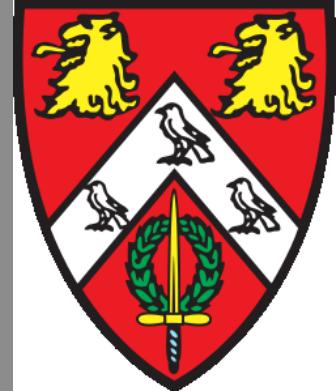


# 13<sup>th</sup> Oxford School on Neutron Scattering

St. Anne's College, University of Oxford

2 - 13 September 2013



## Chemical Applications of Neutron Scattering

*Part 1*

*Elastic Scattering and  
Structural Studies*

Alberto Albinati

University of Milan, Department of Chemistry

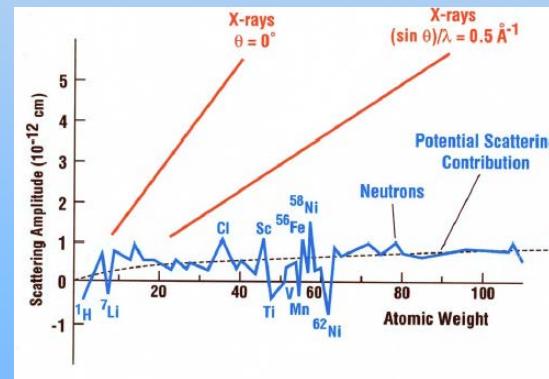
# Neutron Properties

- @ 293.6 K for a “neutron gas”:

$$E = 25 \text{ meV}, \quad v \approx 2.2 \text{ km s}^{-1}, \quad v \approx 200 \text{ cm}^{-1} \approx 6 \times 10^{12} \text{ Hz}, \quad \lambda = h/mv \approx 1.8 \text{ \AA}$$

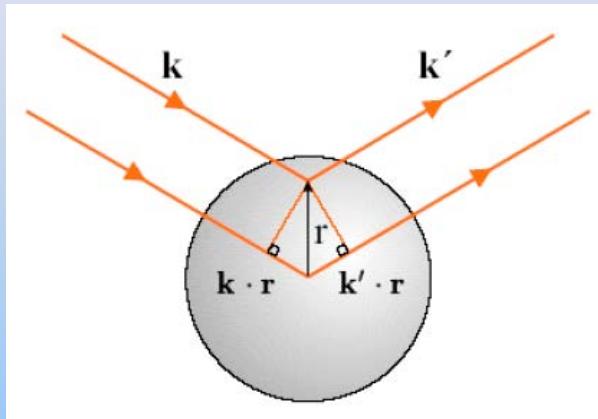
- Zero Electric Charge: *Negligible absorption, scattering from the bulk.*

- Neutron Scattering is a Nuclear Process:



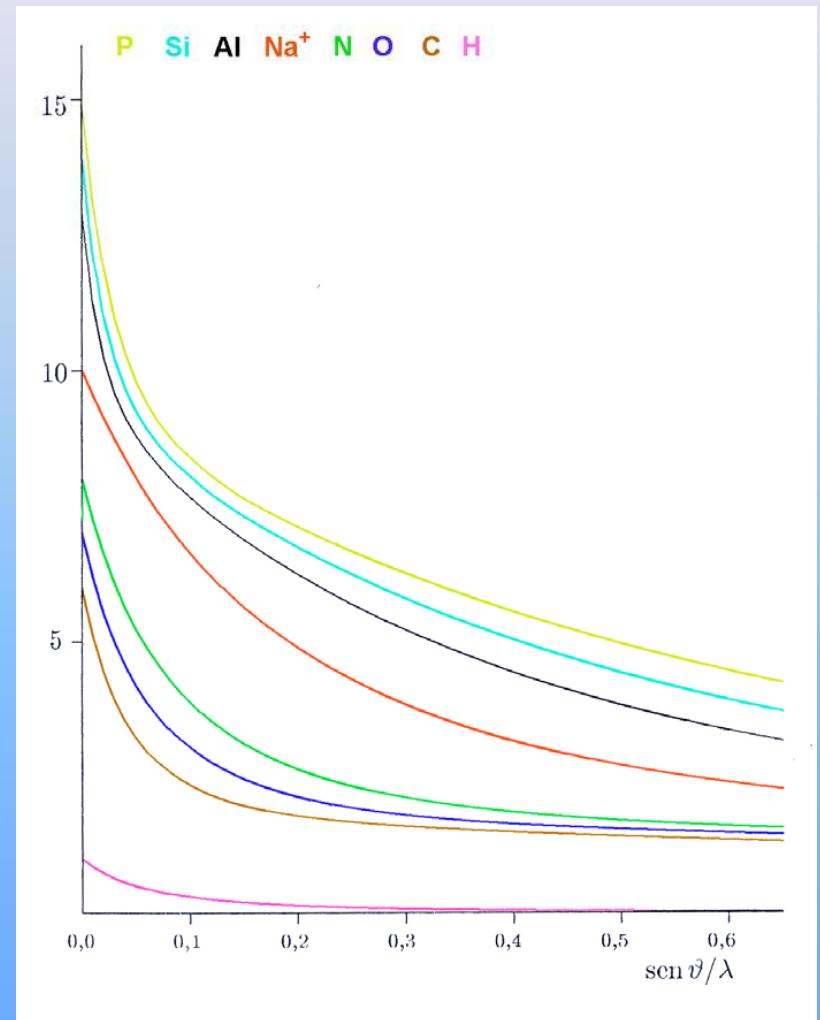
- De Broglie Wavelength is Comparable with Interatomic Distances: *Bragg Scattering.*
- The Energy of Thermal Neutrons is Comparable with the Energy of Molecular and Lattice Vibrations (Phonons): *Inelastic Coherent & Incoherent Scattering Can Probe Lattice and Molecular Vibrations*

# Atomic Scattering Factors for X-rays

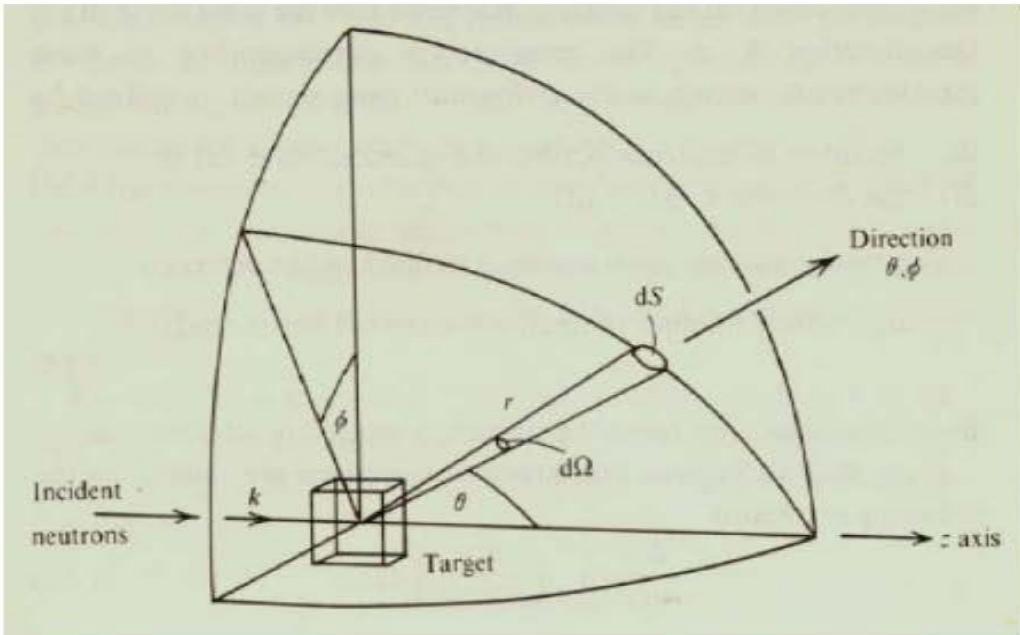


$$f(\mathbf{S}) = \int \rho(\mathbf{r}) \exp 2\pi i (\mathbf{S} \cdot \mathbf{r}) d\mathbf{r}$$

$$f(S) = 4\pi \int_0^{\infty} r^2 \rho(r) \frac{\sin 2\pi(Sr)}{2\pi(Sr)} dr$$



# Cross Sections

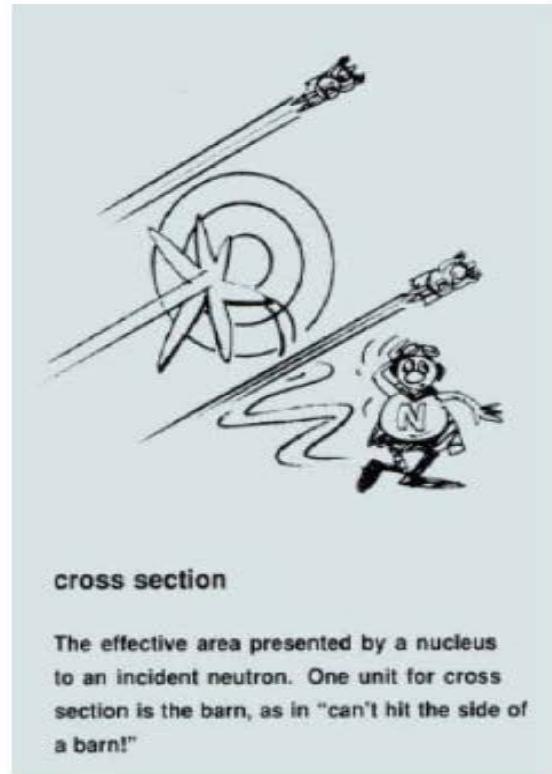


$\Phi$  = number of incident neutrons per  $\text{cm}^2$  per second

$\sigma$  = total number of neutrons scattered per second /  $\Phi$

$$\frac{d\sigma}{d\Omega} = \frac{\text{number of neutrons scattered per second into } d\Omega}{\Phi d\Omega}$$

$$\frac{d^2\sigma}{d\Omega dE} = \frac{\text{number of neutrons scattered per second into } d\Omega \& dE}{\Phi d\Omega dE}$$



$\sigma$  measured in barns:

$$1 \text{ barn} = 10^{-24} \text{ cm}^2$$

$$\text{Attenuation} = \exp(-N\sigma t)$$

N = # of atoms/unit volume

t = thickness

# COHERENT and INCOHERENT NEUTRON SCATTERING

“ $b_R$ ” varies from isotope to isotope

$$\frac{d\sigma}{d\Omega} = \left| \sum_{\mathbf{R}} b_{\mathbf{R}} \exp(i\mathbf{Q} \cdot \mathbf{R}) \right|^2 = \sum_{\mathbf{R}} \sum_{\mathbf{R}'} b_{\mathbf{R}} b_{\mathbf{R}'} \exp[i\mathbf{Q} \cdot (\mathbf{R} - \mathbf{R}')] = \\ = \sum_R b_R^2 + \sum_{\mathbf{R}, \mathbf{R}'} b_{\mathbf{R}} b_{\mathbf{R}'} \exp[i\mathbf{Q} \cdot (\mathbf{R} - \mathbf{R}')] \quad \text{(in red)}$$

$$\sum_{\mathbf{R}} b_{\mathbf{R}}^2 = N \langle b_{\mathbf{R}}^2 \rangle \quad \text{(in red)}$$

assuming no correlation  $\langle b_{\mathbf{R}} b_{\mathbf{R}'} \rangle = \langle b_{\mathbf{R}} \rangle \langle b_{\mathbf{R}'} \rangle = \langle b_{\mathbf{R}} \rangle^2$

$$\sum_{RR'} b_{\mathbf{R}} b_{\mathbf{R}'} \exp[i\mathbf{Q} \cdot (\mathbf{R} - \mathbf{R}')] = \\ = N \langle b_{\mathbf{R}} \rangle^2 \sum_{RR'} \exp[i\mathbf{Q} \cdot (\mathbf{R} - \mathbf{R}')] = -N \langle b_{\mathbf{R}} \rangle^2 + N \langle b_{\mathbf{R}} \rangle^2 \sum_R \sum_{R'} \exp[i\mathbf{Q} \cdot (\mathbf{R} - \mathbf{R}')] \quad \text{(in blue)}$$

$$\frac{d\sigma}{d\Omega} = N \left( \langle b^2 \rangle - \langle b \rangle^2 \right) + N \langle b \rangle^2 \left| \sum_R \exp(i\mathbf{Q} \cdot \mathbf{R}) \right|^2 \quad \text{(in magenta)}$$

$$N \left( \langle b^2 \rangle - \langle b \rangle^2 \right) = N \left\langle (b - \langle b \rangle)^2 \right\rangle \quad \text{Incoherent Scattering Cross Section}$$

# Neutron Scattering Cross Sections

## COHERENT CROSS SECTIONS

$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

Coherent Scattering depends on the correlation between the positions of the **same nucleus at different times** and the positions of different nuclei at different times.

**STRUCTURAL INFORMATION.**

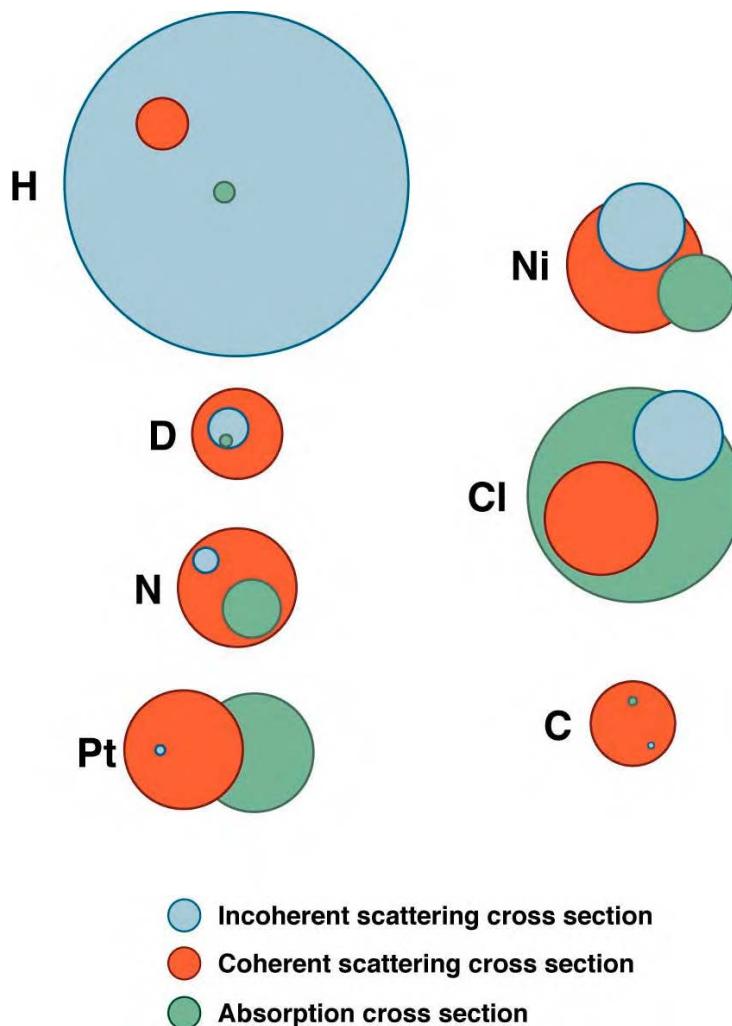
## INCOHERENT CROSS SECTIONS

$$\sigma_{incoh} = 4\pi \left( \langle b^2 \rangle - \langle b \rangle^2 \right)$$

Incoherent Scattering arises from the random distribution of different isotopes with different scattering lengths. Incoherent Scattering depends on the correlation between the positions of the same nucleus at different times.

**SPECTROSCOPY**

# Neutron Incoherent Cross-Sections



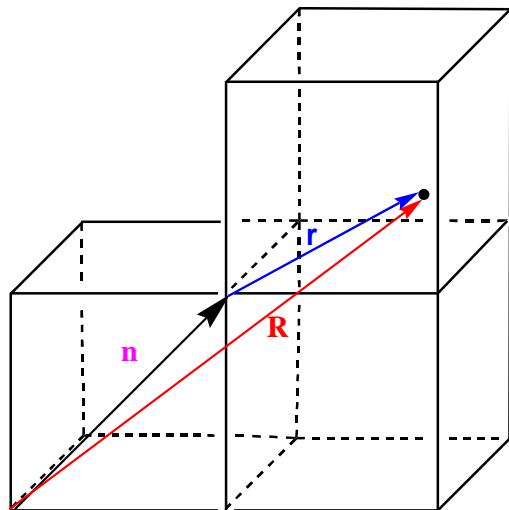
# Bragg Scattering in Crystals

$$\mathbf{R} = \mathbf{n} + \mathbf{r}$$

$\mathbf{R} \equiv$  atomic position

$\mathbf{n} \equiv$  lattice vector  $\equiv n_1 \mathbf{a}_1 + n_2 \mathbf{a}_2 + n_3 \mathbf{a}_3$

$\mathbf{r}_{(xyz)}$   $\equiv$  atomic position in the unit cell



$$|\mathbf{Q}| = \frac{4\pi \sin \theta}{\lambda}$$

$$|\mathbf{k}| = \frac{2\pi}{\lambda}$$

$$I = \left( \frac{d\sigma}{d\Omega} \right) = \left| \sum_{\mathbf{R}} b_R e^{i\mathbf{Q} \cdot \mathbf{R}} \right|^2 = \left| \sum_{\mathbf{n}} \exp i(\mathbf{Q} \cdot \mathbf{n}) \sum_{\mathbf{r}} b_r \exp i(\mathbf{Q} \cdot \mathbf{r}) \right|^2 = \left| \sum_{\mathbf{n}} \exp i(\mathbf{Q} \cdot \mathbf{n}) \right|^2 \left| \sum_{\mathbf{r}} b_r \exp i(\mathbf{Q} \cdot \mathbf{r}) \right|^2$$

# Bragg Scattering in Crystals

$$\left| \sum_{\mathbf{n}} \exp i(\mathbf{Q} \cdot \mathbf{n}) \right|^2 = \frac{\sin^2 h N_1 \pi}{\sin^2 h \pi} \frac{\sin^2 k N_2 \pi}{\sin^2 k \pi} \frac{\sin^2 l N_3 \pi}{\sin^2 l \pi}$$

if  $\mathbf{Q} = \mathbf{H}$

$$I = \left( \frac{d\sigma}{d\Omega} \right) = N_1 N_2 N_3 \frac{(2\pi)^3}{V_{cell}} \sum_{\mathbf{H}} \delta(\mathbf{Q} \cdot \mathbf{H}) |F_{\mathbf{H}}|^2$$

$$F_{\mathbf{H}} = \sum_{\mathbf{r}} b_r^{coh} \exp i(\mathbf{H} \cdot \mathbf{r}) \quad F_{\mathbf{H}} \equiv \text{Structure Factor}$$

In Crystallography:

$$F_{hkl} = \sum_j b_j^{coh} \exp 2\pi i (\mathbf{S} \cdot \mathbf{r}_j) = \sum_j b_j^{coh} \exp 2\pi i (hx_j + ky_j + lz_j) \quad \mathbf{S} = h\mathbf{a}_1^* + k\mathbf{a}_2^* + l\mathbf{a}_3^*$$
$$\mathbf{r} = x\mathbf{a}_1 + y\mathbf{a}_2 + z\mathbf{a}_3$$
$$|\mathbf{S}| = \frac{2\pi \sin \theta}{\lambda}$$

# Bragg Scattering in Crystals

and finally atoms do "move"....

$$F_{\mathbf{H}} = \sum_{\mathbf{r}} b_r \exp i(\mathbf{H} \cdot \mathbf{r}) T_r(\mathbf{Q})$$

where " $T_r(\mathbf{Q})$ " is the **temperature factor** of atom "r".

for an *harmonic crystal*

$$T_r(\mathbf{Q}) = \exp(-2\pi\langle \mathbf{Q} \cdot \mathbf{r} \rangle^2)$$

## The Temperature Factor and the P.D.F. (Probability Density Function)

$$T(Q) = \exp\{-\frac{1}{2}[\langle(Q \cdot u)^2\rangle]\}$$

p.d.f. of an atom ( $p_\kappa(u)$ ) is the probability of finding an atom in the volume element  $d^3u$  when it is displaced by  $u$  from its rest position.

If  $\rho_0(u)$  ≡ scattering density

$$\rho_\kappa(u) = \rho_0(u) * p_\kappa(u)$$

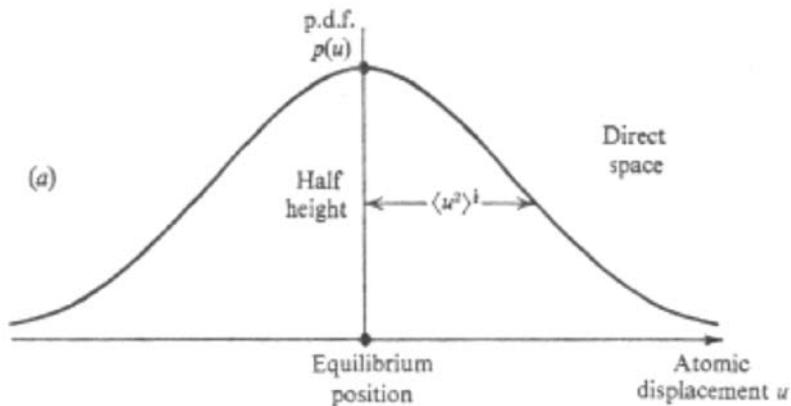
$$\text{F.T. } (\rho_\kappa(u)) = \text{FT}(\rho_{0,\kappa}(u)) \times \text{FT}(p_\kappa(u))$$

$$T(Q) = \int p_\kappa(u) \exp(i(Q \cdot u)) d^3u$$

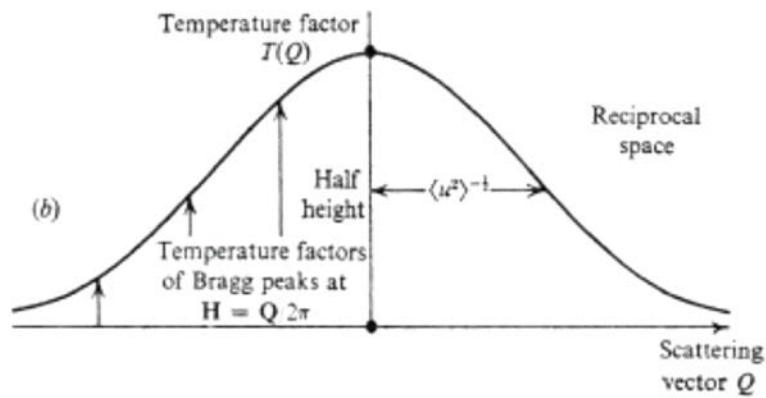
$$p_\kappa(u) = (2\pi^3)^{-1} \int T(Q) \exp(-i(Q \cdot u)) d^3Q$$

p.d.f. for a SHO is gaussian

$$p(u) = (2\pi\langle u^2 \rangle)^{\frac{1}{2}} \exp\left(-\frac{u^2}{2\langle u^2 \rangle}\right)$$



## Probability Density Function



**Temperature factor for an isotropically vibrating atom.**

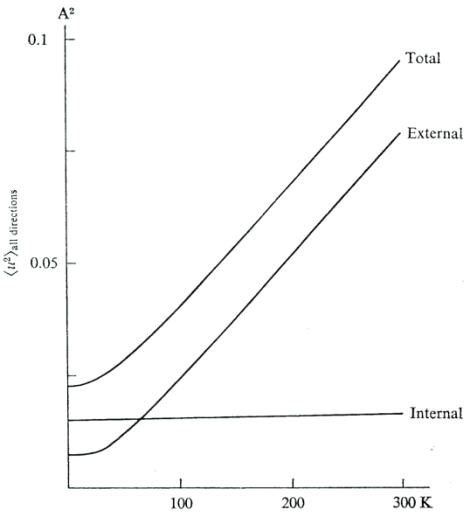
Both curves are Gaussians in the harmonic approximation

# Temperature Factors and Atomic Vibrations

$$T_{\kappa}^{anis}(\mathbf{H}) = \exp\left\{-\left[2\pi\mathbf{H}^T \mathbf{U} \mathbf{H}\right]\right\}$$

$$T_{\kappa}^{iso}(\mathbf{H}) = \exp\left\{-\left[2\pi\mathbf{H}^T \mathbf{H} \langle u(\kappa)^2 \rangle\right]\right\} = \exp\left(-\frac{4\pi^2 \sin^2 \vartheta}{\lambda^2} \langle u(\kappa)^2 \rangle\right) =$$

$$\exp\left(-B \frac{\sin^2 \vartheta}{\lambda^2}\right)$$



$$\langle u_i^2 \rangle = \frac{h}{8\pi^2 \mu \nu} \coth\left(\frac{h\nu}{2kT}\right)$$

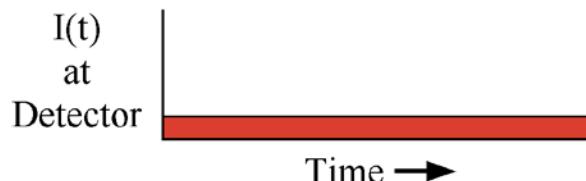
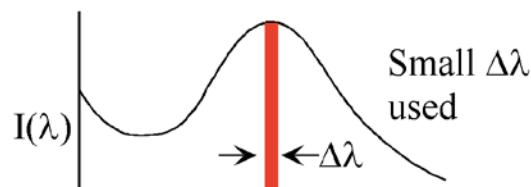
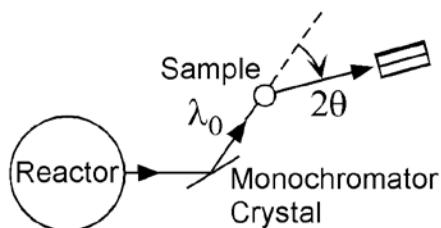
$$\mathbf{U} = \mathbf{t} + \mathbf{l} \wedge \mathbf{r}$$

# NEUTRON DIFFRACTION

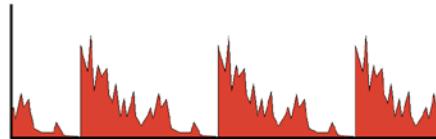
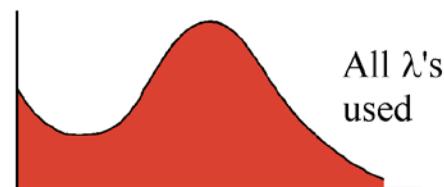
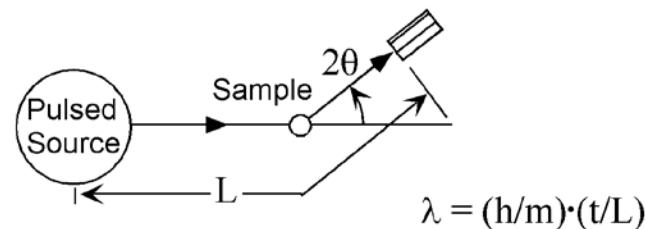
MEASURE  $F(d)$

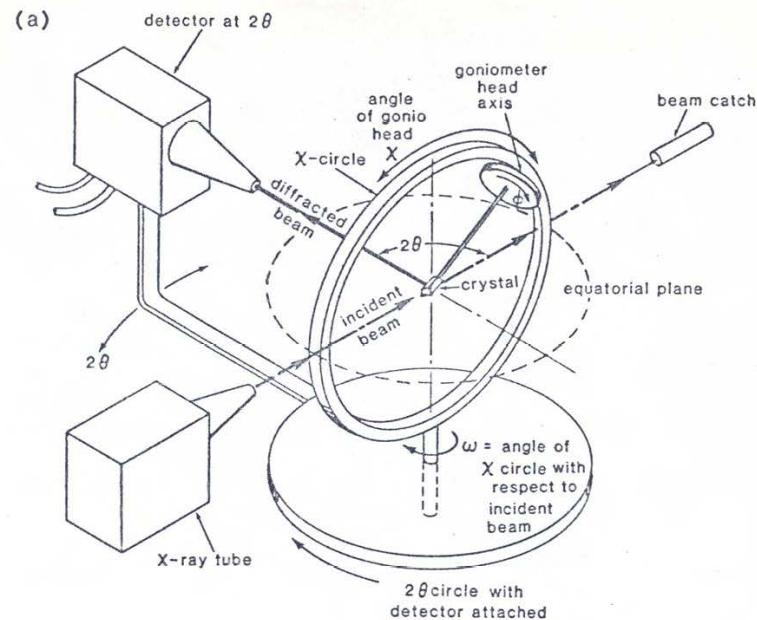
$$d = \frac{\lambda}{2\sin\theta}$$

STEADY STATE TECHNIQUE



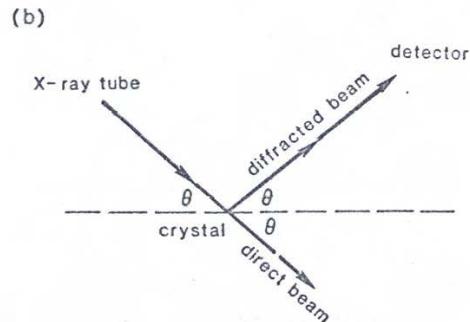
TIME OF FLIGHT TECHNIQUE



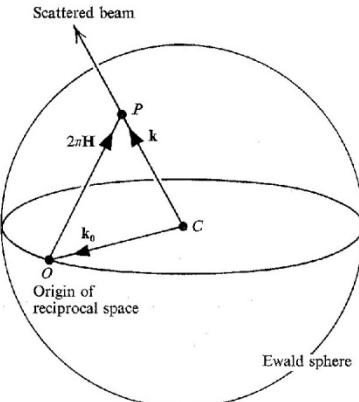


$\phi$  = spindle axis of goniometer head  
 $2\theta$  = angle between directions of incident and diffracted beams  
 $=$  angle detector has to be rotated to intercept diffracted beam

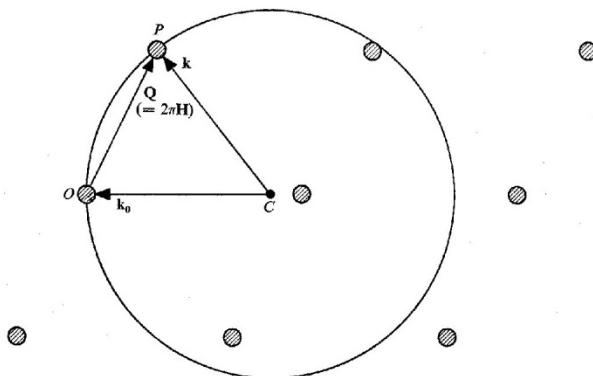
$\omega$  = angle between diffracted vector and plane of  $X$ -circle  
 $X$  = angle between  $\phi$  axis (gonio. head) and diffractometer axis (equatorial plane)



## The Ewald Sphere



Ewald sphere of reflection. Bragg scattering takes place when the reciprocal-lattice point  $2\pi\mathbf{H}$  lies on the sphere. The radius of the sphere is  $2\pi/\lambda$ .



# Disadvantages of neutron scattering

**There are not enough of them!**

**The fluxes are too low and the experiments take too long!**

**The samples needed are too big!**

**For powder diffraction at least 100 mg;**

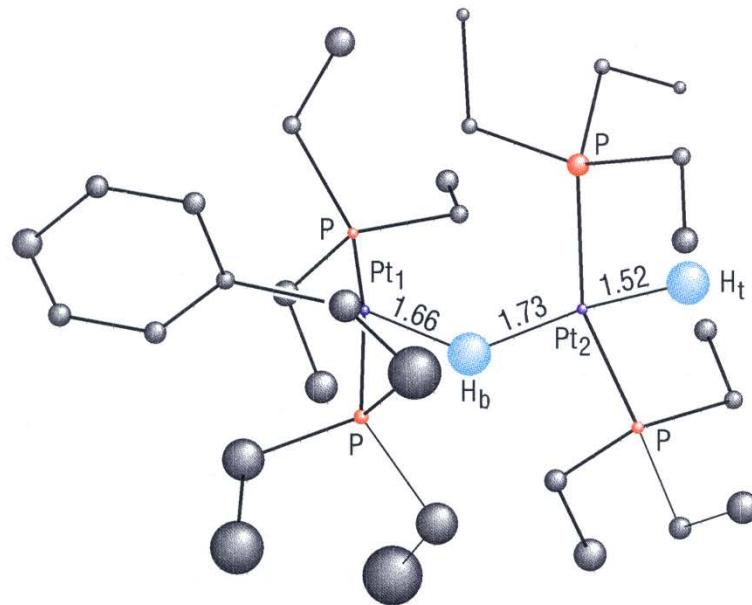
**For single crystal elastic a few mg**

**For single crystal inelastic 100 -1000 mg.**



BNL Single XXI Diffractometer on H6

$[H_2Pt_2Ph(PEt_3)_4]^+[BPh_4]^-$   
13K Neutron Diffraction Study

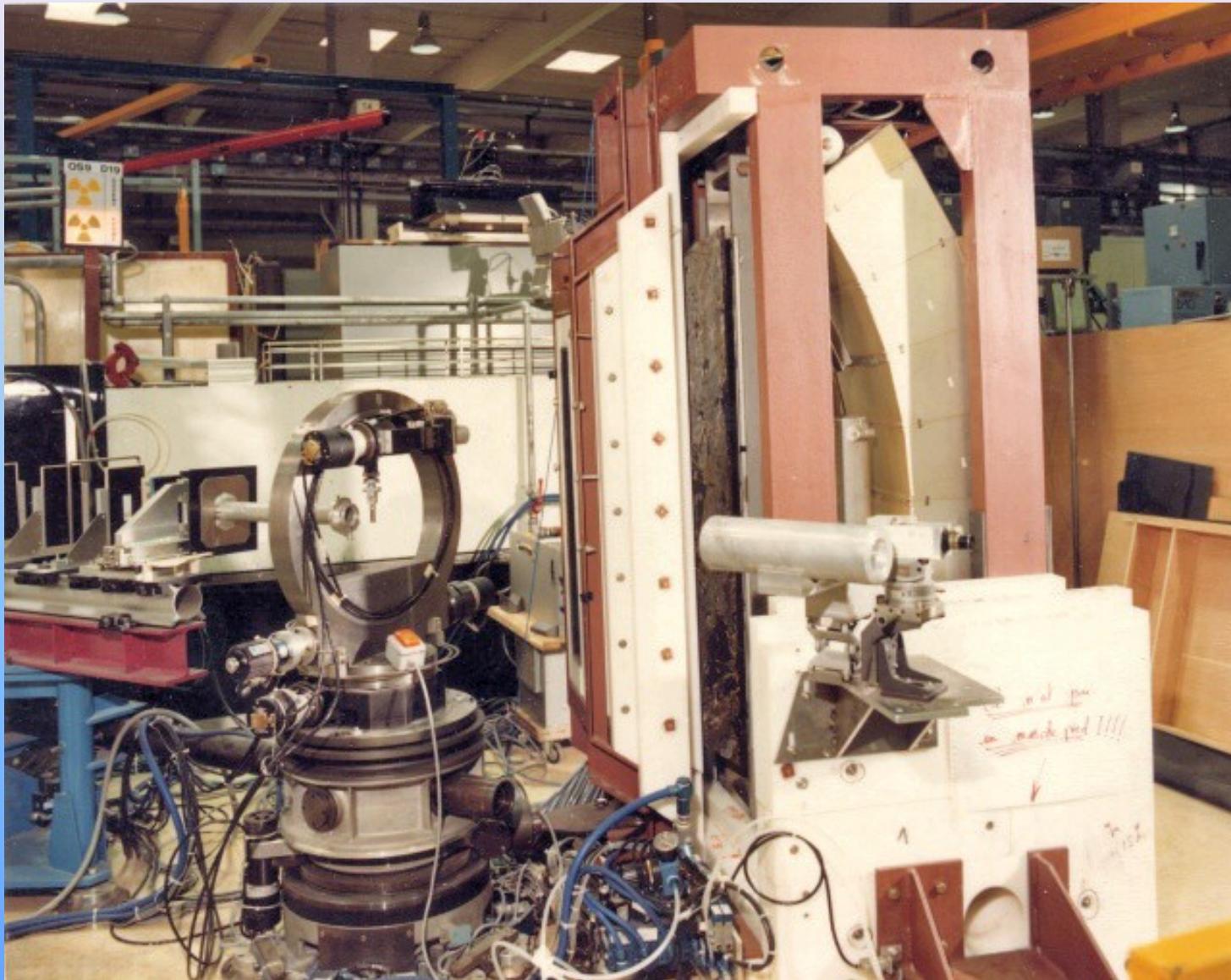


Pt<sub>1</sub>···Pt<sub>2</sub> 3.05(1) Å  
 $\angle$  Pt<sub>1</sub>-H<sub>b</sub>-Pt<sub>2</sub> 128(2)<sup>°</sup>

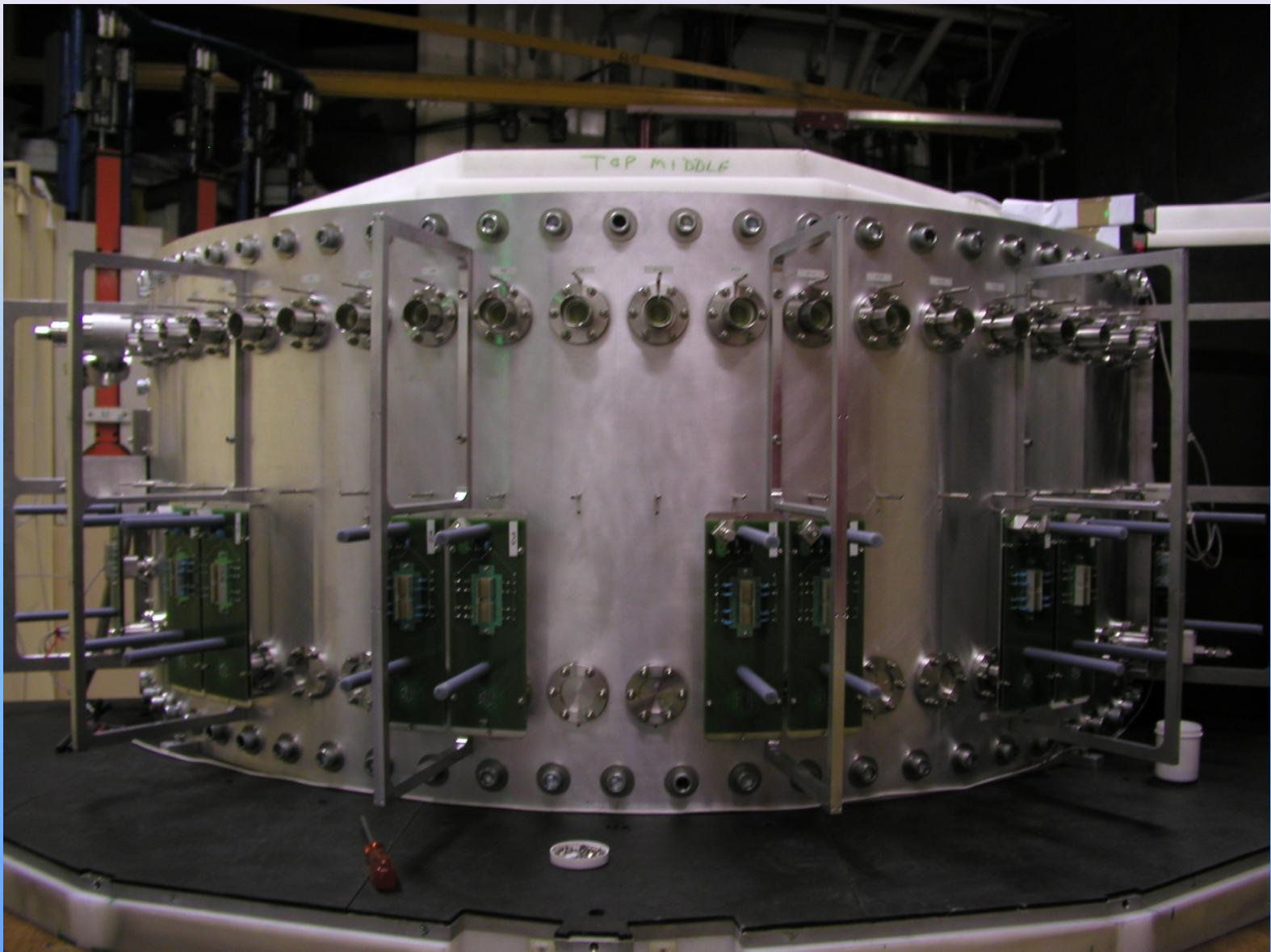
Ricci, Albinati and Koetzle (1995)

**ADVANCES IN COORDINATION  
CHEMISTRY ARE LINKED TO  
ADVANCES IN  
INSTRUMENTATION**

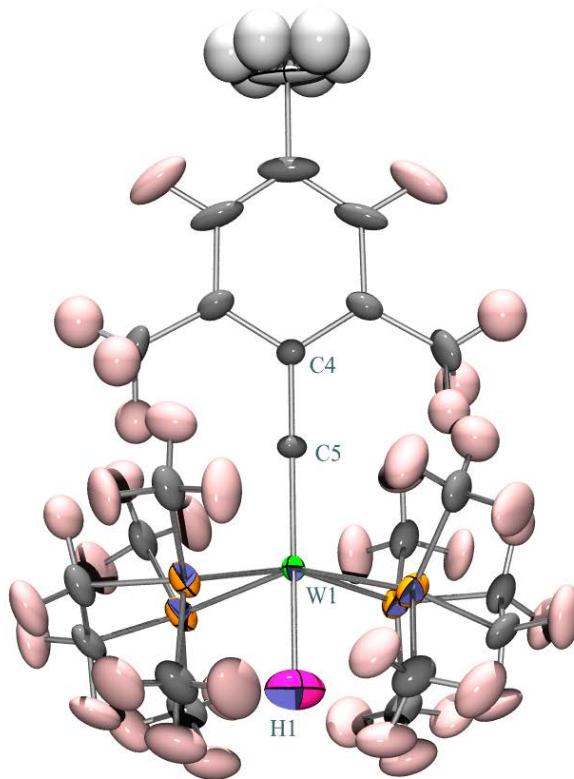
## D19 “Old Banana Detector”



## D19 New Detector 2007



*trans* - W(C-Mesityl)(dmpe)<sub>2</sub>H



W – H1 1.84 (2) Å

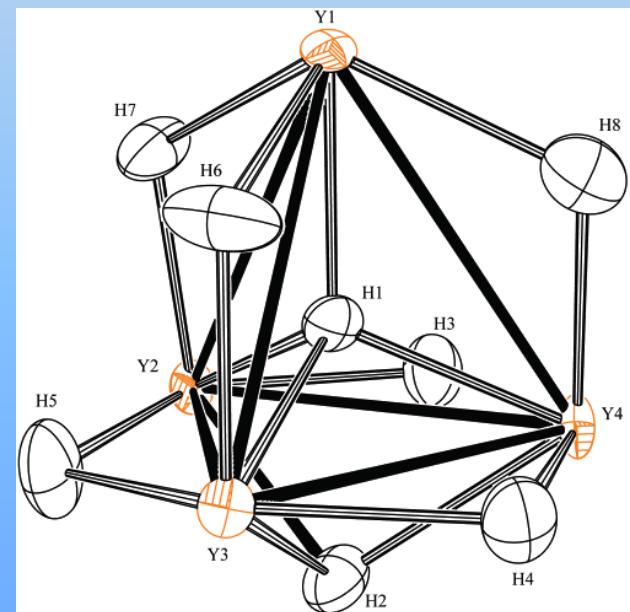
W – C5 1.868 (9)

C4 – C5 1.461 (8)

D19 @20K

V = 2667 Å<sup>3</sup>; C2/c

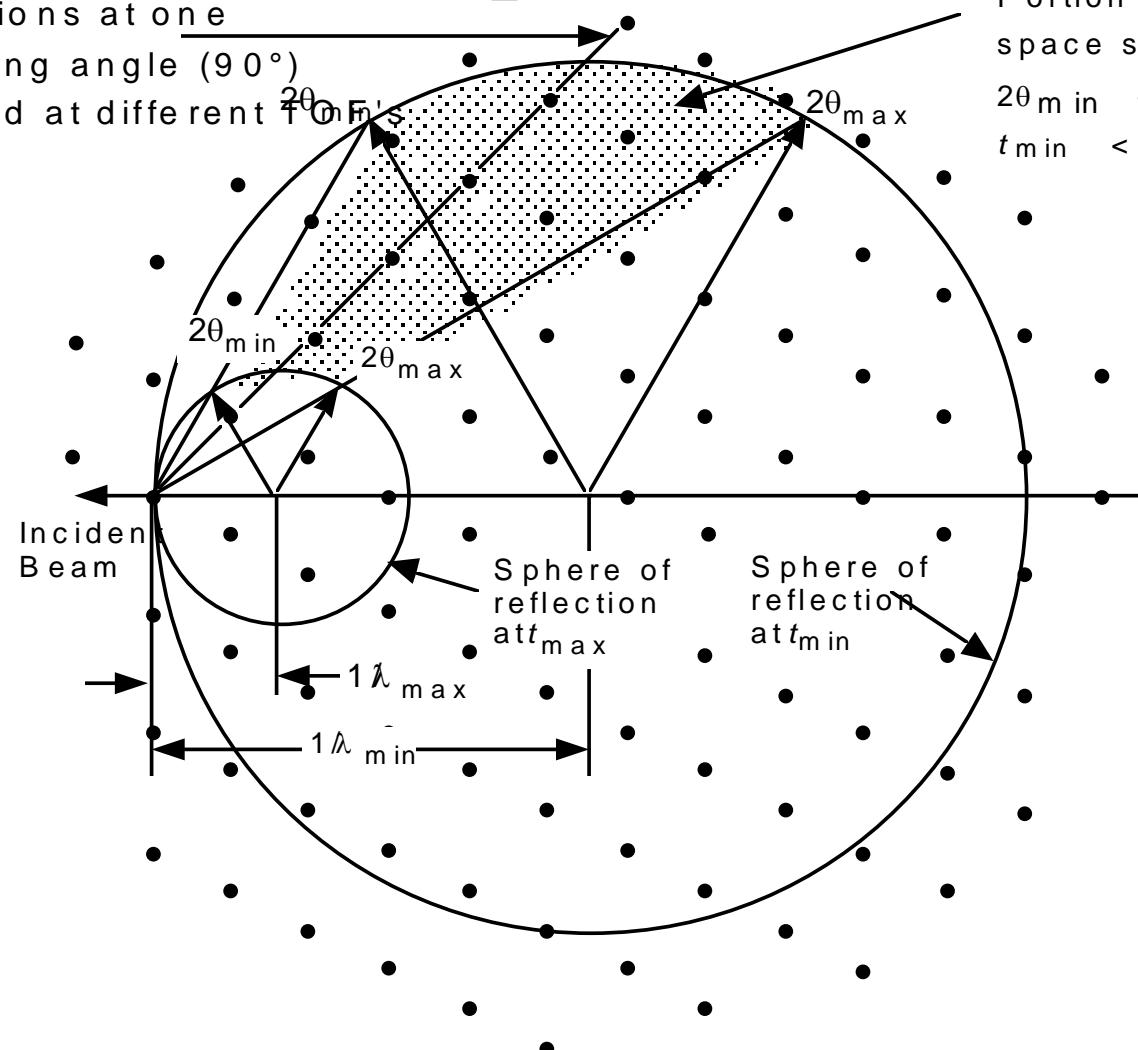
# VIVALDI @ ILL



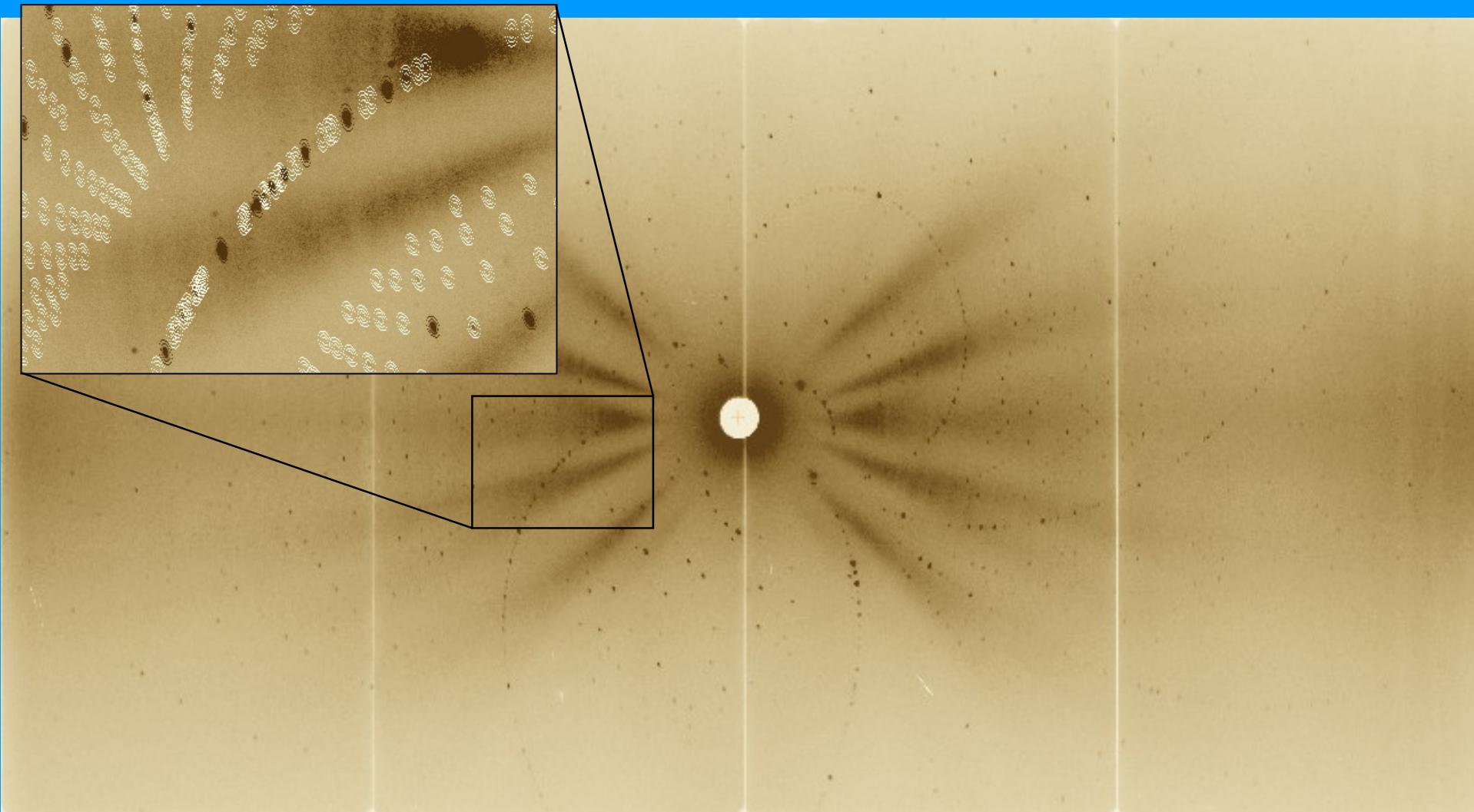
# Ewald Spheres

Reflections at one scattering angle ( $90^\circ$ ) resolved at different  $t_{OF}$ 's

Portion of reciprocal space sampled for  $2\theta_m \in \theta \leq \theta_{max}$  and  $t_{min} < t_{OF_{max}}$

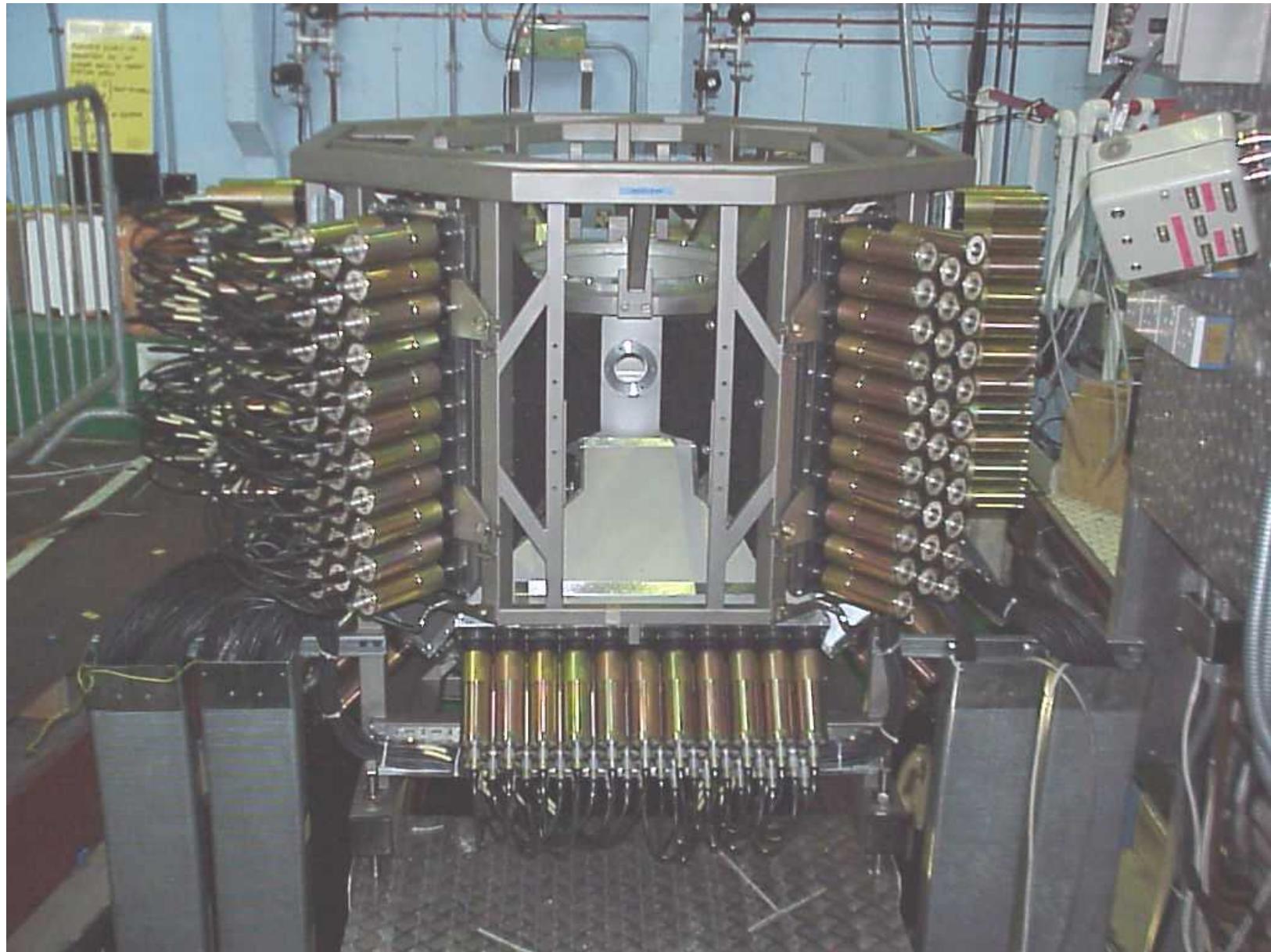


# VIVALDI

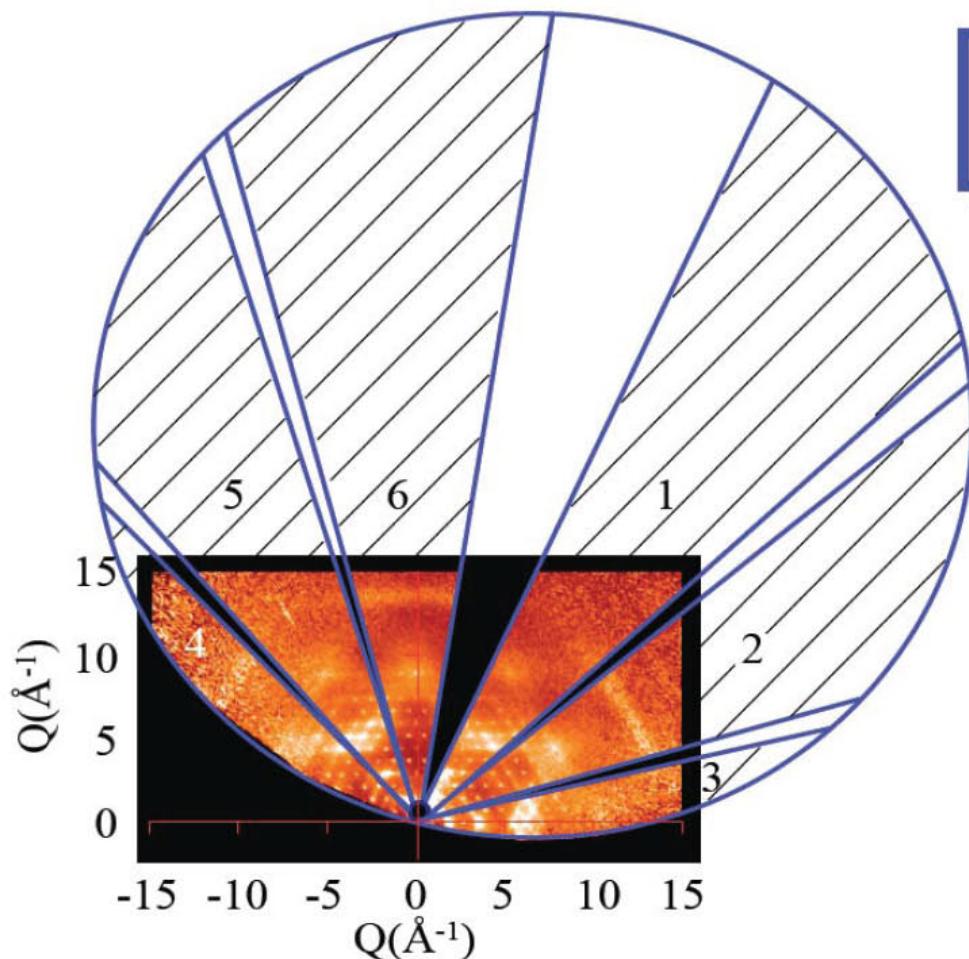


Typical neutron Laue pattern from  $(C_5Me_4(SiMe_3) - Y)_4H_{11}[C_5Me_4(SiMe_3)^*W]$ , Stewart, Bau *et al.*

# SXD - ISIS

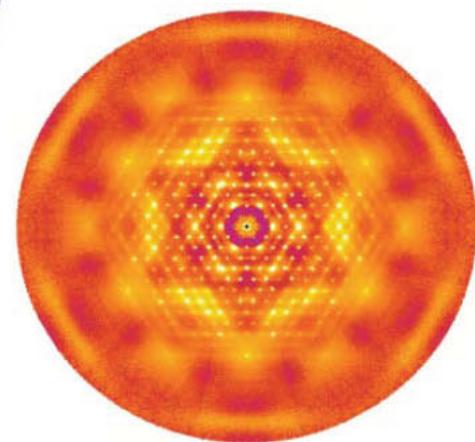


# SXD in practice

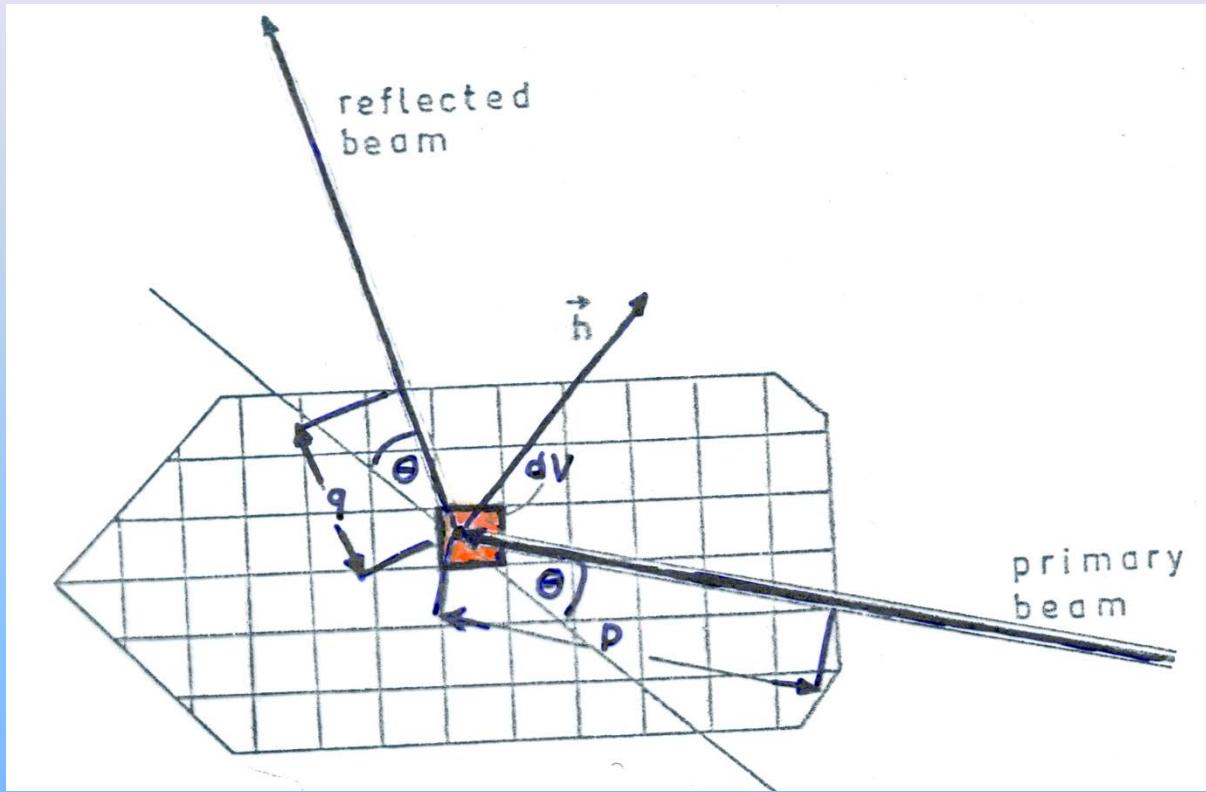


AN SXD DATA  
‘FRAME’

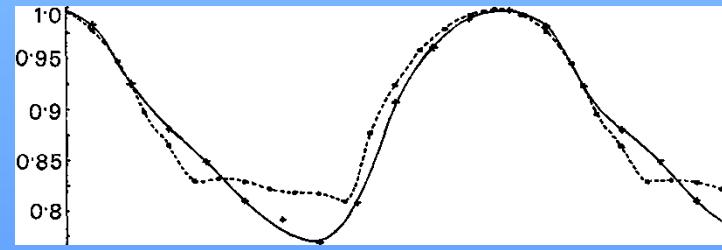
BENZIL on SXD  
001 vertical, one  
crystal setting, six  
equatorial detectors.



# Absorption

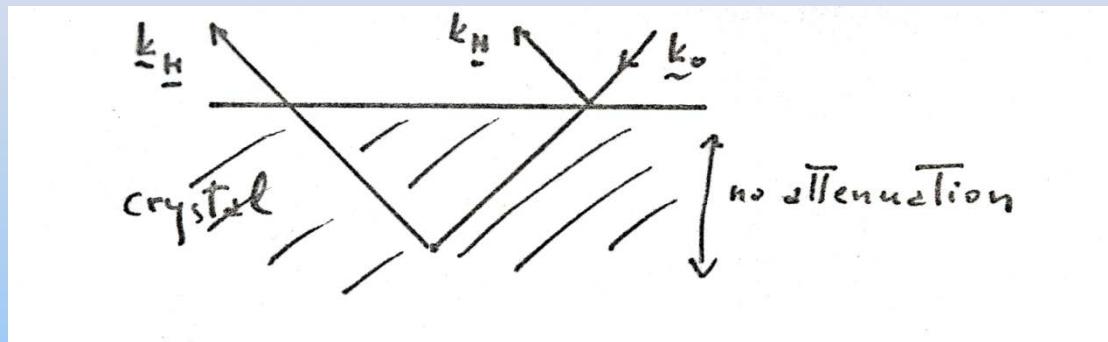


$$A = \frac{1}{V} \int_V e^{-\mu(p+q)} dV$$

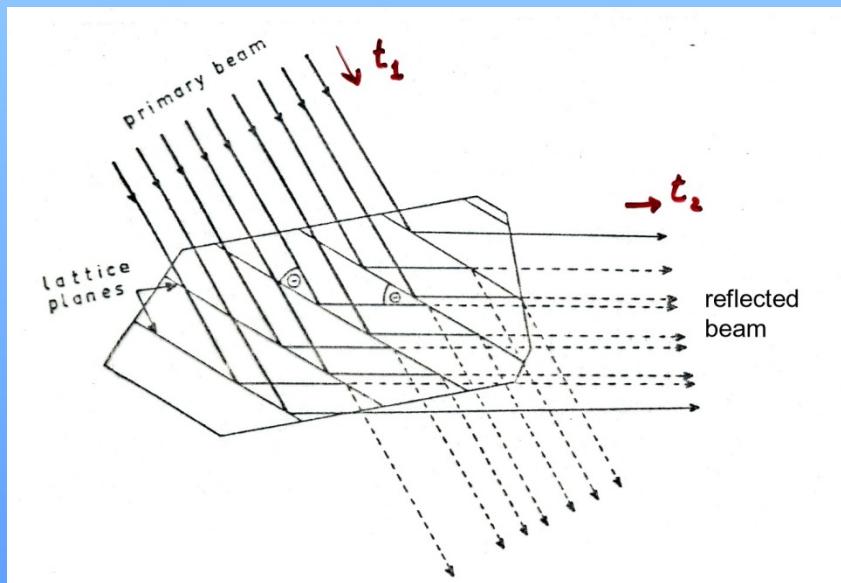


# Extinction I

- Reduction of Intensity by scattering not by absorption



- Primary Extinction: Weakening of Intensity by Multiple Reflections in the Crystal



# Extinction I

Zachariasen 1967

- 1) Spherical Crystal
- 2) Ideal Perfect Crystal

$$\frac{\partial I_0}{\partial t_1} = -\sigma I_0 + \sigma I$$

$\sigma \equiv$  diffracting power

$$\frac{\partial I}{\partial t_2} = -\sigma I + \sigma I_0$$

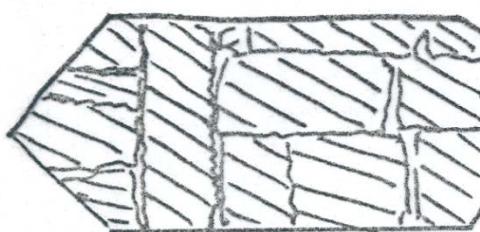
Intensity reduction  $\phi(\sigma) = \frac{1}{1 + \sigma t}$

# Extinction II

- Secondary Extinction: Weakening of the beam due to the **shielding** of the **inner planes** by the **outer planes**
- Most important for **strong reflections at low  $\sin\theta/\lambda$**

## Mosaic Crystal

Zachariasen 1967



Zachariasen 1967

### Type I

Depends mainly on  $g$   
(mosaic spread)

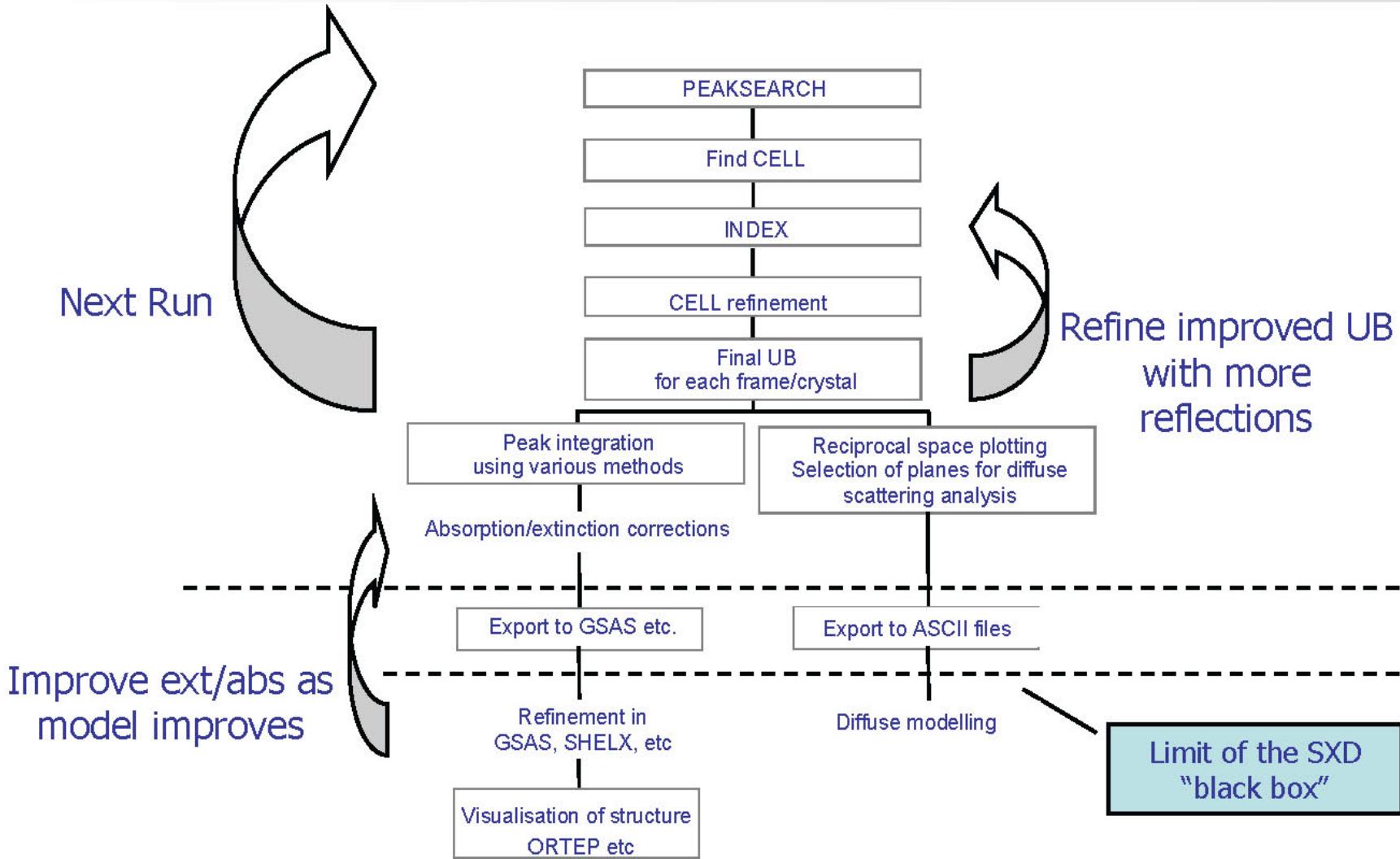
### Type II

Depends mainly on  $t$   
(radius of mosaic blocks)

$$I = I_{\text{obs}}(1 + 2gI_{\text{calc}})$$

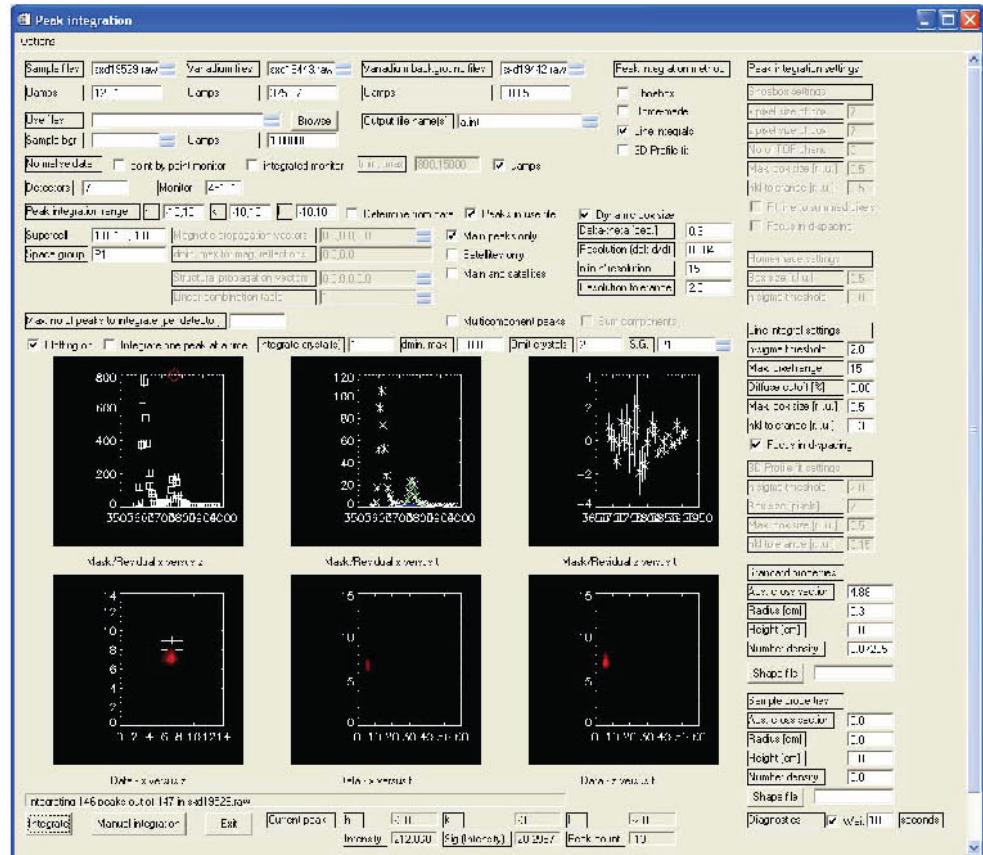
$$F_{\text{crted}} = F_o \{k[1 + 0.001|F_c|^2 \lambda^3 / \sin 2\theta]^{-1/4}\} \quad (\text{Shelxl})$$

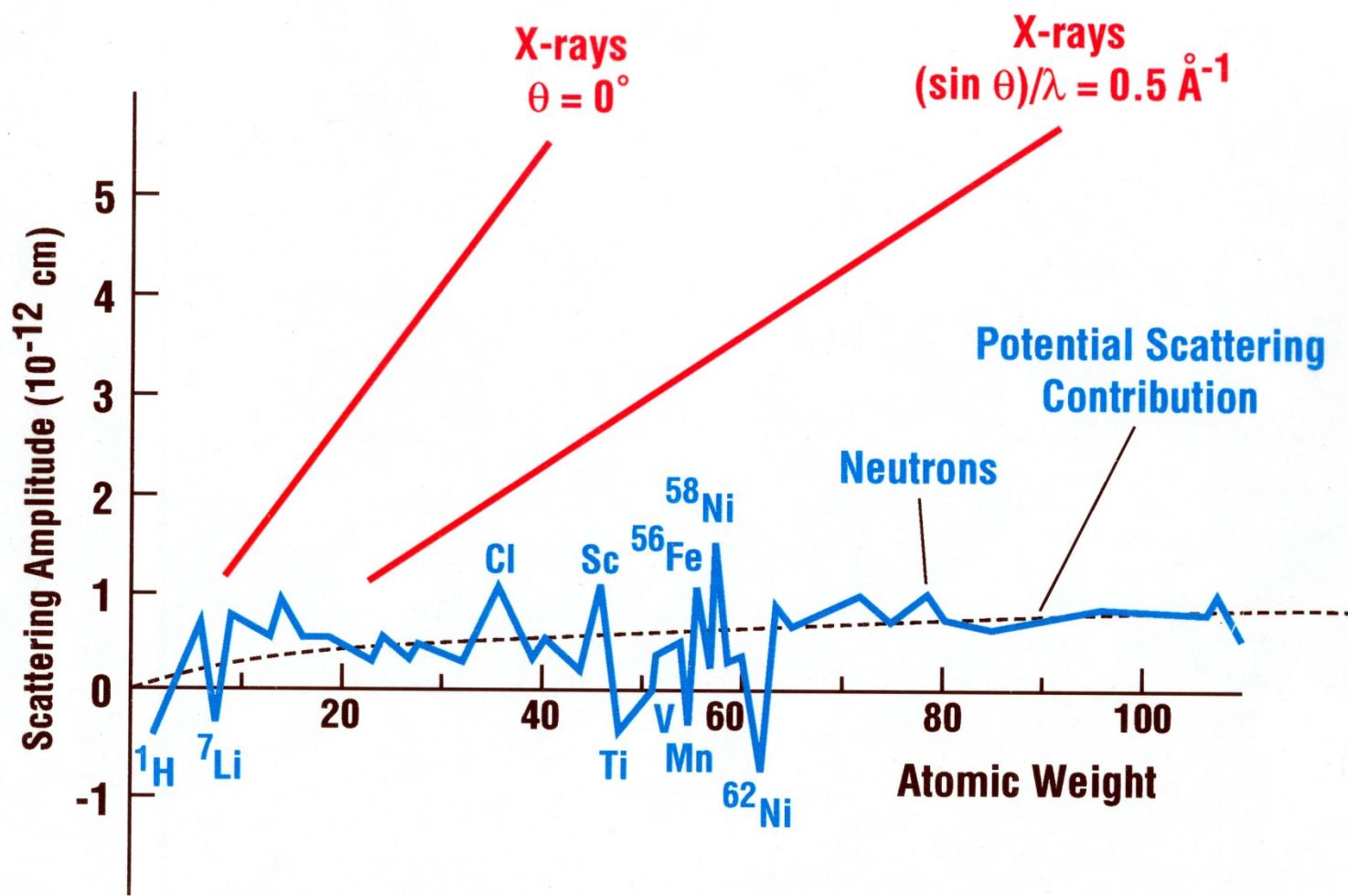
# SXD data processing

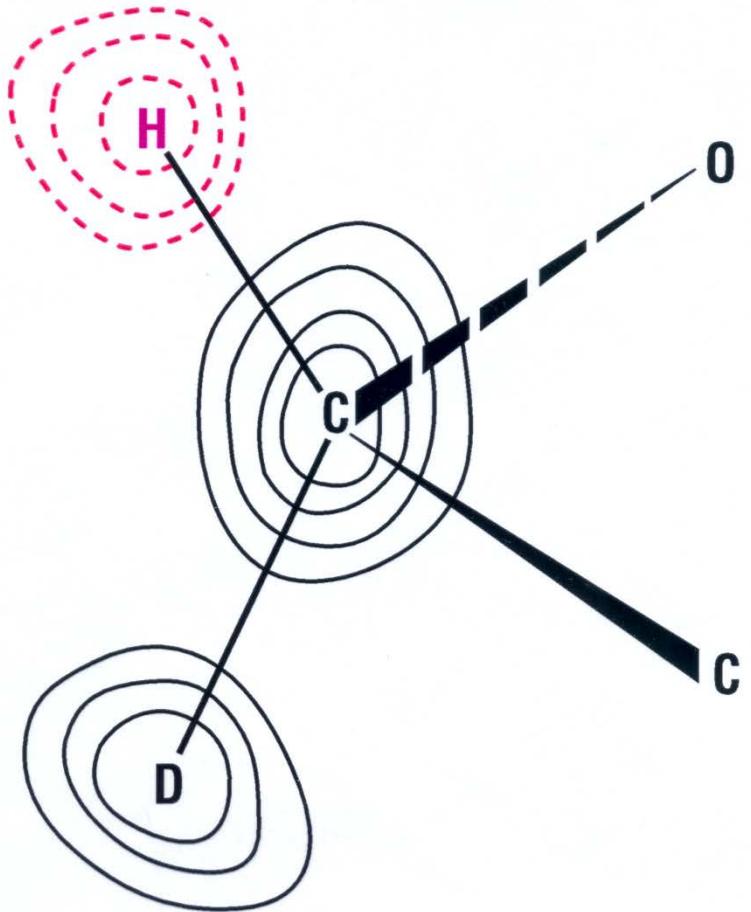


# Peak integration

- Choice of three algorithms
  - Shoebox
  - Dynamic box
  - 3D Gauss ellipsoid
- Manual integration
- Information about resolution used
- Gives directly  $F^2$
- Propagation vectors can be used

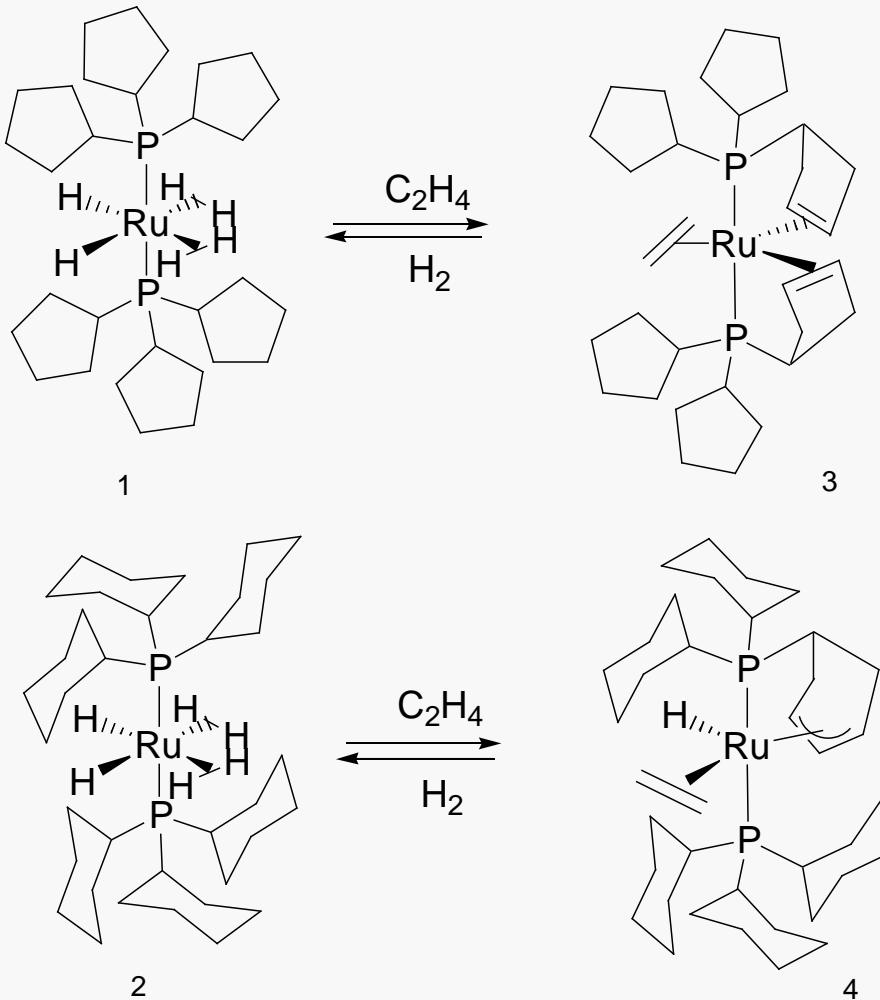




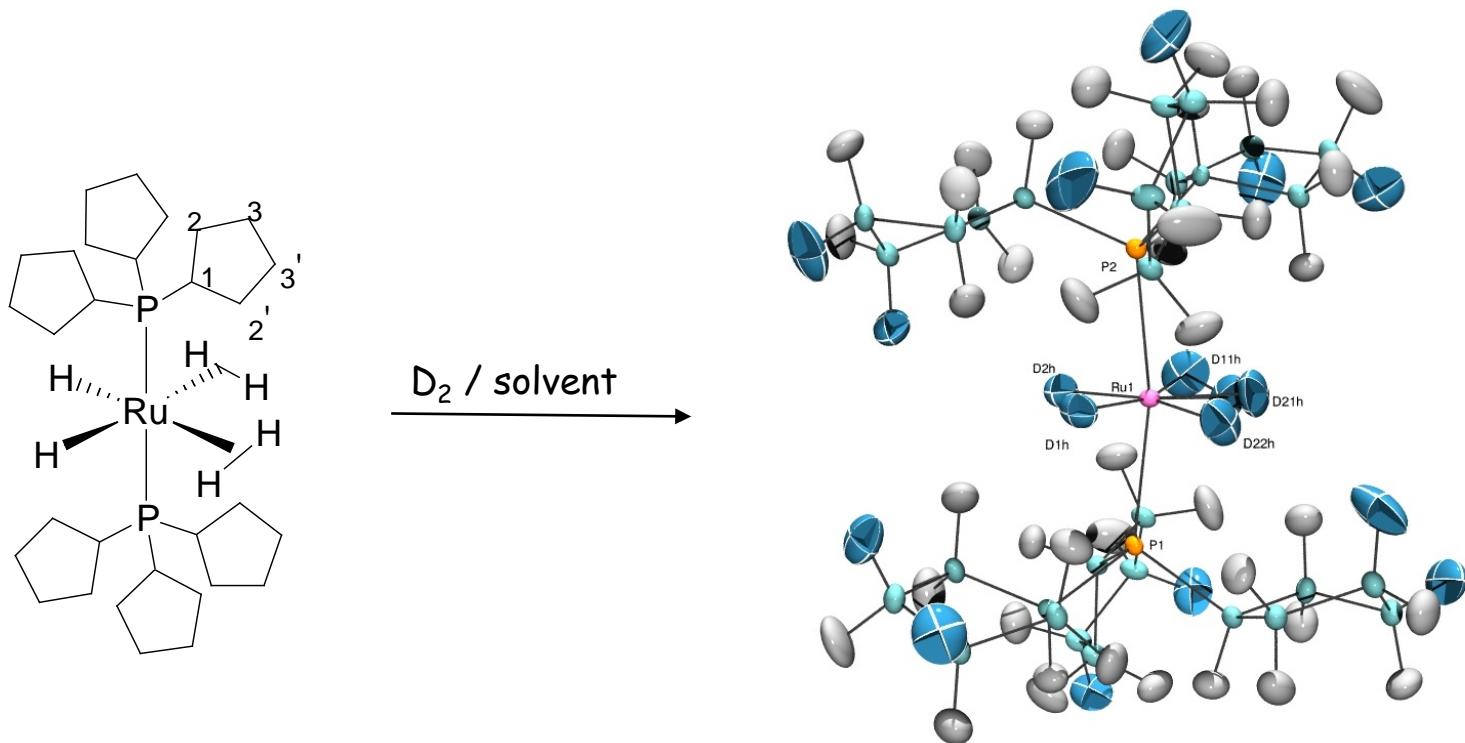


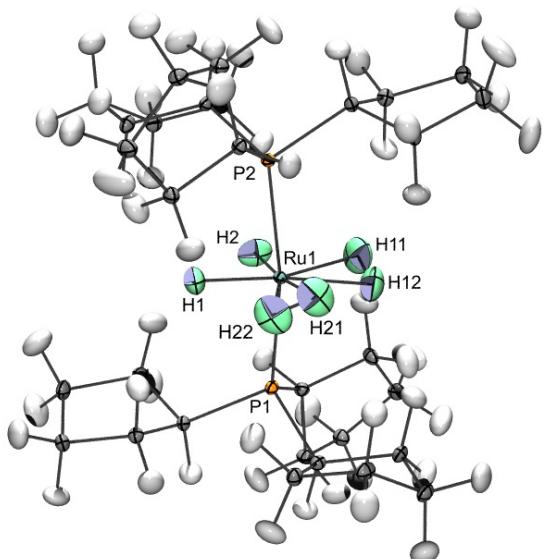
+ neopentyl-1-*d* alcohol (S)  
(yeast alcohol dehydrogenase)

# Reactivity of "Ru(PCyp)<sub>3</sub>" vs. "Ru(PCy)<sub>3</sub>"

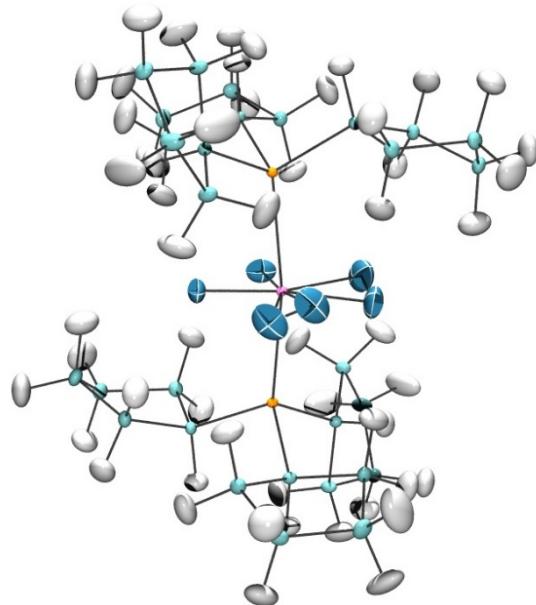


# Selective Deuteration of "RuH<sub>6</sub>"





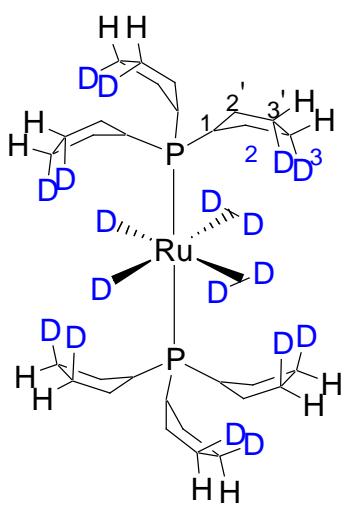
pentane  
3bar D<sub>2</sub>



M-D≈91%

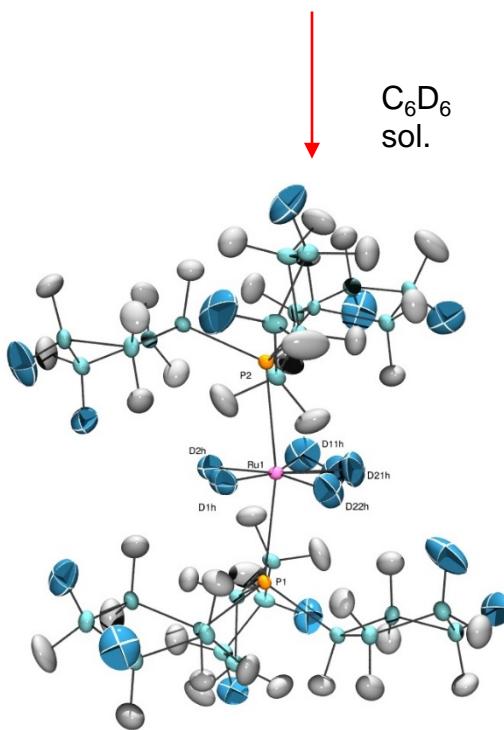
C<sub>6</sub>D<sub>6</sub>

Pentane/stirring  
3 d  
3bar D<sub>2</sub>  
3h

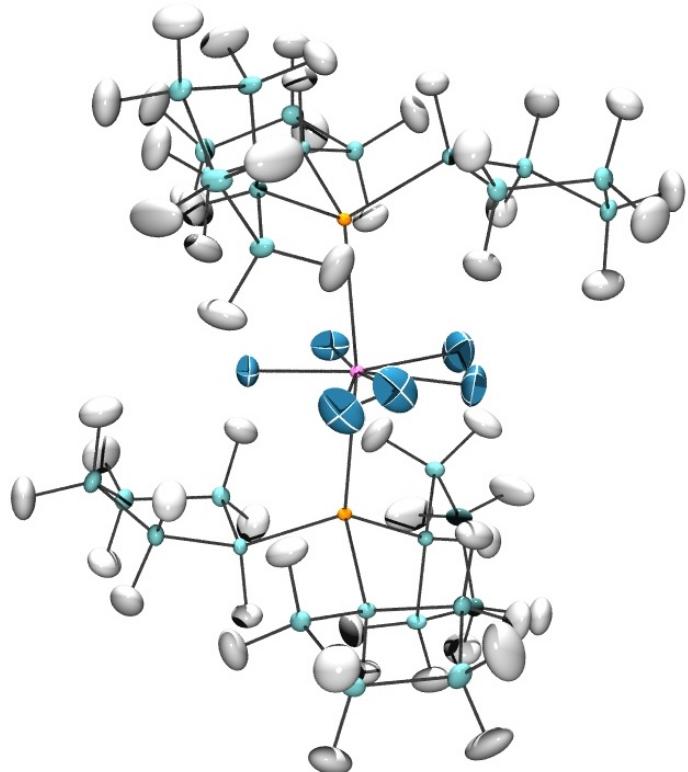


C<sub>6</sub>D<sub>6</sub>

M-D≈91%  
C-D≈32%



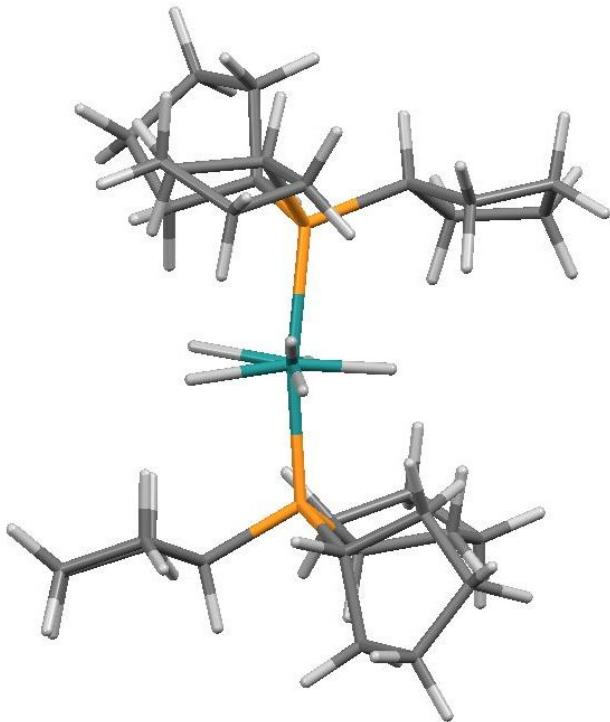
# " RuH<sub>2</sub>(H<sub>2</sub>)<sub>2</sub>((Cyp)<sub>3</sub>)<sub>2</sub>" vs. " RuD<sub>2</sub>(D<sub>2</sub>)<sub>2</sub>((Cyp)<sub>3</sub>)<sub>2</sub>"



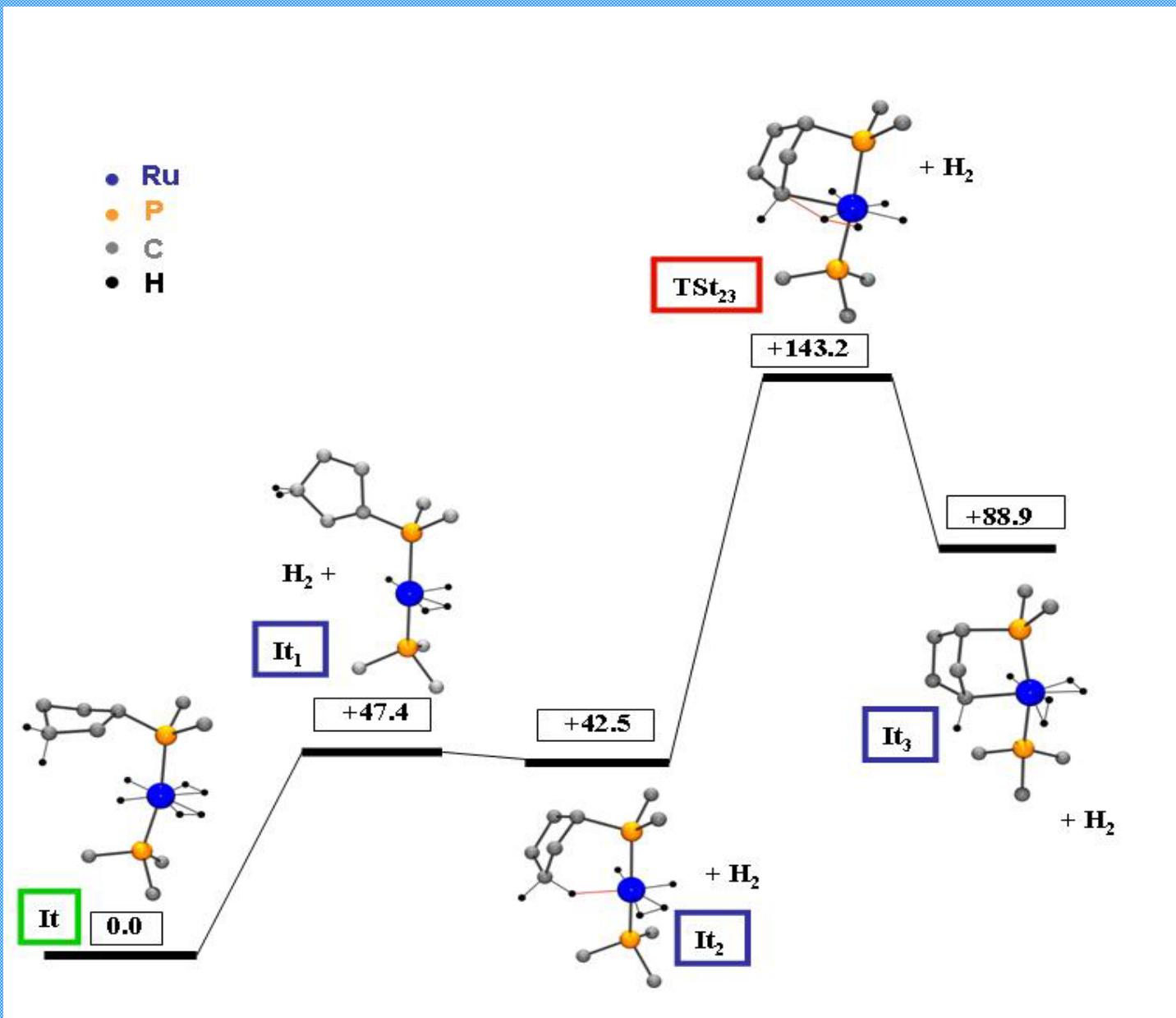
MD  
0.840(2) Å  
0.840(2) Å

MH  
0.825(7) Å  
0.835(7) Å

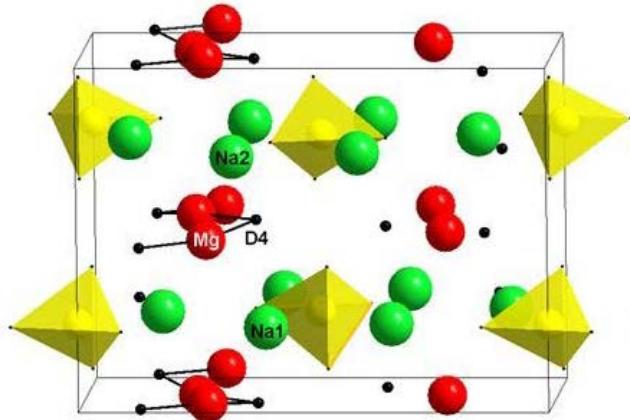
P-Ru-P    168.87(3)°    168.9(1)°



# C-H Activation Pathway in "RuH<sub>6</sub> – trans-P"

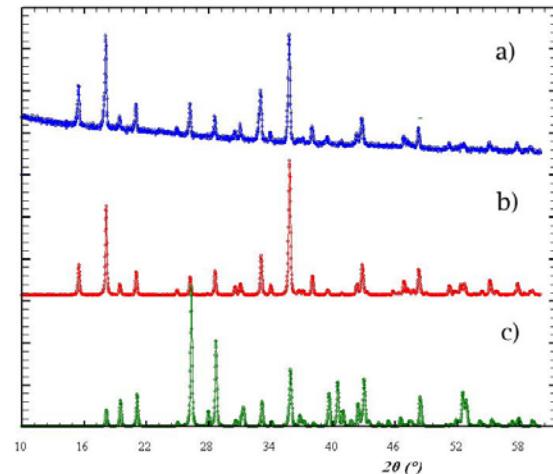
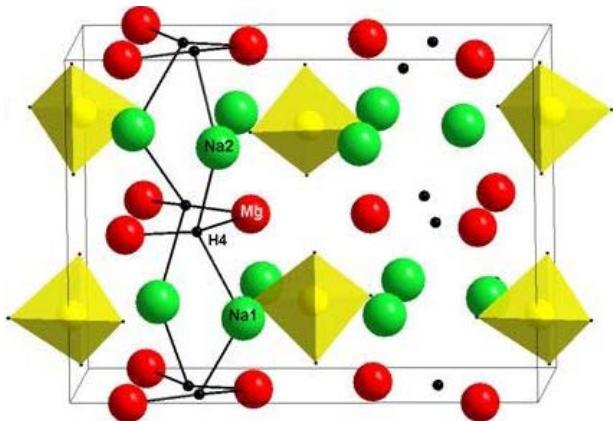


# The Structure of $\text{Na}_2\text{Mg}_2\text{NiH}_6$

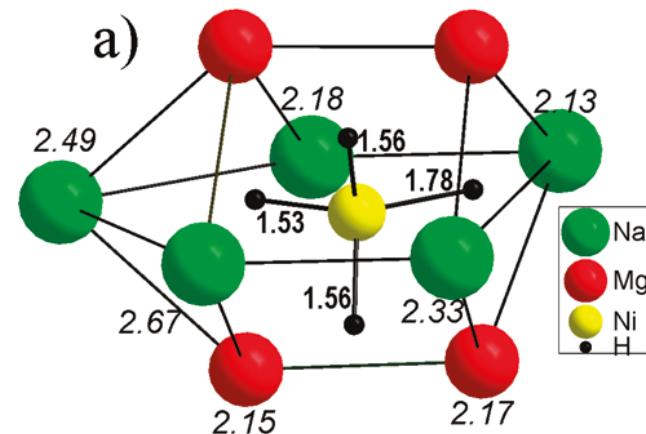


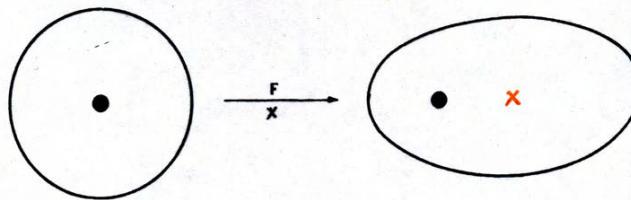
$\text{Na} - \text{Mg} = 2.66 \text{ \AA}$   
 $\text{Na} - \text{H} \sim 3.3 \text{ \AA}$   
 $\text{Mg} - \text{H} > 2.72 \text{ \AA}$

$$\begin{aligned} b_{\text{coh}}(\text{Mg}) &= 5.375(4) \\ b_{\text{coh}}(\text{D}) &= 6.674(6) \end{aligned}$$



a) exp. XRD pattern b) correct model c) Mg/D

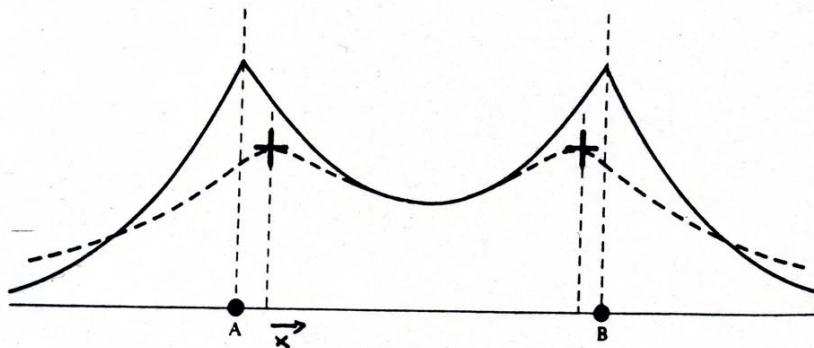




$\bar{x}$  = shift from the nucleus . . .  $F \equiv$  electric field

$$\bar{x} = \frac{q F}{e}$$

$$\bar{x} = 4.5 q^3 / R^2 \quad \left\{ \begin{array}{l} \bar{x} = 0.03 \text{ \AA} \\ R = 5.0 \text{ \AA} \end{array} \right.$$



Charge density along the axis for vibrating and non-vibrating  $\text{H}_2^+$ .  
The crosses denote 'apparent' positions of the nuclei.

for each electron:

$$\chi = (\phi_A + \phi_B) / \sqrt{2(1+s)}$$

$$\rho = \chi^2 = \phi_A^2 + \phi_B^2 - \frac{s}{1+s} (\phi_A^2 + \phi_B^2) + \frac{2}{1+s} \phi_A \phi_B$$

$$\bar{x} = \int x \rho d\tau$$

## **X-ray vs. Neutron Diffraction**

### **Systematic Differences in Observed Bond Separations**

**C – H**                    **– 0.096 (7) Å**

**O – H**                    **– 0.155 (10) Å**

**C = C**                    **– 0.005 (1) Å**

**C = C**                    **– 0.008 (2) Å**    *(benzenoid)*

**C – OH**                    **0.005 (1) Å**

**C – O**                    **0.008 (2) Å**

# Electron Density in Crystals: X-Ray Scattering

$$F_{calc}(\mathbf{H}) = \sum_j f_j \exp(2\pi i \mathbf{H} \cdot \mathbf{r}_j) T_j(\mathbf{H})$$

where  $f_j$ 's are calculated **for spherical atoms**

## In real molecules:

- Atoms participate in charge transfers
- Atoms deviate from spherical symmetries (lone pairs, incomplete shells)
- Atoms form covalent bonds sharing electron densities

X-rays see the **Centroid of Charges** displaced in the direction of the bond

**Positional and thermal parameters** from *X-ray diffraction may be biased*

**Positional and thermal parameters** from *neutron diffraction (nuclear scattering)*

**are not biased**

# Electron Density in Crystals: X - N Maps

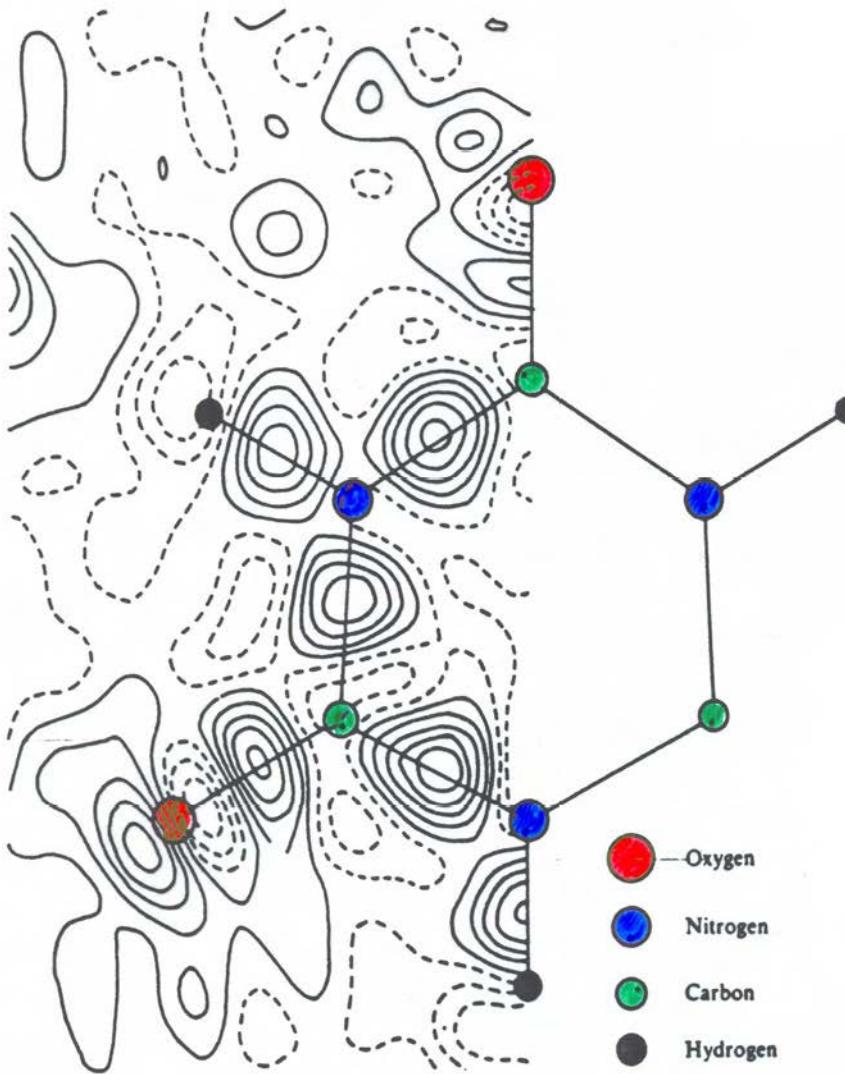
$$F_{calc,\text{N}}(\mathbf{H}) = \sum_{atoms} f_{j,\text{X}} \exp(2\pi i \mathbf{H} \cdot \mathbf{r}_{j,\text{N}}) T_{j,\text{N}}$$

$\mathbf{r}_{j,\text{N}} \equiv \mathbf{r}(x_j, y_j, z_j)$  neutron values - **unbiased**

$F_{calc,\text{N}}$  *ideal molecule* built from spherical atoms and neutron unbiased parameters → **correct geometry and no electron redistribution**

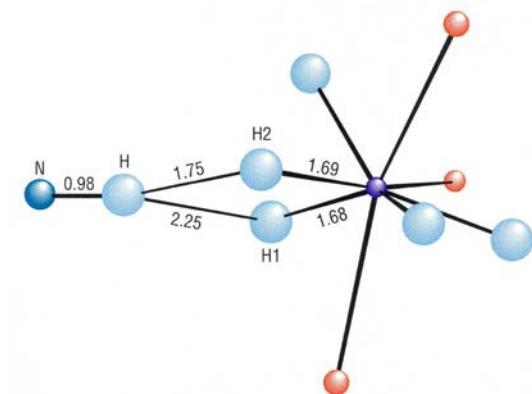
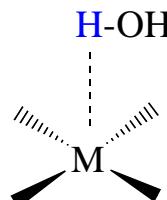
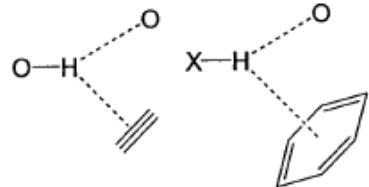
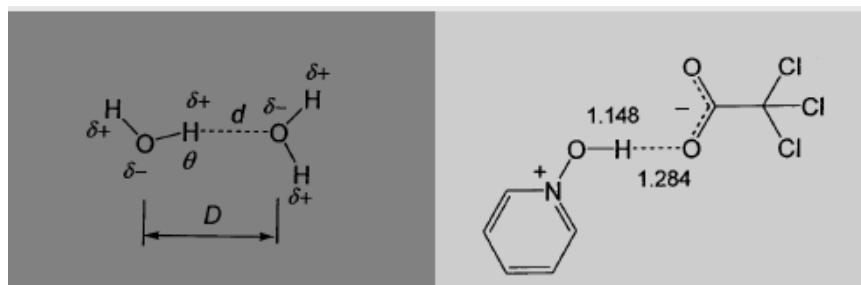
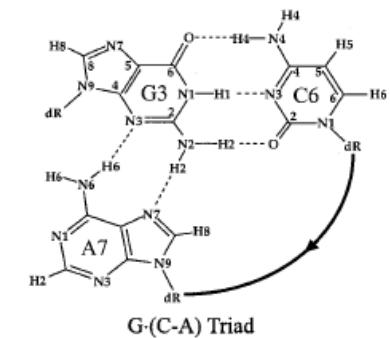
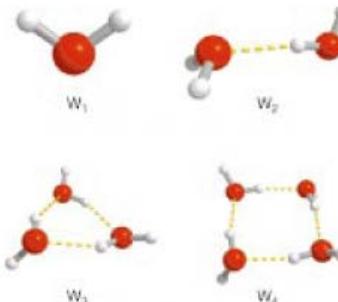
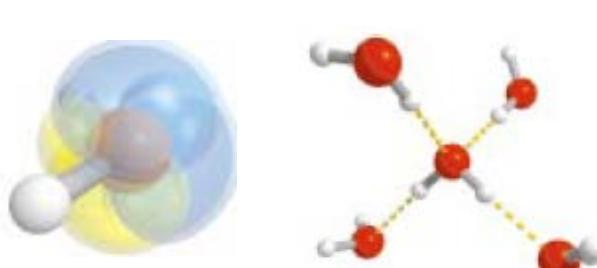
$F_{obs,\text{X}}$  obtained from X-ray diffraction are the F.T. of the **actual electron distribution**

$$\rho_{deformation}^{\text{X-N}}(\mathbf{r}) = \frac{1}{V} \sum_{\mathbf{H}} (F_{obs,\text{X}} - F_{calc,\text{N}}) \exp(-2\pi i \mathbf{H} \cdot \mathbf{r})$$

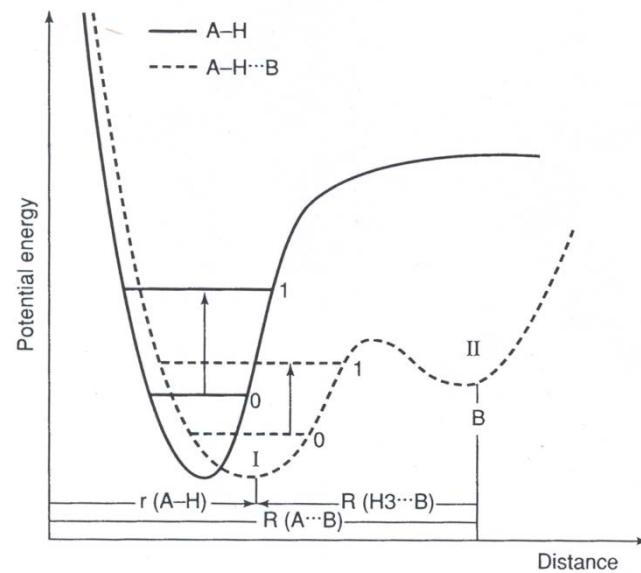
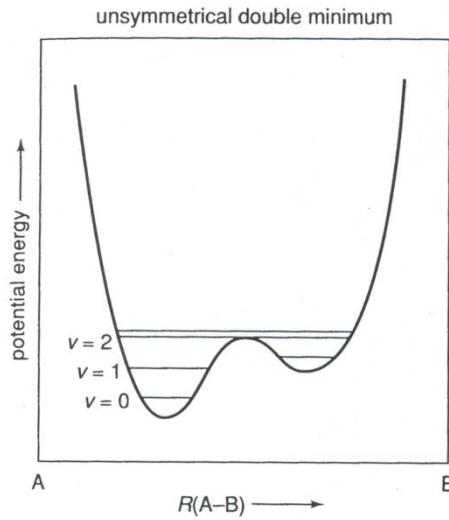
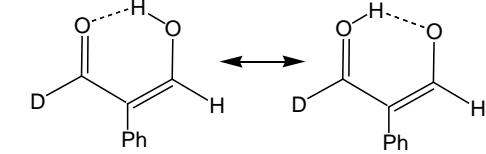
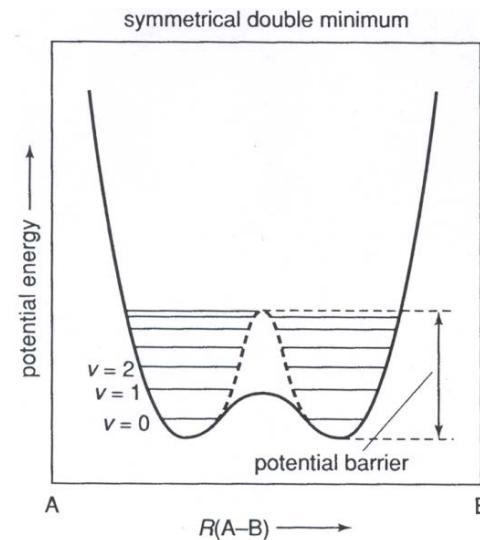
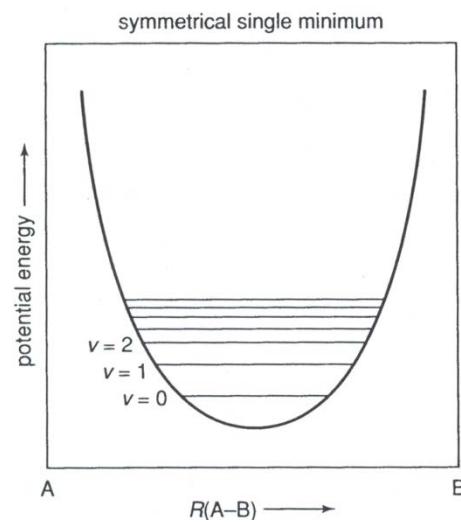
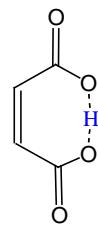


X-N map through the cyanuric acid molecule. Contours at  $0.10 e$   
 $\text{\AA}^{-3}$ . Zero and negative contours dotted.

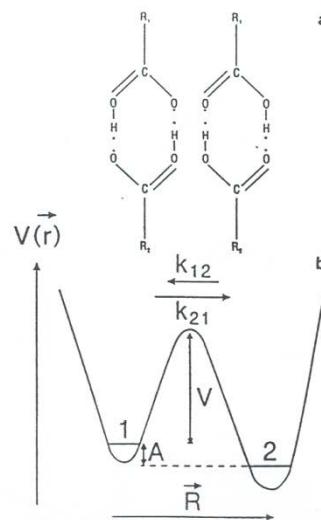
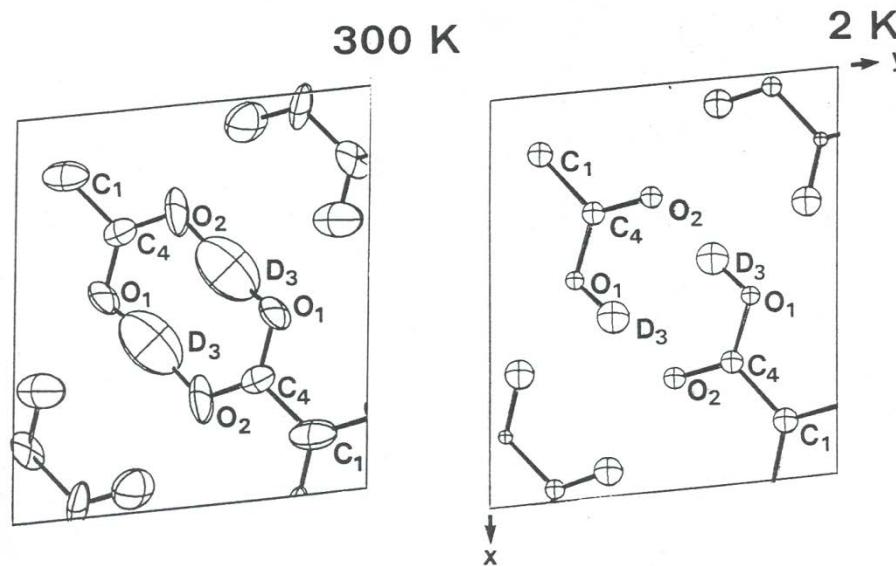
# A Few Types of Hydrogen-Bonds



# Potential Energy Landscapes for Hydrogen Bonds



# Proton Transfer in Terephthalic Acid



# $\text{KH}_2\text{PO}_4$ (KDP) – Ferroelectric Phase Transition

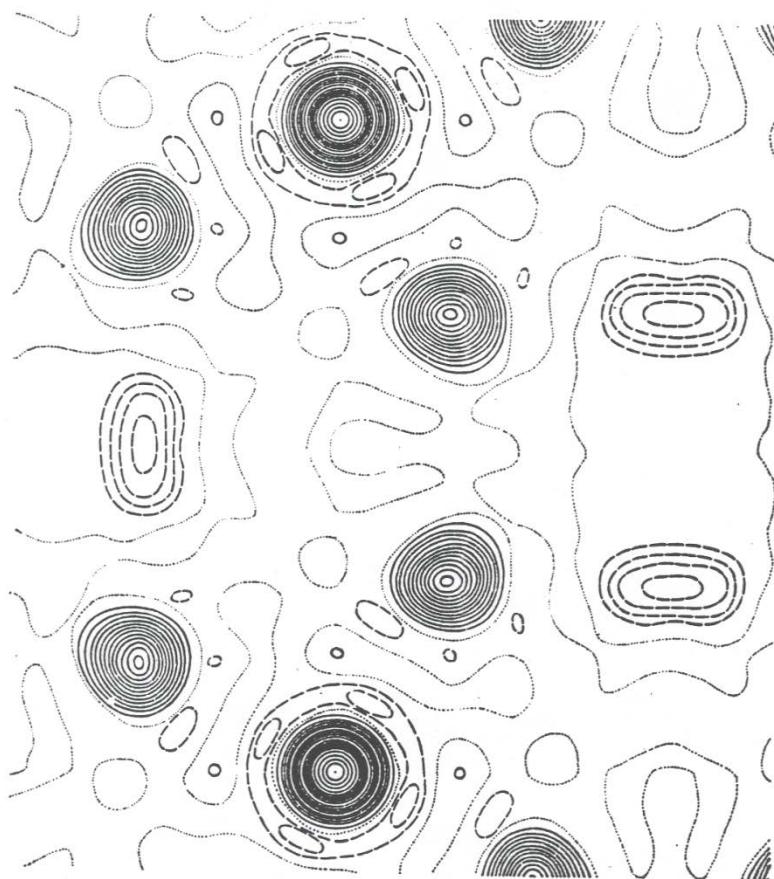
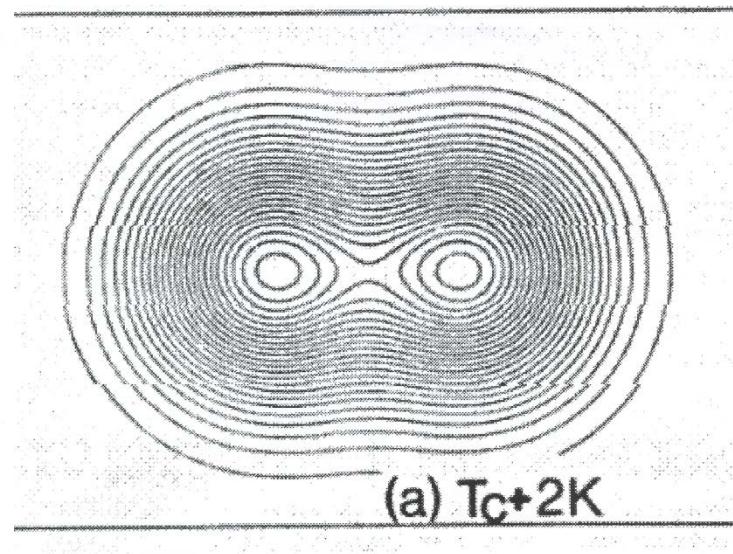
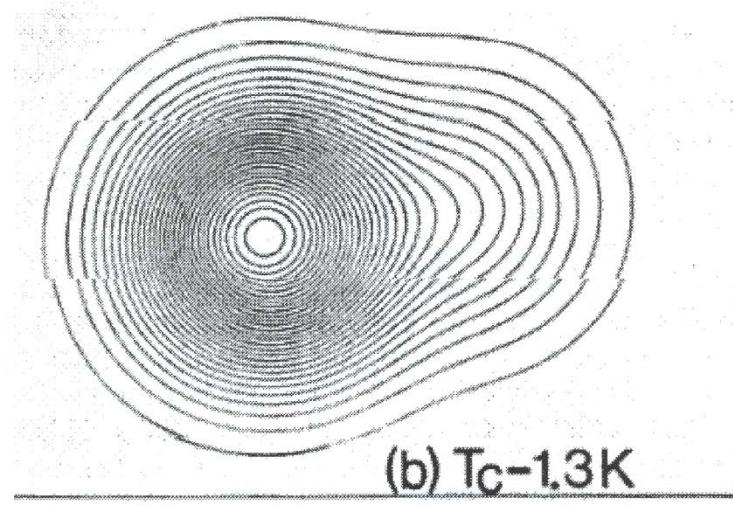


FIG. 107. A Fourier projection of the scattering density on the (001) plane of  $\text{KH}_2\text{PO}_4$  at room temperature. Contours are at intervals of 50 units, with additional contours at  $-75$ ,  $-125$ . Full lines are positive, broken lines are negative, and dotted lines are zero contours. The most intense peaks are superimposed K, P. The other positive peaks are O: the negative peaks H. (Bacon and Pease, 1953.)

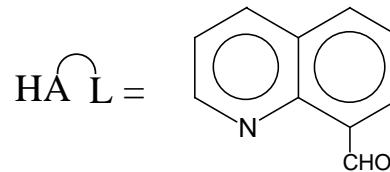
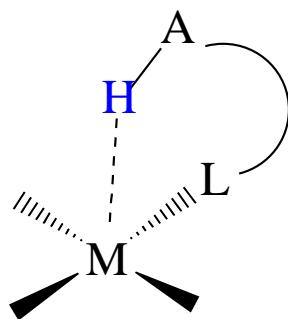
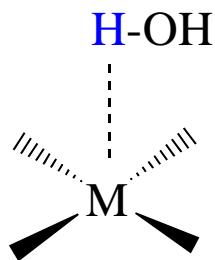
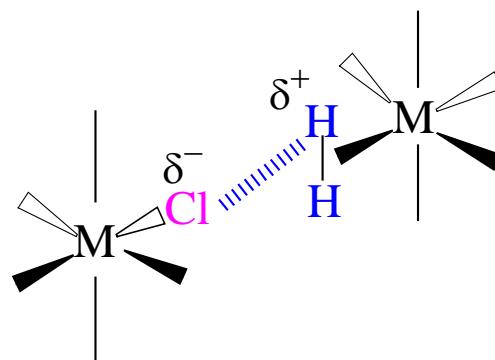


(a)  $T_c + 2\text{ K}$

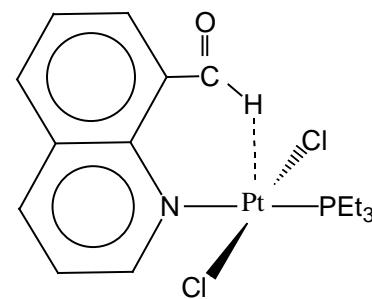


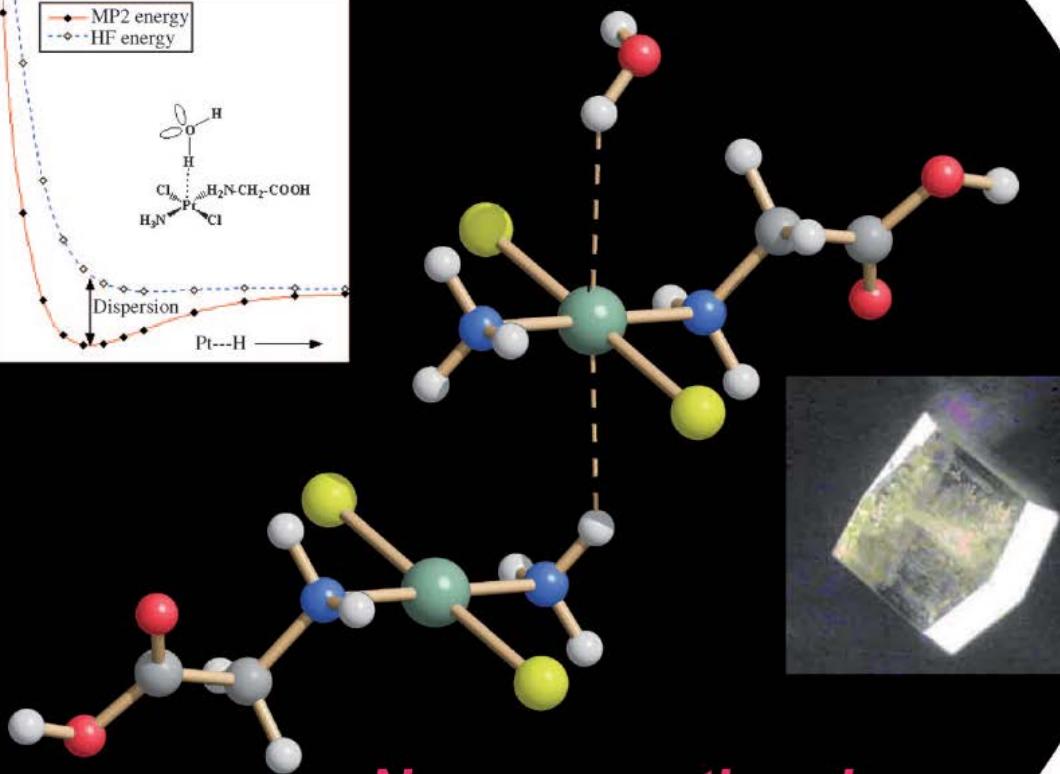
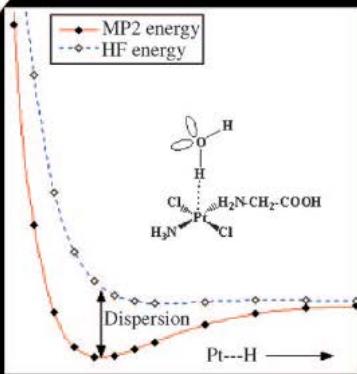
(b)  $T_c - 1.3\text{ K}$

# Non-Classical "Weak" Interactions



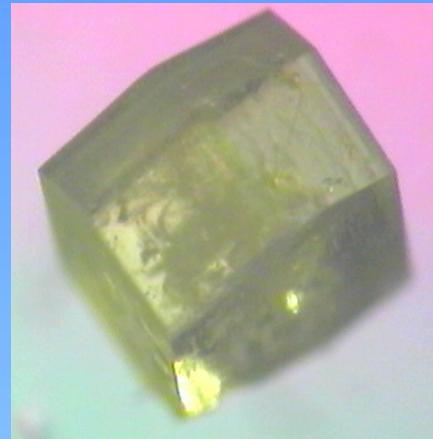
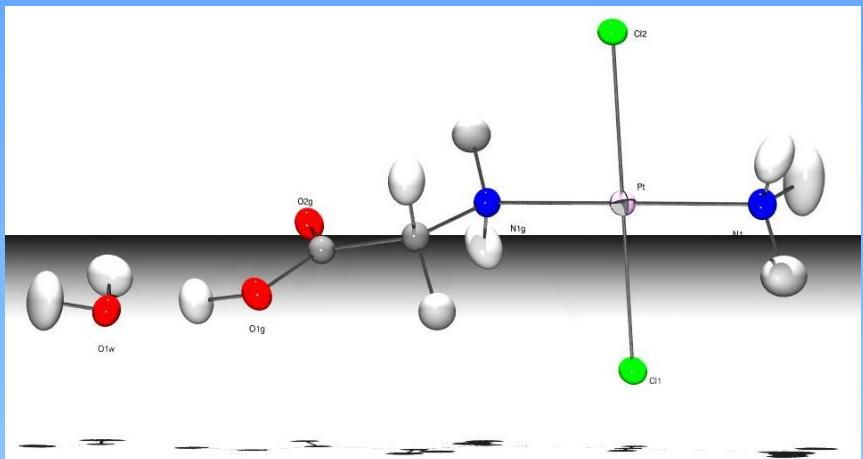
$d_{(M \dots H)} \approx 2.5 - 2.9 \text{ \AA}$



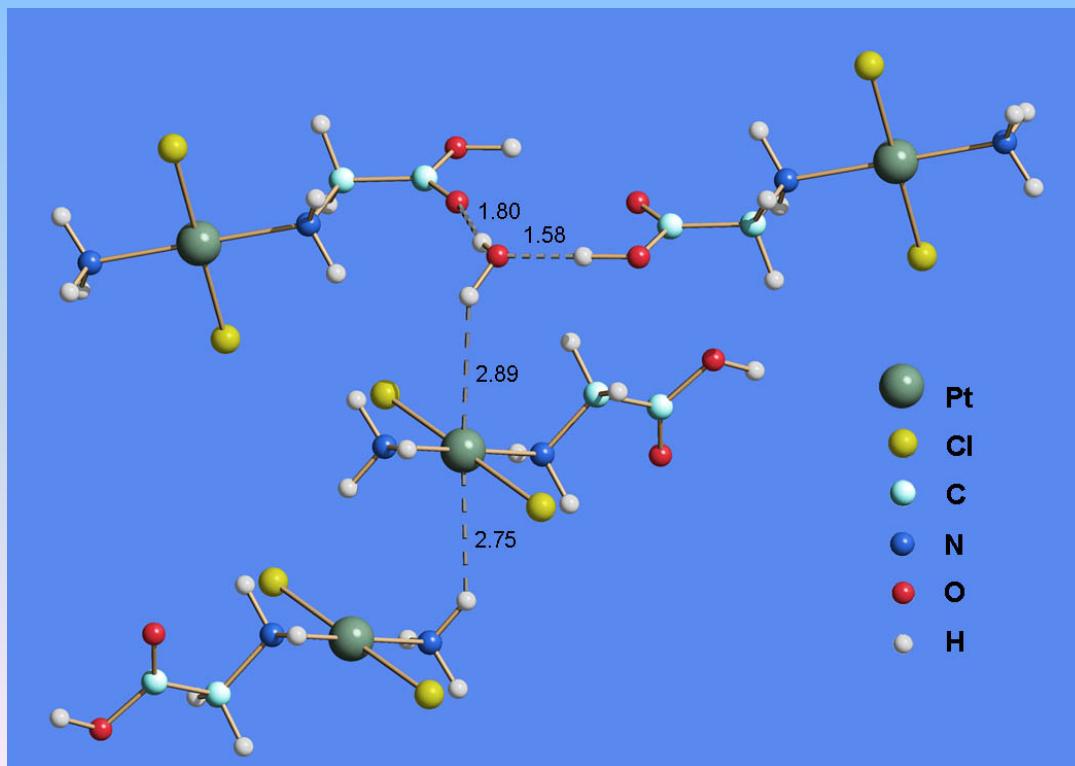


## *Nonconventional hydrogen bonding*

# $\text{Pt}(\text{NH}_3)(\text{gly})\text{Cl}_2 \cdot (\text{H}_2\text{O})$



$1.30 \times 0.50 \times 0.50 \text{ mm}$



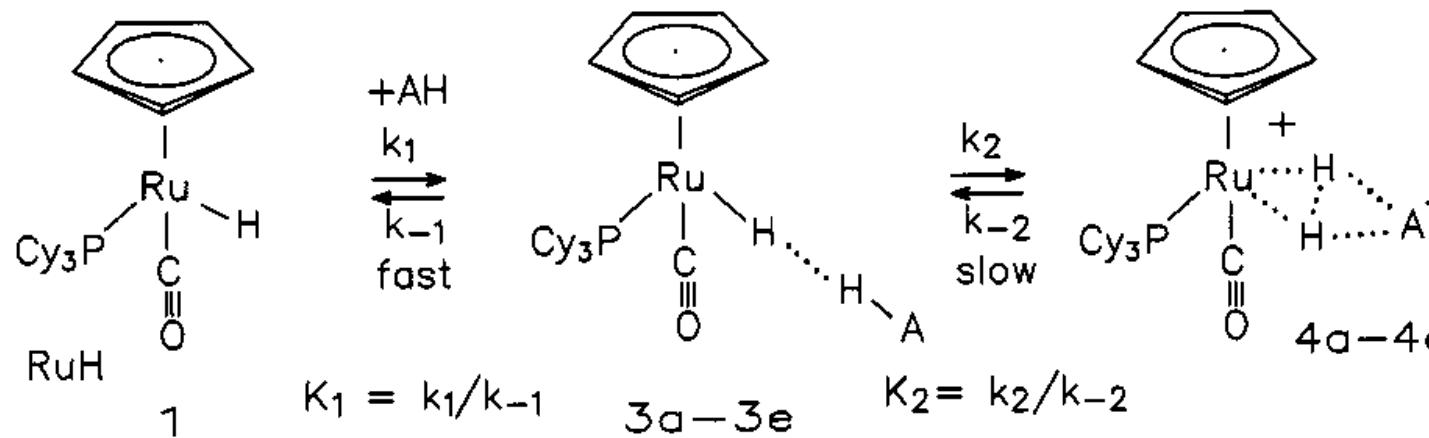
D19 @ 20K

Pt – N1  $2.040(2)$  Å

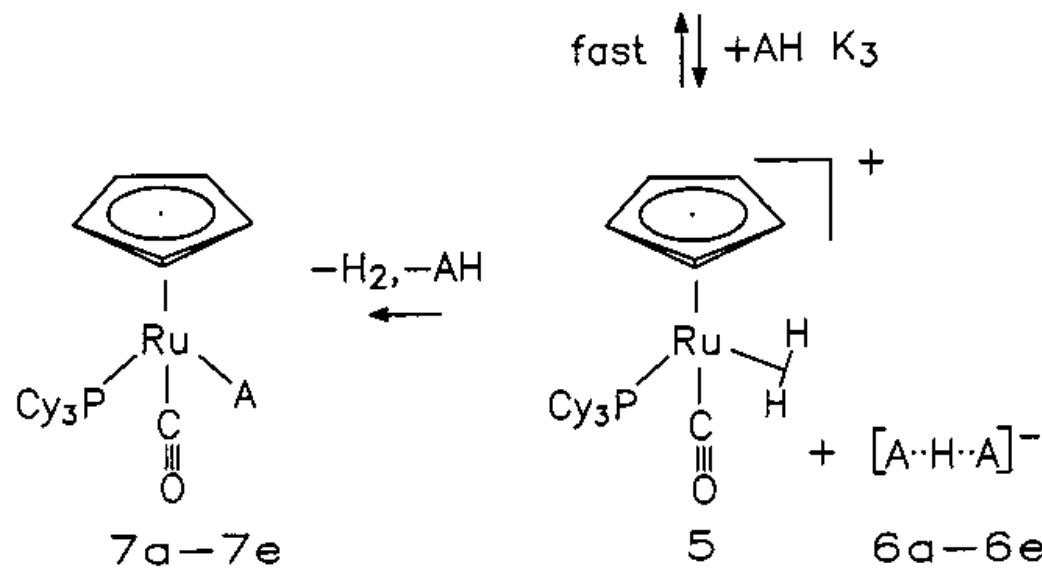
Pt – H1  $2.89(1)$  Å

N – C  $1.480(3)$  Å

# Proton Transfer to CpRuH(CO)(PCy<sub>3</sub>)



- 2 AH
- 2a  $\text{CF}_3\text{CH}_2\text{OH}$
- 2b  $(\text{CF}_3)_2\text{CHOH}$
- 2c  $(\text{CF}_3)_3\text{COH}$
- 2d  $\text{CF}_3\text{COOH}$
- 2e  $\text{HBF}_4^-$

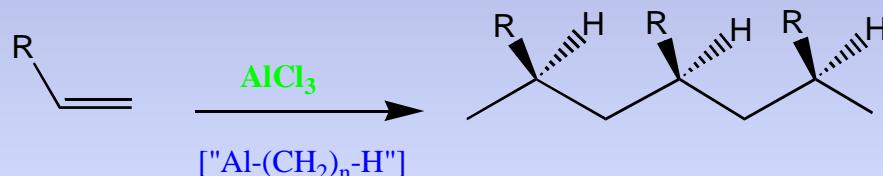


# Reactions Involving Hydrido-Complexes

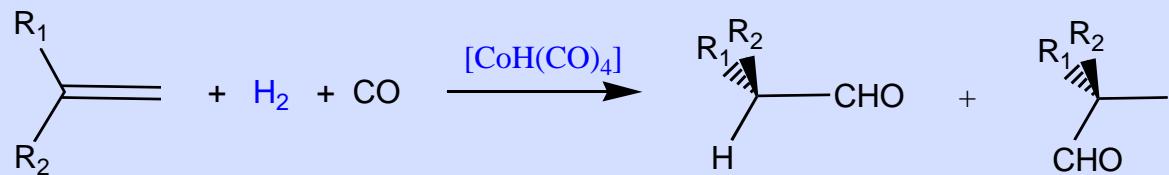
Homogeneous Hydrogenation



Olefin Polymerization

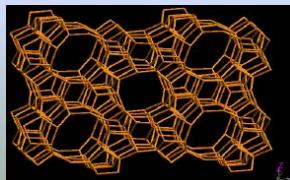
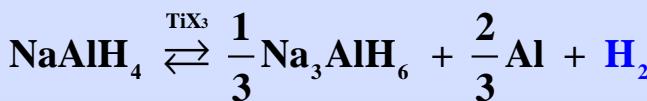


Olefin Hydroformylation



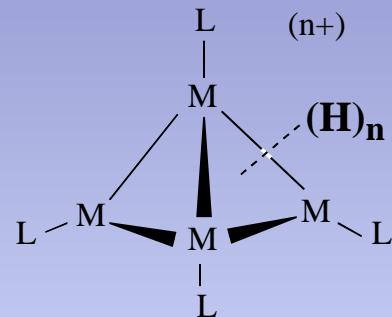
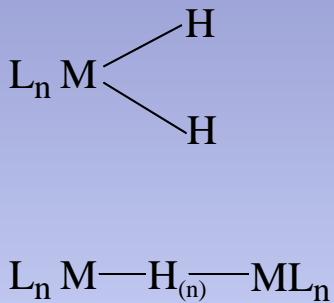
Materials for Hydrogen Storage

(new entry)

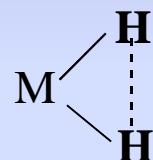


# Types of Metal - Hydrogen Interactions.

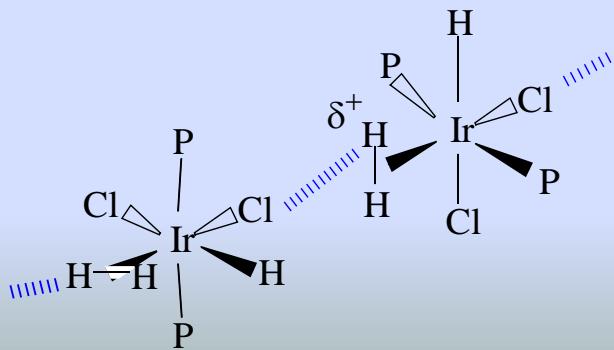
*Classical Metal Hydrides:*



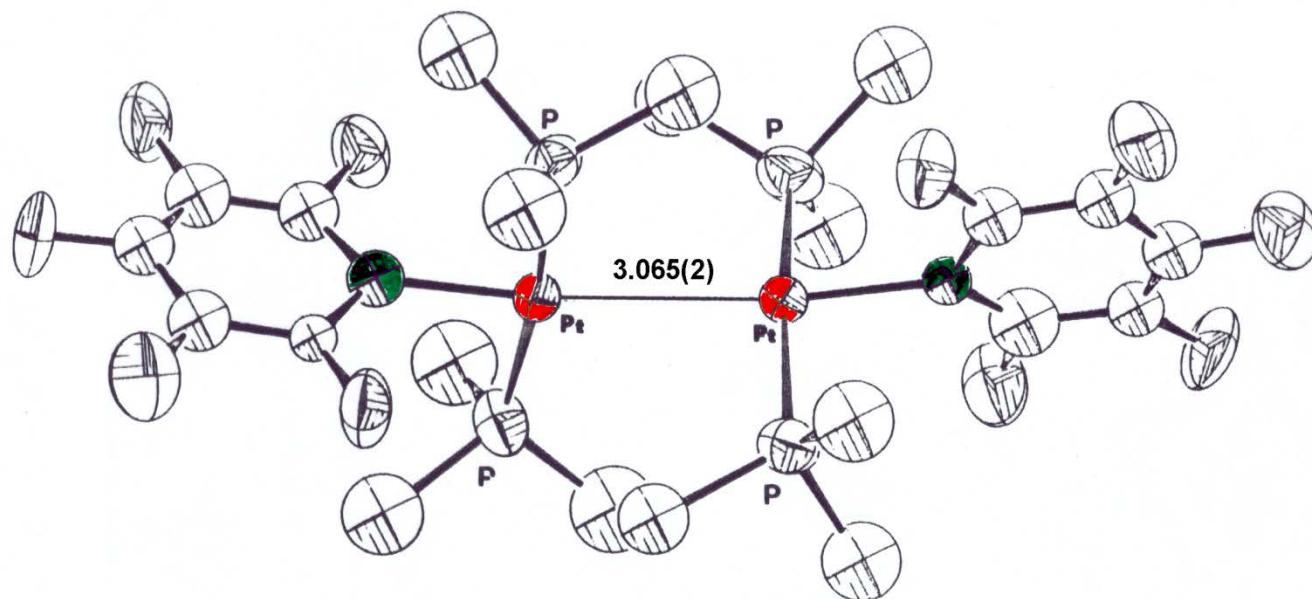
*"Non-Classical" Metal Hydrides:*



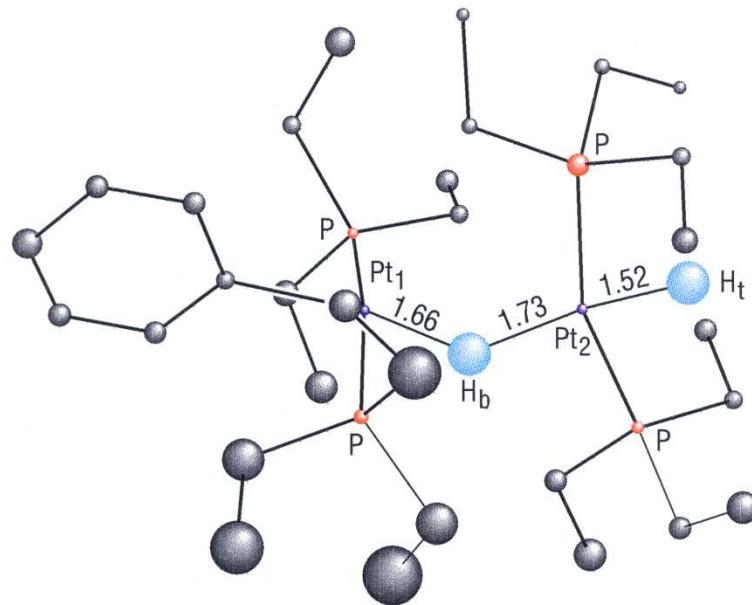
*Non-Classical "Weak" Interactions:*



# X-ray Structure of $[(\text{Cl}_5\text{C}_6)(\text{PMet}_3)_2\text{Pt-H-Pt}(\text{PMet}_3)_2(\text{C}_6\text{Cl}_5)]^+$



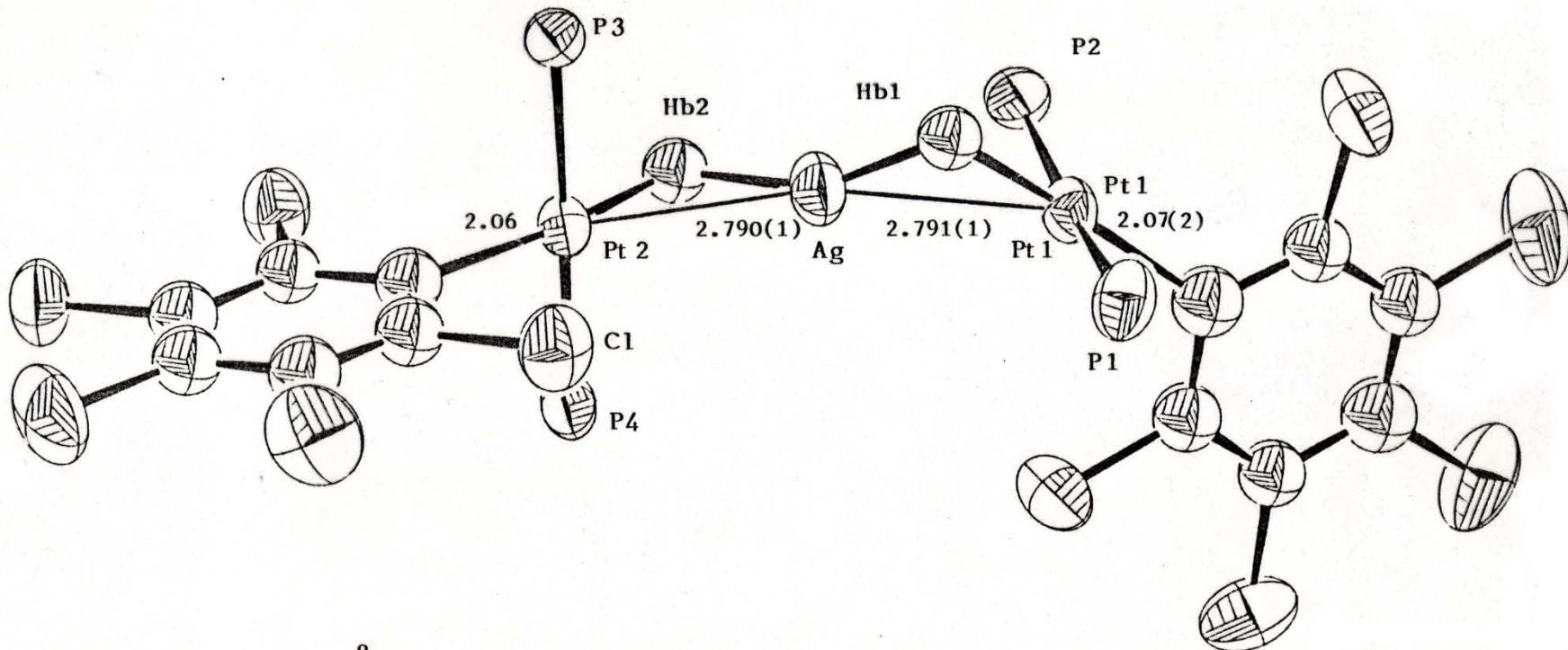
$[H_2Pt_2Ph(PEt_3)_4]^+[BPh_4]^-$   
13K Neutron Diffraction Study



Pt<sub>1</sub>···Pt<sub>2</sub> 3.05(1) Å  
 $\angle$  Pt<sub>1</sub>-H<sub>b</sub>-Pt<sub>2</sub> 128(2)<sup>°</sup>

Ricci, Albinati and Koetzle (1995)

# X-ray Structure of $[(\text{Cl}_5\text{C}_6)\text{L}_2\text{Pt} - \text{H} - \text{Ag} - \text{H} - \text{Pt}(\text{C}_6\text{Cl}_5)]^+$ (R.T. data)



Pt1 - Hb1 1.70(8) Å

Pt2 - Hb2 1.80(10)

Ag - Hb1 1.79(11)

Ag - Hb2 1.62(14)

Pt1 - P1 2.311(5)

Pt1 - P2 2.330(5)

Pt2 - P3 2.334(6)

Pt2 - P4 2.306(6)

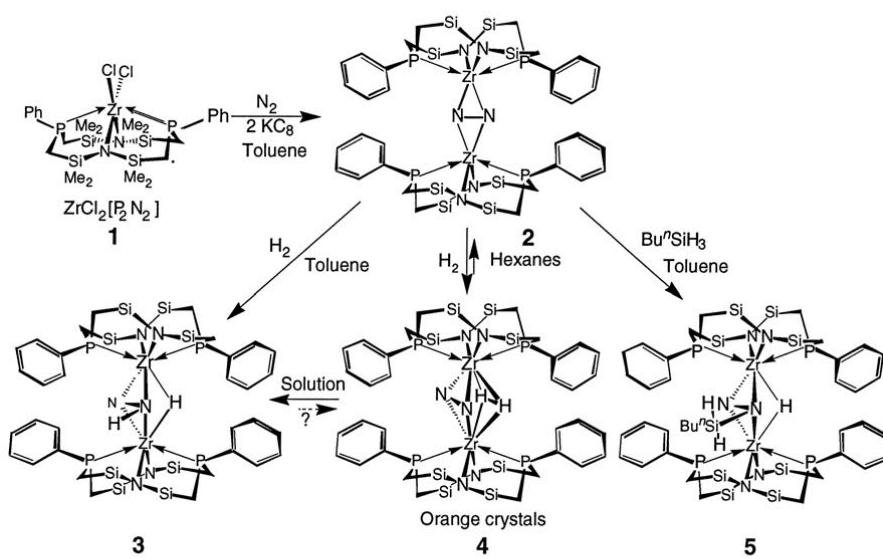
Pt1 - Ag - Pt2 166.04(4)°

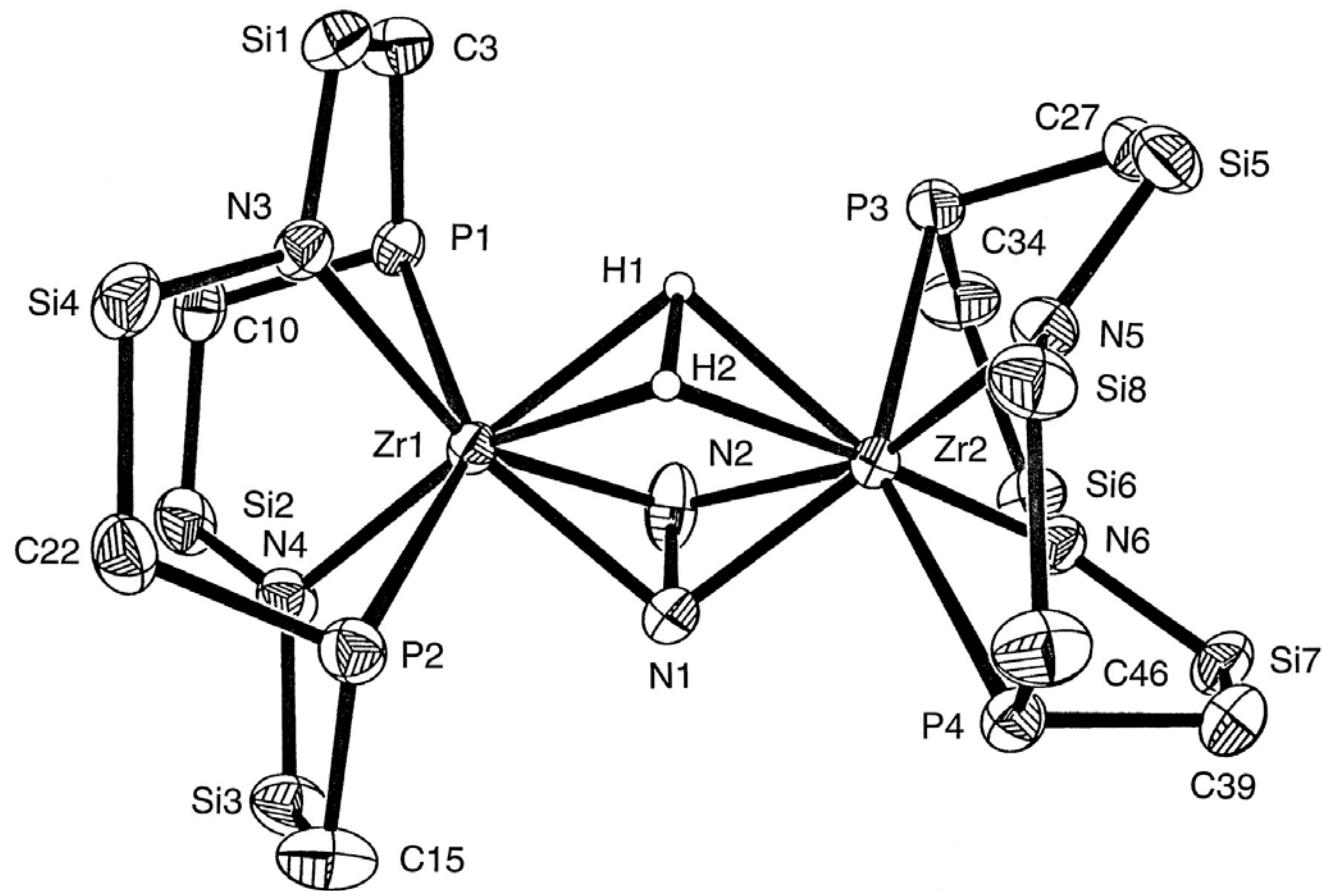
Ag - Pt - C<sub>av</sub> 148.5 (7)°

Pt - Hb - Ag 107 (1)°

Hb1 - Ag - Hb2 152 (2)°

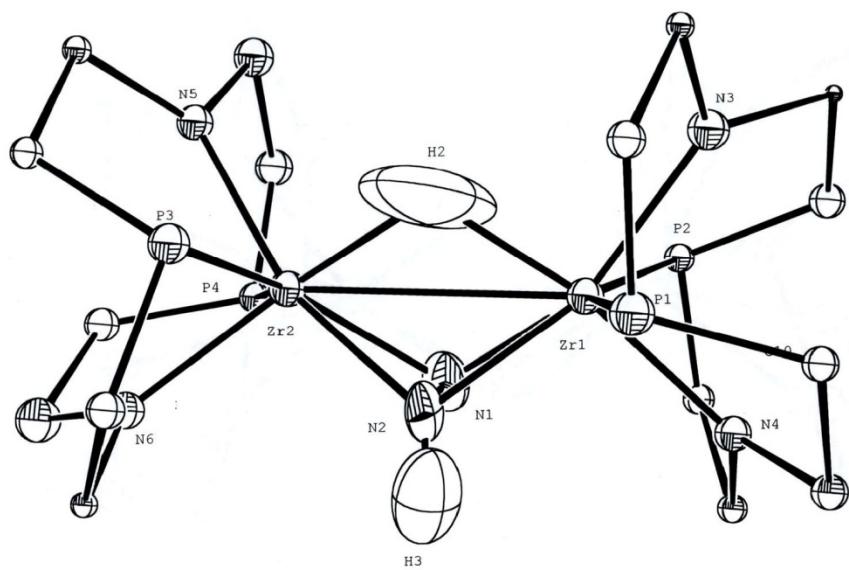
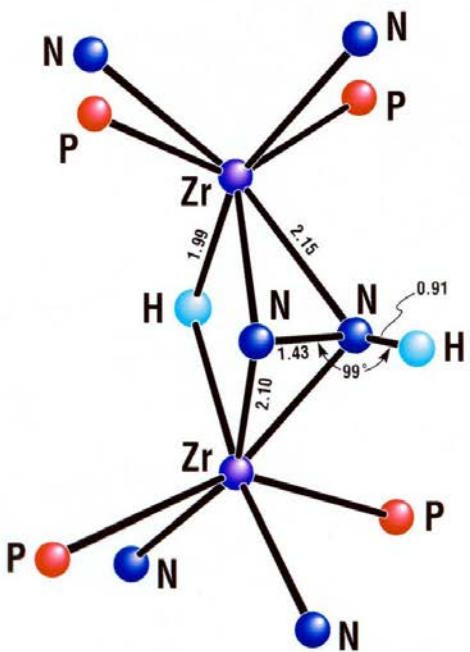
## Scheme for the Preparation of N<sub>2</sub> and “H<sub>2</sub>“ Complexes from ZrCl<sub>2</sub>[P<sub>2</sub>N<sub>2</sub>]







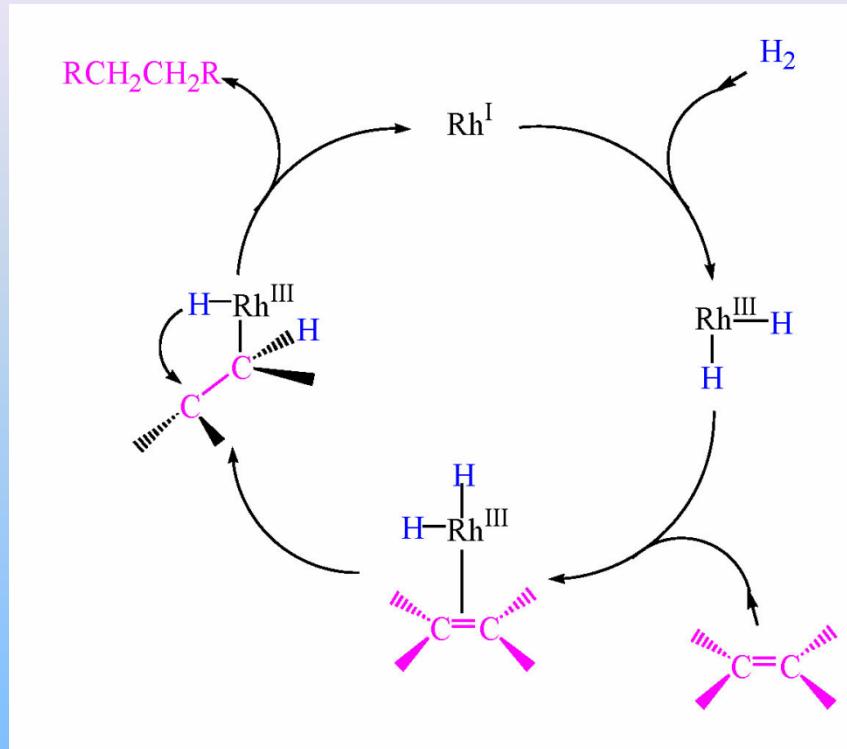
Neutron Structure at 25 K



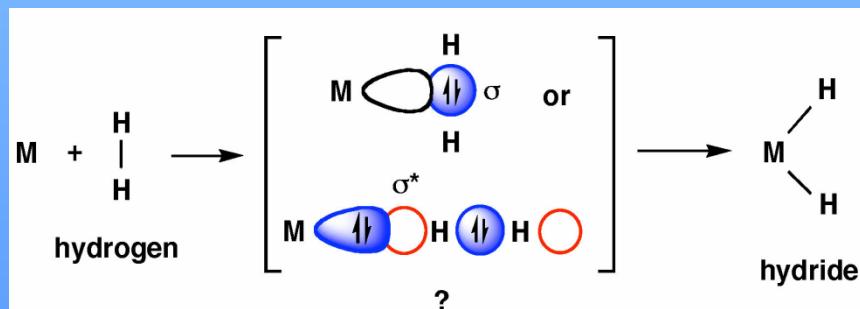
Albinati, Fryzuk, Klooster, Koetzle and Mason (1998)

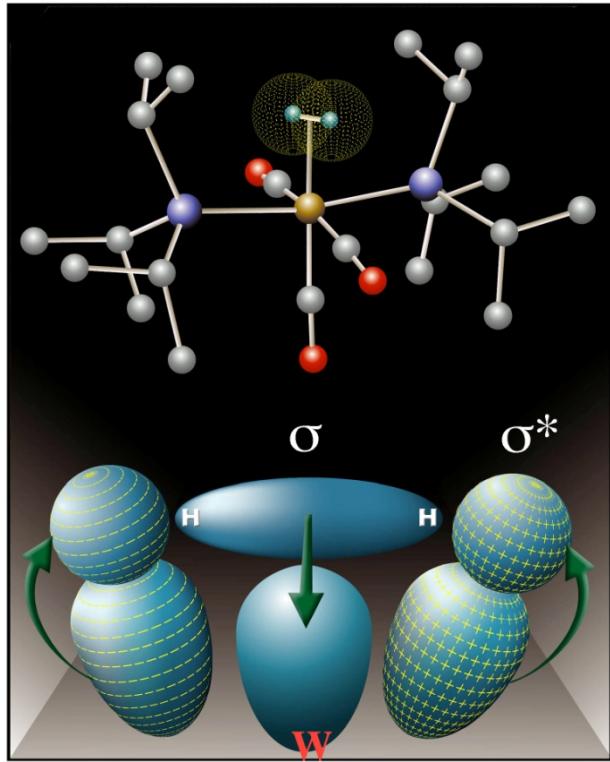


# Schematic Catalytic Cycle for Homogeneous Hydrogenation

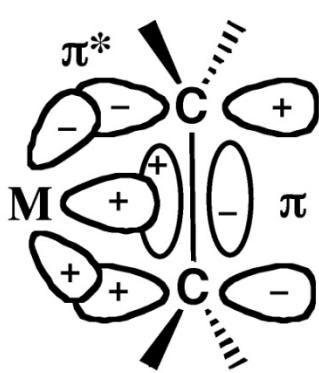


Metals break H-H bonds but HOW?



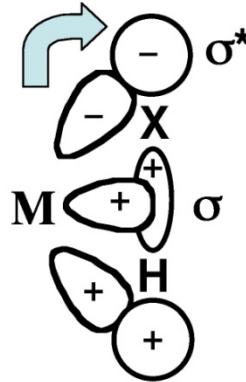


**backdonation is critical in stabilizing  
H<sub>2</sub> and other σ complexes**



M–π bond

olefin complex



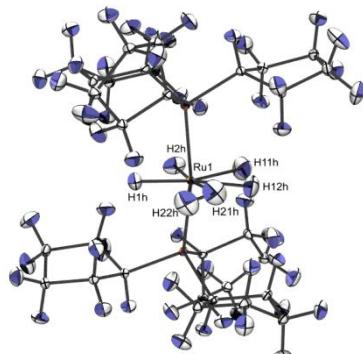
M–σ bond

X = H, C, Si, etc

Bonding in W(CO)<sub>3</sub>(P*i*Pr<sub>3</sub>)<sub>2</sub>(H<sub>2</sub>):  
*donation of the bonding σ electrons*  
*in H<sub>2</sub>* to a filled metal d orbital and  
backdonation to the antibonding  
orbital (σ\*) of H<sub>2</sub>

Hay, P. J. *Chem. Phys. Lett.* **1984**, *103*, 466.

# H-H Distances ( $\text{\AA}$ ) from Single XXL Neutron Diffraction



$[\text{Ru}(\text{dppe})_2(\text{H})(\text{H}_2)]^+$

0.82(3)

$[\text{Ru}(\text{cp}^*)(\text{dppm})_2(\text{H}_2)]^+$

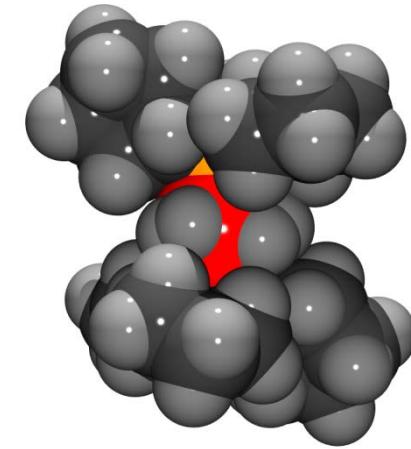
1.08(3)

$\text{Ir}(\text{P}^i\text{Pr}_3)_3\text{Cl}_2(\text{H})(\text{H}_2)$

1.11(3)

$\text{Ir}(\text{P}^i\text{Pr}_3)_2\text{I}(\text{H})_2(\text{H}_2)$

0.856(9)



$[\text{Os}(\text{en})_2\text{I}(\text{H}_2)]^+$

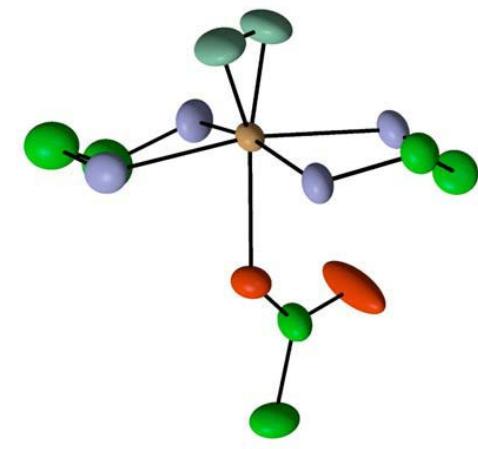
1.224(7)

$[\text{Os}(\text{en})_2(\text{OAc})(\text{H}_2)]^+$

1.34(2)

$[\text{Os}(\text{PMe}_2\text{Ph})_3(\text{H})_5]^+$

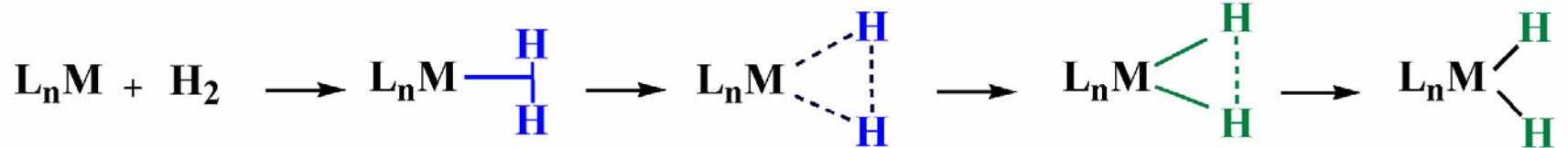
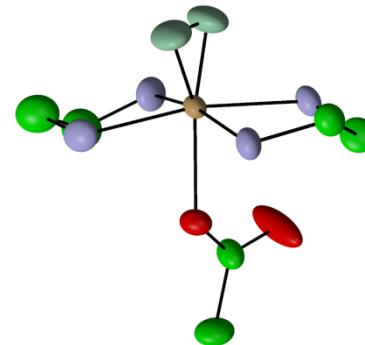
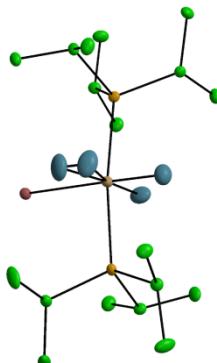
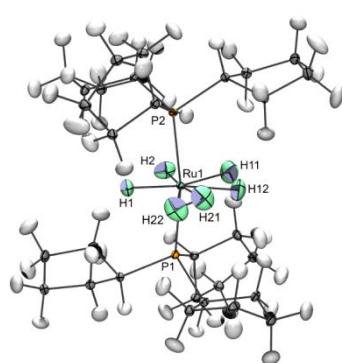
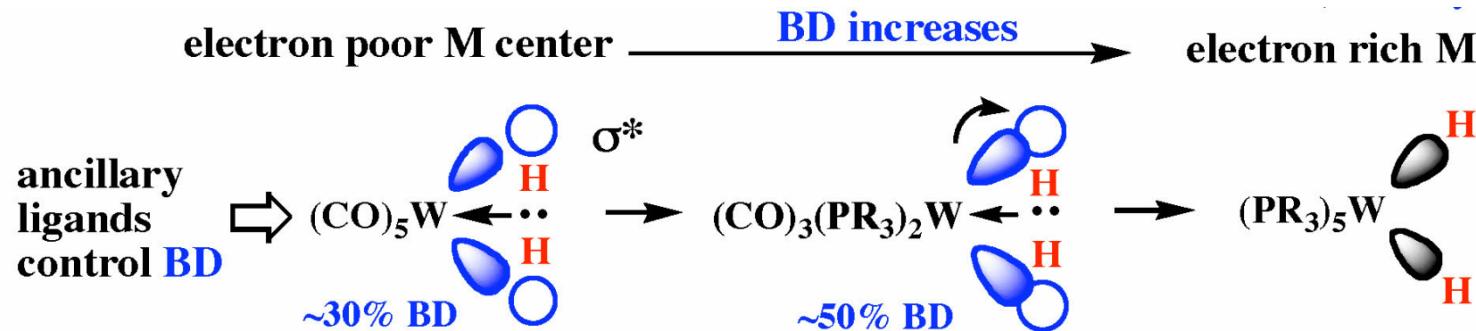
$\geq 1.49(4)$



$\text{Ir}(\text{P}^i\text{Pr}_3)_2\text{Br}(\text{H})_2(\text{H}_2)$

0.819(8)

# Backdonation is critical for the stability of M-H<sub>2</sub> Complexes



$\eta^2$  0.8 – 0.9 Å

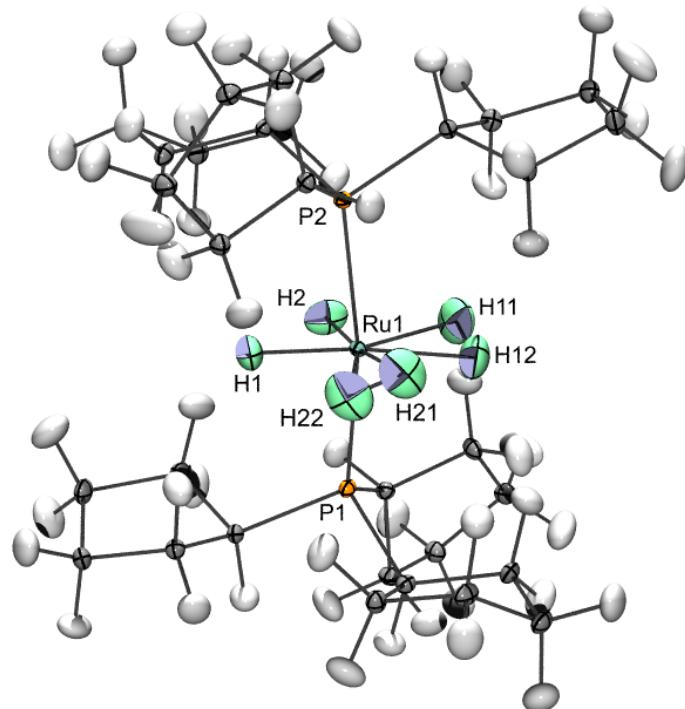
stretched 1.0 – 1.2 Å

weak interactions  $\geq 1.3 \text{ \AA}$

$\geq 1.7 \text{ \AA}$

# The Structure of RuH<sub>2</sub>(H<sub>2</sub>)<sub>2</sub>((Cyp)<sub>3</sub>)<sub>2</sub>

D19 @20K



H – H      0.825(7)      0.835(7) Å

P - Ru - P      168.9(1)°

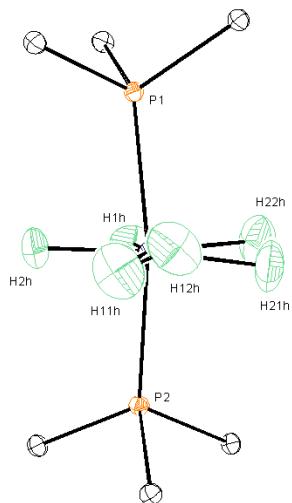
Ru - (H1)    1.730(5)      1.753(5) Å

Ru - (H2)    1.745(5)      1.764(5) Å

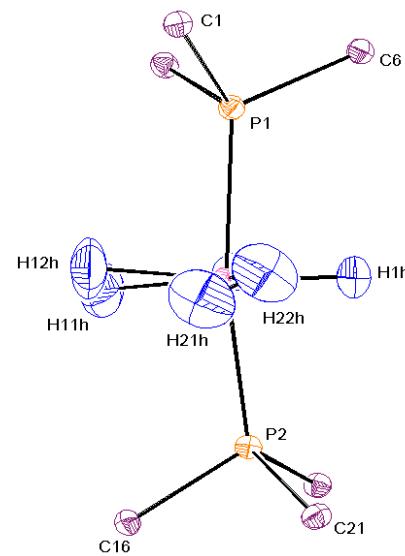
Ru - H      1.628(4)      1.625(4) Å

# Ru(H)<sub>2</sub>(H<sub>2</sub>)(P(cyp)<sub>3</sub>)<sub>2</sub>: a Parametric Study on D19

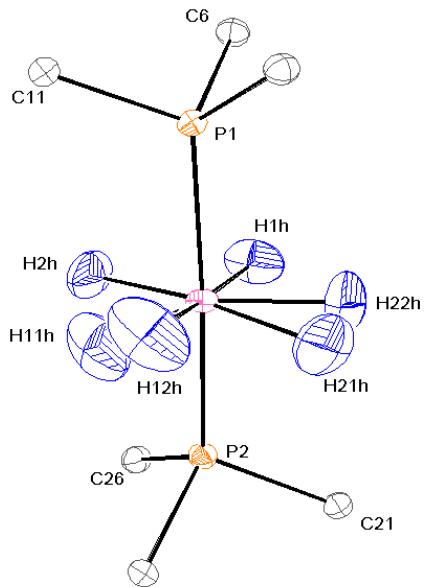
20K



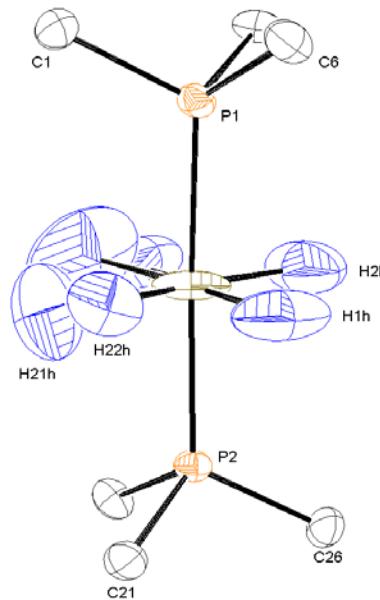
60K



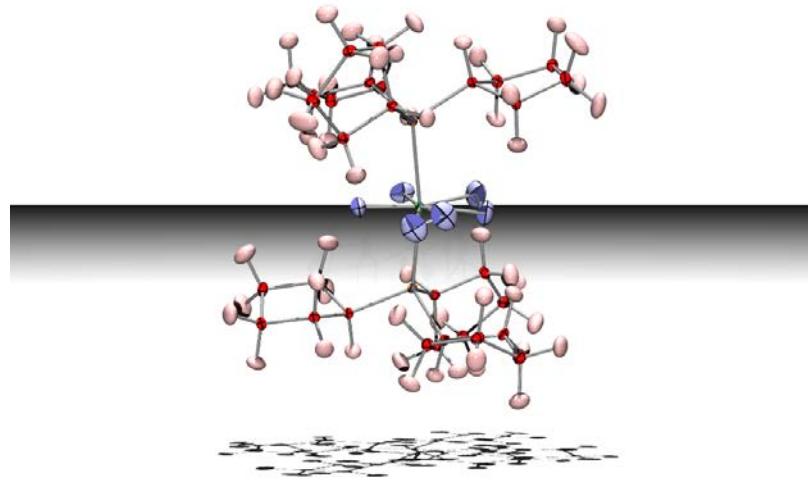
100K



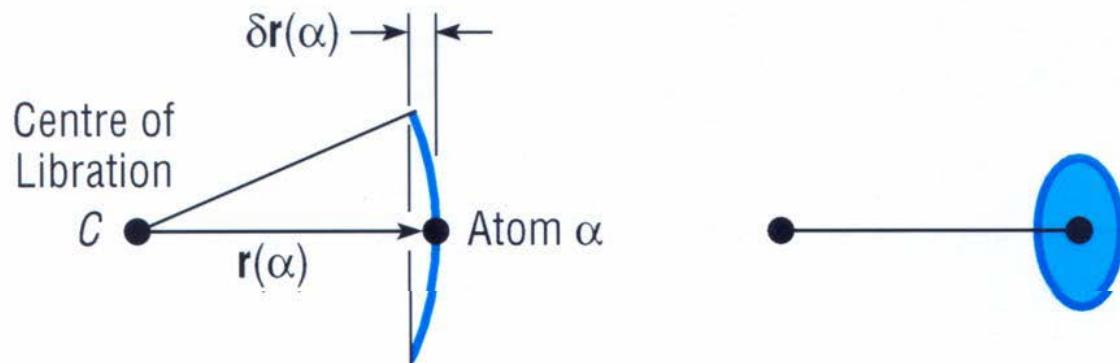
180K



# **RuH<sub>2</sub>(H<sub>2</sub>)<sub>2</sub>(P(Cyp)<sub>3</sub>)<sub>2</sub>**

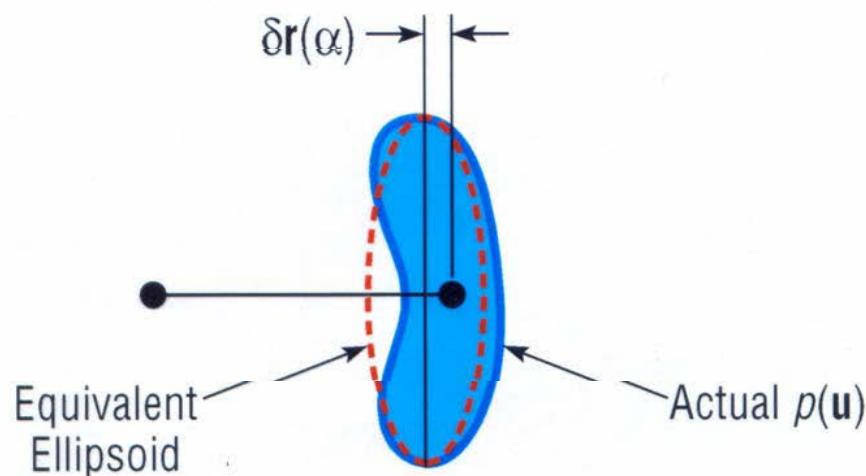


	<b>20K</b>	<b>60K</b>	<b>100K</b>	<b>180K</b>
<b>H11 – H12</b>	<b>0.825(8)</b>	<b>0.796(8)</b>	<b>0.794(8)</b>	<b>0.73(5)</b>
<b>Ru – H1</b>	<b>1.628(4)</b>	<b>1.618(4)</b>	<b>1.626(4)</b>	<b>1.69(2)</b>
<b>Ru – P1</b>	<b>2.307(3)</b>	<b>2.309(3)</b>	<b>2.310(2)</b>	<b>2.325(8)</b>
<b>P1 – Ru – P2</b>	<b>168.9(1)</b>	<b>168.8(1)</b>	<b>168.72(9)</b>	<b>178.0(5)</b>



P.D.F. from Libration

P.D.F. from Translation



Total P.D.F. from  
Translation and Vibration

## Rigid Molecule – The TLS Model

$$\mathbf{B}^{atom}(\kappa) = \left\langle \mathbf{u}(\kappa) (\mathbf{u}(\kappa))^T \right\rangle$$

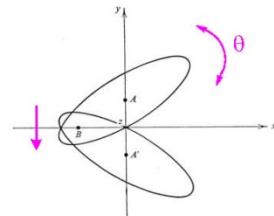
*(3x1) matrix;  $u_i = inst. disp$*

*time average*

for a rigid molecule

$$\mathbf{B}^{mol}(\kappa) = \begin{pmatrix} \mathbf{T} & \mathbf{S} \\ (\mathbf{S}^*)^T & \mathbf{L} \end{pmatrix}$$

$$\begin{aligned}\mathbf{T}(\kappa) &= \left\langle \mathbf{u}(\kappa) (\mathbf{u}(\kappa))^T \right\rangle \\ \mathbf{L}(\kappa) &= \left\langle \boldsymbol{\theta}(\kappa) (\boldsymbol{\theta}(\kappa))^T \right\rangle \\ \mathbf{S}(\kappa) &= \left\langle \mathbf{u}(\kappa) (\boldsymbol{\theta}(\kappa))^T \right\rangle\end{aligned}$$



*distance of atom  $\kappa$  from center of libration*

$$\boxed{\delta \mathbf{r}(\kappa) = -\frac{1}{2} [(\text{trace} \mathbf{L}) \mathbf{r}(\kappa) - \mathbf{L} \mathbf{r}(\kappa)]}$$

## H - H Distances (Å) and TLS Correction

	Uncorrected	TLS	Solid State NMR
Fe(PEtPh <sub>2</sub> ) <sub>3</sub> (H) <sub>2</sub> (H <sub>2</sub> )	<b>0.82 (1)</b>	<b>1.05</b>	
[Fe(dppe) <sub>2</sub> (H)(H <sub>2</sub> )][BF <sub>4</sub> ]	<b>0.82 (2)</b>	<b>0.85</b>	<b>0.90</b>
[Os(dppe) <sub>2</sub> (H)(H <sub>2</sub> )][PF <sub>6</sub> ]	<b>0.79 (2)</b>	<b>0.96</b>	<b>0.99</b>
[Os(dppe) <sub>2</sub> (Cl)(H <sub>2</sub> )][PF <sub>6</sub> ]	<b>1.15 (3)</b>	<b>1.24</b>	<b>1.19</b>
[Os(en) <sub>2</sub> (I)(H <sub>2</sub> )]	<b>1.224 (7)</b>	<b>1.272</b>	
[Ru(cp*)(dppm) <sub>2</sub> (H <sub>2</sub> )][BF <sub>4</sub> ]	<b>1.08 (3)</b>	<b>1.09</b>	<b>1.02</b>

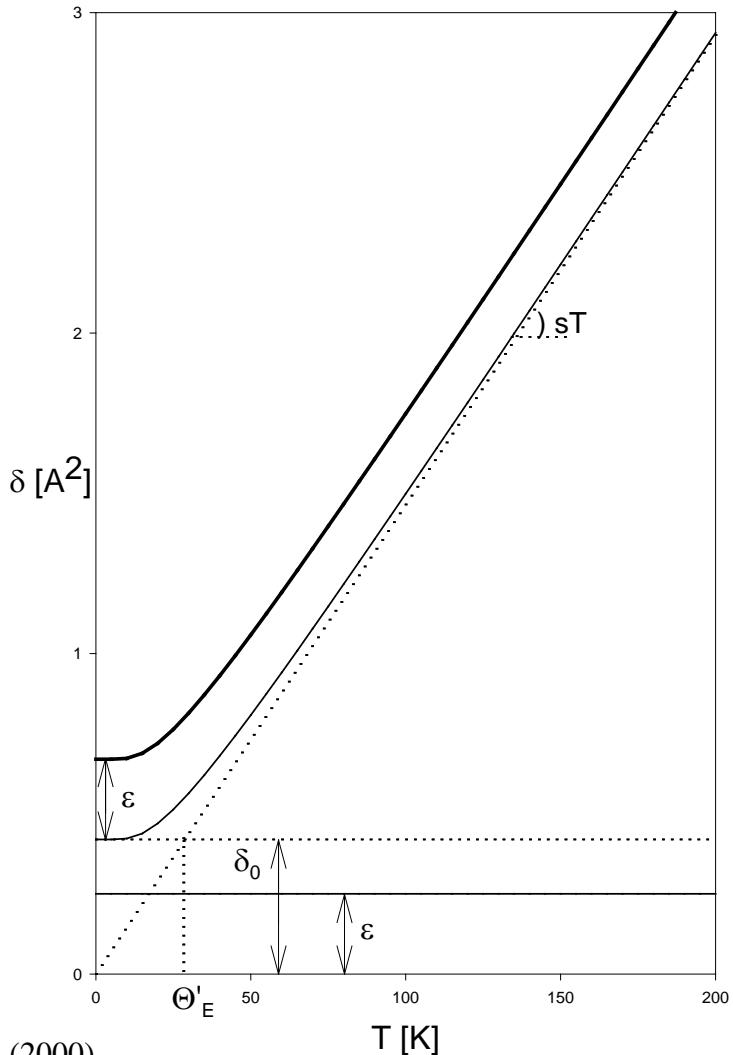
# Molecular Motion from the Temperature Dependence of ADP's

$$\Sigma^x(T) = A g V \delta(1/\omega, T) V' g' A' + \varepsilon$$

ADPs,  
determined  
experimentally  
at several  
temperatures

Intramolecular,  
**high frequency  
vibrations and  
disorder ( $\varepsilon$ )**  
(~temperature  
independent)

Low frequency, **large-amplitude vibrations**  
( $\omega$ ), e.g. librations,  
translations and  
deformations ( $V$ )



# Molecular Motion from the Temperature Dependence of the ADPs

$$\sum^X(\textcolor{blue}{T}) = \mathbf{A}\mathbf{g}\mathbf{V}\boldsymbol{\delta}(\omega^{-1}, T)\mathbf{V}^T\mathbf{g}^T\mathbf{A}^T + \boldsymbol{\varepsilon}$$

$\Sigma(T)$  observed atomic mean square displ. at multiple T

$\mathbf{A}$  transformation matrix from atomic to normal modes coords

$\mathbf{V}$  eigenvectors matrix

$\boldsymbol{\varepsilon}$  3x3 upper triangular tensor

contribution to ADPs from high v small amplitude vibrations

$\boldsymbol{\delta}(T)$  diagonal matrix of mean square normal modes displacements

For each normal mode j

$$\boldsymbol{\delta}(T) = \left( \frac{h}{2\omega_j} \right) \coth \left( \frac{h\omega_j}{2k_b T} \right)$$

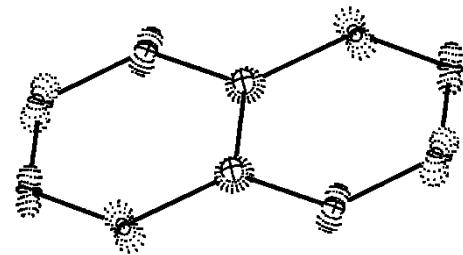
Considering only T and L

$$\sum^X = \mathbf{A} \begin{pmatrix} \mathbf{T} & \mathbf{S} \\ \mathbf{S}^T & \mathbf{L} \end{pmatrix} \mathbf{A}^T + \boldsymbol{\varepsilon}$$

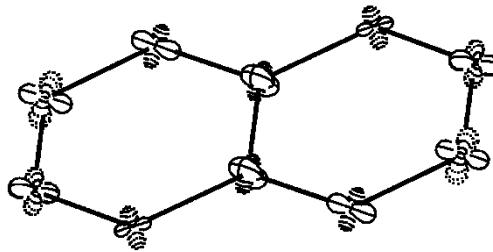
# Naphthalene : ADP's and rms difference surfaces

X-ray (SR)

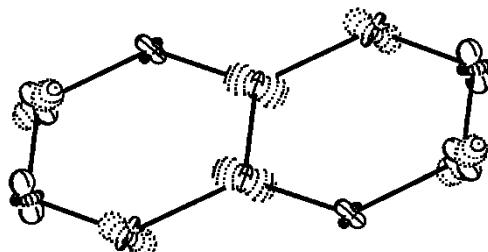
10 K



30 K

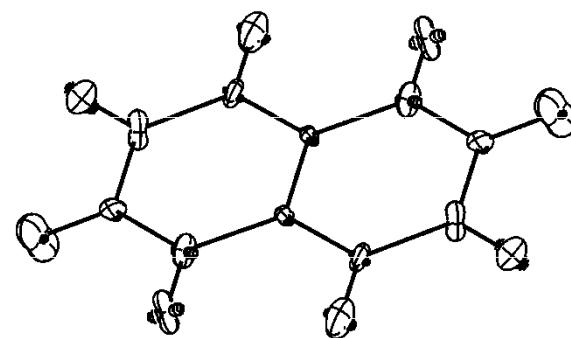
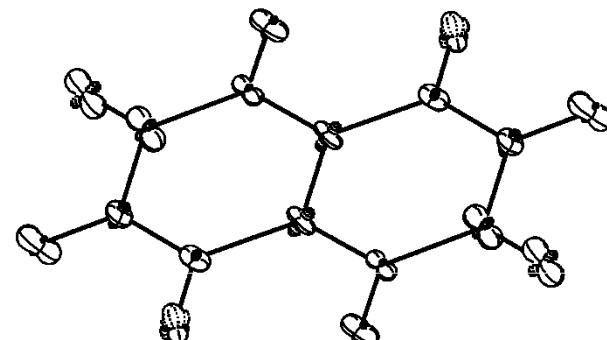
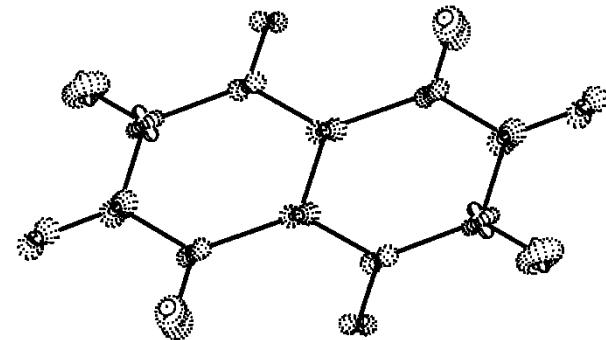


150 K



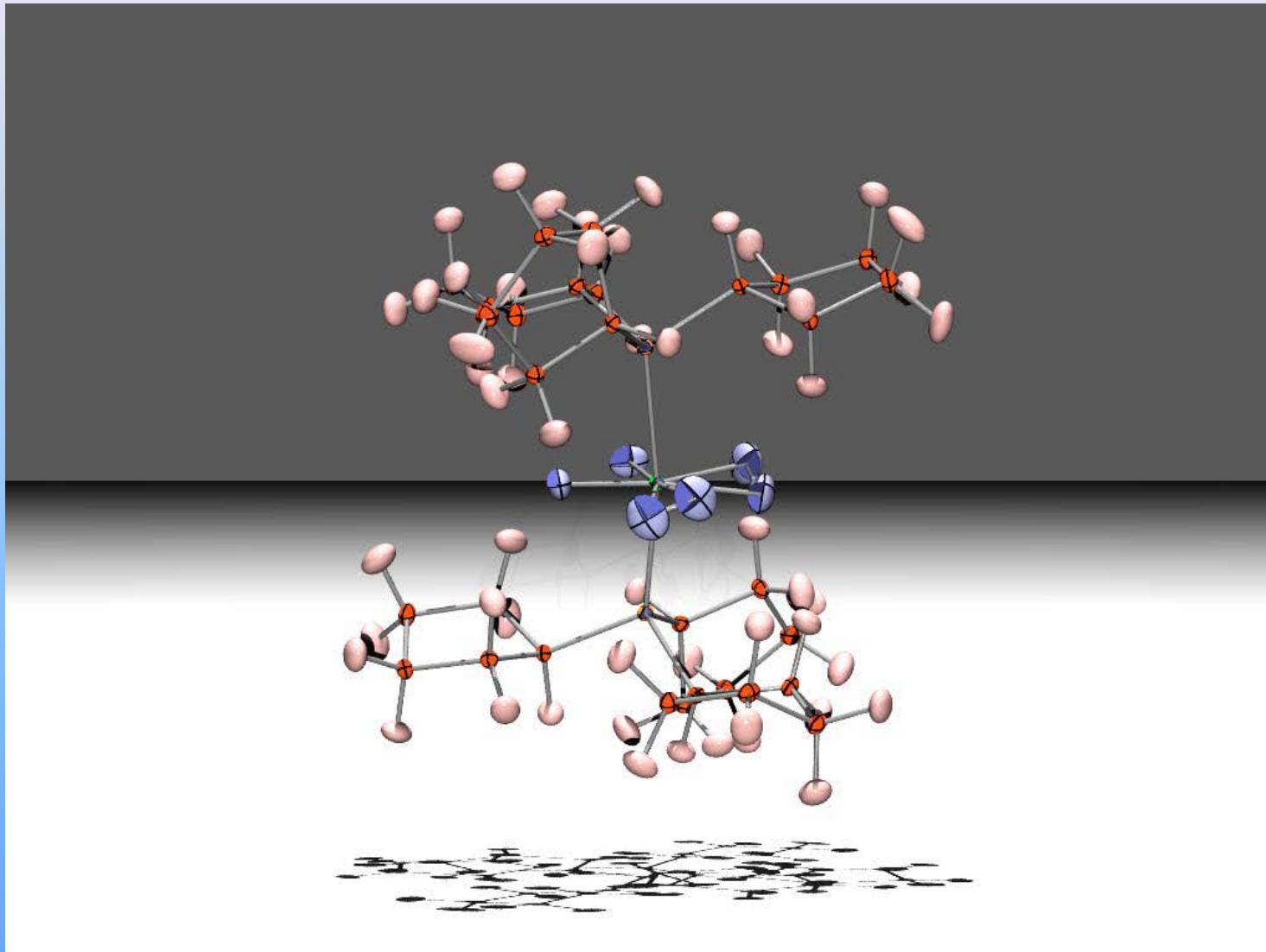
Neutron

5 K



# The Structure of $\text{RuH}_2(\text{H}_2)_2((\text{Cyp})_3)_2$

D19 @20K



# MODELS OF MOTION USED TO ANALYSE THE 20, 60 AND 100K NEUTRON DATA (with both THMA and ADP-analysis)

## 1) Rigid-body description:

the motion of the molecule as a whole is defined by 6 degrees of freedom :

$$L_x, L_y, L_z, T_x, T_y, T_z$$

→ Serious problems in describing the out-of-plane displacements of the H atoms involved in the di-hydrogen groups

$$\begin{array}{cccc} \varepsilon(H) \rightarrow & 0.017(2) & 0.001(1) & 0.005(1) \\ & 0.019(2) & -0.006(1) & \\ & & \textcolor{red}{0.058(2)} & \end{array}$$

## 2) Rigid-body + internal rotations of the di-hydrogen groups:

8 degrees of freedom:

$$L_x, L_y, L_z, T_x, T_y, T_z, U_1, U_2$$

→ Significant improvement in the description of the out-of-plane displacements in the di-hydrogen groups

$$\begin{array}{cccc} \varepsilon(H) \rightarrow & 0.018(2) & 0.000(1) & 0.000(1) \\ & 0.018(2) & 0.005(1) & \\ & & \textcolor{red}{0.028(2)} & \end{array}$$

→ normal mode frequencies for di-hydrogen rotations: 104 and 170  $\text{cm}^{-1}$

# Models used for the ADP's analysis

Model	T (K)	R (%)	Goof	Obs	Par	Obs/Par ratio	Degrees of freedom
A	20 60 100	41 36 33	5.16 5.01 5.82	54	20	2.70	6 librations, 6 translations, 8 coupling terms
B	20 60 100	23 19 17	3.83 3.39 3.83	54	32	1.69	6 librations, 6 translations, 8 coupling terms + 6 components of the librations of the di-hydrogen groups
C	all	12	1.62	162	28	5.79	1 librational frequency, 3 translational frequencies, 24 components of the 4 temperature-independent $\epsilon$ -tensors
D	all	9	1.16	162	32	5.06	3 librational frequencies, 3 translational frequencies, 2 eigenvector components, 24 components of the 4 temperature-independent $\epsilon$ -tensors

# Bond Distances (Å) Corrected for Libration

T (K)	RU-H11		RU-H12		RU-H21		RU-H22	
	TLS	ADP	TLS	ADP	TLS	ADP	TLS	ADP
20	1.7397	1.7475	1.7618	1.7709	1.7732	1.7819	1.7537	1.7610
60	1.7322	1.7412	1.7606	1.7710	1.7699	1.7786	1.7488	1.7562
100	1.7435	1.7540	1.7543	1.7682	1.7673	1.7755	1.7491	1.7580

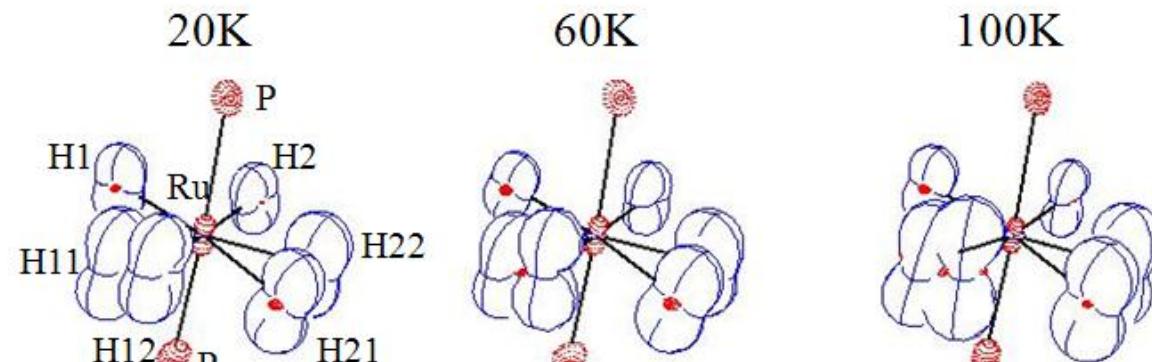
T (K)	H11-H12			H21-H22		
	TLS	ADP	Exp	TLS	ADP	Exp
20	0.831	<b>0.899</b>	<b>0.825(8)</b>	0.840	<b>0.911</b>	<b>0.835(8)</b>
60	0.804	<b>0.887</b>	0.796(8)	0.826	<b>0.902</b>	0.820(8)
100	0.806	<b>0.916</b>	0.794(8)	0.820	<b>0.910</b>	0.812(8)

THMA: rigid-body model

$R_{20K} = 41\%$  Goof = 5.16

$R_{60K} = 36\%$  Goof = 5.01

$R_{100K} = 33\%$  Goof = 5.82

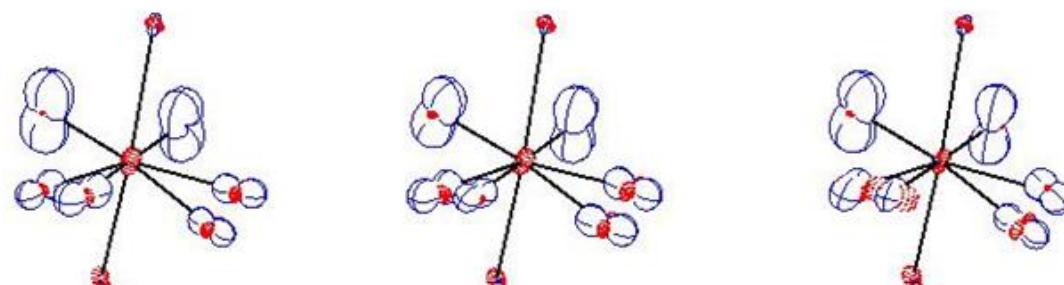


THMA: rigid-body model +  $H_2$  librations

$R_{20K} = 23\%$  Goof = 3.83

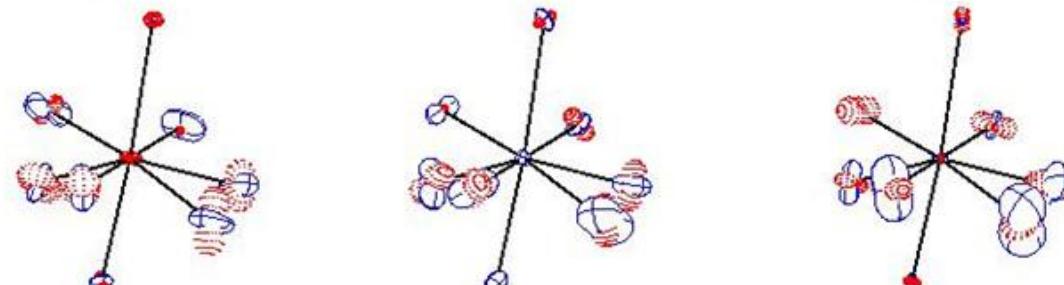
$R_{60K} = 19\%$  Goof = 3.39

$R_{100K} = 17\%$  Goof = 3.83



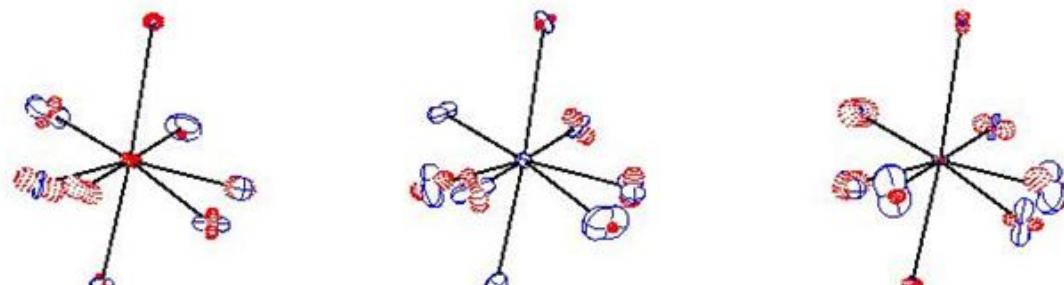
ADP-analysis:  
rigid-body model

$R_{all\_T} = 12\%$  Goof = 1.62



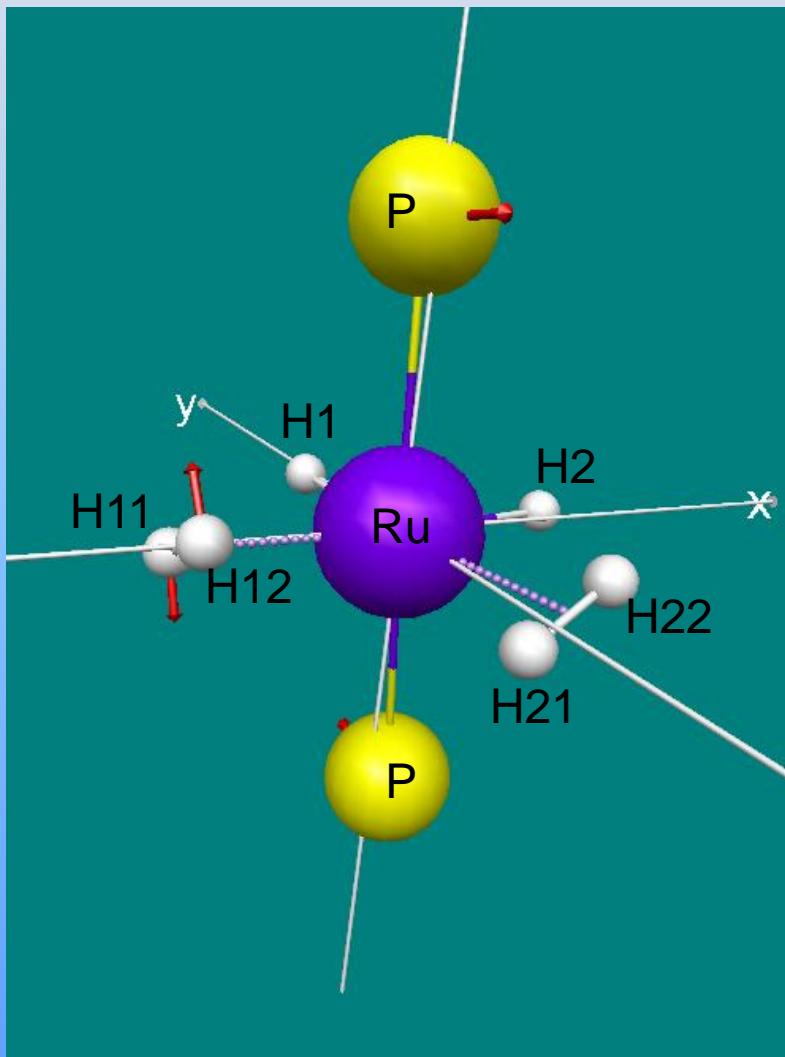
ADP-analysis: rigid-body  
model +  $H_2$  librations

$R_{all\_T} = 9\%$  Goof = 1.16

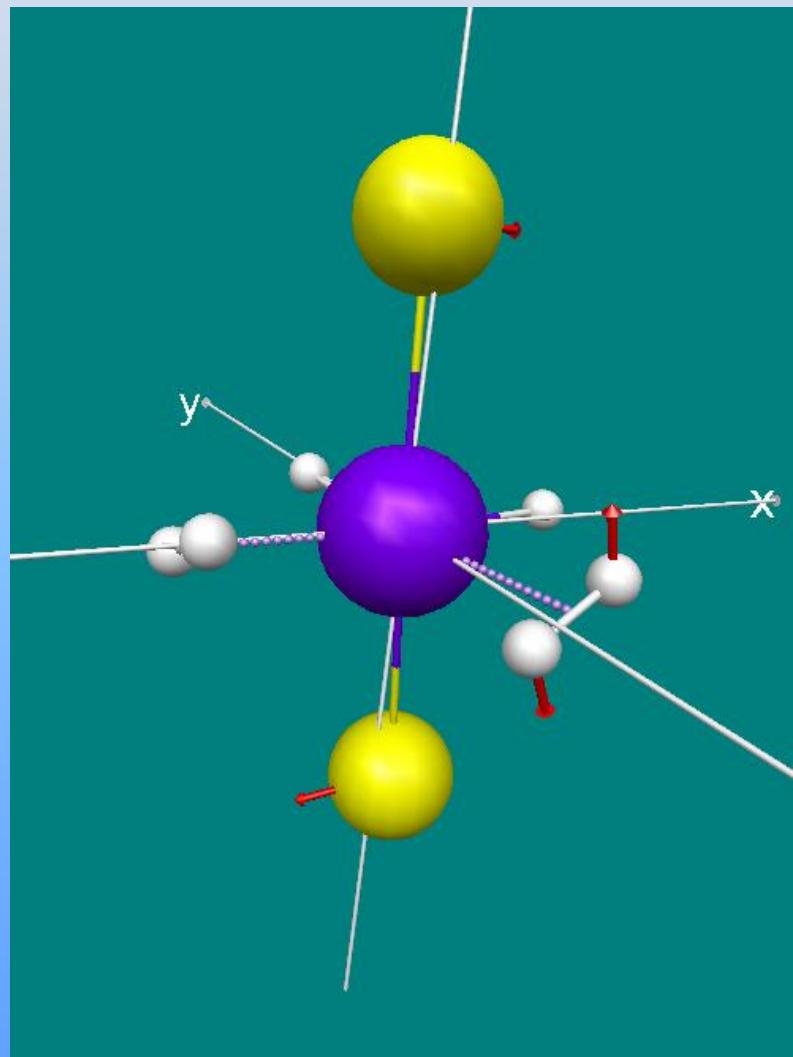


# Dihydrogen Ligands: Low Frequency Normal Modes

104.5  $\text{cm}^{-1}$



170.3  $\text{cm}^{-1}$



## **H-H Distances ( $\text{\AA}$ ) vs. $J^{\text{HD}}$ (Hz)**

	$d_{(\text{H-H})}$	$J^{\text{HD}}$
$\text{W}(\text{CO})_3(\text{P}^{\text{i}}\text{Pr}_3)_2(\text{H}_2)$	<b>0.82 (1)</b>	<b>34.0</b>
$[\text{Fe}(\text{dppe})_2(\text{H})(\text{H}_2)]^+$	<b>0.82 (2)</b>	<b>30.5</b>
$[\text{Ru}(\text{dppe})_2(\text{H})(\text{H}_2)]^+$	<b>0.82 (3)</b>	<b>32.0</b>
$[\text{Os}(\text{dppe})_2(\text{H})(\text{H}_2)]^+$	<b>0.79 (2)</b>	<b>25.5</b>

**We need more information.....**

**Another technique**