

Diffraction from Crystalline* Materials

Paul Henry

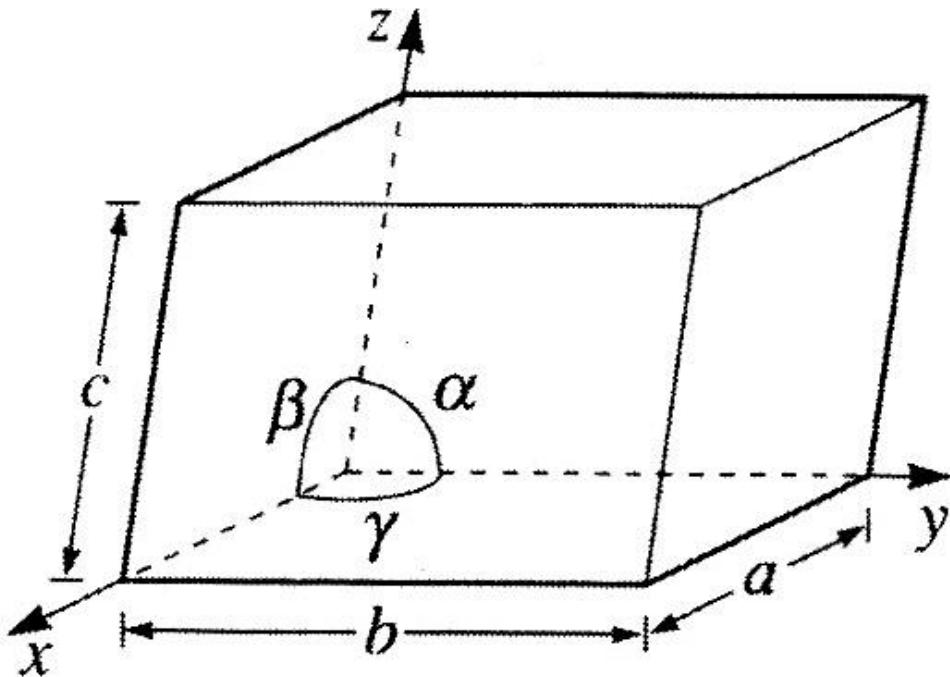
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Outline

- Key concepts
- Instrumentation for neutron diffraction
- Extra information 1: Uses of diffraction
- Extra information 2: Structure solution & refinement
- Extra Information 3: Diffraction suites of ILL and ISIS

Key concept 1a: lattice & unit cell



Conventions

- cell parameters are in Å or pm
- Angles are in °

The unit cell has lattice parameters defined by the cell length a , b , and c , and the cell angles α , β , and γ :

γ is angle between a and b

β is angle between a and c

α is angle between b and c

Atomic positions are given as xyz coordinates:

x is fraction of a axis

y is fraction of b axis

z is fraction of c axis

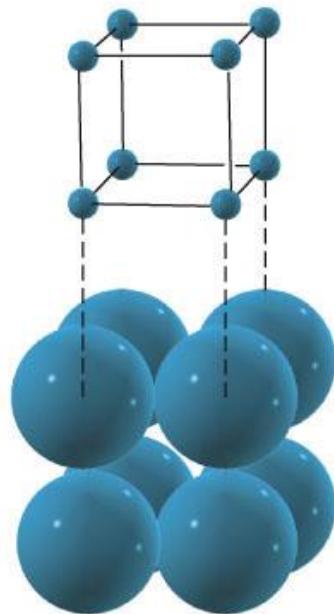
Key concept 1b: crystal systems

Triclinic	$a \neq b \neq c$	$\alpha \neq \beta \neq \gamma \neq 90^\circ$
Monoclinic	$a \neq b \neq c$	$\alpha = \gamma = 90^\circ$ $\beta \neq 90^\circ$
Orthorhombic	$a \neq b \neq c$	$\alpha = \beta = \gamma = 90^\circ$
Rhombohedral	$a = b = c$	$\alpha = \beta = \gamma \neq 90^\circ$
Hexagonal	$a = b \neq c$	$\alpha = \beta = 90^\circ$ $\gamma = 120^\circ$
Tetragonal	$a = b \neq c$	$\alpha = \beta = \gamma = 90^\circ$
Cubic	$a = b = c$	$\alpha = \beta = \gamma = 90^\circ$

Key concept 1c: centring

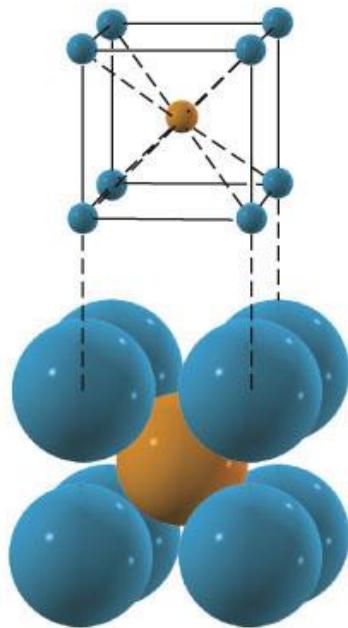
Primitive

P



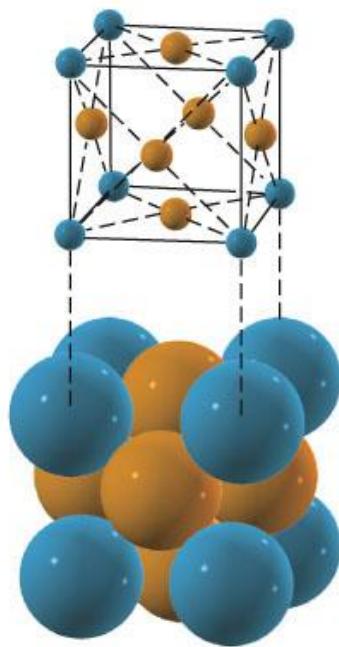
Body centred

I



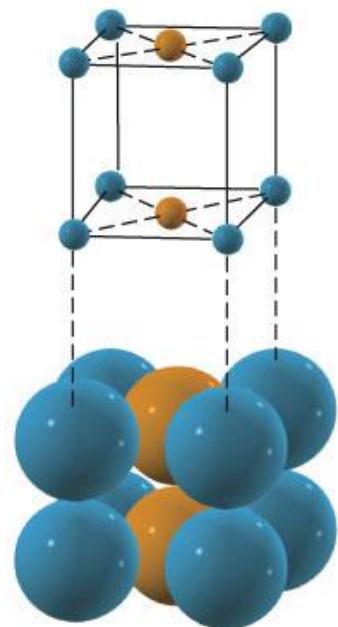
Face centred

F



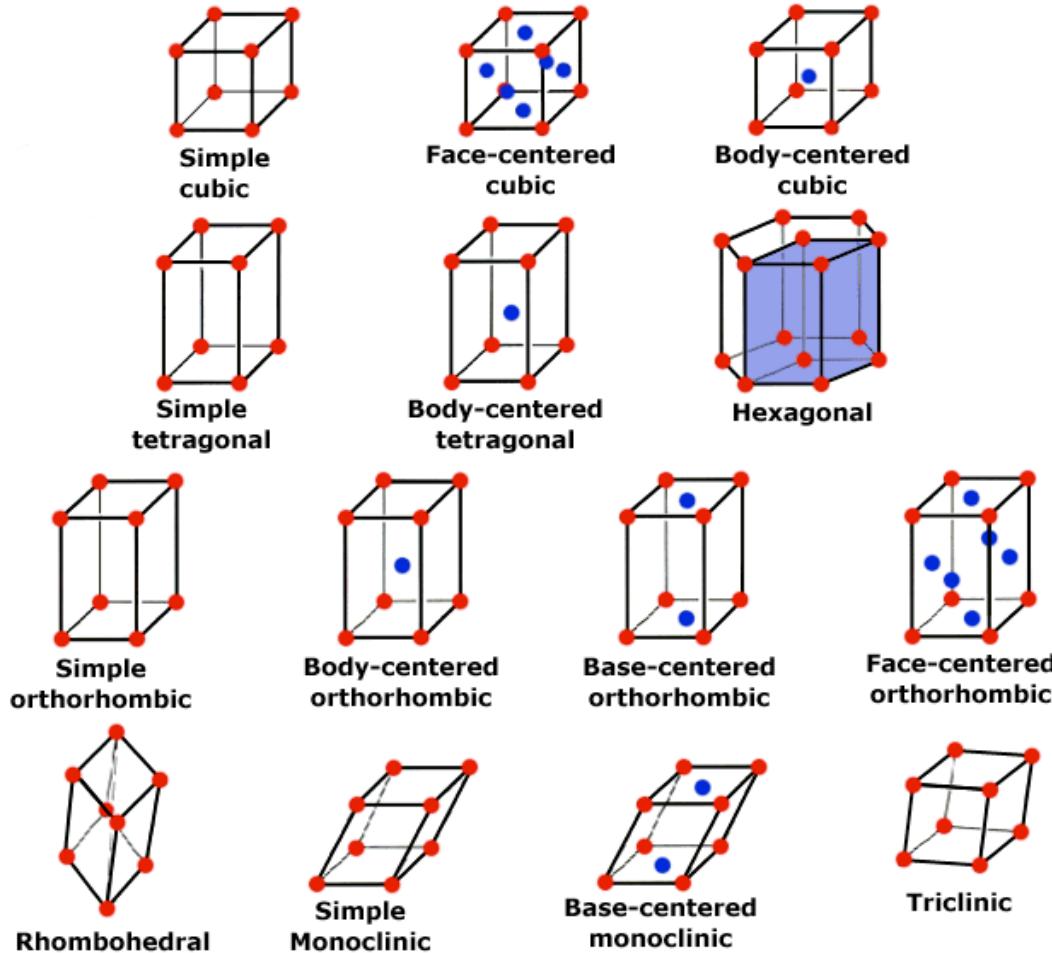
Side centred

C (A/B)



NB: Atom types are identical even though coloured differently

Key concept 1: Summary



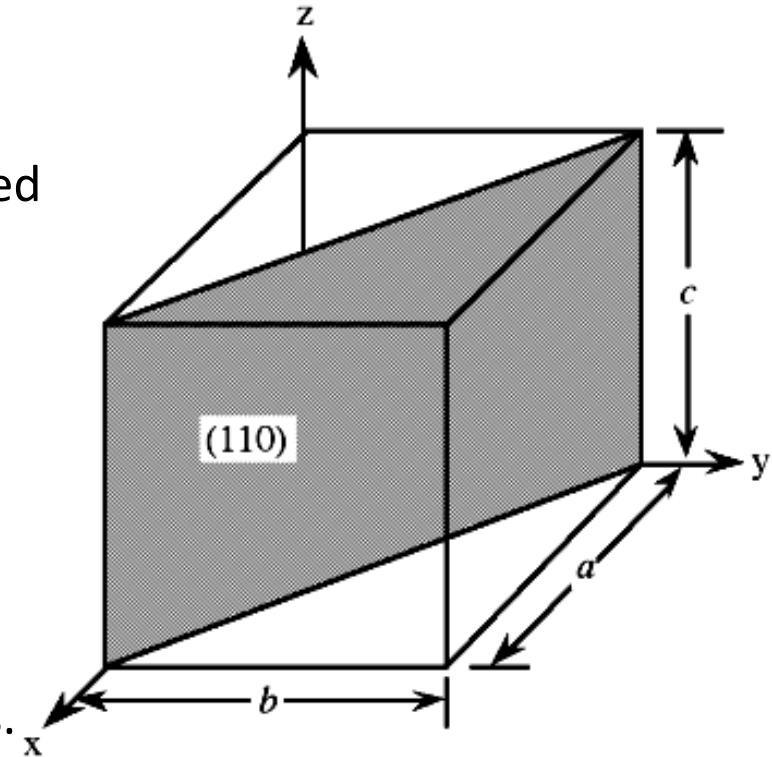
7 crystal classes
14 Bravais Lattice types
230 space groups

Key concept 2: Miller indices / planes

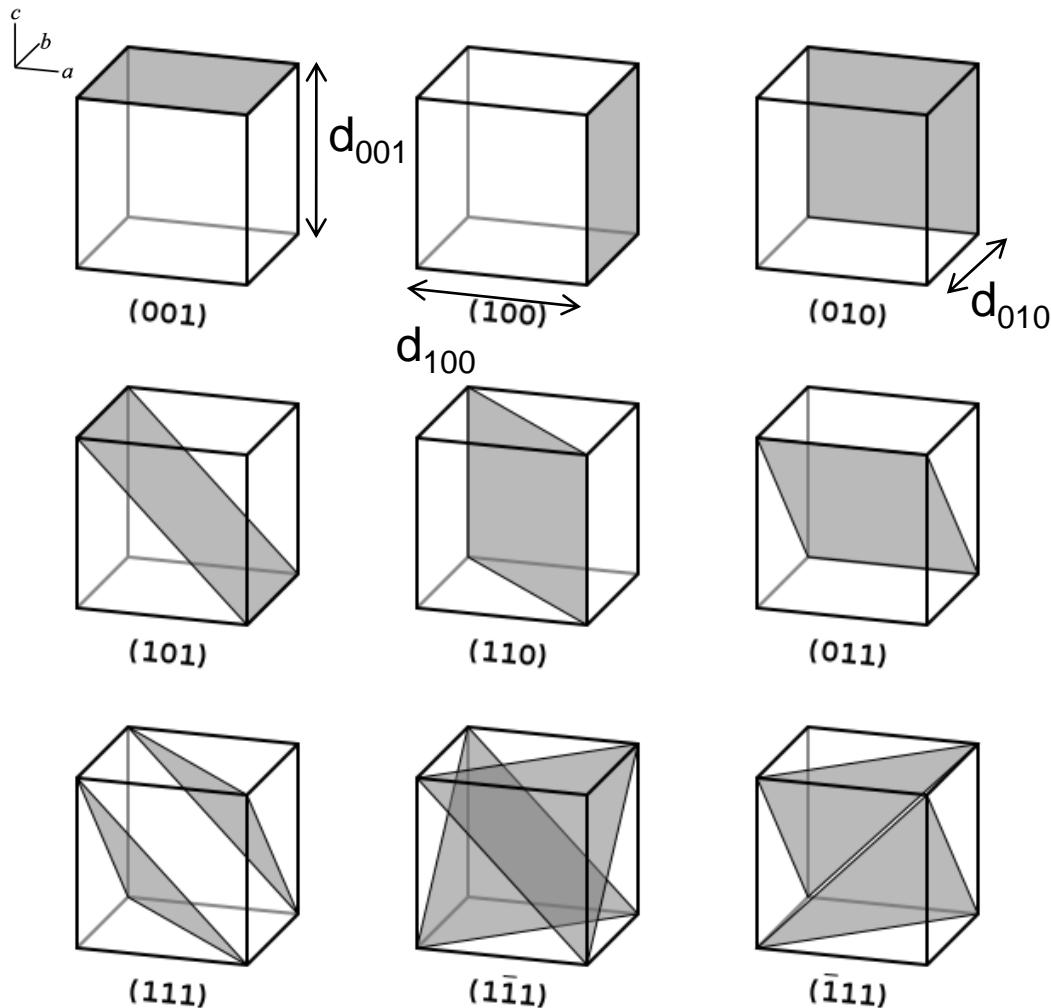
Unit cell planes can be defined by a notation called a Miller index (hkl).

To obtain the Miller indices of a given plane requires the following steps:

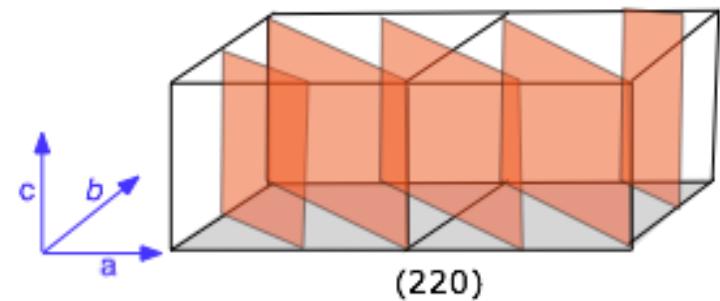
1. The plane in question is placed on a unit cell.
2. Find its intercepts with each of the crystal axes.
3. The reciprocal of the intercepts are taken.
4. Multiply by a scalar to get a ratio of integers.



Key concept 2: Miller indices



The higher the Miller index the less distance there is between equivalent planes, dividing the unit cell into ever smaller slices



For higher symmetry cells interplane distances can be identical
 $d_{001} = d_{010} = d_{100}$ for cubic

Key concept 2: Miller planes

d-spacings in different crystal systems

Crystal system d_{hkl} as a function of Miller indices and lattice parameters

Cubic
$$\frac{1}{d^2} = \frac{h^2 + k^2 + l^2}{a^2}$$

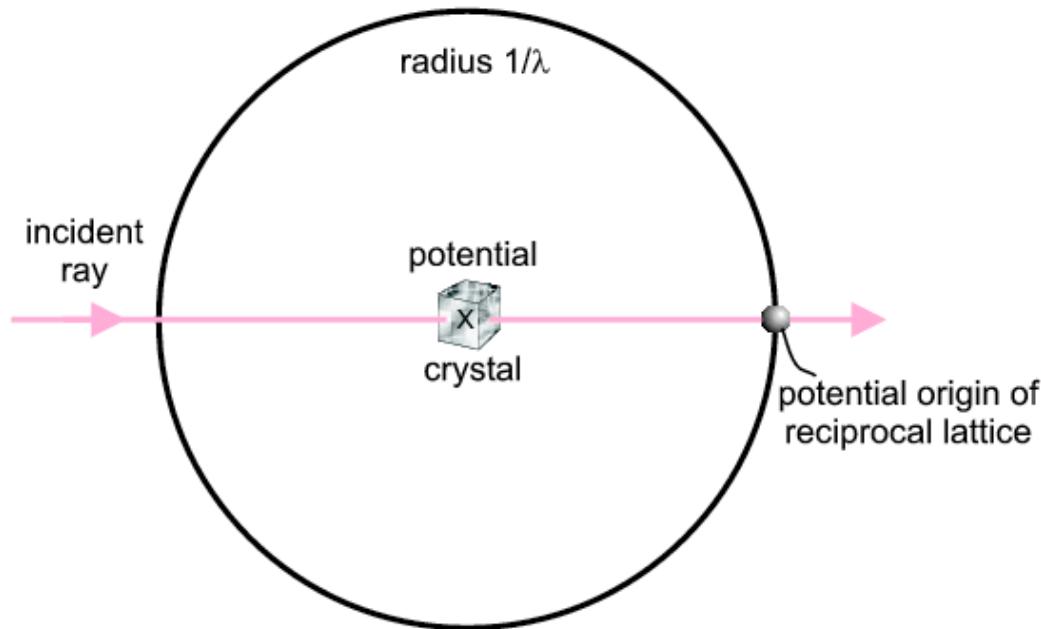
Tetragonal
$$\frac{1}{d^2} = \frac{h^2 + k^2}{a^2} + \frac{l^2}{c^2}$$

Orthorhombic
$$\frac{1}{d^2} = \frac{h^2}{a^2} + \frac{k^2}{b^2} + \frac{l^2}{c^2}$$

Hexagonal
$$\frac{1}{d^2} = \frac{4}{3} \left(\frac{h^2 + hk + k^2}{a^2} \right) + \frac{l^2}{c^2}$$

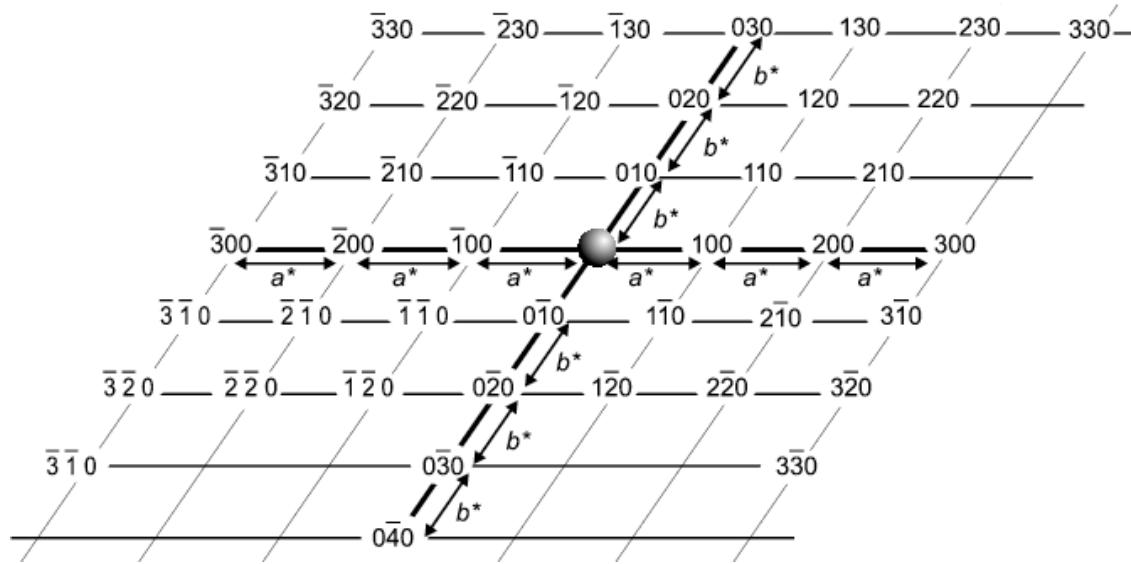
Monoclinic
$$\frac{1}{d^2} = \frac{1}{\sin^2 \beta} \left(\frac{h^2}{a^2} + \frac{k^2 \sin^2 \beta}{b^2} + \frac{l^2}{c^2} - \frac{2hl \cos \beta}{ac} \right)$$

Key concept 3: Ewald sphere



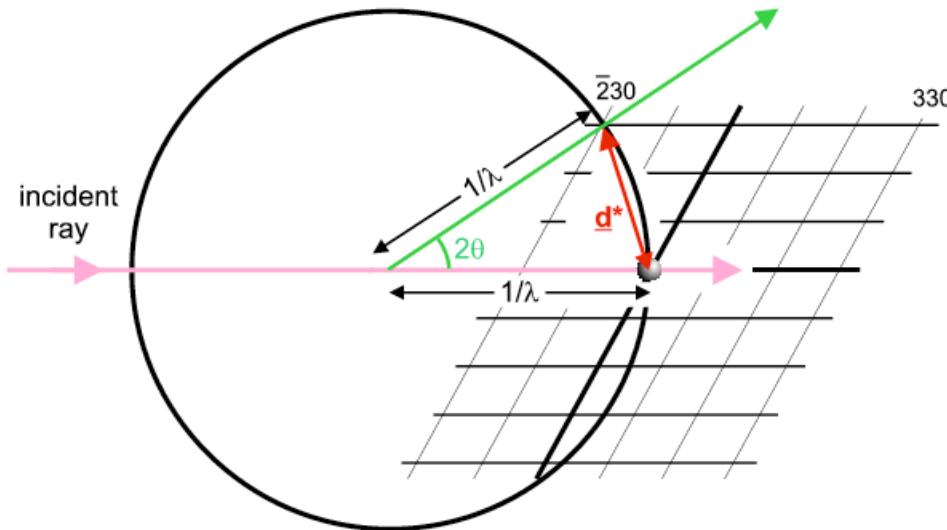
- A sphere of radius $1/\lambda$ (2-D projection shown above)
- Diffracted X-rays/neutrons can be along any radius from the centre of the sphere to the circumference

Key concept 4: Reciprocal lattice



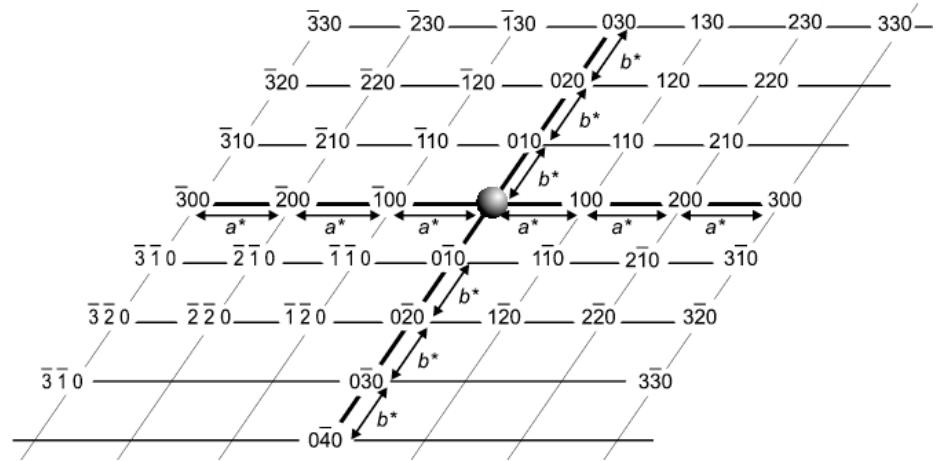
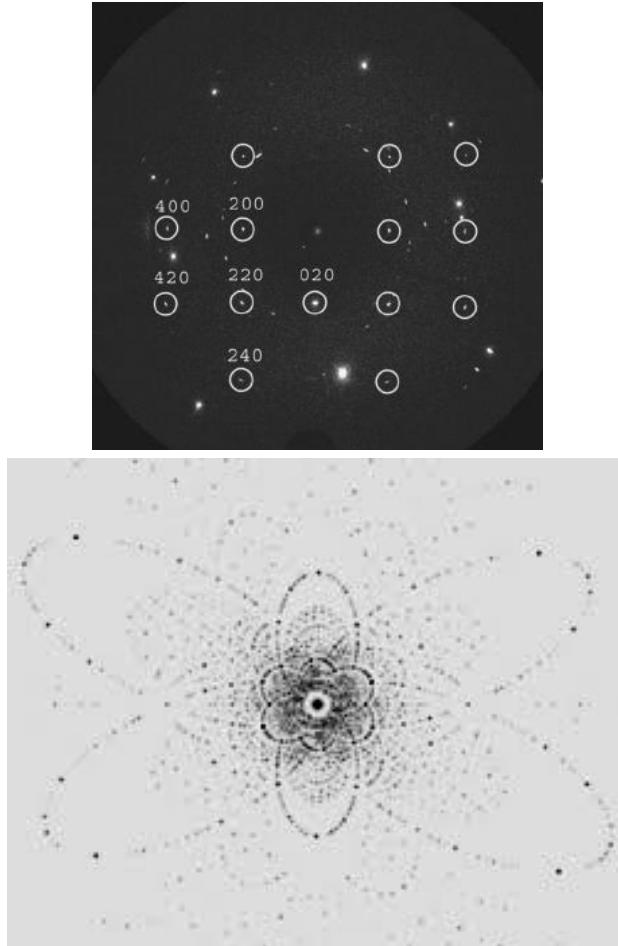
- The reciprocal lattice consists of points which represent diffraction possibilities
- Each point can be labeled with a Miller index
- The units are a^* , b^* and c^* and any point can be reached using the vector equation $d^* = ha^* + kb^* + lc^*$

Key concept 5: Bragg diffraction



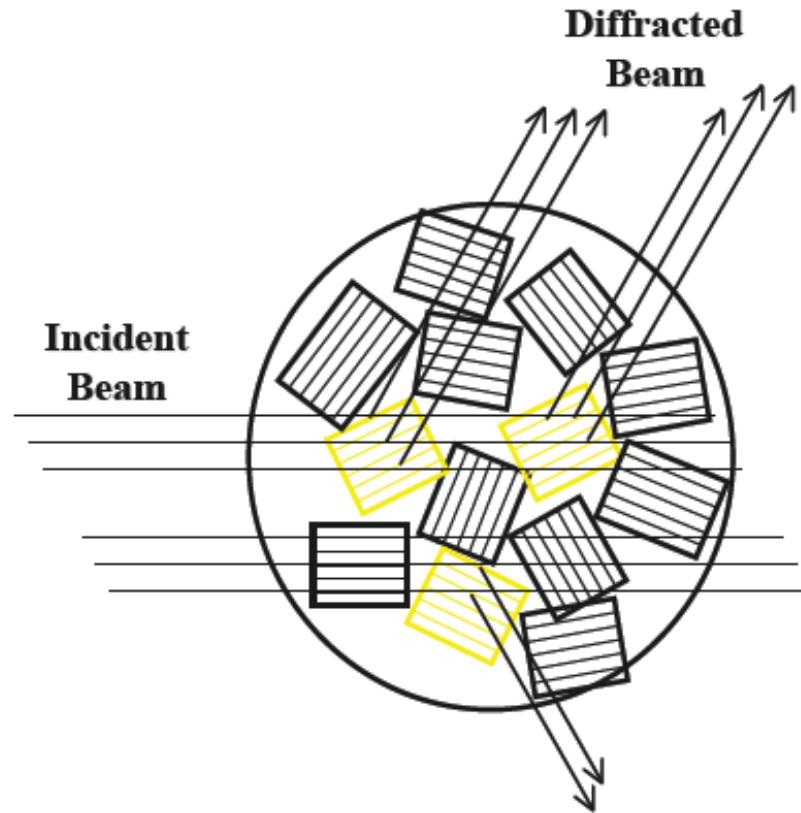
- Diffraction observed when a reciprocal lattice point intersects Ewald sphere
- Crystal rotation brings other lattice points into contact with Ewald sphere
- The vector from origin to lattice point is d^* (reciprocal lattice spacing) is red – it is exactly equal to $1/d$ and its direction is perpendicular to the hkl plane
- The direction of the diffracted ray is indicated in green

Single Crystal Diffraction

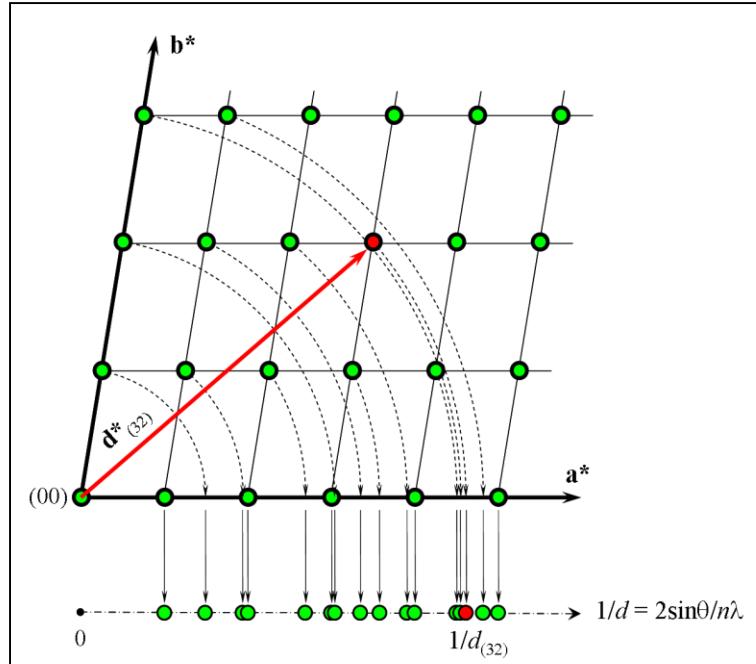


Direct observation of the reciprocal lattice

Powder diffraction

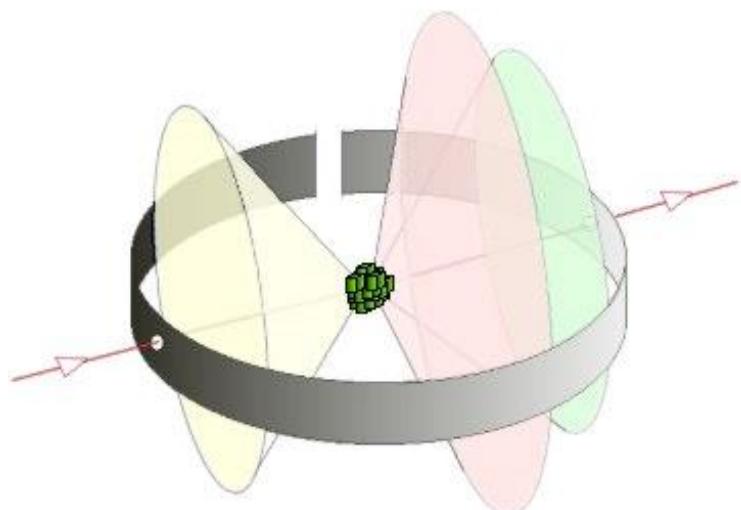
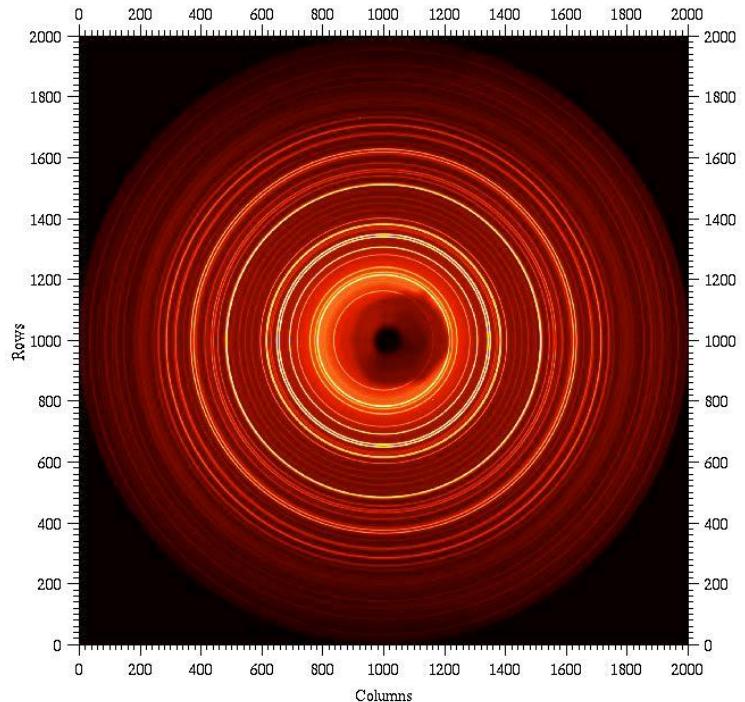


Powder diffraction

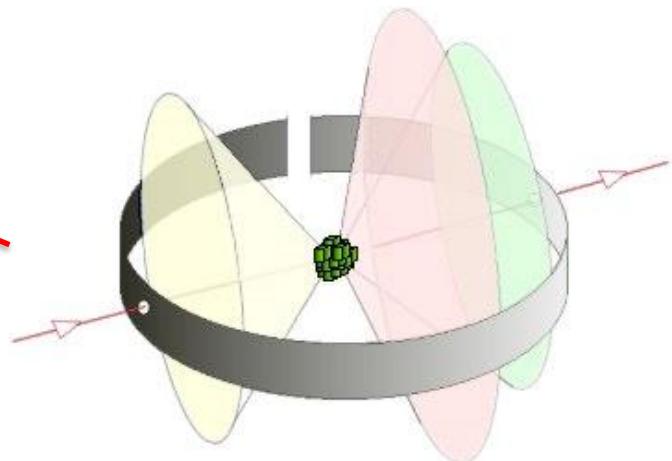
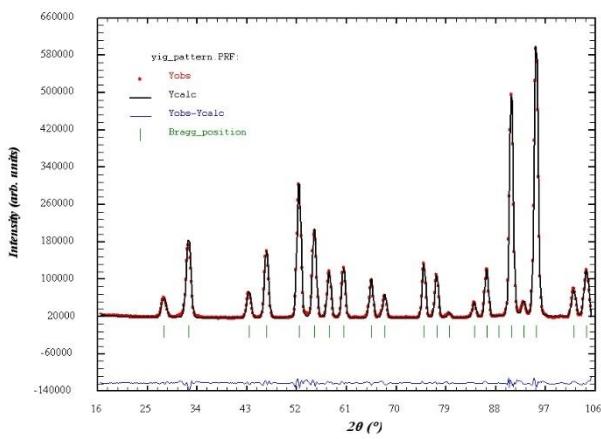
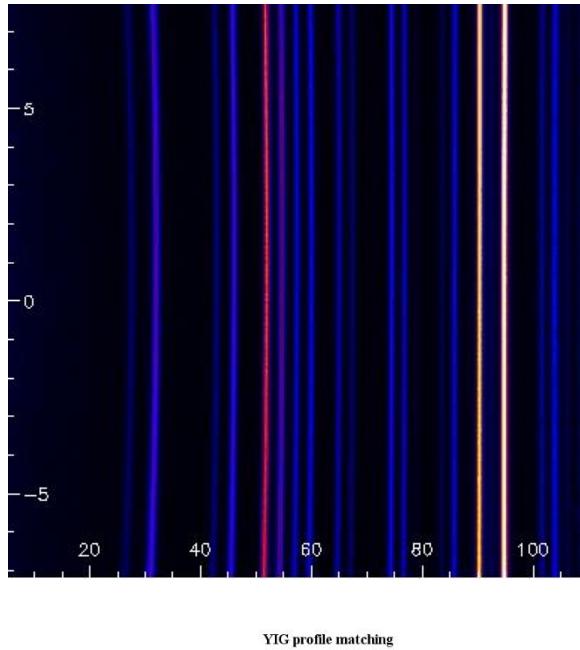


- Many crystallites with random orientation mean that each reciprocal lattice point will occur in every orientation possible, broadening into the surface of a sphere with radius d^*
- The intersection of the Ewald sphere and the reciprocal lattice becomes a cone (intersection of 2 spheres)
- The directions of the vectors are lost and only the lengths of the reciprocal lattice vectors are measurable with powder diffractometers
- 3-D information collapsed into 1-D

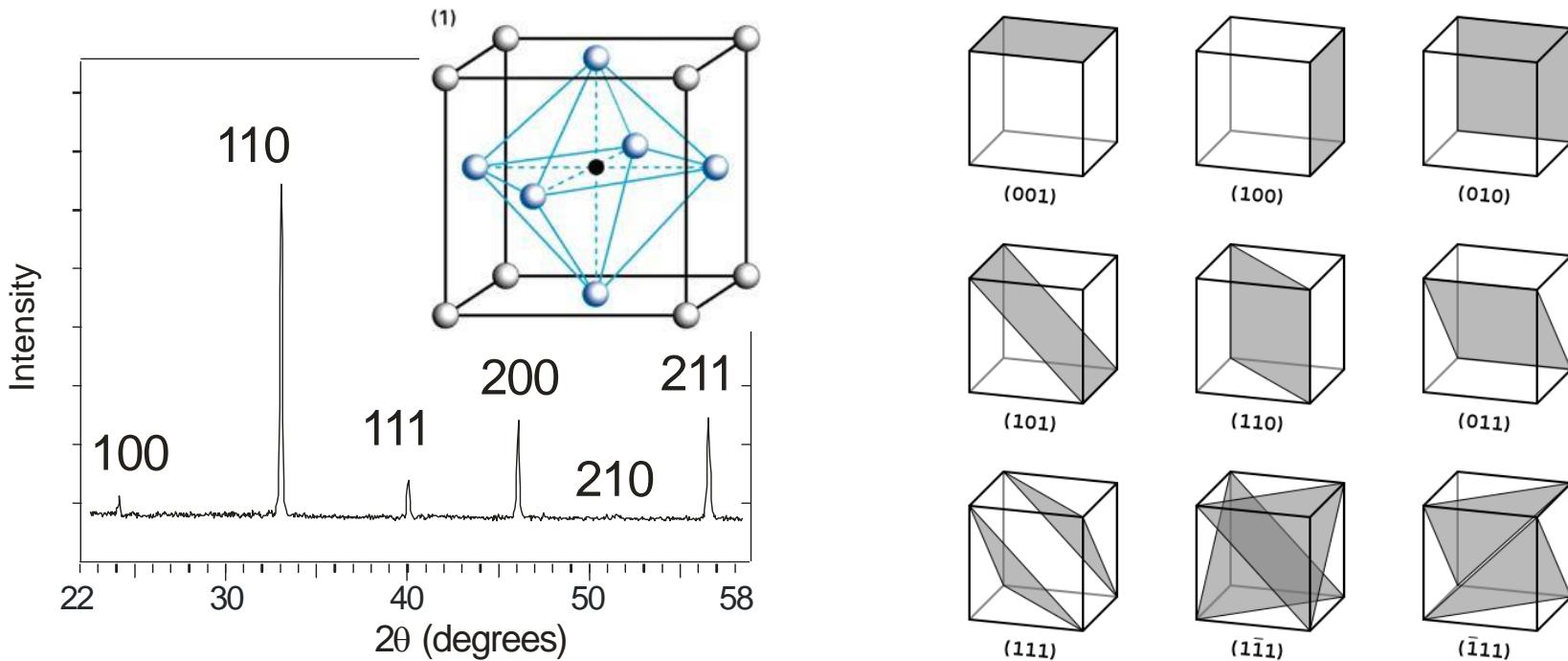
Powder diffraction



Powder diffraction



Miller plane equivalence in powder diffraction

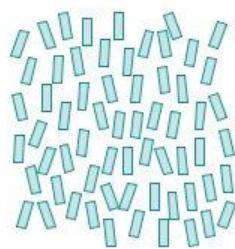
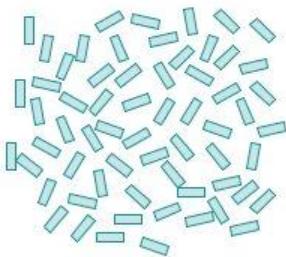
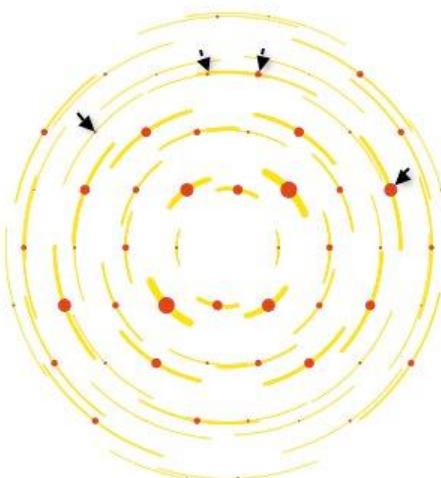
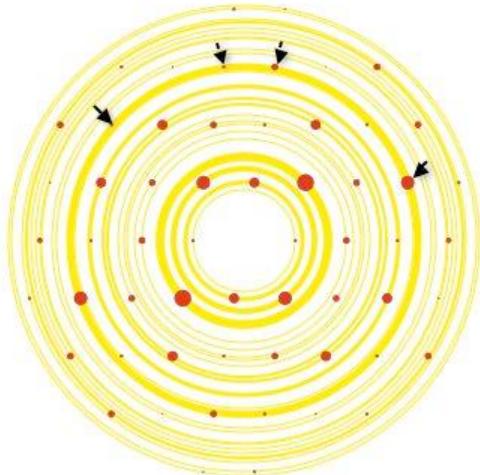


All equivalent planes occur at same scattering angle

All planes separated by the same distance occur at one scattering angle in powder diffraction

e.g. (511) and (333) occur at same 2θ for a cubic material

Not enough crystallites or a non-powder average



When number of crystals is too small, the pattern becomes “grainy”
-- diffraction from individual crystals dominate.

- Increase sample size
- Grind the sample to decrease domain size
- Oscillate or rotate the sample
- Use area detection & integrate the entire Debye-Scherrer ring

Intensity and structure factor

$$I_{hkl} \propto |F_{hkl}|^2$$

Measured intensity proportional to F_{hkl}^2 and so we cannot tell whether F_{hkl} is positive or negative – the Phase problem

$$F_{hkl} \propto \sum f_i \exp[2\pi i(hx_i + ky_i + lz_i)] \exp(-U_i Q^2/2)$$

f_i is the scattering power (form factor of the ith site i.e. (x_i, y_i, z_i) and includes fractional occupancy

Contribution of the ith site to the F_{hkl} in question

Atomic displacement of the ith atom site

The phase problem

$$I_{hkl} \propto |F_{hkl}|^2$$

In diffraction we measure the magnitudes and not the phase. The phases contain the bulk of the information. This is why crystallography is hard....

...but not impossible. We can recover phase information from:

- Related or isostructural materials
- Knowledge of atom positions (heavy atoms from X-rays)
- Known motifs (molecules)
- Brute force

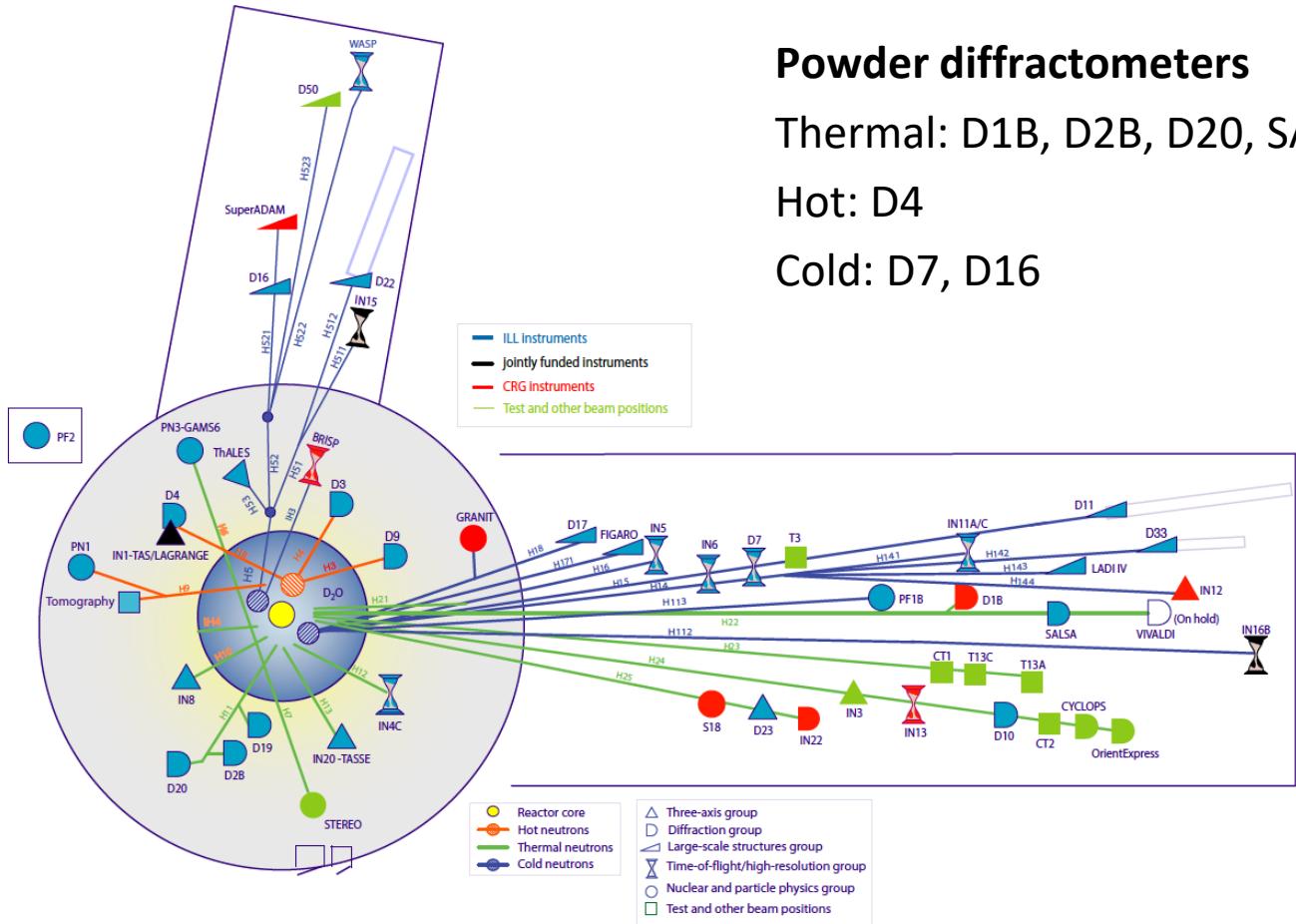
Instrumentation for neutron diffraction



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Diffraction at a continuous source: ILL



Powder diffractometers

Thermal: D1B, D2B, D20, SALSA

Hot: D4

Cold: D7, D16

Single crystal diffractometers

Thermal: D19, D10, CYCLOPS, VIVALDI, D23, OrientExpress

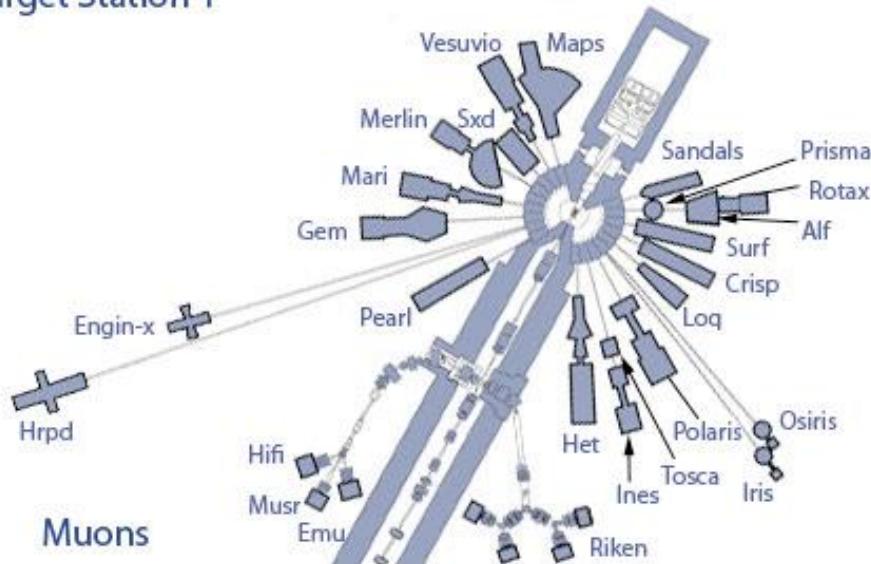
Hot: D3, D9

Cold: LADI-III, D7

Half of instrument suite are diffractometers

Diffraction at a pulsed source: ISIS

Target Station 1



HRPD

ENGIN-X

GEM

SXD

INES

PEARL

POLARIS

WISH

IMAT

(LMX)

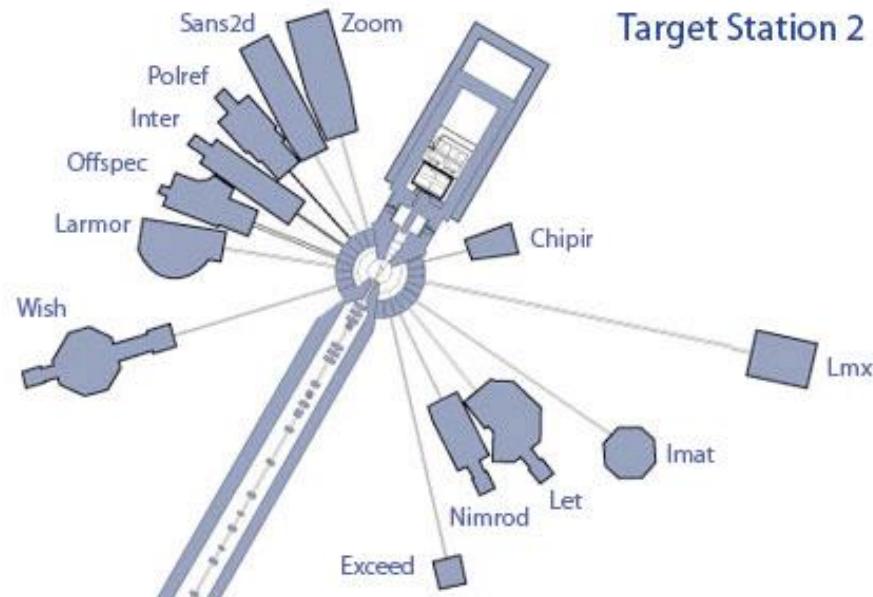
(EXCEED)

SANDALS

NIMROD

OSIRIS

Target Station 2



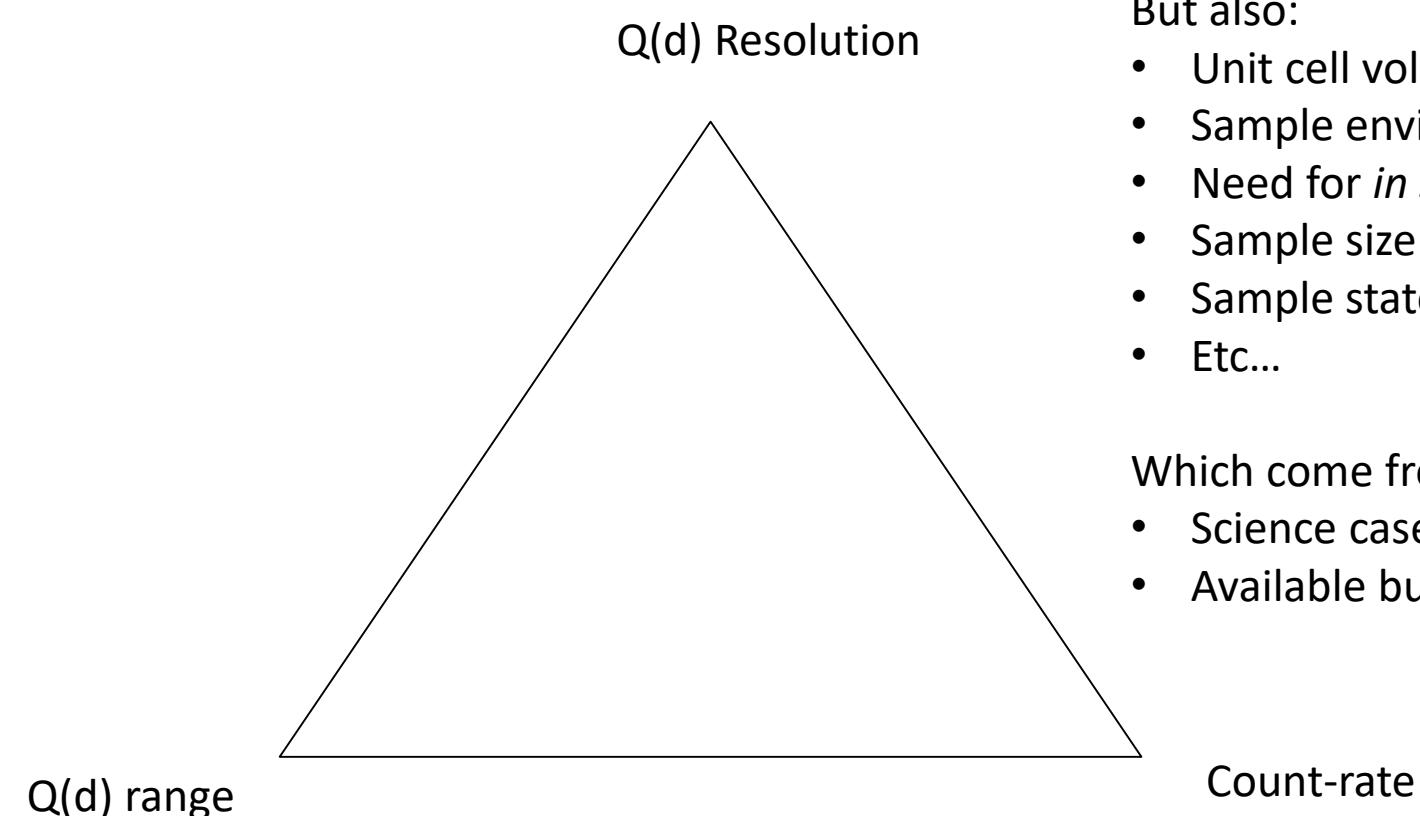
Almost half of the instrument suite are diffractometers or carry significant diffraction capability



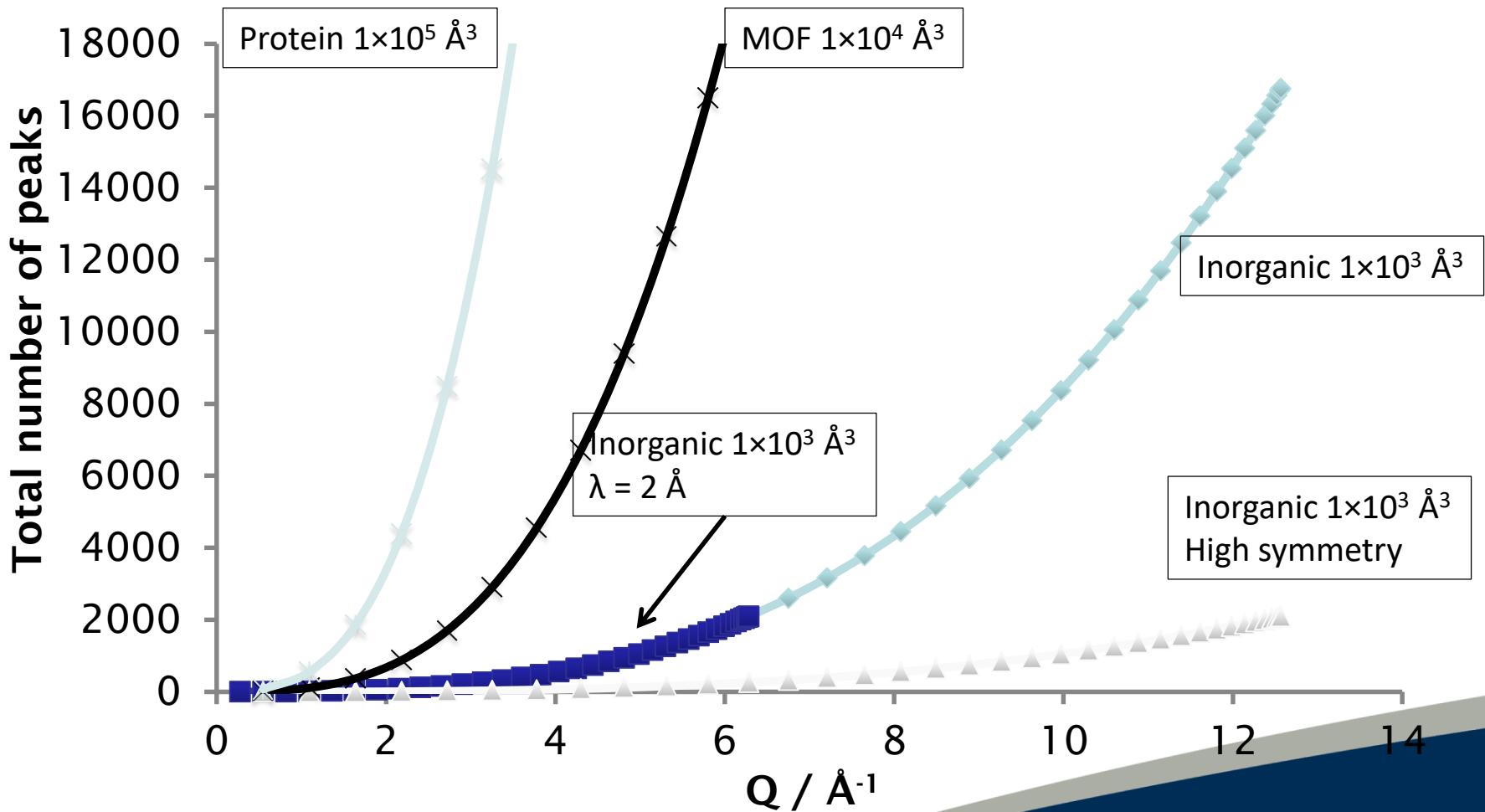
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Why so many diffractometers?



Number of possible reflections



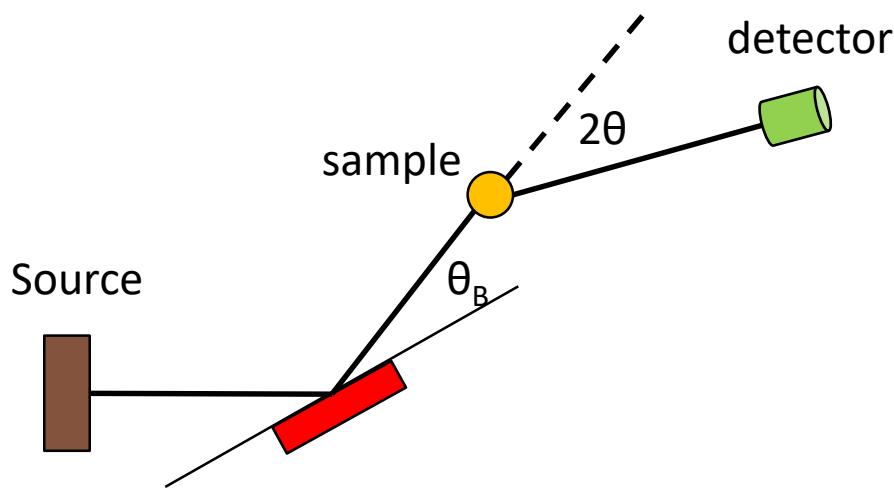
Types of diffractometer

$$\lambda = 2d\sin\theta$$

- Monochromatic (CW)
 - Fix wavelength and scan detector angle
 - Multiple 2θ required to cover $Q(d)$ spacing range
 - $Q(d)$ spacing limit $4\pi/\lambda$ ($2\pi/d$)
 - Instrumental count rate factors: Source power, monochromator reflectivity, detector coverage and efficiency, etc
- TOF
 - Fix detector angle and scan wavelength
 - Single 2θ covers range of $Q(d)$ space
 - $Q(d$ range) determined by λ_{\max} , λ_{\min} and θ
 - Instrumental count rate factors: Source power, moderator performance, beam transport efficiency, detector coverage and efficiency, etc

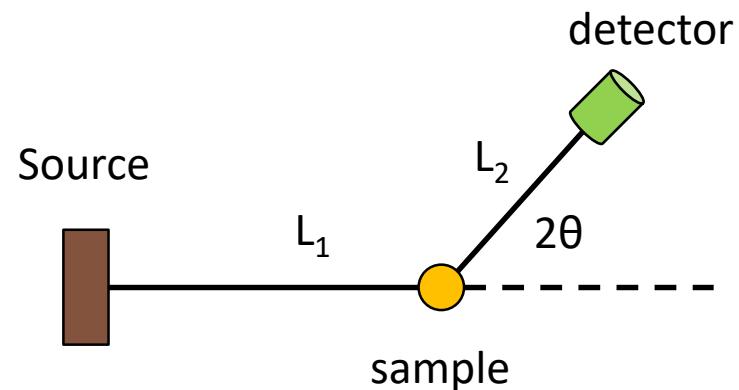
Instrument layouts

CW



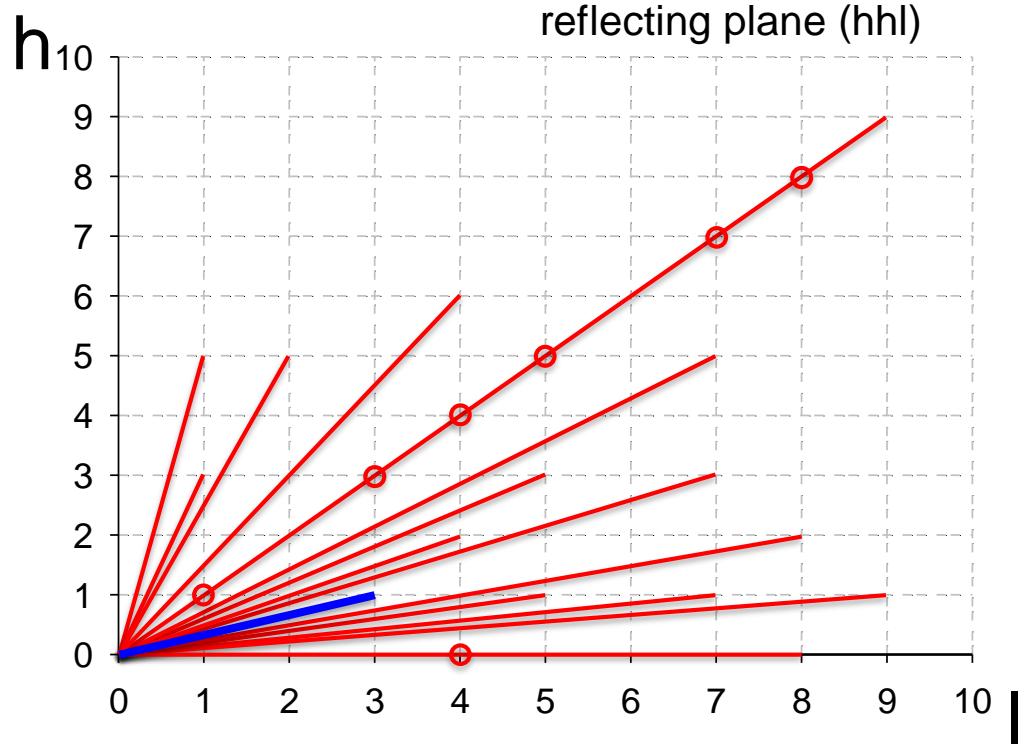
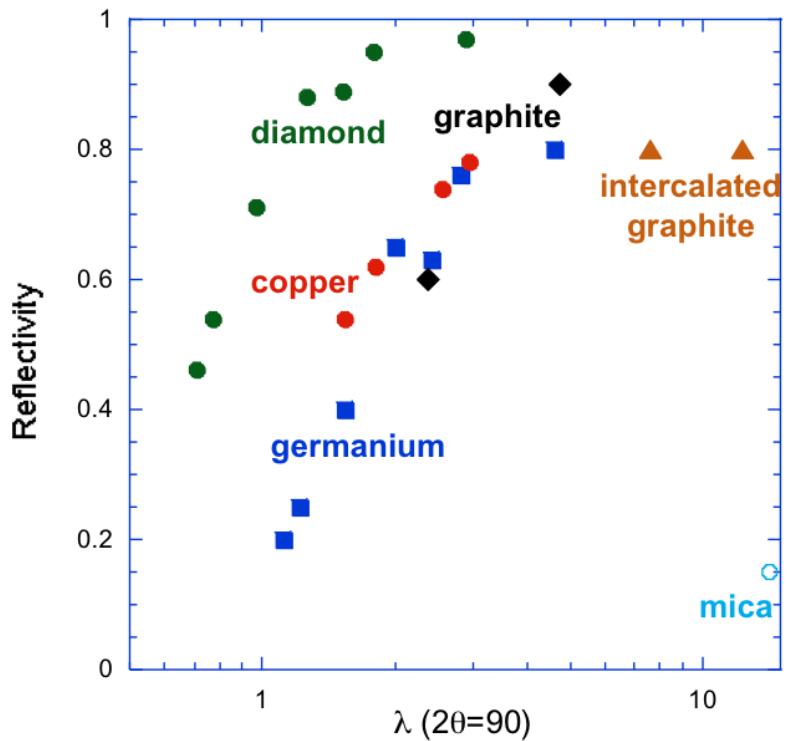
$$\lambda = \frac{2d_c \sin \theta_B}{n}$$

TOF



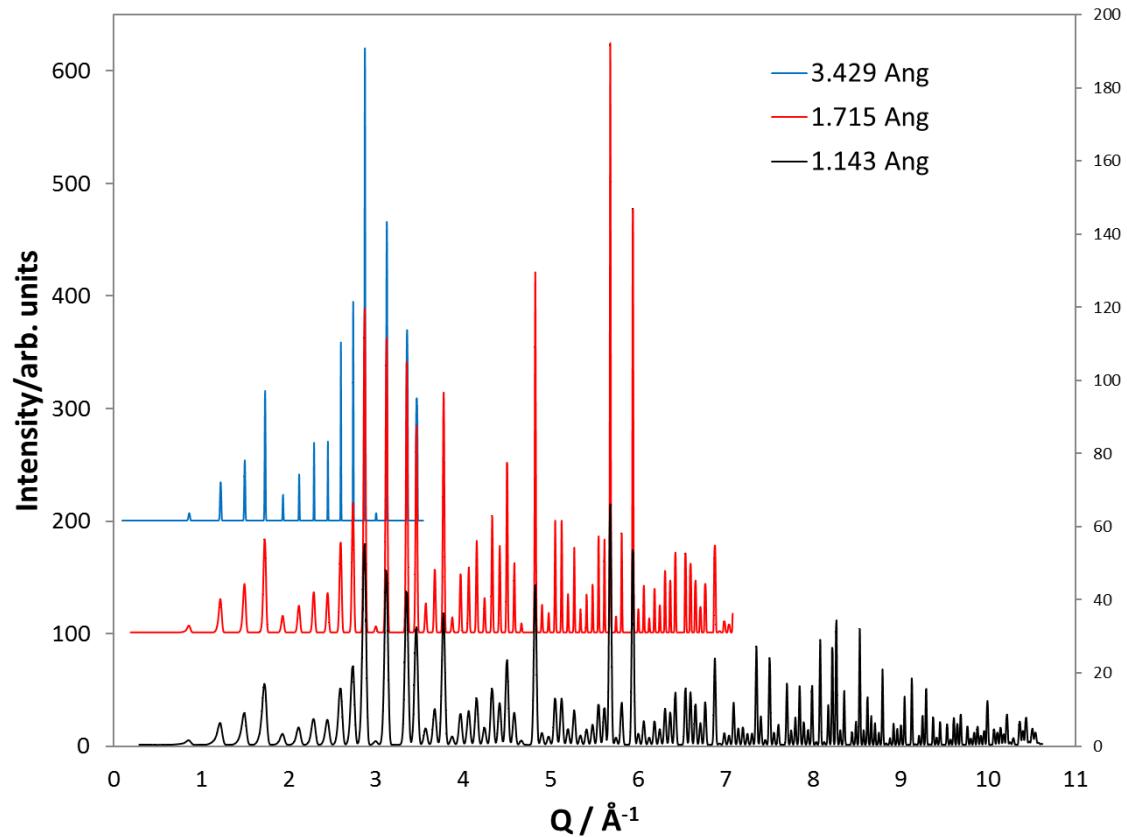
$$\lambda = \frac{3956}{v} = \frac{3956(t-t_0)}{L_1+L_2}$$

Monochromator materials



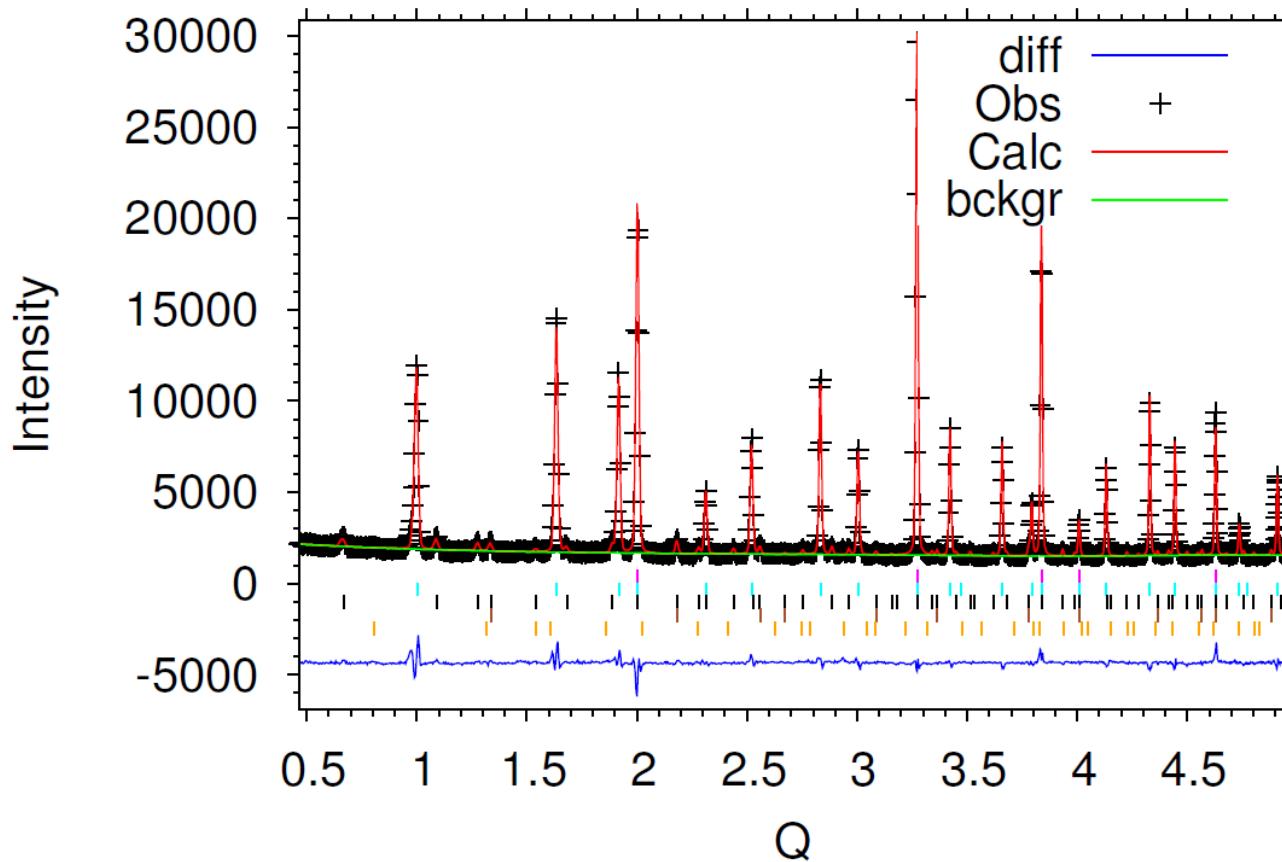
Choose material and plane to access necessary Q-range
Reflectivity falls as wavelength decreases

Q range with monochromators



Shorter wavelengths access higher Q but have lower reflectivity

Higher reflection order contamination

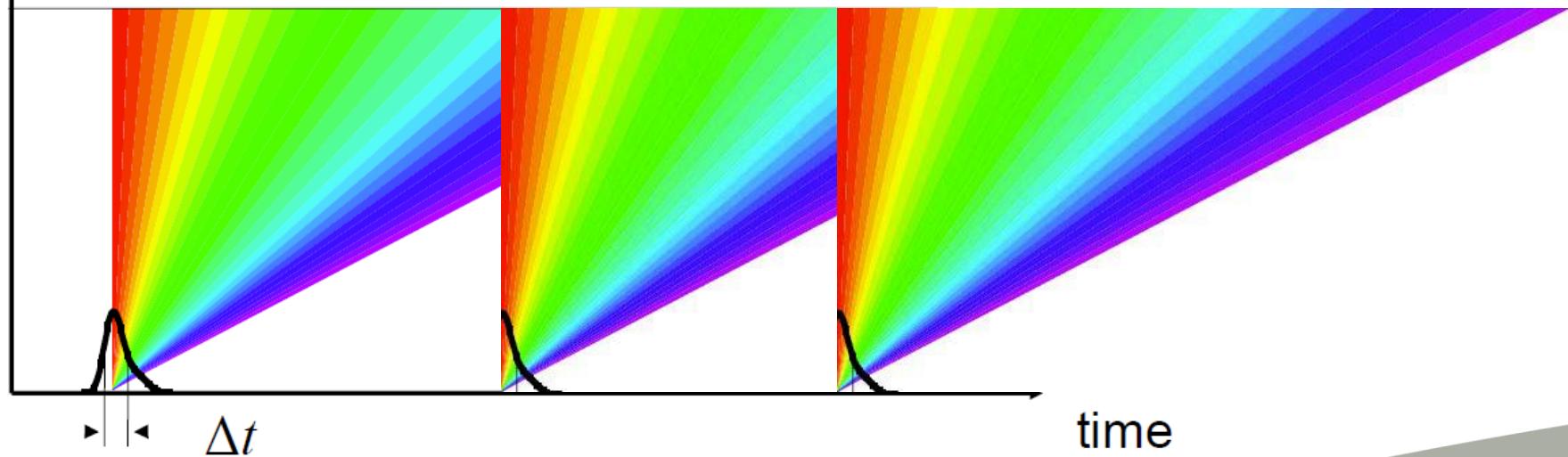


High reflection order contamination complicates analysis with CW data

TOF method

distance

- Use distance to separate wavelengths
- Need choppers to prevent frame overlap
- Moderator pulse-width and distance determine resolution
- Source repetition rate determines available time-window
- Beam transport crucial



CW or TOF: Q-range summary

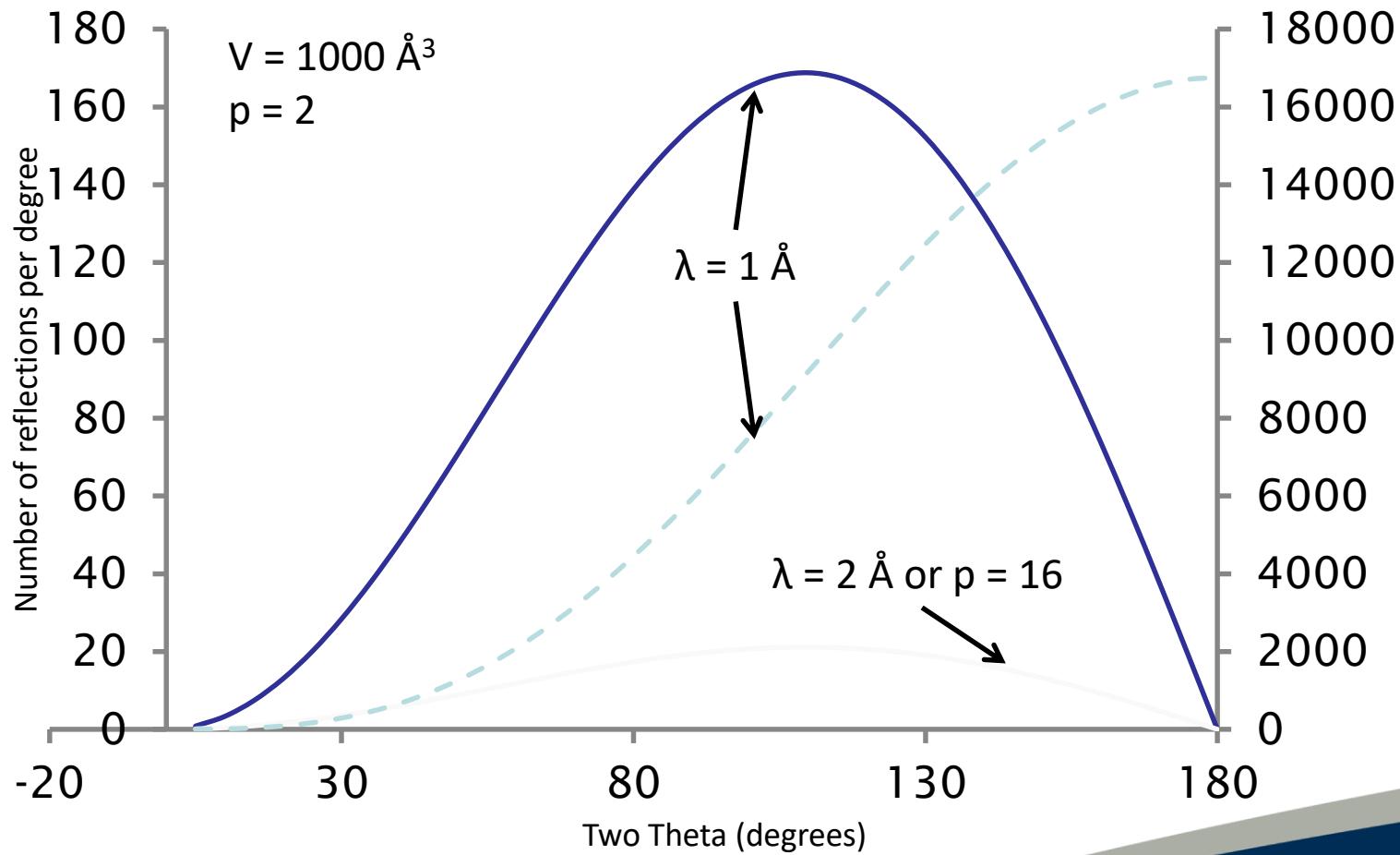
For CW:

- For monochromatic instruments the Q_{\max} is $4\pi/\lambda$ i.e. when $\sin\theta = 1$, $\theta = 90^\circ$, $2\theta = 180^\circ$
- If a high Q_{\max} is required a shorter wavelength must be used.
- Shorter wavelengths are produced by higher order hkl planes
- Reflectivity is lower for shorter wavelengths
- Realistic Q_{\max} of around 25 \AA^{-1}

For TOF:

- Q_{\max} depends on λ_{\min} and detector θ .
- λ_{\min} can be much lower than for the CW case allowing $Q_{\max} > 100 \text{ \AA}^{-1}$
- λ_{\min} determined by the moderator, transport characteristics of the guide and which frame the instrument is working in

Reflection density for CW

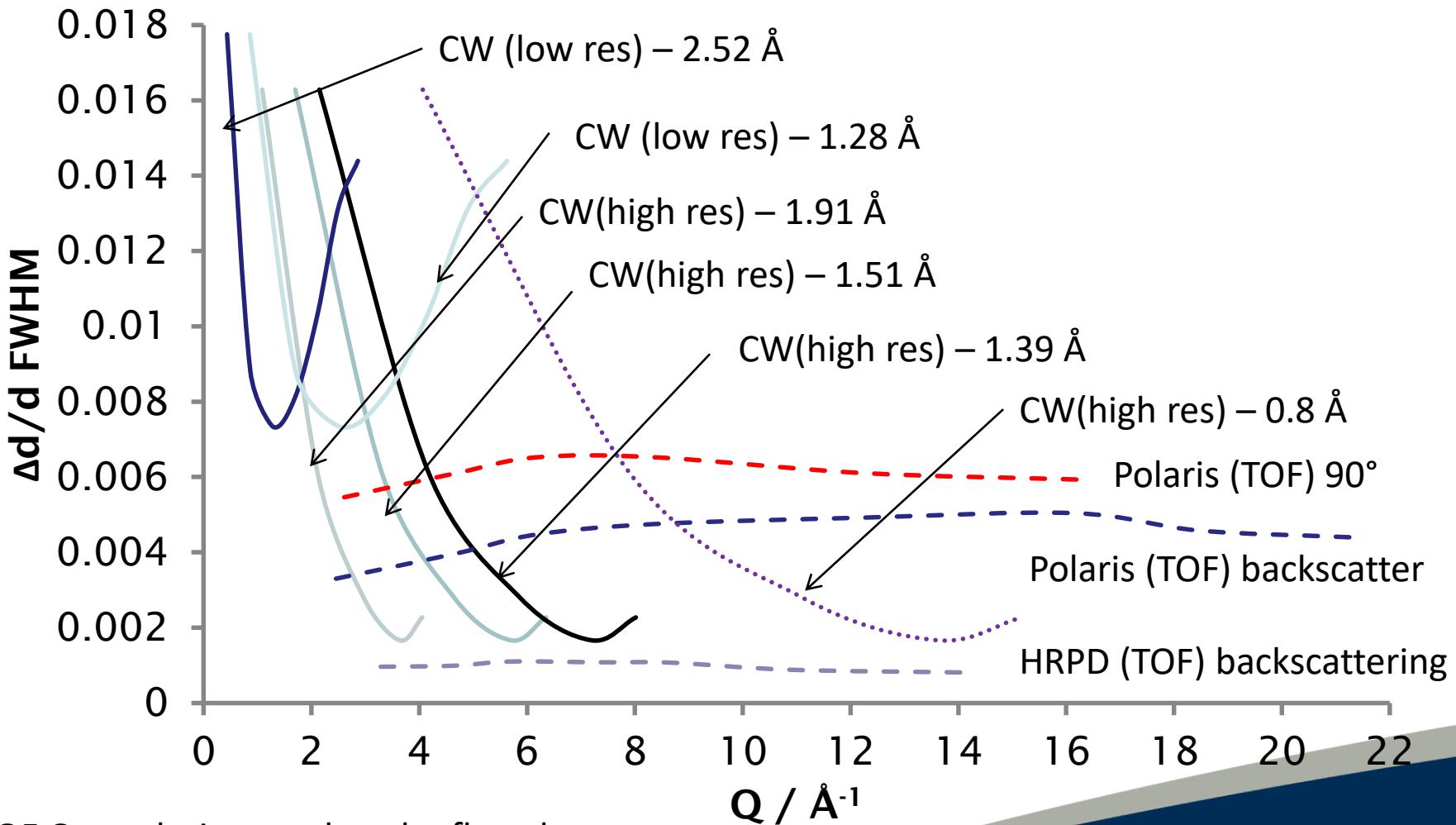


CW instruments designed to have best resolution at highest peak density

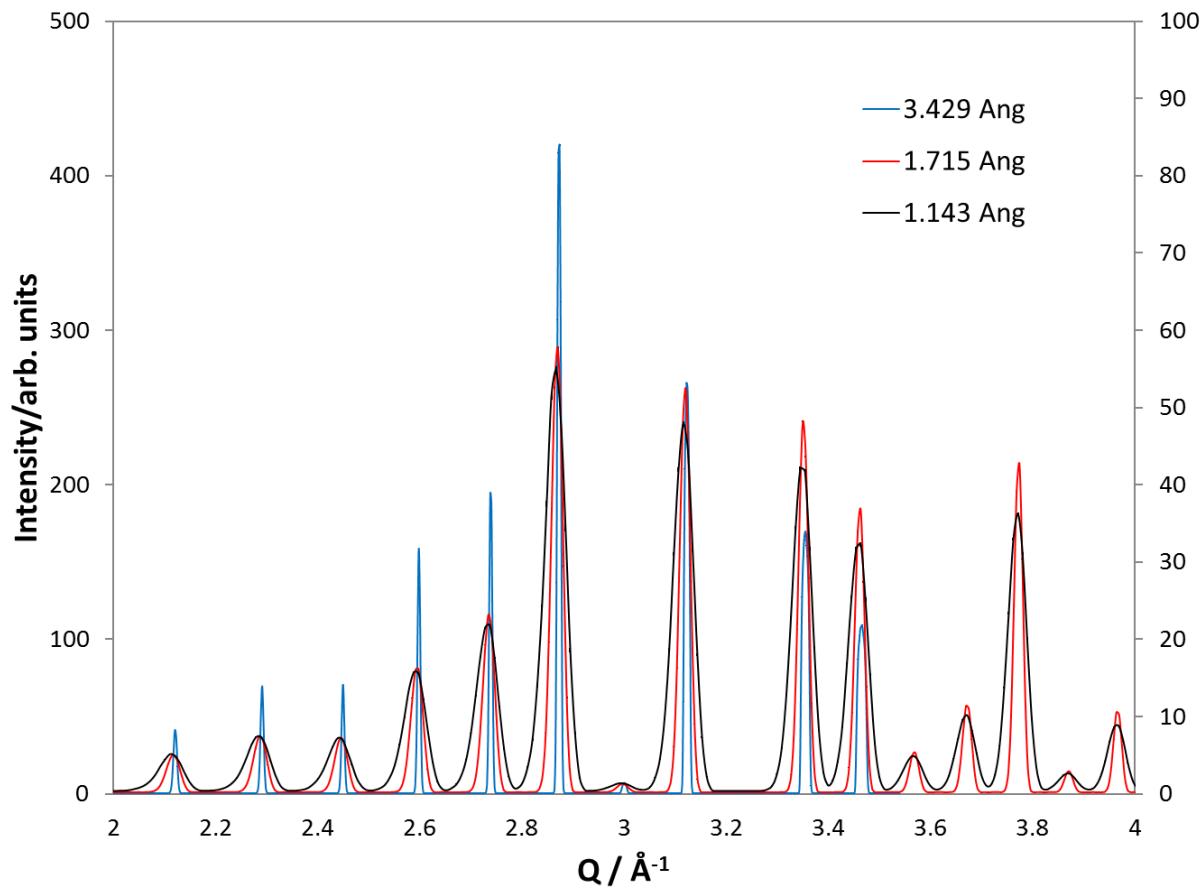


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Resolution functions CW v TOF



CW Q resolution example



Choose wavelength to match Q resolution required
by science in a given Q range



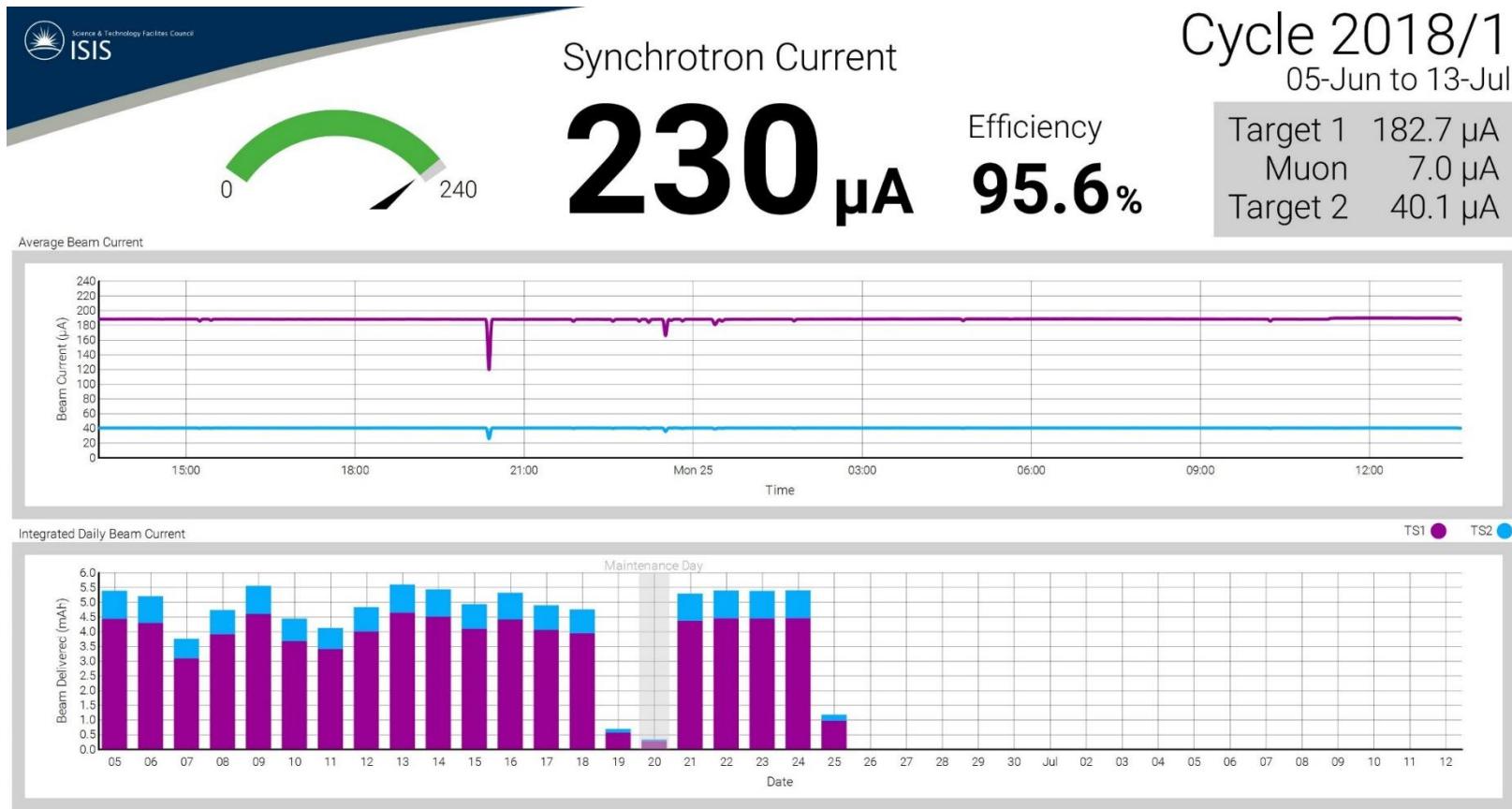
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CW or TOF: Q resolution

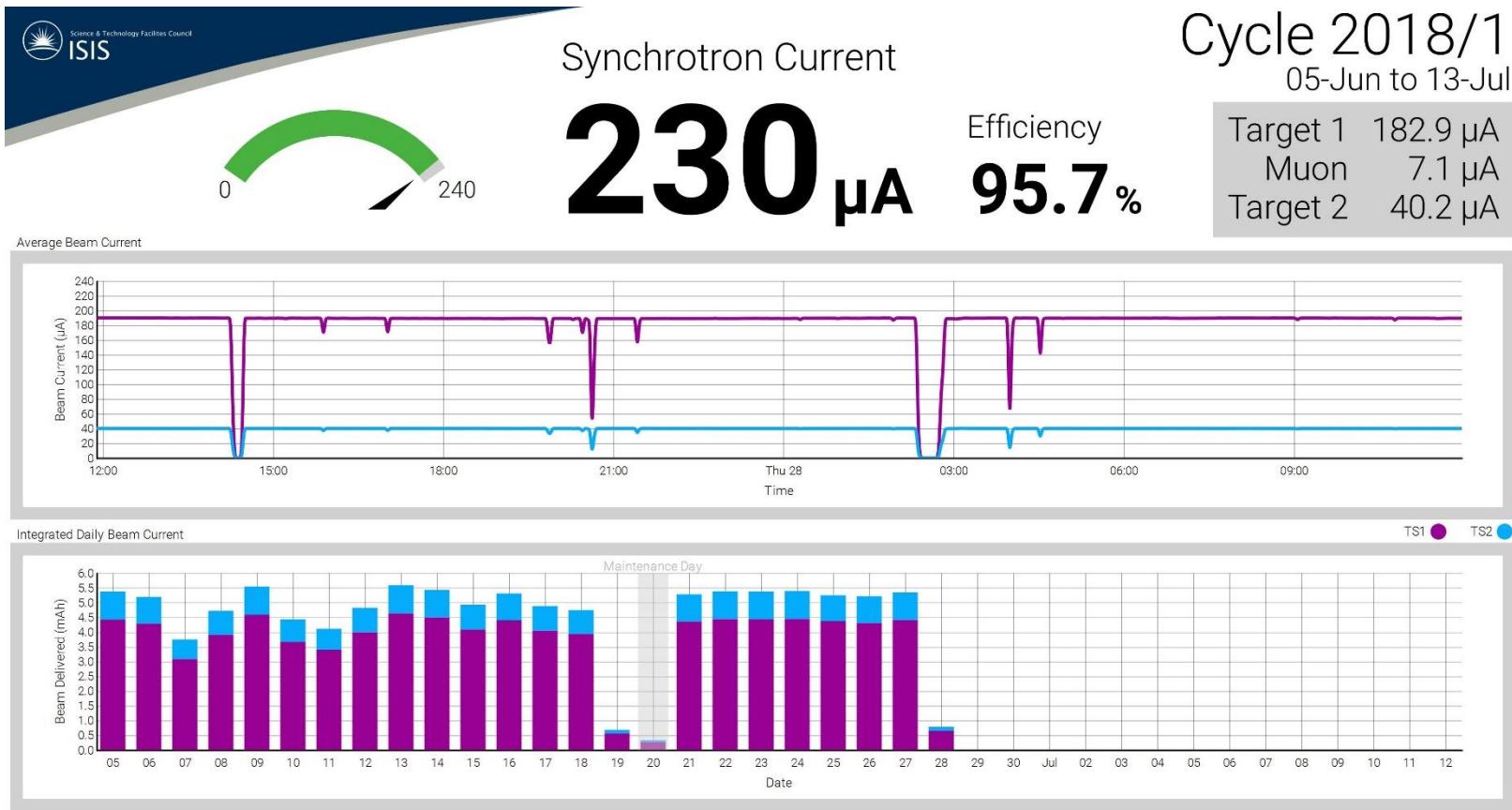
- CW:
 - Simple, symmetric peakshape function
 - Best resolution where diffraction peak density is highest in scattering angle
 - Different wavelength can be used to give Q resolution where required
 - Different takeoff angle can be used to change resolution function and wavelength
 - Instrument can be high Q resolution but with very limited Q range
- TOF:
 - Complex asymmetric peakshape related to moderator characteristics
 - Instrument length and moderator give wavelength band and overall resolution
 - Q resolution almost constant for a given detector bank so increasing peak density with Q can be an issue
 - Q resolution improved by moving to higher scattering angle detector bank
 - Q range determined by scattering angle of detector bank

Pulsed source availability



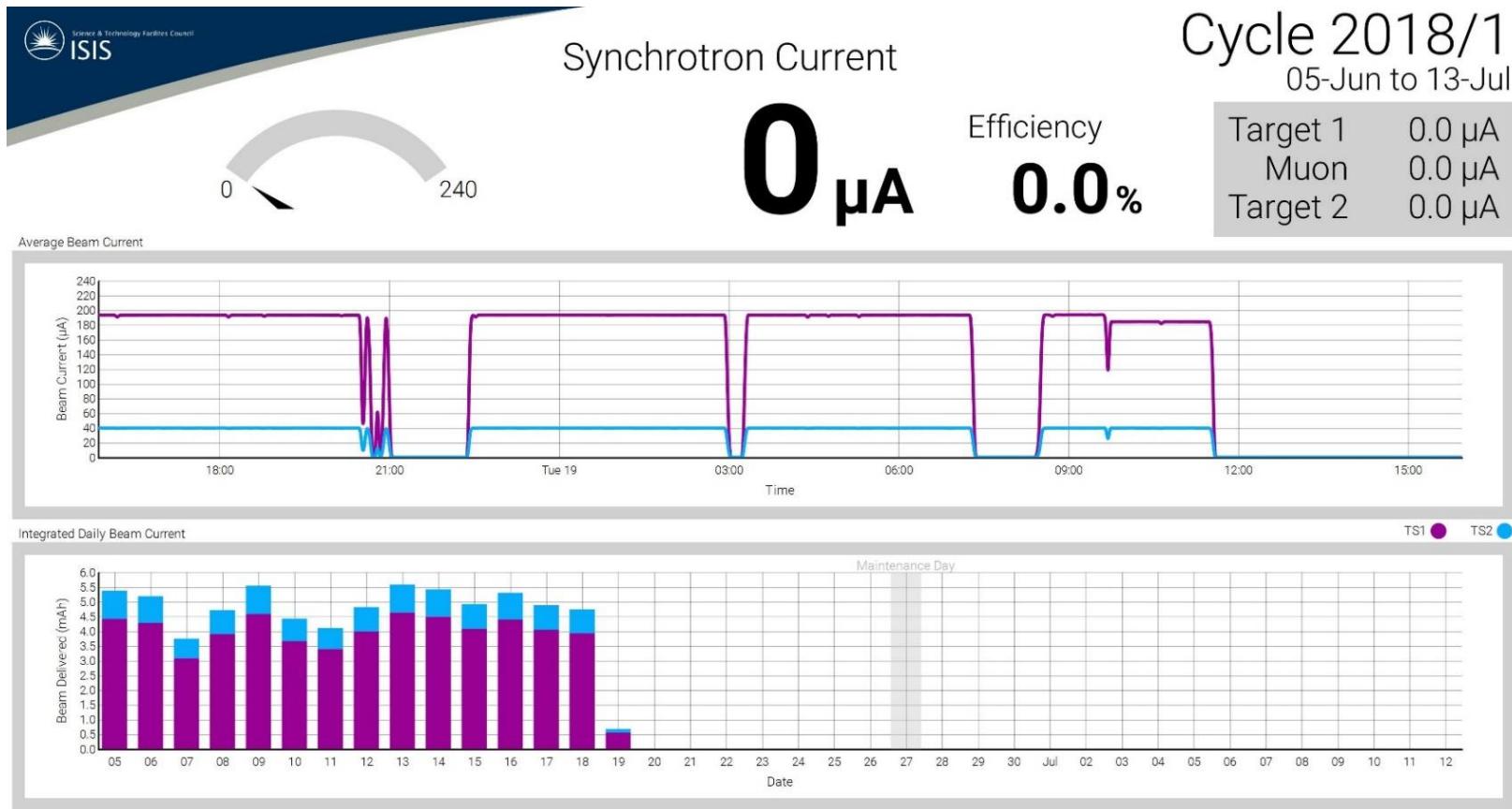
A very good day in terms of beam
All experiment types possible

Pulsed source availability



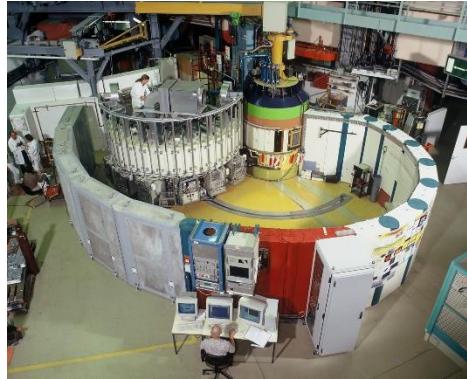
A good day in terms of beam – still possible issues with *in situ* and time resolved experiments

Pulsed source availability

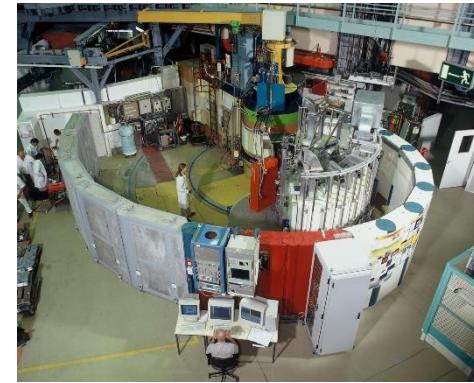
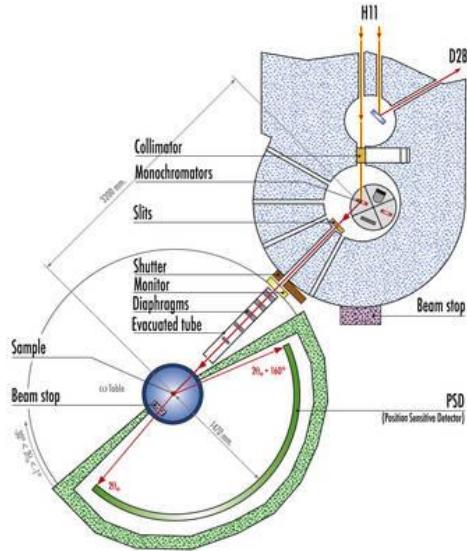


A bad day – any time resolved experiment is compromised

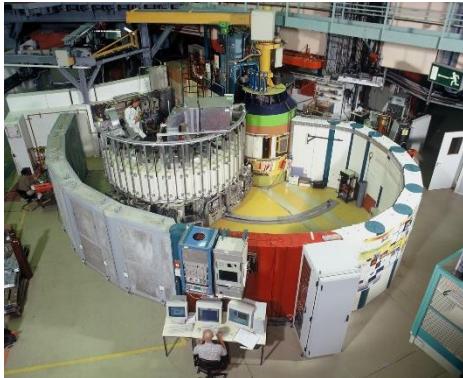
CW, variable resolution, thermal powder diffractometer: D20



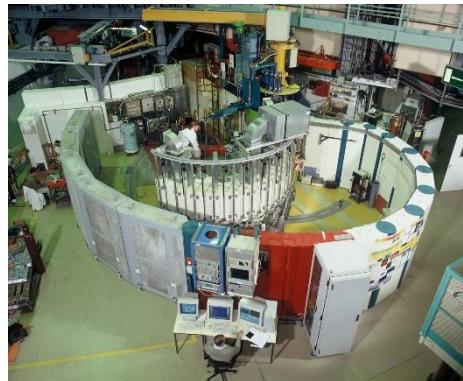
120°



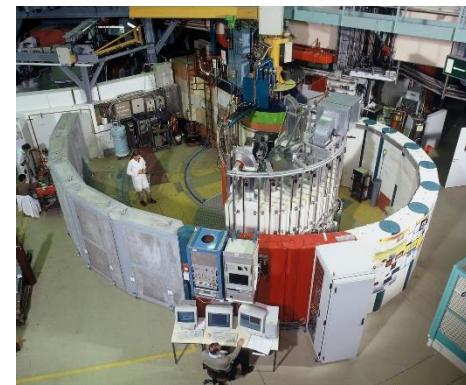
28°



90°



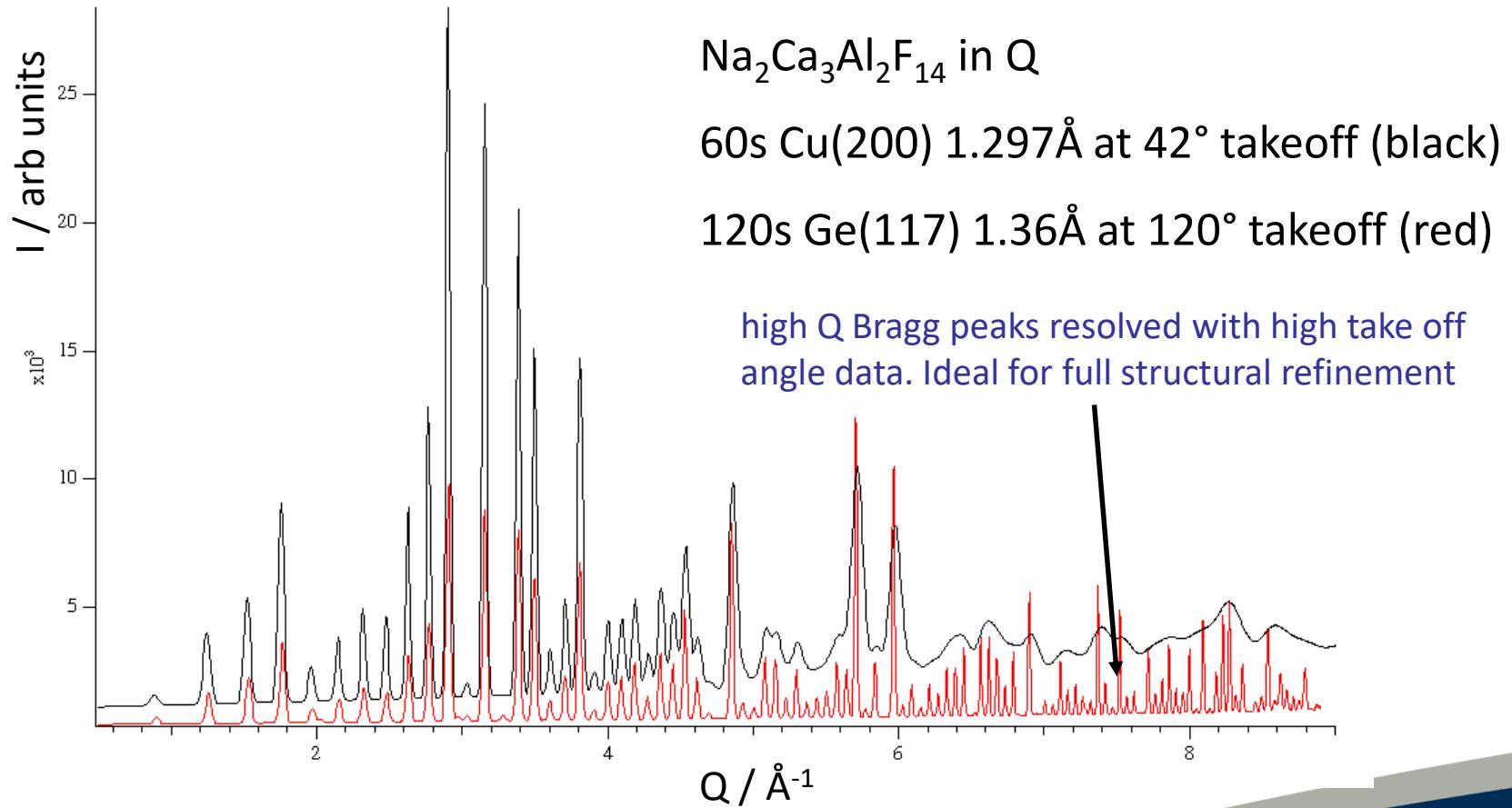
65°



42°

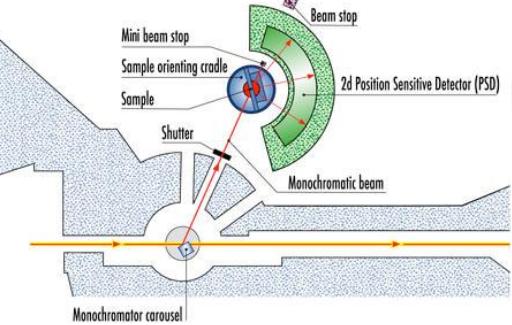
Photos courtesy T. Hansen ILL

Low θ_B v high θ_B : Q resolution v count-rate

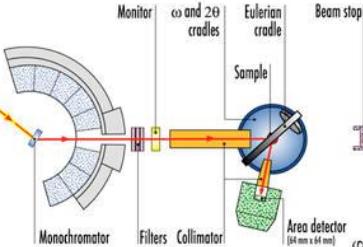


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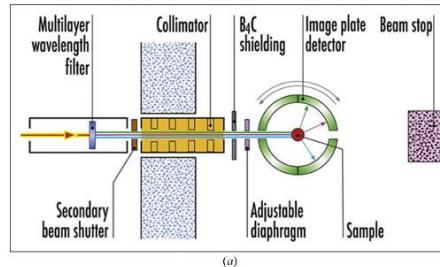
CW single crystal diffraction



D19



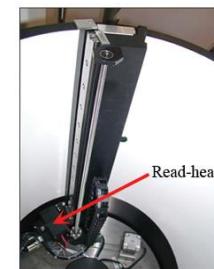
D9



(a)



(c)

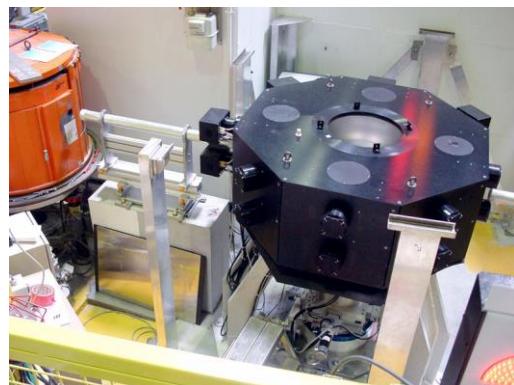


(d)



↑
LADI-III

CYCLOPS

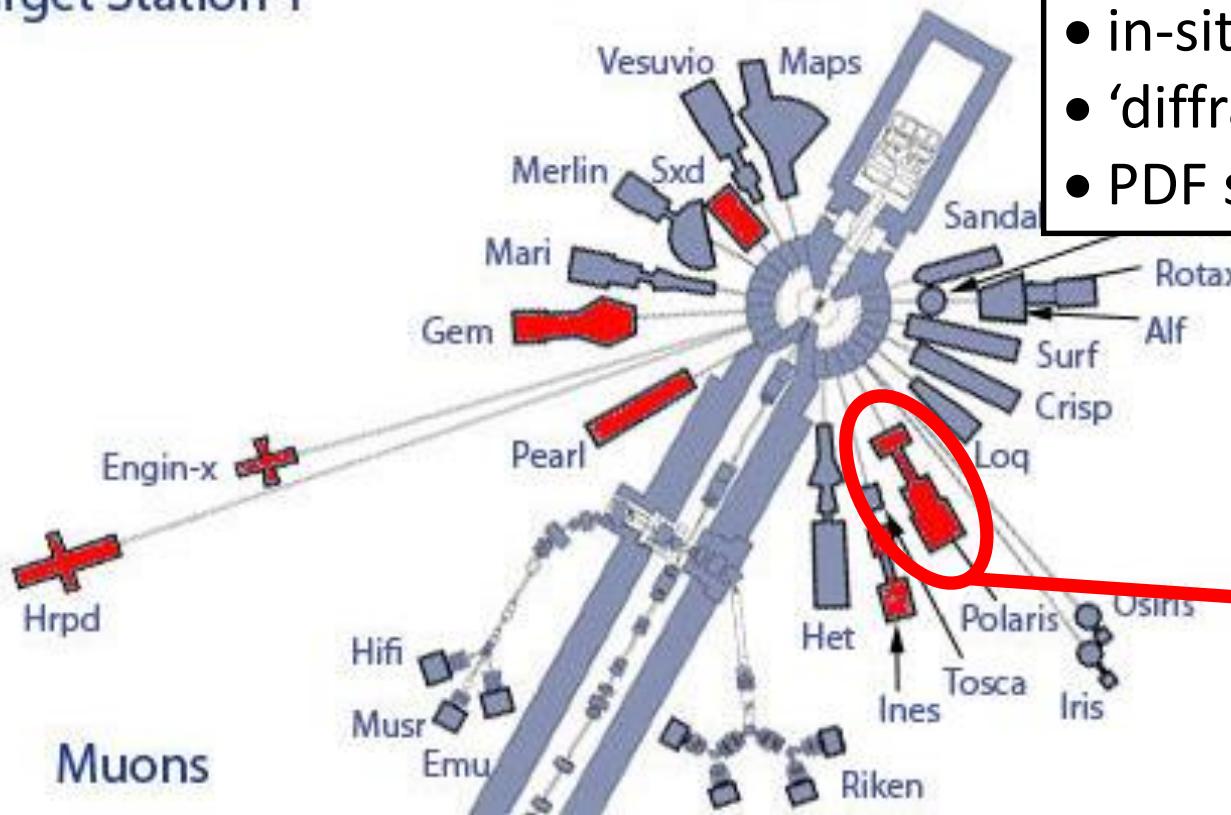


- Monochromatic and Laue type instruments are represented
- Q-range of interest and unit cell volume determine whether hot, thermal or cold neutron spectrum required for both instrument types

<http://www.ill.eu/instruments-support/instruments-groups/>

TOF: high intensity powder diffraction

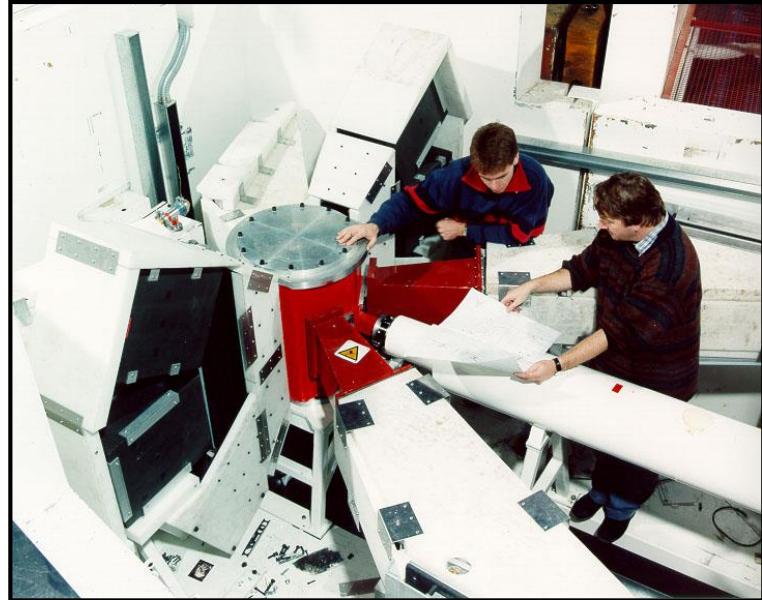
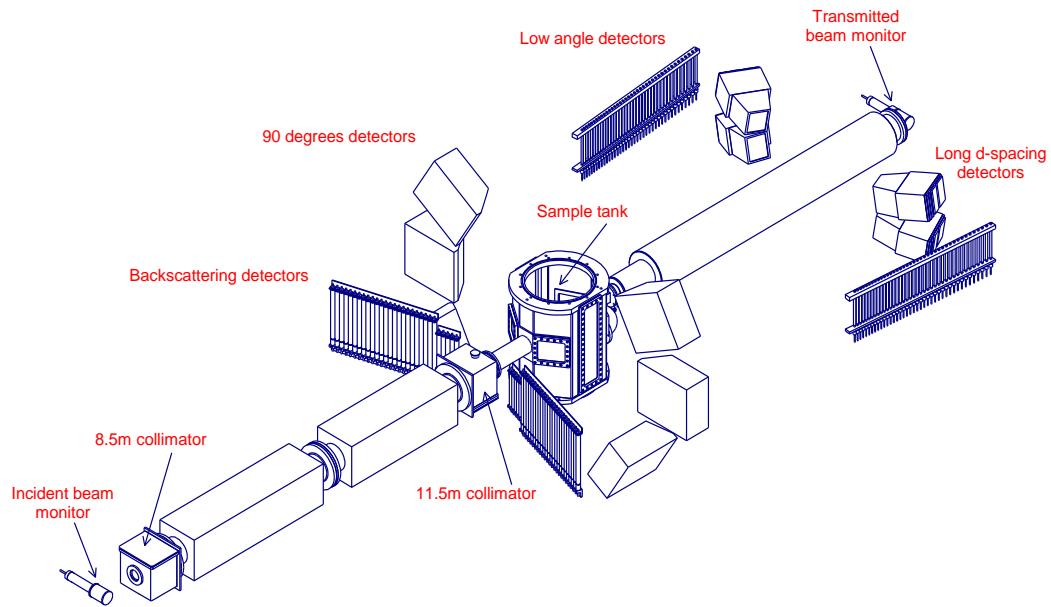
Target Station 1



- chemical crystallography
- in-situ studies
- 'diffraction plus'
- PDF studies

Polaris

Polaris: old configuration



Compare with GEM:

- Higher sample flux
- Wider bandwidth
- Lower resolution
- Hotter spectrum
- Lower detector coverage

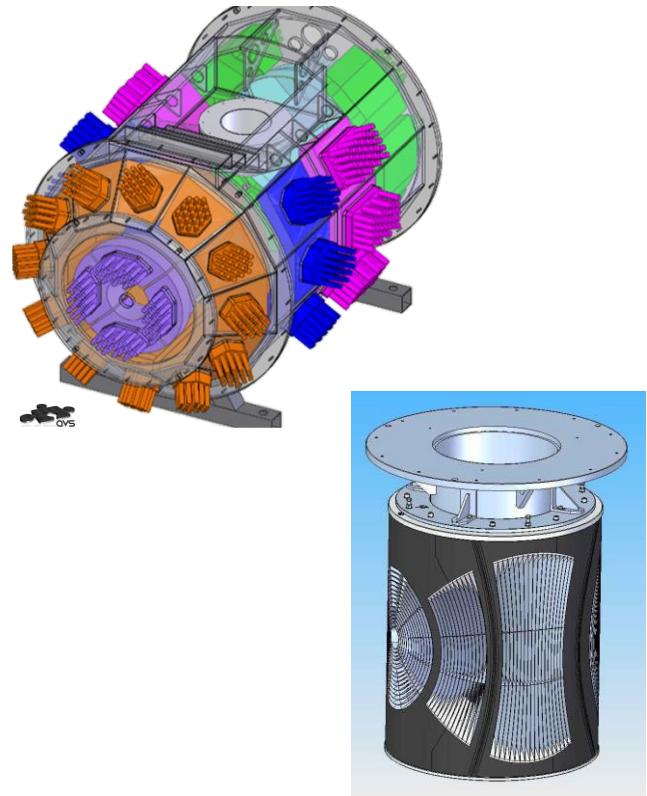
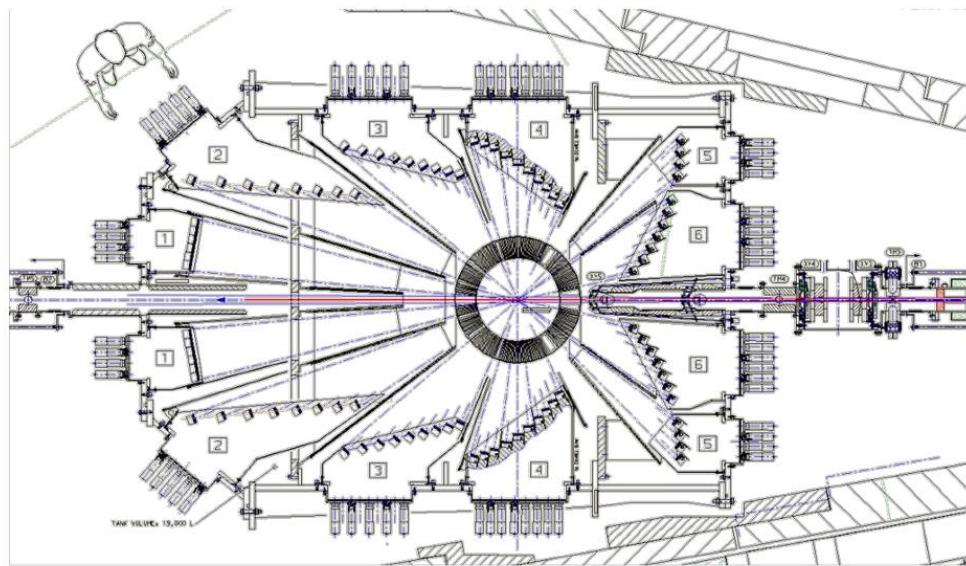
<http://www.isis.stfc.ac.uk/instruments/polaris/polaris4643.html>

Polaris: old configuration

bank position (label)	low angle (A)	low angle (B)	backscattering (C)	90 degrees (E)
detector type	^3He	ZnS	^3He	ZnS
no. of elements	$2 \times 40 = 80$	$4 \times 20 = 80$	$2 \times 29 = 58$	$6 \times 36 = 216$
L_2 (m)	1.72 - 2.65	~2.2	0.65 - 1.35	~0.80
2θ range	$28^\circ < 2\theta < 42^\circ$	$13^\circ < 2\theta < 15^\circ$	$130^\circ < 2\theta < 160^\circ$	$83^\circ < 2\theta < 97^\circ$
Ω (ster)	0.046	0.009	0.29	0.48
$\Delta d/d$	$\sim 1 \times 10^{-2}$	$\sim 3 \times 10^{-2}$	$\sim 5 \times 10^{-3}$	$\sim 7 \times 10^{-3}$
d -range (\AA)	0.5 - 8.3	0.5 - 21.6	0.2 - 3.2	0.2 - 4.0
Q-range (\AA^{-1})	0.75 - 12.6	0.3 - 12.6	2.0 - 31.4	1.5 - 31.4

- Good workhorse instrument for powder diffraction
- High Q accessible for disordered materials investigation using the PDF method
- Some in situ capability but limited by count-rate
- Compatible with a wide range of restricted geometry sample environment

Polaris upgrade



- Increase primary flight path to 14 m
- Optimise each detector bank to give constant resolution
- Increase detector coverage
- Design a collimator to reduce background and parasitic scattering

<http://www.isis.stfc.ac.uk/instruments/polaris/polaris-upgrade-poster11575.pdf>

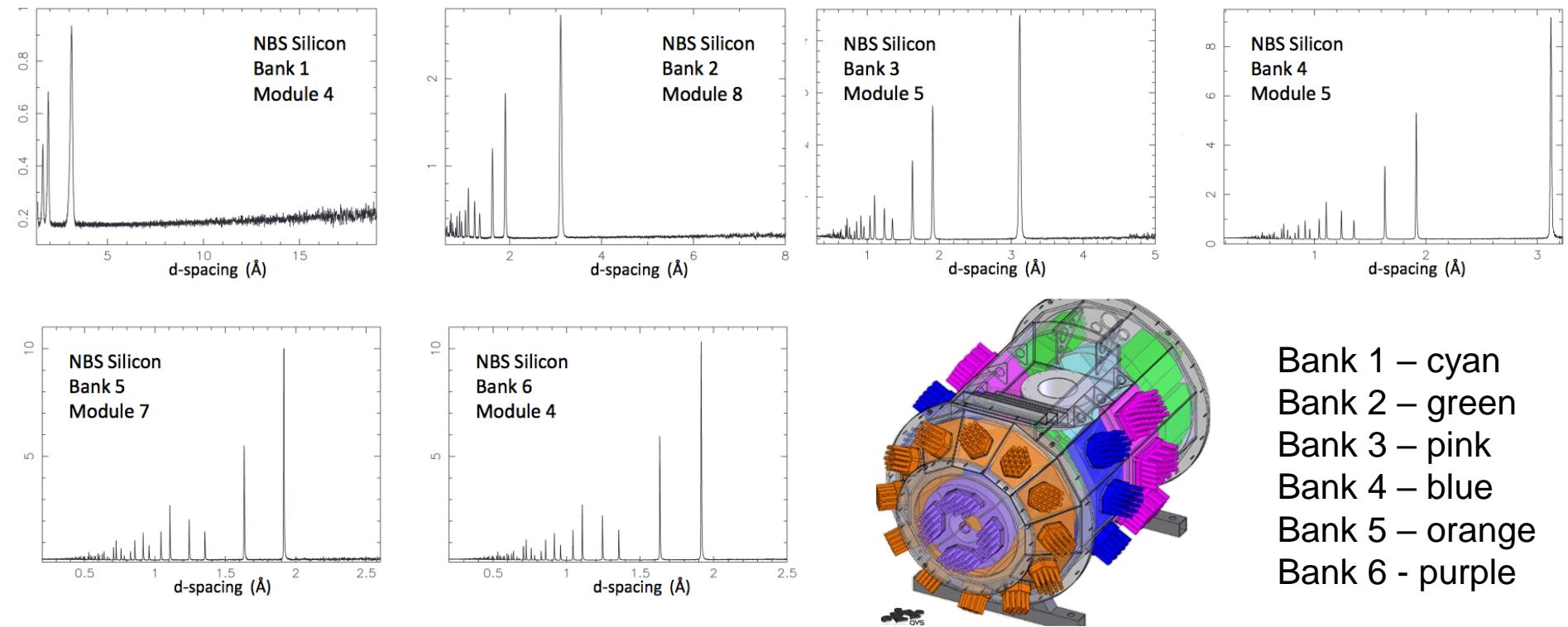
Polaris upgrade



<http://www.isis.stfc.ac.uk/instruments/polaris/polaris4643.html>

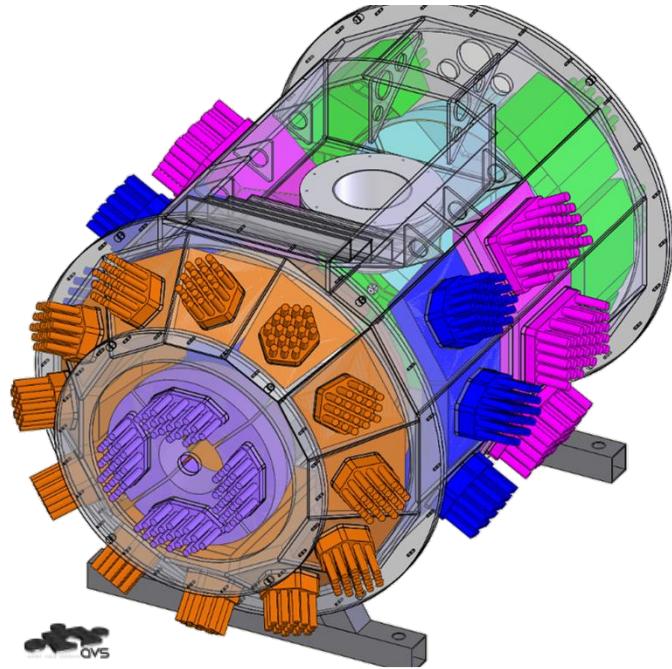
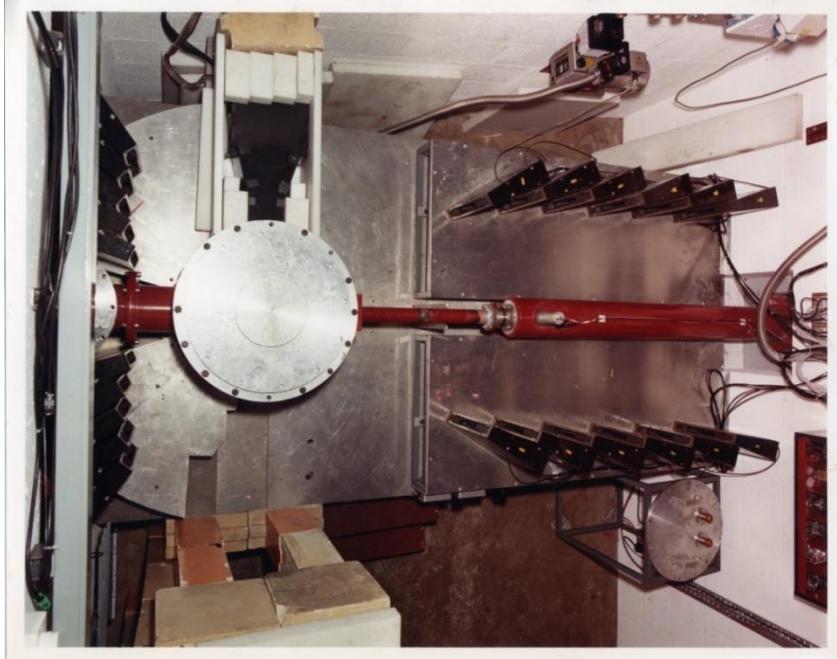
Current Polaris

<http://www.isis.stfc.ac.uk/instruments/polaris/polaris-upgrade---first-diffraction-pattern12763.pdf>



- Increased count rate $\times 3$ at high scattering angle to >20 for low angle banks
- Resolution improvement e.g. bank 5 and 6 of 3×10^{-3} cf. 5×10^{-3}
- Improvement in data at high Q

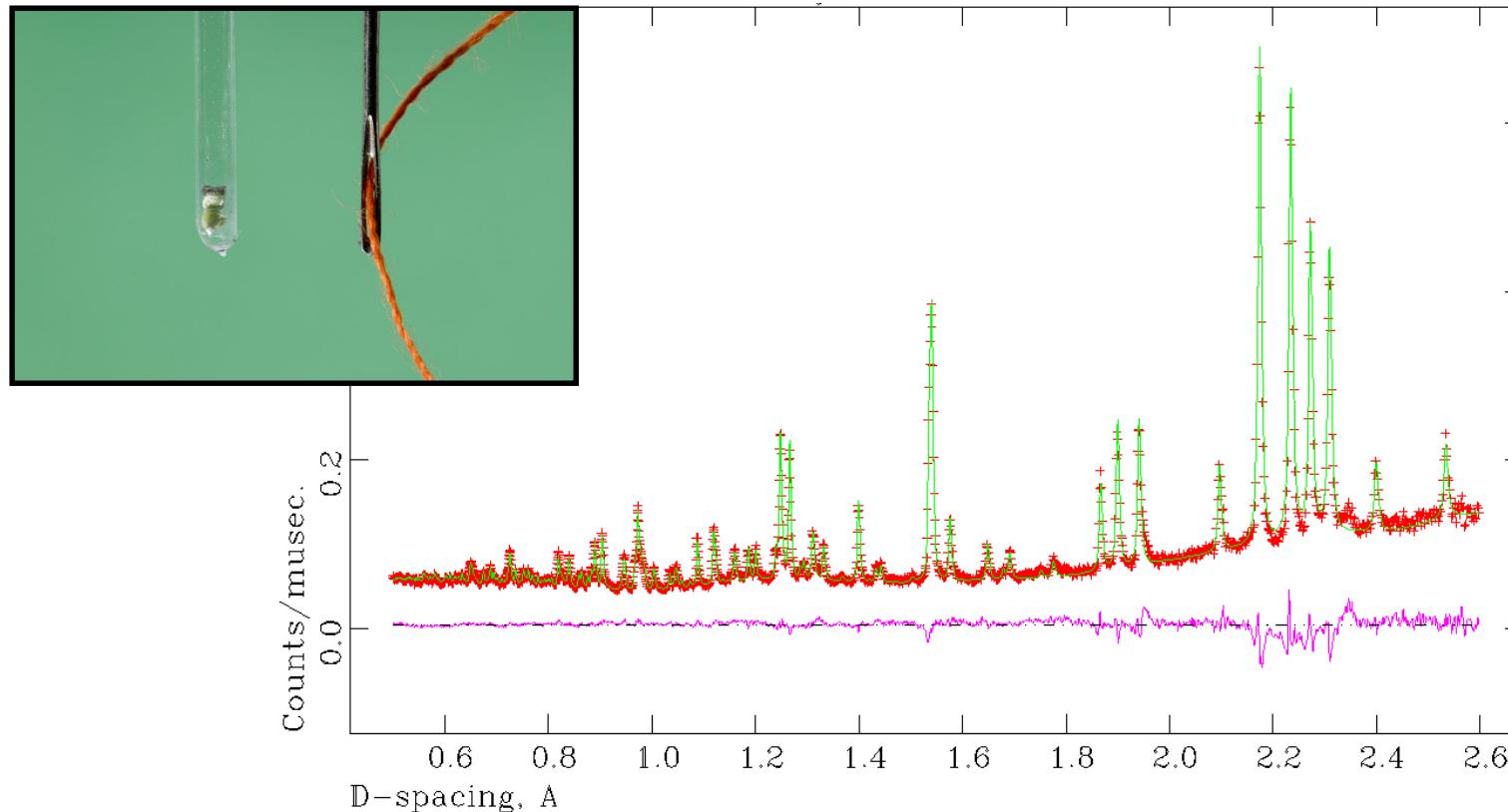
Polaris 1995 v 2013



- 1995 500 mg 24+ hrs
- 2013 500 mg 15-20 minutes with increased Q-range

Contemporary instruments NOVA and iMateria (J-PARC), POWGEN-3 (SNS)

Polaris: pushing boundaries in sample size



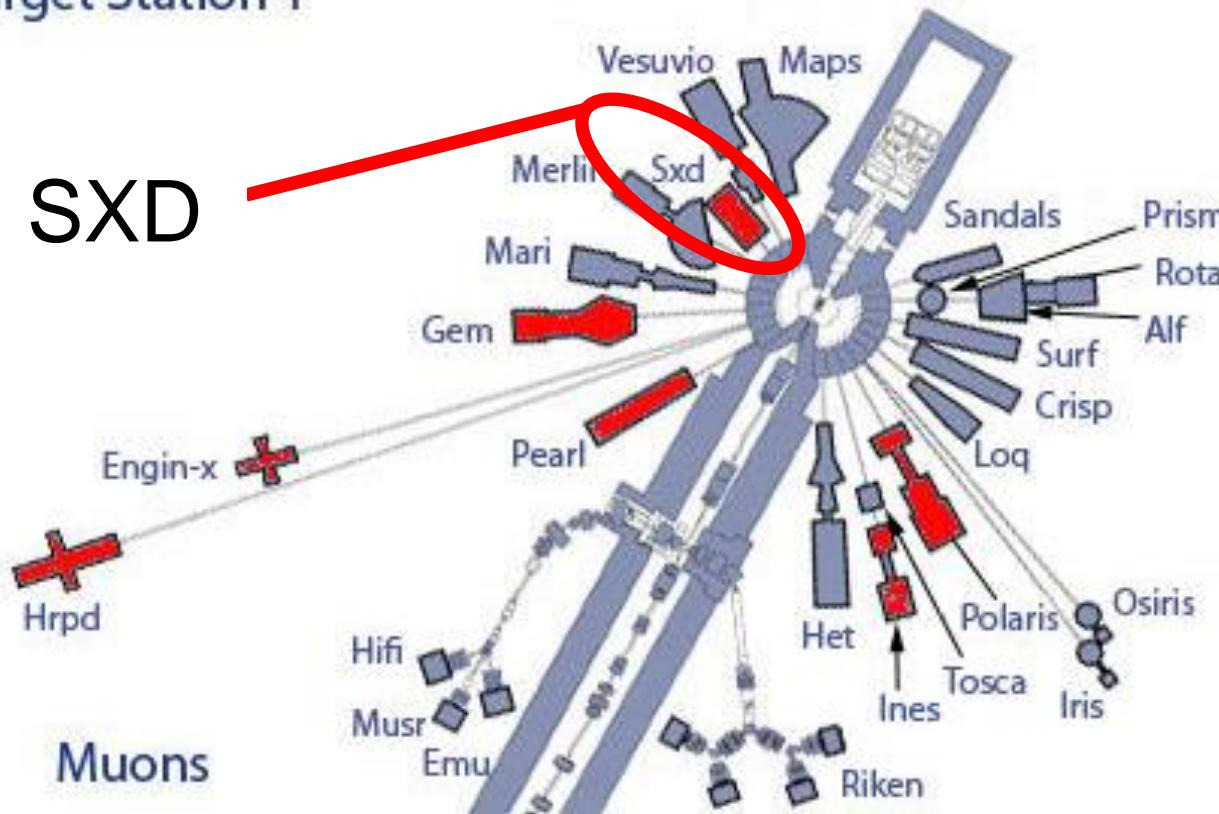
~1mm³ sample of NaNiF₃ phase prepared at high *p* + high T

Lindsay-Scott *et al*, J. Appl. Cryst., **47**, 1939 (2014)

SXD: single crystal diffraction

Target Station 1

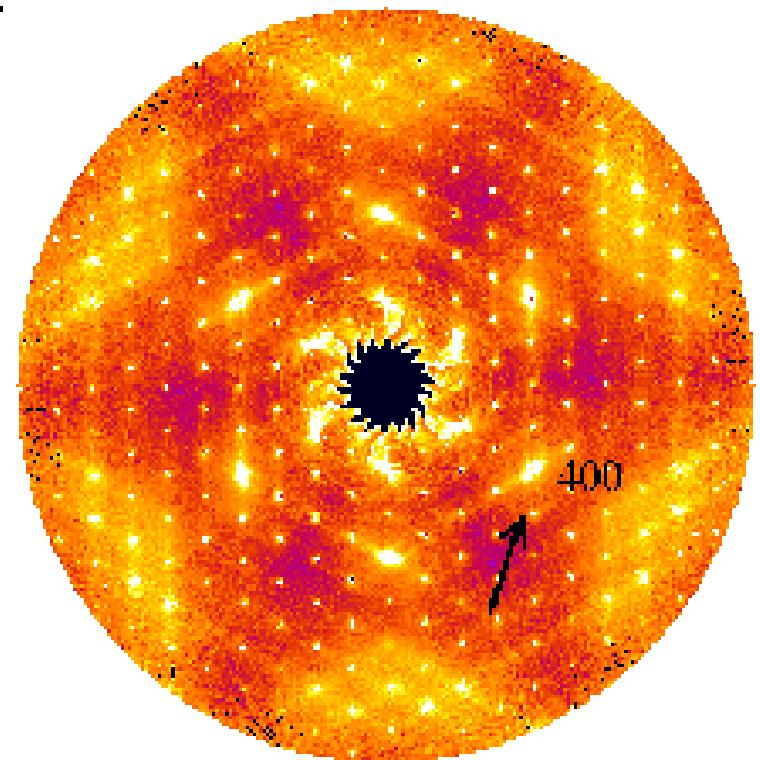
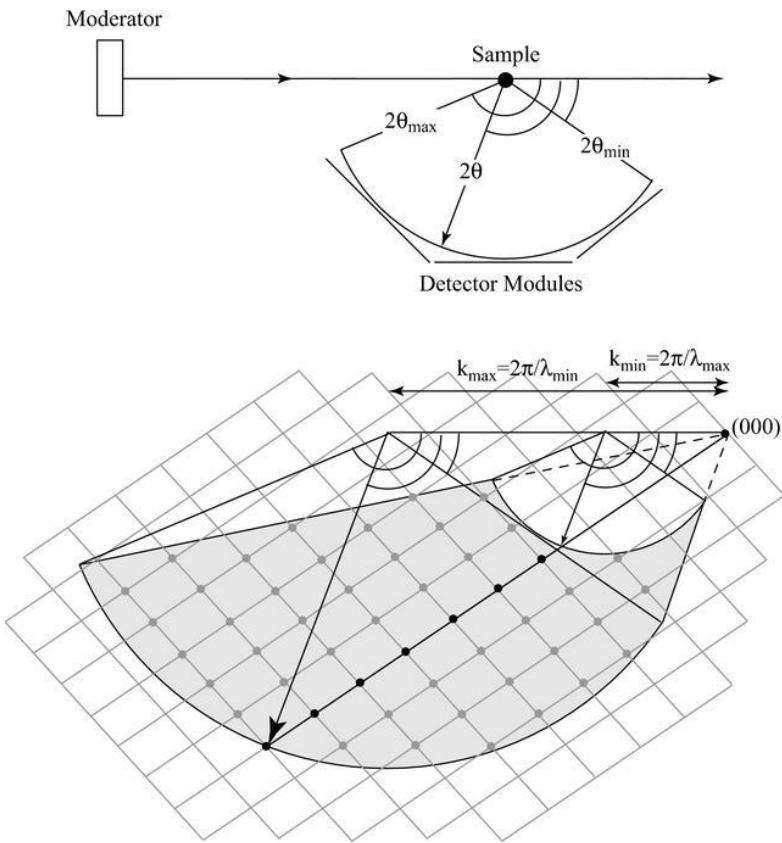
SXD



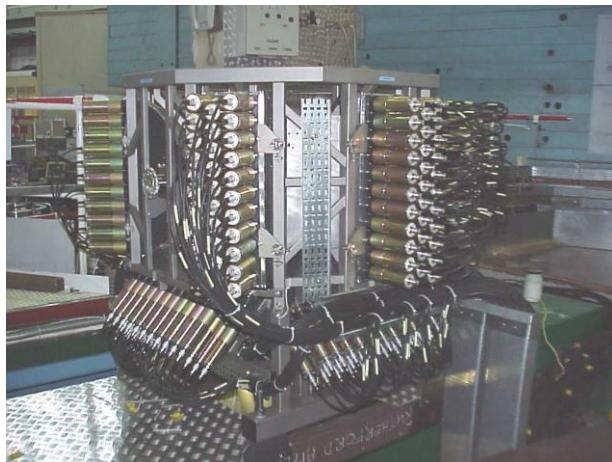
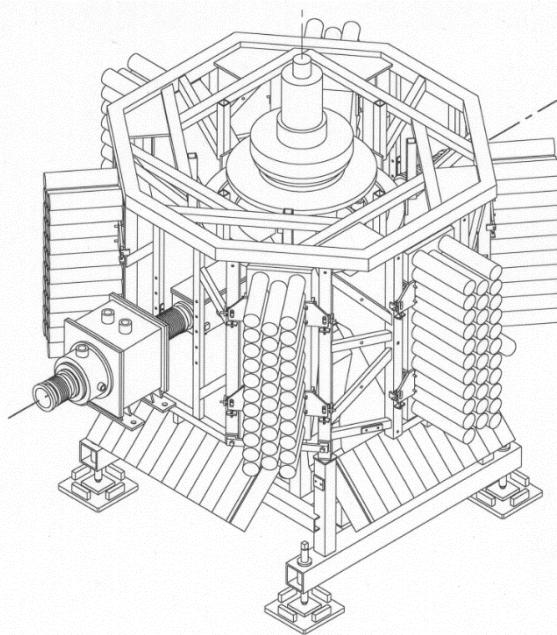
- chemical crystallography
- Q space mapping
- incommensurate structures
- high pressure

TOF Laue method

SXD uses the ‘time-of-flight Laue’ method to scan a large volume of reciprocal space at each crystal orientation.

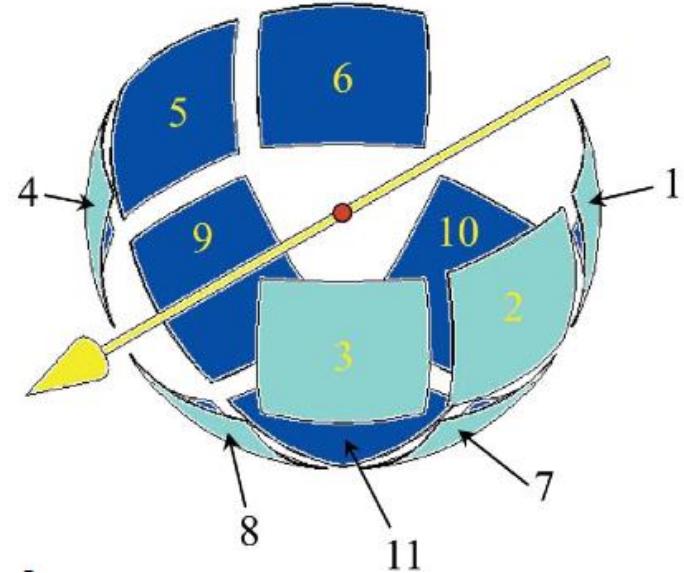


SXD: single crystal diffraction



SXD: single crystal diffraction

- H₂O moderator poisoned at 2 cm
- 0.2 – 10 Å wavelength band
- Primary flight path 8.3 m
- Beam size < 15 mm
- Eleven 192 × 192 mm² detectors
(3 × 3 mm² resolution)



Unlike an image plate set-up the detectors are continuously read-out as a function of TOF allowing spatial overlap to be resolved in the TOF channel while minimising background

Keen *et al.* J. Appl. Cryst. (2006), **39**, 714-722) <http://www.isis.stfc.ac.uk/instruments/sxd/sxd4813.html>

Contemporary instruments: Topaz (SNS), Senju
(J-Parc), Mandi (SNS)

Summary: Diffraction instruments

- Reactors build CW instruments*
 - Low peak brilliance, high time-average brilliance
 - Variable reflectivity from monochromators limit low λ use
 - High Q not easily reached
 - Match moderator and monochromator take-off angle to Q range and resolution
 - Beam always on

*Except when significantly restricted geometry constraints from science case necessitate use of TOF

- Pulsed sources build TOF instruments#
 - High peak brilliance, low time-averaged brilliance
 - Require efficient beam transport
 - High Q possible
 - Increase instrument length to improve resolution at expense of bandwidth
 - Variable Q range and resolution from detector angles
 - Beam availability can compromise science

#Remains to be seen for long pulse sources

Extra Information 1: Uses of diffraction

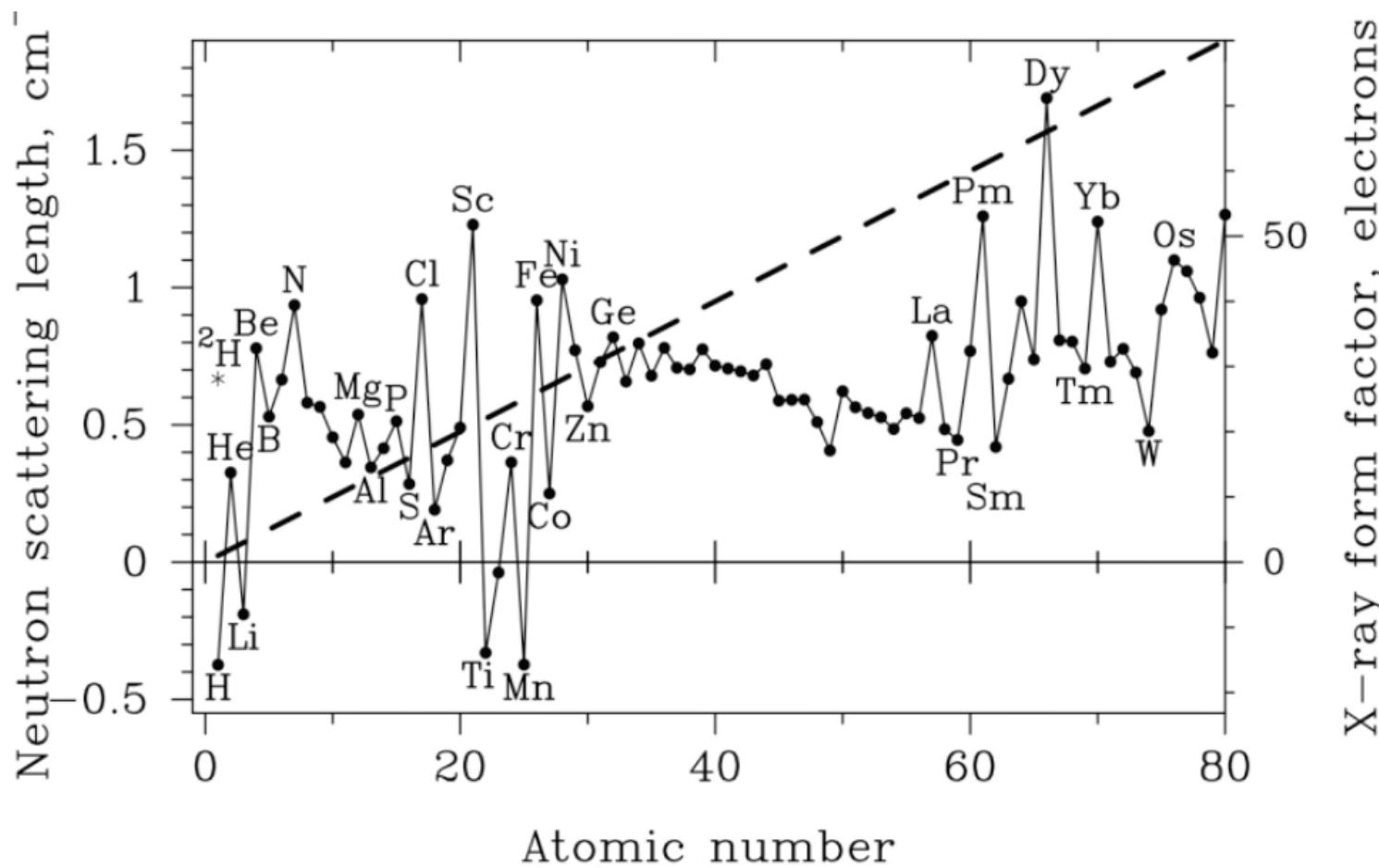


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Uses of diffraction

- Check purity of a sample
- Identify known phases
- Identify new phases
- Collect data for structural analysis
- Follow phase transitions
- Construct phase diagrams (T, P, B, etc)
- Study chemical processes *in situ*
- Monitor particle sizes
- Analyse residual stress within materials
- Process control
- Etc...

X-ray vs neutron scattering power



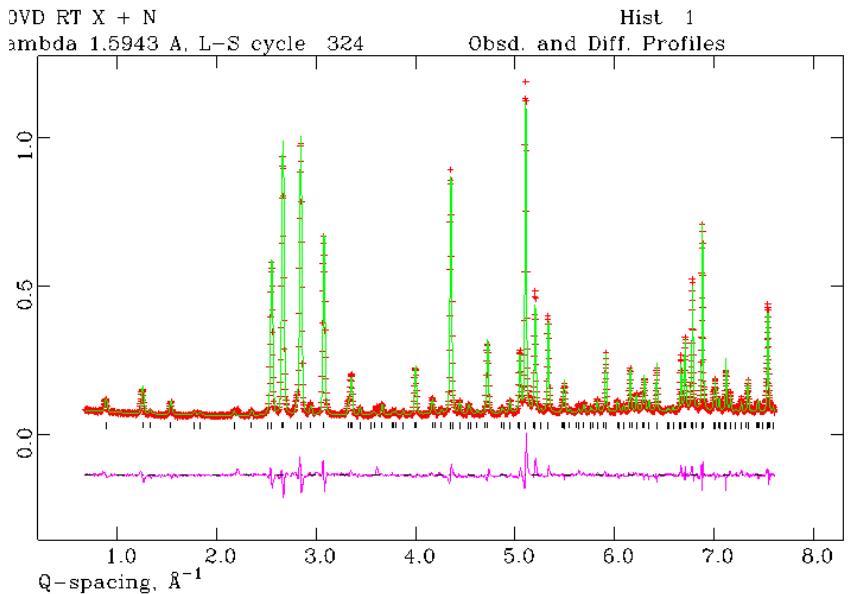
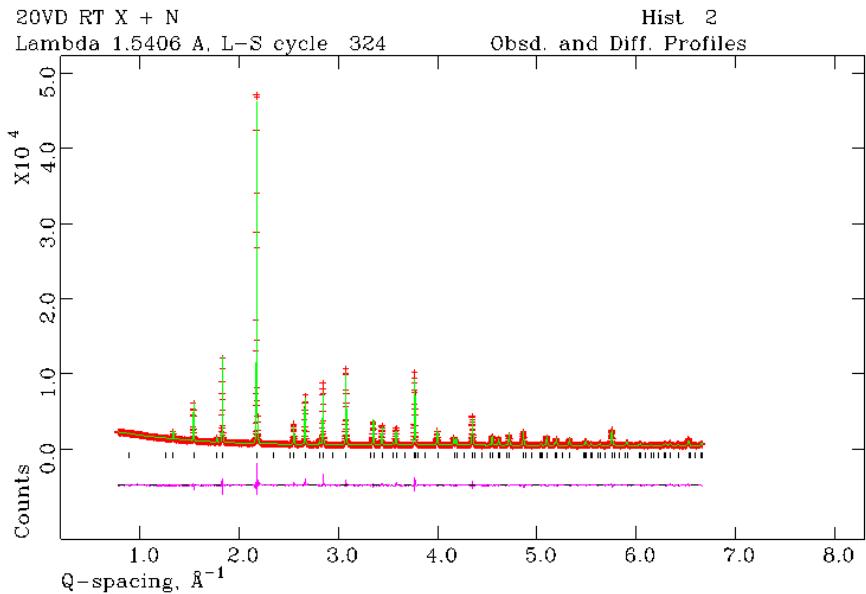
<https://www.ncnr.nist.gov/resources/n-lengths/>



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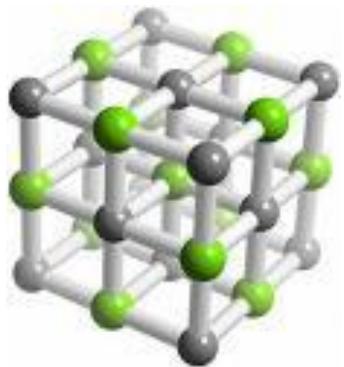
Complementarity of neutrons and X-rays



Note peak intensity differences – complementary information

Main peak in X-ray data (left) almost zero intensity in neutron pattern (right)
Scattering to higher Q in neutron data (form factor)

Neighbouring Element Discrimination



KCl

Fm-3m $a = 6.29 \text{ \AA}$

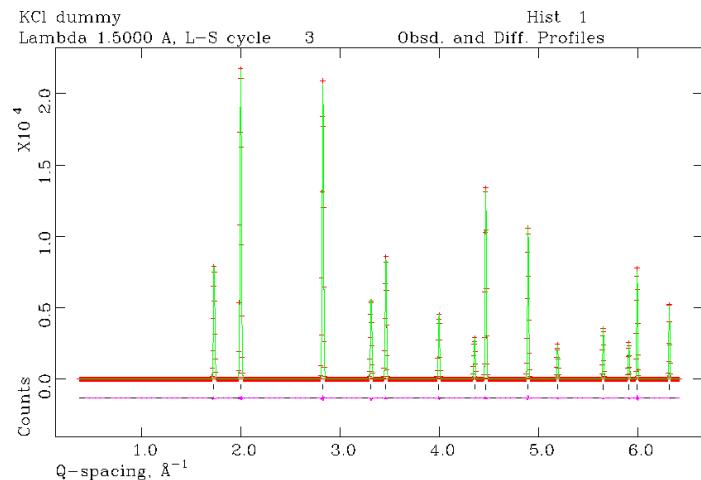
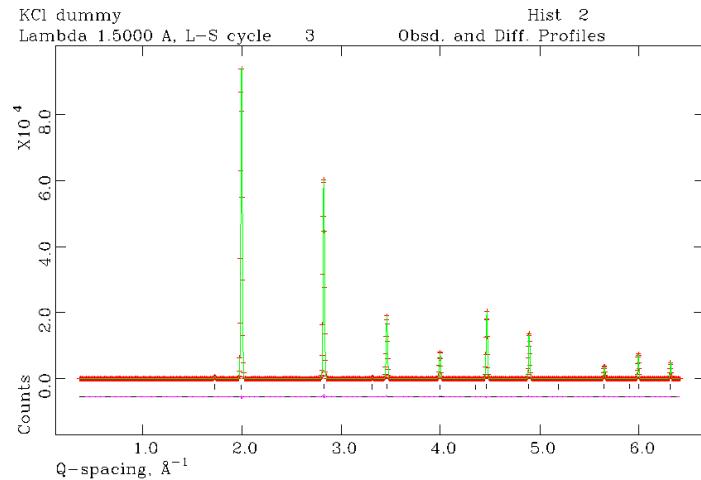
K $z = 19$

Cl $z = 17$

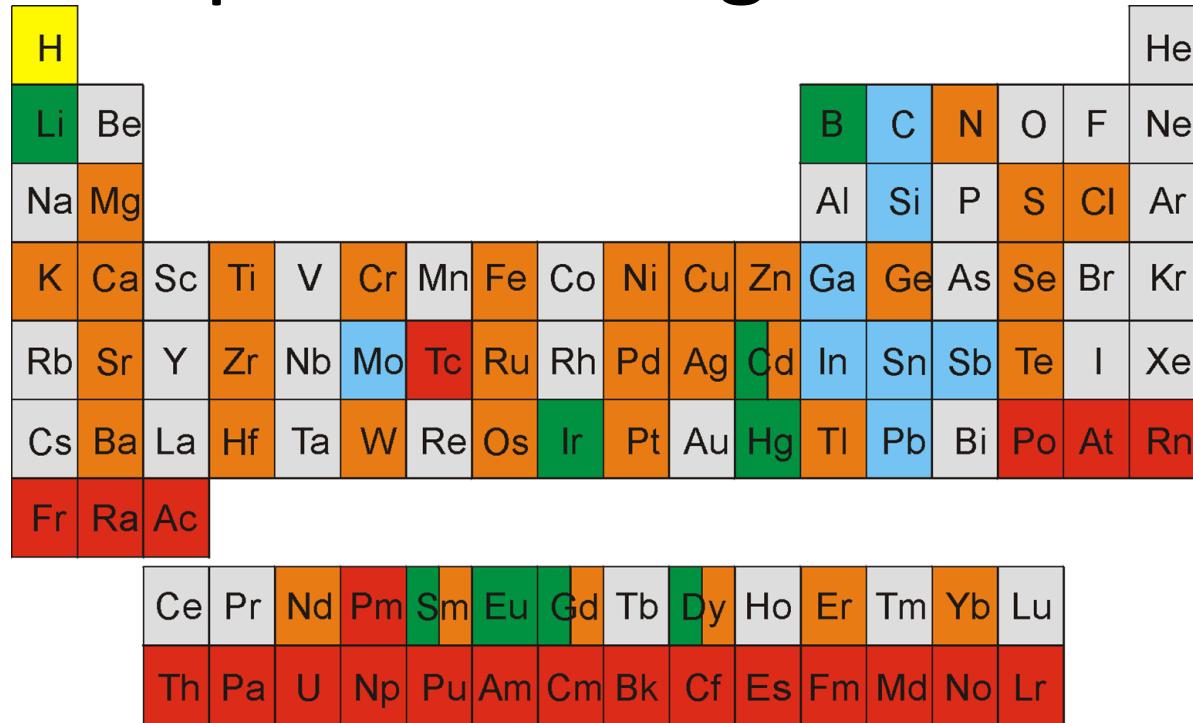
But $\text{K}^+ = \text{Cl}^- = 18 \text{ e}^-$

Without care KCl indexes from X-ray data on a cell that is $\frac{1}{2}$ that from neutron data as elements are identical to X-rays as both have 18 e^-

The non-linear relationship of neutron scattering length between neighbouring elements is crucial



Isotope scattering contrast



> 20% scattering length contrast isotopes

5-20% scattering length contrast isotopes

Non-absorbing isotopes available

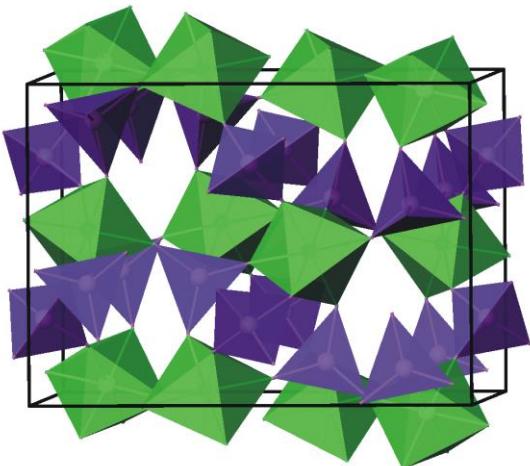
Non-incoherent scattering isotope available

Mono-isotopic elements

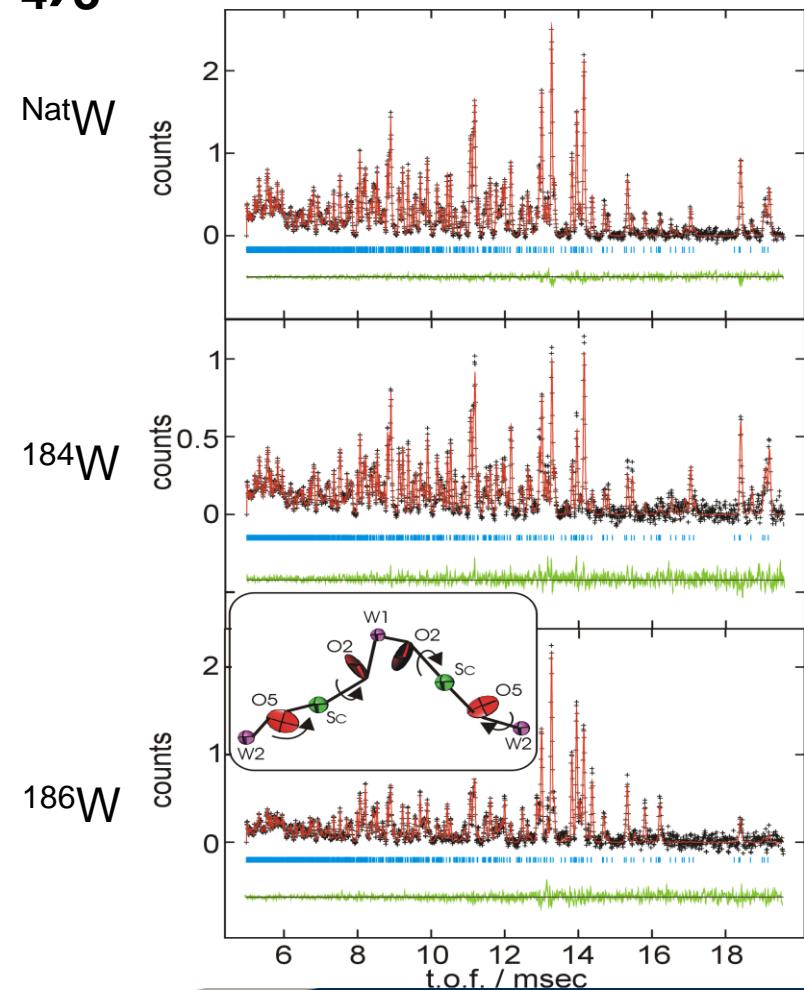
Radioactive elements

Isotopes have different scattering properties for neutrons

Example: Origin of Negative thermal expansion in $\text{Sc}_2(\text{WO}_4)_3$



- Resolved co-operative atomic displacements
- First example using PND
- 10 years later same experiment possible without ISND on upgraded instrument HRPD



M.T. Weller, P.F. Henry, C.C. Wilson. *J. Phys. Chem. B* 2000, **105**(51), 12224-12229

Single crystal or powder?

Depends on the scientific problem:

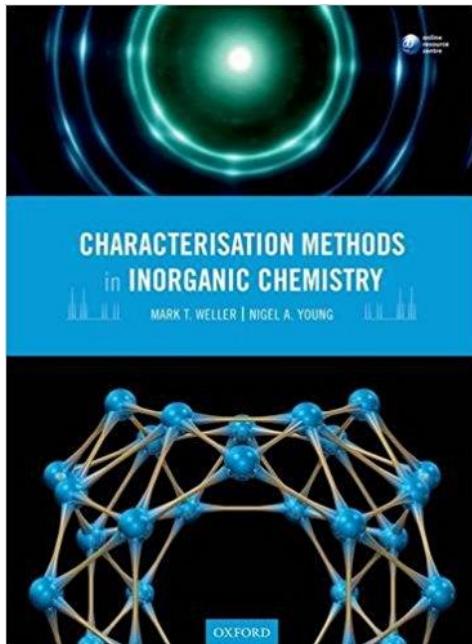
- Unambiguous structure determination – single crystal
 - Beware extinction and absorption issues
- In situ studies – powders
 - Generally the only practical option
- Fast measurements – powders
 - Larger samples
- High background materials (such as incoherent scattering) – single crystal
 - BUT is possible with powder
- Multi-component systems investigations – powder
- Structural phase transitions – powder
 - Crystals tend to shatter
- Real systems – powders
- Very small samples – single crystal
 - Can become difficult to get a powder average

CW or TOF?

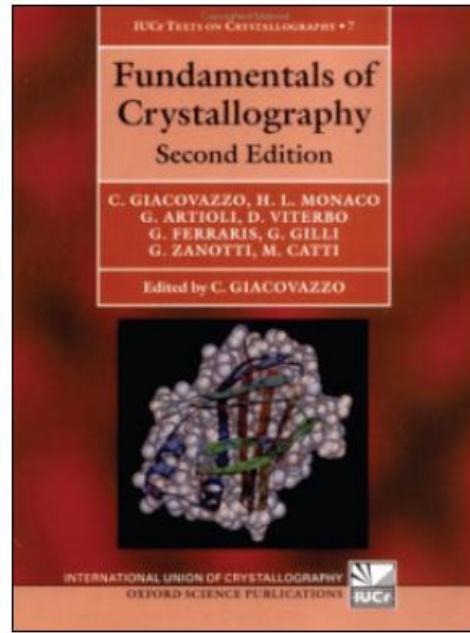
Depends on the scientific problem:

- Structure refinement – either but TOF preferred as complexity increases
- *In situ*, time resolved studies – CW
- Parametric studies – either but CW probably faster for T, P, B mapping
- Fast measurements – either (flux for CW, detector coverage for TOF)
- Small samples – either (lowest background instrument)
- PDF studies – TOF to access high Q
- Magnetic structures – CW preferred but TOF catching up
- Polarised neutron work – traditionally CW but TOF developing
- Hydrogenous materials – CW still preferred but TOF developing
- Large unit cells – either source type, but use Laue methods (quasi-TOF)
- High pressure – TOF offers wider Q range, CW higher flux
- Engineering applications – either, pick depending on Q range
- Texture – either but requires specialized instrument

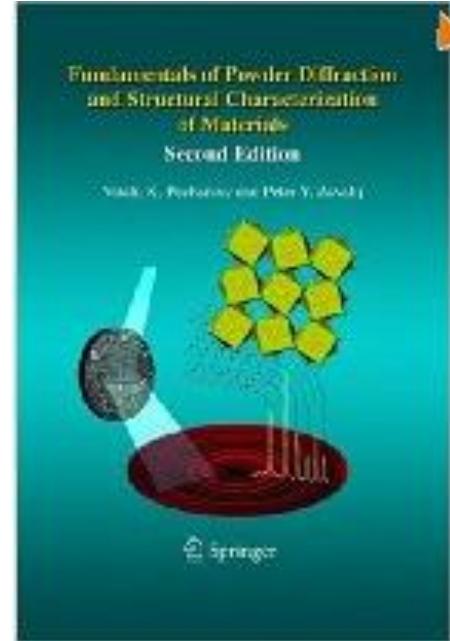
Summary: further reading



Weller & Young
Chapter 2



Giacovazzo *et al.*
Chapter 1



Pecharsky & Zavalij
Chapters 1-9

Extra Information 2

(Structure solution & refinement)



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Intensity and structure factor

$$I_{hkl} \propto |F_{hkl}|^2$$

Measured intensity proportional to F_{hkl}^2 and so we cannot tell whether F_{hkl} is positive or negative – the Phase problem

$$F_{hkl} \propto \sum f_i \exp[2\pi i(hx_i + ky_i + lz_i)] \exp(-U_i Q^2/2)$$

f_i is the scattering power (form factor of the ith site i.e. (x_i, y_i, z_i) and includes fractional occupancy

Contribution of the ith site to the F_{hkl} in question

Atomic displacement of the ith atom site

The phase problem

$$I_{hkl} \propto |F_{hkl}|^2$$

In diffraction we measure the magnitudes and not the phase. The phases contain the bulk of the information. This is why crystallography is hard....

...but not impossible. We can recover phase information from:

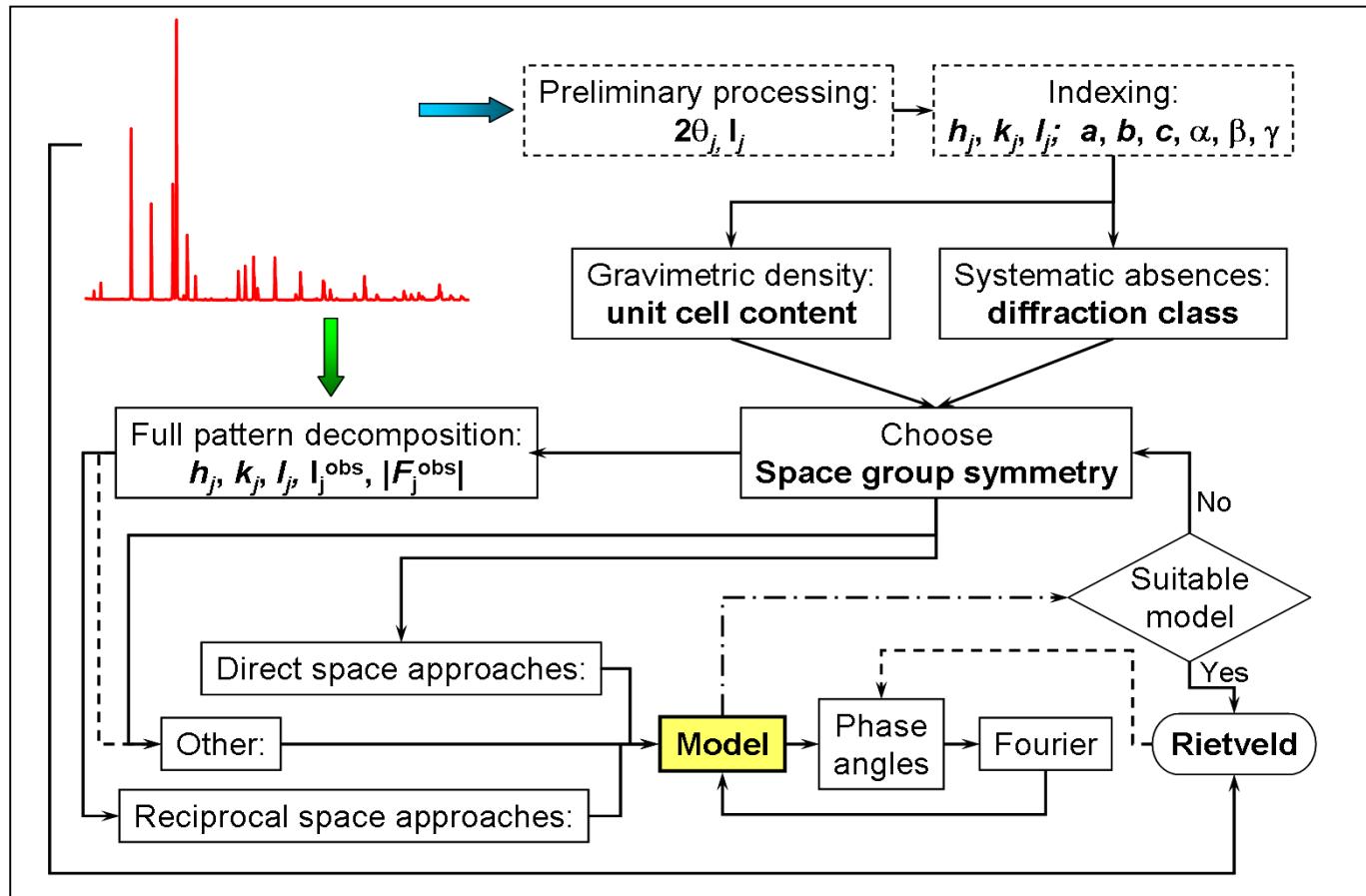
- Related or isostructural materials
- Knowledge of atom positions (heavy atoms from X-rays)
- Known motifs (molecules)
- Brute force

Other factors contributing to measured intensity

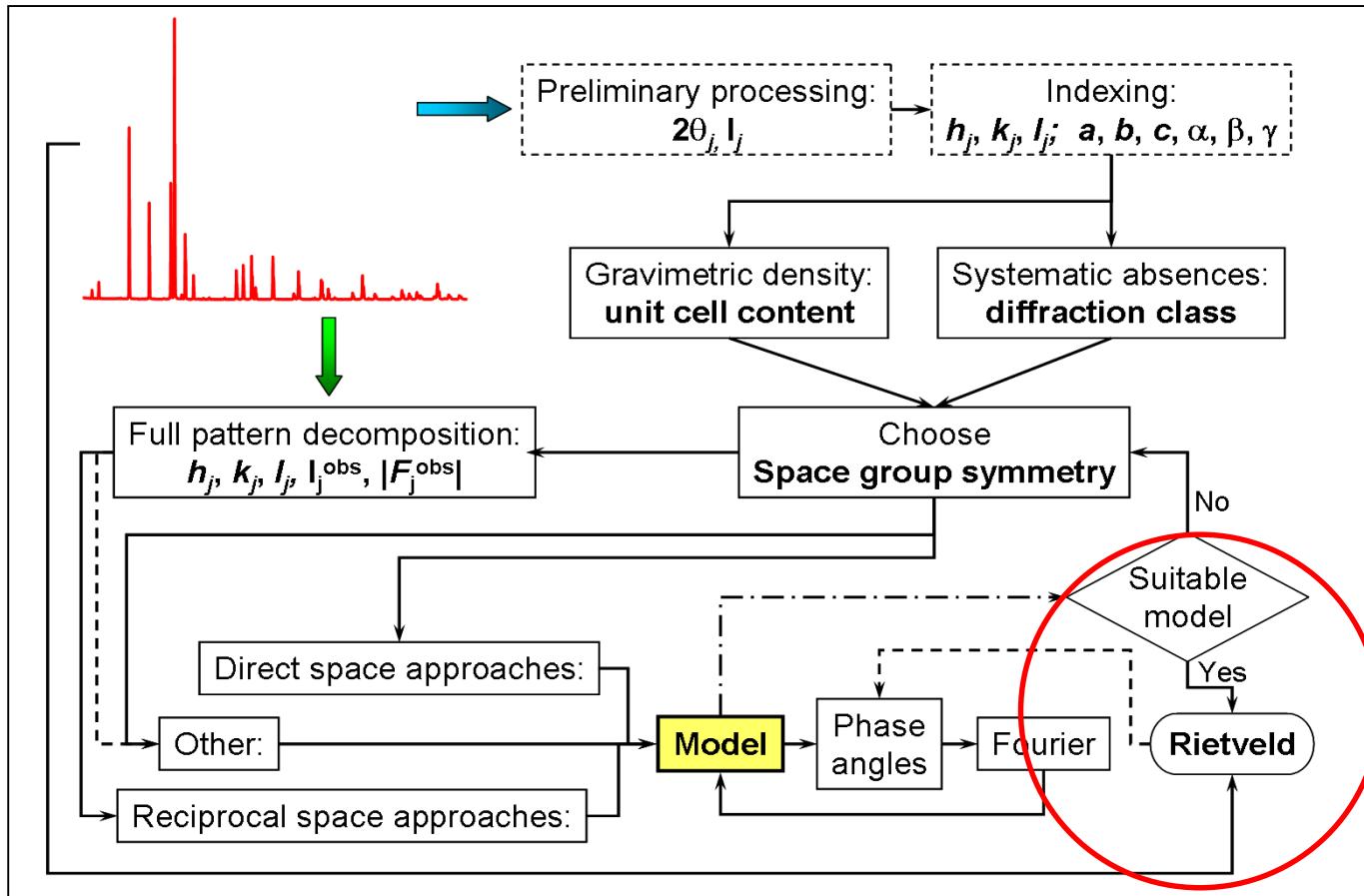
$$I_K = S M_K L_K |F_K|^2 P_K A_K E_K$$

- S is an arbitrary scale factor
 - used to adjust the relative contribution of individual phases to the overall diffraction pattern
- M is the multiplicity of the reflection
 - accounts for the fact that some observed diffraction peaks are actually the product of multiple equivalent planes diffracting at the same position 2θ (for example, (001) (100) (010) etc in cubic)
 - automatically calculated based on the crystal structure
- L is the Lorentz polarization factor
- P is the modification of intensity due to preferred orientation
- A is the absorption correction
- E is the extinction correction
- F is the structure factor, which is the amplitude of the scattering due to the crystal structure

Process for structure solution



Structure refinement



Structure solution/refinement

To unambiguously solve a structure the ratio of resolved and observed unique Bragg reflections to crystallographically independent atoms should be at least 10.

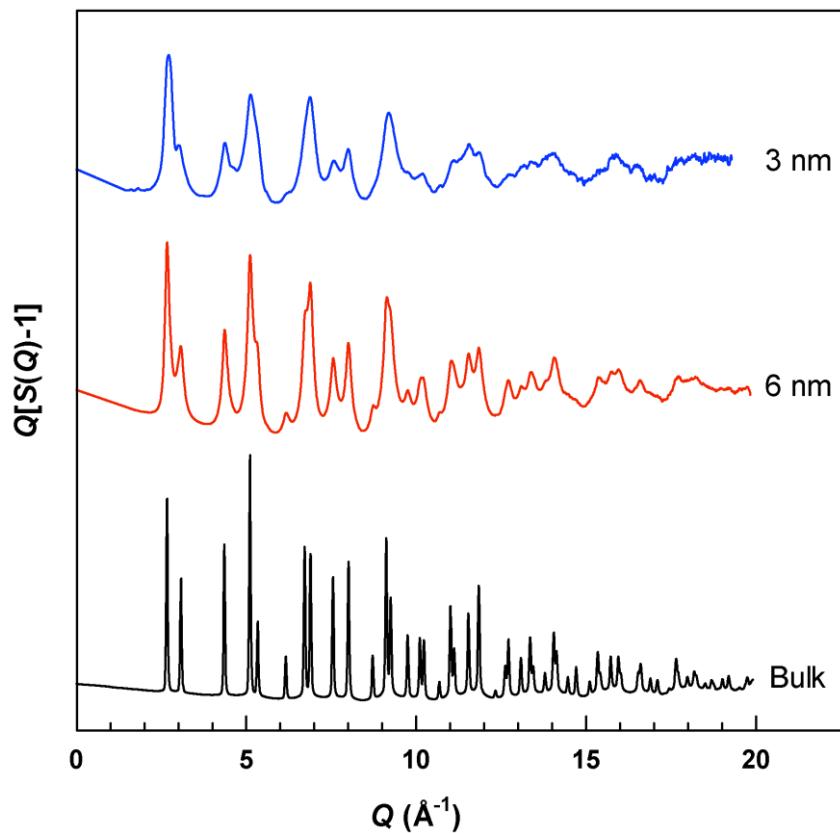
Crystal / powder sample requirements:

- Diffract out to beyond Q_{\max} of the instrument
- Be of adequate particle size
- Preferably not contain any local ordering

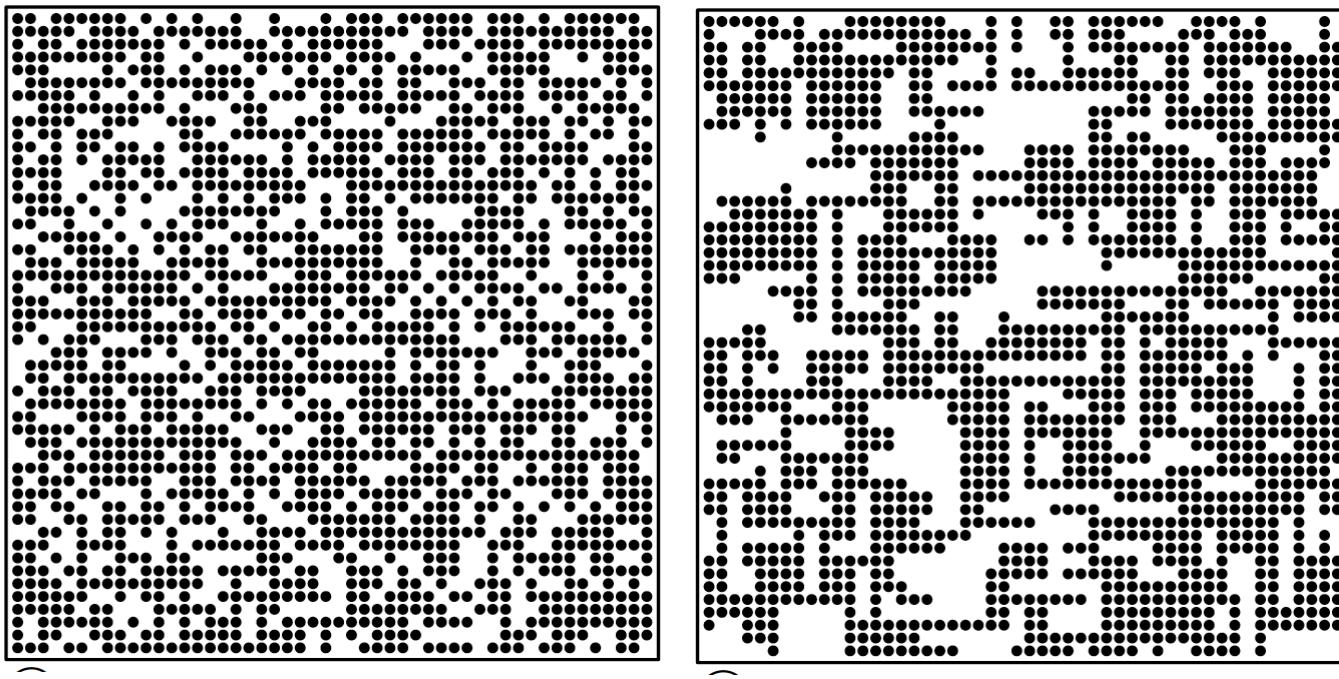
Instrument requirements:

- Access a Q_{\max} such that ratio of accessible Bragg reflections to crystallographically independent atoms is met
- Have a Q-resolution to resolve these Bragg reflections
- Be able to accurately measure the Bragg reflection intensity above the background

Particle size effect

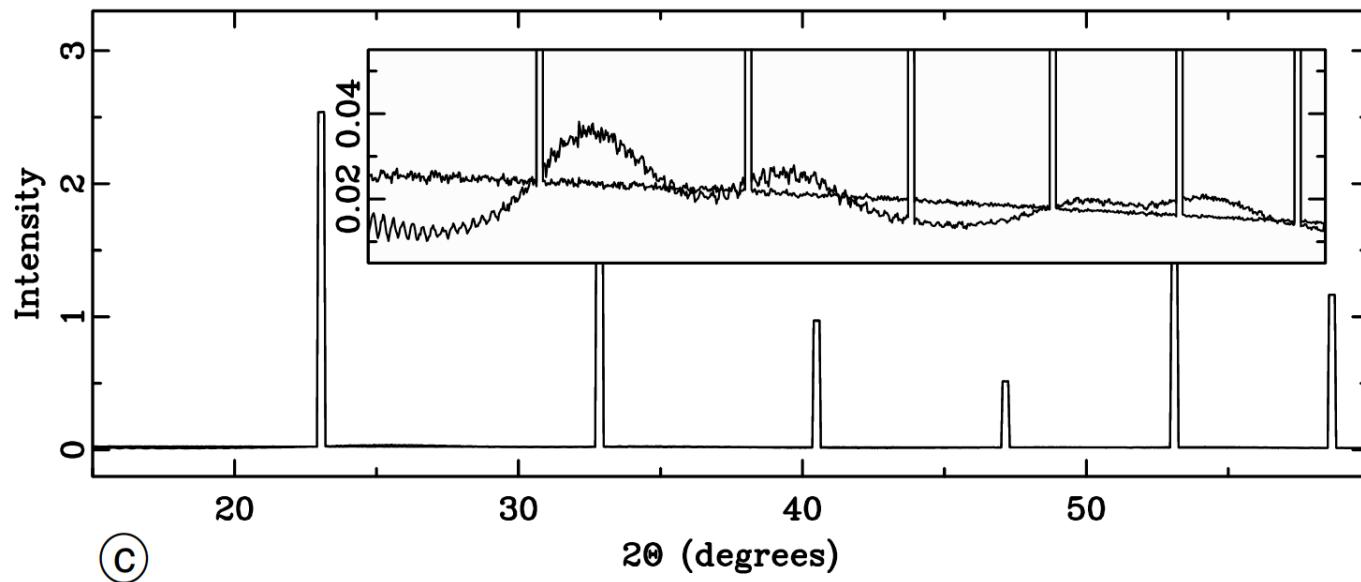


Local Order



(reproduced from Proffen *et al.* (2003). *Z. Kristallogr.* **218**, 132-143)

Local order

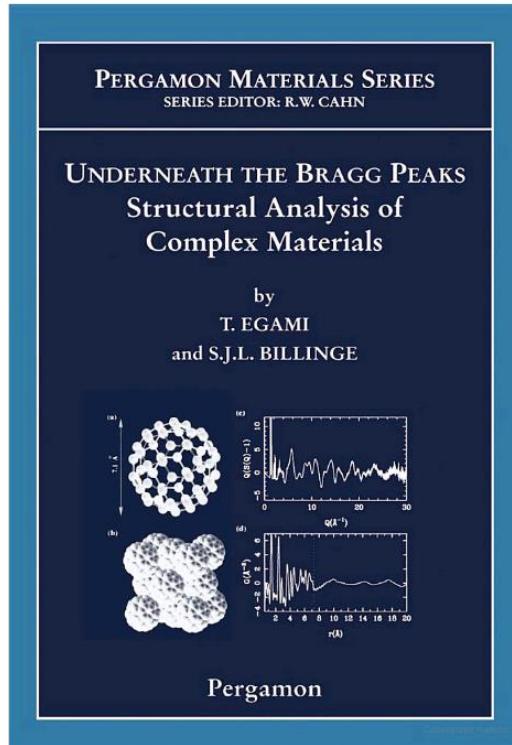


(c)

Diffuse scattering contributions appear between the Bragg reflections. This is ignored by the standard crystallographic approach, which only yields the long range average structure.

(reproduced from Proffen *et al.* (2003). *Z. Kristallogr.* **218**, 132-143)

Pair distribution function (PDF) analysis



The PDF is obtained from the powder diffraction data via a sine Fourier transform of the normalized scattering intensity $S(Q)$:

$$\begin{aligned} G(r) &= 4\pi r[\rho(r) - \rho_0] \\ &= \frac{2}{\pi} \int_0^{\infty} Q[S(Q) - 1] \sin(Qr) dQ, \end{aligned} \quad (1)$$

where $\rho(r)$ is the microscopic pair density, ρ_0 is the average number density and Q is the magnitude of the scattering vector. For elastic scattering $Q = 4\pi \sin(\theta)/\lambda$ with 2θ being the scattering angle and λ the wavelength of the radiation used.

<http://www.totalscattering.org>



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The Rietveld method

J. Appl. Cryst. (1969), **2**, 65

A Profile Refinement Method for Nuclear and Magnetic Structures

BY H. M. RIETVELD

Reactor Centrum Nederland, Petten (N.H.), The Netherlands

(Received 29 November 1968)

A structure refinement method is described which does not use integrated neutron powder intensities, single or overlapping, but employs directly the profile intensities obtained from step-scanning measurements of the powder diagram. Nuclear as well as magnetic structures can be refined, the latter only when their magnetic unit cell is equal to, or a multiple of, the nuclear cell. The least-squares refinement procedure allows, with a simple code, the introduction of linear or quadratic constraints between the parameters.

- Originally written to analyse neutron powder diffraction data
- Both nuclear and magnetic structure refinement
- Adapted for X-ray methods in 1977 by Young
- Thousands of publications per year published using the method
- It is the reason powder crystallography is so successful!!



Hugo M. Rietveld 1932-2016

“NASA would never have sent an X-ray powder diffractometer to Mars without the Rietveld method” (David Blake, 2012)

Rietveld refinement software

Many programs out there. Well used examples include:

- GSAS
- GSAS-II
- Fullprof
- Topas
- Jana
- Maud
- Reitan
- BGMN
- Etc...

The Rietveld method

What it is not:

- A phase identification tool – use databases for this
- For structure solution – it refines a given structural model to data

What it can tell us:

- Phase fractions
- Unit cell dimensions
- Atomic coordinates / bond lengths / substitutions and vacancies
- Strain and texture effects

What you need:

- Good quality data
- A good starting structural model
- An instrument description file



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The Rietveld method

- The intensity, Y_{ic} , of each individual data point i is calculated using the equation:

$$Y_{ic} = Y_{ib} + \sum_{k=k1}^{k2} G_{ik} I_k$$

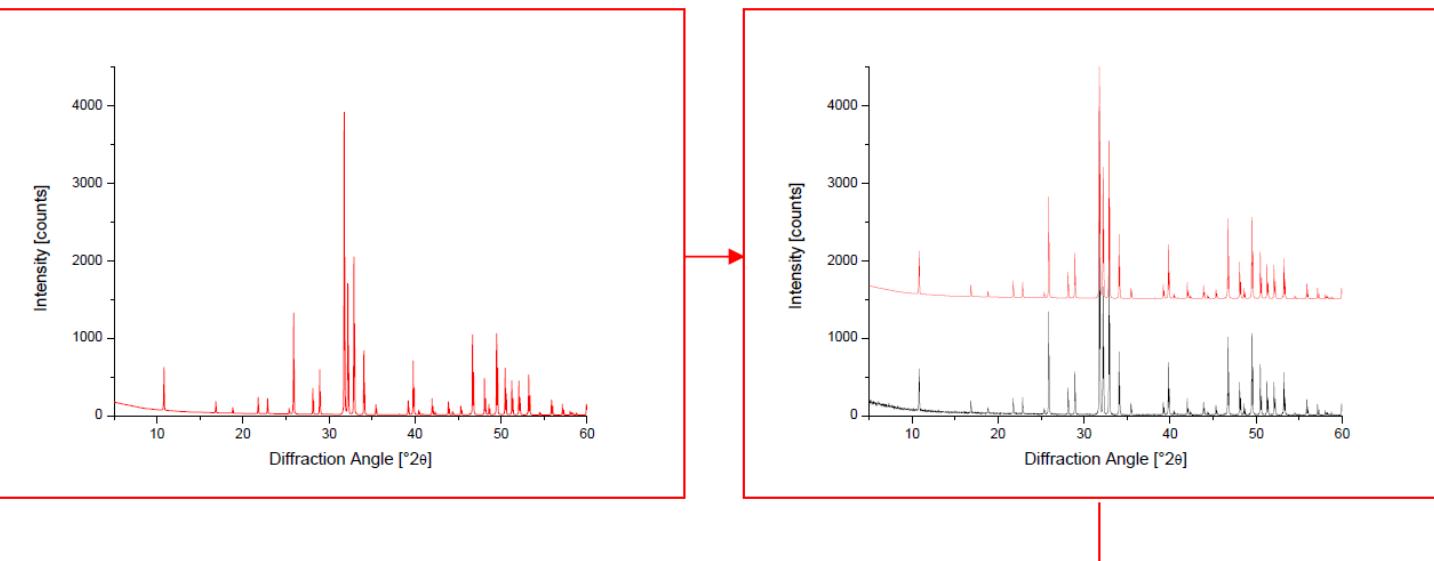
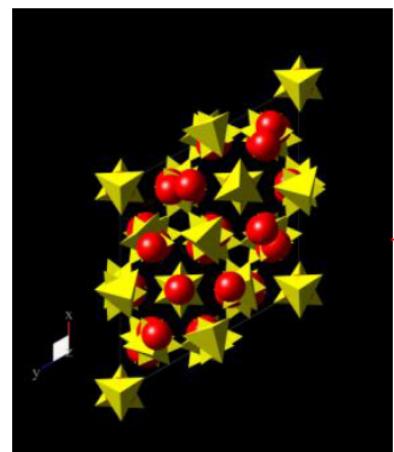
- We already know how to calculate I_k , the intensity of the Bragg diffraction peak k :
 $I_k = S M_k L_k |F_k|^2 P_k A_k E_k$
- Y_{ib} is the intensity of the background at point i in the pattern
- $k1 - k2$ are the reflections contributing to data point i in the pattern
 - sometimes multiple Bragg diffraction peaks overlap, resulting in multiple contributions to the observed intensity at a single data point
- G_{ik} is the peak profile function
 - this describes how the intensity of the diffraction peak is distributed over a range of 2θ rather than at a single point
 - this profile is due to instrument broadening, sample broadening, etc

The Rietveld method

Known structure model

Calculate theoretical diffraction pattern

Compare with measured pattern



Optimize structure model, repeat calculation

Basic refinement procedure

Experimental diffraction pattern

Starting crystal structure (.cif, ICSD)

Instrument file (.inst, LaB₆ standard)

Refine:

- Background
- Lattice parameters
- Peak intensities
- Peak shapes
- Peak positions
- Phase fractions



bad peak position



poor peak shape



peak intensities are off

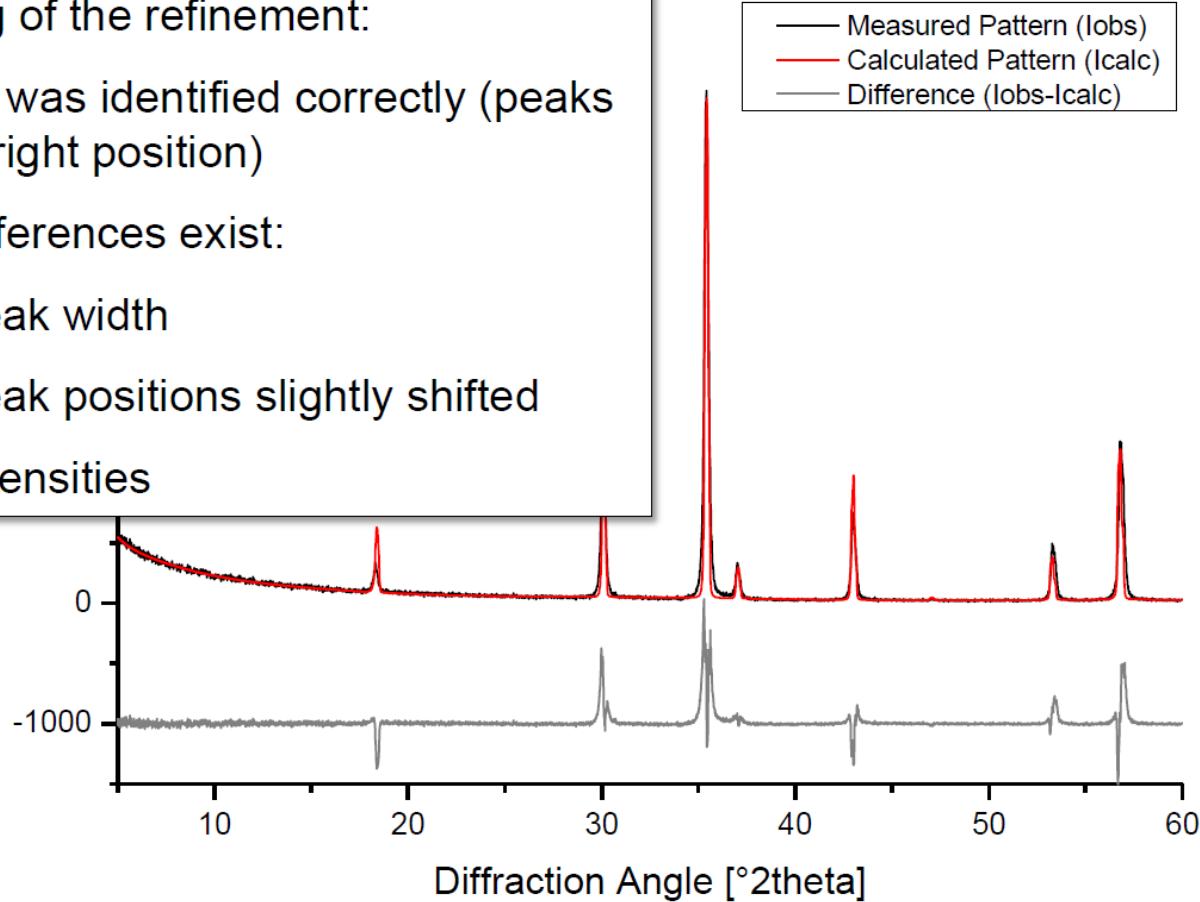
Assess:

- Goodness of fit/R factors
- Impurity phases
- Peak/background shapes
- Difference pattern

Initial model

Beginning of the refinement:

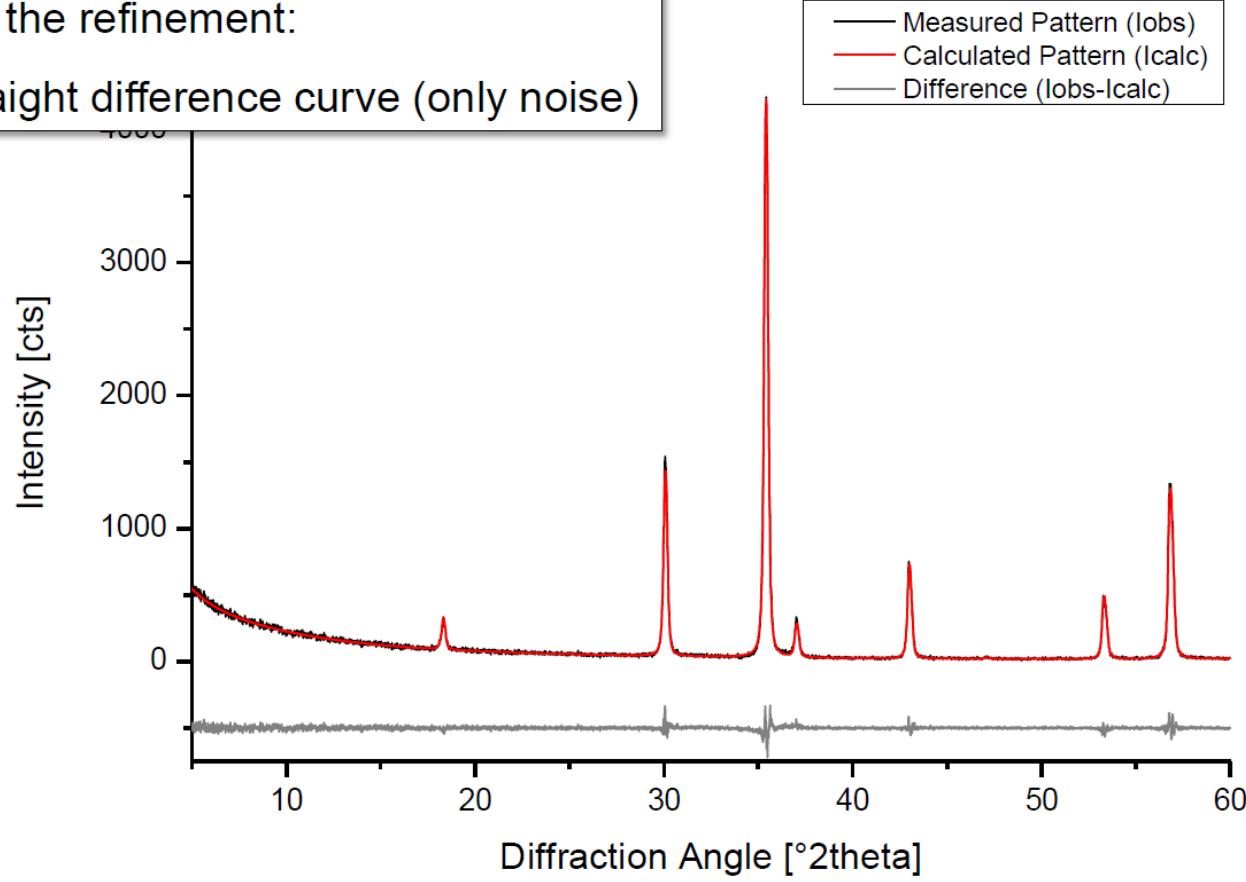
- Phase was identified correctly (peaks at the right position)
- But differences exist:
 - Peak width
 - Peak positions slightly shifted
 - Intensities



Final refined model

After the refinement:

- Straight difference curve (only noise)



Some common challenges/problems

- Incorrect starting crystal structure
- Poor quality data!
- False minimas
- Refinement diverges (“blows up”)
- Over interpretation
- Refine unnecessary variables
- Parameter correlation
- Which goodness of fit to choose? R vs. Chi sq?
- Preferred orientation
- High background
- Ignores non-Bragg diffraction peak information

Summary: complete guide for Rietveld refinement

36

J. Appl. Cryst. (1999). **32**, 36–50

Rietveld refinement guidelines

L. B. McCUSKER,^{a*} R. B. VON DREELE,^b D. E. COX,^c D. LOUËR^d AND P. SCARDI^e

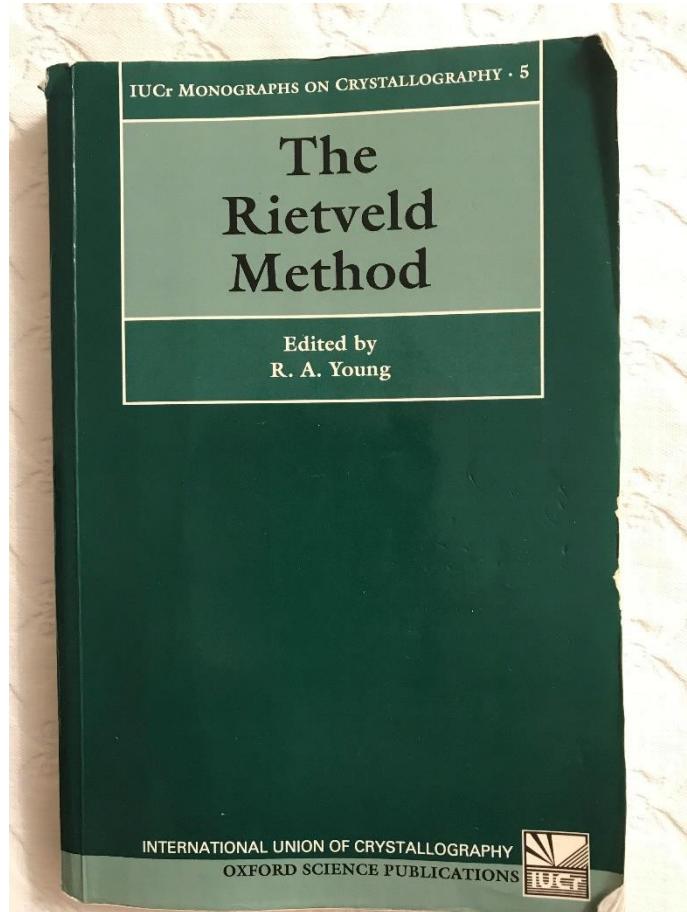
^a*Laboratorium für Kristallographie, ETH, Zürich, Switzerland*, ^b*LANSCE, Los Alamos National Laboratory, Los Alamos, NM, USA*, ^c*Physics Department, Brookhaven National Laboratory, Brookhaven, NY, USA*, ^d*Laboratoire de Chimie du Solide et Inorganique Moléculaire (UMR 6511 CNRS), Université de Rennes I, Rennes, France*, and

^e*Dipartimento di Ingegneria dei Materiali, Università di Trento, 38050 Mesiano (TN), Italy.*

E-mail: lynne.mccusker@kristall.erdw.ethz.ch

(Received 23 February 1998; accepted 22 July 1998)

Background and theory



Extra Information 3

(CW v TOF for resolution)
(ILL and ISIS diffraction suites)

CW or TOF: Q resolution

Monochromatic

$$\frac{\Delta d}{d} = \frac{1}{2} \sqrt{U \cot^2(\theta) + V \cot(\theta) + W}$$

- U, V and W are functions of the collimation and U, V also takeoff angle to the monochromator
- Resolution minimum found near the takeoff angle of the monochromator $2\theta_M$
- Higher takeoff angle gives higher resolution for identical wavelength
- Wavelength produced by monochromator is takeoff angle dependent for any particular hkl plane

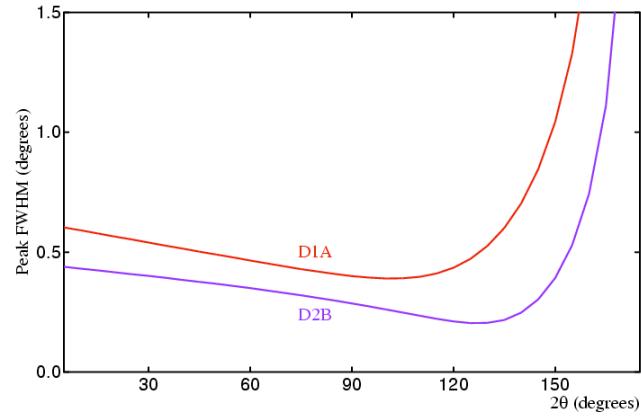
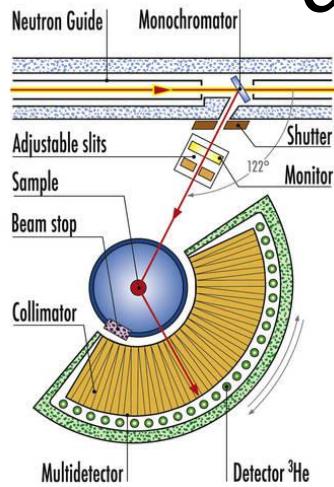
Time-of-flight

$$\frac{\Delta d}{d} = \left[\Delta\theta^2 \cot^2 \theta + \left(\frac{\Delta t}{t} \right)^2 + \left(\frac{\Delta L}{L} \right)^2 \right]^{\frac{1}{2}}$$

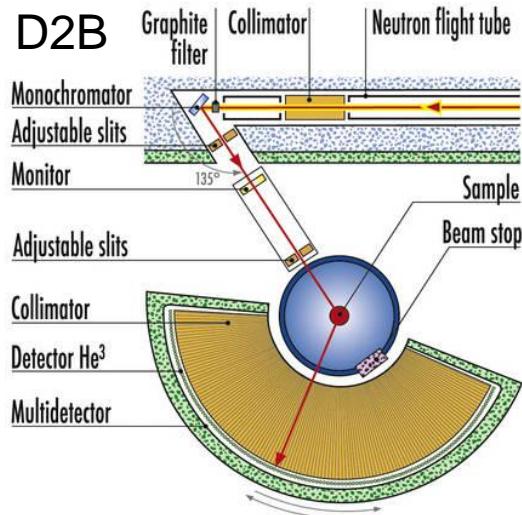
- $\Delta\theta$ is the angular uncertainty
- The main component of Δt is the moderation time of the neutron
- ΔL is the flight path uncertainty of the neutron mainly due to the finite width of the moderator
- First term can be minimised by moving to higher scattering angle
- Second and third terms minimised by increasing instrument length

CW, High resolution, thermal powder diffractometers: D1A, D2B

D1A



D2B



Science: Inorganics, small molecule, magnetism,
 Q range up to about 12 \AA^{-1}

<http://www.ill.eu/instruments-support/instruments-groups/>

CW, High resolution powder diffractometers: D1A, D2B

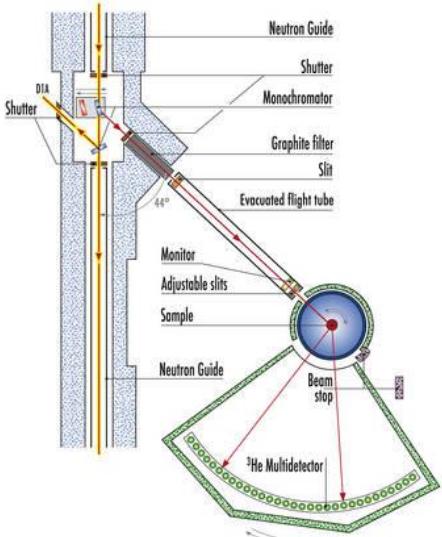
Instrument	D1A	D2B
Takeoff angle / °	122	135
Flux / n cm ⁻² s ⁻¹	10^6	10^6 to 10^7
Beam (h × w) / mm	30×20	50×20
Detectors	$25 \text{ } ^3\text{He} \times 10 \text{ cm h}$	$128 \text{ } ^3\text{He} \times 30 \text{ cm h}$
Wavelengths	Ge(hhl)	Ge(hhl)
Δd/d Resolution	$2-3 \times 10^{-3}$	Min 5×10^{-4}
Background	Very Low (60 m)	Low (15 m)
Average data collection time	3-24 hrs	0.25-4 hrs

Similar instruments at all continuous sources:

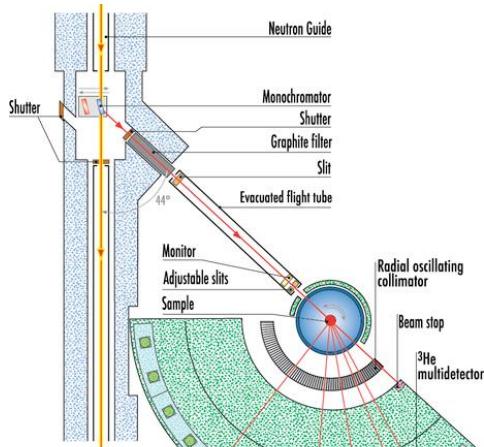
Echidna (ANSTO), Spodi (FRM-II), BT-1 (NIST), 3T2 (LLB), HB-2A (HFIR) etc...

CW, High flux, thermal powder diffractometers: D1B, D20

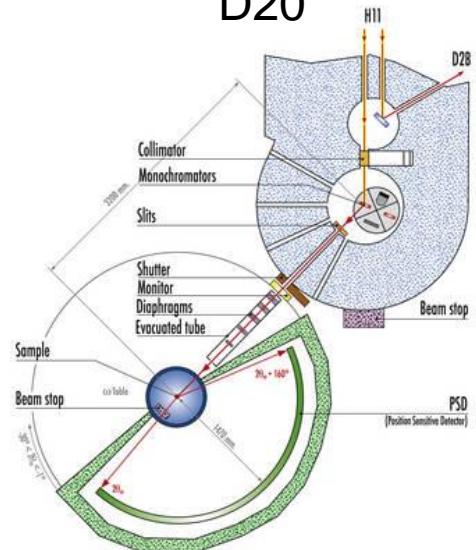
D1B



newD1B



D20



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CW, High flux, thermal powder diffractometers: D1B, D20

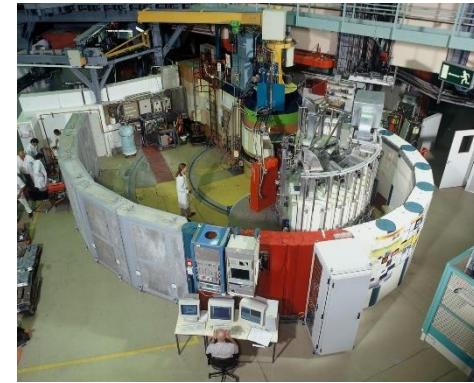
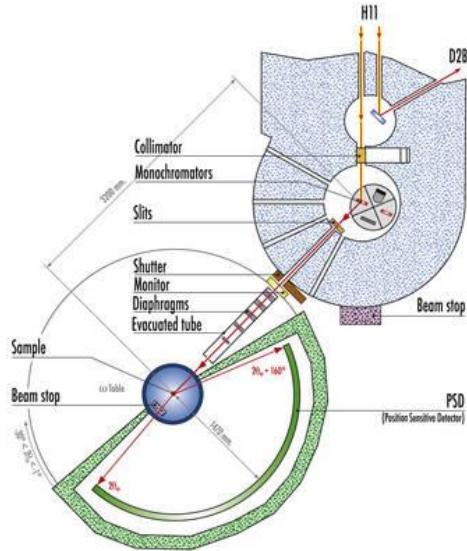
Instrument	D1B (old) / D1B (new)	D20
Takeoff angle / °	44	28, 42 (± 2)
Flux / $n\text{ cm}^{-2}\text{ s}^{-1}$	6.5×10^6 HOPG(002) 0.4×10^6 Ge(311)	4.2×10^7 HOPG(002) 9.8×10^7 Cu(200) 42° 3.2×10^7 Cu(200) 28°
Beam (h × w) / mm	50 × 20	50 × 20
Detectors	80° multi-wire / 128° multi-wire 0.2° separation / 0.1° separation 400 channels / 1280 channels	153.6° micro-strip detector 0.1° separation 1536 channels
Wavelengths / Å	2.52, 1.28	2.42, 1.30, 0.87
Δd/d Resolution	$> 1 \times 10^{-2}$	$> 1 \times 10^{-2}$
Background	Medium (low with ROC)	Medium/High (low with ROC)
Average data collection time	5-10 mins / 1-5 mins	<1 min

Fewer comparable instruments: Wombat (ANSTO), G4.1 (LLB), HB-2C (HFIR)

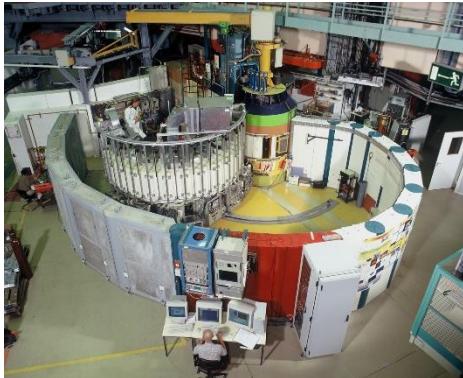
CW, variable resolution, thermal powder diffractometer: D20



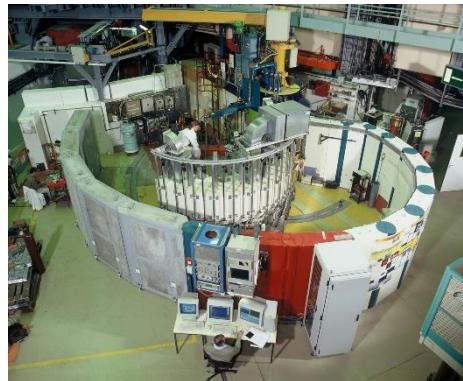
120°



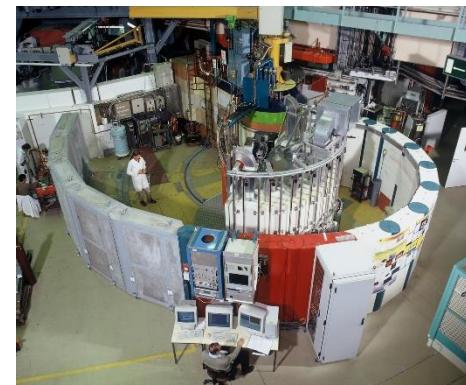
28°



90°



65°



42°

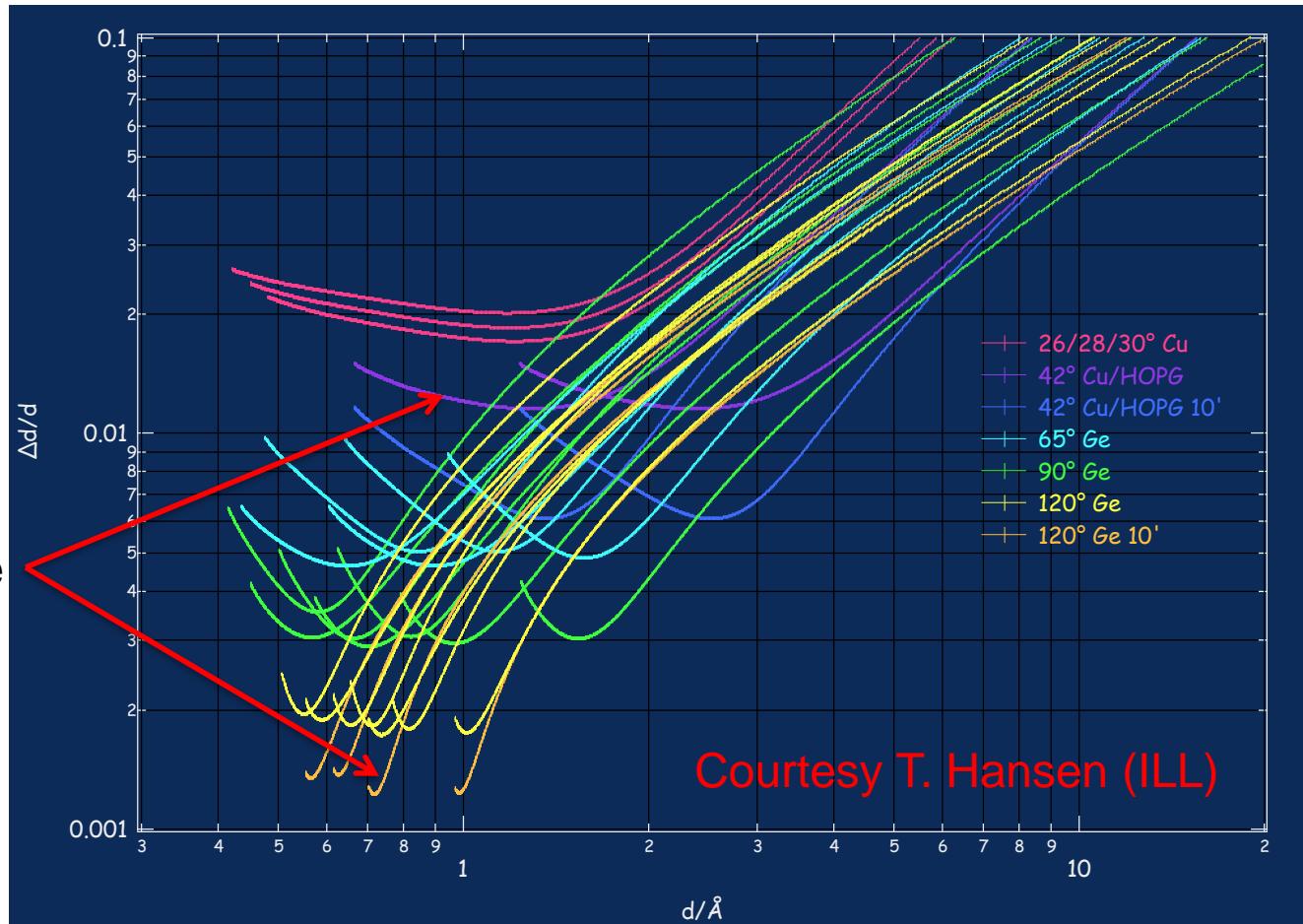
Photos courtesy T. Hansen ILL

CW, variable resolution, thermal powder diffractometer: D20

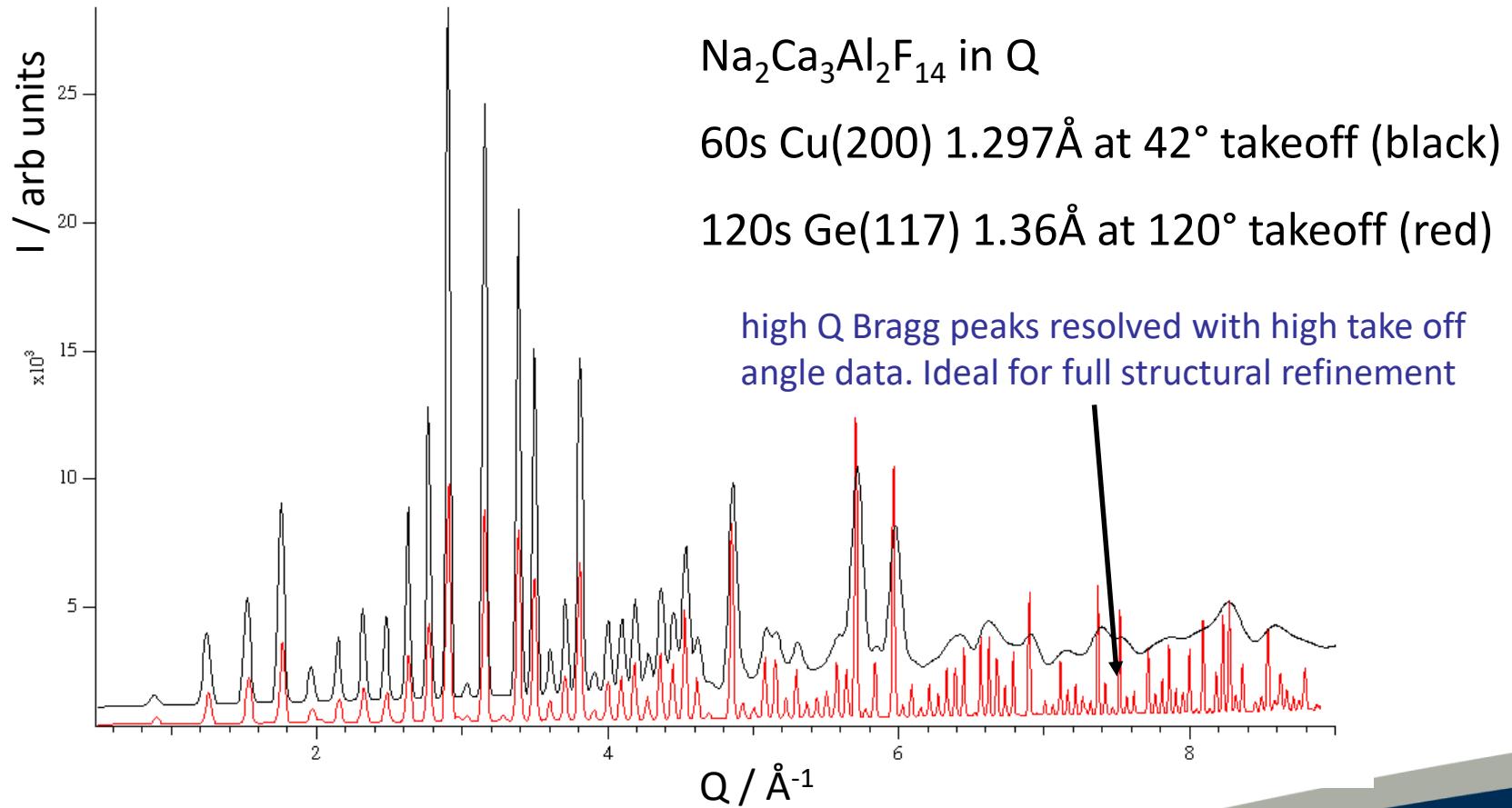
Instrument	D20 (high flux)	D20 (high takeoff angle)
Takeoff angle / °	28, 42 (± 2)	65, 90, 120 (± 2)
Flux / $n\text{ cm}^{-2}\text{ s}^{-1}$	4.2×10^7 HOPG(002) 9.8×10^7 Cu(200) 42° 3.2×10^7 Cu(200) 28°	8.0×10^6 Ge(115) 7.5×10^6 Ge(117) 4.0×10^6 Ge(119)
Beam (h × w) / mm	50 × 20	50 × 20
Detectors	153.6° micro-strip detector 0.1° separation 1536 channels	153.6° micro-strip detector 0.1° separation 1536 channels
Wavelengths / Å	2.42, 1.30, 0.87	variable 0.8-3 Ge(hhl/00l/hhh)
Δd/d Resolution	$> 1 \times 10^{-2}$	See next slide
Background	Medium/High (low with ROC)	Medium/High (low with ROC)
Average data collection time	<1 min	5-15 mins (30 times faster than similar counting statistics on D2B)

Few contemporary instruments: HRPT (PSI), Wombat (ANSTO has potential)

Tune resolution using θ_B

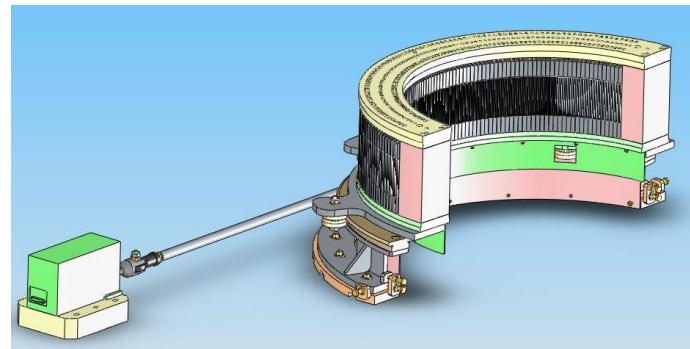
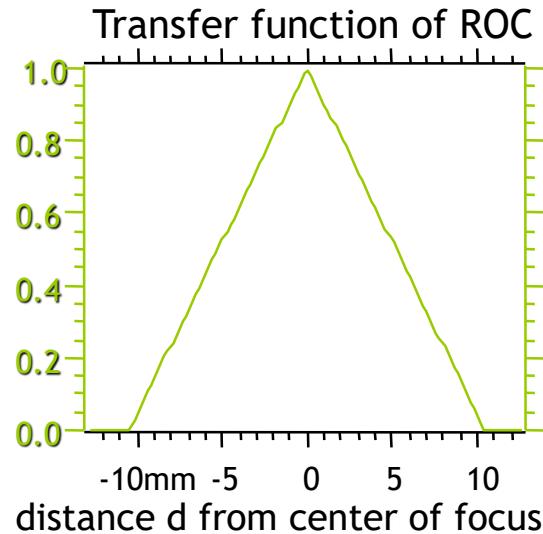
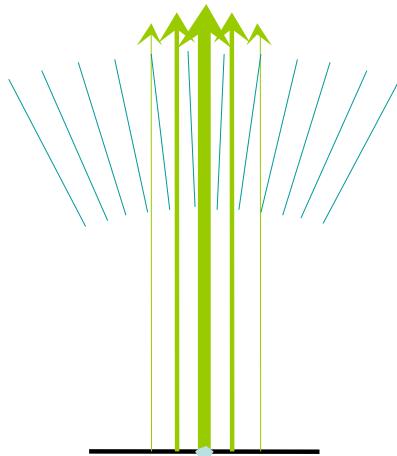


Low θ_B v high θ_B : Q resolution v count-rate

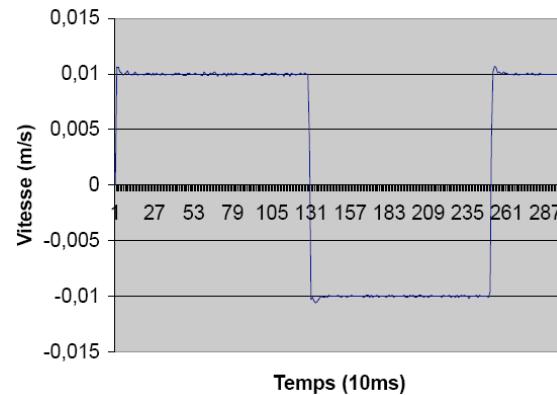
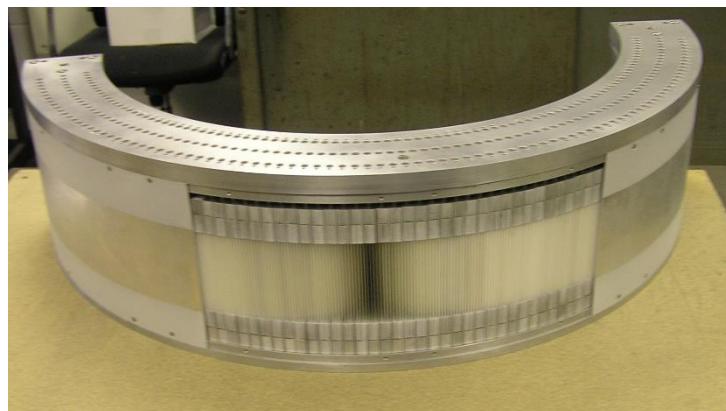


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Parasitic scattering on instruments with area detectors: D20 collimator

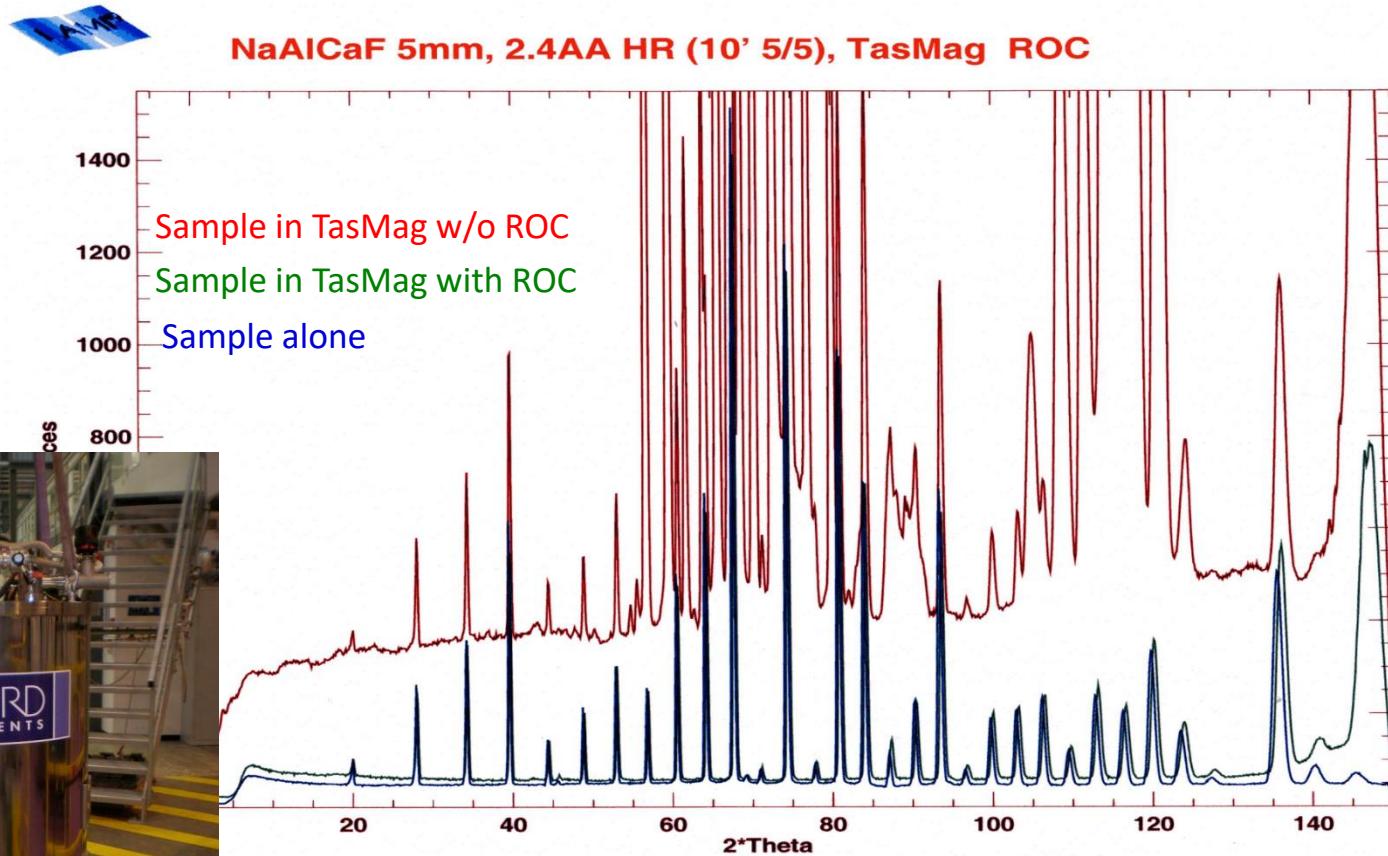


Photos Courtesy ILL



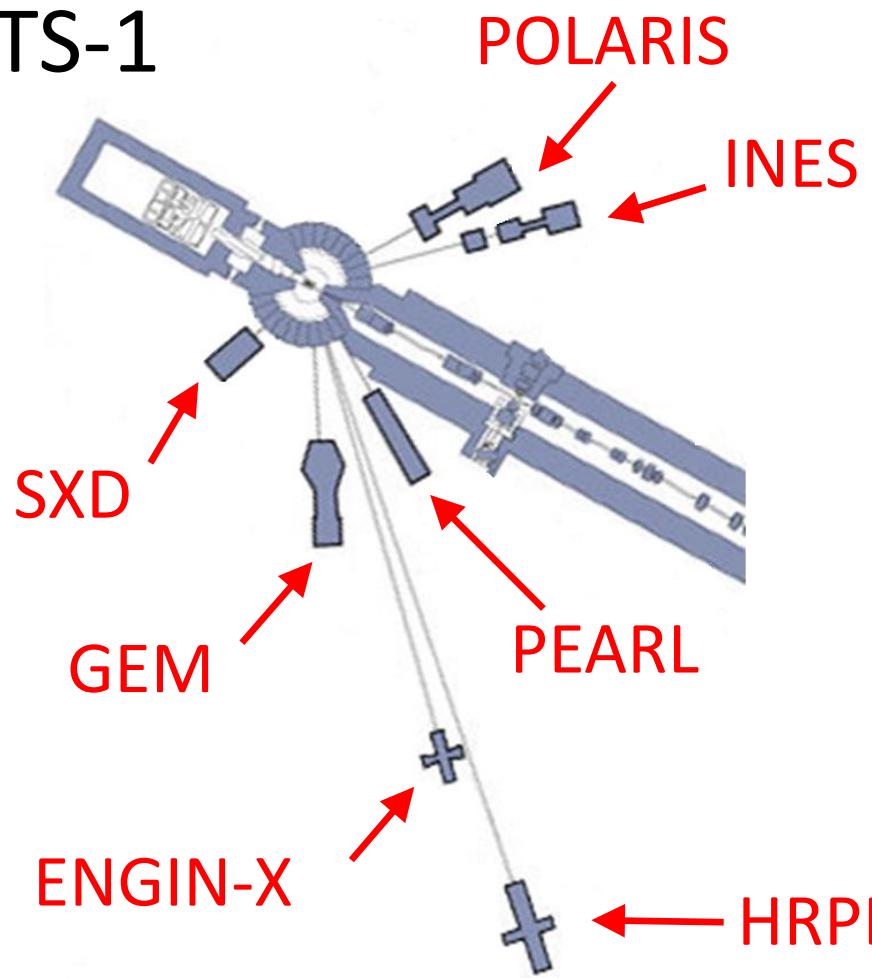
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Parasitic scattering on instruments with area detectors: D20 example

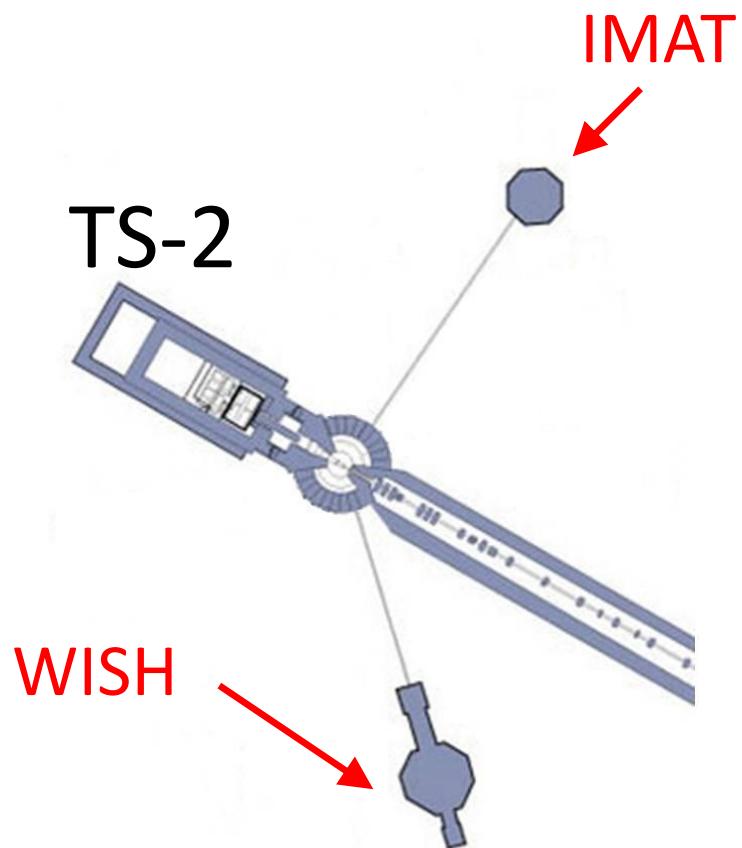


Diffraction at a pulsed source: ISIS

TS-1



TS-2

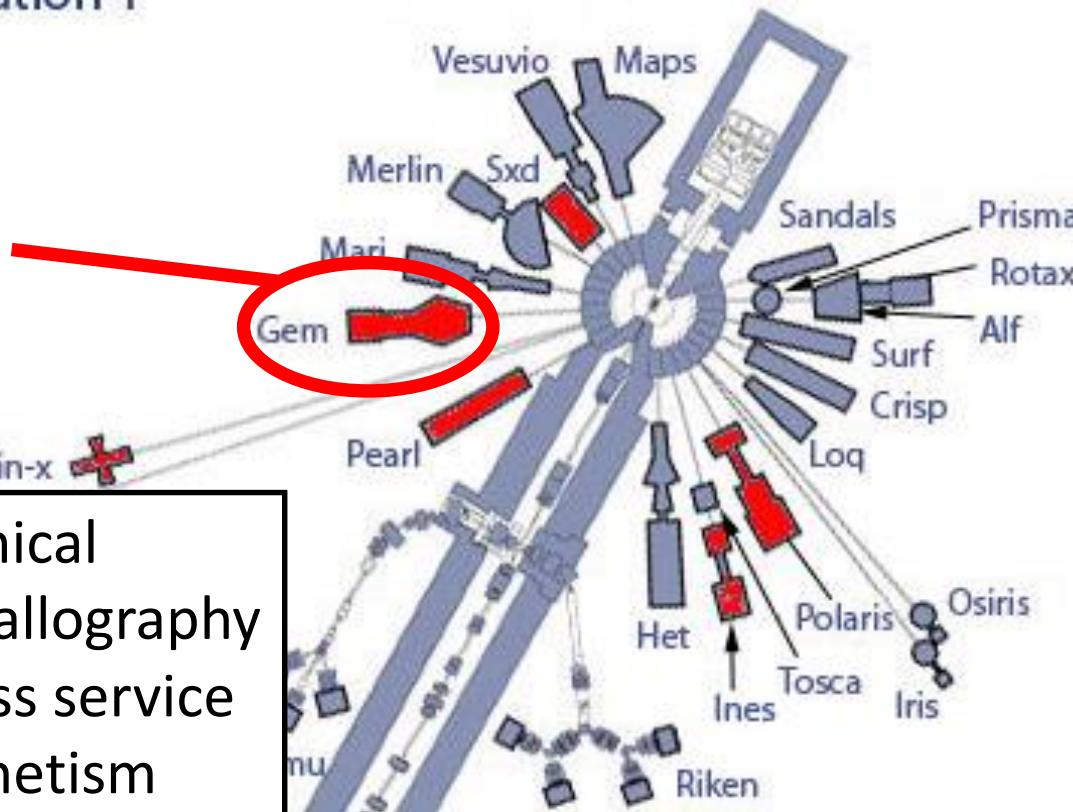


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GEM: high intensity powder diffraction

Target Station 1

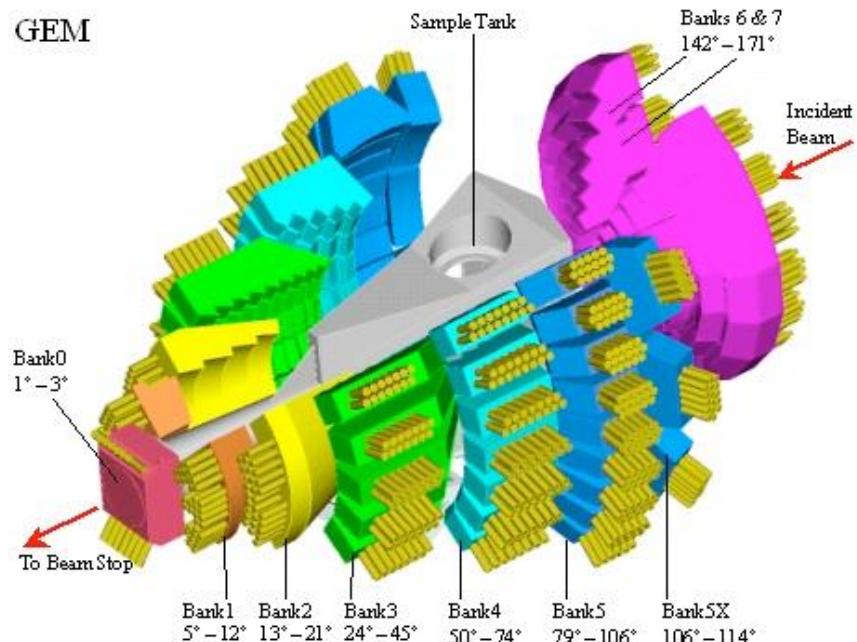
GEM



- chemical crystallography
- Xpress service
- magnetism
- PDF studies
- glasses

GEM: high intensity powder diffraction

GEM

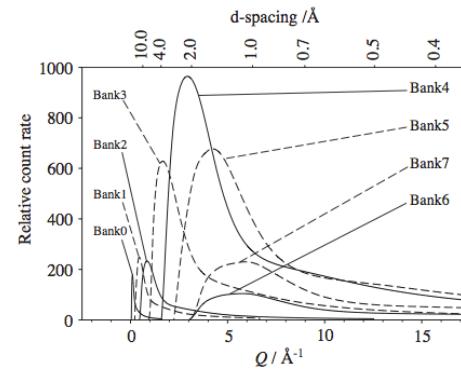
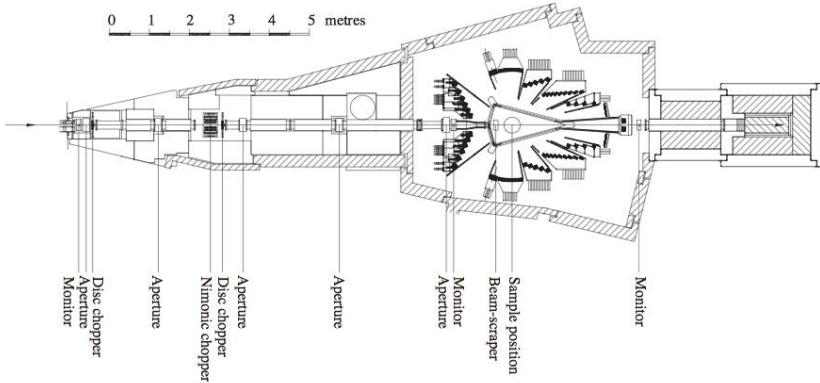


Initially constructed in the late 1990s this powder/liquids diffractometer hybrid changed the way TOF diffraction instruments were designed and built

http://wwwisis2.isis.rl.ac.uk/disordered/gem/gem_home.htm

GEM: high intensity powder diffraction

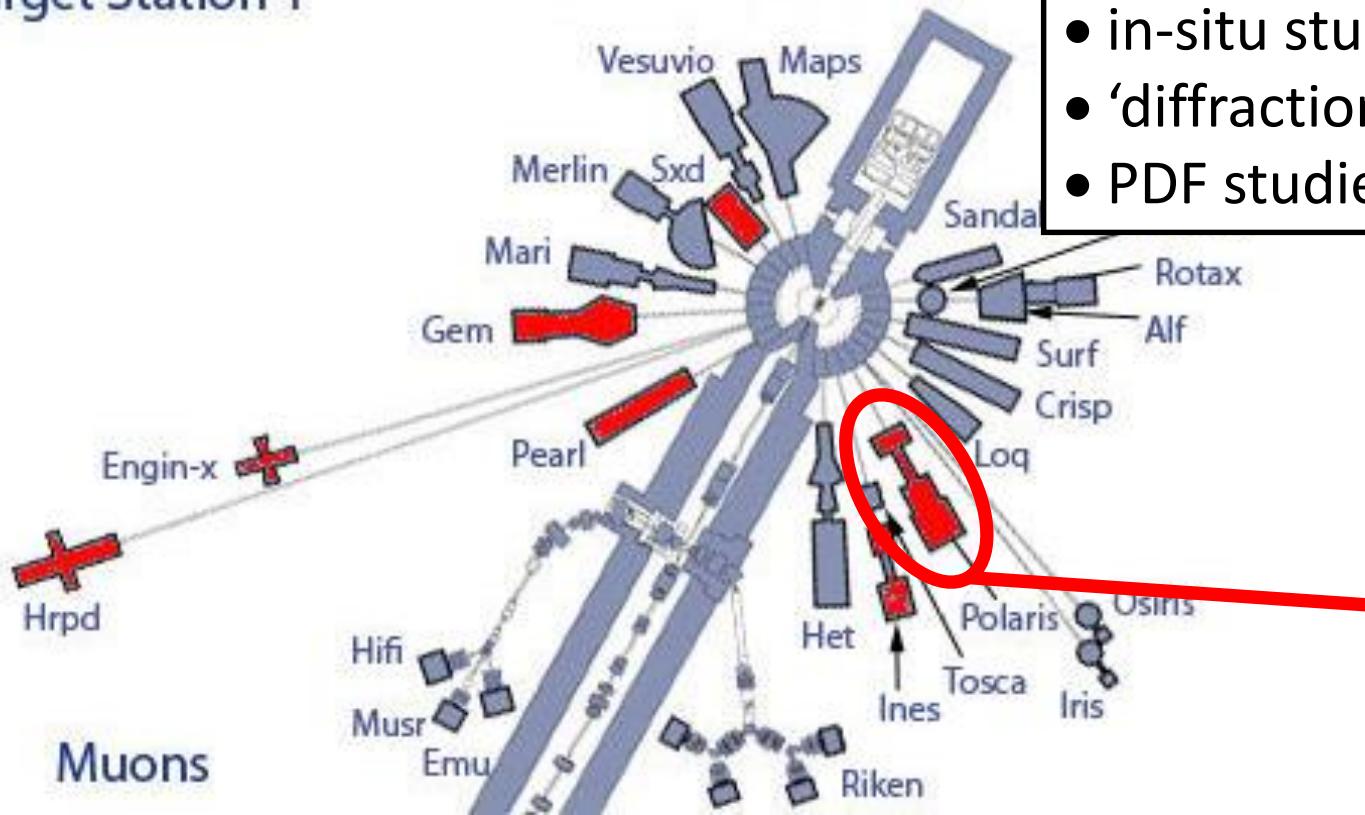
Detector Bank	Scattering angle 2θ (deg)	Range in azimuthal angle ϕ (deg)	Secondary flight path L_2 (m)	Number of detector elements/modules	Solid angle Ω (sr)	Resolution $\Delta Q/Q(\%)$	Minimum accessible momentum transfer Q_{\min} (\AA^{-1})
Bank0	1.21–3.18	± 90.0	2.757–2.767	80/4	0.008	5–10	0.04
Bank1	5.32–12.67	± 45.0	2.365–2.376	330/6	0.056	4.7	0.17
Bank2	13.44–21.59	± 43.4	1.477–2.100	320/4	0.093	2.4	0.43
Bank3	24.67–45.61	± 42.5	1.077–1.893	900/10	0.478	1.7	0.79
Bank4	50.07–74.71	± 44.4	1.028–1.436	1400/14	0.988	0.79	1.56
Bank5	79.07–106.60	± 44.5	1.376–1.383	2160/18	1.135	0.51	2.35
Bank5X	106.02–114.19	± 42.7	1.377–1.387	720/18	0.378	0.5	2.95
Bank6	142.50–149.72	± 69.3	1.544–1.738	560/14	0.280	0.34	3.50
Bank7	149.98–171.40	± 66.6	1.035–1.389	800/10	0.443	0.35	3.57



From Hannon, NIMA (2005), 551, 88-107

Polaris: high intensity powder diffraction

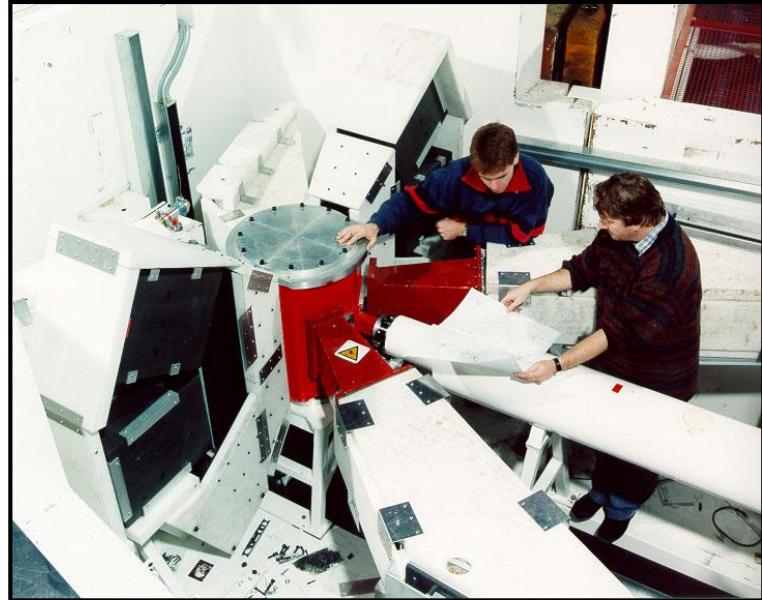
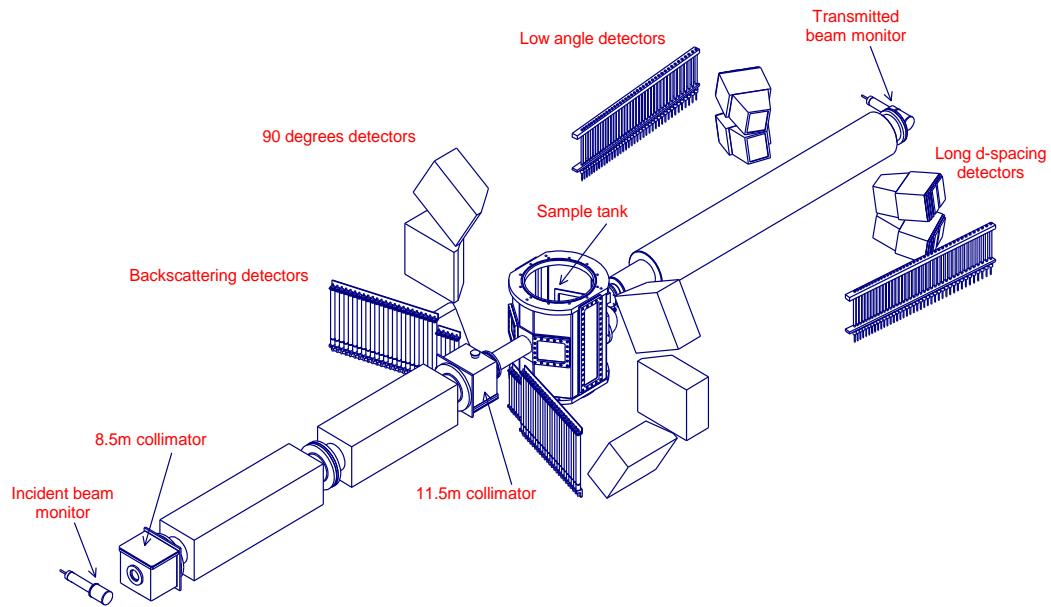
Target Station 1



- chemical crystallography
- in-situ studies
- 'diffraction plus'
- PDF studies

Polaris

Polaris: old configuration



Compare with GEM:

- Higher sample flux
- Wider bandwidth
- Lower resolution
- Hotter spectrum
- Lower detector coverage

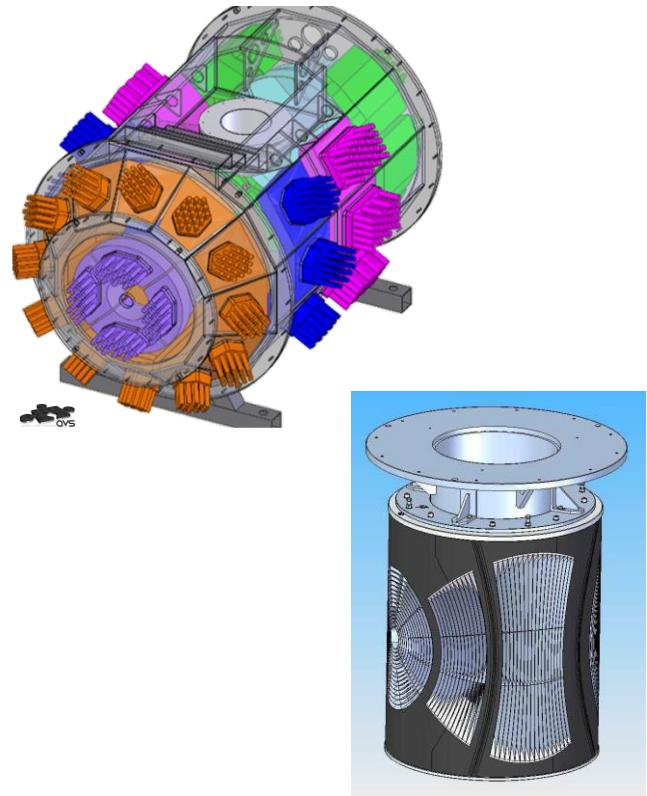
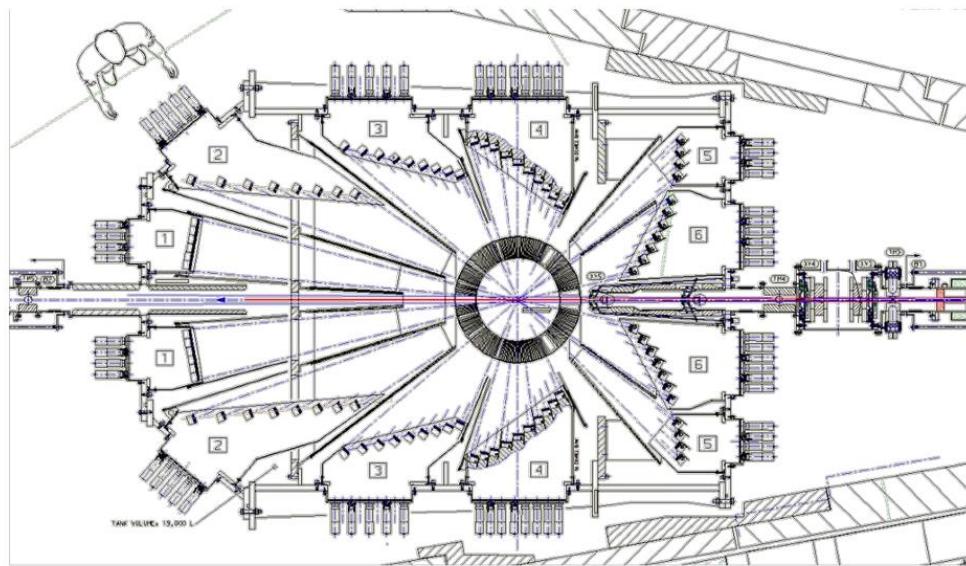
<http://www.isis.stfc.ac.uk/instruments/polaris/polaris4643.html>

Polaris: old configuration

bank position (label)	low angle (A)	low angle (B)	backscattering (C)	90 degrees (E)
detector type	^3He	ZnS	^3He	ZnS
no. of elements	$2 \times 40 = 80$	$4 \times 20 = 80$	$2 \times 29 = 58$	$6 \times 36 = 216$
L_2 (m)	1.72 - 2.65	~2.2	0.65 - 1.35	~0.80
2θ range	$28^\circ < 2\theta < 42^\circ$	$13^\circ < 2\theta < 15^\circ$	$130^\circ < 2\theta < 160^\circ$	$83^\circ < 2\theta < 97^\circ$
Ω (ster)	0.046	0.009	0.29	0.48
$\Delta d/d$	$\sim 1 \times 10^{-2}$	$\sim 3 \times 10^{-2}$	$\sim 5 \times 10^{-3}$	$\sim 7 \times 10^{-3}$
d -range (\AA)	0.5 - 8.3	0.5 - 21.6	0.2 - 3.2	0.2 - 4.0
Q-range (\AA^{-1})	0.75 - 12.6	0.3 - 12.6	2.0 - 31.4	1.5 - 31.4

- Good workhorse instrument for powder diffraction
- High Q accessible for disordered materials investigation using the PDF method
- Some in situ capability but limited by count-rate
- Compatible with a wide range of restricted geometry sample environment

Polaris upgrade



- Increase primary flight path to 14 m
- Optimise each detector bank to give constant resolution
- Increase detector coverage
- Design a collimator to reduce background and parasitic scattering

<http://www.isis.stfc.ac.uk/instruments/polaris/polaris-upgrade-poster11575.pdf>

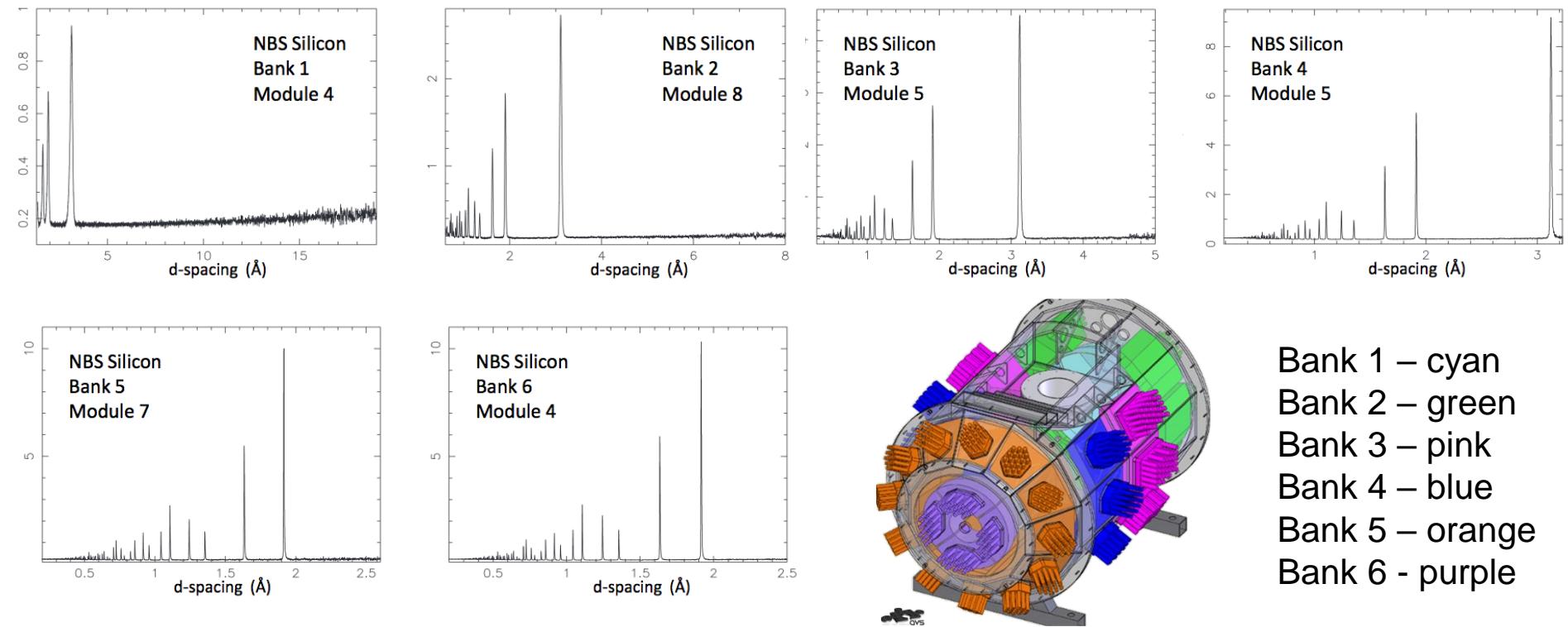
Polaris upgrade



<http://www.isis.stfc.ac.uk/instruments/polaris/polaris4643.html>

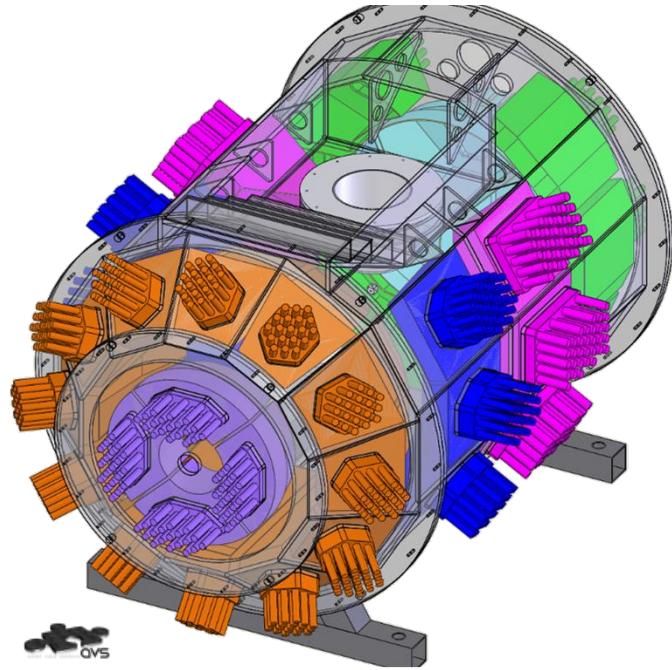
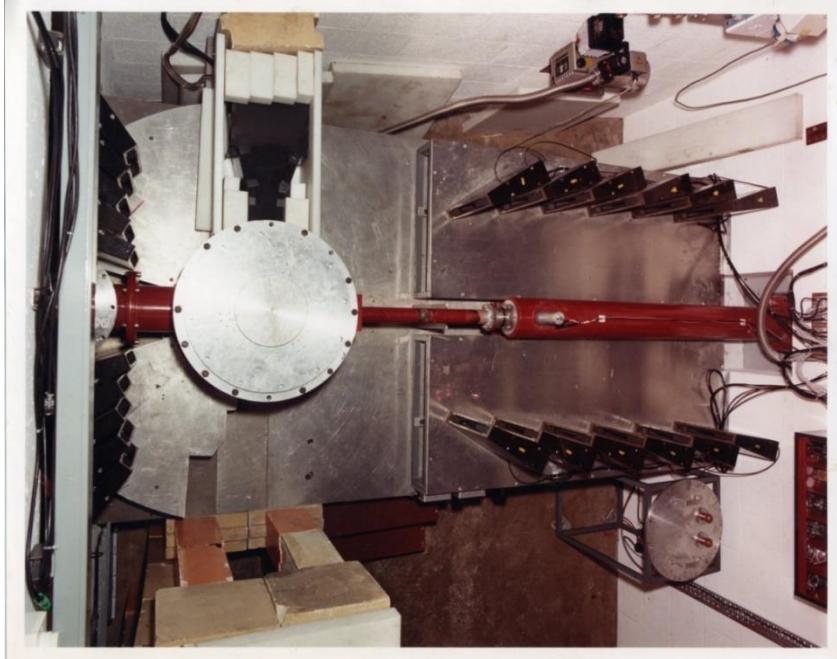
Current Polaris

<http://www.isis.stfc.ac.uk/instruments/polaris/polaris-upgrade---first-diffraction-pattern12763.pdf>



- Increased count rate $\times 3$ at high scattering angle to >20 for low angle banks
- Resolution improvement e.g. bank 5 and 6 of 3×10^{-3} cf. 5×10^{-3}
- Improvement in data at high Q

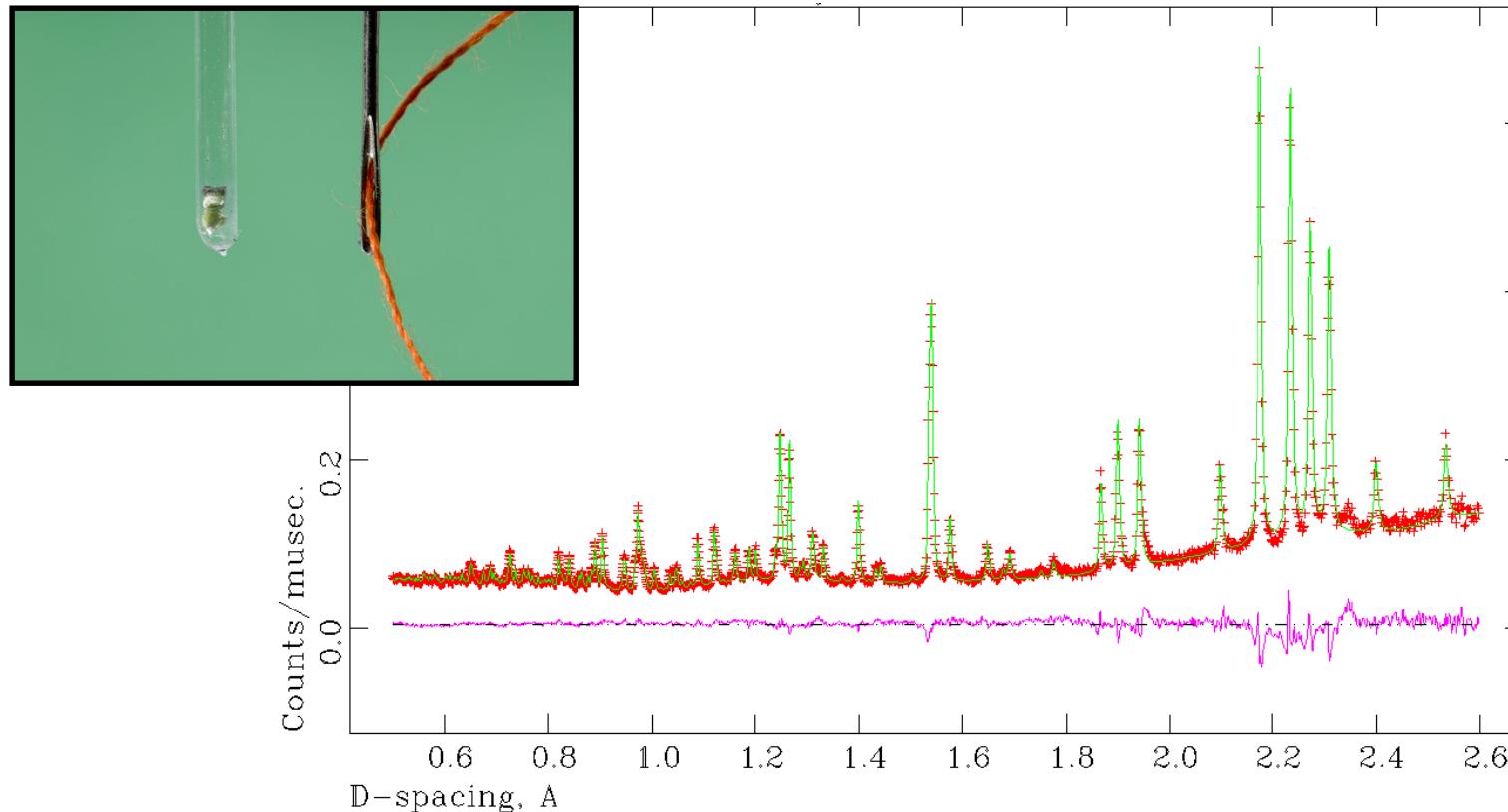
Polaris 1995 v 2013



- 1995 500 mg 24+ hrs
- 2013 500 mg 15-20 minutes with increased Q-range

Contemporary instruments NOVA and iMateria (J-PARC), POWGEN-3 (SNS)

Polaris: pushing boundaries in sample size

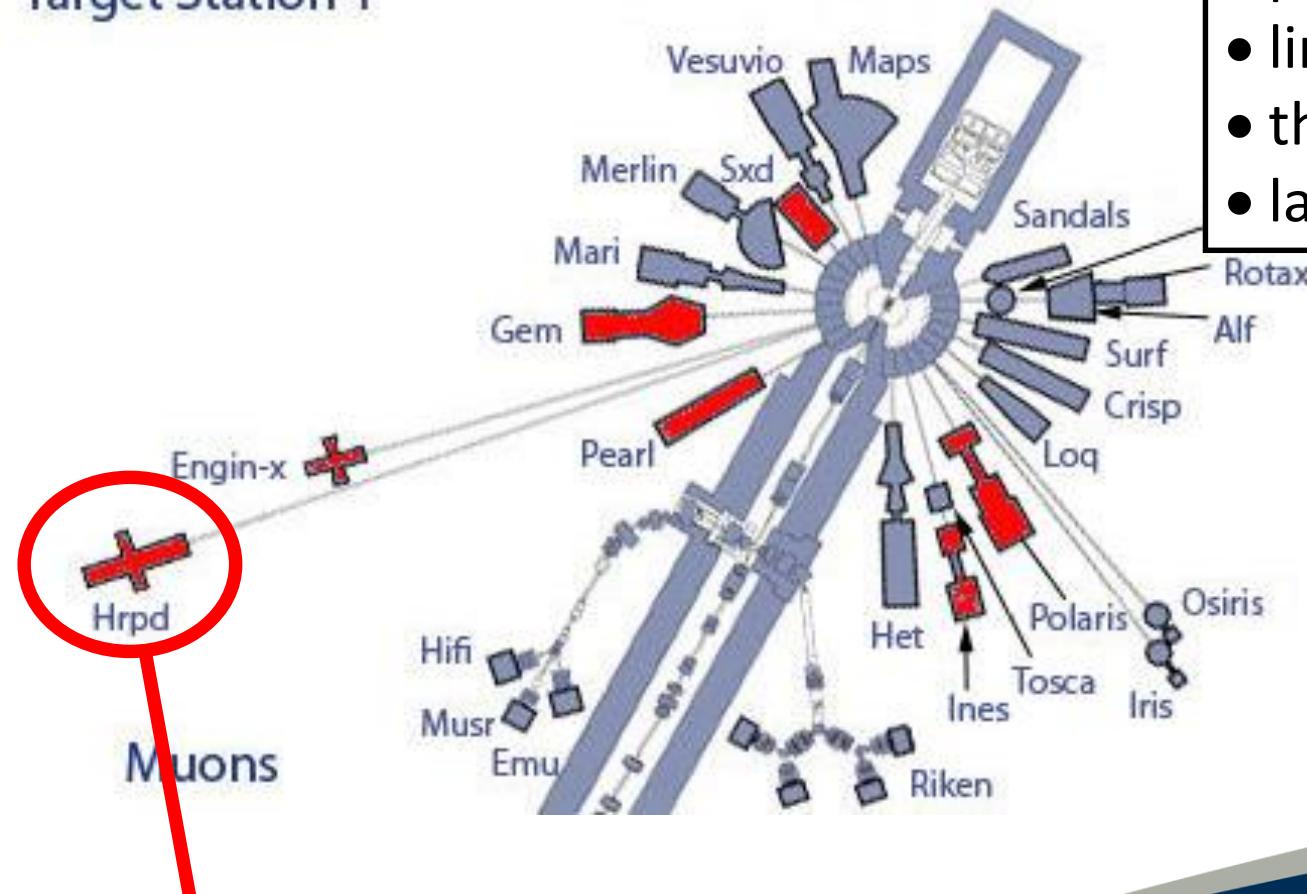


~1mm³ sample of NaNiF₃ phase prepared at high *p* + high T

Lindsay-Scott *et al*, J. Appl. Cryst., **47**, 1939 (2014)

HRPD: high resolution powder diffraction

Target Station 1

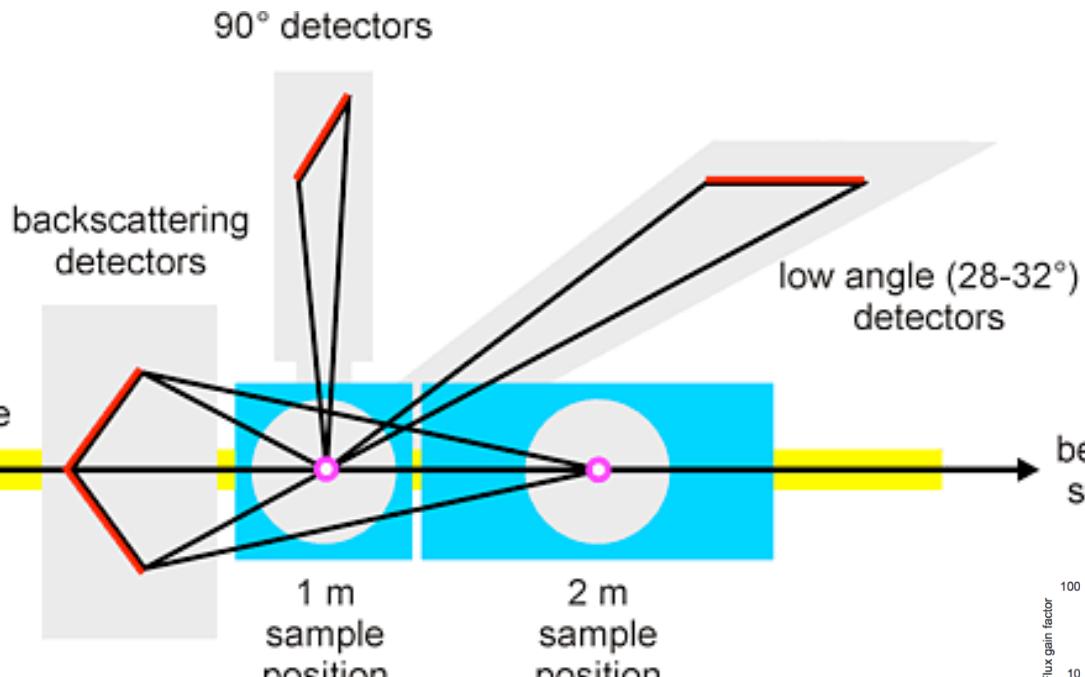


- phase transitions
- line broadening
- thermal expansion
- large unit cells

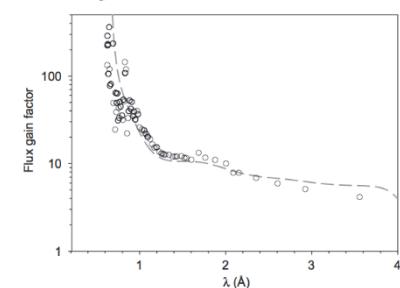
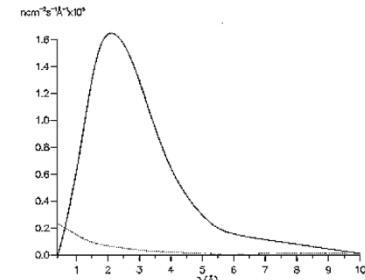


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HRPD: high resolution powder diffraction



- Approximately 100 m primary flight path
- 100K methane moderator – peak flux around 2 \AA
- Operates at 5 or 10 Hz cf. TS-1 50 Hz



<http://www.isis.stfc.ac.uk/instruments/hrpd/hrpd.html>

HRPD: high resolution powder diffraction

Table 1. HRPD Detector Bank Details

	Backscattering	90°	Low Angle
Detector Specification	ZnS scintillator	ZnS scintillator	½" 10atm He ³ gas tubes
Geometry	60 rings: $7 < r_1 < 8.5\text{cm}$ $35.5 < r_{60} < 37\text{cm}$ 8 Octants: 4147cm^2	Slab: 20 x 20cm 66 x 3mm elements 6 Modules: 2400cm^2	72 tubes: (20cm active length) 8 tubes/module 9 Modules: 1800cm^2
Fixed Scattering Angle	$160^\circ < 2\theta < 176^\circ$ (1m)	$87^\circ < 2\theta < 93^\circ$	$28^\circ < 2\theta < 32^\circ$
Solid Angle (Ω)	0.41 ster (1m)	0.08 ster	0.01 ster
Resolution ($\Delta d/d$)	$\sim 4.5 \times 10^{-4}$	$\sim 2 \times 10^{-3}$	$\sim 2 \times 10^{-2}$
d-spacing range (30-230ms)	$\sim 0.6 - 4.6\text{\AA}$ $0.25 - 4.6\text{\AA}$	$\sim 0.9 - 6.6\text{\AA}$ $0.4 - 6.6\text{\AA}$	$\sim 2.2 - 16.5\text{\AA}$ $1.0 - 16.5\text{\AA}$

- Large backscattering detector to minimise $\cot^2\theta$ in resolution term
- High resolution at intermediate Q
- Long flight path to reduce $\Delta t/t$ and $\Delta L/L$ uncertainties
- Pulse skipping to increase bandwidth (@50 Hz 0.4\AA)
- Good combination of parameters for a high resolution TOF diffractometer

Contemporary instrument sHRPD (J-PARC), no current analogue at SNS