

# Diffraction from non-crystalline materials



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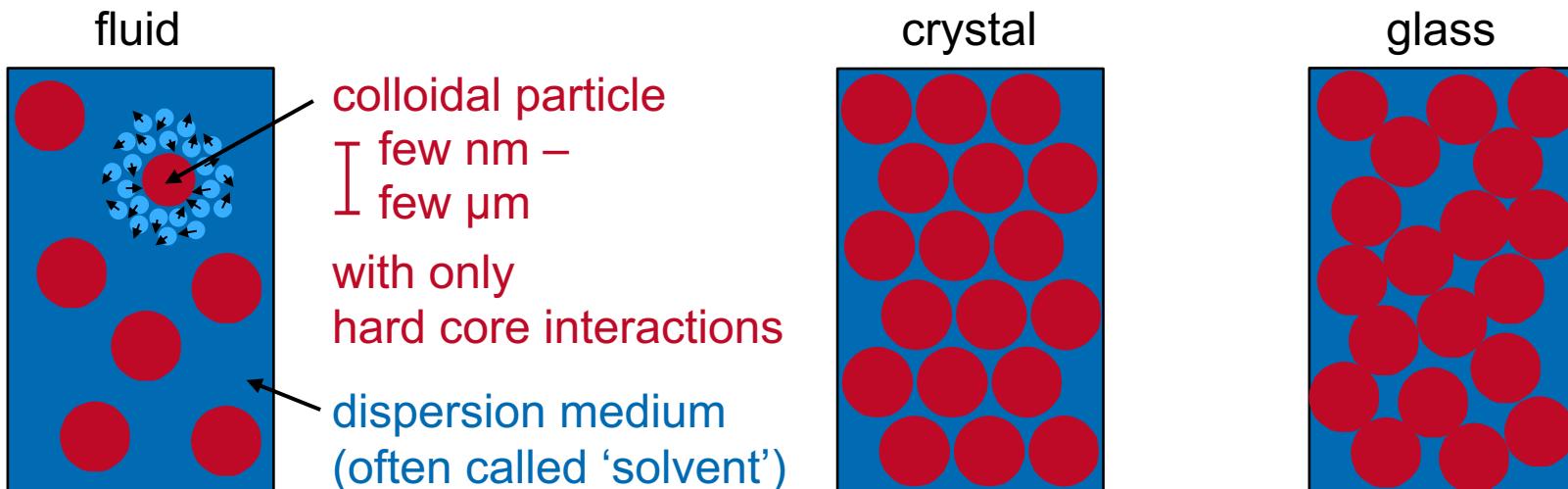
examples of non-crystalline materials

characterization of non-crystalline materials

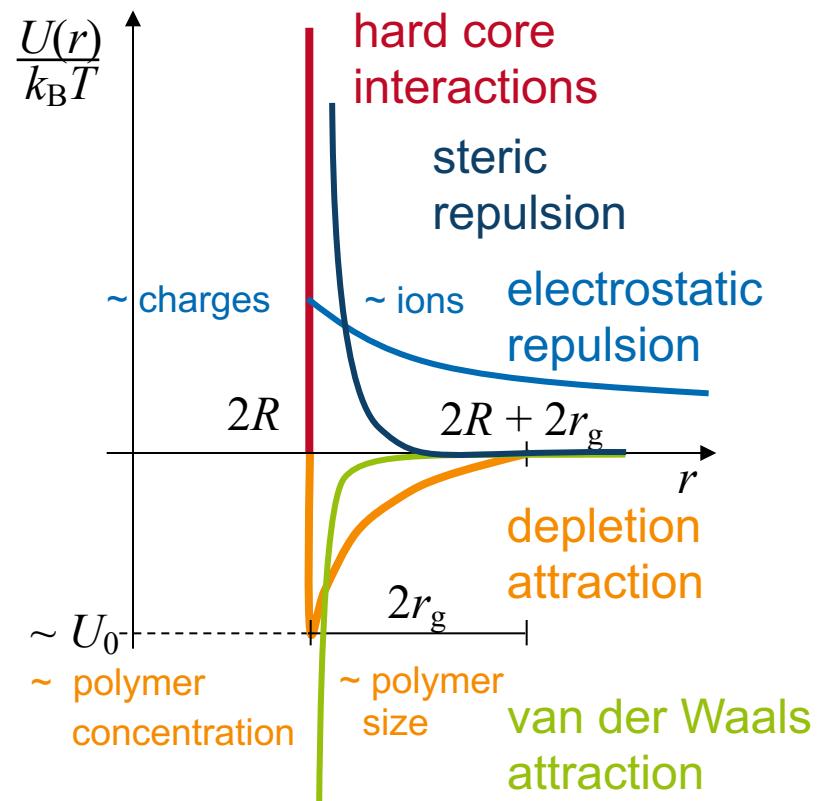
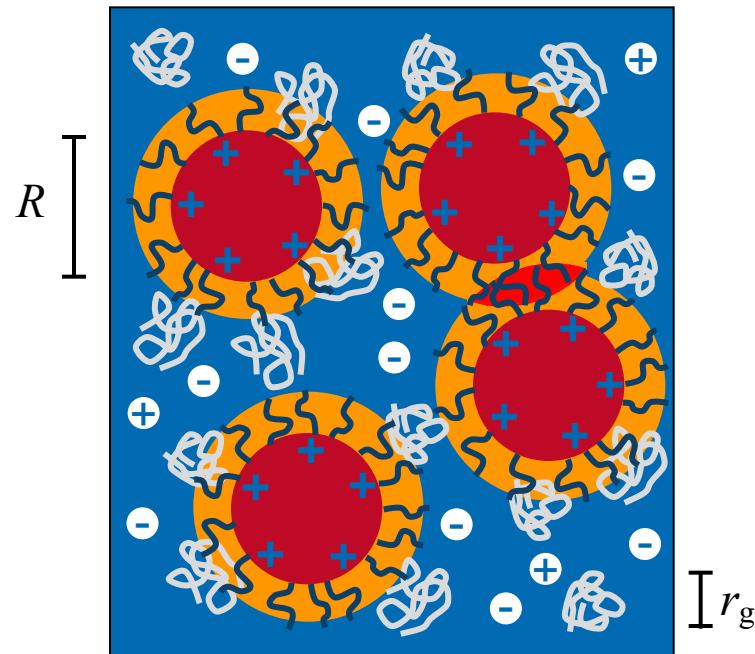
diffraction from non-crystalline materials

- general concept
- form factor – particle properties
- structure factor – particle arrangements
- ‘prefactor’ – contrast variation

# Colloidal Suspensions

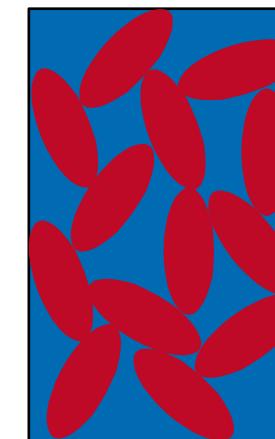
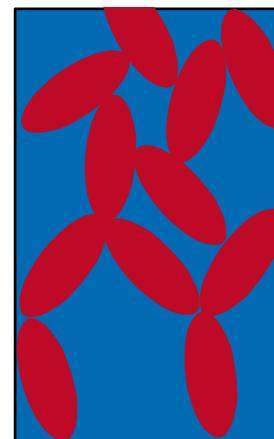
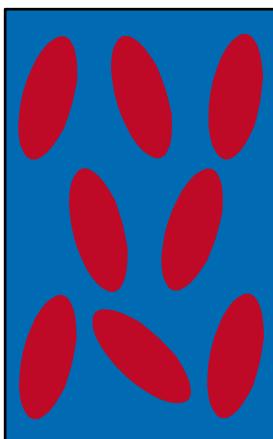
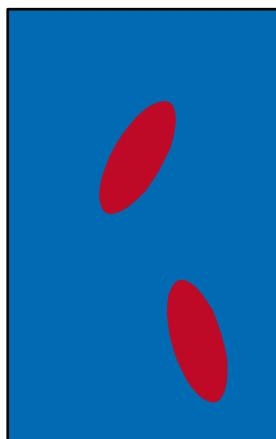
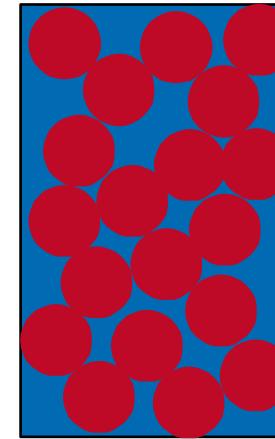
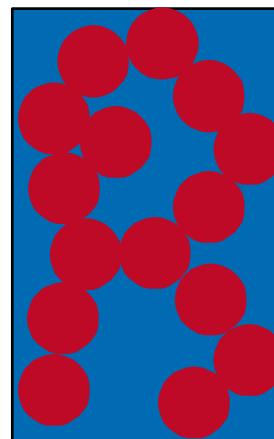
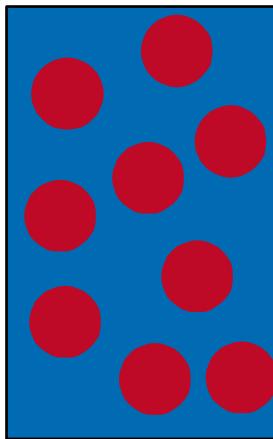
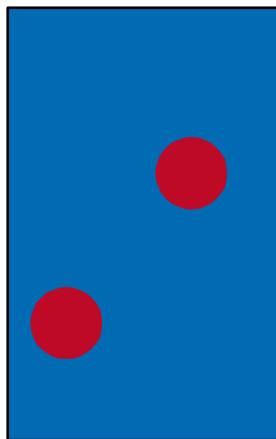


# Particle interactions

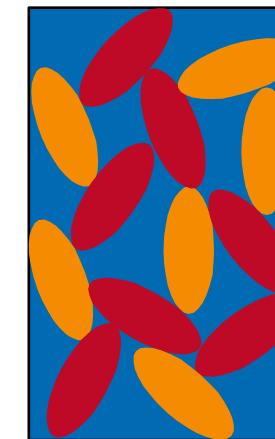
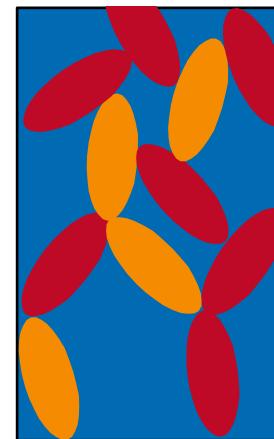
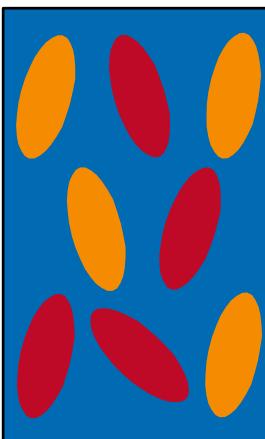
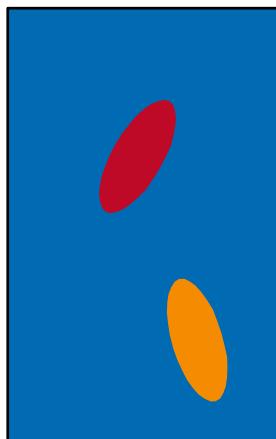
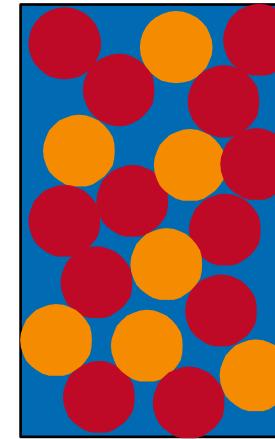
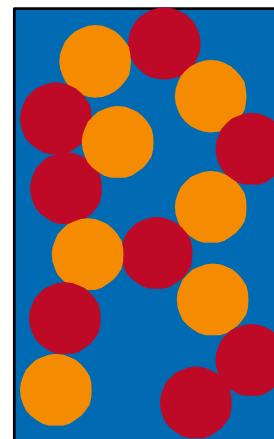
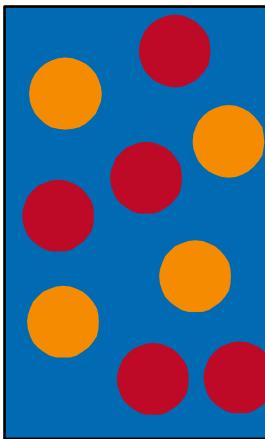
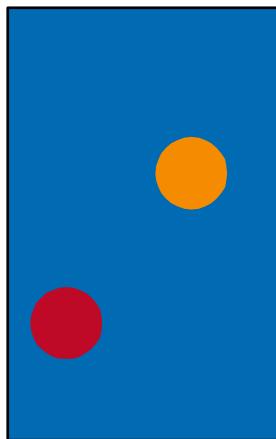


as well as anisotropic interactions,  
(e.g. dipol or 'patchy' interactions)

# Non-crystalline materials



# Non-crystalline materials



examples of non-crystalline materials

characterization of non-crystalline materials

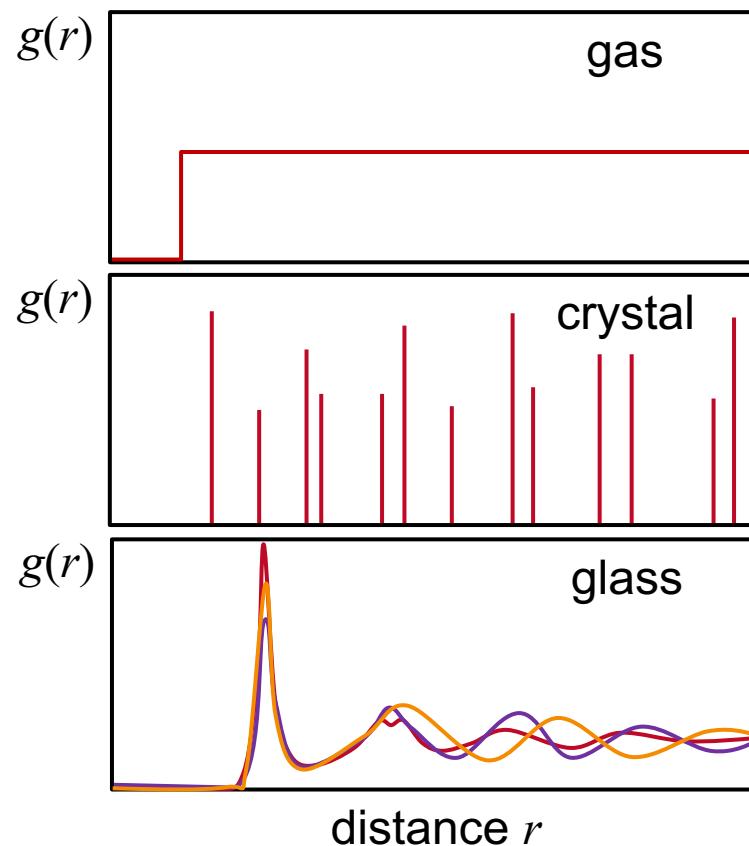
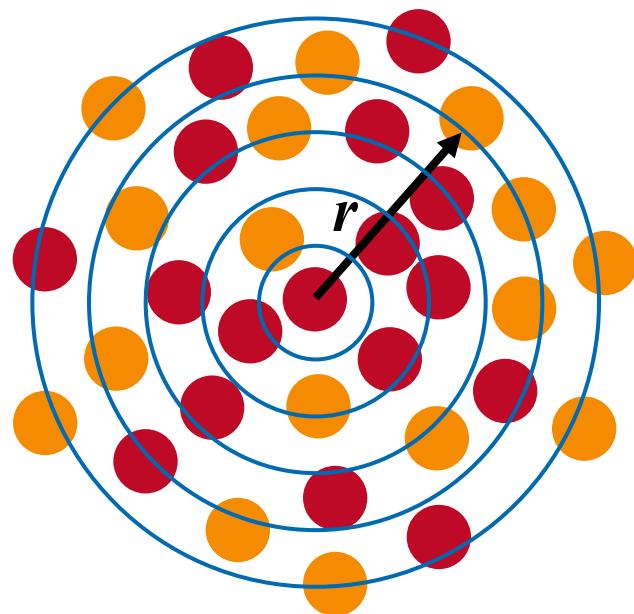
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- structure factor – particle arrangements
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# Characterization - arrangement

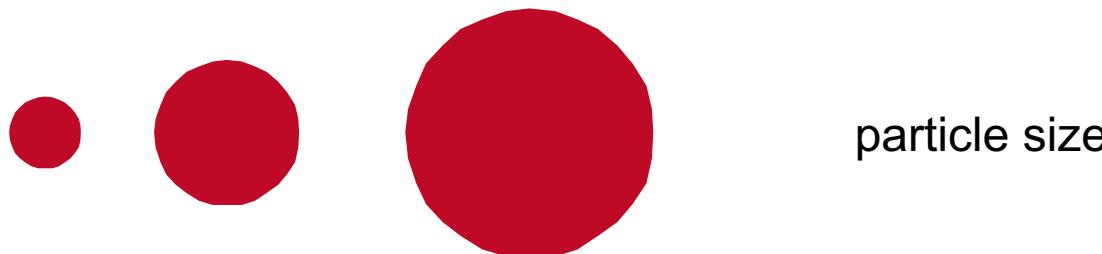
## radial distribution function $g(r)$

$(N/V) g(r) d^3r$  is the number of particles in  $d^3r$  at  $r$



contains information on interactions

# Characterization – particle



particle size



particle shape



particle structure

examples of non-crystalline materials

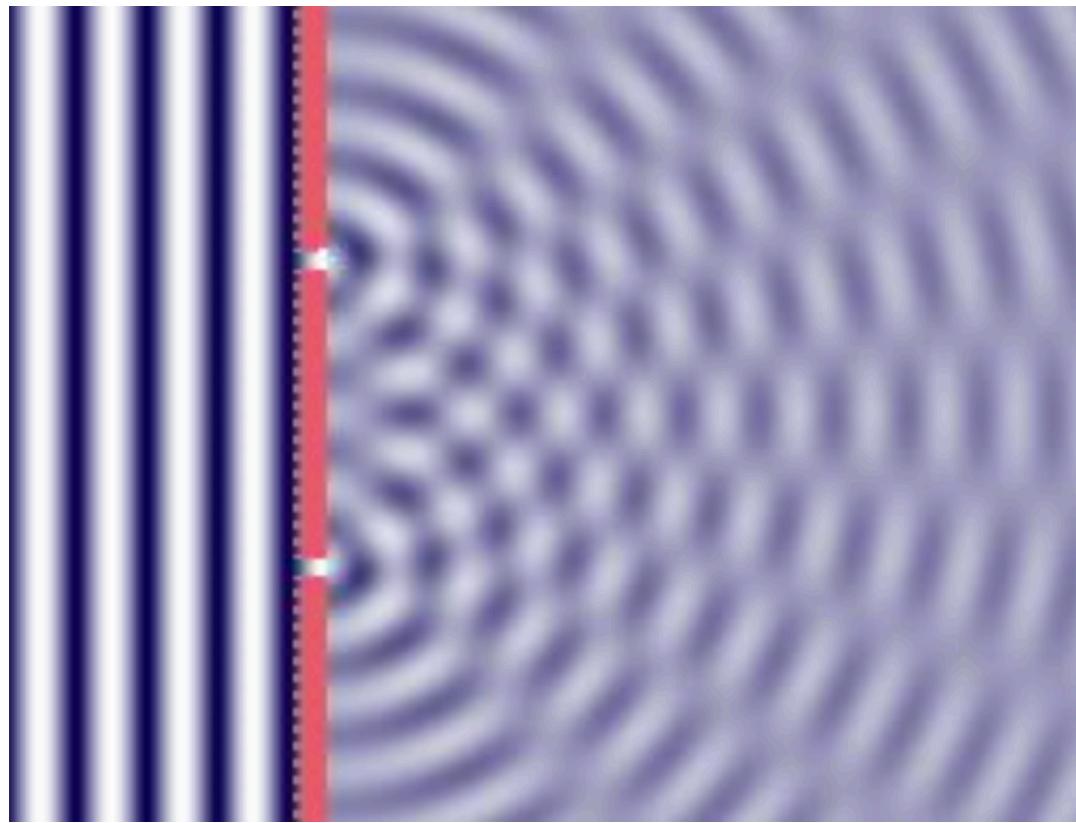
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# Two-slit experiment

hhu



# Scattering by a particle

particle = ensemble of small volume elements  $dV$  (using Born approx.)

$$\psi_p^{sc} = \int_V e^{i\delta\phi} d\psi^{sc} \sim \int_V e^{i\mathbf{Q}\cdot\mathbf{r}} \Delta b(\mathbf{r}) dV$$

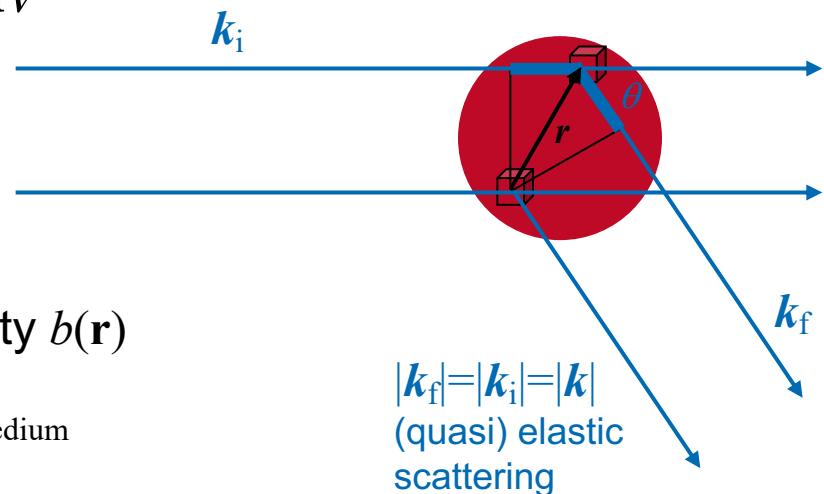
Fourier transform

$$\delta\phi = \frac{2\pi}{\lambda} \delta L = \mathbf{k}_i \cdot \mathbf{r} - \mathbf{k}_f \cdot \mathbf{r} = \mathbf{Q} \cdot \mathbf{r}$$

$d\psi^{sc} \sim \Delta b(\mathbf{r}) dV$  scatt. length density  $b(\mathbf{r})$

$$\Delta b(\mathbf{r}) = b_{\text{part}}(\mathbf{r}) - b_{\text{medium}}$$

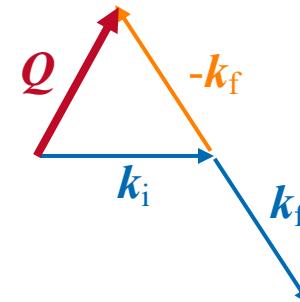
$$\delta L = \delta L_1 + \delta L_2$$



with scattering vector  $\mathbf{Q} = \mathbf{k}_i - \mathbf{k}_f$

$$|\mathbf{Q}| = 2k \sin \frac{\theta}{2} = \frac{4\pi}{\lambda} \sin \frac{\theta}{2} = \frac{4\pi n}{\lambda_0} \sin \frac{\theta}{2}$$

$$|\mathbf{Q}| \sim (\text{length scale})^{-1}$$



# Scattering by many particles

many particles = ensemble of particles  $j=1..N$  (using Born approx.)

$$\psi^{\text{sc}} = \sum_j \psi_{\text{p},j}^{\text{sc}} e^{i\mathbf{Q}\cdot\mathbf{R}_j}$$

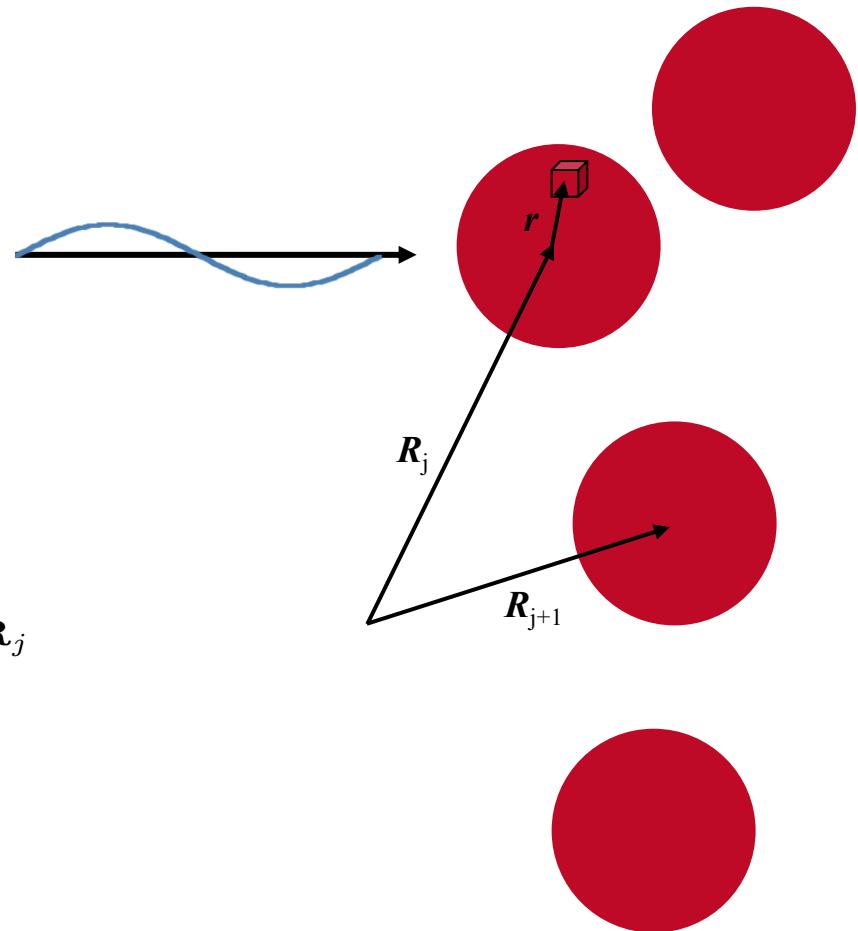
'Fourier transform'

$$\psi_{\text{p},j}^{\text{sc}} \sim \int_{V_j} \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV$$

Fourier transform

$$\psi^{\text{sc}} \sim \sum_j \int_{V_j} \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV e^{i\mathbf{Q}\cdot\mathbf{R}_j}$$

$$\sim \sum_j \int_{V_j} \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot(\mathbf{r}+\mathbf{R}_j)} dV$$



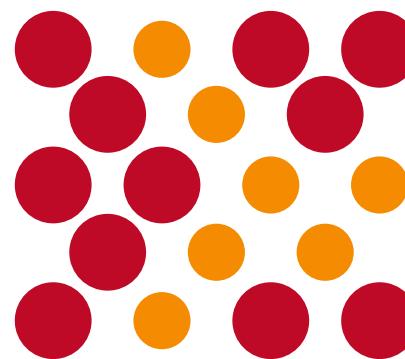
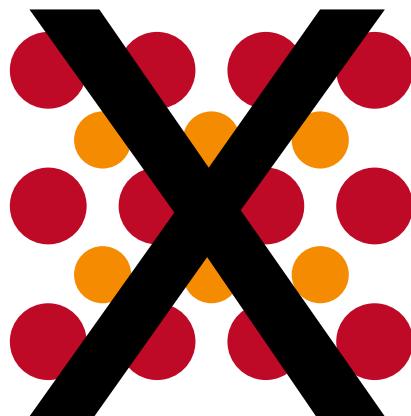
# Differential cross section

time-averaged (or ensemble-averaged) differential cross section is detected

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi^{\text{sc}}|^2 \sim \sum_k \sum_j \left\langle \psi_{p,j}^{\text{sc}}(\mathbf{Q}) \psi_{p,k}^{\text{sc}\star}(\mathbf{Q}) e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \right\rangle$$

assumption: all particles are identical, i.e.  $\psi_{p,j}^{\text{sc}}(\mathbf{Q}) = \psi_{p,k}^{\text{sc}}(\mathbf{Q}) = \psi_p^{\text{sc}}(\mathbf{Q})$

(important: particle properties must not be linked to their positions)



# Differential cross section

time-averaged (or ensemble-averaged) differential cross section is detected

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi^{\text{sc}}|^2 \sim \sum_k \sum_j \left\langle \psi_{p,j}^{\text{sc}}(\mathbf{Q}) \psi_{p,k}^{\text{sc}*}(\mathbf{Q}) e^{i\mathbf{Q}\cdot(\mathbf{R}_j - \mathbf{R}_k)} \right\rangle$$

assumption: all particles are identical, i.e.  $\psi_{p,j}^{\text{sc}}(\mathbf{Q}) = \psi_{p,k}^{\text{sc}}(\mathbf{Q}) = \psi_p^{\text{sc}}(\mathbf{Q})$

(important: particle properties must not be linked to their positions)

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi_p^{\text{sc}}(\mathbf{Q})|^2 \sum_k \sum_j \left\langle e^{i\mathbf{Q}(\mathbf{R}_j(t) - \mathbf{R}_k(t))} \right\rangle$$

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim N |\psi_p^{\text{sc}}(0)|^2 \frac{|\psi_p^{\text{sc}}(\mathbf{Q})|^2}{|\psi_p^{\text{sc}}(0)|^2} \frac{1}{N} \sum_k \sum_j \left\langle e^{i\mathbf{Q}(\mathbf{R}_j(t) - \mathbf{R}_k(t))} \right\rangle$$

‘prefactor’  
scattering  
by  $N$  particles  
(random walk)

$= P(\mathbf{q})$   
form factor  
(intraparticle)

$= S(\mathbf{q})$   
structure factor  
(interparticle)

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# Form factor $P(Q)$

**intraparticle interference**  
depends on particle size and shape

$$P(\mathbf{Q}) = \frac{|\psi_p^{\text{sc}}(\mathbf{Q})|^2}{|\psi_p^{\text{sc}}(0)|^2}$$

with  $\psi_p^{\text{sc}}(\mathbf{Q}) \sim \int_V \Delta b(\mathbf{r}) e^{i\mathbf{Q} \cdot \mathbf{r}} dV$

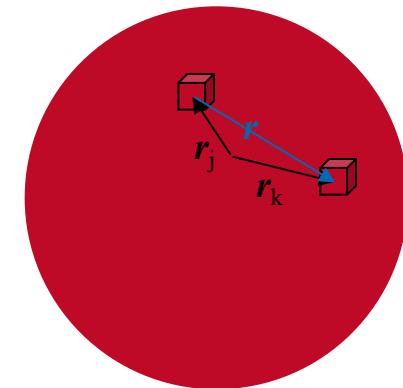
⋮

$$|\psi_p^{\text{sc}}(\mathbf{Q})|^2 = 4\pi \int_0^\infty p(r) \frac{\sin(qr)}{qr} dr$$



pair distance distribution function

$$p(\mathbf{r}) = r^2 \int_V \Delta b(\mathbf{r}') \Delta b(\mathbf{r} - \mathbf{r}') d^3 \mathbf{r}$$



any shape – Guinier approximation

for small  $Q$  ( $Q < R_g^{-1}$ ):

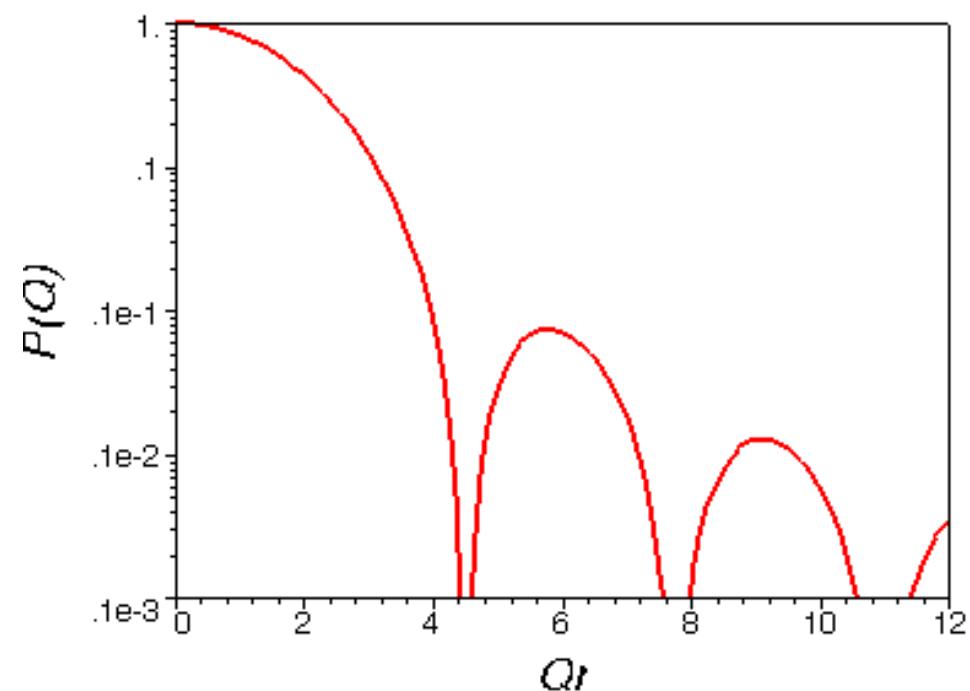
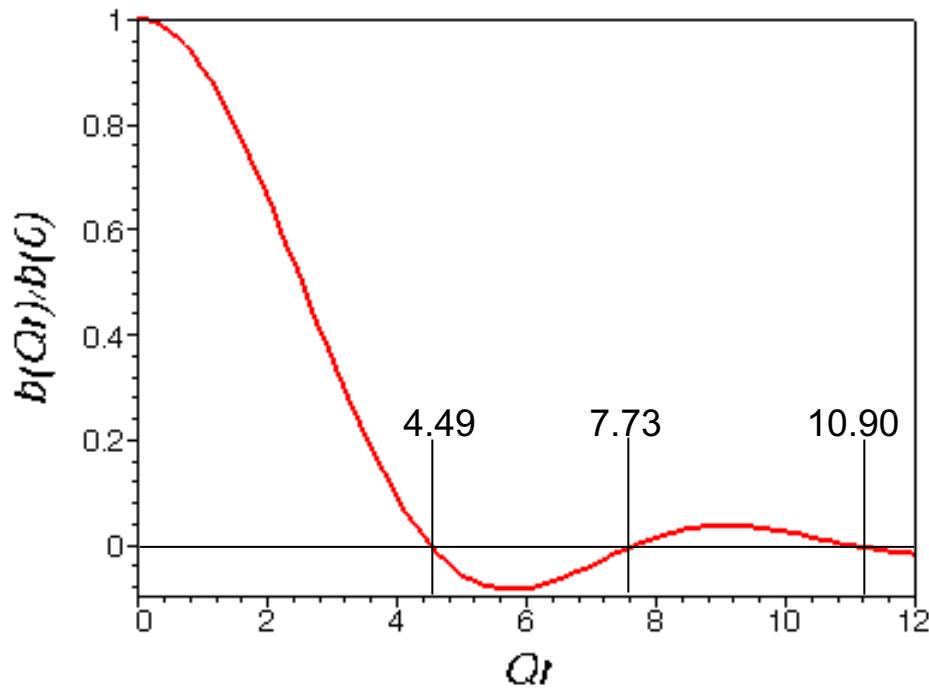
$$P(Q) = 1 - \frac{1}{3}(QR_g)^2$$

with radius of gyration  $R_g$

$$R_g^2 = \int_V \rho(\mathbf{r}) (\mathbf{r} - \mathbf{r}_{\text{cm}})^2 d\mathbf{r}$$

## homogenous sphere

$$\begin{aligned}\psi_p^{\text{sc}}(\mathbf{Q}) &= \int \Delta b(\mathbf{r}) e^{i\mathbf{Q} \cdot \mathbf{r}} dV = \Delta b \int_{V_{\text{sphere}}} e^{i\mathbf{Q} \cdot \mathbf{r}} dV = \Delta b \int d\phi \int r^2 dr \int e^{iQr \cos \theta} \sin \theta d\theta \\ &\vdots \\ &\sim \frac{3}{(QR)^3} (\sin(QR) - QR \cos(QR))\end{aligned}$$



## polydisperse homogenous spheres

radius of gyration

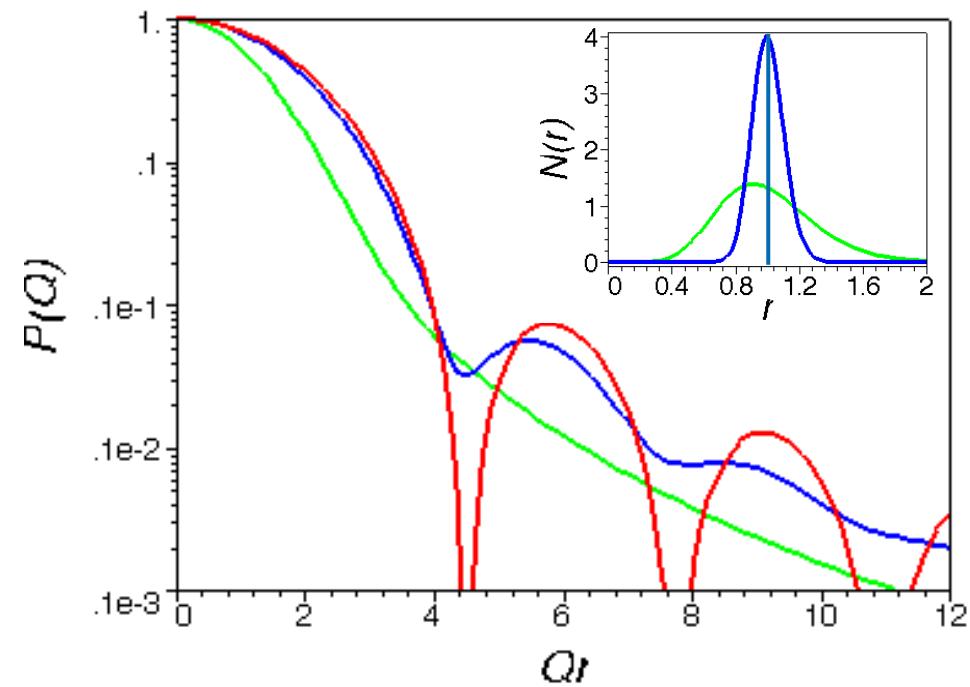
$$\langle R_g^2 \rangle = \frac{\int_0^\infty N(r) M^2(r) R_g^2 dr}{\int_0^\infty N(r) M^2(r) dr}$$

form factor

$$\langle P(\mathbf{q}) \rangle = \frac{\int_0^\infty N(r) M^2(r) P(\mathbf{q}, r) dr}{\int_0^\infty N(r) M^2(r) dr}$$

molar mass

$$\langle M \rangle = \frac{\int_0^\infty N(r) M^2(r) dr}{\int_0^\infty N(r) M(r) dr}$$



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# Structure factor $S(Q)$

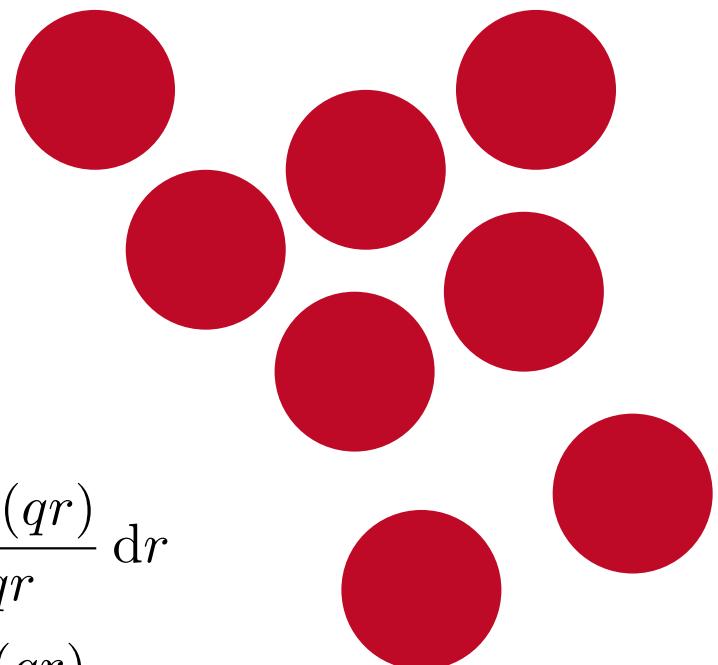
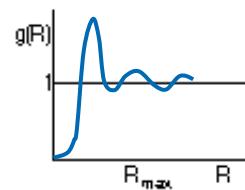
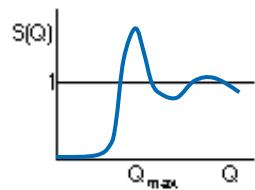
interparticle interference  
depends on particle arrangement

$$S(\mathbf{Q}) = \frac{1}{N} \sum_j \sum_k \left\langle e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \right\rangle$$

⋮

$$S(\mathbf{Q}) = 1 + 4\pi \frac{N}{V} \int (g(\mathbf{r}) - 1) r^2 \frac{\sin(qr)}{qr} dr$$

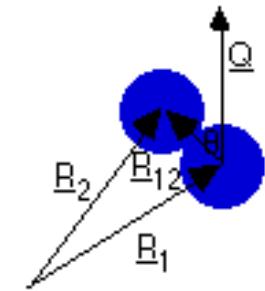
$$S(\mathbf{Q}) - 1 = 4\pi \frac{N}{V} \int (g(\mathbf{r}) - 1) r^2 \frac{\sin(qr)}{qr} dr$$



# Structure factor $S(Q)$

dumbbell

$$\begin{aligned} S(\mathbf{Q}) &= \frac{1}{N} \sum_j \sum_k \left\langle e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \right\rangle = \frac{1}{2} \sum_{j=1}^2 \sum_{k=1}^2 \left\langle e^{i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_k)} \right\rangle \\ &= \frac{1}{2} \left\langle 1 + e^{i\mathbf{Q} \cdot (\mathbf{R}_1 - \mathbf{R}_2)} + e^{i\mathbf{Q} \cdot (\mathbf{R}_2 - \mathbf{R}_1)} + 1 \right\rangle \\ &= \langle 1 + \cos(\mathbf{Q} \cdot \mathbf{R}_{12}) \rangle \end{aligned}$$



↓ spherical averaging

$$\begin{aligned} S(\mathbf{Q}) &= 1 + \frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi \cos(\mathbf{Q} \cdot \mathbf{R}_{12}) \sin \theta d\theta d\phi \\ &= 1 + \frac{\sin(2QR_{12})}{2QR_{12}} \end{aligned}$$

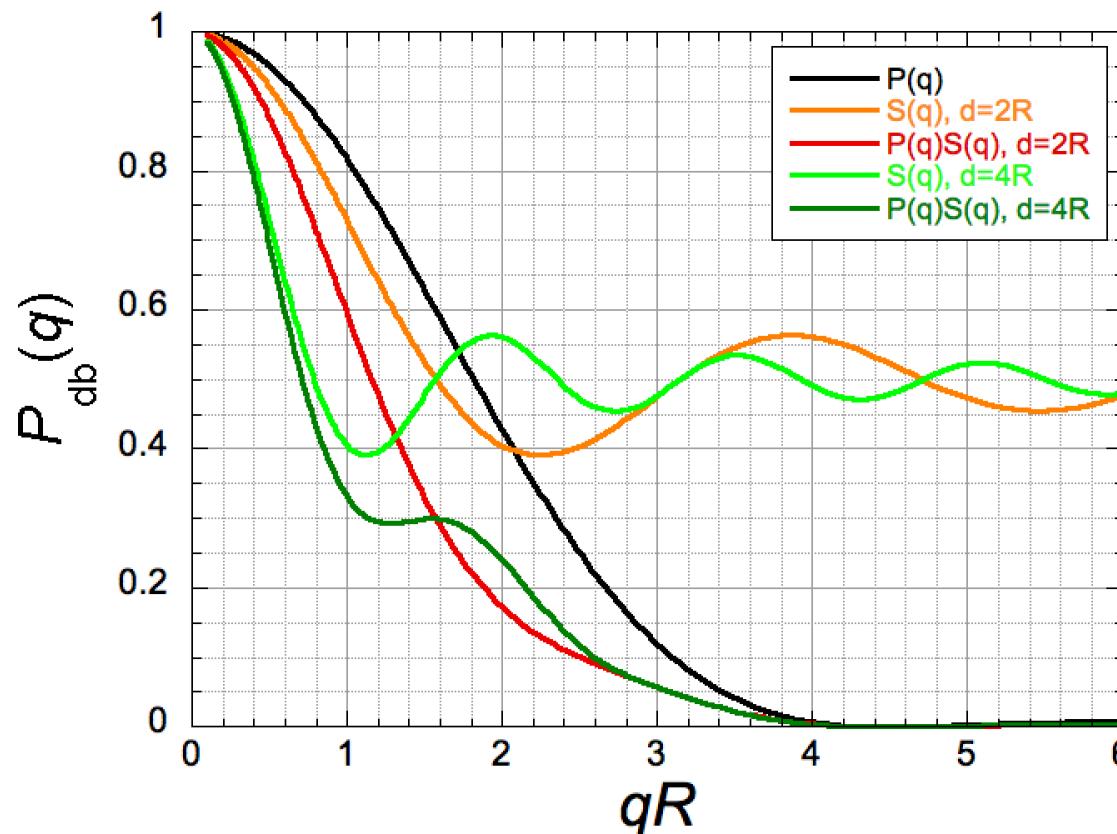
P( $q$ ) of a sphere  
dumbbell = 2 spheres

.....

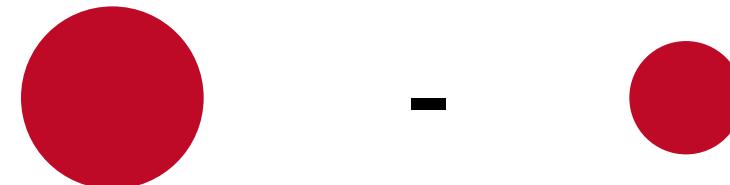
# Structure factor $S(Q)$

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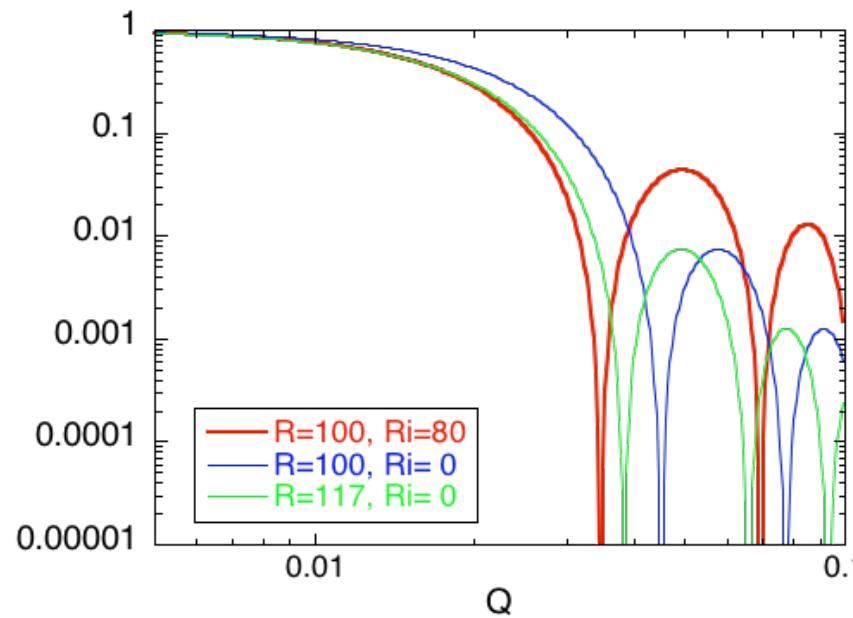
dumbbell



inhomogenous sphere (core-shell

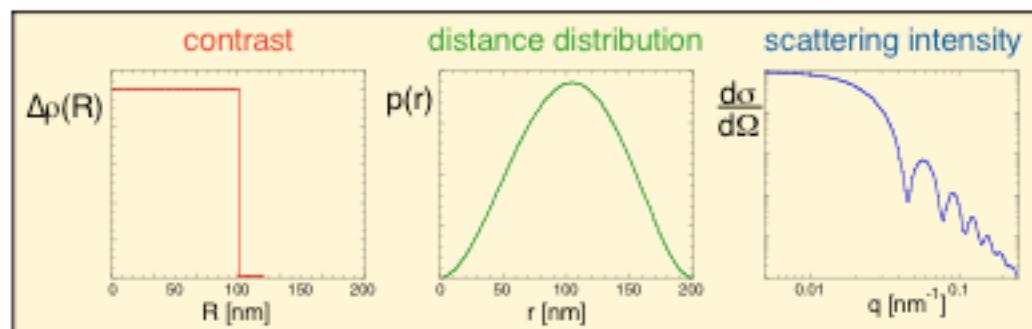


$$P(q) = \left( \frac{3}{(QR_s)^3 - (QR_c)^3} \{ (\sin(QR_s) - QR_s \cos(QR_s)) - (\sin(QR_c) - QR_c \cos(QR_c)) \} \right)^2$$



# Data analysis

indirect Fourier transformation



model fitting (analytical, numerical, simulations)

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# 'Prefactor'

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim N|\psi_p^{sc}(0)|^2$$

$$\psi_p^{sc}(\mathbf{Q}) = \int \Delta b(\mathbf{r}) e^{i\mathbf{Q}\cdot\mathbf{r}} dV \rightarrow \psi_p^{sc}(0) \sim \Delta b V \text{ with } \Delta b = \langle \Delta b(\mathbf{r}) \rangle \\ \sim (b_{\text{part}} - b_{\text{medium}}) V$$

$$\frac{d\sigma}{d\Omega}(0) \sim (b_{\text{part}} - b_{\text{medium}})^2$$

consider that solvent consists of a mixture of solvent A ( $b_A$ ) and solvent B ( $b_B$ ) with volume fraction  $\phi_B$  of solvent B:

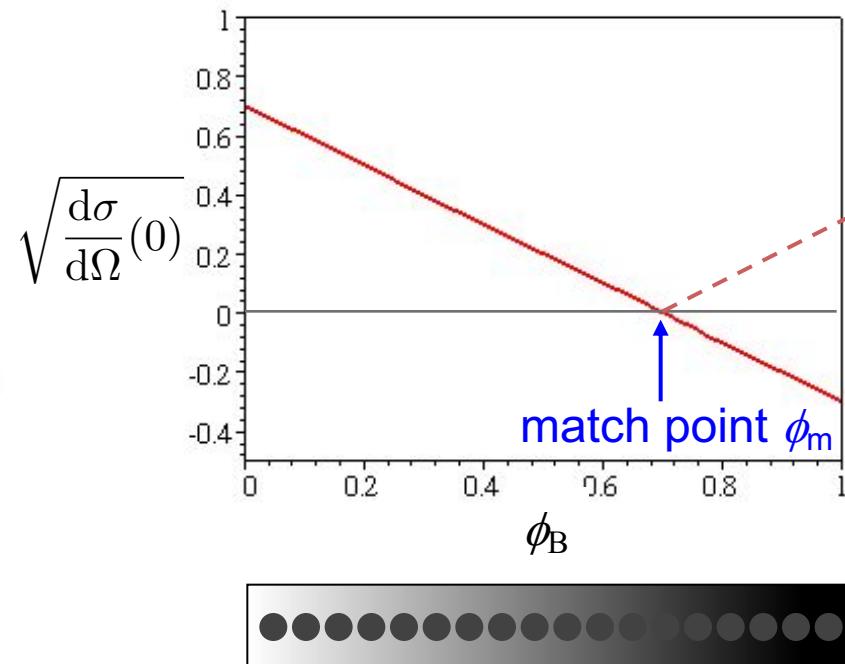
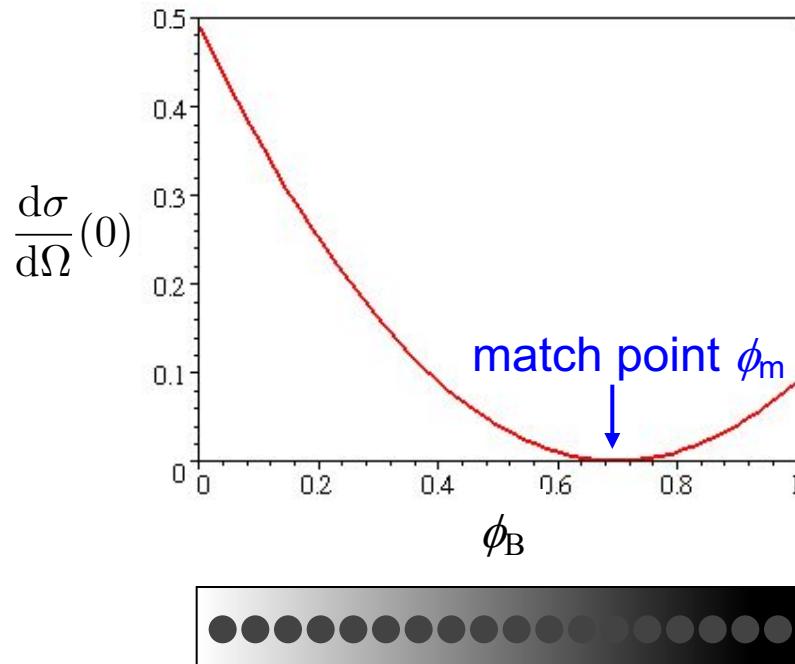
$$b_{\text{medium}} = (1 - \phi_B)b_A + \phi_B b_B = b_A + \phi_B(b_B - b_A)$$



$$\frac{d\sigma}{d\Omega} \sim (b_{\text{part}} - b_{\text{medium}})^2 = (b_{\text{part}} - \{b_A + \phi_B(b_B - b_A)\})^2$$

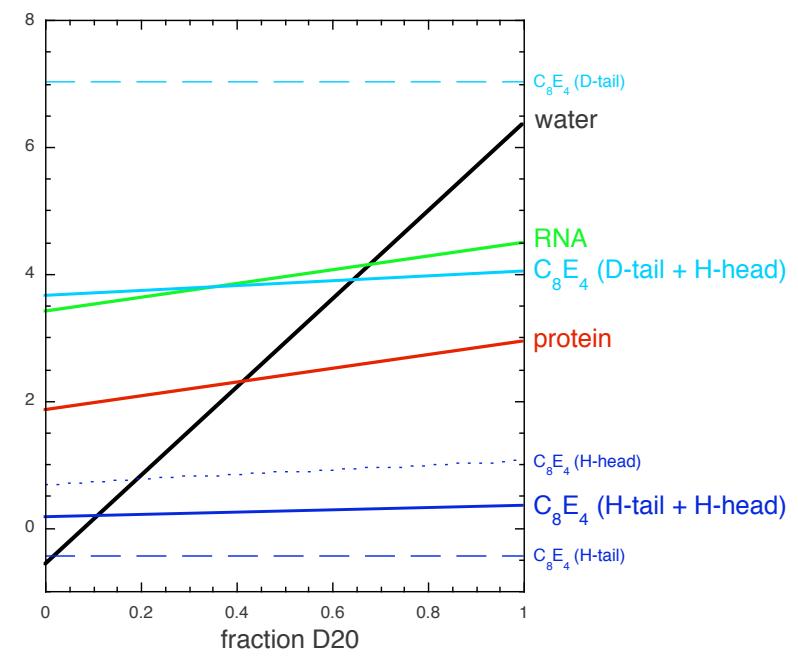
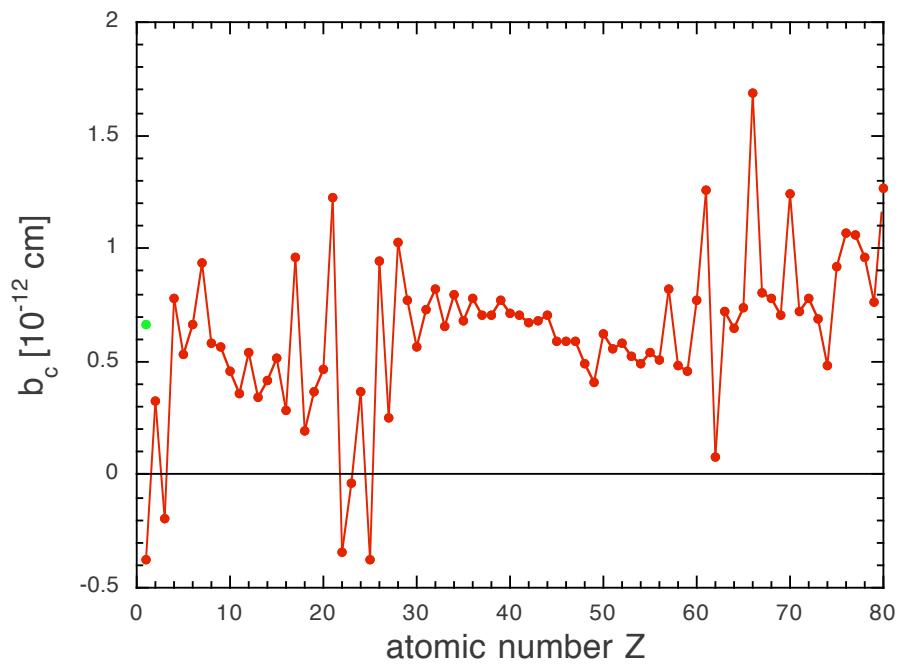
$$\sim (b_B - b_A)^2 (\phi_m - \phi_B)^2, \quad \phi_m = \frac{b_{\text{part}} - b_A}{b_B - b_A}$$

# (Solvent) contrast variation



scattering length depends on:

- nucleus, i.e. isotope (non-systematic)
- nuclear spin (magnetic scattering)





- increase signal
  - improve statistics
- reduce signal
  - reduce multiple scattering
- reduce ‘noise’
  - reduce background due to solvent

# Optimize measurement

**example:** dilute polystyrene spheres in cyclohexane  
with  $R = 10 \text{ nm}$ ,  $c = 10 \text{ mg/ml} (\approx 1\%)$

substance	formula	molar mass $M$ (g/mol)	density $\bar{\nu}^{-1}$ (g/cm <sup>3</sup> )	scatt. length dens. $b$ (10 <sup>10</sup> cm <sup>-2</sup> )
polystyrene	$(C_8H_8)_n$	$(104.2)_n$	1.04	1.44
	$(C_8D_8)_n$	$(112.2)_n$	~1.12	6.44
cyclohexane	$C_6H_{12}$	84.2	0.779	- 0.24
	$C_6D_{12}$	96.2	0.893	6.74

# Optimize measurement

**example:** dilute polystyrene spheres in cyclohexane  
with  $R = 10 \text{ nm}$ ,  $c = 10 \text{ mg/ml} (\approx 1\%)$

## excess spheres ('signal')

**H in H:**  $10.8 \text{ cm}^{-1}$

**H in D:**  $113.1 \text{ cm}^{-1}$

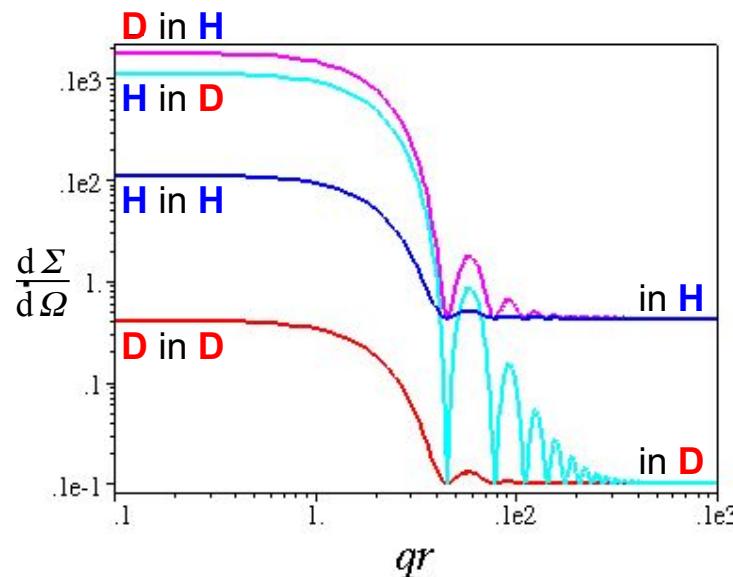
**D in H:**  $179.7 \text{ cm}^{-1}$

**D in D:**  $0.4 \text{ cm}^{-1}$

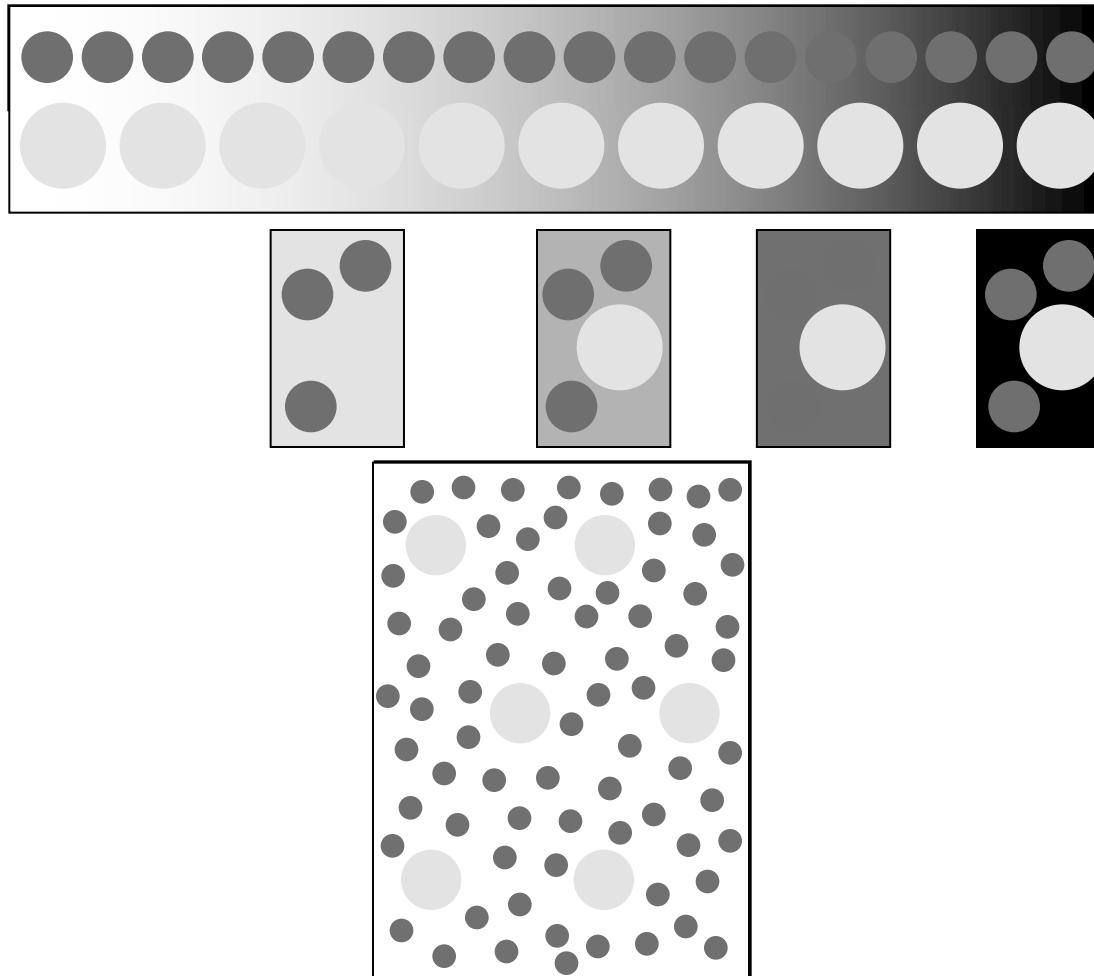
## solvent ('noise')

in **H:**  $0.43 \text{ cm}^{-1}$

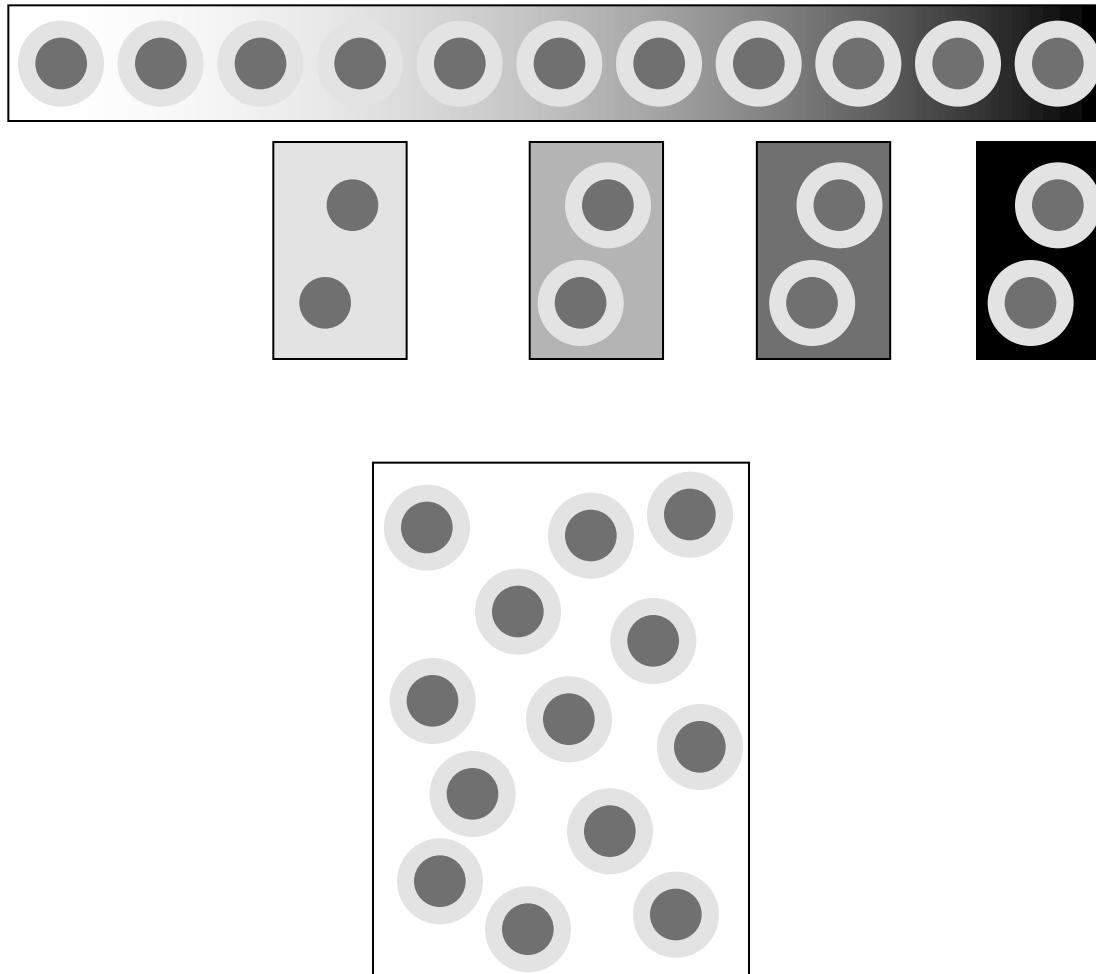
in **D:**  $0.01 \text{ cm}^{-1}$



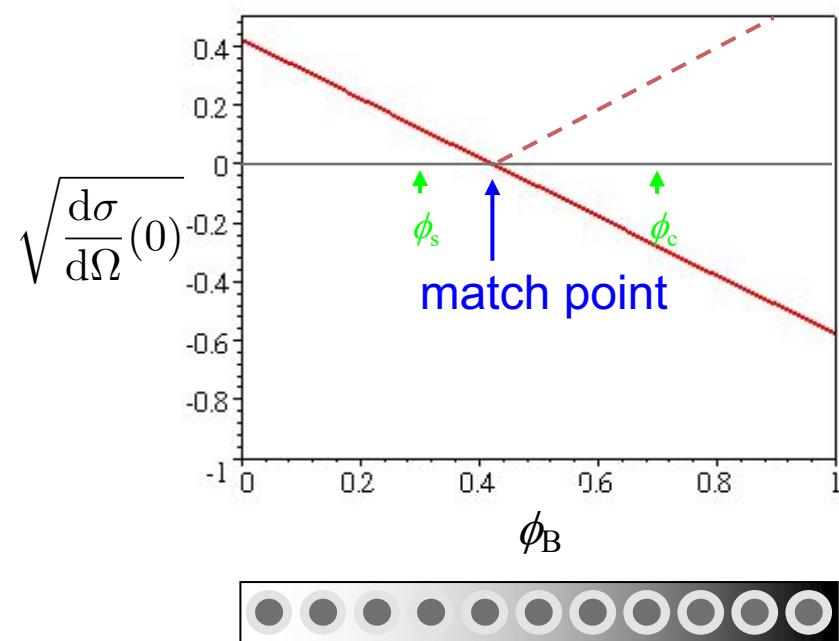
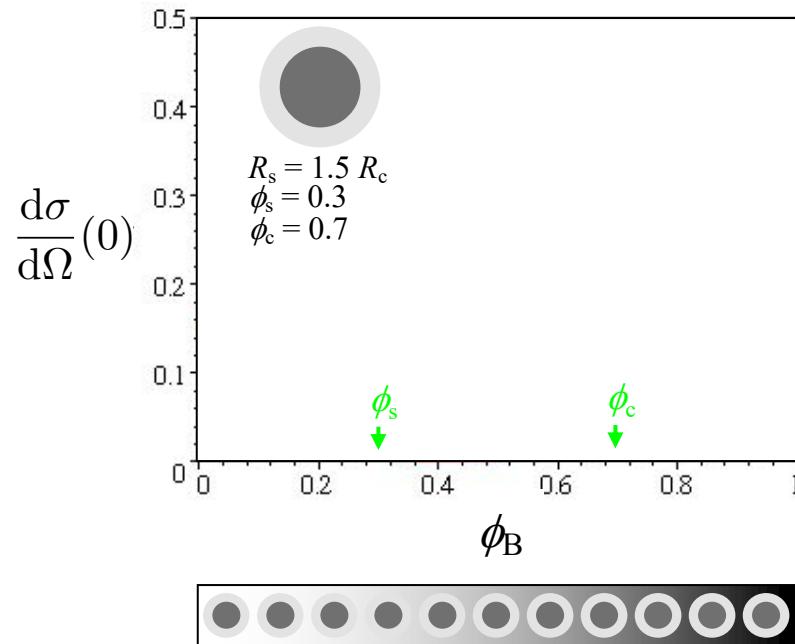
# Heterogeneous samples



# Heterogeneous particles

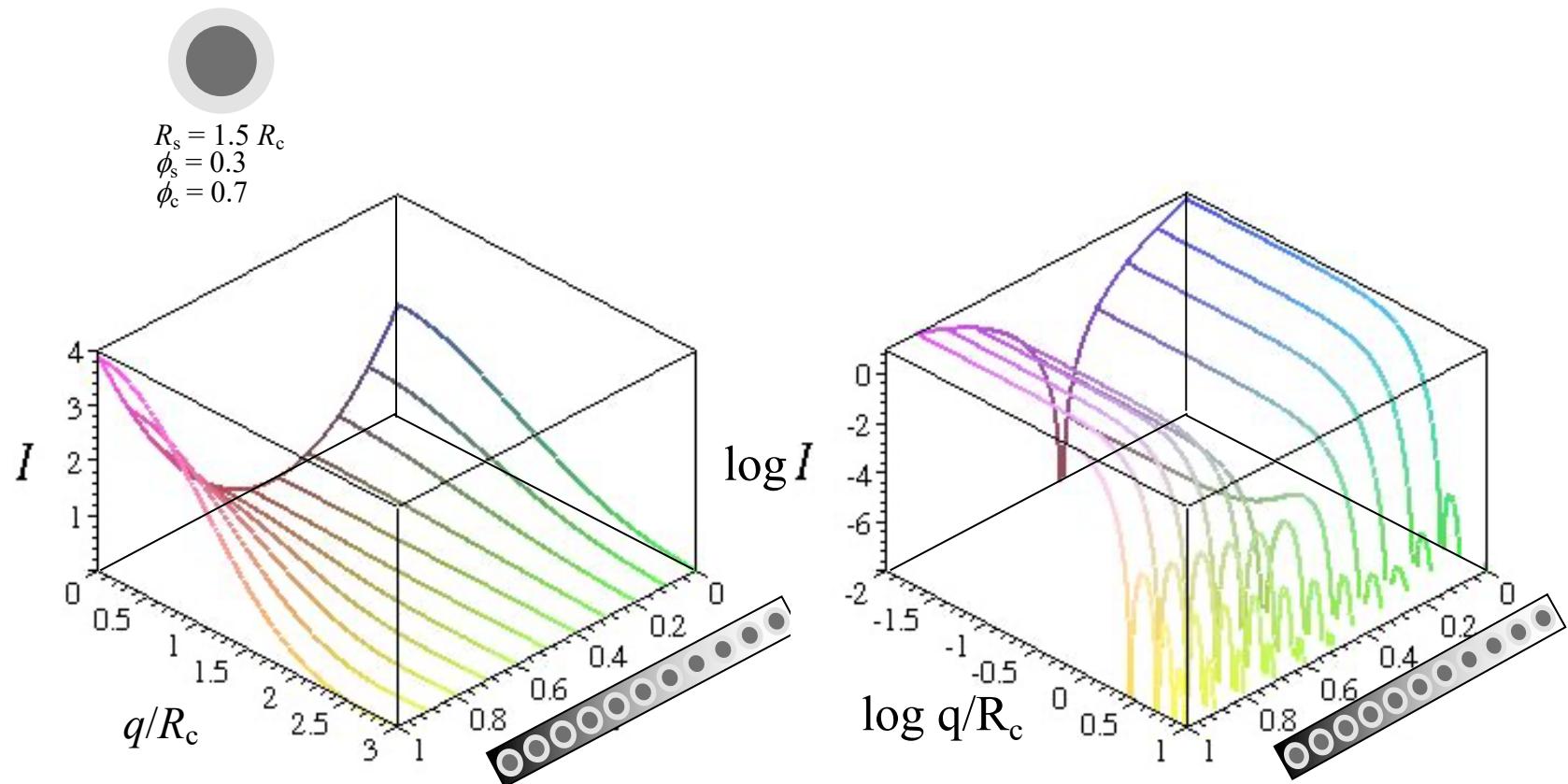


# (Solvent) contrast variation



$$\phi_{part} = \frac{\phi_s(V_s - V_c) + \phi_c V_c}{V_s} = \frac{0.3 \times (1.5^3 - 1^3) + 0.7 \times 1^3}{1.5^3} = 0.42$$

# (Solvent) contrast variation



# Differential cross section

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim |\psi^{\text{sc}}|^2 \sim \sum_k \sum_j \left\langle \psi_{p,j}^{\text{sc}}(\mathbf{Q}) \psi_{p,k}^{\text{sc}\star}(\mathbf{Q}) e^{i\mathbf{Q}\cdot(\mathbf{R}_j - \mathbf{R}_k)} \right\rangle$$

↓ identical scatterers

$$\frac{d\sigma}{d\Omega}(\mathbf{Q}) \sim N |\psi_p^{\text{sc}}(0)|^2 \frac{|\psi_p^{\text{sc}}(\mathbf{Q})|^2}{|\psi_p^{\text{sc}}(0)|^2} \frac{1}{N} \sum_k \sum_j \left\langle e^{i\mathbf{Q}(\mathbf{R}_j(t) - \mathbf{R}_k(t))} \right\rangle$$

but  $b_j, b_k$  depend on nucleus (isotope) and spin

thus for same isotope

$$j \neq k : \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \langle b_j \rangle \langle b_k \rangle \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \langle b \rangle^2 \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle$$

$$j = k : \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \langle b_j^2 \rangle = \langle b^2 \rangle$$

↓

$$\frac{d\sigma}{d\Omega} \sim \sum_j \sum_k \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle = \sum_j \sum_{k \neq j} \langle b_j b_k e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle + \sum_j \langle b_j b_j e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle$$

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_{k \neq j} \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle + N \langle b^2 \rangle$$

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_k \langle e^{i\mathbf{Q}\cdot\mathbf{r}_{jk}} \rangle + N \{ \langle b^2 \rangle - \langle b \rangle^2 \}$$

## coherent cross section

$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

- contains structural information on particles & arrangement ( $e^{iQr}$  term)
- collective properties of particles ( $S(Q)$ ,  $D_c$ )

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_k \langle e^{i\mathbf{Q} \cdot \mathbf{r}_{jk}} \rangle$$

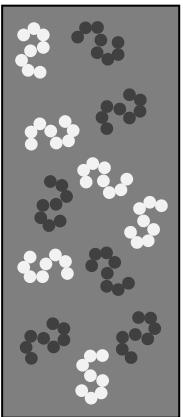
## incoherent cross section

$$\sigma_{incoh} = 4\pi (\langle b^2 \rangle - \langle b \rangle^2)$$

- contains no structural information on arrangement (no interference term, missing phase relation)
- properties of individual particles ( $D_s$ )

$$+ N \{ \langle b^2 \rangle - \langle b \rangle^2 \}$$

# Zero average contrast



identical particles, but hydrogenated ( $b_H$ ) and deuterated ( $b_D$ ) species in H/D-solvent such that  $\Delta b_H = -\Delta b_D$

$$\Rightarrow \langle b \rangle = \frac{1}{2}(\Delta b_H + \Delta b_D) = 0$$

$\Rightarrow$  coherent contribution = 0  
(i.e. structural contribution)

$$\langle b^2 \rangle = \frac{1}{2}(\Delta b_H^2 + \Delta b_D^2) = \Delta b_H^2 = \Delta b_D^2$$

incoherent contribution  $\neq 0$   
(i.e. individual particle contribution)

## coherent cross section

$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

- contains structural information on particles & arrangement ( $e^{iQr}$  term)
- collective properties of particles ( $S(Q)$ ,  $D_c$ )

$$\frac{d\sigma}{d\Omega} = \langle b \rangle^2 \sum_j \sum_k \langle e^{i\mathbf{Q} \cdot \mathbf{r}_{jk}} \rangle$$

## incoherent cross section

$$\sigma_{incoh} = 4\pi (\langle b^2 \rangle - \langle b \rangle^2)$$

- contains no structural information on arrangement (no interference term, missing phase relation)
- properties of individual particles ( $D_s$ )

$$+ N \{ \langle b^2 \rangle - \langle b \rangle^2 \}$$

## solvent isotope substitution can cause

- shift in cmc (generally higher in H<sub>2</sub>O)
- shift in Θ temperature ( $\pm 4^\circ\text{C}$  for PS in cyclohexane) or critical point
- shift in melting point (6°C for polyethylene)
- selective adsorption
- exchange of particle H/D with solvent (e.g. OH, NH, COOH groups, scattering length density becomes a function of the solvent H/D ratio)
- ...

## ‘particle’ isotope substitution can cause

- phase separation (e.g. H- and D-polymer)
- ...

# Conclusions

non-crystalline materials are wonderful

K. Paul, V. Thomas, *Oxford University Press* (1987)

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