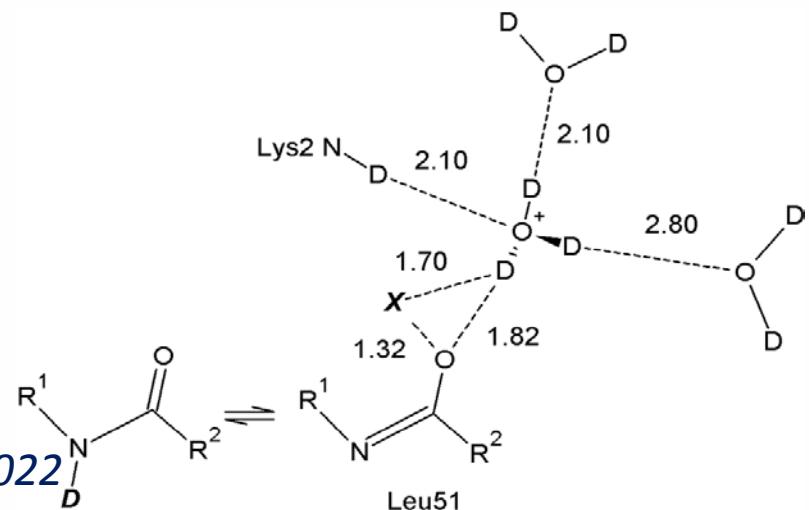


Complementarity: neutrons

- Structure of reduced and oxidised form measured on D19 at ILL

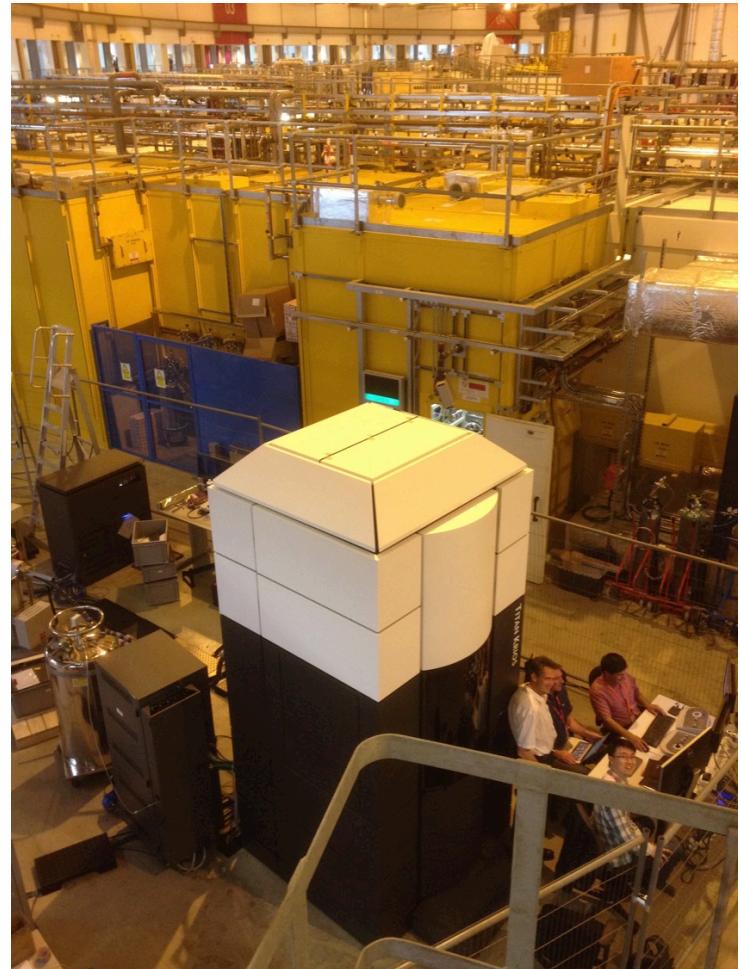
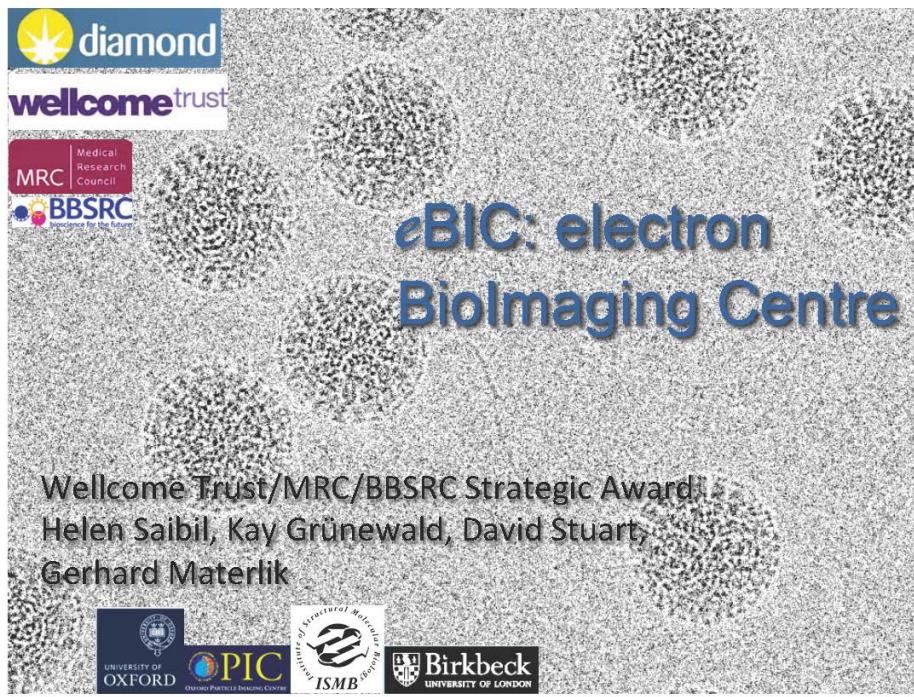


Observation of hydronium (D) ions and of tautomeric shifts following the change from the oxidised form to the reduced form



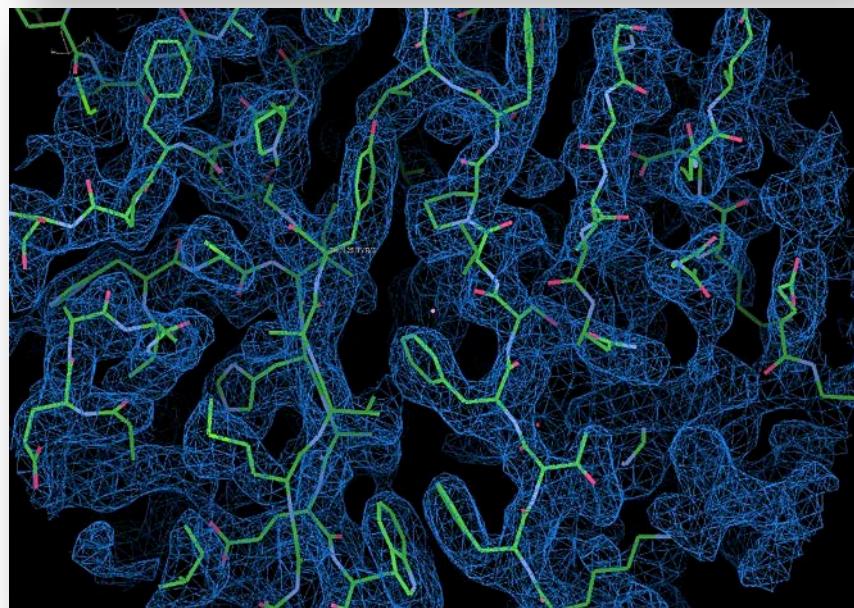
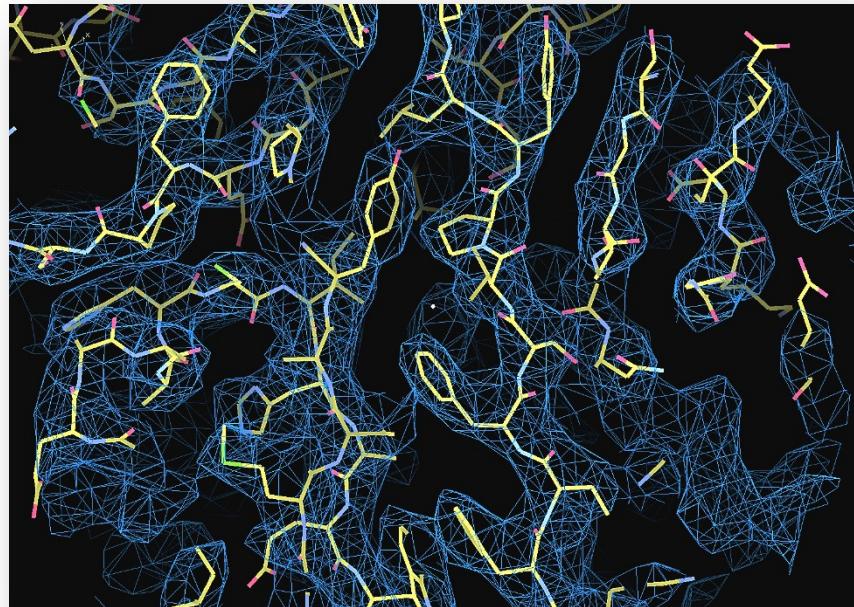
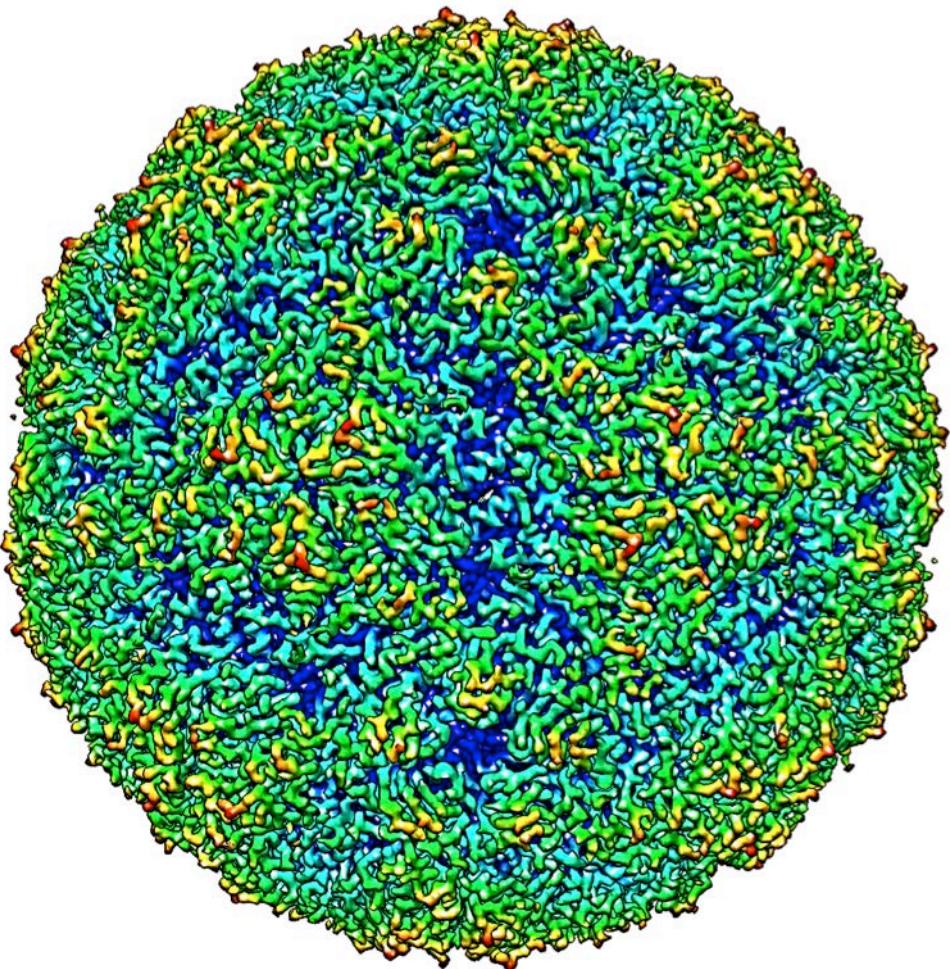
Complementarity: the EM revolution

- Wonderful new detectors that collect movies and allow particle movements to be tracked in software to reveal atomic detail
- Native particles frozen in a water ‘glass’



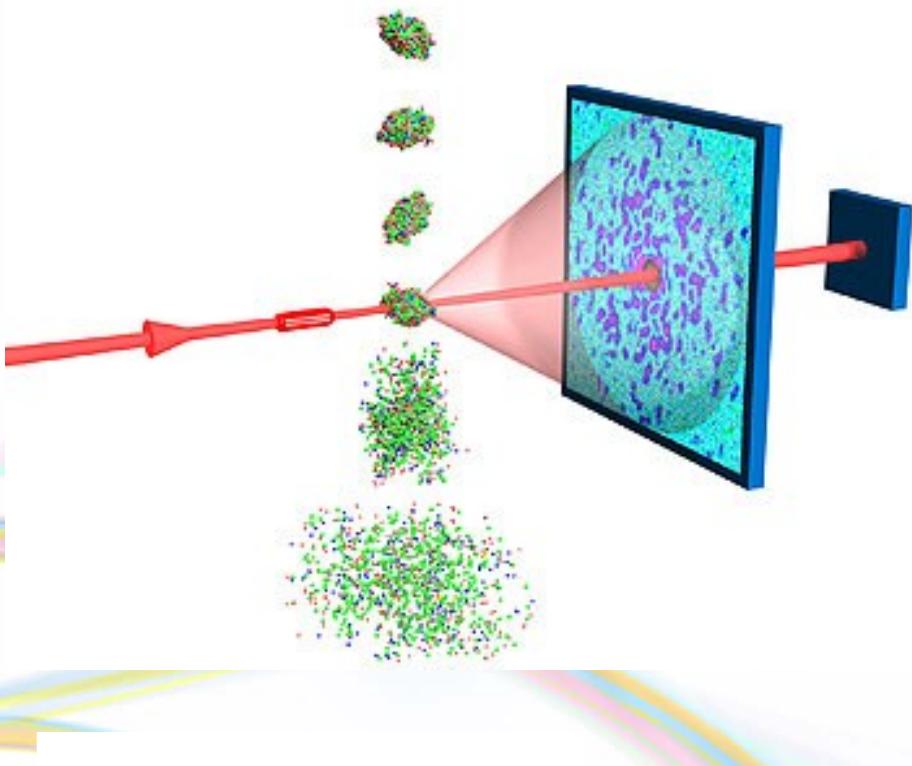
Complementarity: the EM revolution

- Example from foot-and-mouth disease virus by electron microscopy and synchrotron X-rays



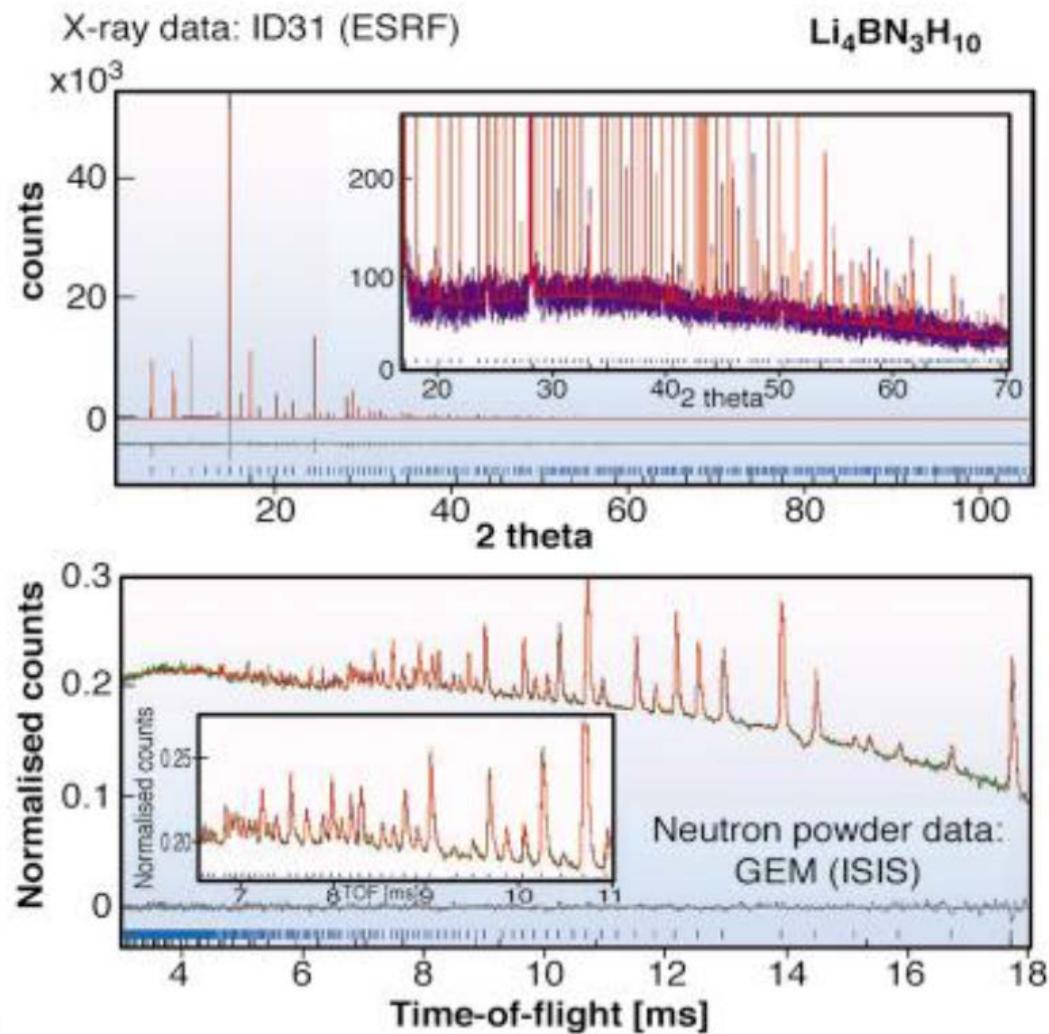
Complementarity: FELs

- X-rays pulses in a free electron laser (FEL) are ultrafast (fs) and ultrabright ($\times 10^{10}$) and capture structure before the sample is perturbed by the beam – ‘diffract and destroy’ !
- Challenge of sample injection, image processing...



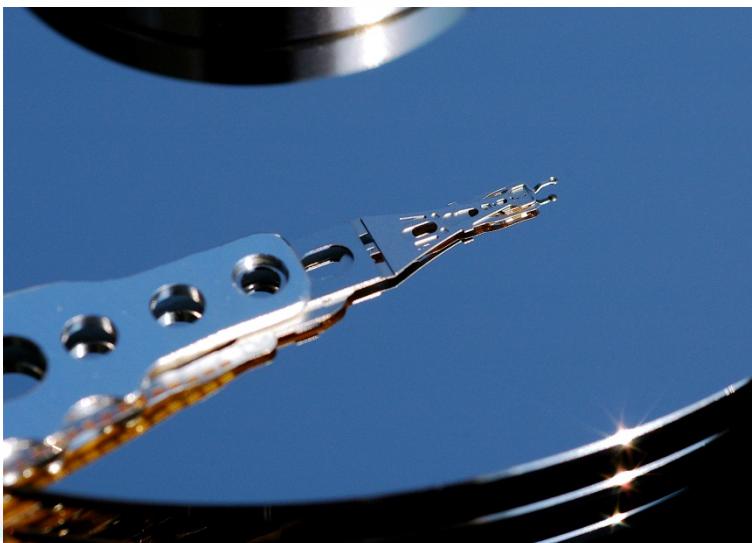
'Hard' materials

- Many important materials contain hydrogen or other light atoms that are crucial to function – e.g. for hydrogen storage
- Combine X-rays for rapid, high-resolution survey/ study then neutrons to locate H, Li, etc....



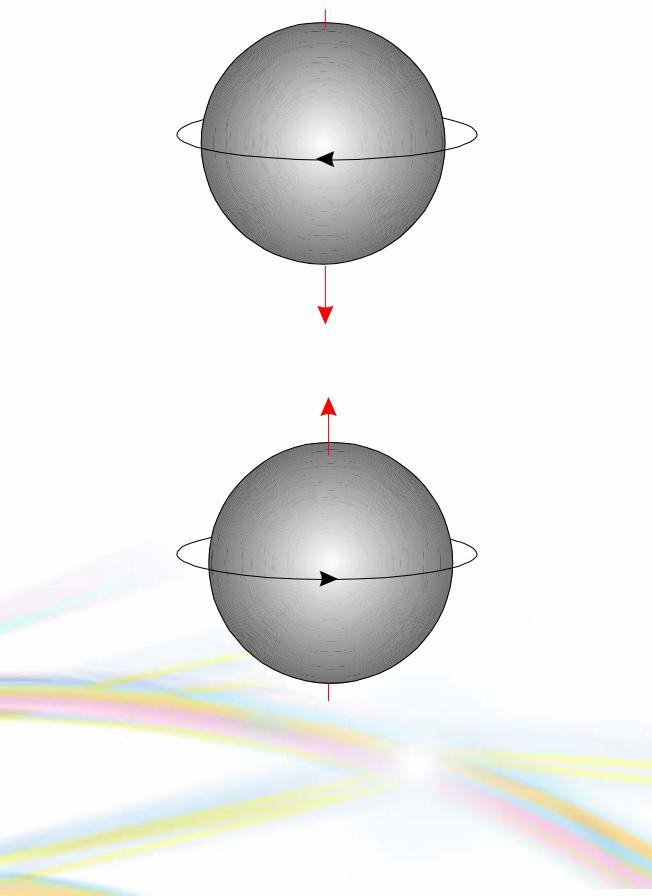
Magnetic structures

- Need to measure magnetic (spin) structures and excitations at an atomic level to understand and predict magnetic properties
- Important for a wide range of functional materials for sensors, actuators, information storage, nanomedicine, superconductors...



Magnetic structures: neutrons

- Scattering cross-section relatively simple and comparable in size to the nuclear cross-section



PHYSICAL REVIEW VOLUME 76, NUMBER 8 OCTOBER 19, 1949

Letters to the Editor

Detection of Antiferromagnetism by Neutron Diffraction*

C. G. SHULL

Oak Ridge National Laboratory, Oak Ridge, Tennessee

AND

J. SAMUEL SMART

Naval Ordnance Laboratory, White Oak, Silver Spring, Maryland

August 29, 1949

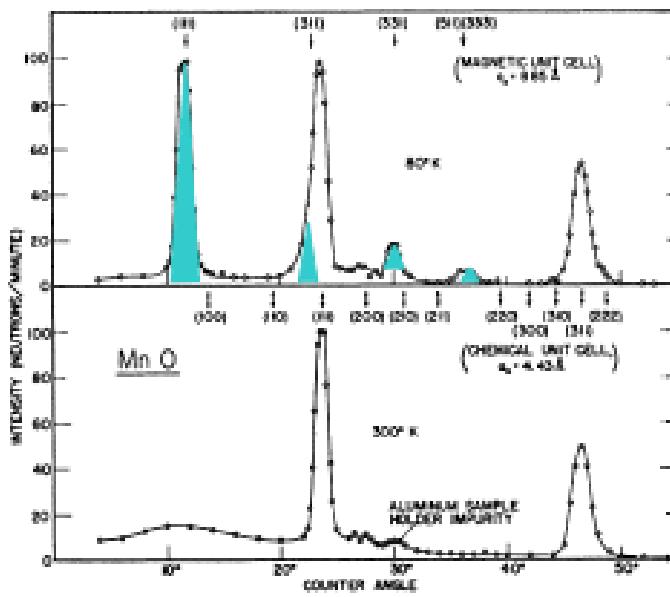
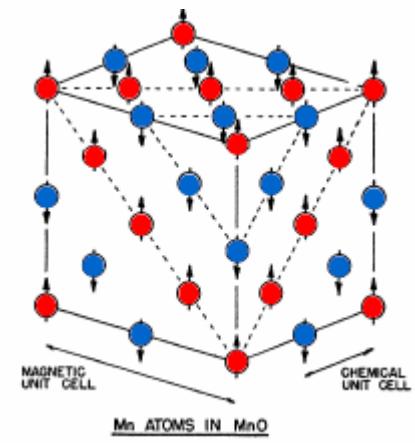
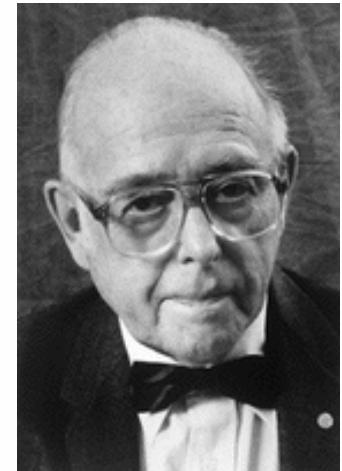
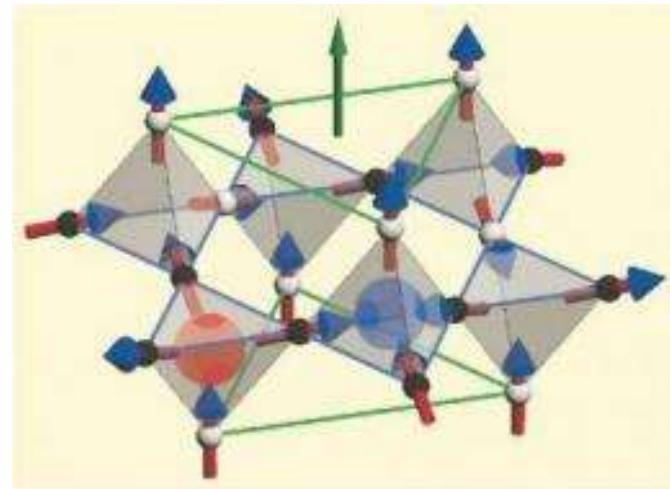
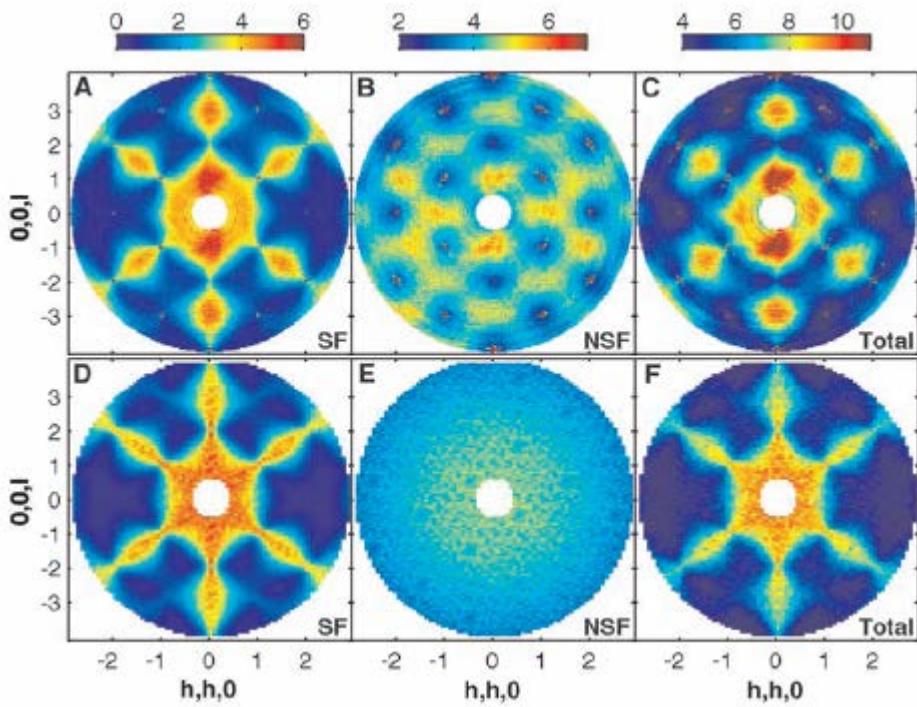


FIG. 1. Neutron diffraction patterns for MnO at room temperature and at 80°K .



Magnetic structures: neutrons

- Polarisation techniques can provide all details of spin structure (zero-field polarimetry) and map out diffuse scattering from short-range order
- Penetrating ability of neutrons enables design of sophisticated sample environment and cryogenics to nK, but large samples are required



Magnetic structures: X-rays

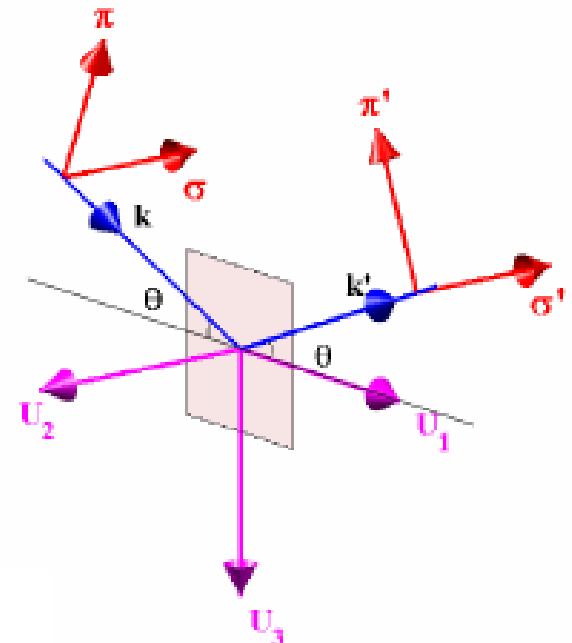
- Scattering cross-section generally very weak – typically 10^6 weaker than the nuclear (Thomson) terms ($(\hbar\omega/mc^2) \sim 10^{-2} \rightarrow |f^{\text{mag}}|/|f^{\text{nuc}}| \sim 10^{-6}$)
- Polarisation analysis can be used to separate spin (S) and orbital (L) contributions
- Magnetic scattering amplitude $f^{\text{mag}}(Q)$:

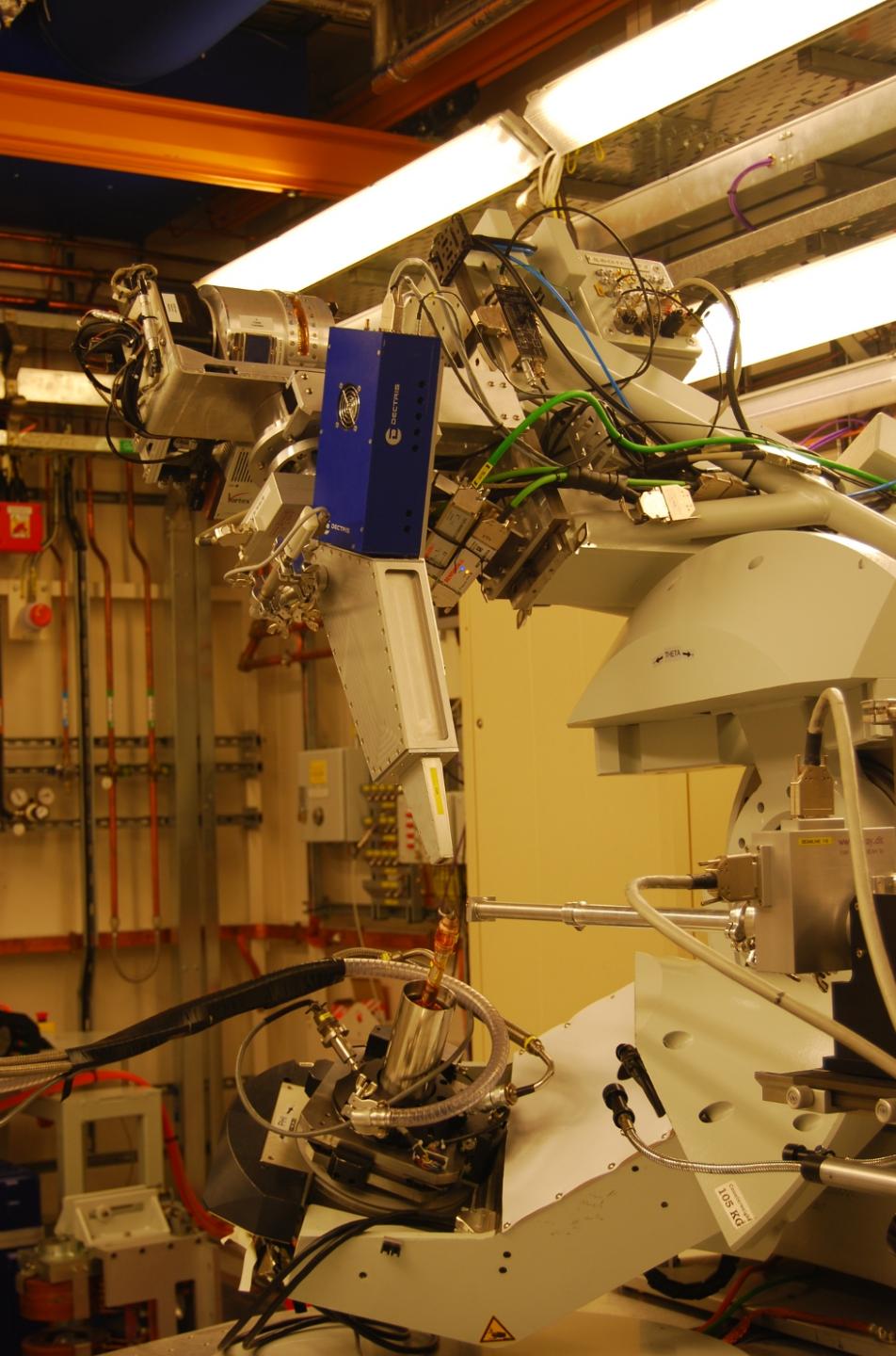
$$f^{\text{mag}}(Q) = ir_0 \left(\frac{\hbar\omega}{mc^2} \right) \left[\frac{1}{2} A'' \cdot L(Q) + B \cdot S(Q) \right]$$

$L(Q)$ and $S(Q)$ are the Fourier transforms of the atomic and spin magnetisation densities; A'' and B depend on the relative orientation of k , k' , ε , ε'

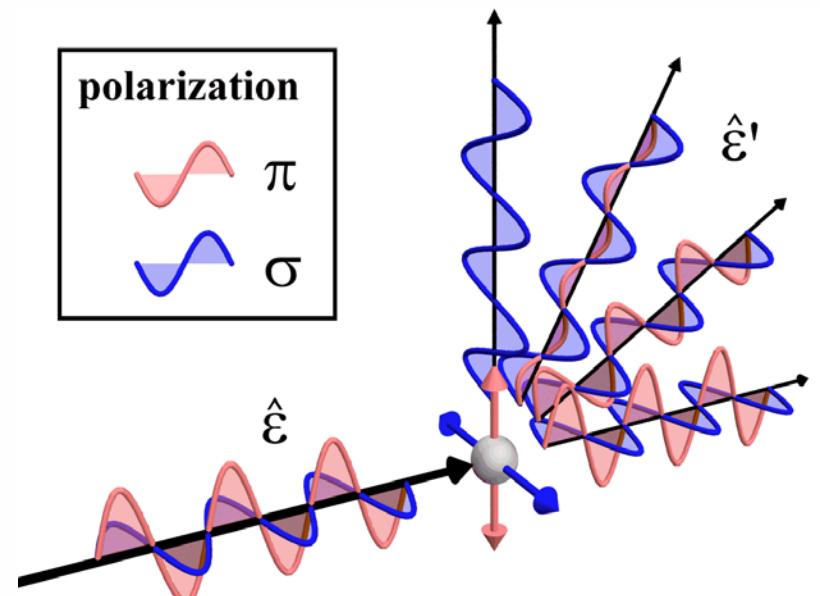
$$f^{\text{mag}}(Q) = ir_0 \left(\frac{\hbar\omega}{mc^2} \right) \times$$

	$\hat{\varepsilon}_\perp \equiv \sigma$	$\hat{\varepsilon}_\parallel \equiv \pi$
$\hat{\varepsilon}'_\perp$	$\sin 2\theta S_2$	$-2\sin^2 \theta [(L_1 + S_1)\cos\theta - S_3 \sin\theta]$
$\hat{\varepsilon}'_\parallel$	$2\sin^2 \theta [(L_1 + S_1)\cos\theta - S_3 \sin\theta]$	$\sin 2\theta [2\sin^2 \theta L_2 + S_2]$





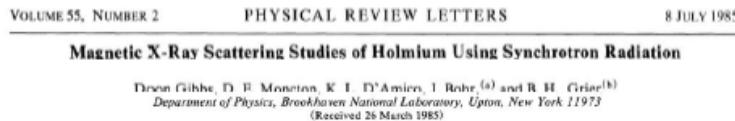
- Versatile diffractometer: polarisation analysis and azimuthal angle
- Energy tunable from 1-15 keV to get good atomic resolution



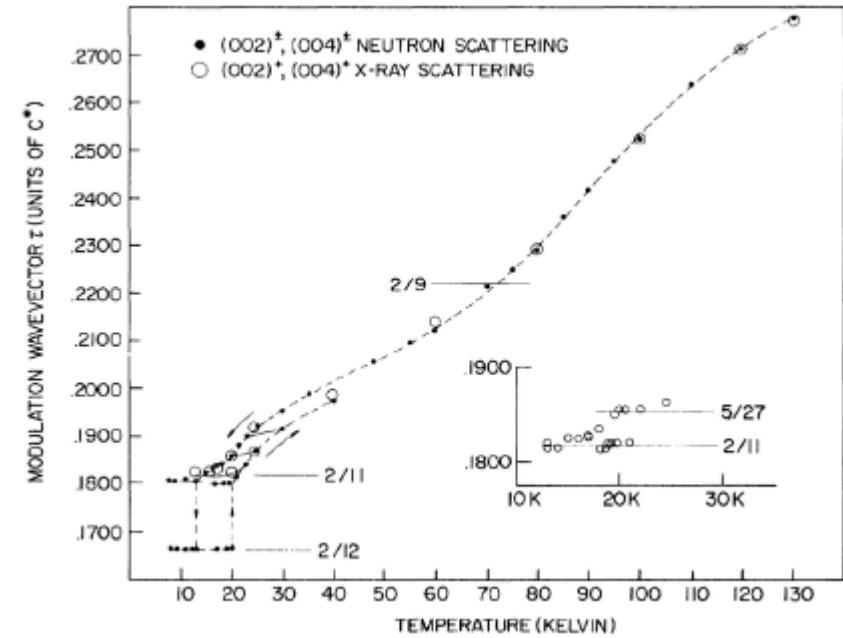
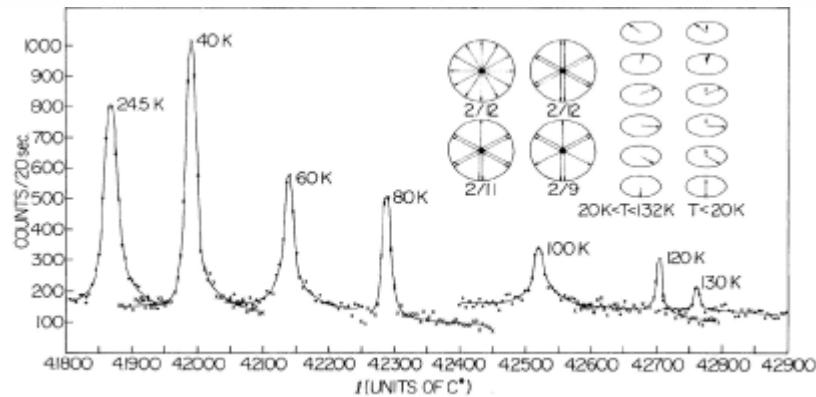
 diamond

Magnetic structures: X-rays

- First magnetic structures studied at a synchrotron in 1985 - spiral spin structure of Ho mapped out as a function of temperature



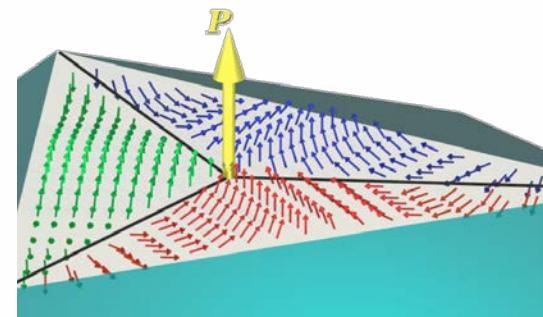
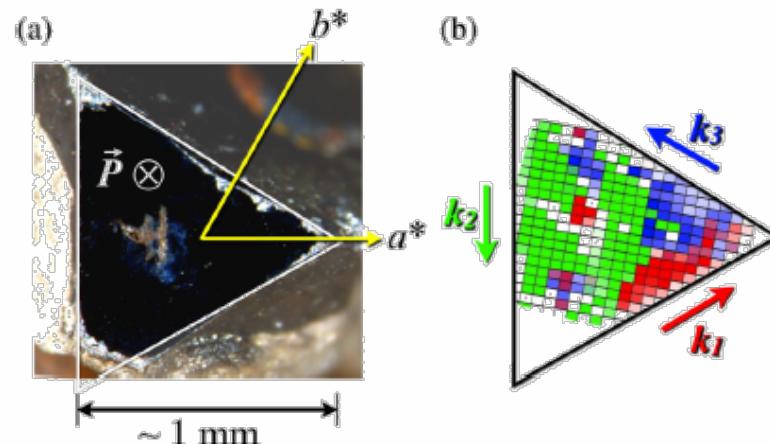
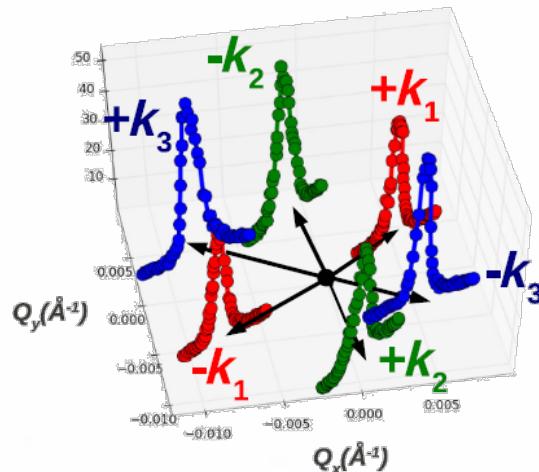
(1985) First Synchrotron Radiation Studies of Magnetism



- Advantages of (non-resonant) X-ray magnetic scattering
 - Orbital and spin magnetisation density may be separated
 - Highly-focused beams enable very small samples to be studied
 - High resolution in Q – e.g. for subtle changes in spin structure
 - **However**, cryogenics are challenging !

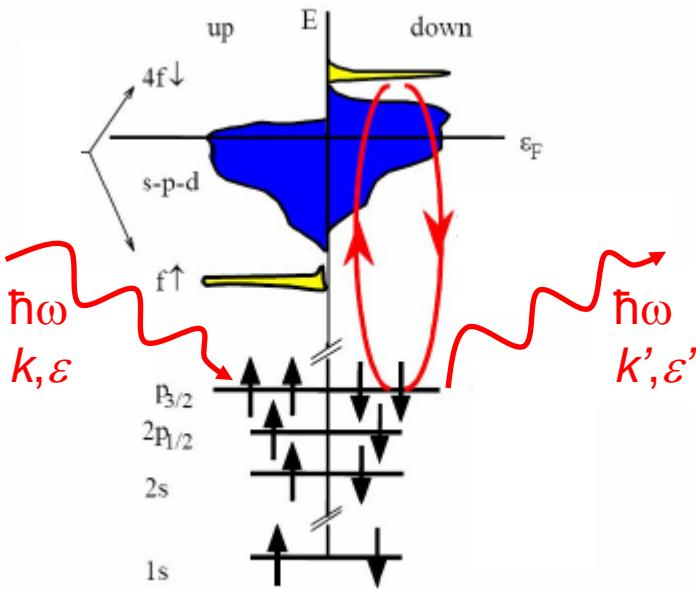
Magnetic structures: X-rays

- BiFeO₃ has room-temperature coupling between magnetic and ferroelectric (FE) order.
- Map domain distribution with 50 μm spot (LD) then probe helicity by CD: same helicity for single FE domain



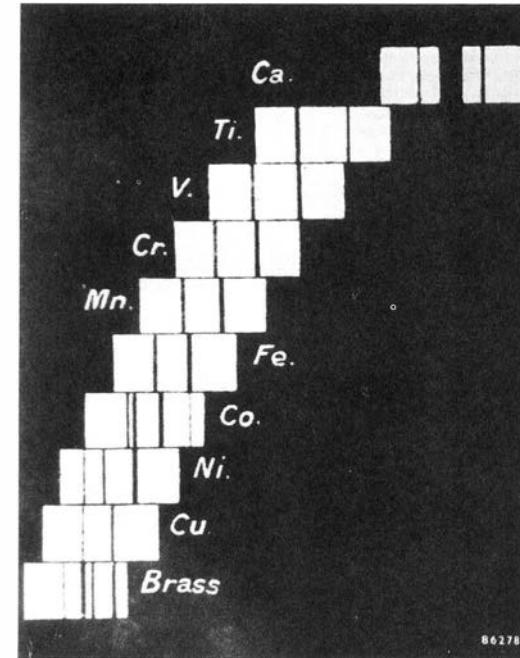
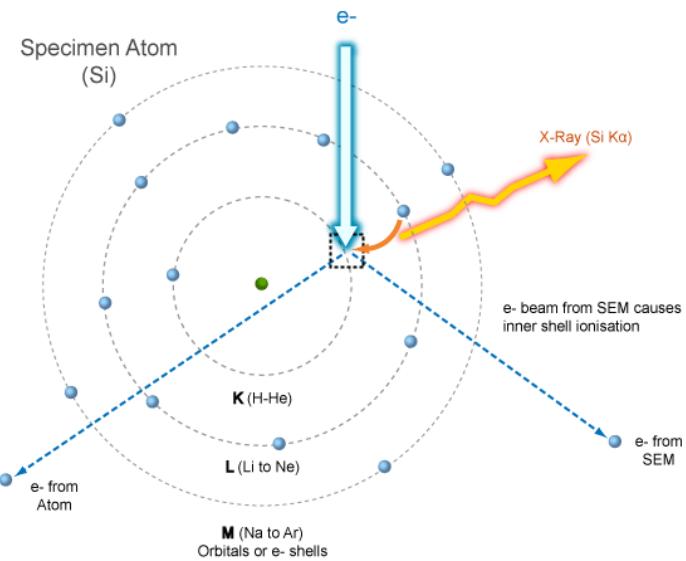
Magnetic structures: resonant scattering

- Tune energy to a core-(empty) valence state transition and exploit resonant scattering to greatly enhance intensity ($\times 10^3$)
 - Sensitive to specific element and oxidation state
 - Dipole transitions ($\Delta L = 1$) generally stronger than quadrupole transitions ($\Delta L = 2$) and these will also be polarisation dependent; transitions from s to d or p states generally weaker because of lower overlap of wavefunctions
 - Sensitive to spin (S) though spin-orbit coupling
 - However, mostly limited to softer X-rays so will not give atomic resolution – tend to use for larger-scale structures (unit cell, periodic structures)



elements	edge	transition	energy range [keV]	resonance strength	comment
3d	K	$1s \rightarrow 4p$	5 - 9	weak	small overlap
3d	L _I	$2s \rightarrow 3d$	0.5 - 1.2	weak	small overlap
3d	L _{II} , L _{III}	$2p \rightarrow 3d$	0.4 - 1.0	strong	dipolar, large overlap, high spin polarisation of 3d
4f	K	$1s \rightarrow 5p$	40 - 63	weak	small overlap
4f	L _I	$2s \rightarrow 5d$	6.5 - 11	weak	small overlap
4f	L _{II} , L _{III}	$2p \rightarrow 5d$ $2p \rightarrow 4f$	6 - 10	medium	dipolar quadrupolar
4f	M _I	$3s \rightarrow 5p$	1.4 - 2.5	weak	small overlap
4f	M _{II} , M _{III}	$3p \rightarrow 5d$ $3p \rightarrow 4f$	1.3 - 2.2	medium to strong	dipolar quadrupolar
4f	M _{IV} , M _V	$3d \rightarrow 4f$	0.9 - 1.6	strong	dipolar, large overlap, high spin polarisation of 4f
5f	M _{IV} , M _V	$3d \rightarrow 5f$	3.3 - 3.9	strong	dipolar, large overlap, high spin polarisation of 5f

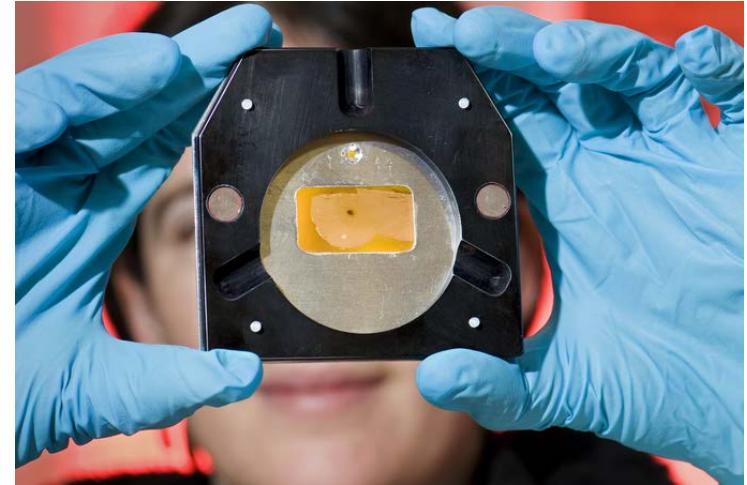
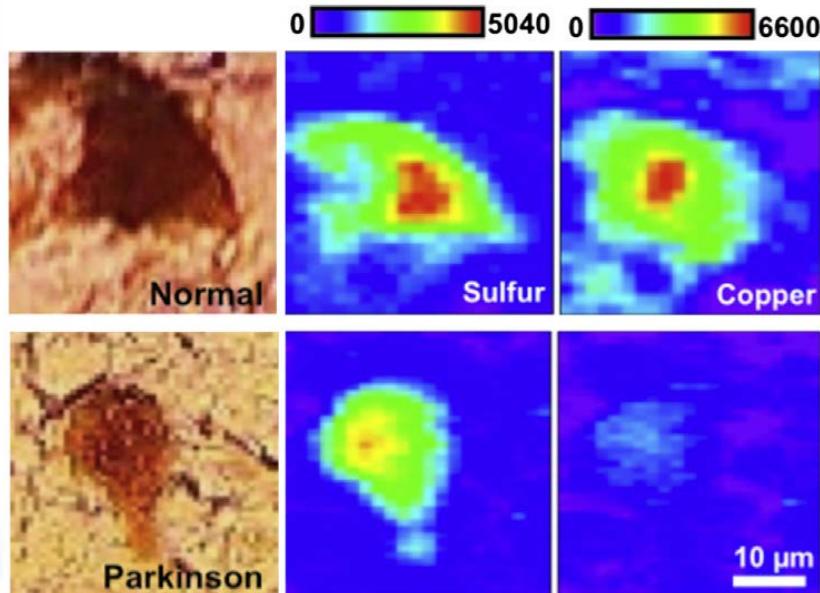
Origins of X-ray spectroscopy



Ce Pr Nd Sm Eu Gd Tb Dy Ho Er Tm Yb Lu

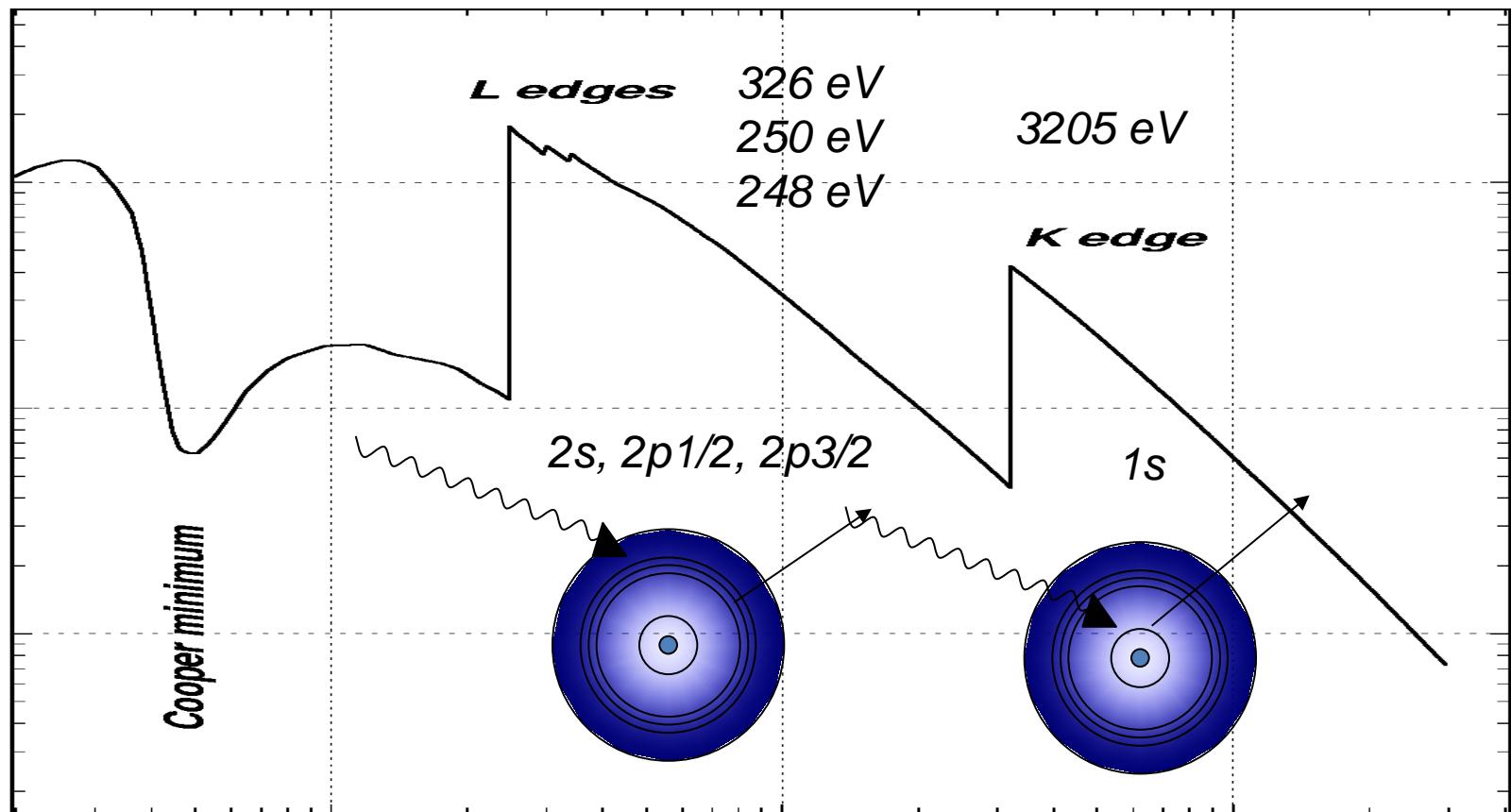


Mapping elements – but what about chemistry ?

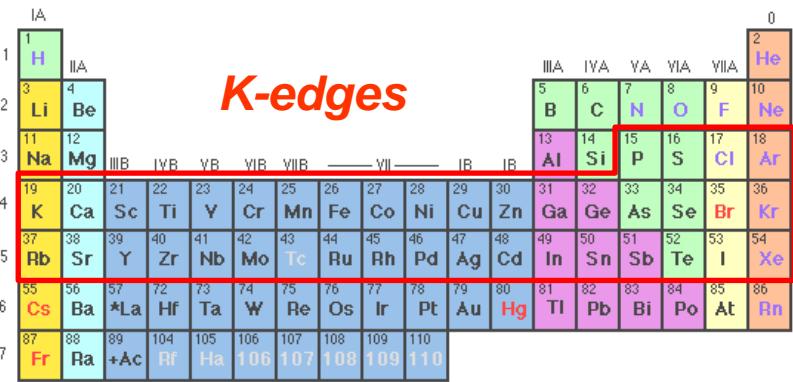


Exploring electronic structure

- Atomic levels - e.g. Argon – $1s^2 2s^2 2p^6 3s^2 3p^6$

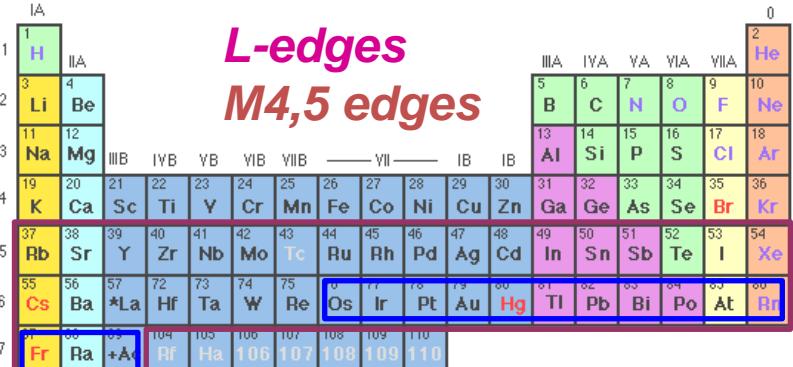


Sensitivity to chemical environment



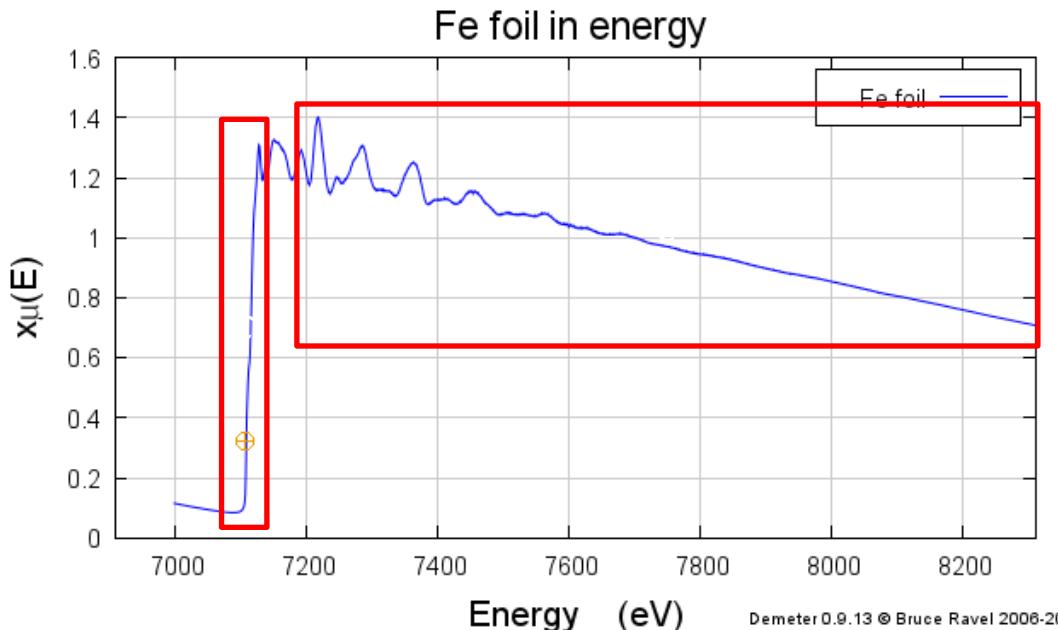
K-edges

* Lanthanide Series	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
+ Actinide Series	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr



L-edges
M4,5 edges

* Lanthanide Series	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
+ Actinide Series	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr

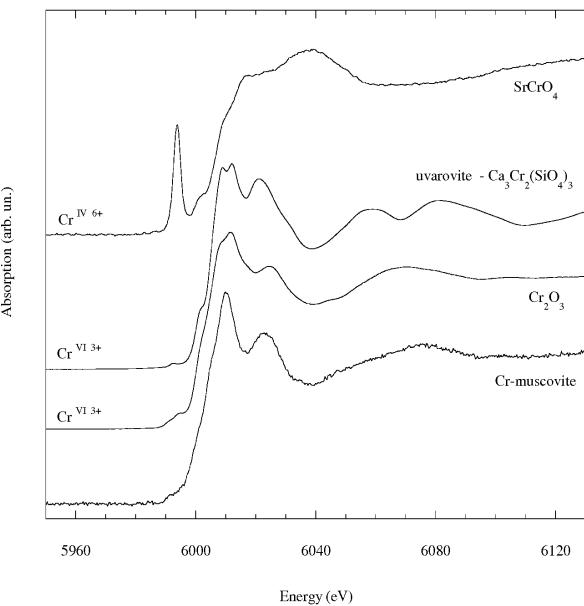
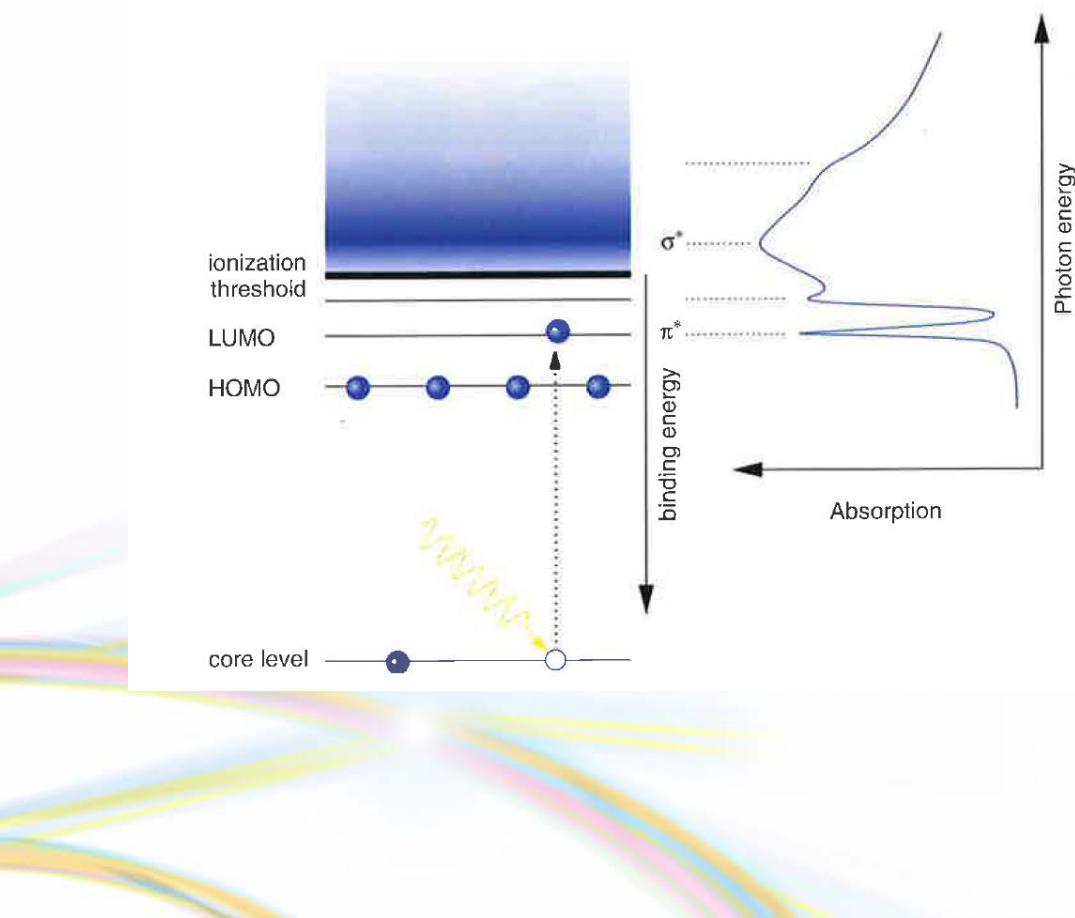


Demeter 0.9.13 © Bruce Ravel 2006-201



XANES

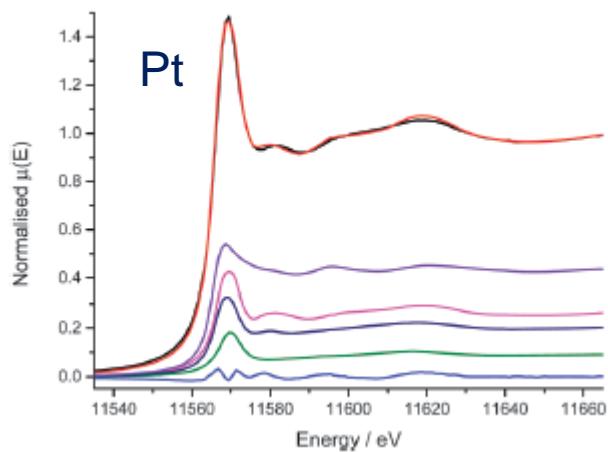
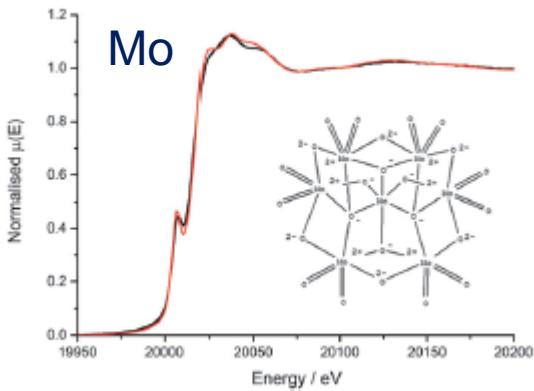
- X-ray Absorption Near Edge Spectroscopy: element-specific oxidation state and bonding in molecules and solids



Cr K absorption edge

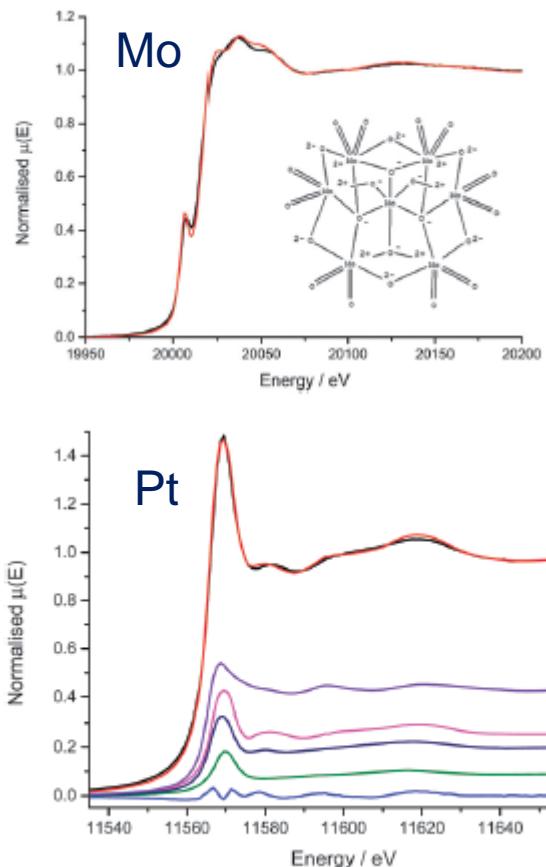


Unravelling chemical clusters

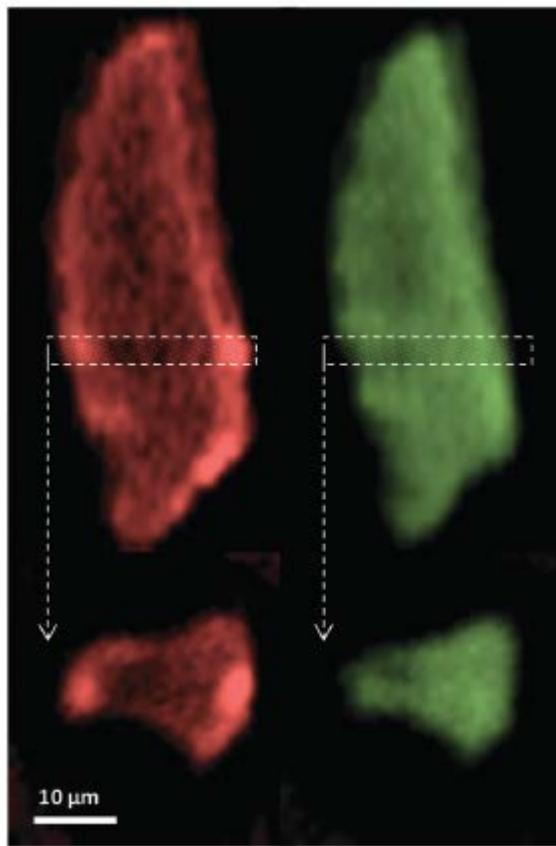


Species identification

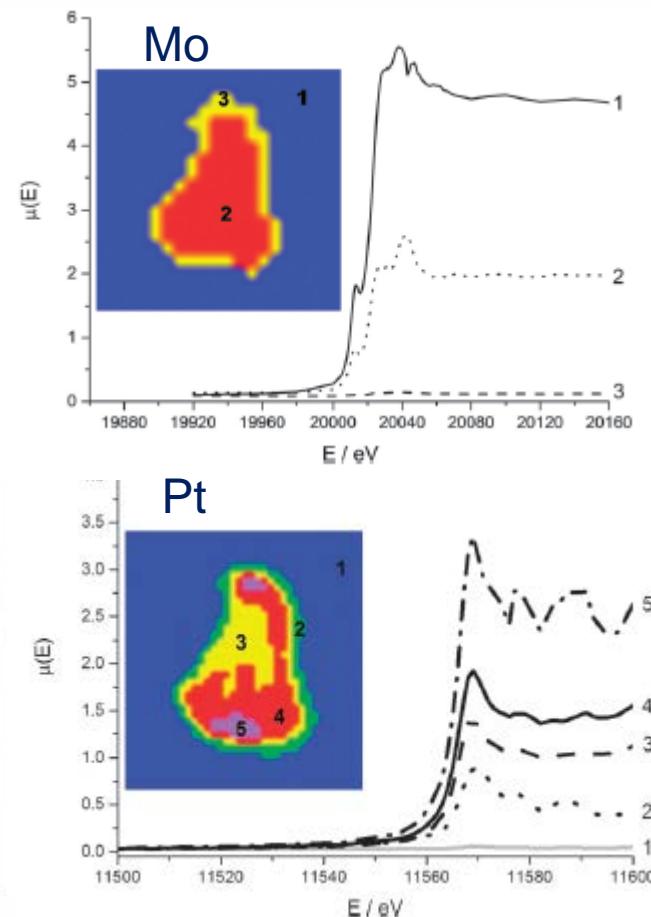
...and mapping where they are



Species identification



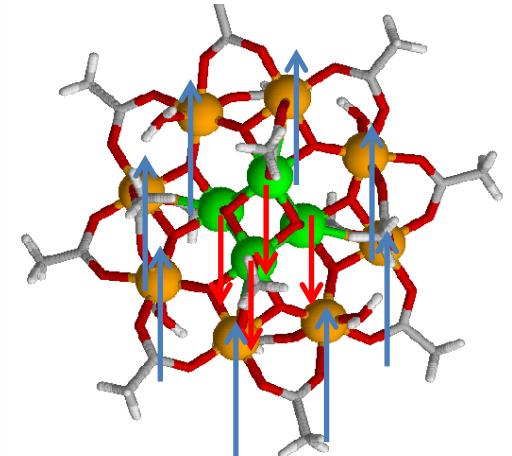
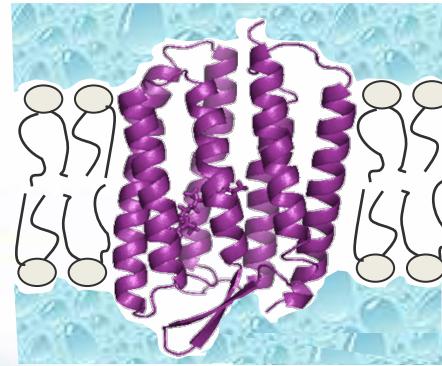
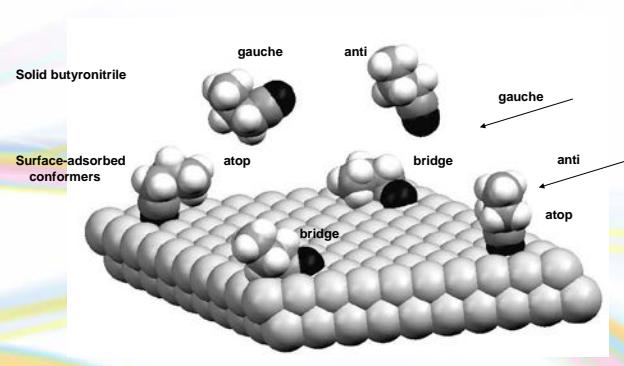
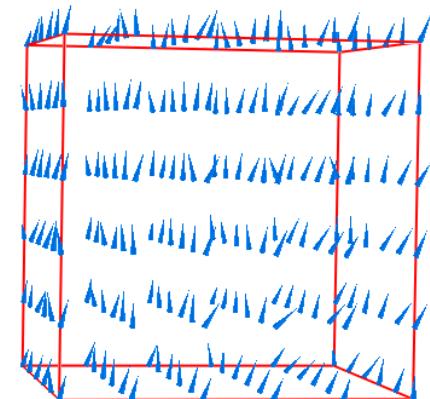
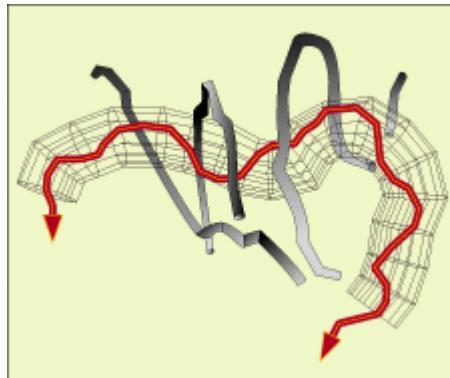
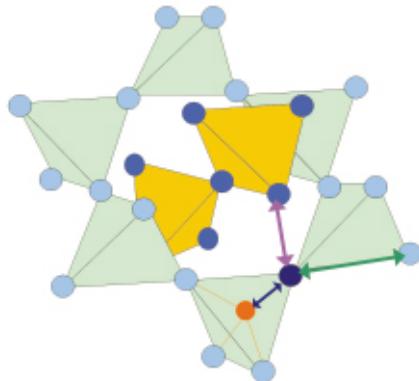
XRF tomography –
Mo (green) and Pt (red)



Species distribution

Higher resolution spectroscopy: atomic and molecular motion in materials

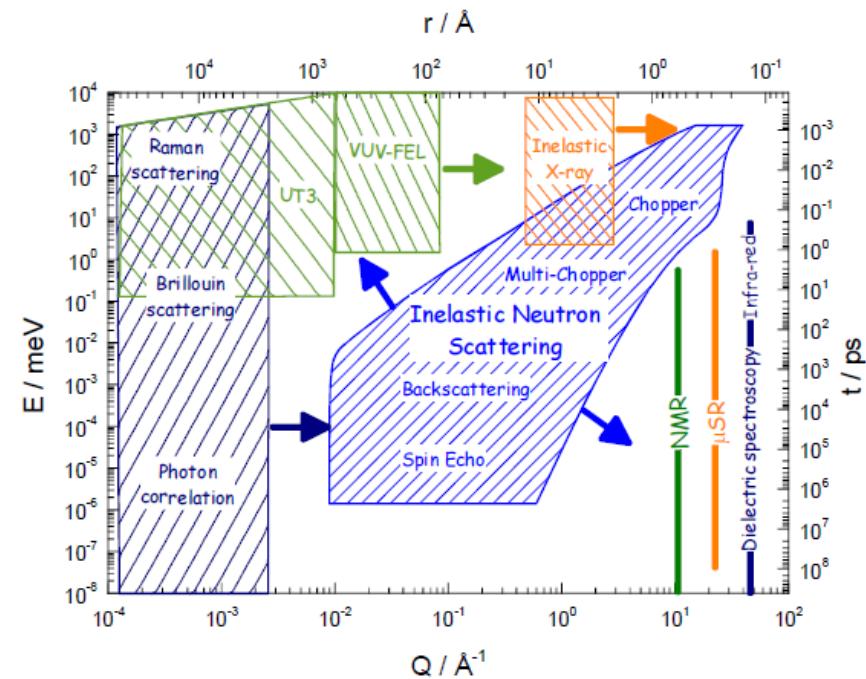
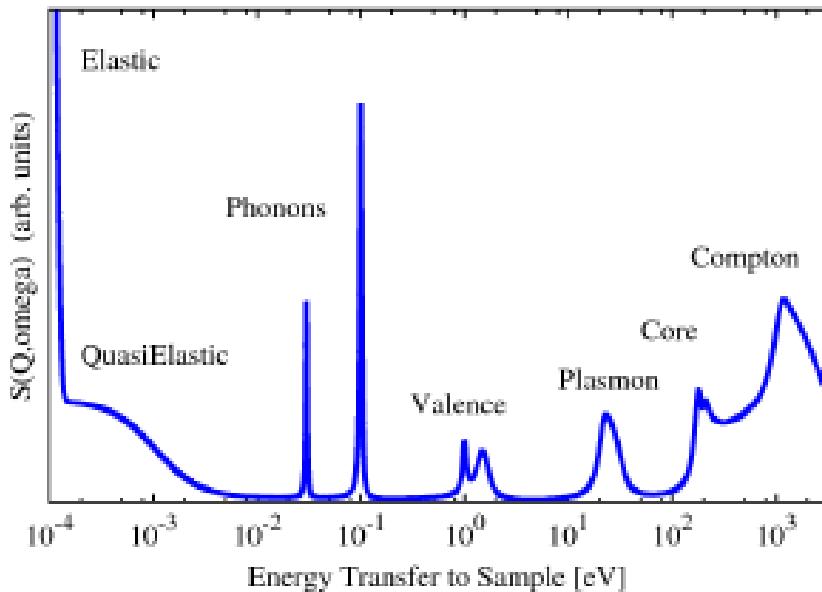
- Vibrations and magnetic excitations in molecules and solids are essential to key processes e.g. thermal conductivity and expansion, polymer processing and properties, biochemical processes, magnetic storage media, superconductors –



Atomic and molecular motion in materials

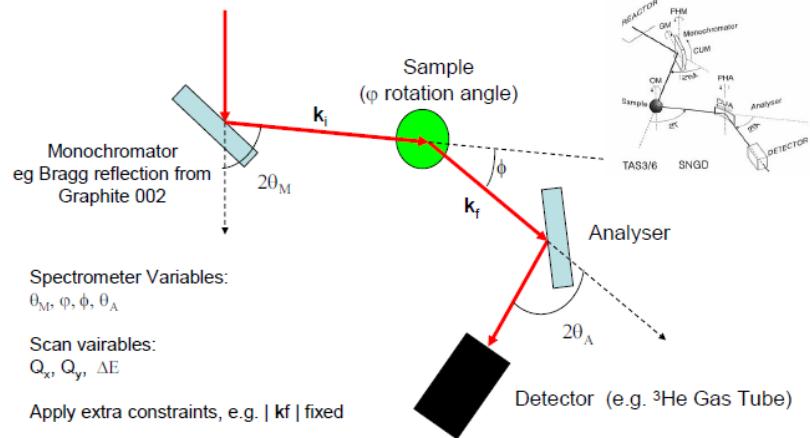
"When a crystal is irradiated with X-rays, the processes of photoelectric absorption and fluorescence are no doubt accompanied by absorption and emission of phonons. The energy changes involved are however so large compared with phonon energies that information about phonon spectrum of the crystal cannot be obtained in this way. The same is true for Compton scattering."¹ W.Cochran, 1966.

- Typically wish to measure vibrations and magnetic excitations in the range 1-100 meV as a function of q ($1/\lambda$) so the wavelength of the neutrons or X-rays should typically be of the order of 1-2 Å (0.1-0.2 nm) [Q – measured momentum transfer; q – wavevector of the excitation]
- 1 Å corresponds to 12.4 keV for X-rays and 81.8 meV for neutrons so instrumental resolution is ***much*** more challenging for X-rays (10meV: $\Delta E/E \approx 10^{-6}$ vs 10^{-6} IXS vs INS)

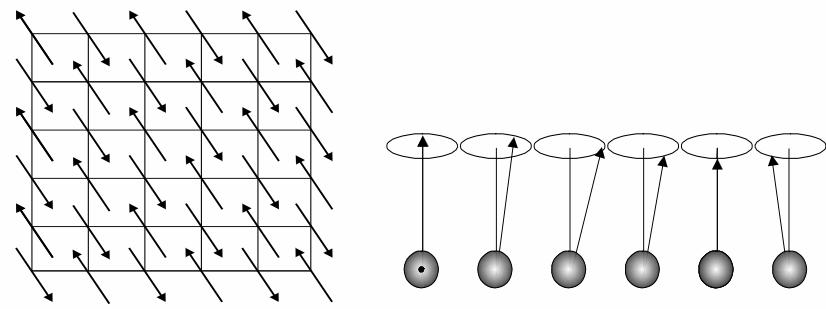


INS: experimental methods

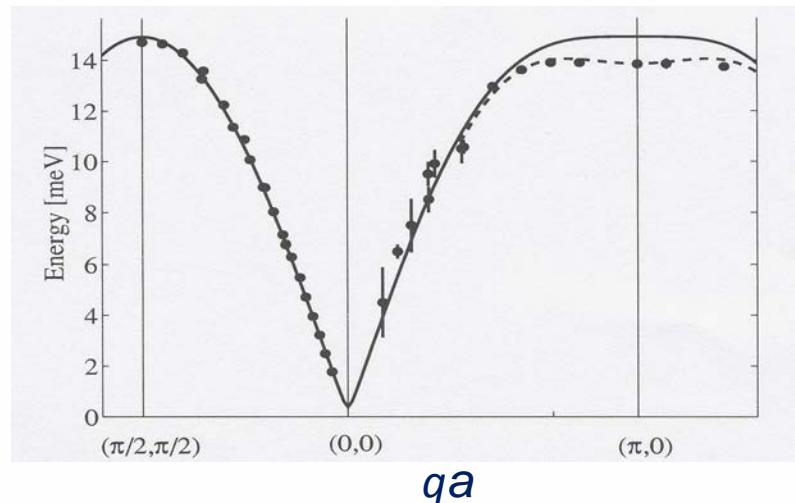
- Triple-axis spectrometer – very versatile means of measuring the change in energy of neutrons scattered by sample as a function of Q (point by point !)



$$Q = k_f - k_i \quad \hbar\omega = E_f - E_i \quad E_{f,i} = \hbar^2 k^2 / 2m_n$$

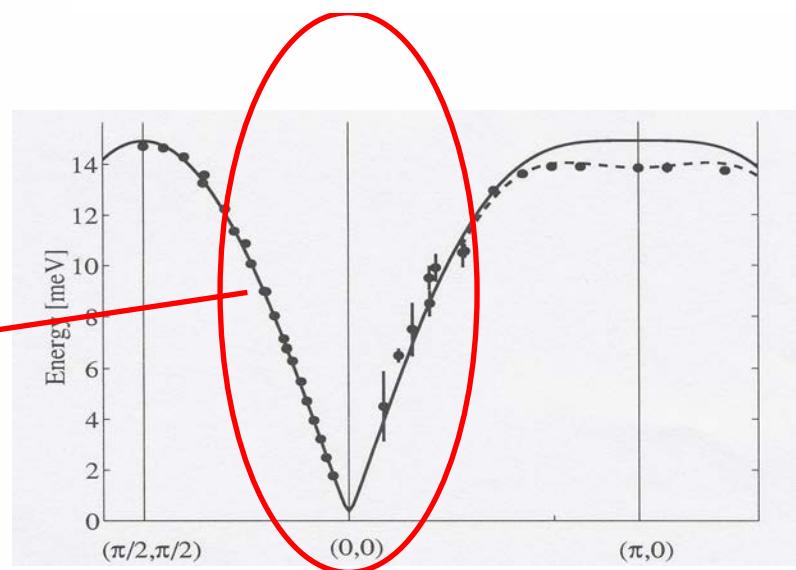
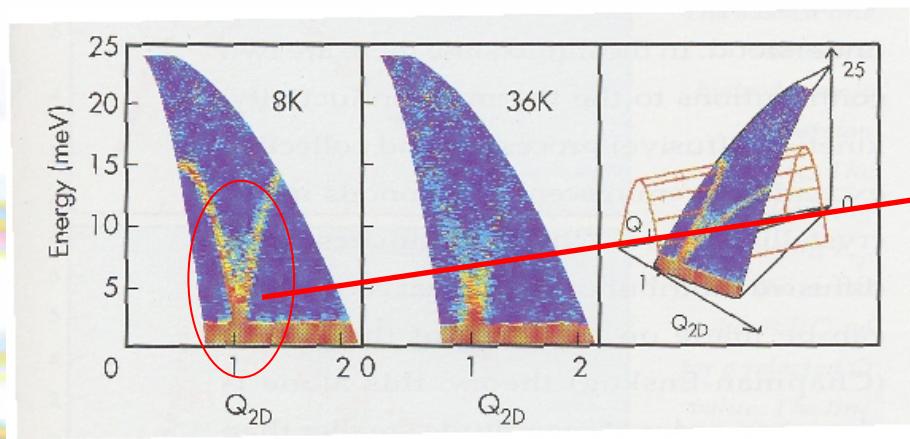
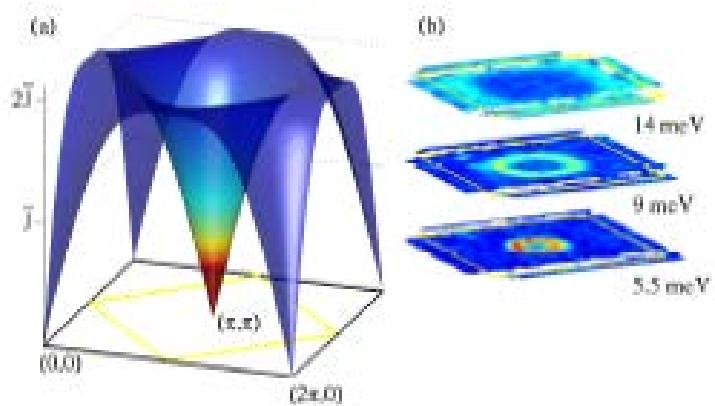
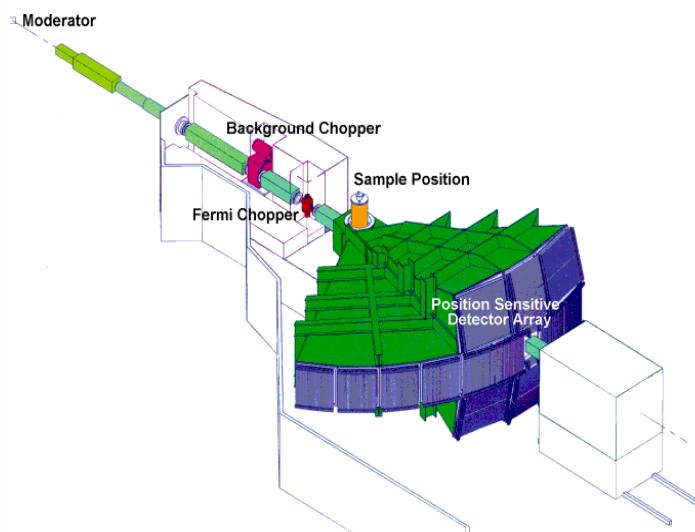


$$\lambda = \infty \quad \lambda = 2a$$



INS: experimental methods

- With time-of-flight methods and large area detector, can map whole regions of Energy and Q in one fell swoop- often revealing unexpected features

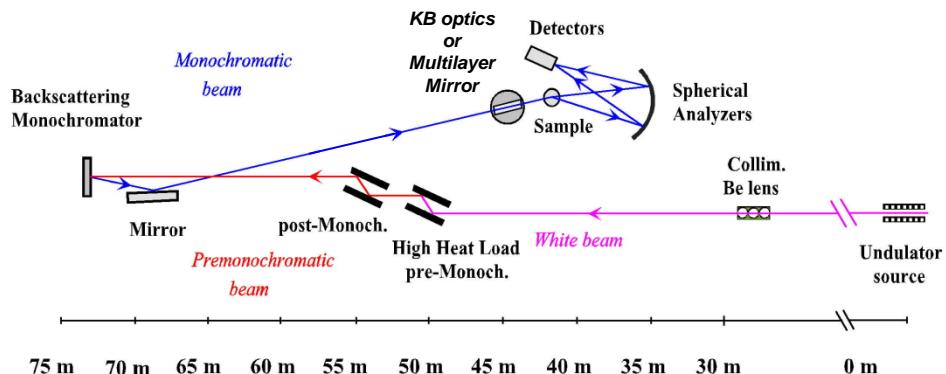


IXS: experimental methods

- Far more challenging to resolve such excitations by inelastic X-ray scattering

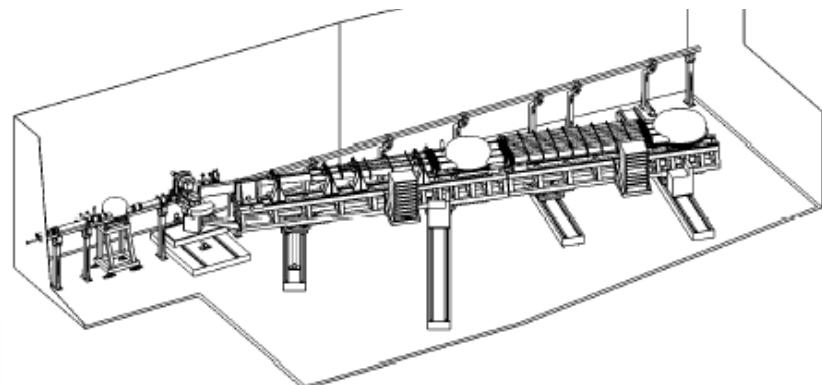
Optimise resolution with monochromator and analyser close to backscattering configurations:

$$d\lambda/\lambda = dE/E = \cot\theta d\theta \rightarrow 0 \text{ as } \theta \rightarrow \pi/2$$



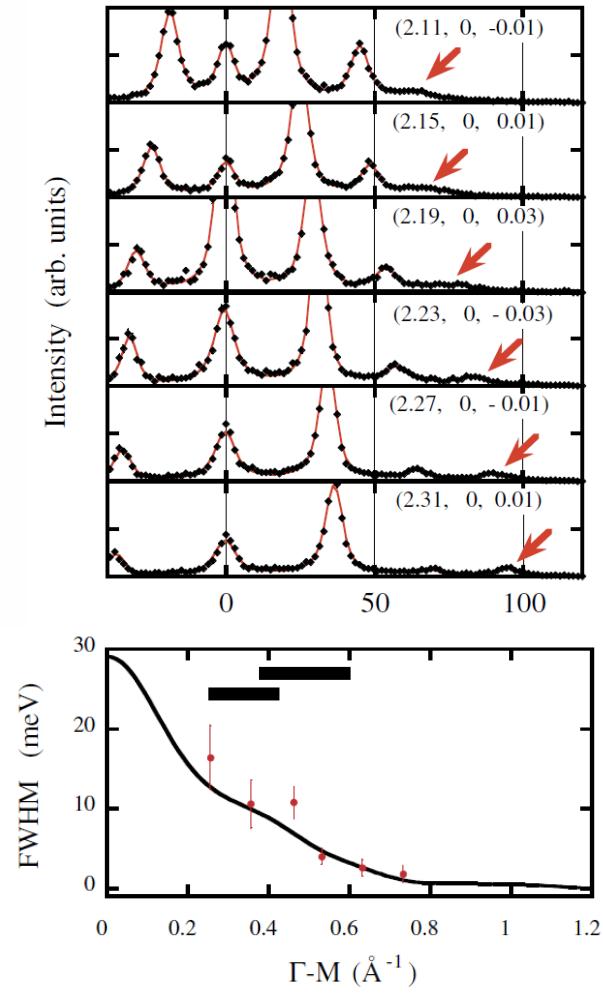
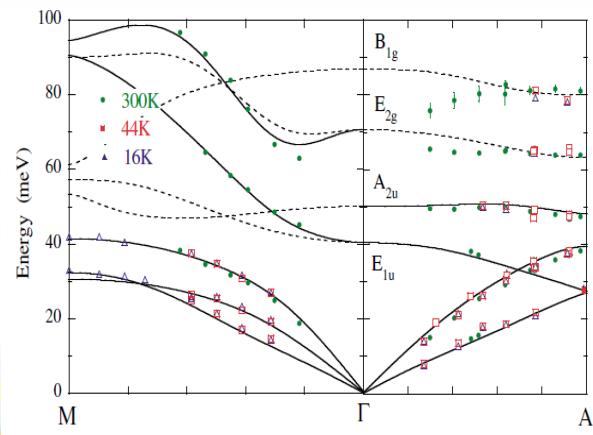
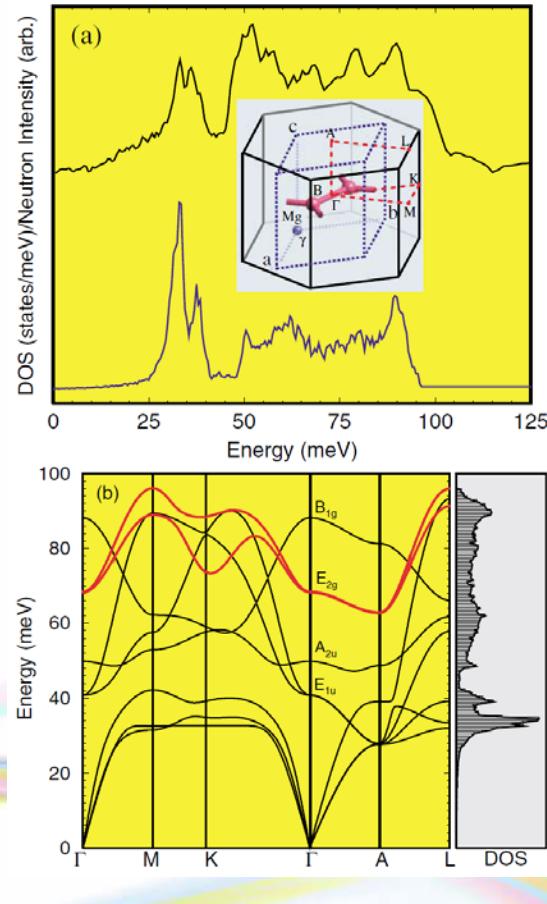
9- analyser crystal spectrometer

Reflection	Energy[keV]	E [meV]	dE/E
Si(7,7,7)	13.840	5.30	$3.8 \cdot 10^{-7}$
Si(8,8,8)	15.817	4.40	$2.8 \cdot 10^{-7}$
Si(9,9,9)	17.794	2.20	$1.2 \cdot 10^{-7}$
Si(11,11,11)	21.747	0.83	$4.7 \cdot 10^{-8}$
Si(12,12,12)	23.725	0.73	$3.1 \cdot 10^{-8}$
Si(13,13,13)	25.704	0.50	$1.9 \cdot 10^{-8}$



IXS: applications

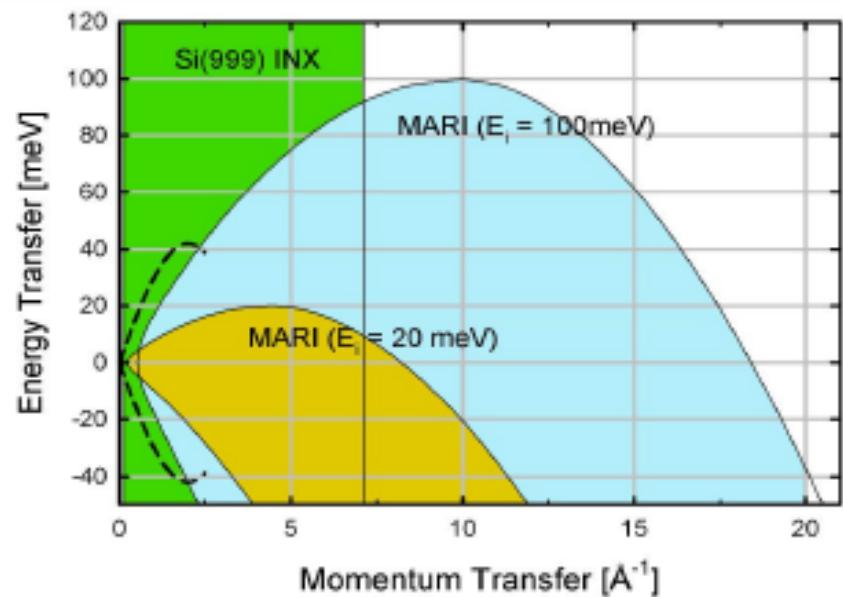
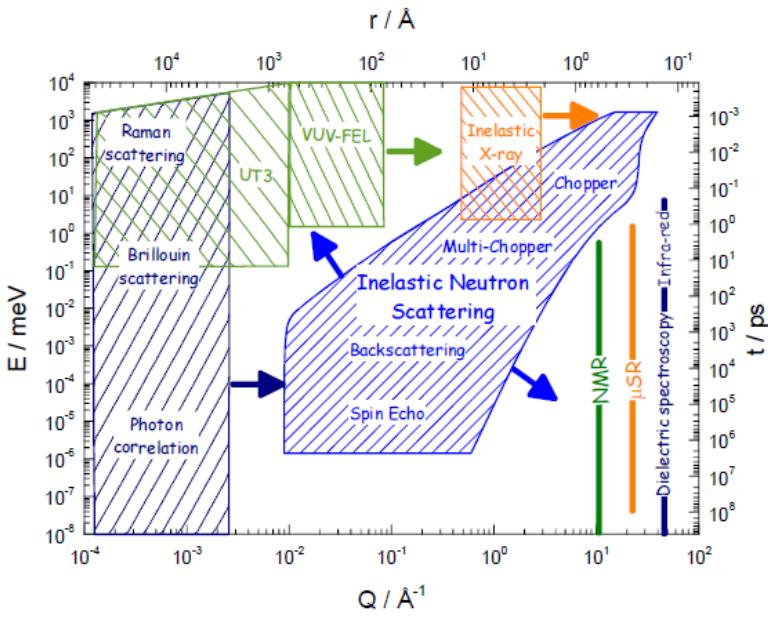
- Role of phonons in superconductivity in MgB₂ – DOS from lattice dynamics and neutron scattering measurements from powders then IXS on 50 x 70 μ² crystals



Yildirim et al, PRL 87 (2001) 37001
Baron et al, PRL 92 (2004) 197004

IXS vs INS

- Advantages of IXS
 - Fine beams - $\sim 10 \mu\text{m}$ - work with very small samples
 - Resolution is independent of energy transfer as ΔE and Q are decoupled (below)
 - Open up new kinematical window in ΔE , Q at small Q and high ΔE –useful for disordered systems
 - No incoherent scattering
- Disadvantages of IXS
 - Very challenging instrumentation to achieve resolution $\Delta E/E \sim 10^{-7}$
 - Cross-section dominated by photoelectric absorption - less sensitive to light atoms
 - Resolution more complex than for neutrons (non Gaussian)

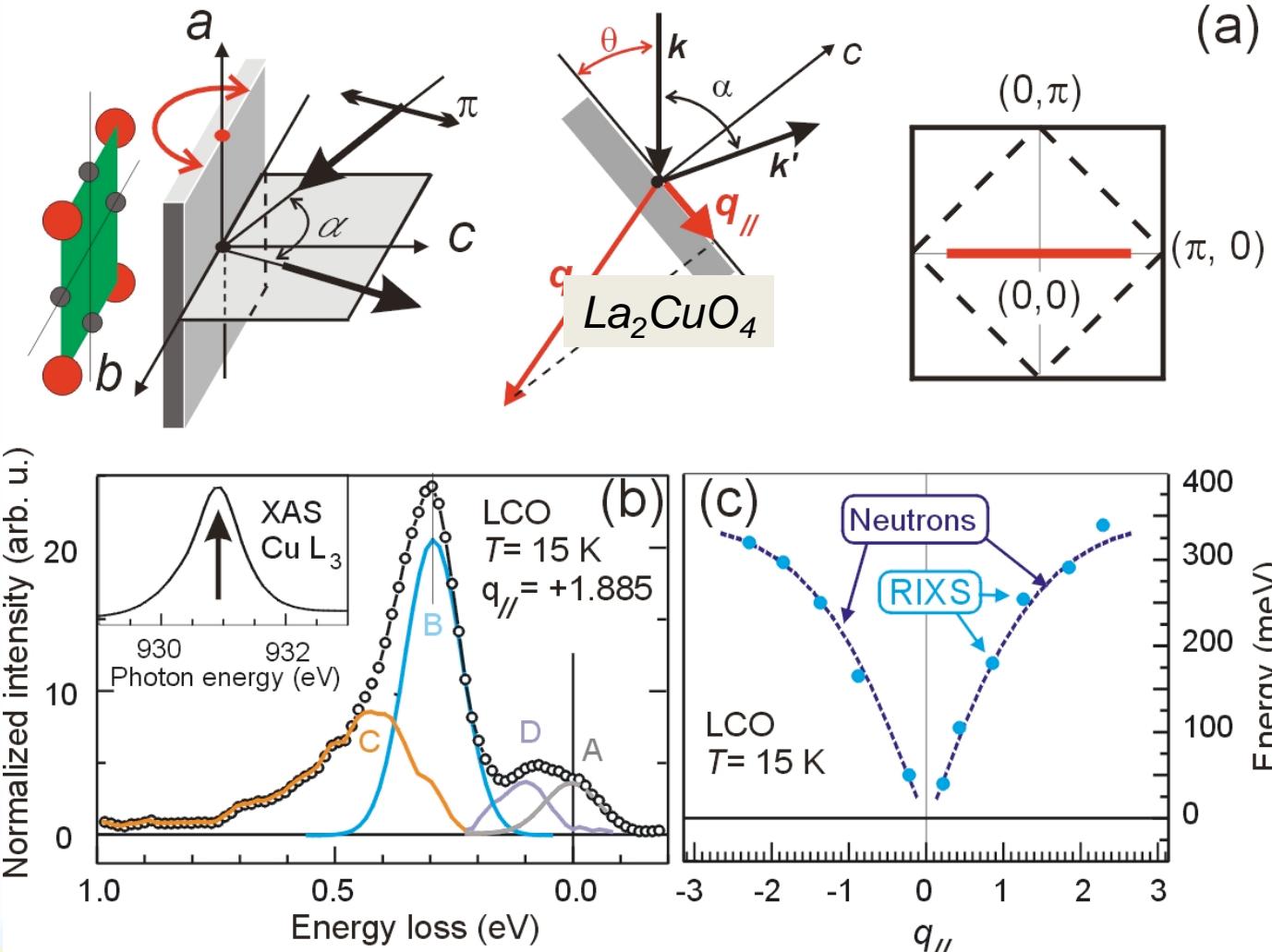


IXS + R \rightarrow RIXS: a complementary method to neutrons for studying magnetic excitations

High sensitivity of soft x-rays allows study of small sample volumes.

Tune to specific transition – element and orbital selective

La_2CuO_4
100nm films



Braicovich et al. PRL 104, 077002 (2010)



RIXS: Analysing spectrometer at the ESRF



2015: ERIXS, ID32 – ESRF/Milano

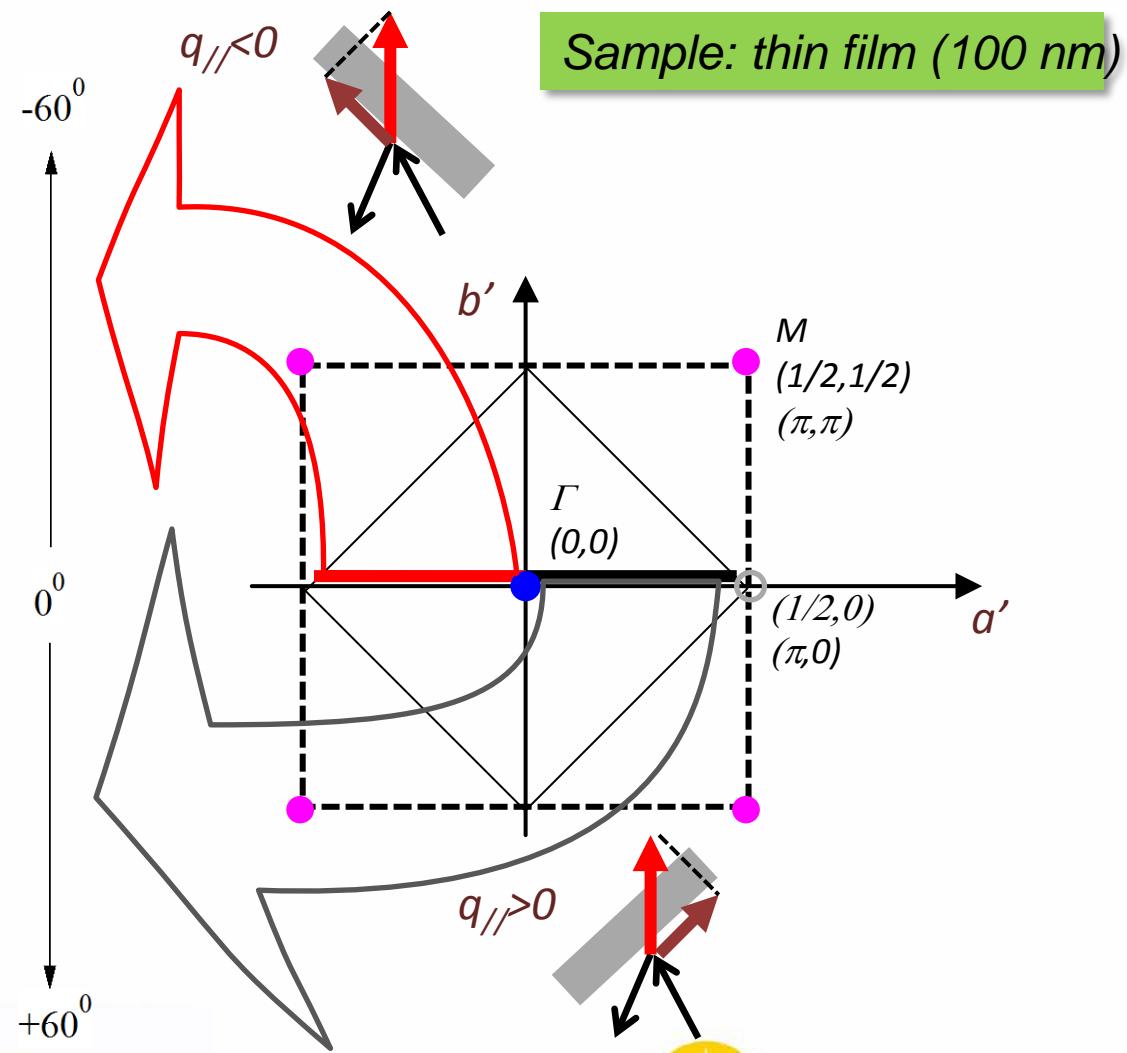
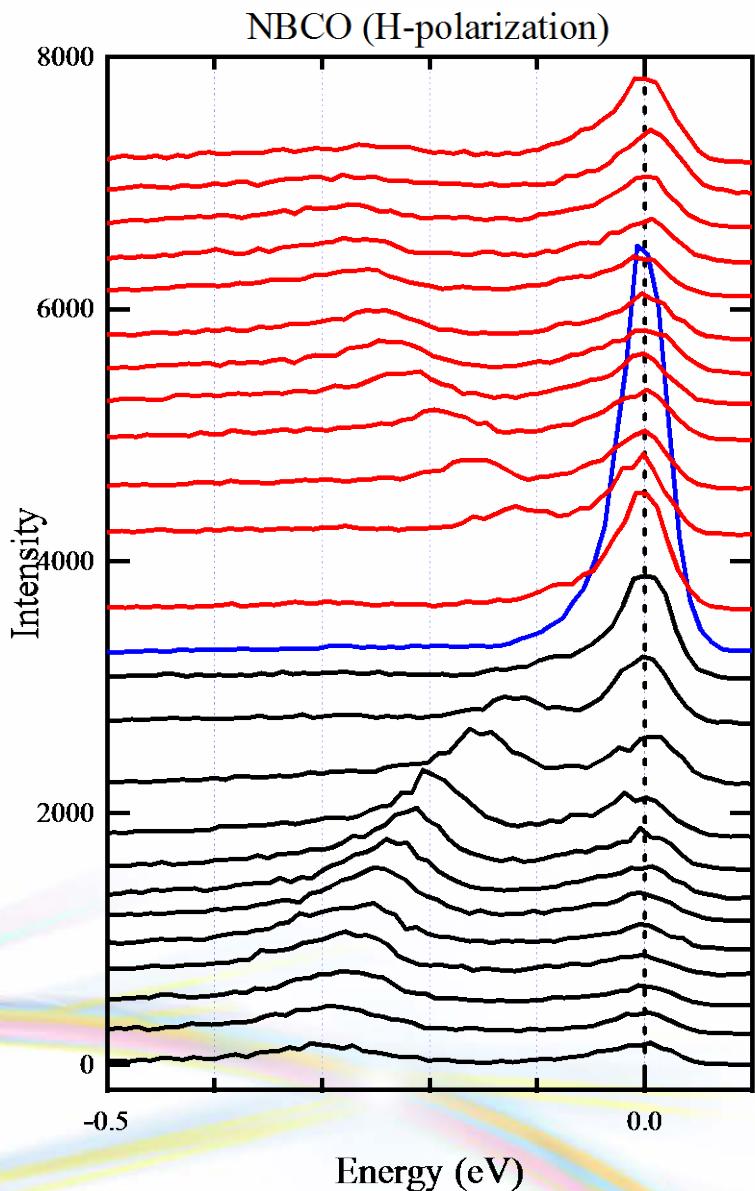
$L = 12.0 \text{ m}$

$\text{Design: } E/\Delta E = 42000 \text{ at Cu-}L_3$



nd

Antiferromagnetic $\text{NdBa}_2\text{Cu}_3\text{O}_{6+\delta}$



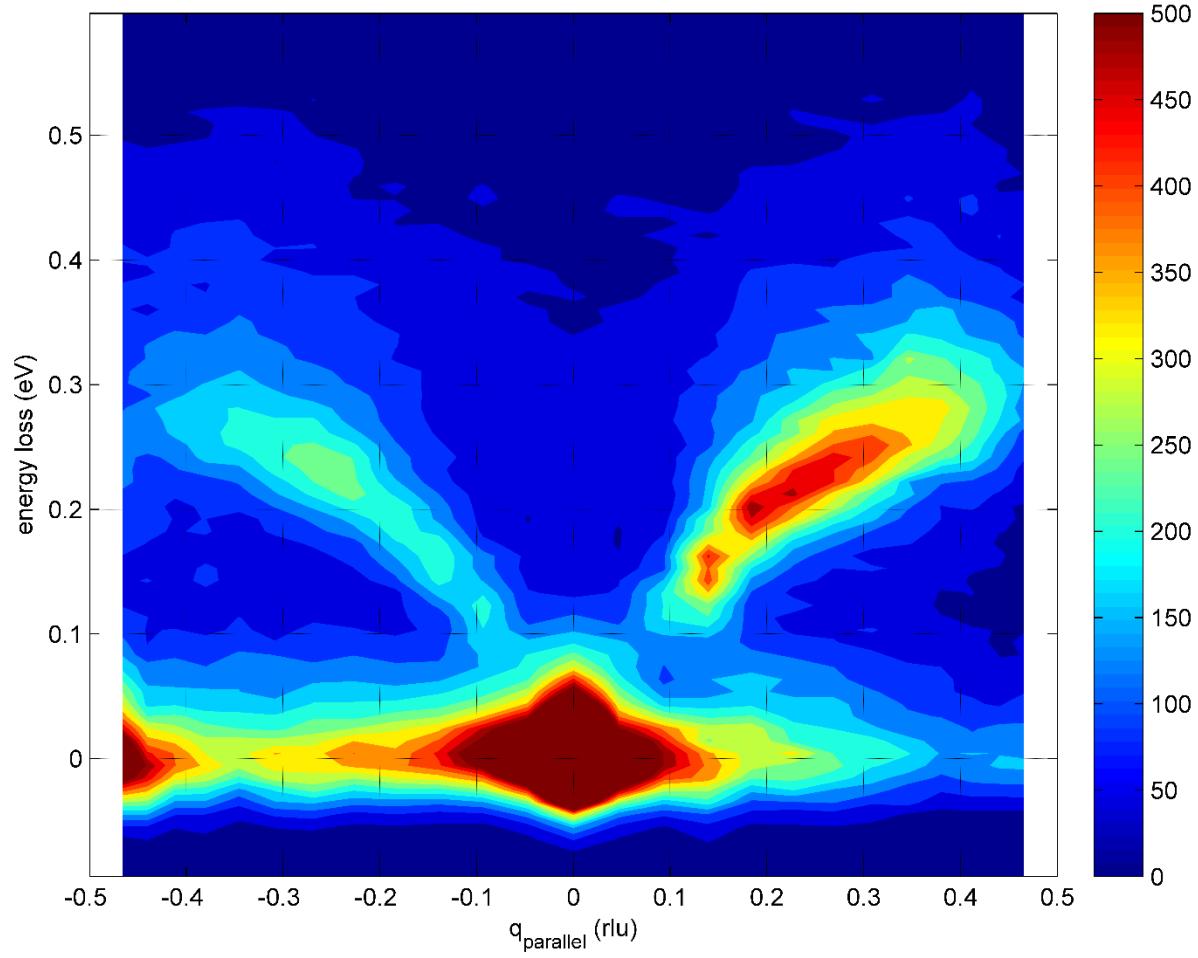
"Low" resolution set up at ID32 – July 2015

Combined BW ≈ 55 meV

diamond

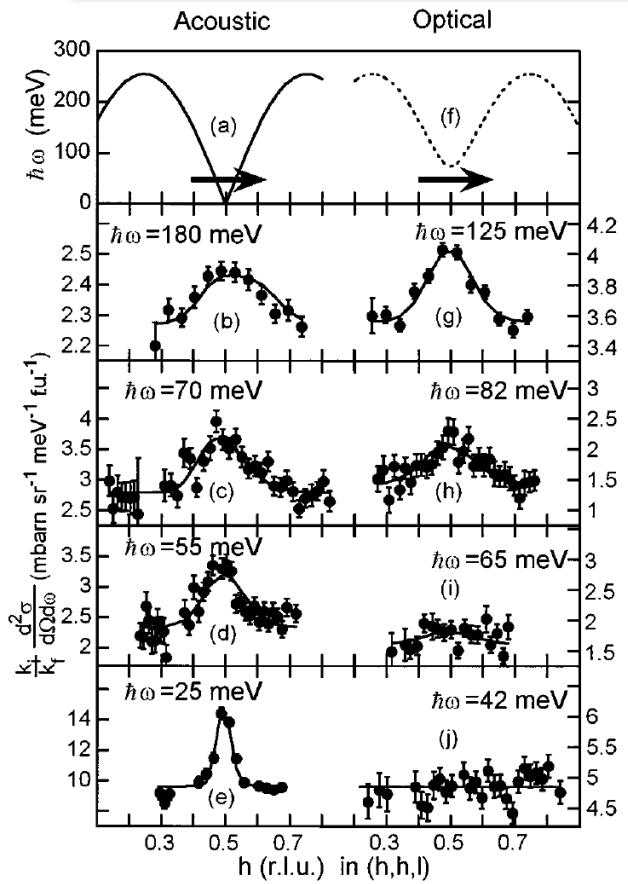
AF Nd₂Ba₂Cu₃O_{6+δ}

Sample: thin film (100 nm)



AF YBCO and NBCO: magnon optical branch

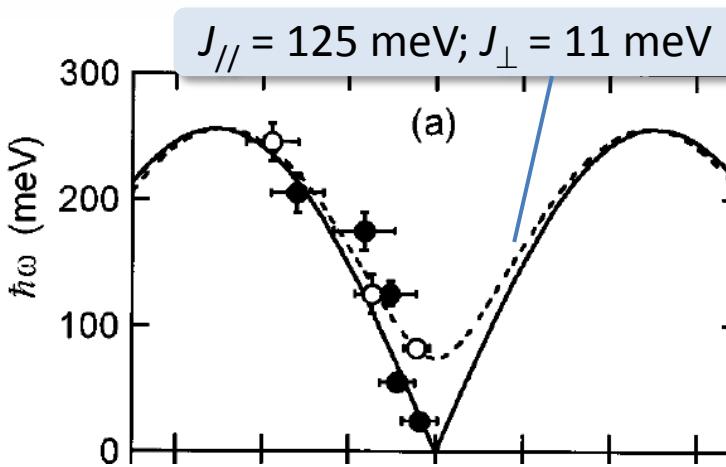
INS: BW ≈ 2 meV



The sample used in the present investigation was a single crystal of $YBa_2Cu_3O_{6.15}$ with mass 96 g.

S. M. Hayden et al, PRB 54 6905 (1996)

$J_{//} = 125$ meV; $J_{\perp} = 11$ meV

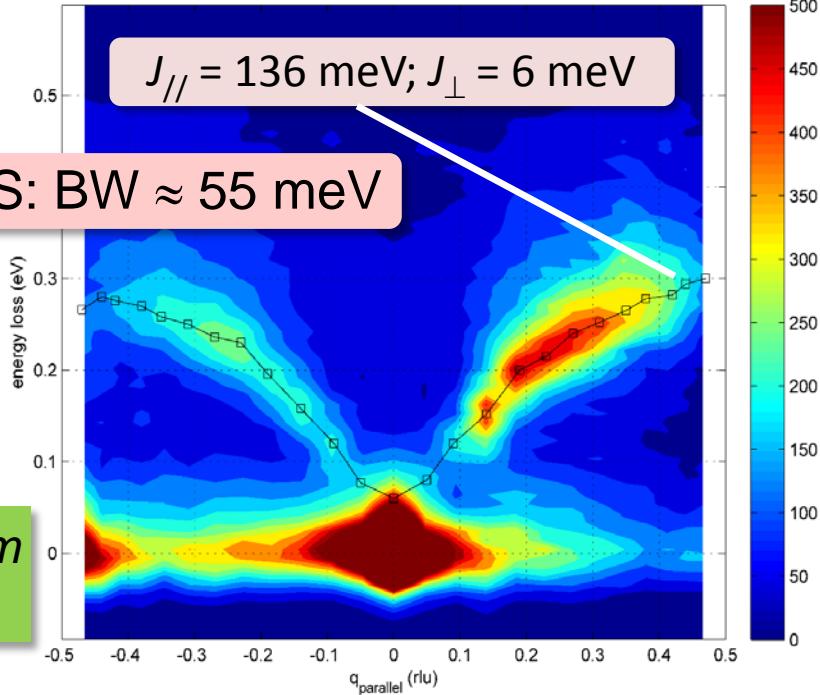


ESRF
ERIXS@ID32



$J_{//} = 136$ meV; $J_{\perp} = 6$ meV

RIXS: BW ≈ 55 meV



Y.Y. Peng, G. Dellea, M. Minola, G.M. De Luca, M. Salluzzo, M. Le Tacon, B. Keimer, L. Braicovich, N.B. Brookes, and G. Ghiringhelli, unpublished

Neutrons and X-rays



Andrew Harrison, Diamond Light Source
14th Oxford School on Neutron Scattering - 2015

Further reading

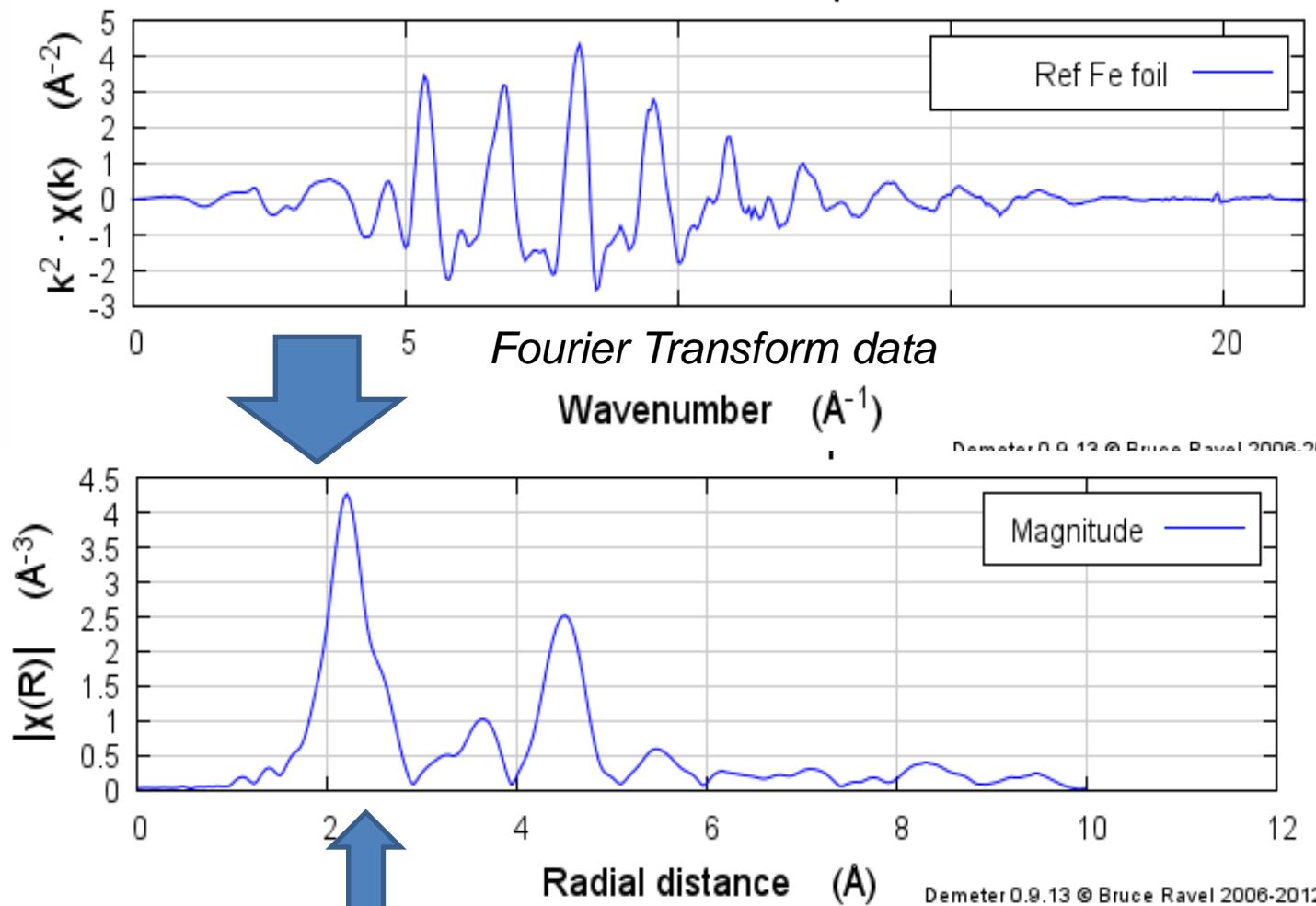
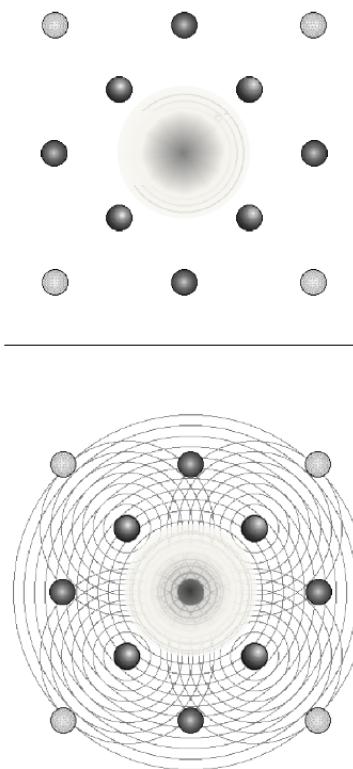
- Introductions
 - P. Willmott, '*An introduction to synchrotron radiation*', Wiley 2011
 - J. Als-Nielsen and D. McMorrow, '*Elements of modern X-ray physics*', Wiley 2011
- Magnetic X-ray scattering
 - S.W.Lovesey and S.P. Collins, '*X-ray scattering and absorption by magnetic materials*', Oxford University Press, 1996
 - G. van der Laan, C.R. Physique 9 (2008) 570
 - Look out the Paolasini review
 - T. Bruckel review
- Inelastic X-ray scattering
 - A.Q.R. Baron, *Introduction to High-Resolution Inelastic X-Ray Scattering*,
[arXiv:1504.01098 \[cond-mat.mtrl-sci\]](https://arxiv.org/abs/1504.01098)
 - Rev. Mod Phys. review

EXAFS

- Extended X-Ray Absorption Fine Structure

EXAFS Oscillations

*come from
photo-electron
scattering
and interference
in the material*

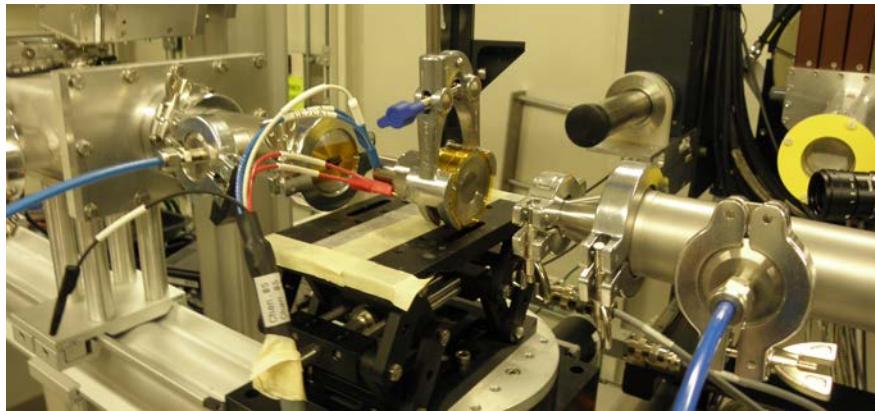


*Detecting Fe atoms AROUND absorber Fe
Position: At about 2.5 Angstrom
Coordination: 8 next neighbours*



B18 – general purpose XAS

- 2.05 – 35 keV – covers absorption edge from P to U with 200 x 250 μm spot size



K-edges

0	He
1	H
2	Li
3	Be
4	Mg
5	Na
6	Al
7	Si
8	P
9	S
10	Cl
11	Ar
12	
13	
14	
15	
16	
17	
18	
19	K
20	Ca
21	Sc
22	Ti
23	Y
24	Cr
25	Mn
26	Fe
27	Co
28	Ni
29	Cu
30	Zn
31	Ga
32	Ge
33	As
34	Se
35	Br
36	Kr
37	Rb
38	Sr
39	Y
40	Zr
41	Nb
42	Mo
43	Tc
44	Ru
45	Rh
46	Pd
47	Ag
48	Cd
49	In
50	Sn
51	Sb
52	Te
53	I
54	Xe
55	Cs
56	Ba
57	*La
58	Hf
59	Ta
60	W
61	Re
62	Os
63	Ir
64	Pt
65	Au
66	Hg
67	Tl
68	Pb
69	Bi
70	Po
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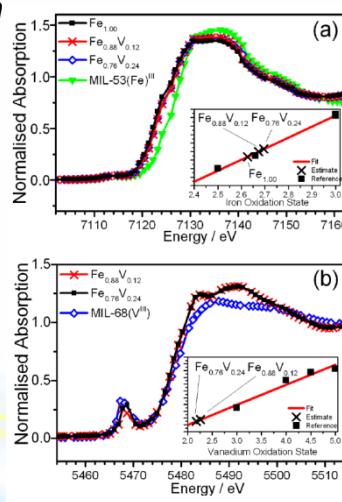
*	Lanthanide Series	58	59	60	61	62	63	64	65	66	67	68	69	70	71	Lu
+	Actinide Series	90	91	92	93	94	95	96	97	98	99	100	101	102	103	U

L-edges
M4,5 edges

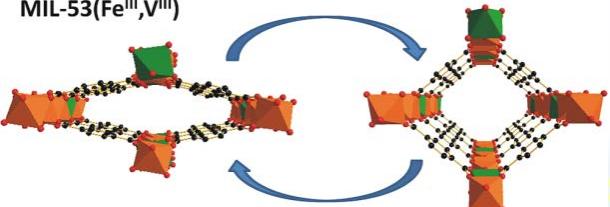
0	He
1	H
2	Li
3	Be
4	Mg
5	Al
6	Si
7	P
8	S
9	Cl
10	Ar
11	
12	
13	
14	
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18	
19	K
20	Ca
21	Sc
22	Ti
23	Y
24	Cr
25	Mn
26	Fe
27	Co
28	Ni
29	Cu
30	Zn
31	Ga
32	Ge
33	As
34	Se
35	Br
36	Kr
37	Rb
38	Sr
39	Y
40	Zr
41	Nb
42	Mo
43	Tc
44	Ru
45	Rh
46	Pd
47	Ag
48	Cd
49	In
50	Sn
51	Sb
52	Te
53	I
54	Xe
55	Cs
56	Ba
57	*La
58	Hf
59	Ta
60	W
61	Re
62	Os
63	Ir
64	Pt
65	Au
66	Hg
67	Tl
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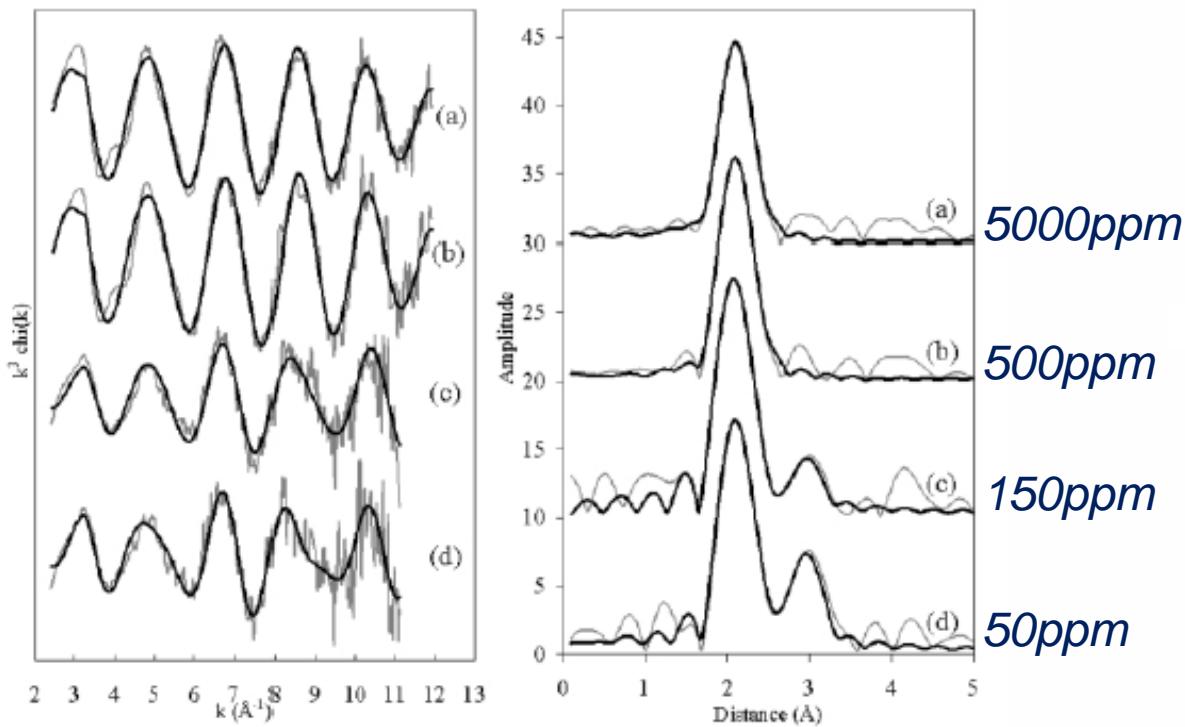
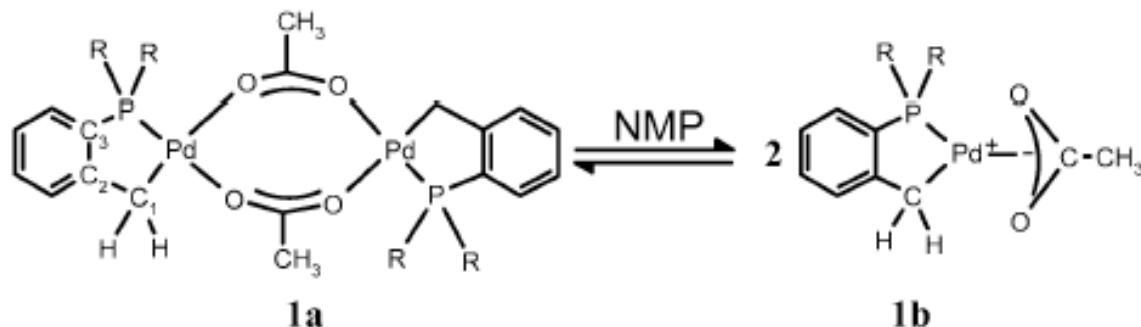
Determining oxidation states of Fe and V in flexible mixed valent MOF by XANES



MIL-53(Fe^{III},V^{IV})



What the Heck ?

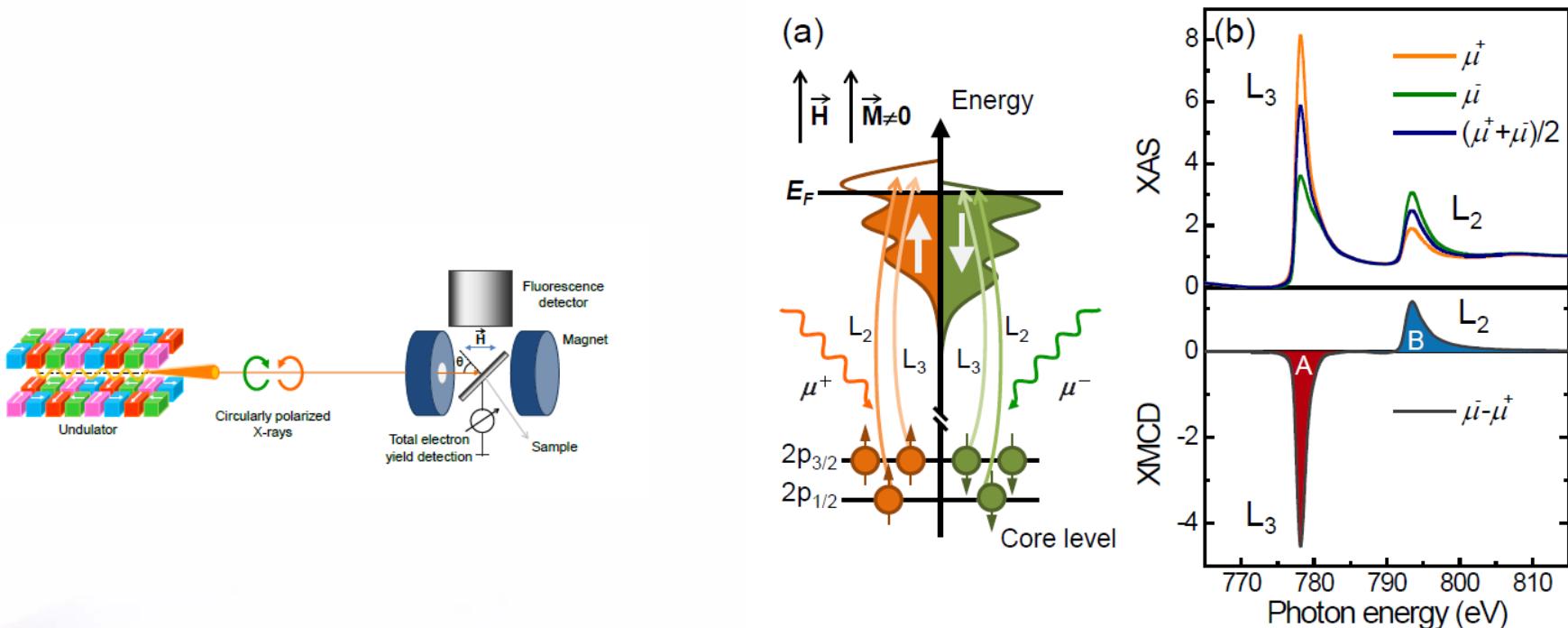


*Future limits ?
I20 at Diamond
will give ppb
of chemical
species*



XAS plus magnetic polarisation

- XAS with circularly polarised (CP) light: LCP-RCP probes polarisation of spin states by selectively exciting transitions with $\Delta m = +/- 1 \rightarrow$ XMCD

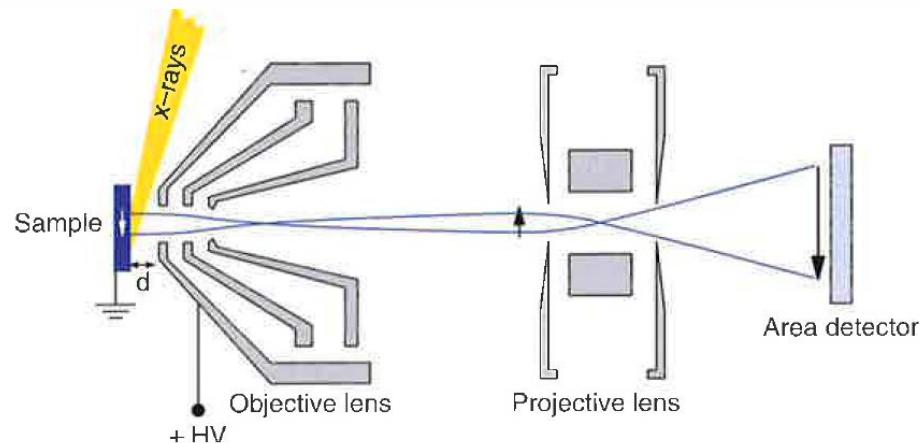
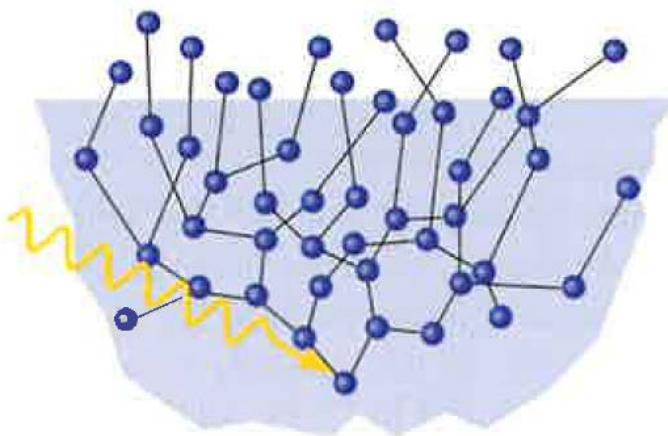


- XAS with LP radiation: $\Delta (\perp-//) \sim \langle M^2 \rangle \rightarrow \text{XMLCD}$

G van der Laan and A.I. Figueroa,
Coord. Chem. Reviews 277-278 (2014) 95

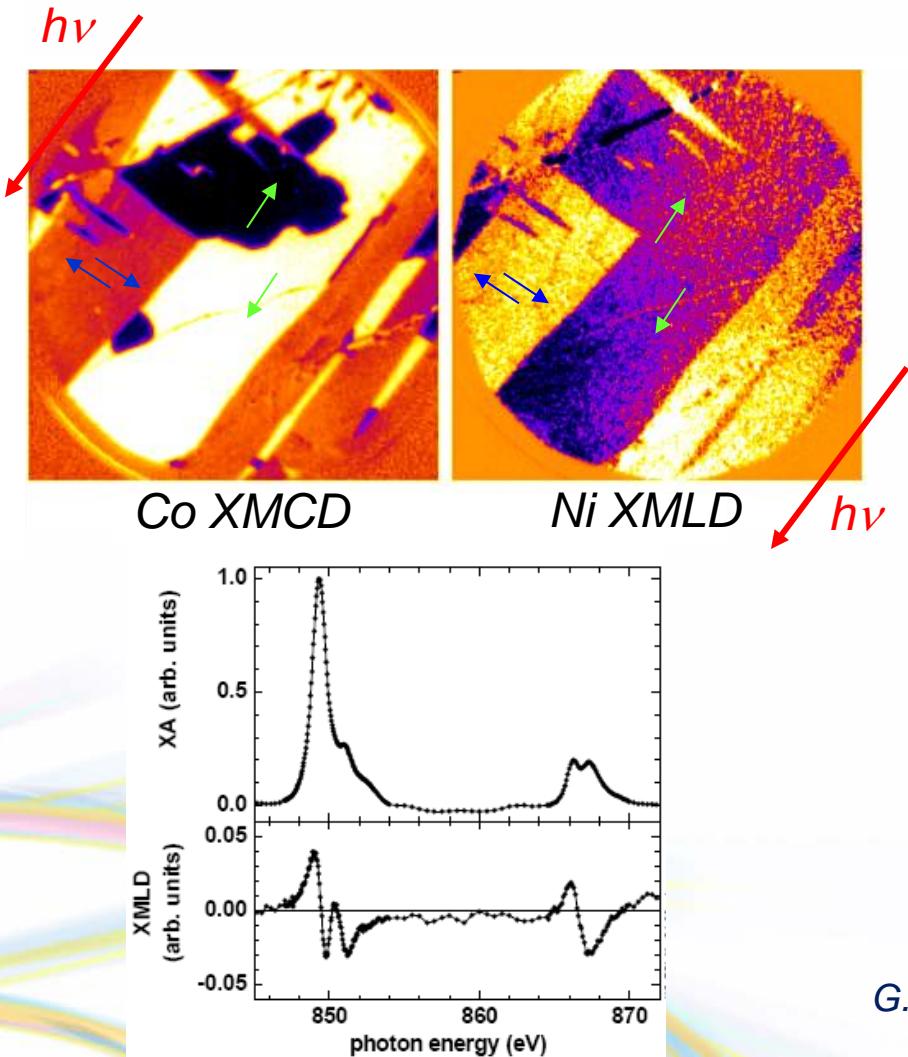
Mapping nanomagnetism

- Resolution to 50 nm with Photoemission Electron Microscopy (PEEM)
- XAS in XANES region produces photoelectrons which create secondary, low-energy (eV) photoelectrons with longer mean-free path. Spatial resolution for imaging these e^- is 10's of nm



Mapping nanomagnetism

- PEEM in combination with XMCD, XMLD to map domains to nm lengthscales



- Co/NiO is prototype exchange bias material – need to understand coupling at interface.
- PEEM at I06 with XMCD for (FM) Co L3 and XMLD for (AFM) Ni L2 shows orientation of the two types of moments at the interface is perpendicular – opposite of what was inferred from less precise measurements

G. van der Laan et al., in prep.

