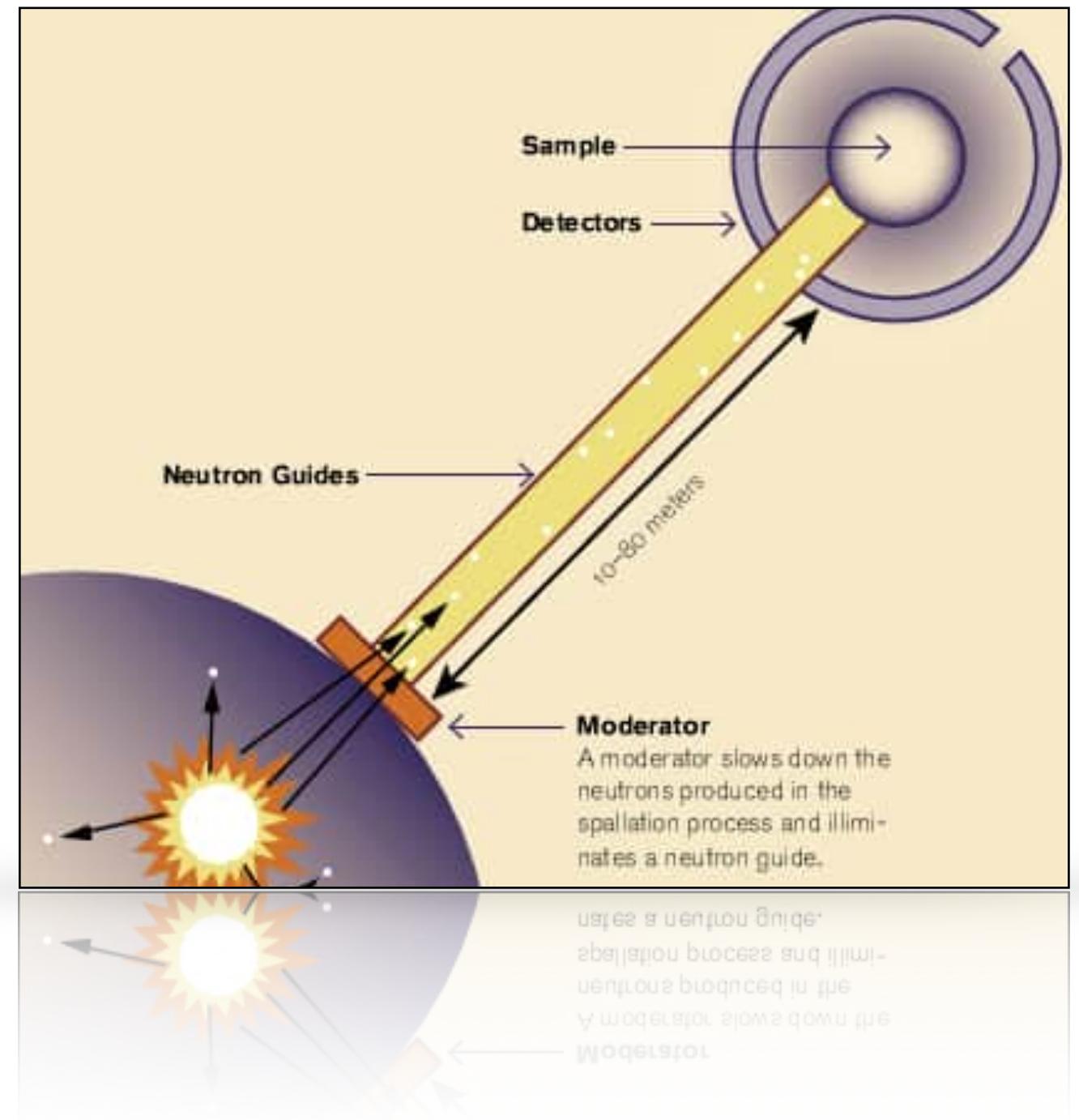


ISIS Neutron and Muon Source

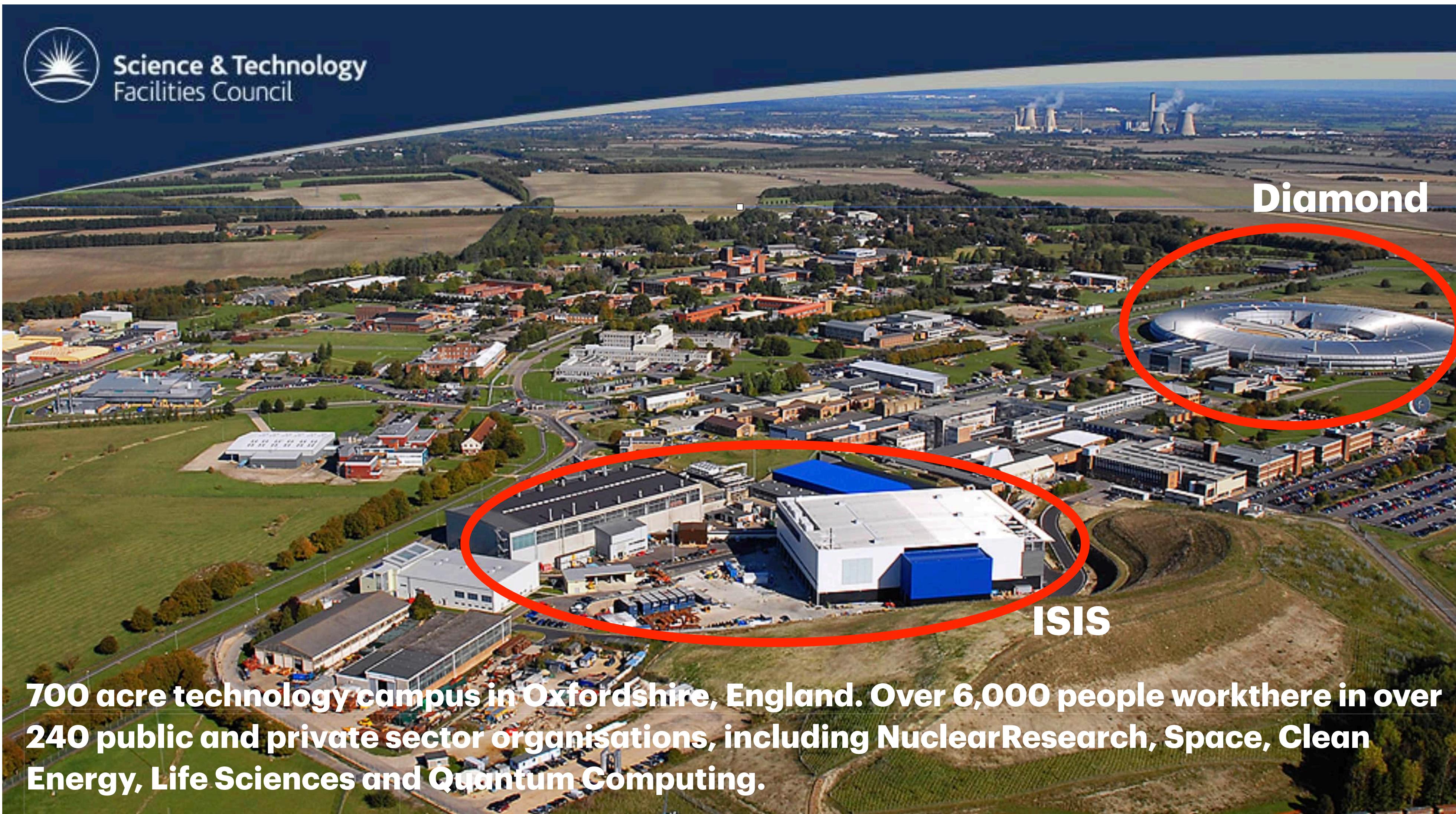


# Neutron Sources

With Focus on Neutron Scattering Facilities

Dr Lina Quintieri

# RAL at Harwell Campus- Oxford



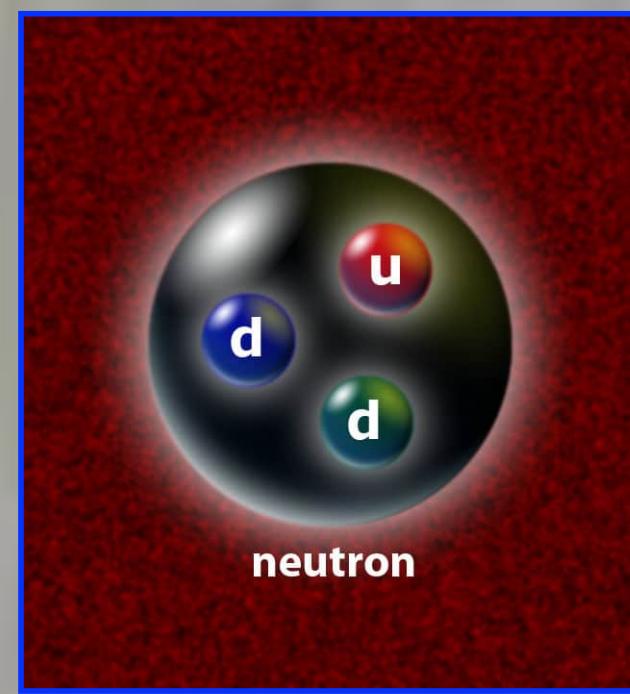
- ISIS
- DIAMOND
- CLFC
- QUANTUM Comp. Center
- Historical remark of Harwell campus: the construction of the first nuclear reactors in Western Europe (PLUTO and DIDO)

# Overview

- Brief introduction to the neutron and its importance in applied and fundamental research
- Neutron production: overview of main mechanisms
- Spallation vs Fission process → Accelerator driven sources vs research nuclear reactors
- Large Neutron Facilities for neutron scattering around the world: a few examples
- The spallation “core” : Target/Moderator/Reflector assembly
- The research reactor core vs NPP
- Moderators in pulsed neutron sources and in research reactors
- Neutron Spectrometry: a few elements of time of flight technique → how reactors compare and contrast to Spallation sources
- short description of main instrument components of beam line of a Large scale scattering facility

# ID of a Neutron

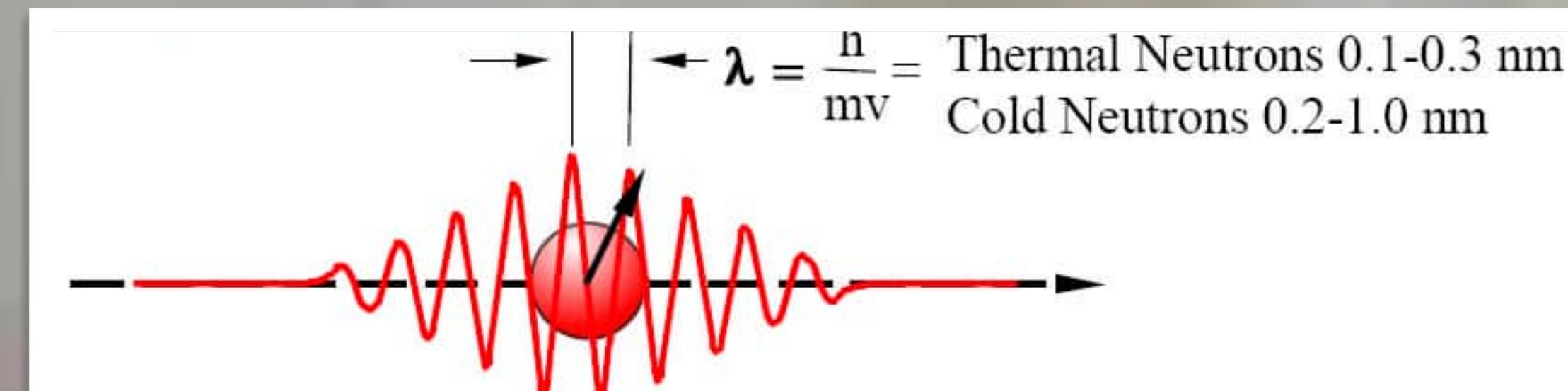
**Neutron mass:**  $1.674927471 \cdot 10^{-27}$  kg  
**Proton mass:**  $1.672621898 \cdot 10^{-27}$  kg  
**Electron mass:**  $9.10938356 \cdot 10^{-31}$  kg



This means that it is energetically possible for the neutron to decay via:

$$n \rightarrow p + e + \bar{\nu}_e$$

**Long life times (almost 15 min)**  $\Rightarrow$  before decaying possibility to interact  $\Rightarrow$  n physics ...



Although neutral, neutron is made of quarks that are electrically charged particles.

The neutron is composed of three quarks, and the magnetic moments of these elementary particles combine to give the neutron its magnetic moment.

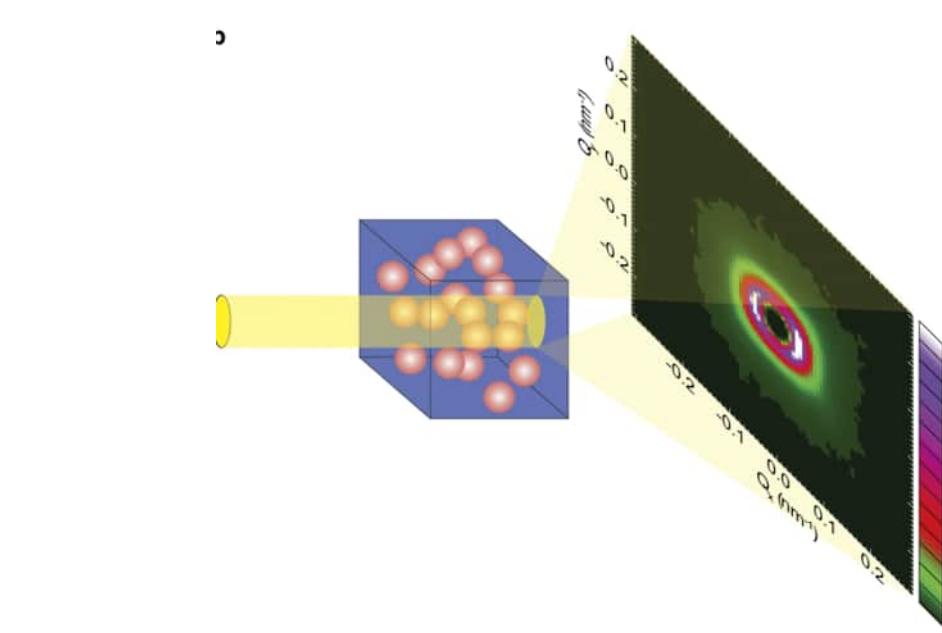
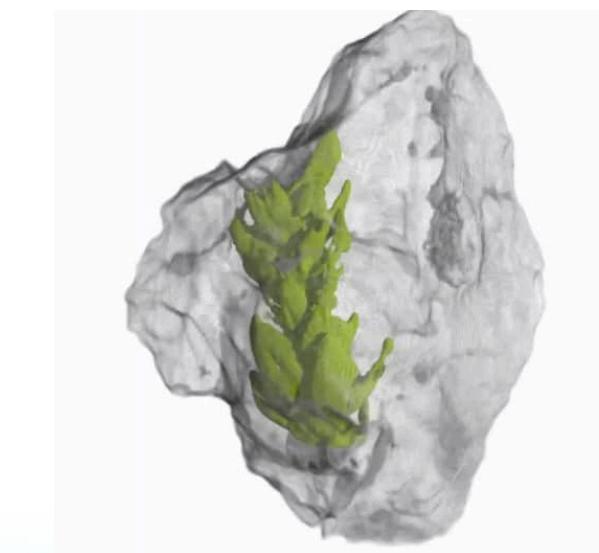
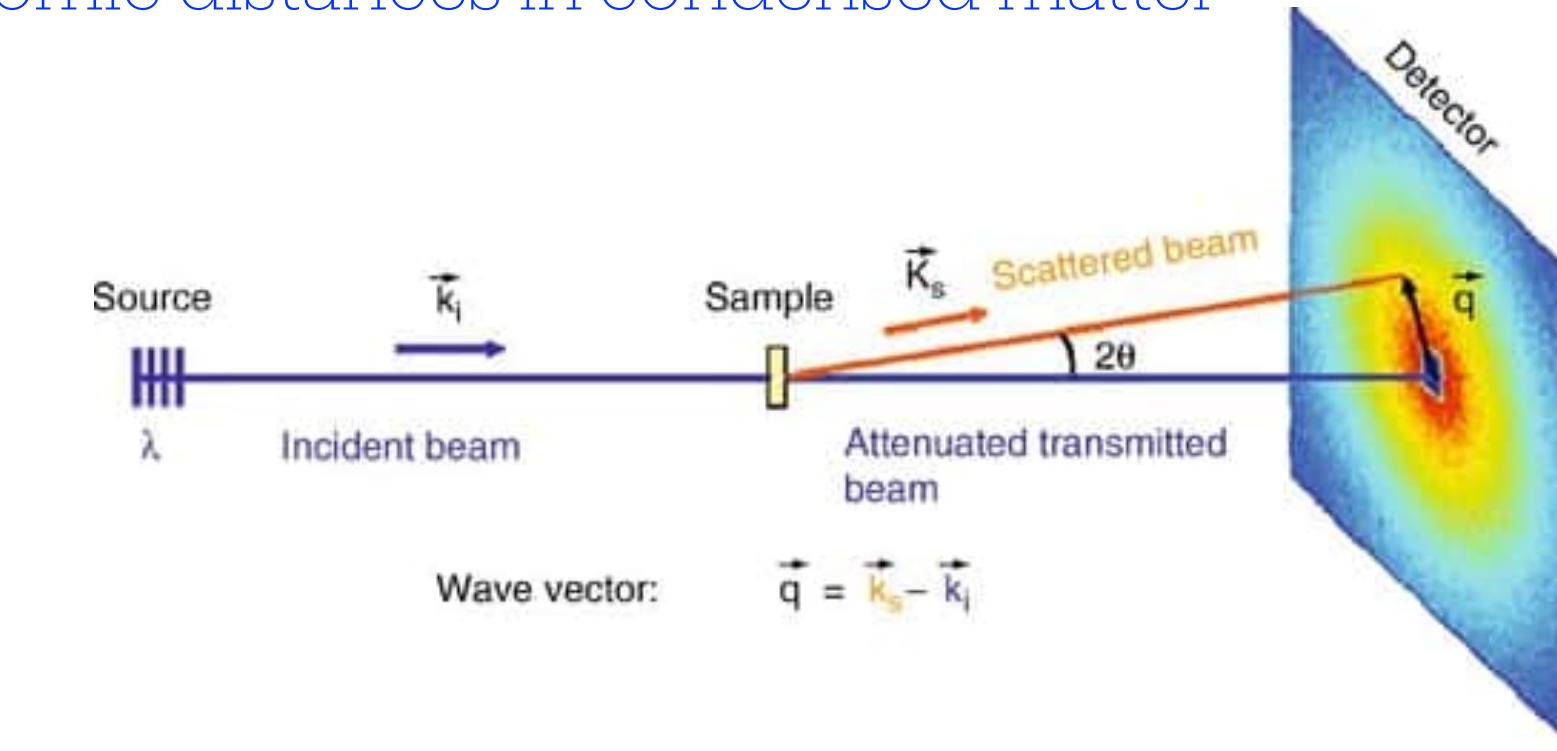


The neutron interacts with matter primarily through the nuclear force and through its magnetic moment.

# Why Neutrons Are So Attractive?

The importance of neutrons is rooted in their unique properties; penetrating, uncharged (weakly interacting with matter), they have a magnetic moment and wavelengths that can be similar to inter-atomic distances in condensed matter

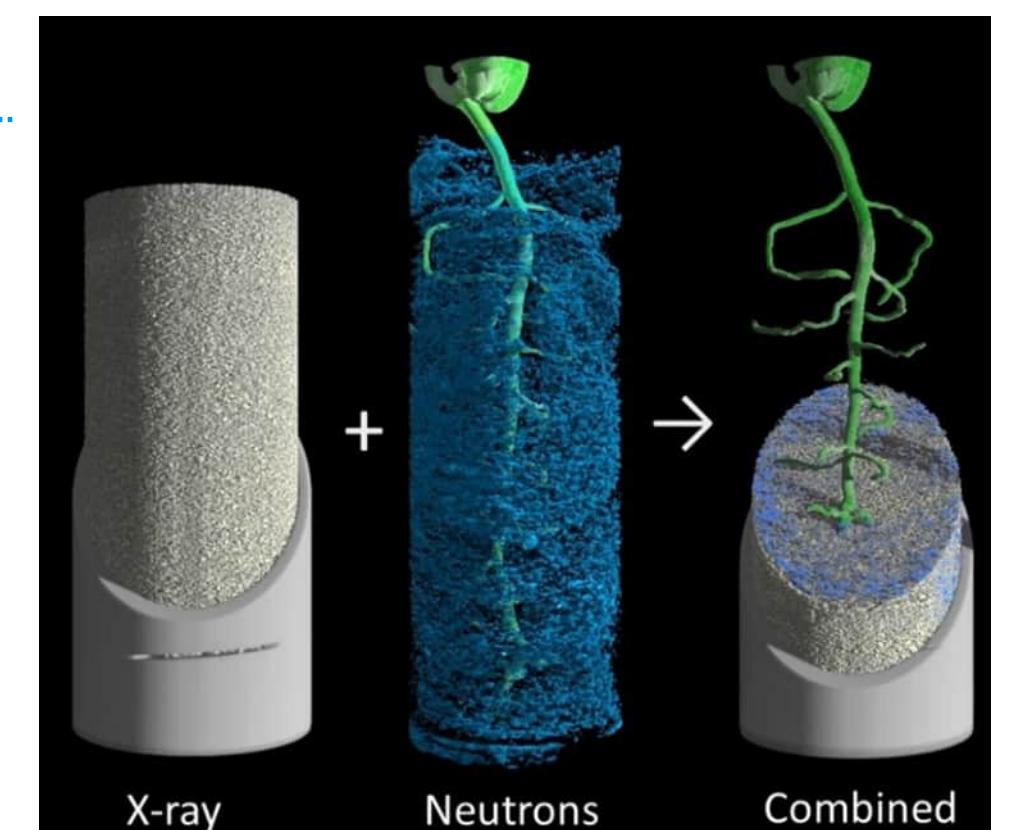
- See the nuclei.
- Interact weakly / penetrate deep into matter
- Have isotopic sensitivity (especially H and D differentiation)
- Thermal neutron have wavelengths similar to inter-atomic distances and energies comparable to lattice vibrations
- Interact via a simple point-like potential amplitudes are straightforward to interpret
- See a completely different contrast to x-rays
- See elementary magnets.



Neutron is also sensitive to the magnetic properties of the sample, and can therefore resolve complex magnetic structures.

By observing the patterns in space of the scattered neutrons, it is possible to establish where the atoms are positioned with exceptional precision. They can provide structural information from Angstrom to microns

Neutrons are highly penetrating and can be used as non-destructive probes, also to study samples in extreme environments (ex. high pressure/ high magnetic field).

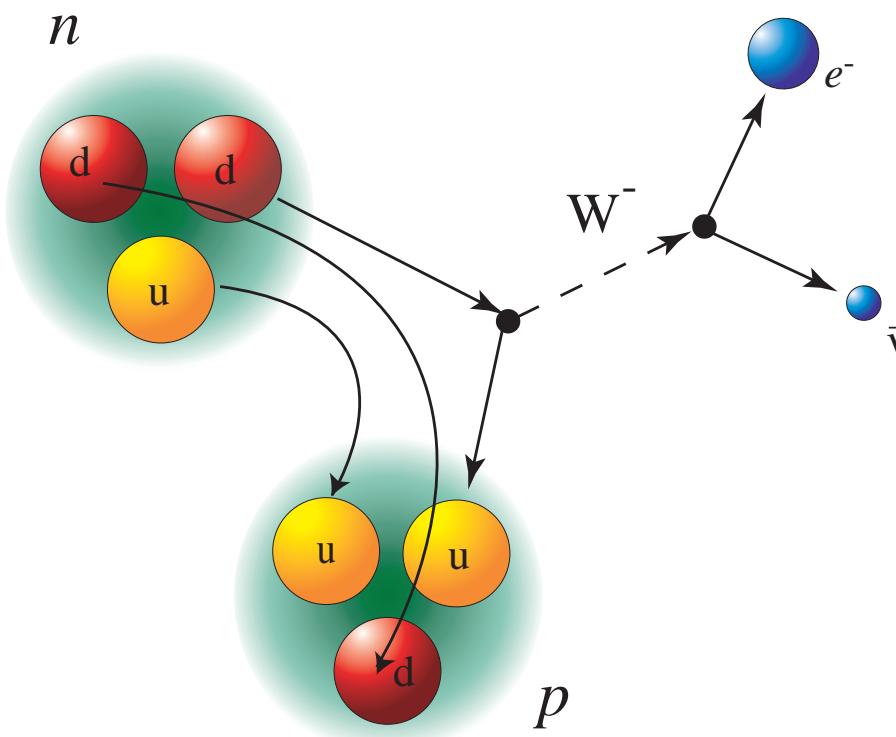


Neutrons can see light nuclei next to heavy ones

# Neutrons and Fundamental Physics Research

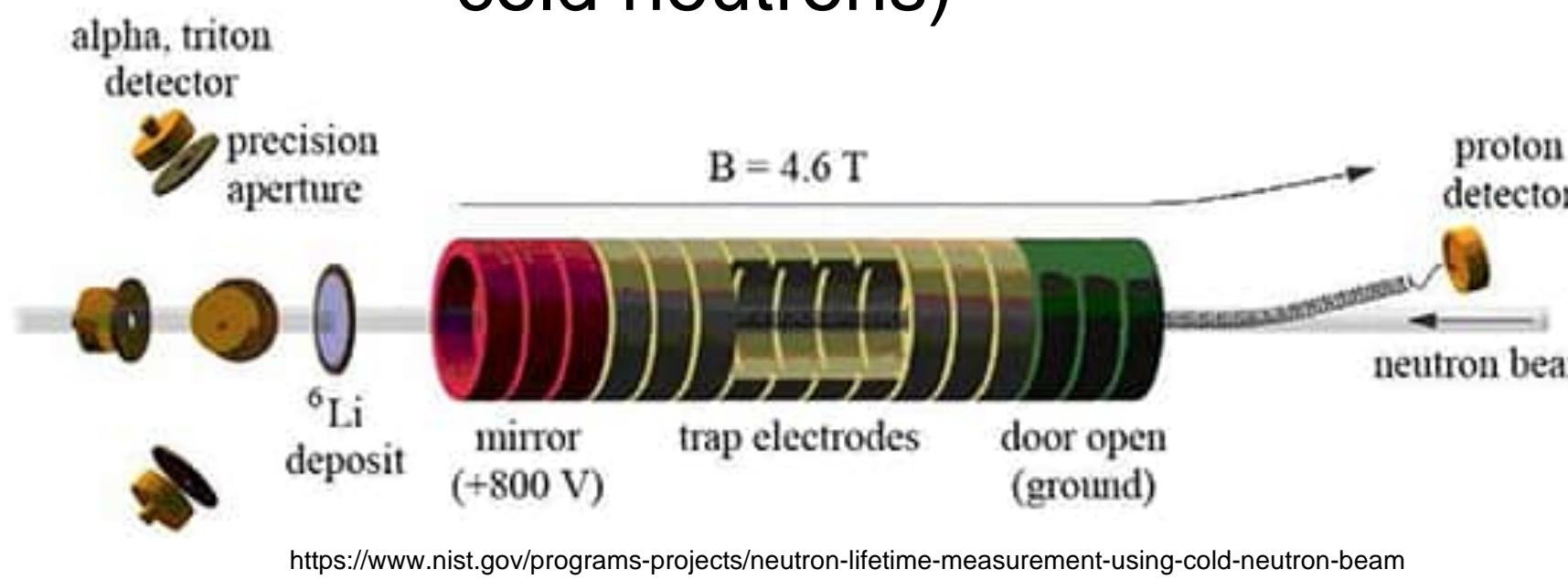
Cold (meV) and ultra cold neutrons (neV) in nuclear experiments: the “enigma” of neutron half-life

## Neutron Decay



$$n \rightarrow p^+ + e^- + \bar{\nu}_e + 782 \text{ keV}$$

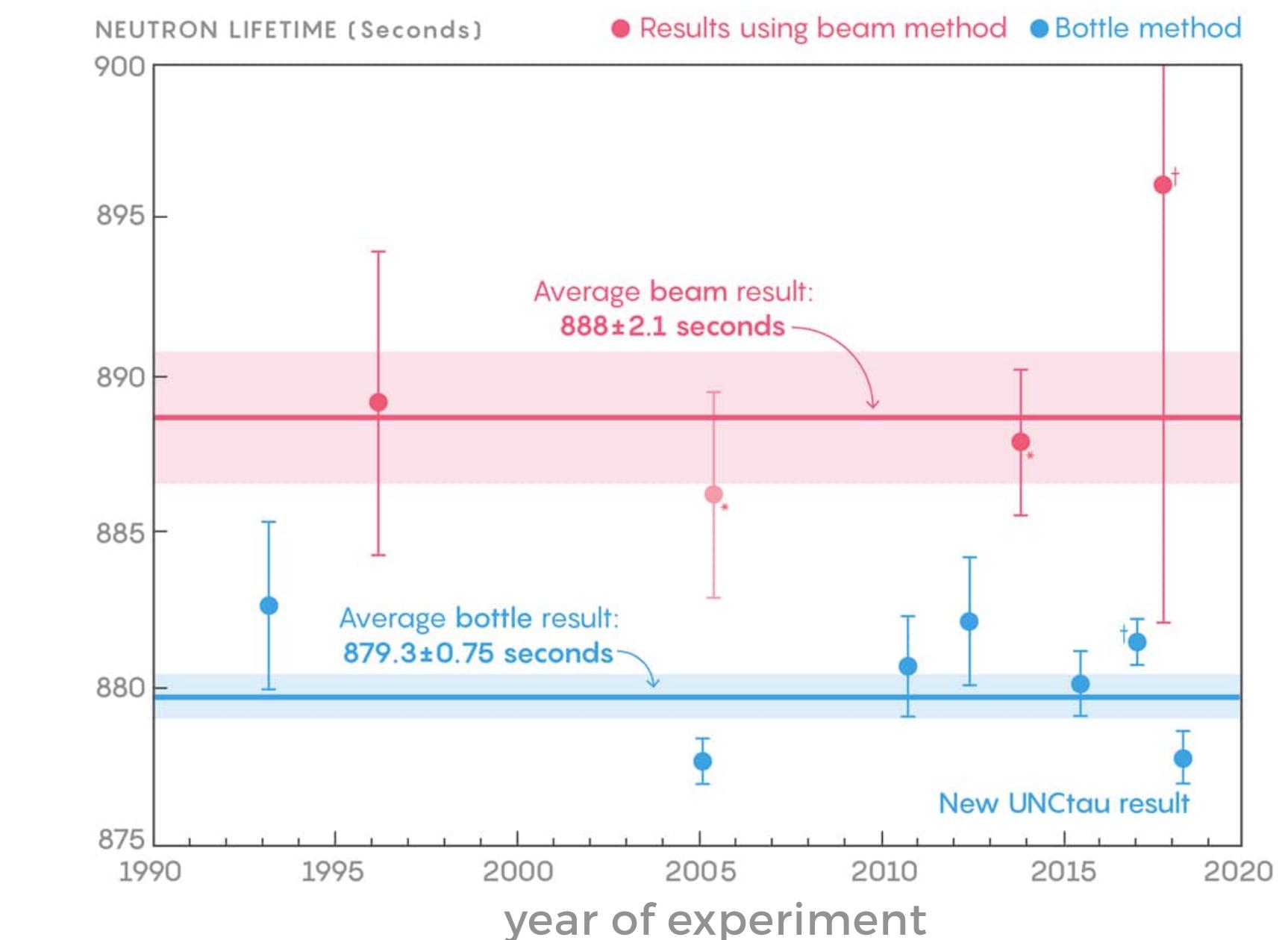
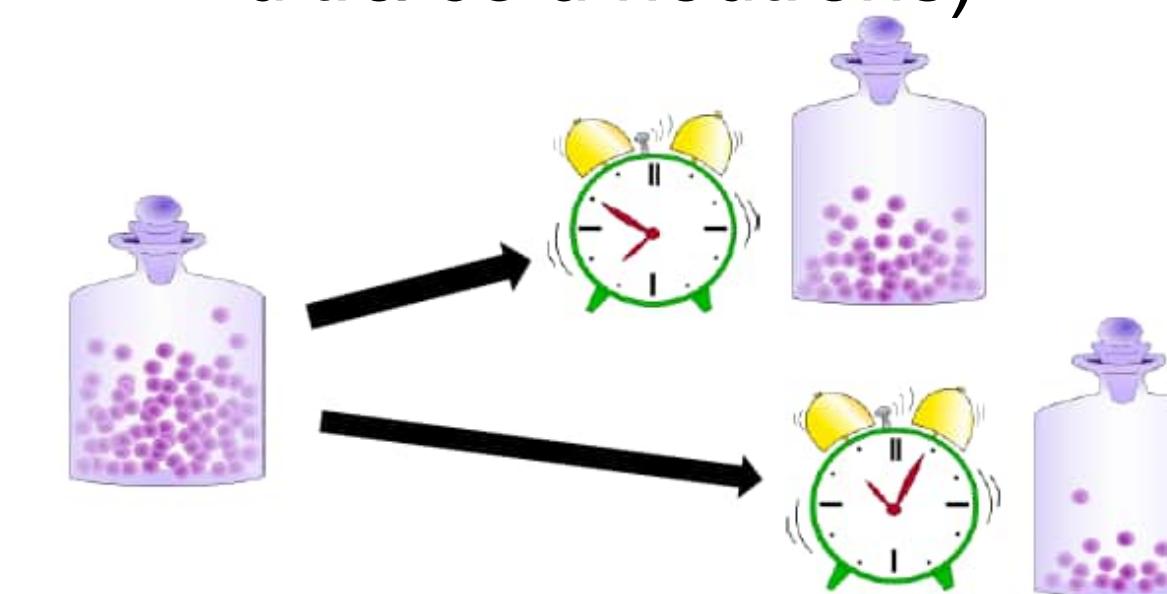
**Beam** experiment  
(2 absolute measurement  
cold neutrons)



The importance of the neutron lifetime knowledge: The accurate assessment of the lifetime of neutrons is the key to understanding the formation of elements after the Big Bang 13.8 billion years ago.

It could influence the standard model of physics that governs our understanding of the formation of the universe. Neutron lifetime is the easiest and most direct way of measuring the weak force, one of four fundamental forces in nature.

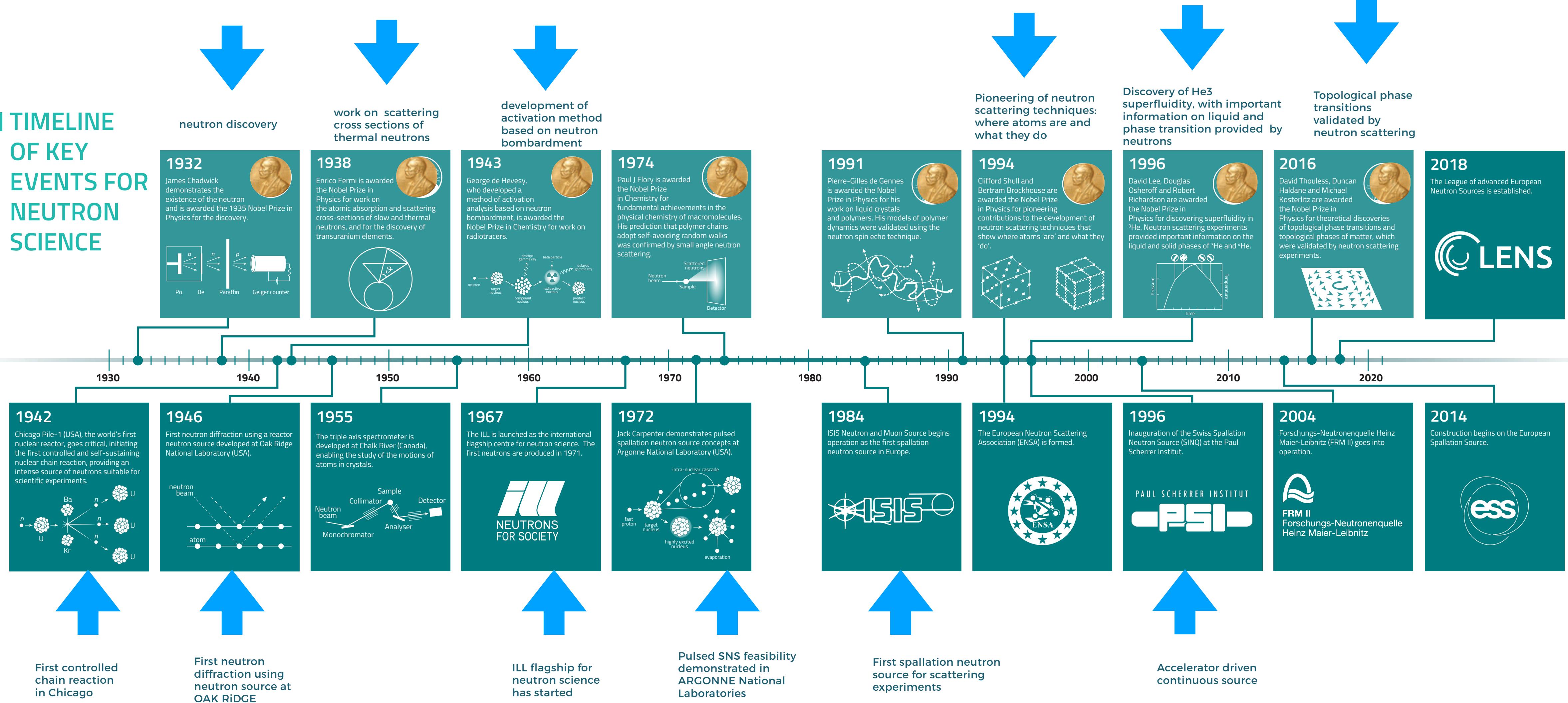
**« Bottle »** experiment  
(2 relative measurement  
ultra cold neutrons)



# Timeline of 92 Years of Neutrons

The discovery of neutron has marked the beginning of modern nuclear physics

## TIMELINE OF KEY EVENTS FOR NEUTRON SCIENCE



# Important Milestones: a Little Bit of History

- 1932: Chadwick discovers the neutron
- 1936: Mitchell and Powers see neutron diffraction (demonstrations of coherent neutron diffraction (Bragg scattering by crystal lattice planes)
- 1947: Zinn measures the first neutron Bragg peak rocking curve



James Chadwick

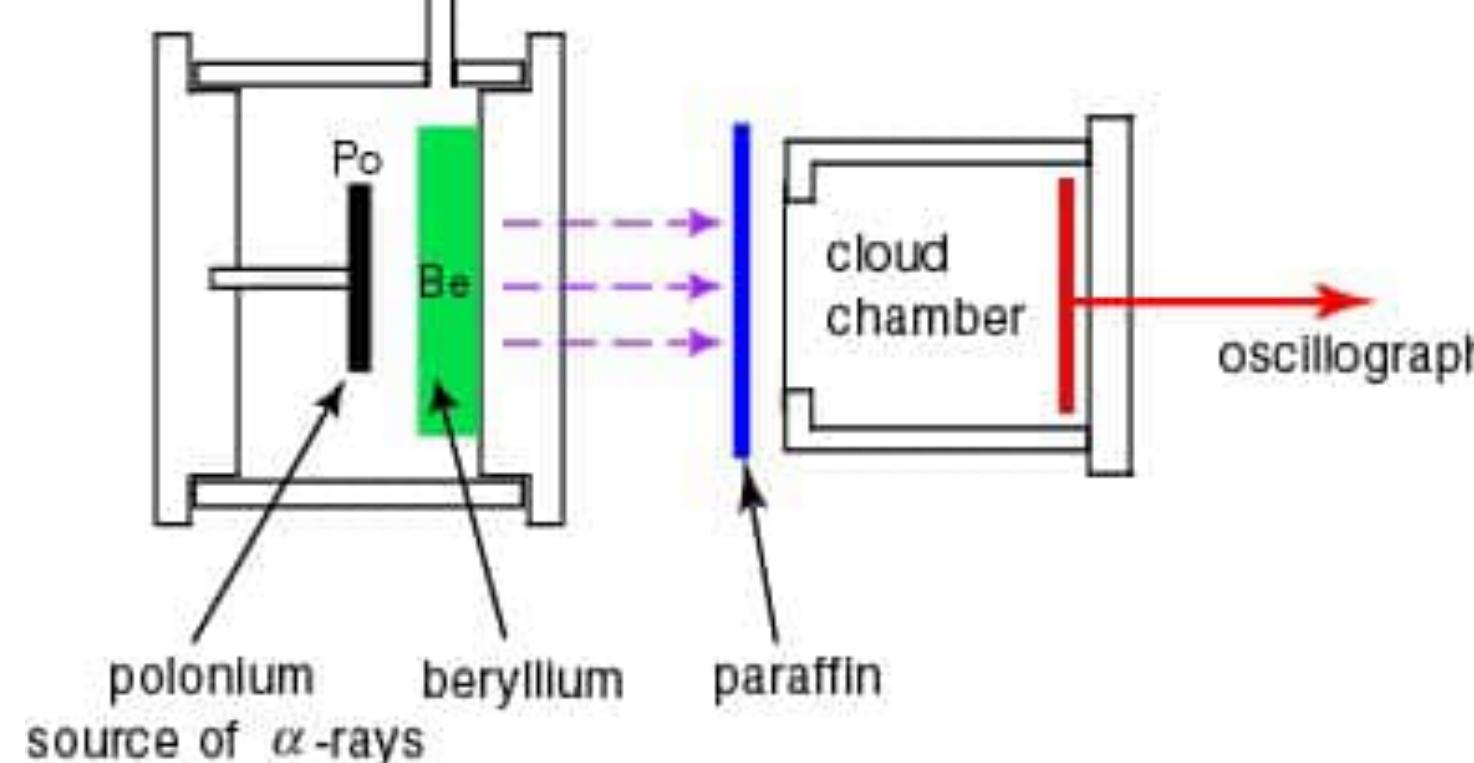
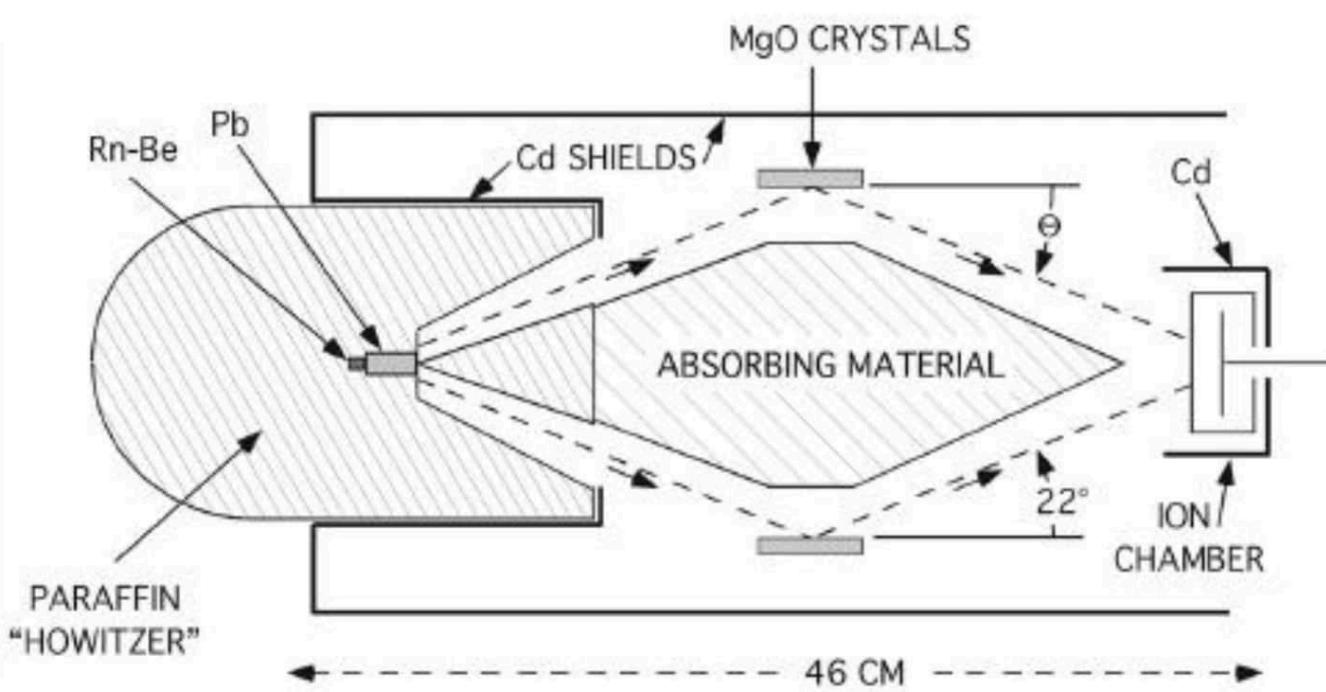


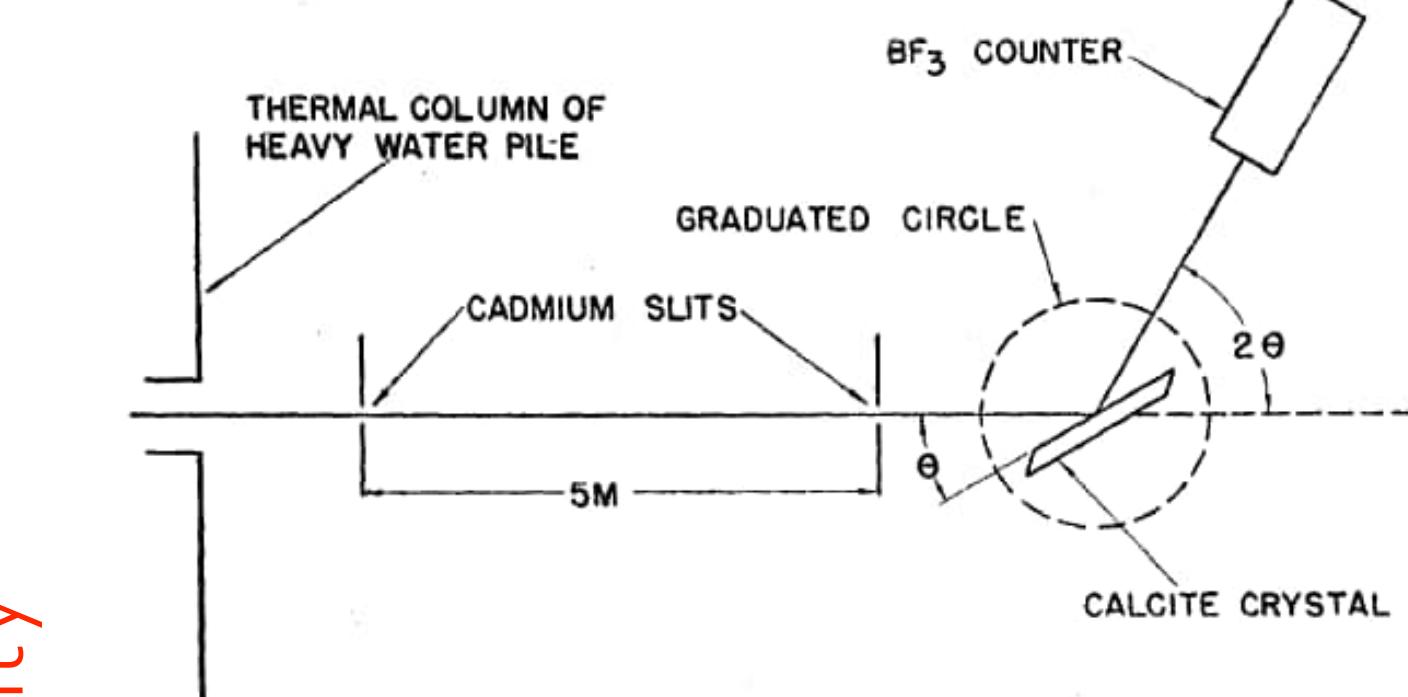
TABLE I. OBSERVED NUMBERS						
Run	Bragg		Crossed		Background	
	Counts $\times 10^{-3}$	Rate $N_B/\text{min.}$	Counts $\times 10^{-3}$	Rate $N_X/\text{min.}$	Counts $\times 10^{-3}$	Rate $N_{Cd}/\text{min.}$
1st	23	$60.5 \pm 4$ $N_B - N_X = 4.9 \pm .6$	21	$55.6 \pm .4$	3.8	$43.3 \pm .7$
2nd	11	$28.8 \pm .3$ $N_B - N_X = 2.3 \pm .4$	3.0	$2.0 \pm .3$	3	$20.9 \pm .4$
3rd	12	$37.6 \pm .3$ $N_B - N_X = .1 \pm .4$	12	$37.7 \pm .3$	6	$28.0 \pm .4$

Laboratory source PoBe, RaBe

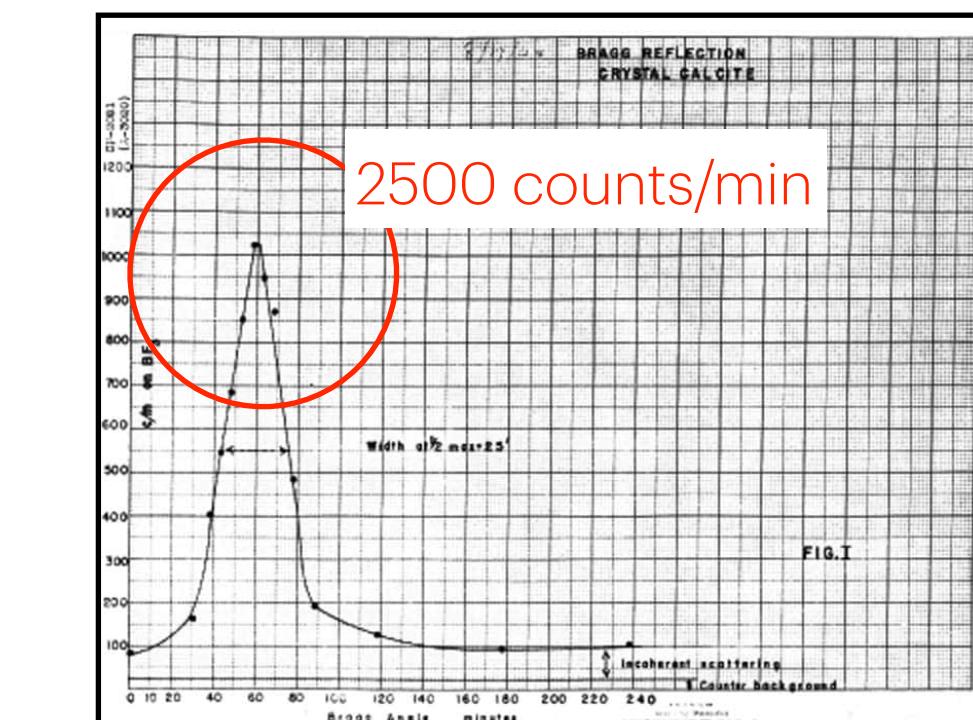


Apparatus of Mitchell and Powers (1936)

Low Intensity      High Intensity



First experimental setup at CP-3  
[W. H. Zinn, Phys. Rev. 71, 752–757 (1946)]

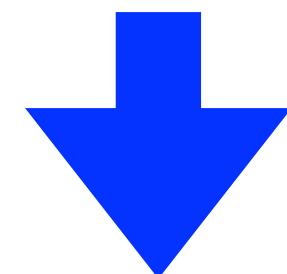


CP-3 reactor

# Intensity Matters

Scattering experiments need of intense neutron beam

- Neutron scattering measures a neutron count rate
- The count rate is the number of counts, N, divided by the measurement time, t.
- The statistical uncertainty is given by the square root of the counts,  $\sqrt{N}$ , i.e. 10 000 counts have an uncertainty of 100, or 1%
- Neutron scattering is generally weak



Scattering experiments need good statistics for significant results

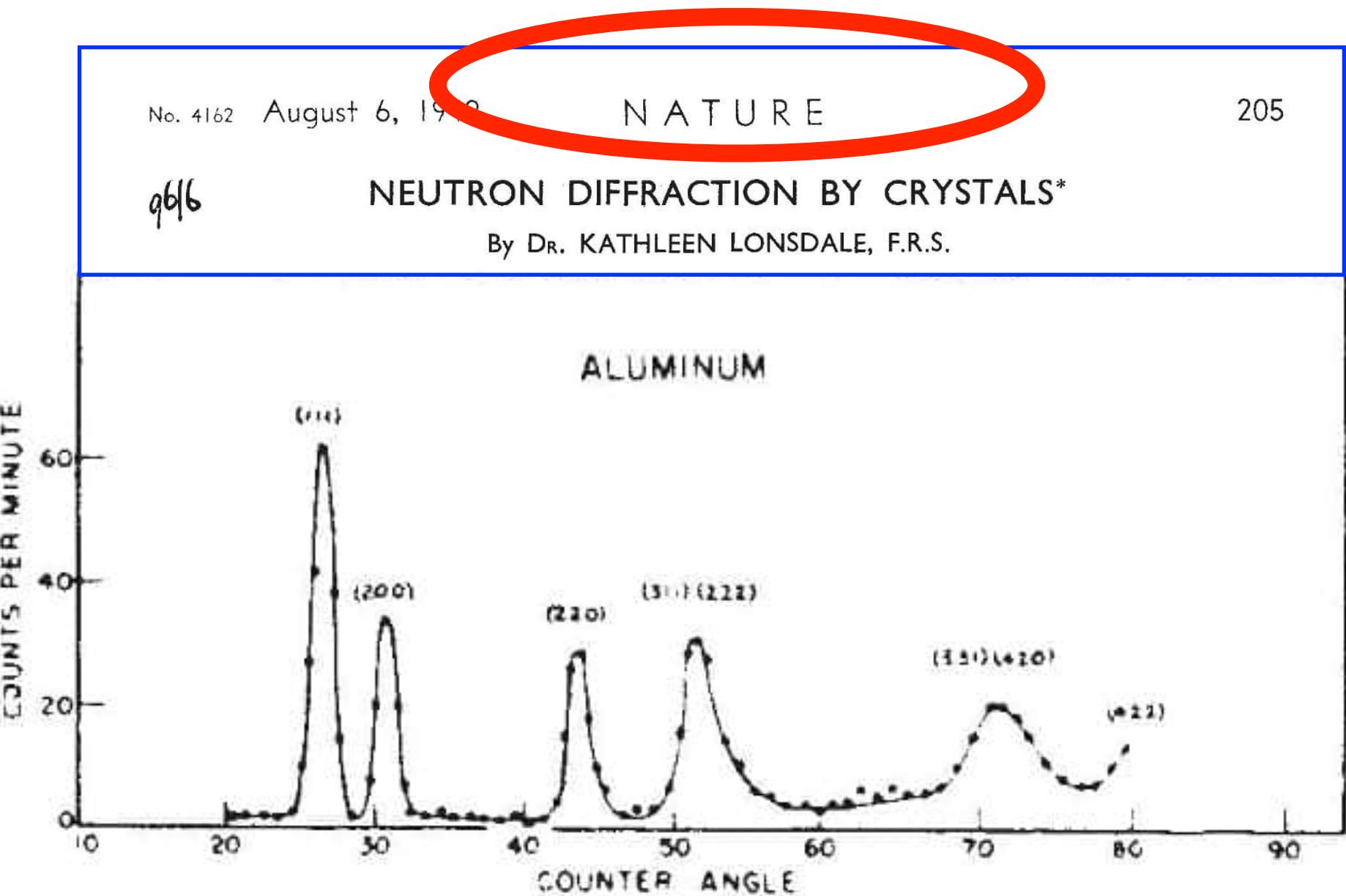


Fig. 1. Diffraction pattern for powdered aluminium in counts per minute versus counter angle ( $2\theta$ ) (E. O. Wollan and C. G. Shull, Oak Ridge)

Wollan's group at Oak Ridge in 1946 laid the foundations for widespread application of neutron diffraction as an important research tool.

# **Neutrons Born Fast**

**(High Energy Physics Jargon)**

# NEUTRONICS: Neutron Transport and Interaction with Matter

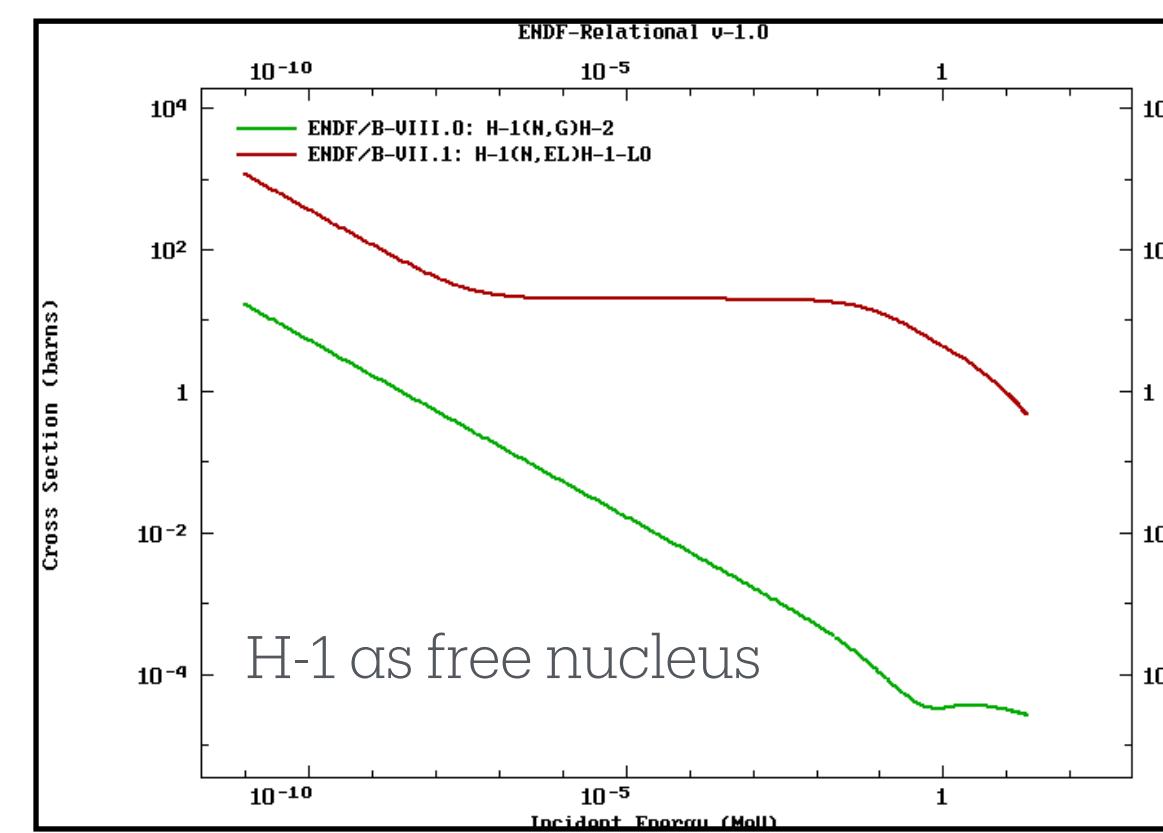
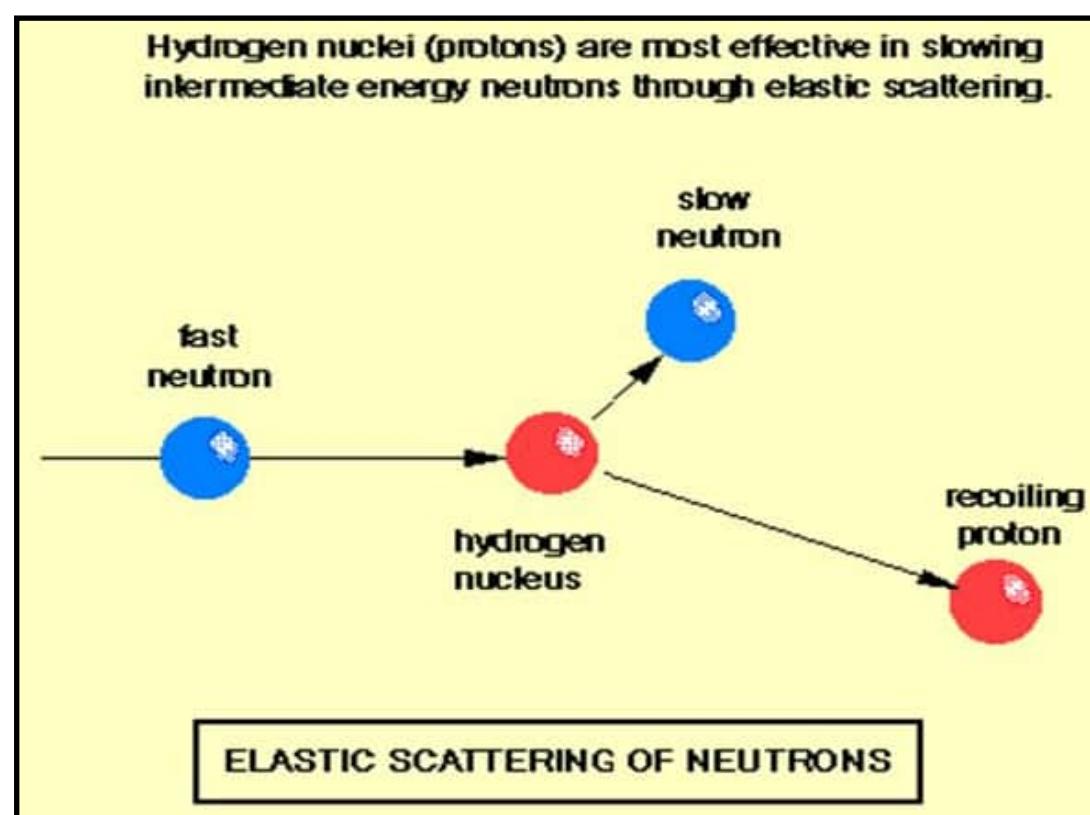
- Wide range of energy —> wide range of mechanism of interactions → different scientific jargon

Elastic collision: conservation of KE and momentum

## Fast neutron down to thermal:

Elastic collisions are the preferred and more efficient way to reduce the neutron energy.

In an elastic collision neutron-nucleus behave as billiard balls → no excitation of the compound nucleus. The total kinetic energy before and after the collision is preserved

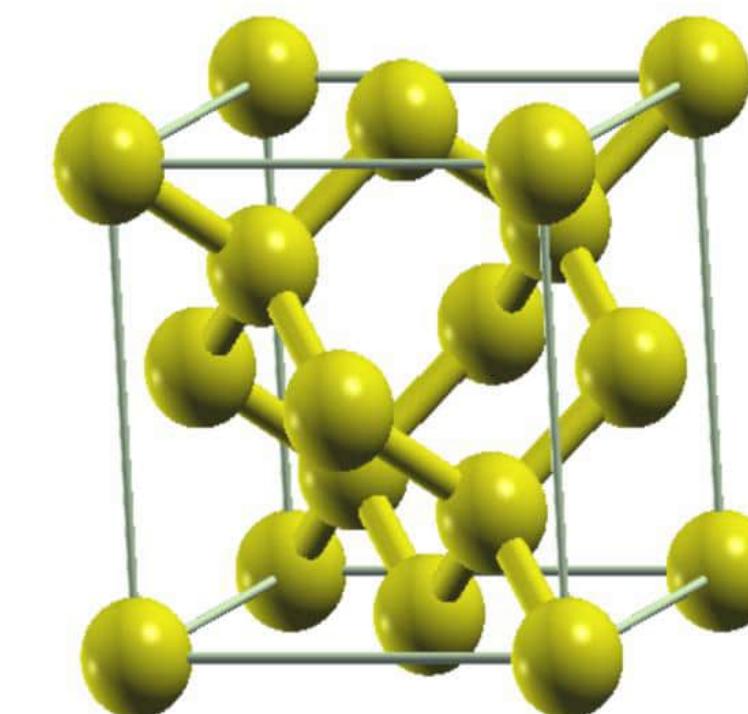
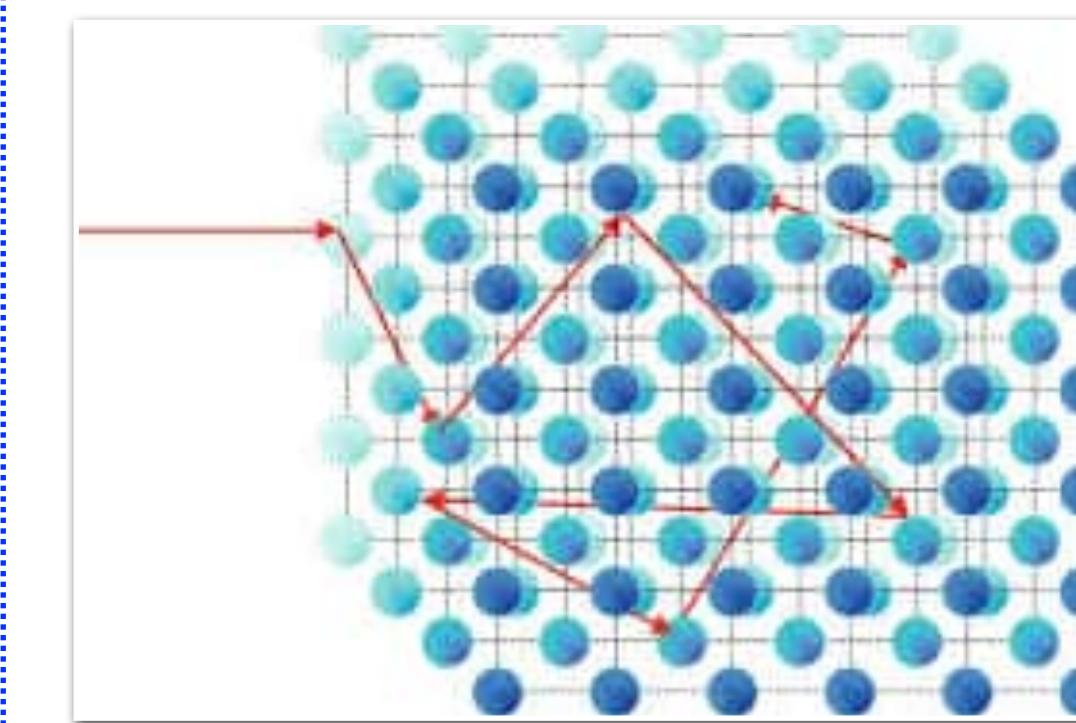


High Energy Range

## In the eV range and below:

Neutrons see nuclei as bound in molecule several vibrational modes can be activated as a mechanism to lose energy.

the scattering nucleus is considered at rest ( $KE=0$ ) neutron scattering elastically only changes direction



Energy range < few eV

# Energy Dependent Interaction with Matter

Designing a neutron source for scattering experiment means optimise the leakage from the source to the instruments.

- Diffraction
- Magnetic scattering
- Elastic Scattering

## Nuclear Reactions:

Radiative Capture:  $A(n,\gamma)$

Other Captures ( $n,p$ ) or ( $n,\alpha$ ):  $A(n,x)$

Inelastic Scattering  $A(n,n')A^*$

Nuclear Fission ( $n,f$ )

total microscopic cross section

$$\sigma_T = \sigma_{el} + \sigma_{inl} + \sigma_\gamma + \sigma_p + \sigma_\alpha + \sigma_f + \dots$$

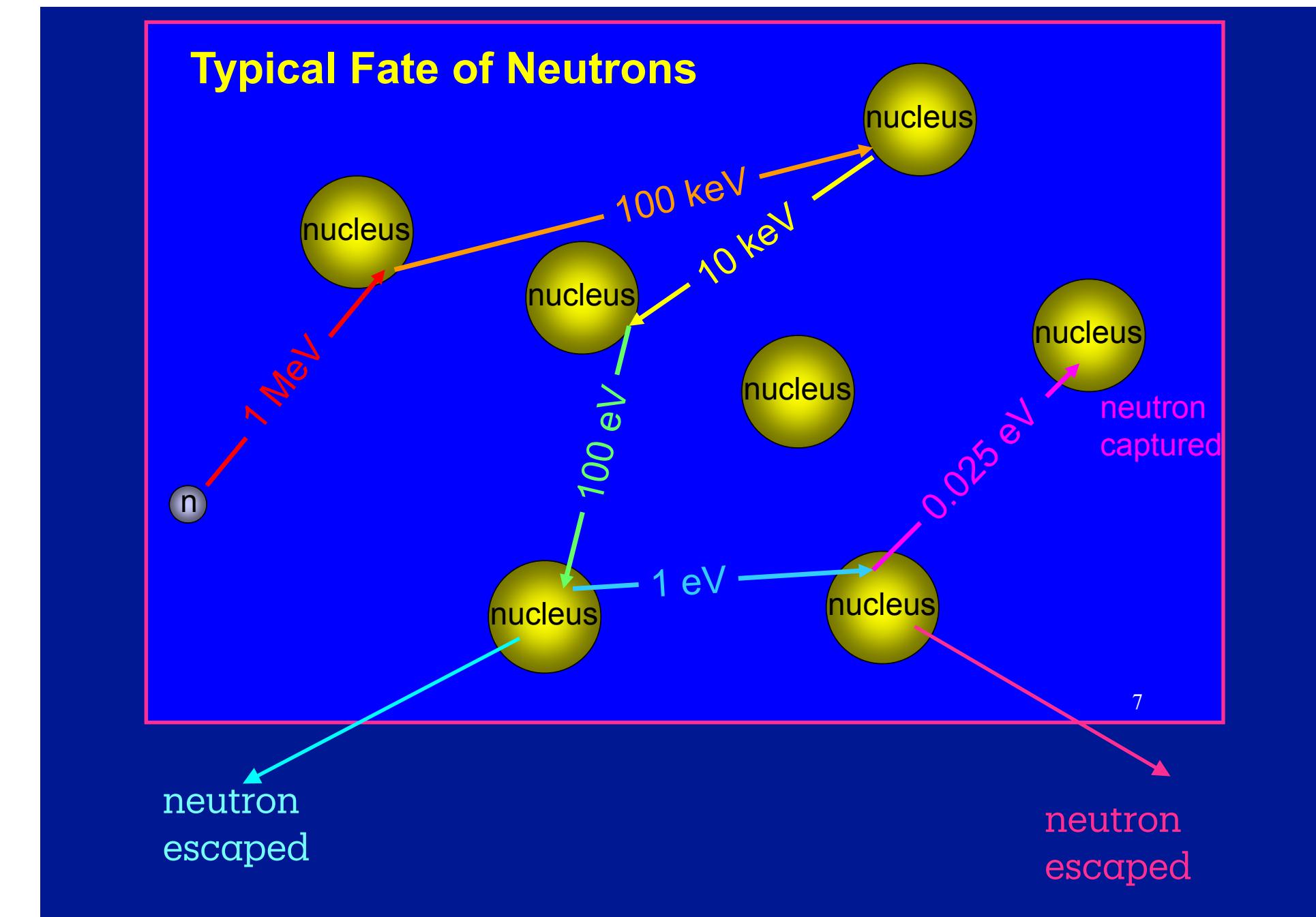
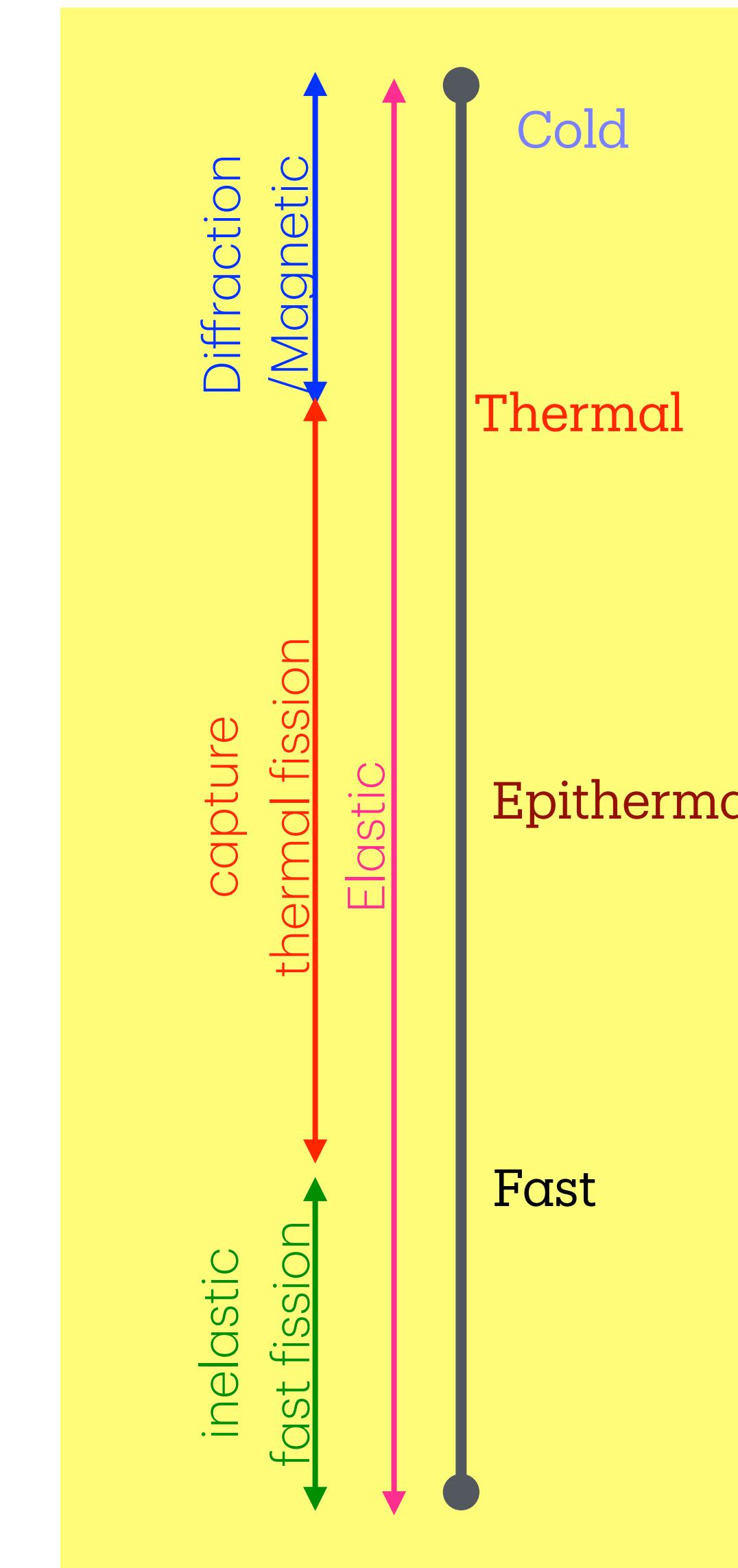
total macroscopic cross section

$$\Sigma_T = N\sigma_T$$

Reaction rate:

$$RR(E) = N\sigma_i(E)\Phi(E) = \Sigma_i(E)\Phi(E)$$

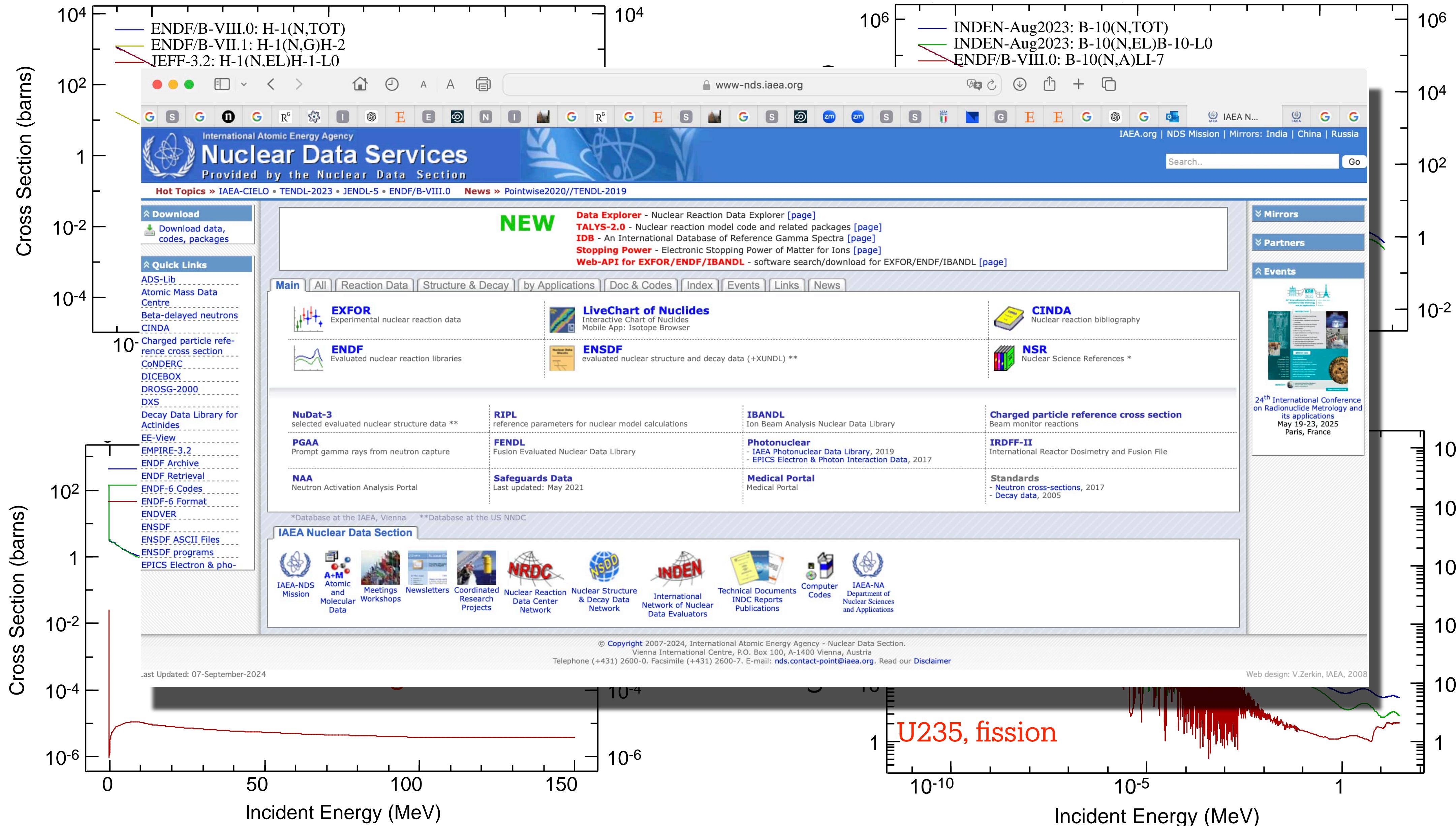
with secondaries



definition	energy range*
ultra cold	< 0.3 $\mu$ eV*
cold	0.3 $\mu$ eV to 0.01eV*
thermal	0.01 eV to 0.3 eV*
epithermal	0.3 eV to 10 keV
fast	10 keV to 20 MeV
Ultra fast	>20 MeV

\* slow neutrons

# Moderators, Absorbers, Fissile Nuclei



# **Neutron Sources**

**(the Physics: Phenomenological Description)**

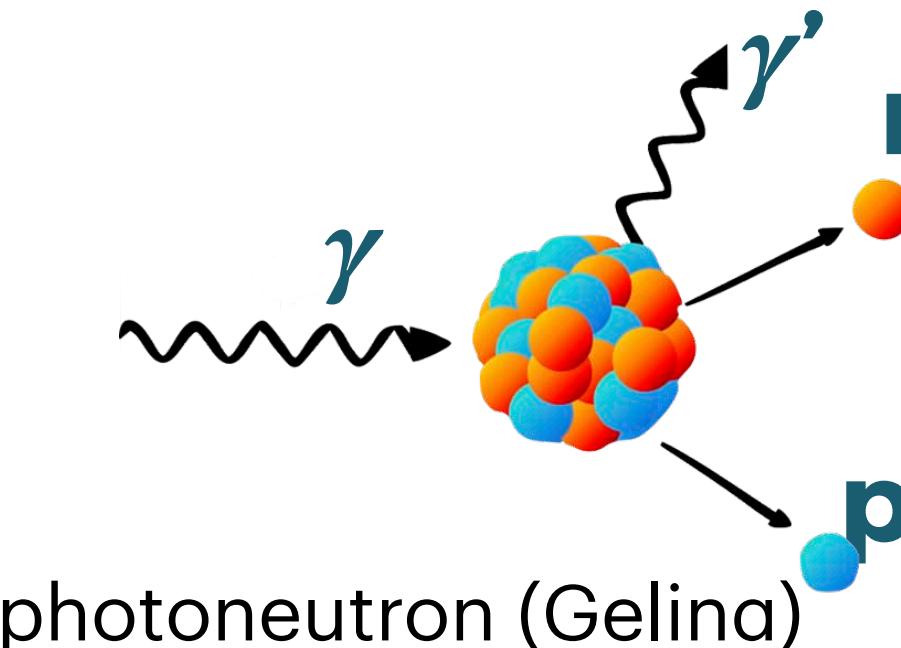
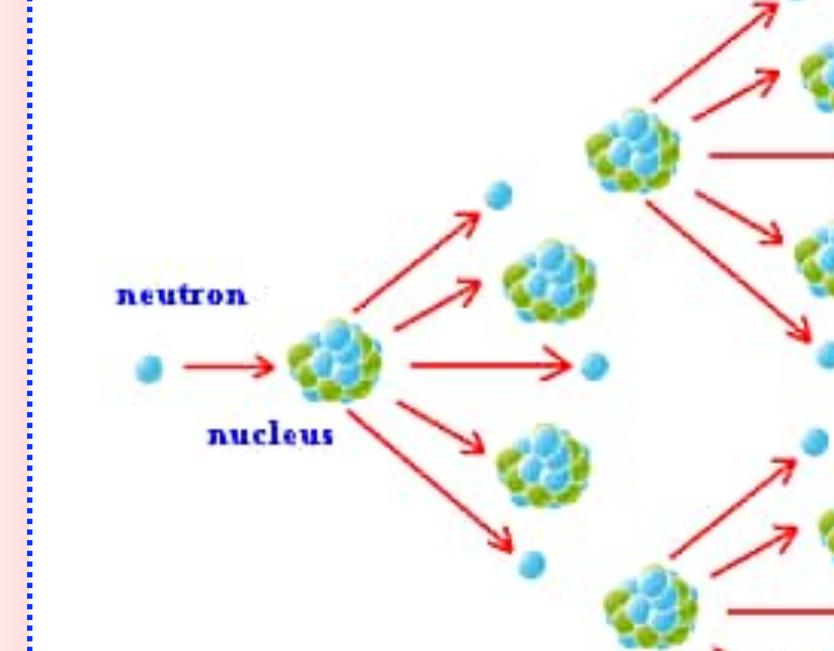
# How To Free Neutrons?

A general overview of the main physics processes

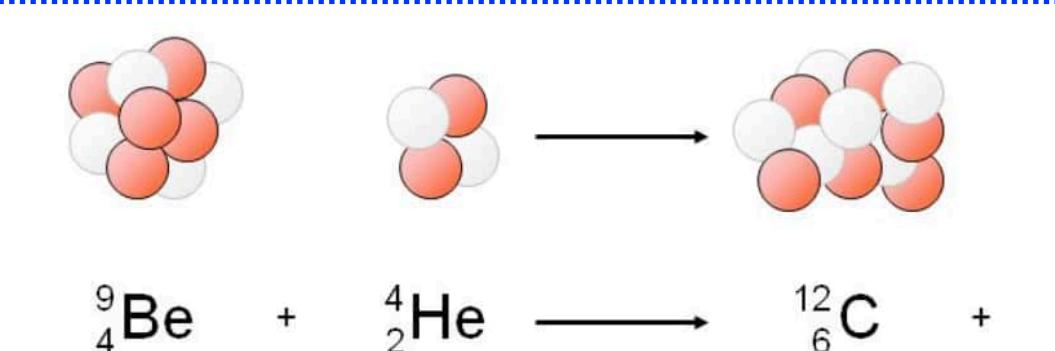
Process	Source Scale
Fission	Reactor
Spallation	Large acc. complex
Photonuclear/ photofission	Large acc. complex
Fusion	intermediate
Laser p acceleration + nuclear reaction	Compact Source (still under study and development)
Radioactive decay + nuclear interaction	Small laboratory Source
Direct radioactive decay	Rare uncontrolled events in large scale facilities (noise source)



## 1. Fission (ILL)

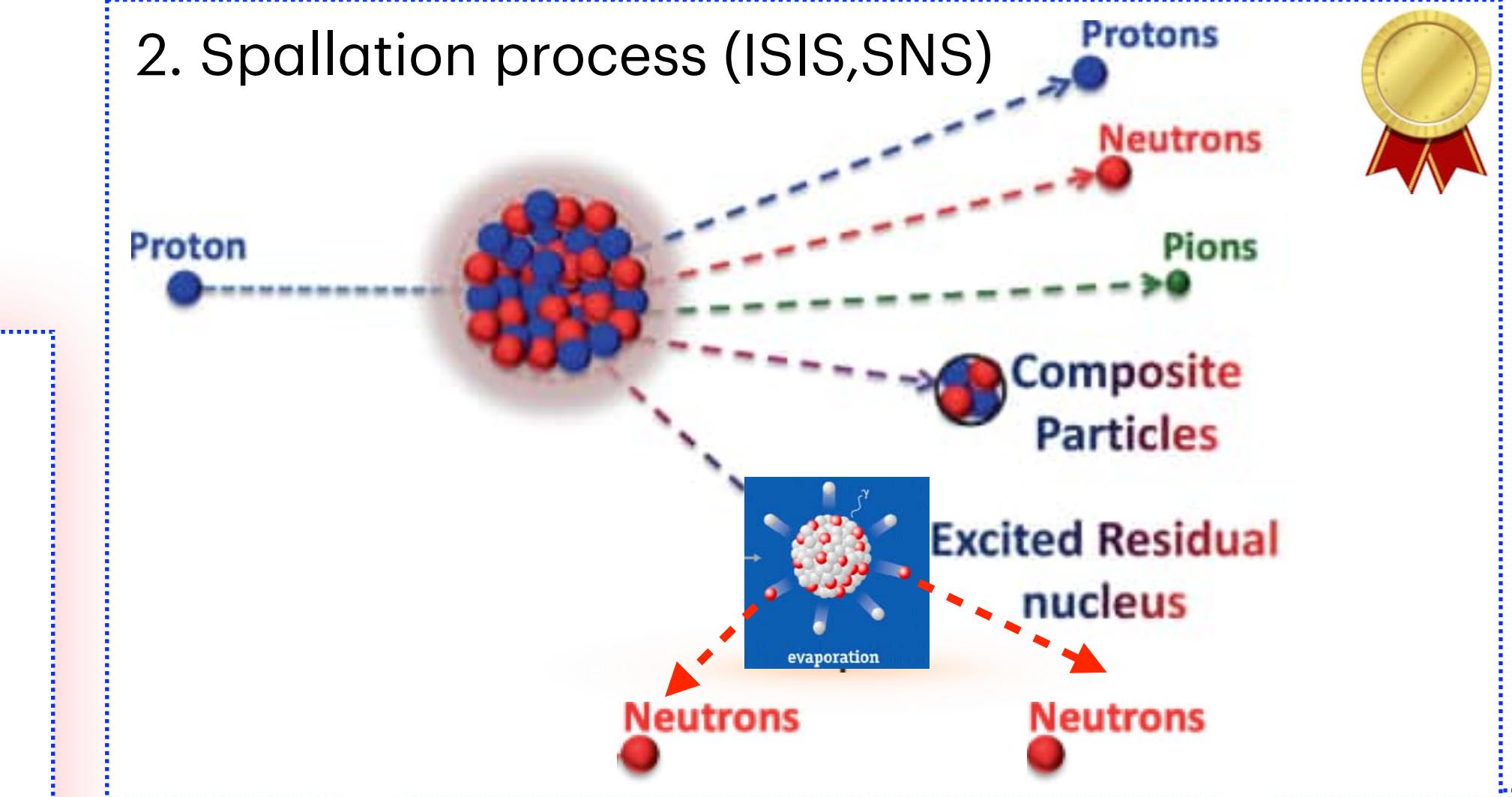


## 3. photoneutron (Gelina)

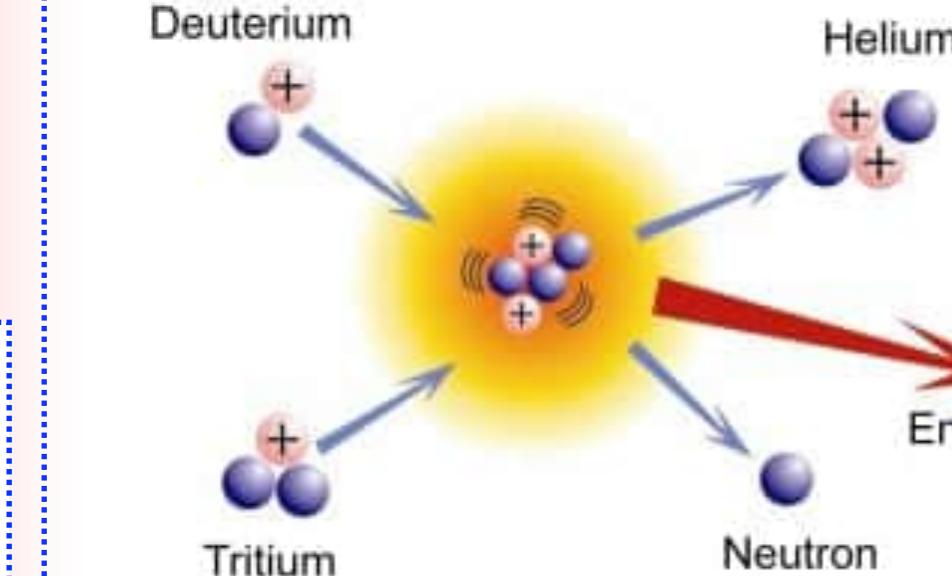


## 6. nuclear reaction

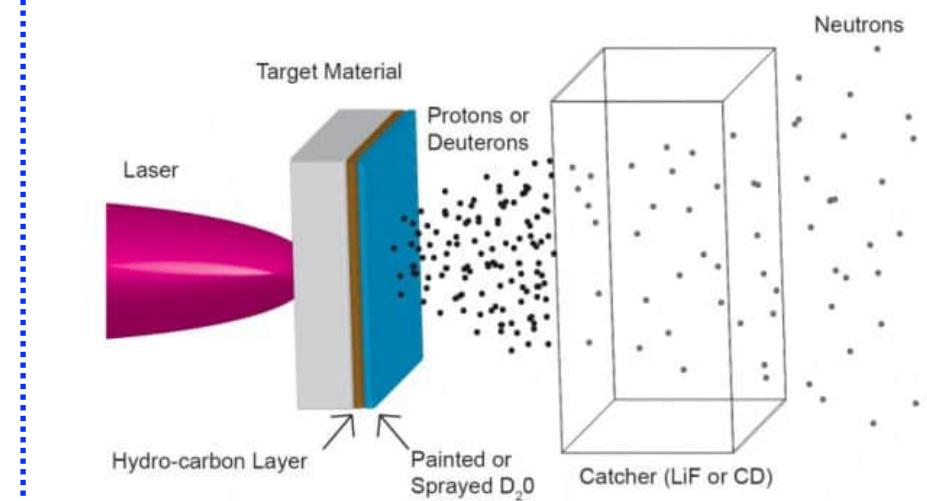
## 2. Spallation process (ISIS,SNS)



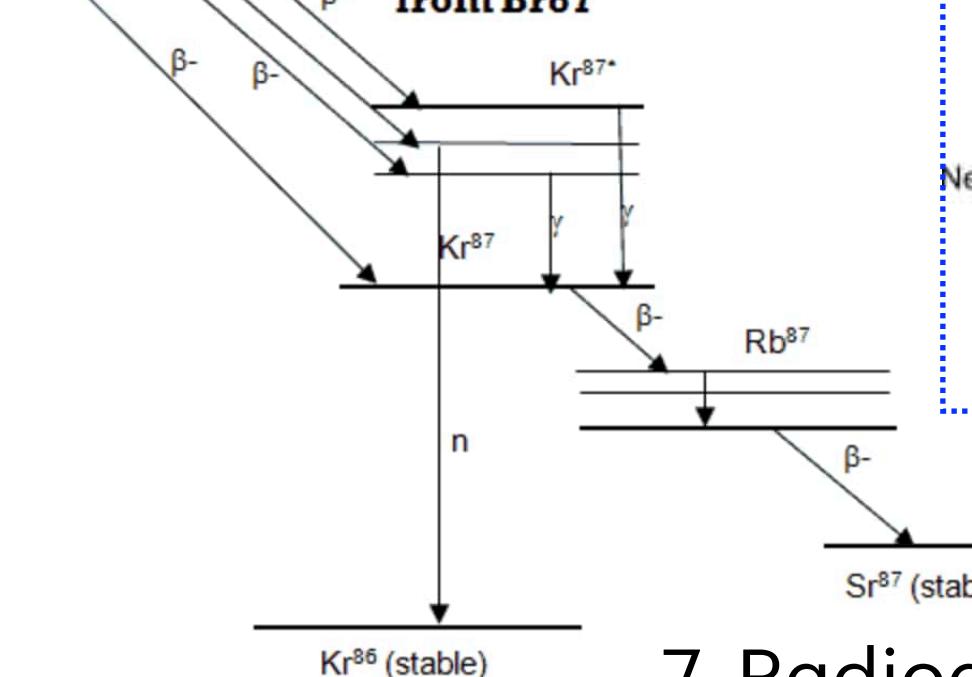
## 4. Fusion



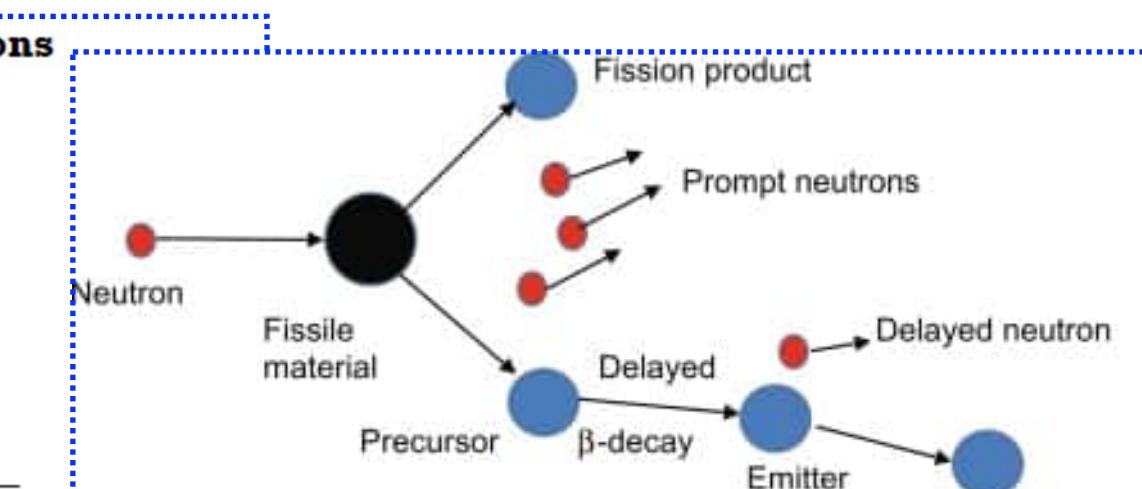
## 5. Laser



## Origin of delayed neutrons from Br87

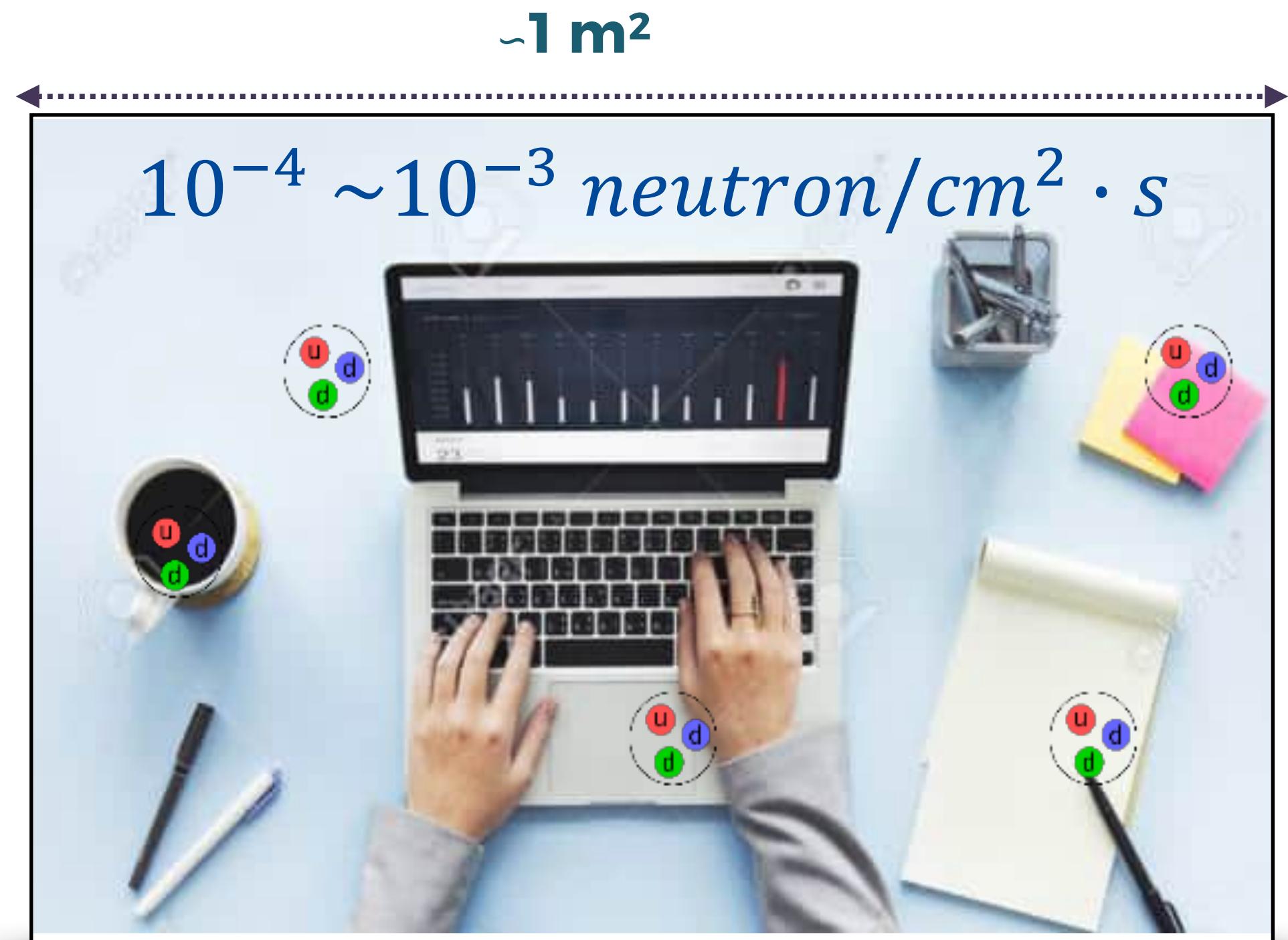
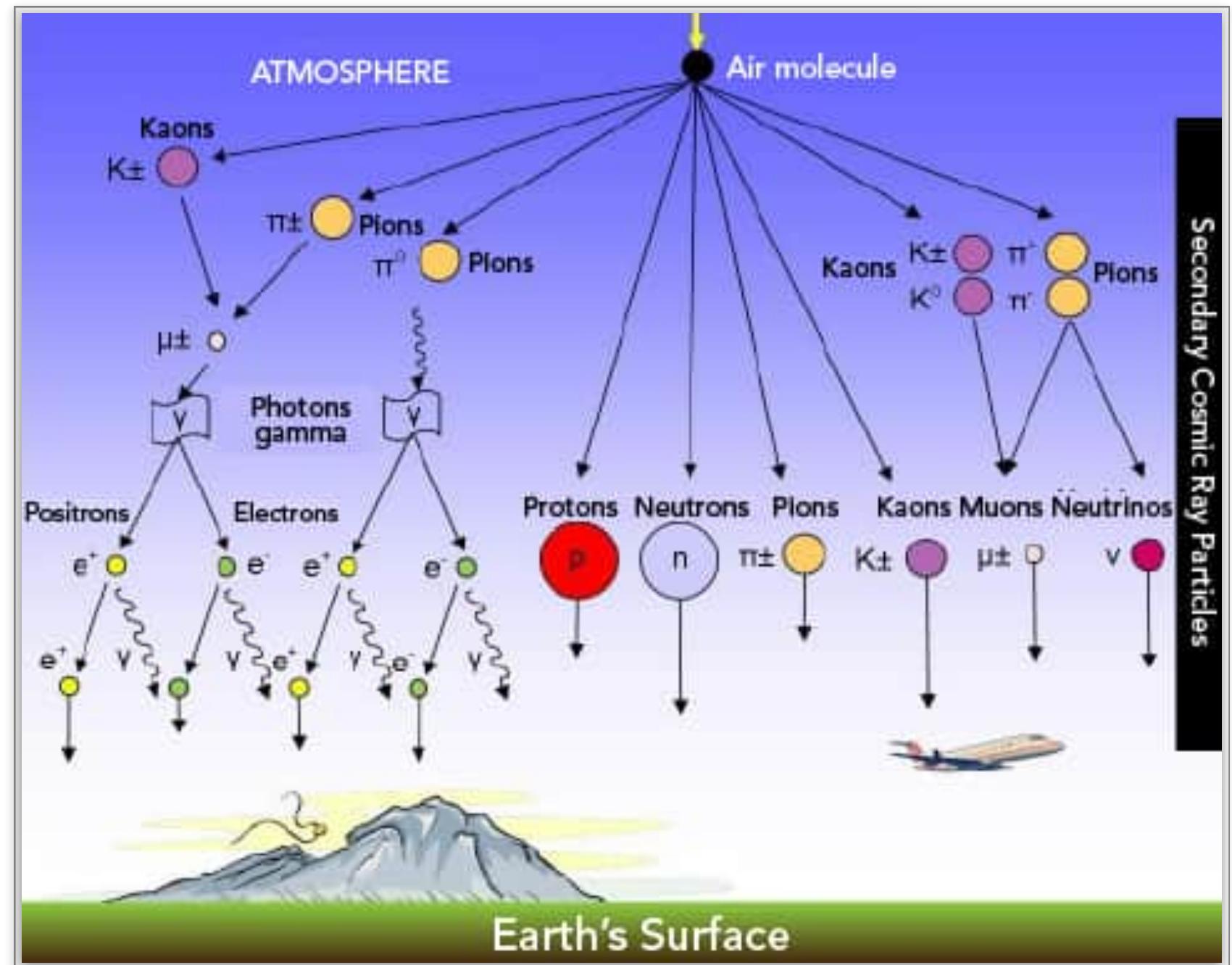


## 7. Radioactive decay



# Neutrons from Cosmic Rays

The “oldest” neutron spallation source on earth is provided by Nature



Neutron background around us is coming mainly from spallation of high energy protons on Oxygen and Nitrogen nuclei in the atmosphere....

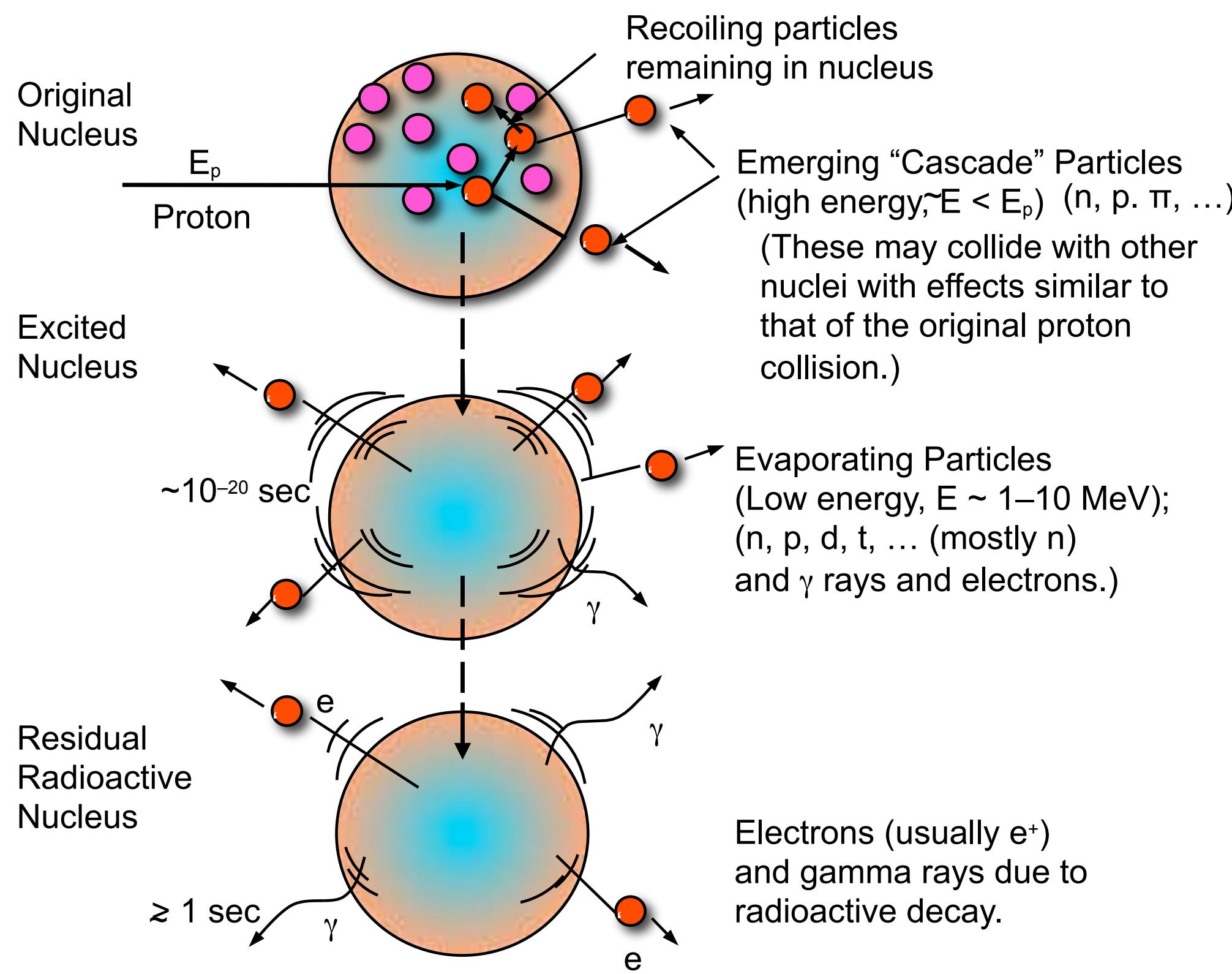
The thermal neutron flux at the Earth's surface is  $\sim 10^{-4}\text{--}10^{-3} \text{ n/cm}^2/\text{s}$ , varying with atmospheric pressure (i.e., weather)

# **Spallation vs Fission**

**(... and Photo-Production)**

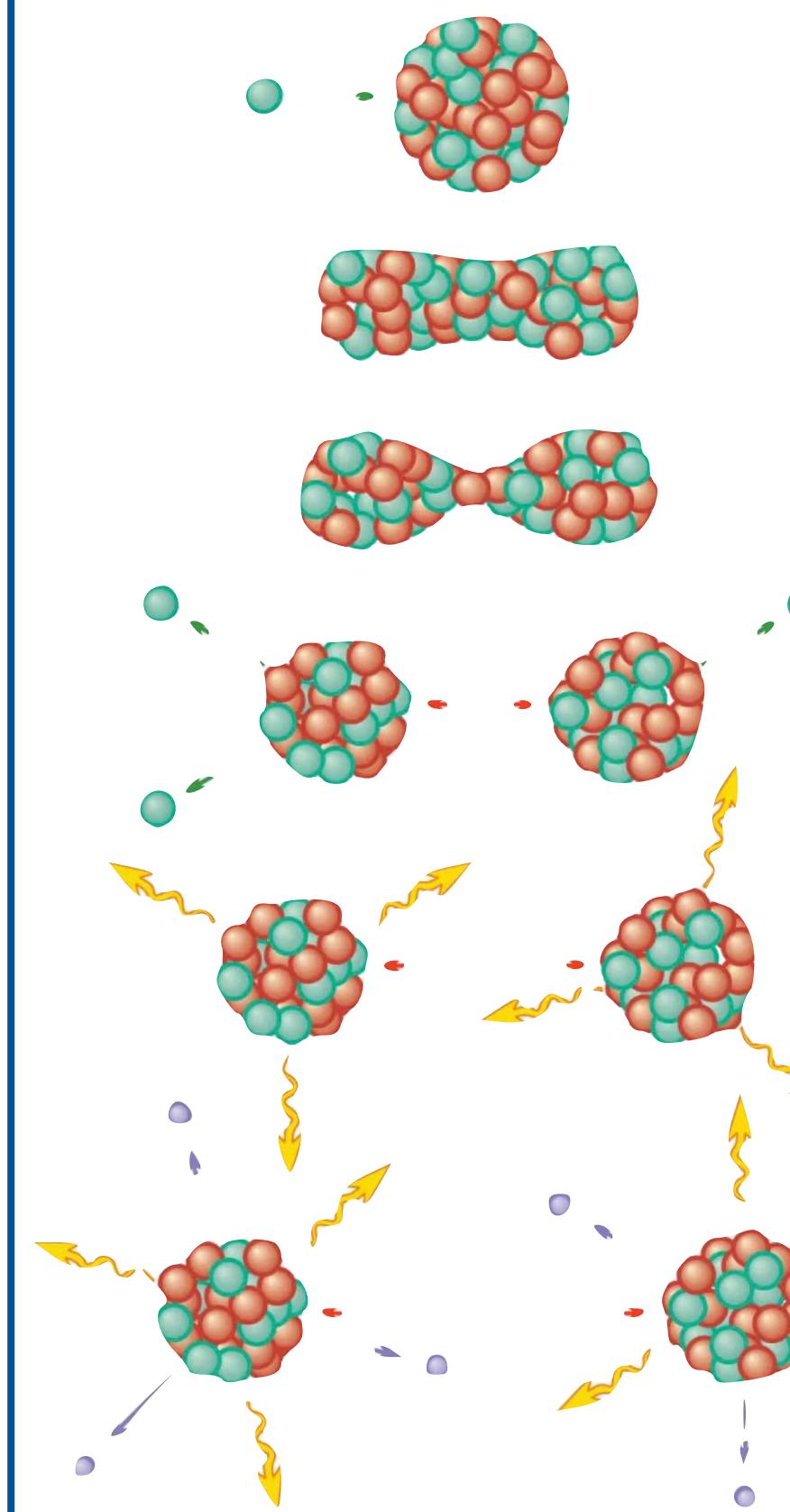
# Mechanisms of Neutron Production

## Spallation



Spallation is an efficient process for releasing neutrons from nuclei. It is an endothermic process and can be in principle triggered in any nucleus but the neutron yield increases with the atomic mass of nucleus

## Fission



The neutron strikes the nucleus and is absorbed.

The absorbed neutron causes the nucleus to undergo deformation.

In about  $10^{-14}$  second, one of the deformations is so drastic that the nucleus cannot recover.

The nucleus fissions, releasing an average of two to three neutrons.

In about  $10^{-12}$  second, the fission fragments loose their kinetic energy and come to rest, emitting a number of gamma rays. Now the fragments are called fission products.

The fission products loose their excess energy by radioactive decay, emitting particles over a lengthy time period (seconds to years).

Fission reactor technology is intrinsically neutrons poor: only few excess neutrons can be made available

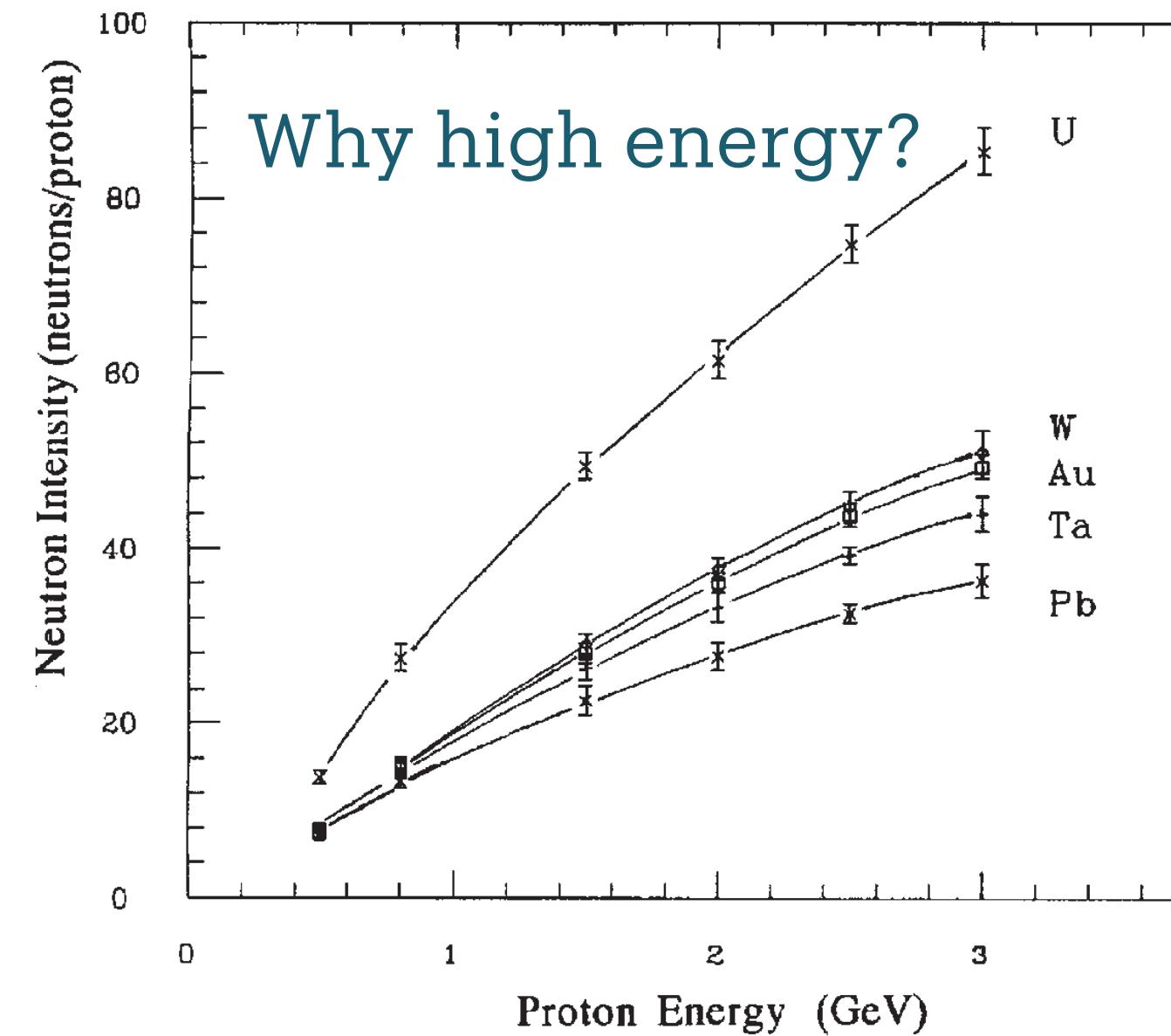
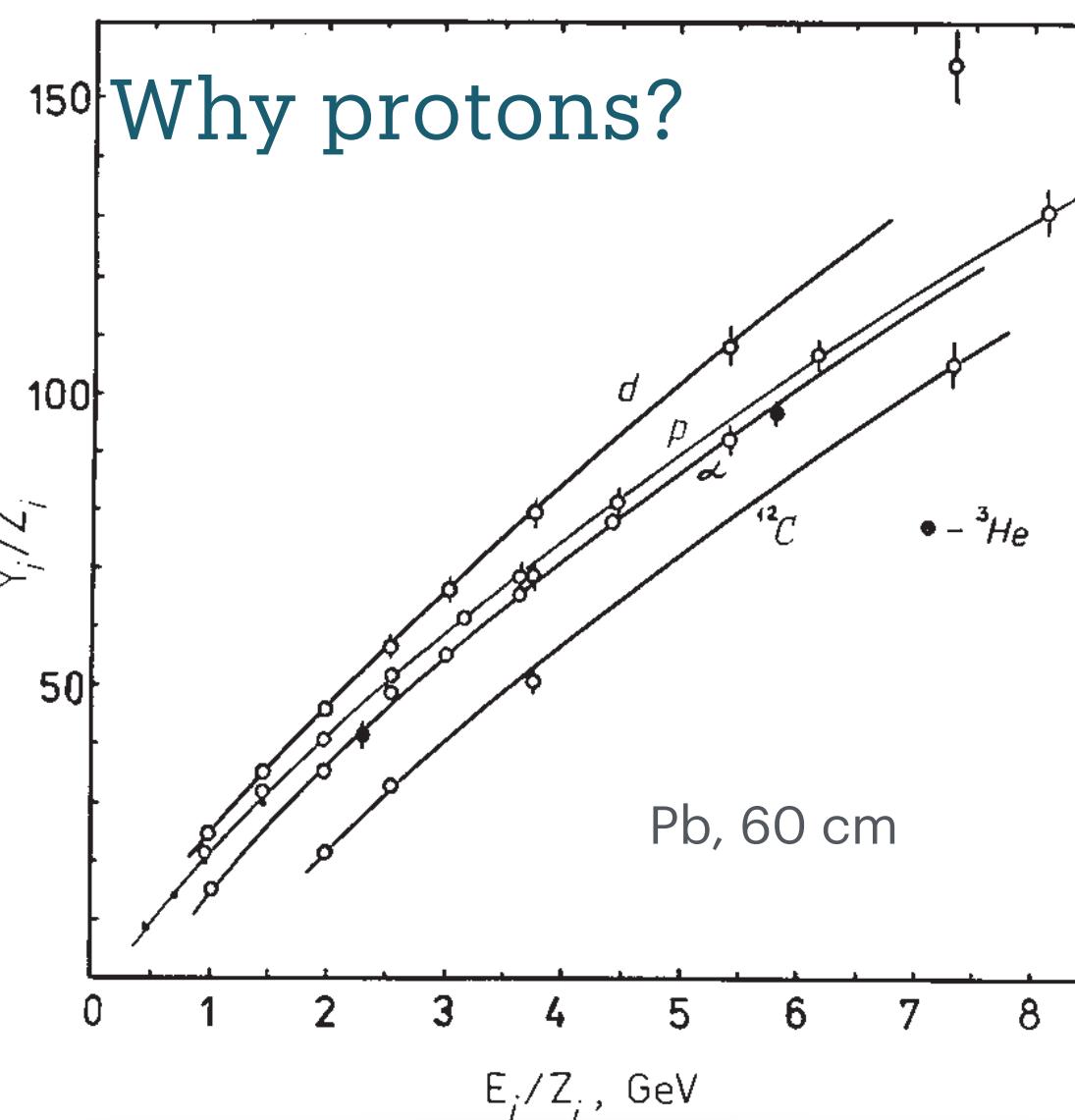
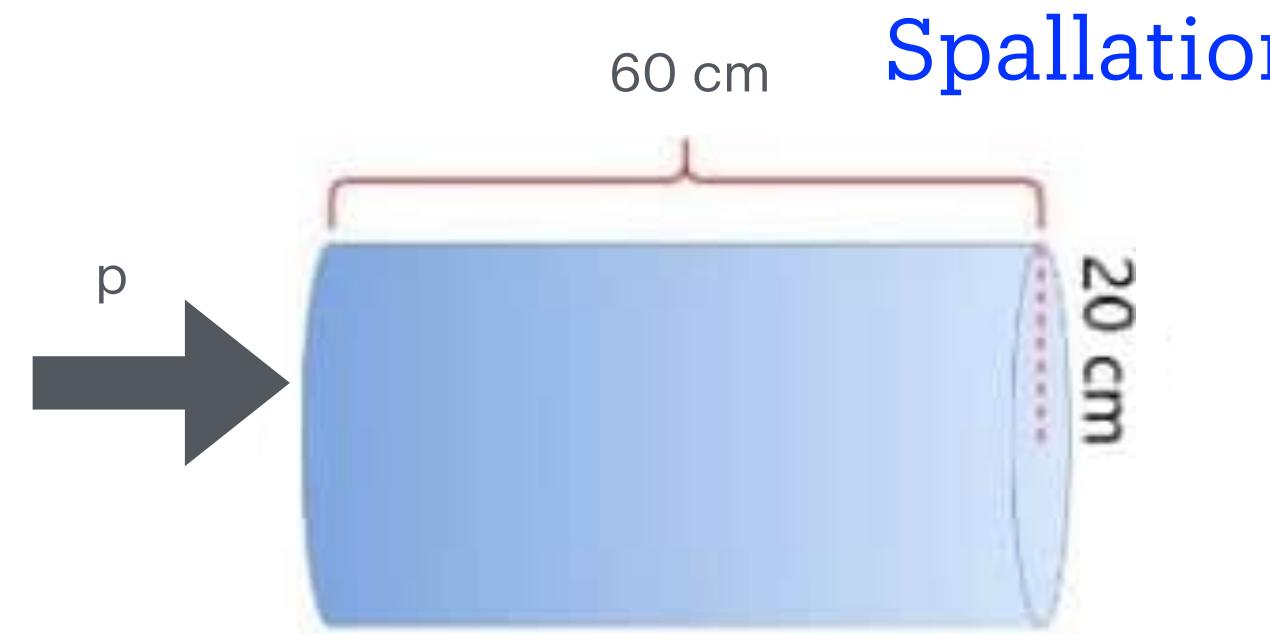
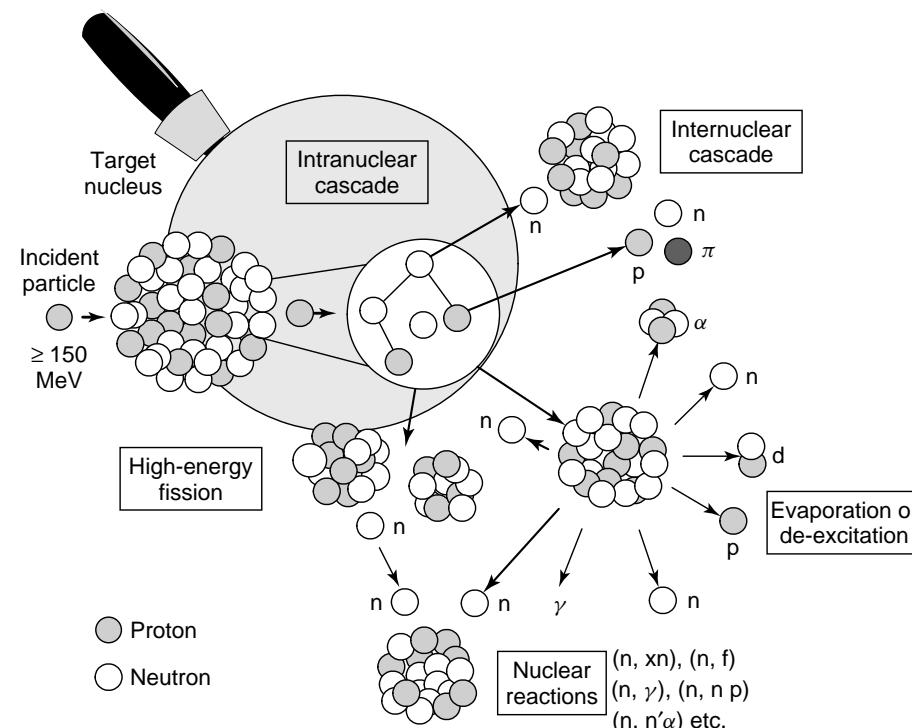
Neutrons    Protons    Beta particles    Gamma rays

Fissile

U233  
U235

Pu239  
Pu241

# Neutron Yield and Leakage



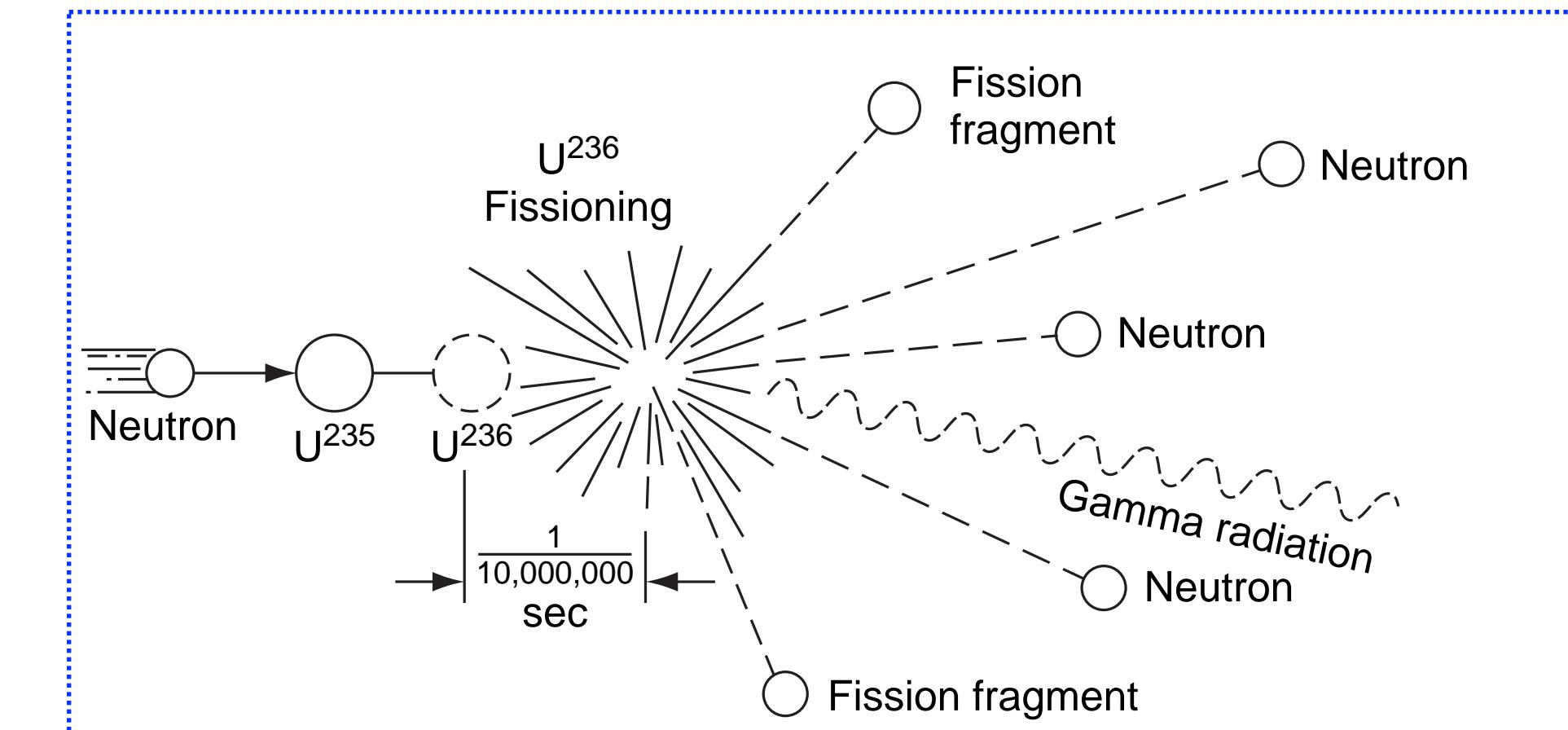
$$Y(E, A) = \{E \text{ (GeV)} - 0.12\} \times \{A + 20\} \times 0.1 \text{ (n/p).}$$

Yield for W target @ 800 MeV:  $Y(0.8 \text{ GeV}, 183) = 14 \frac{n}{p}$

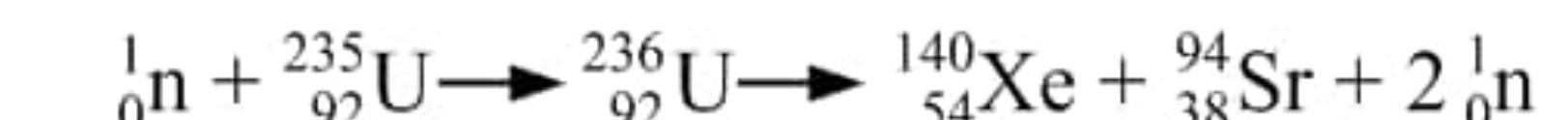
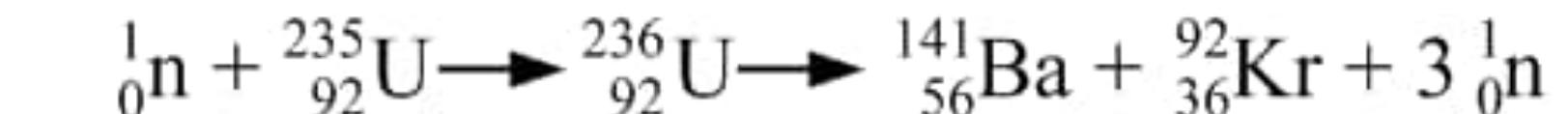
Y = Leakage (neutron absorption inside the target is negligible)

$$\text{ISIS}(200 \mu\text{A}) \rightarrow 1.25 \cdot 10^{15} \frac{p}{s} \rightarrow \text{Leakage}(\text{ISIS}) \simeq 2 \cdot 10^{16} \frac{n}{s}$$

## Fission



**Production:** an average of 2.5 neutrons/fission



**Leakage:** 1 n is needed to sustain chain reaction →

1 neutron/fission available for scattering

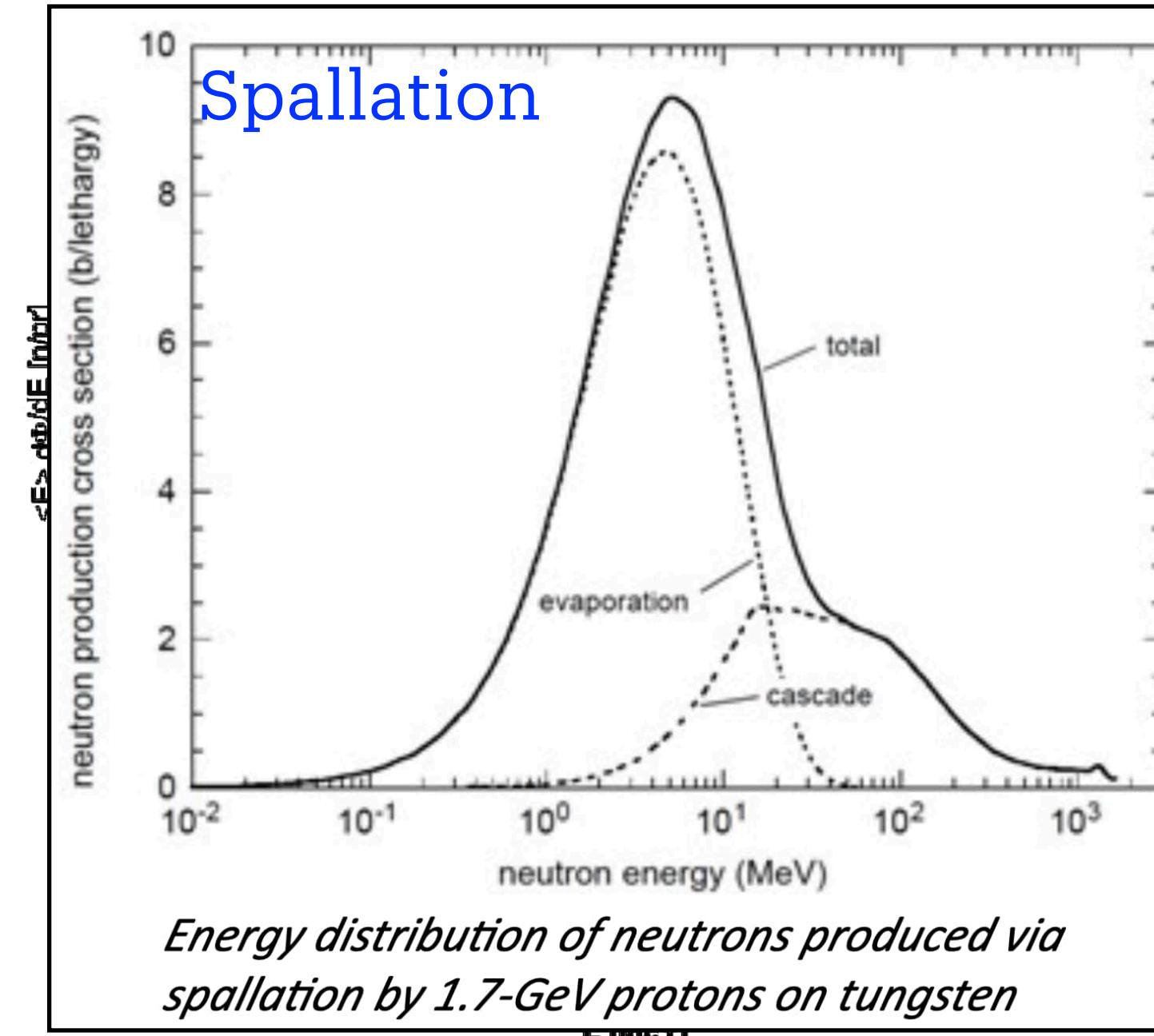
$$\text{Rate of Fissions(ILL)} = \frac{\text{Power}}{\text{energy from one fission}}$$

$$\text{Leakage(ILL)} = \frac{58 \cdot 10^{16}}{200 \cdot 10^6 \cdot 1.6 \cdot 10^{-19}} \simeq 2 \cdot 10^{18} \frac{n}{s}$$

# Neutron Energy Spectrum

Difference in the energy distribution of the emitted neutrons

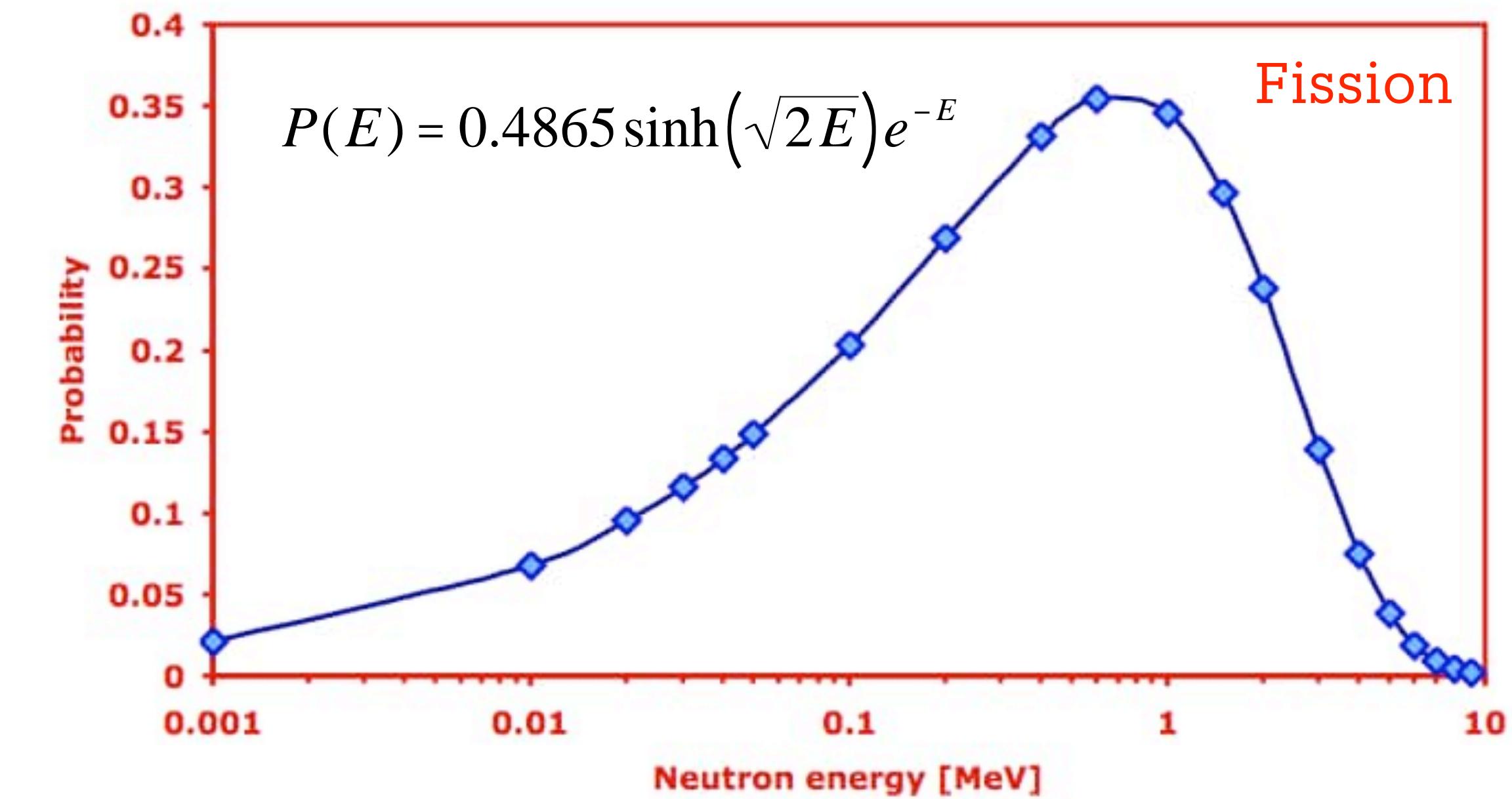
Figure 11: Energy spectrum of neutron escaping from the *bare* target



Spallation neutrons extend to higher energies than fission neutrons:

In a spallation source, high-energy cascade neutrons approach the energy of the incident proton (i.e 800 MeV at ISIS).

Neutrons from evaporation OF HIGH Z NUCLEI are around 1 MeV and are emitted isotropically



Fission spectrum is peaked around 1 MeV and extends up to 10 MeV

Very few neutrons are in the eV range

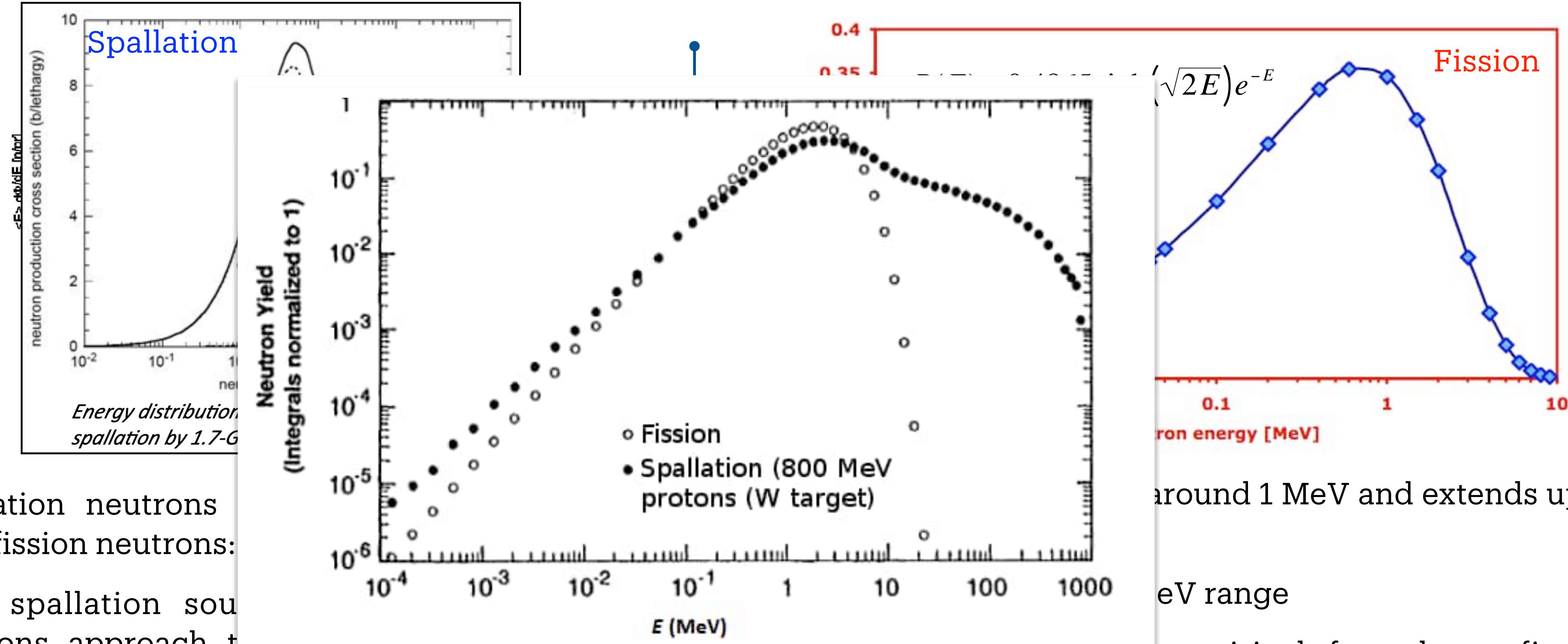
The Watt function is an empirical formula to fit the experimental data

The shape is pretty the same for thermal fissions of different fissile nuclei

# Neutron Energy Spectrum

Difference in the energy distribution of the emitted neutrons

Figure 11: Energy spectrum of neutron escaping from the *bare* target



around 1 MeV and extends up to eV range

The Watt function is an empirical formula to fit the experimental data

The shape is pretty the same for thermal fissions of different fissile nuclei

# Photonuclear (vs Spallation)

Same spectrum shape but different yield efficiency

Bremsstrahlung Photons interact with nuclei causing the emission of nucleons (neutron, protons, pion, etc)

Photon energy levels above the nuclei binding energy (5-15 MeV) are necessary. Photo-nuclear interactions are threshold reactions: threshold is lower for heavy nuclei

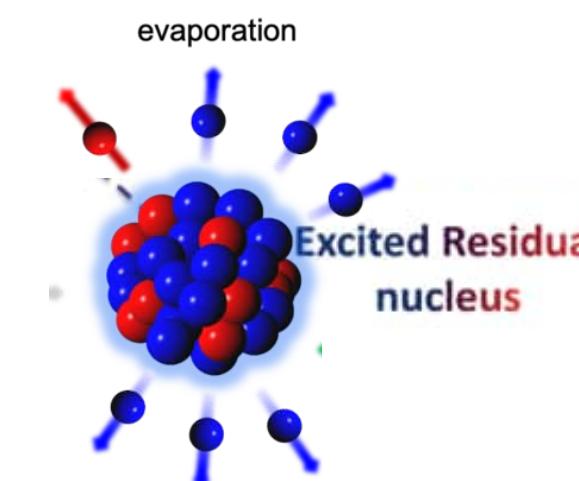
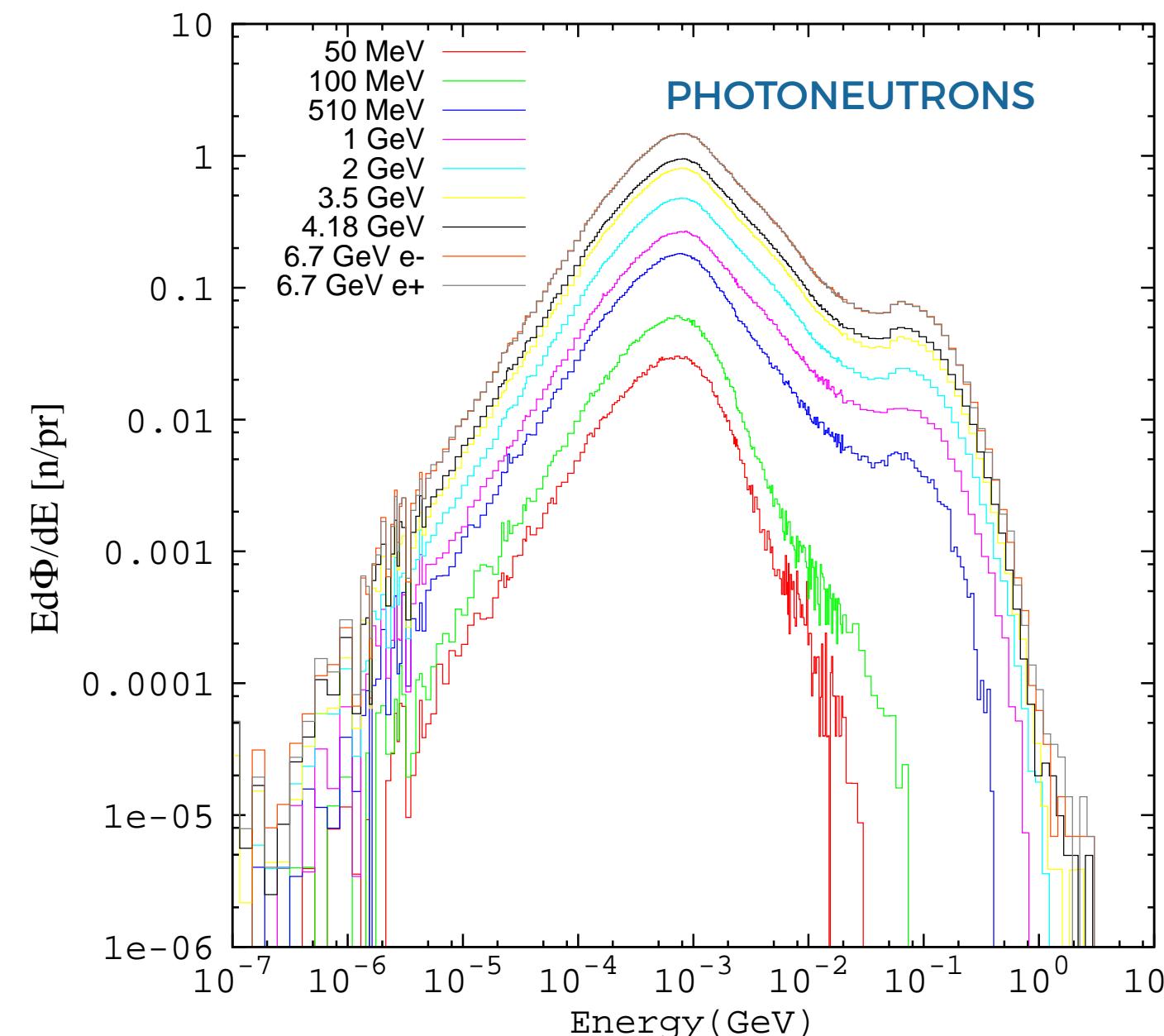


**electrons → gammas → neutrons**



Main mechanism of neutron production: Giant dipole resonance

Neutron flux from the giant resonance process is distributed according to a Maxwellian



$$\frac{d\phi_n}{dE} \propto \frac{E}{T^2} \cdot e^{-\frac{E}{T}}$$

$T = 0.7 \text{ MeV for } W$

Neutron energy spectra from spallation and photonuclear production are pretty similar : they both extend up to the primary  $E$  and 90% of the neutrons are around 1 MeV ( $W$ )....  
but the efficiency of production of neutron is quite different (photo nuclear yield is almost 2 orders of magnitude less than spallation)

Spallation:

**1 GeV p on W target → 14 n/p**

Photonuclear:

**1 GeV e- on W target → 0.7 n/e-**

# Delayed Neutrons

1. Delayed neutrons are emitted by neutron-rich (fission or spallation) fragments that are called delayed neutron precursors.
2. These precursors usually undergo beta decay, but a small fraction of them are excited enough to undergo neutron emission.
3. The emission of neutrons happens orders of magnitude later compared to the emission of the prompt neutrons.

## Relevance in Nuclear reactors

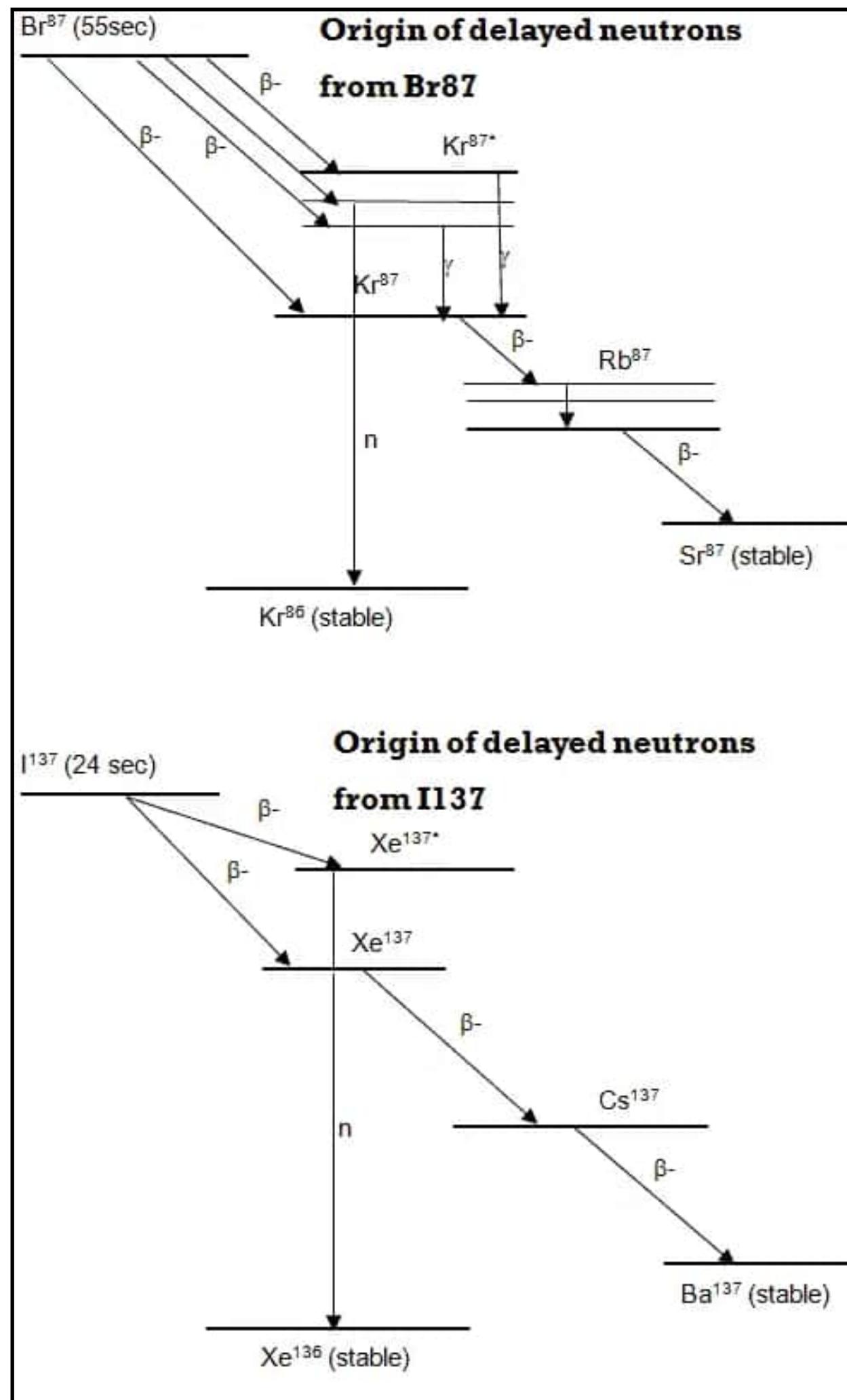
Most of the neutrons produced in fission are prompt neutrons,

"Delayed" neutrons are emitted with half-lives ranging from few milliseconds up to 55 s for the longest-lived precursor (<sup>87</sup>Br)

The presence of delayed neutrons is perhaps the most important aspect of the fission process from reactor control:

Neutron balance for CANDU reactors

Neutrons	Contribution
Prompt Neutrons	99,470,000,000,000
Delayed Neutrons	500,000,000,000
Photoneutrons	30,000,000,000
Spontaneous Fission	1
Total ( $10^{14} \text{ n cm}^{-2} \text{ s}^{-1}$ )	100,000,000,000,000



## Neutronic background in Spallation Facilities

Delayed neutrons from Bromine and Rubidium are produced as a results of spallation process.

Neutrons coming from (alpha, n) reactions with high energy alpha (> 4 MeV) coming from short life alpha decay radionuclides

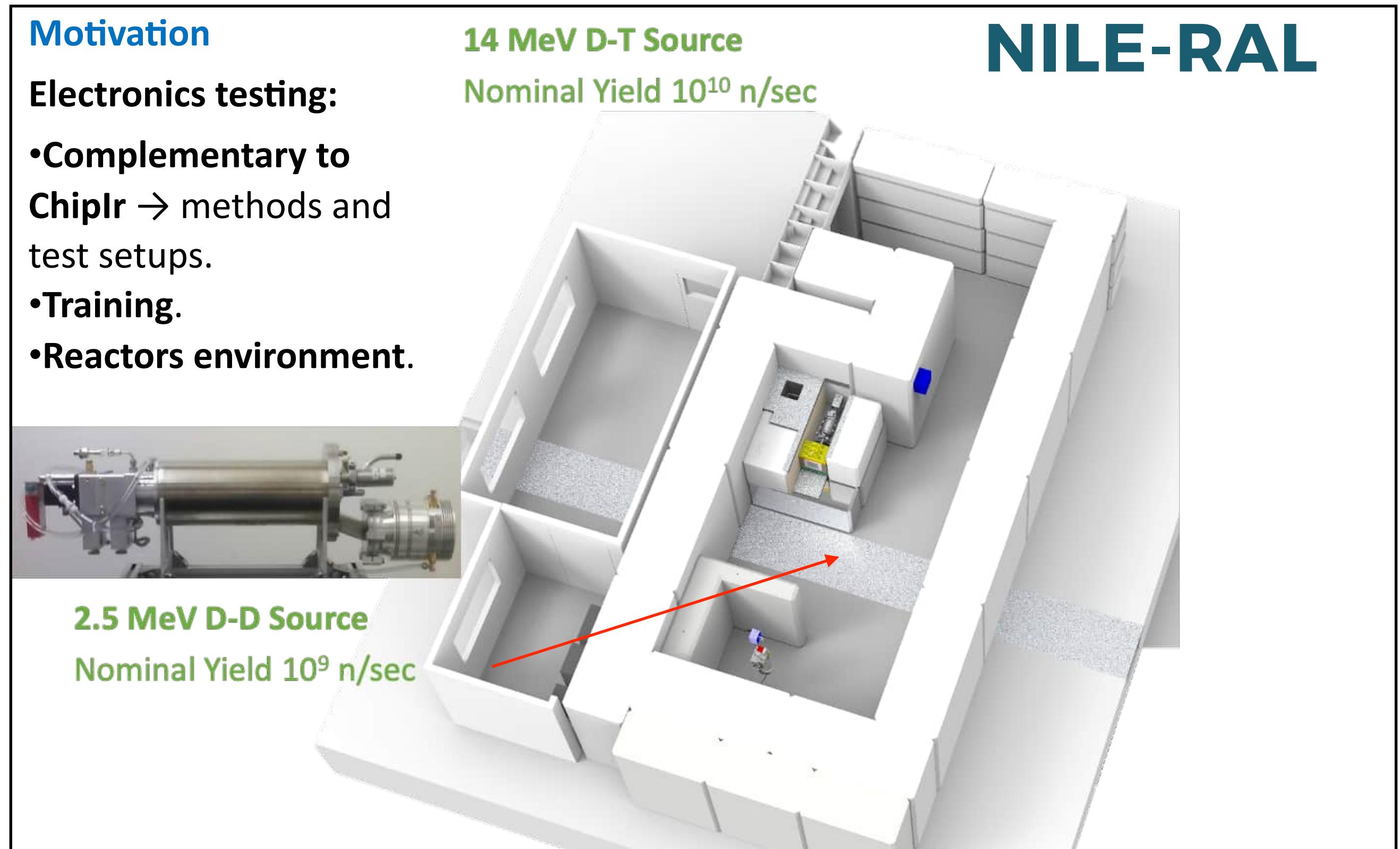
Photoneutron produced in coolant or Be when the target is not irradiated by proton but still kept in loco for cool down purposes

## Delayed neutron precursor's activity at EOI @ ISIS

Z	A	RN	Activity [Bq]	T <sub>1/2</sub>
35	86	Br86	5.06E+09	54-sec
35	87	Br87	1.27E+09	55-sec
35	88	Br88	4.43E+09	15.5 sec

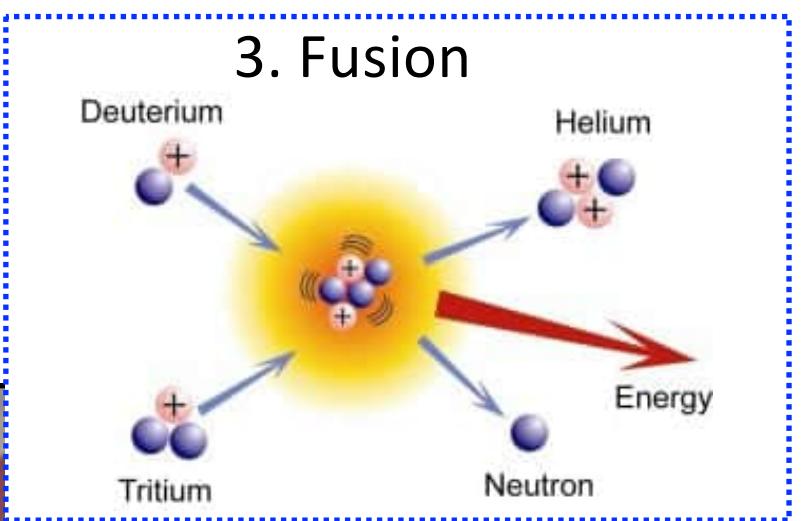
**A Few Words on the Other Production  
Mechanisms of Neutron....**

# Intermediate Flux Sources (Fusion Sources)



Fast-neutrons fields at 14MeV (with a total flux of  $10^{10}$  neutrons/sec) and 2.5MeV (with a total flux of  $10^9$  neutrons/sec).

Neutron D-T & D-D generator(fusion neutrons @ 14 MeV and 2.5 MeV respectively ) up  $10^{11} - 10^8$  n/s



# Low Flux Sources: Laboratory Neutron Sources

- **Transuranic nuclide in low Z matrix:** AmBe  $1 \text{ Ci} \rightarrow 10^6 \text{ n/s}$
- **Photo-neutron source: Sb-Be** Be bombarded by photon with E above the binding energy of nucleons ( $E > 1.7 \text{ MeV}$ ) (short half life) emits neutrons: ( $1 \text{ Ci} \rightarrow 5 * 10^6 \text{ n/s}$ )
- **Spontaneous fission source:** Cf-252



$^{252}\text{Cf}$  (Californium)  $\rightarrow 10^7\text{-}10^8 \text{ n/s}$

- 2.6 year half life; mainly alpha decay but also fission(SF branch ratio about 3.1 percent)
- 1 mg emits  $2.3 \times 10^{11} \text{ n/s}$  with average energy 2.1 MeV
- 3.8 n/fission + 9.7  $\gamma$

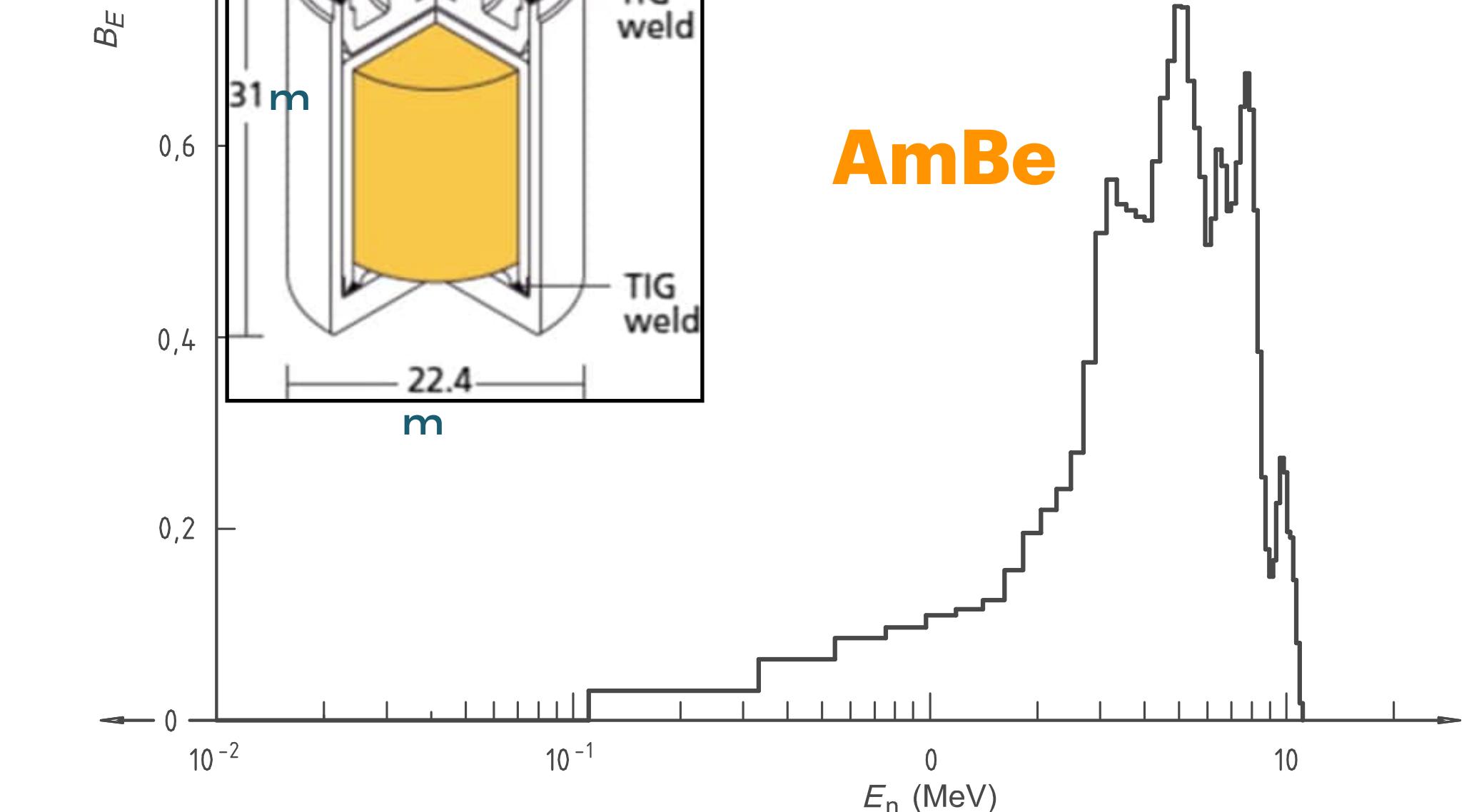
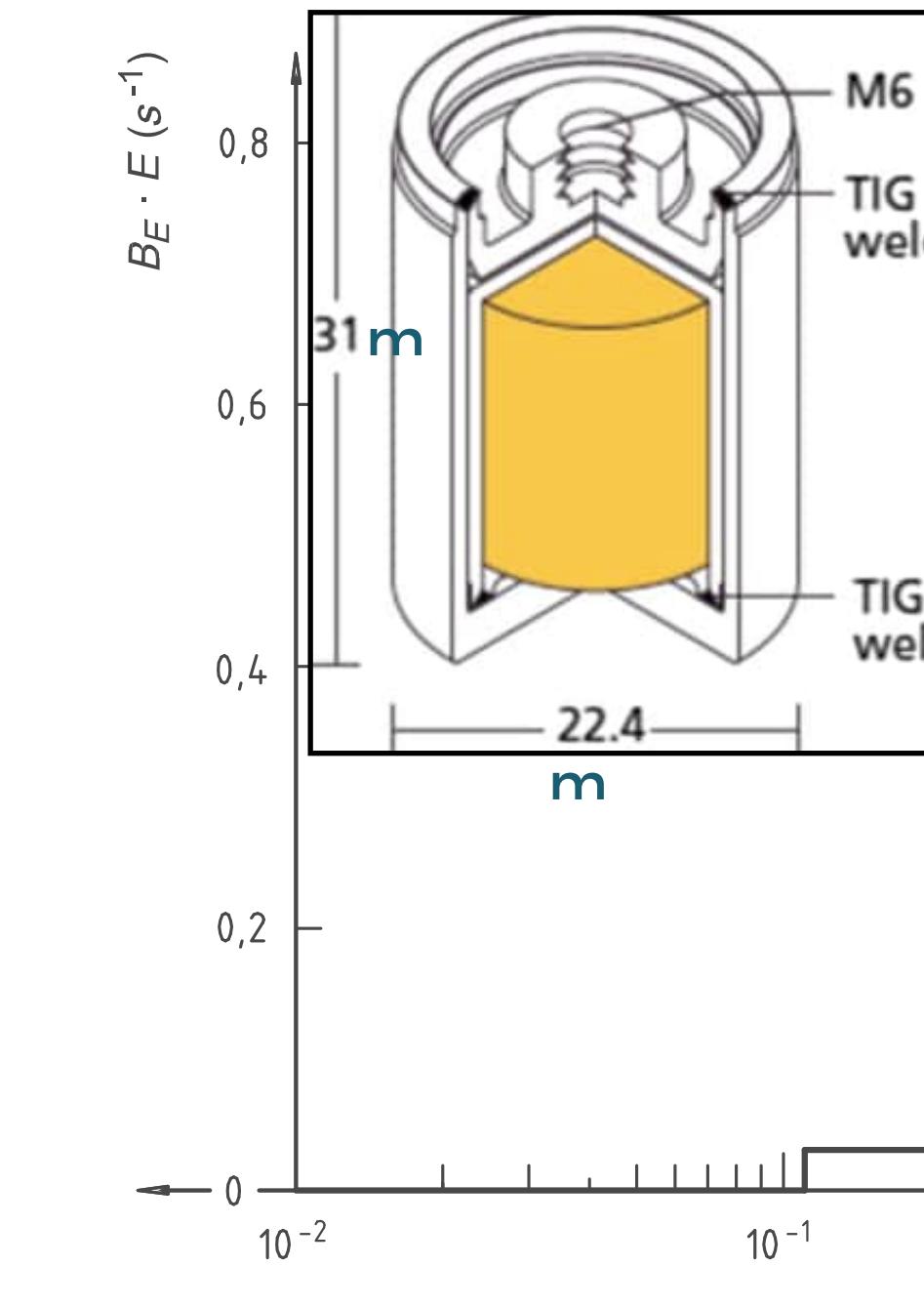
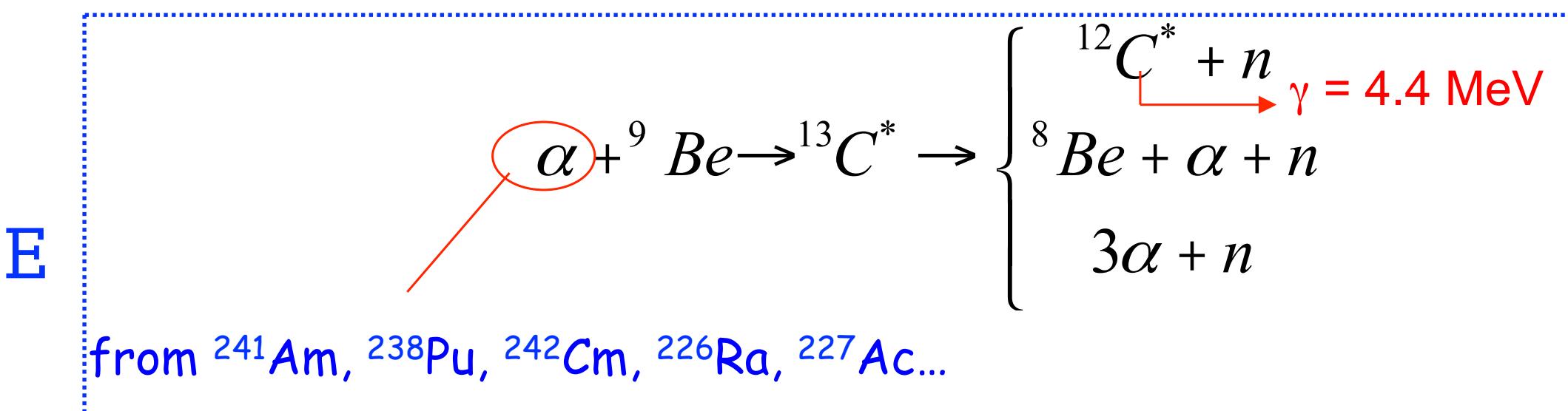
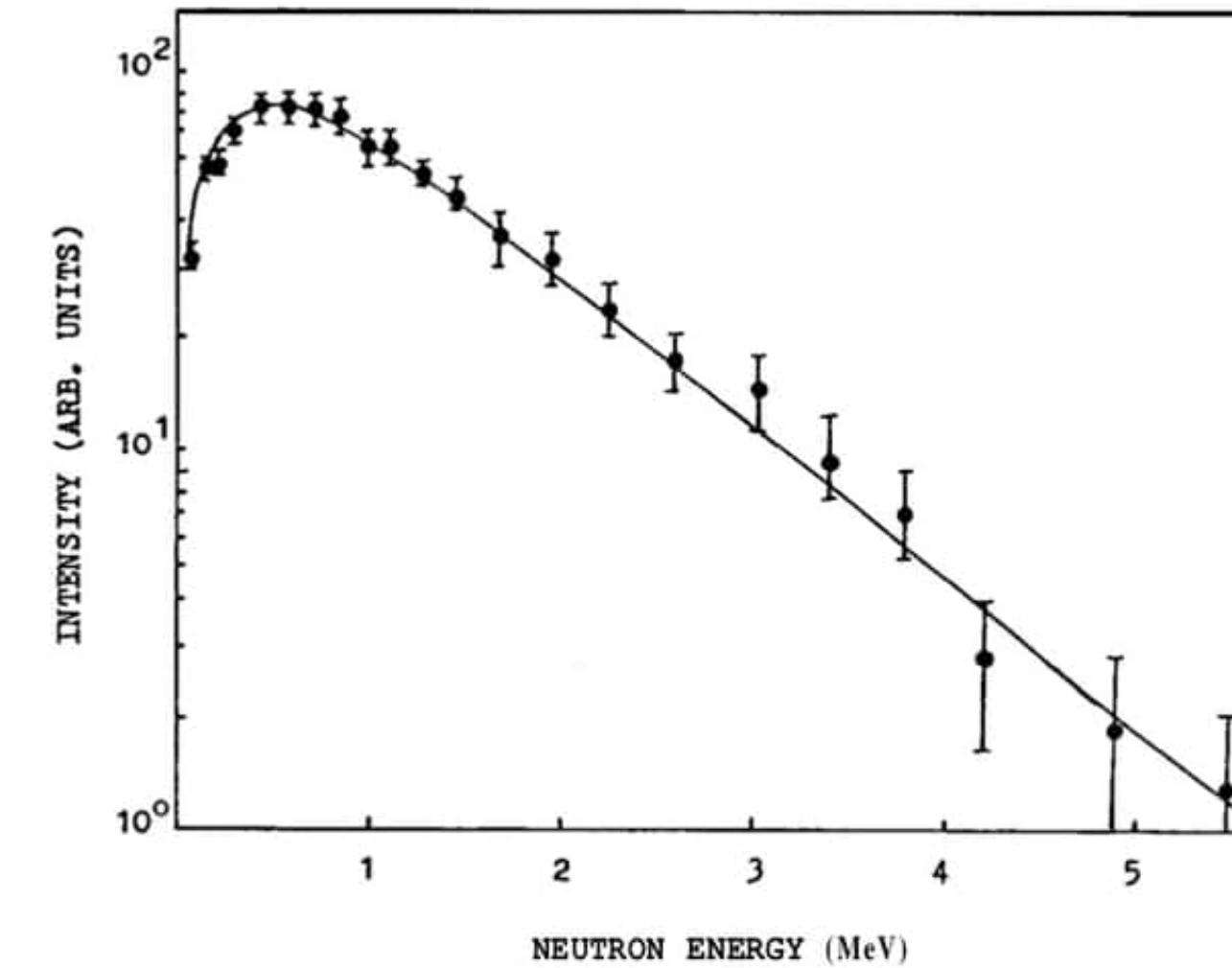
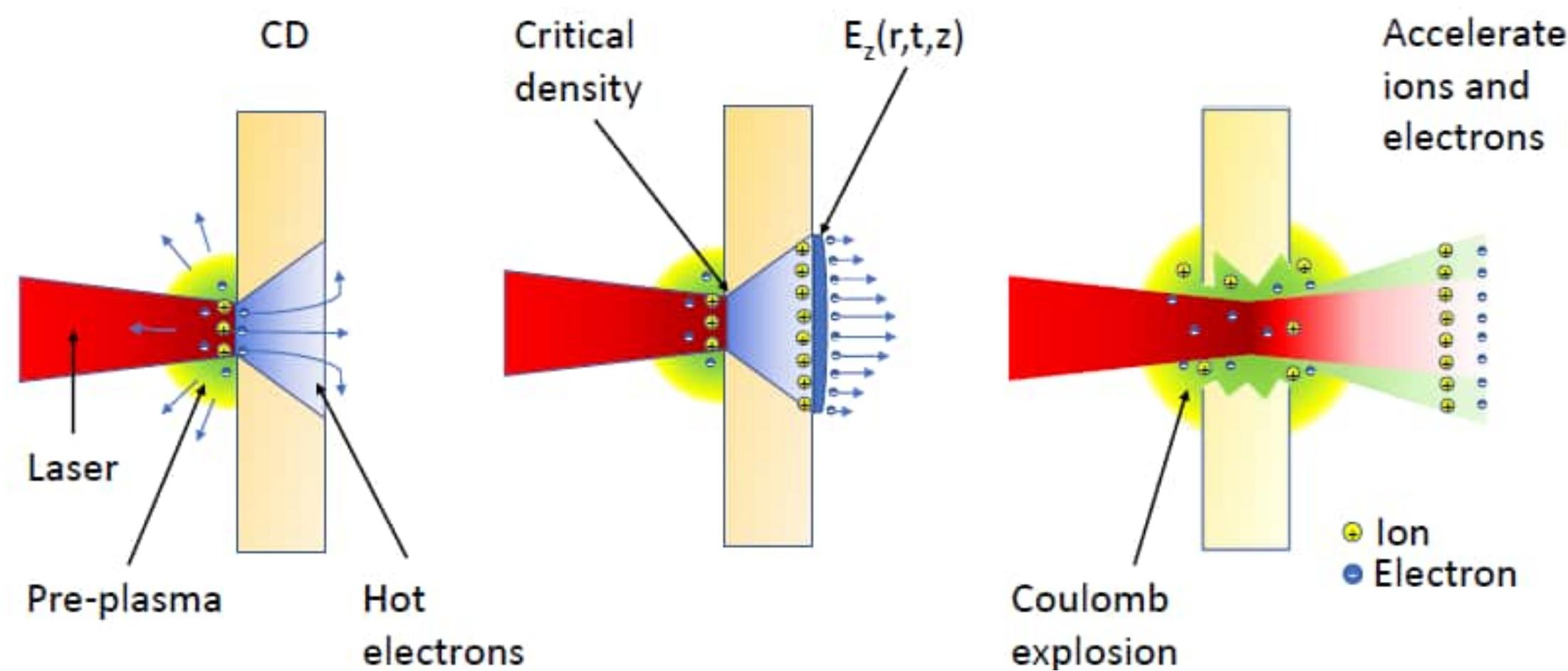


Figure A.4 — Neutron spectrum from a  ${}^{241}\text{Am-Be}(\alpha, n)$  source

Used in metrology for calibration of neutron dose rate meters

# Innovative Neutron Sources: Laser Driven Sources

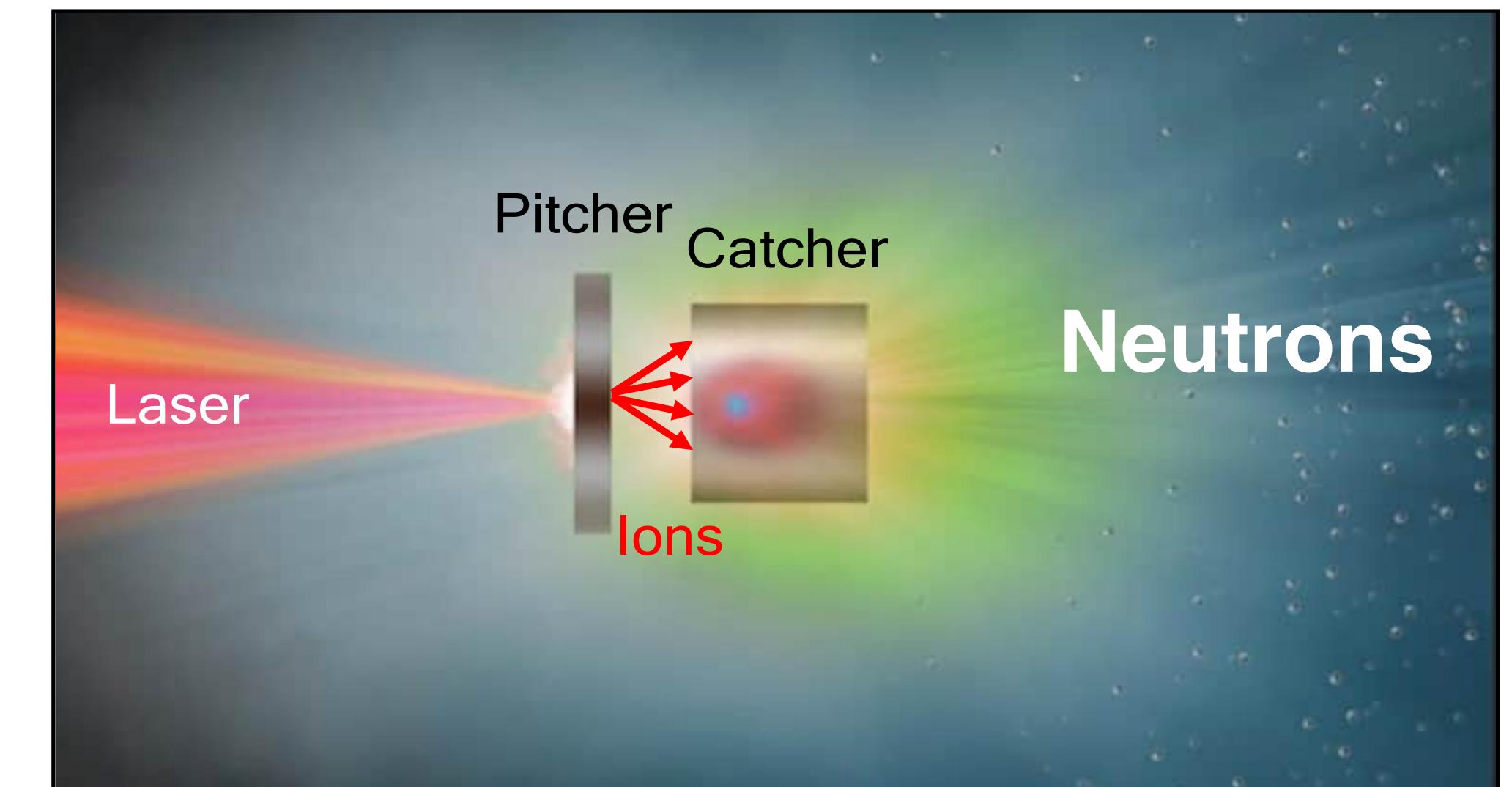
Laser driven neutron sources are undergoing rapid development at big laboratories



Neutron beams with multi-MeV energies can be produced using laser based acceleration mechanisms. Within plasma intense electric fields are generated Two stage process: 1) ions produced in the Pitcher are accelerated and 2) converted into neutrons through nuclear reactions in the Catcher: such as  $^{2d}(d, n)^{3He}$  and  $^{7Li}(p, n)^{7Be}$ . Spallation and laser-driven sources are pulsed owing to the short duration of the driving ion beam,: t is in orders of  $\mu s$  for spallation and ps for laser driven source respectively.

Compared to conventional accelerators, laser based techniques offer the advantage of:

- reduced charged particle acceleration distances, from the meter scale to the millimeter scale due to the high electrical fields supported by plasmas
- high instantaneous neutron production rates, due to the short pulse duration of the laser and the corresponding short acceleration time.



**Proton Accelerator &**

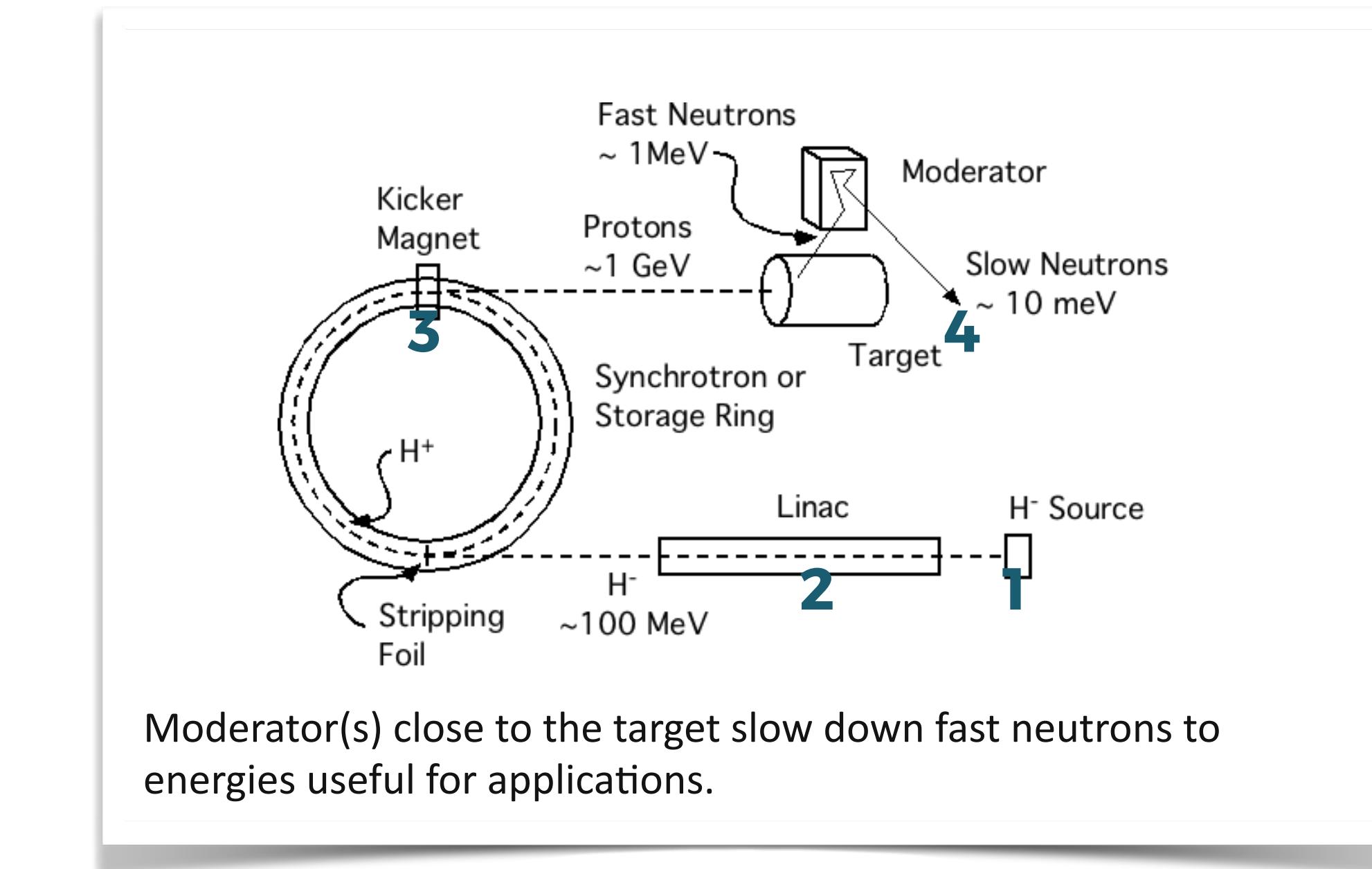
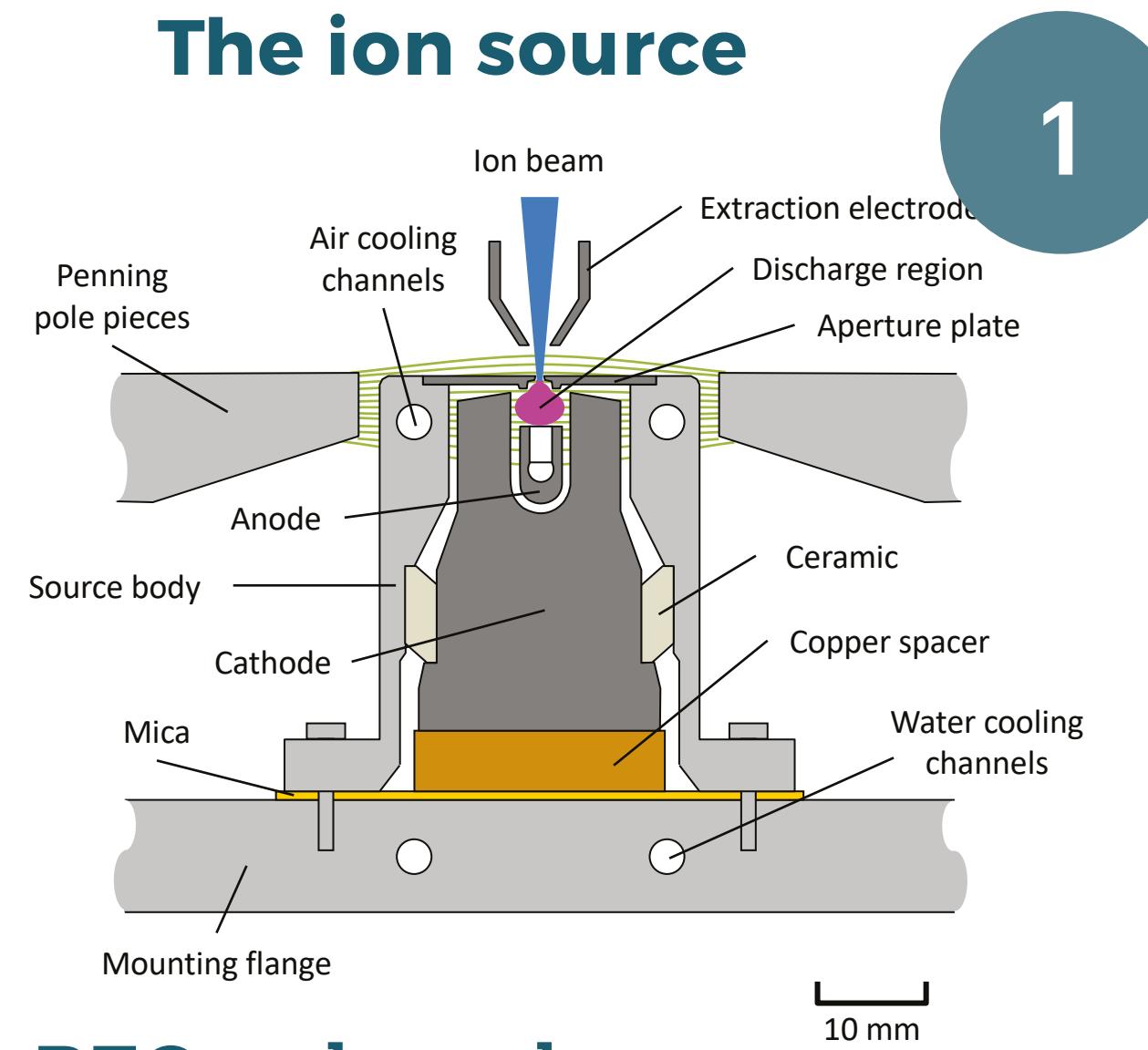
**Design of the TARGET for Spallation Sources**

(Where Primary Protons and Spallation Neutrons Are Produced)

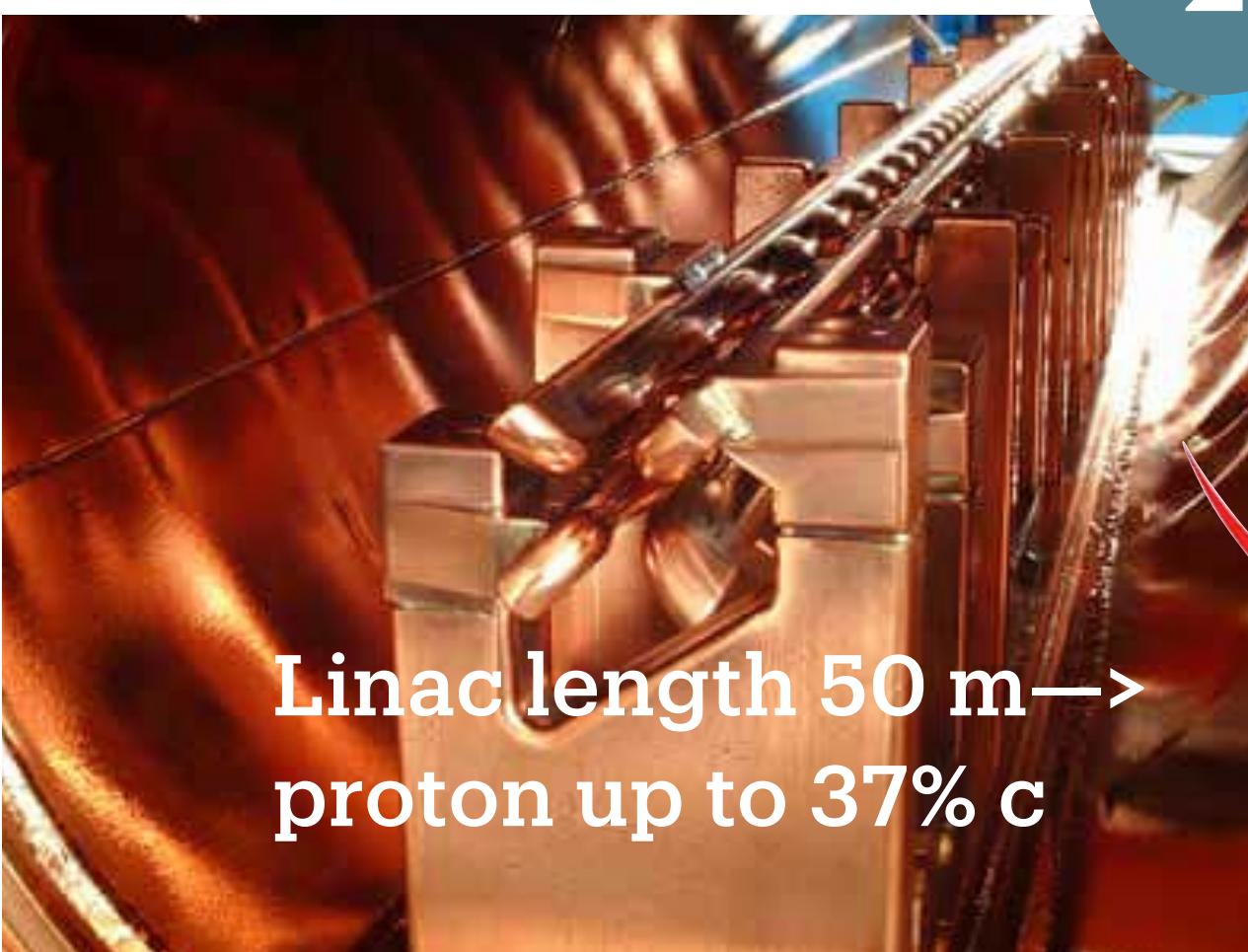
# Anatomy of a Pulsed Spallation Source (ISIS)

Spallation neutron sources can offer the extra advantage of a time structure of the neutron beam that should best match certain experimental techniques (TOF)

## The ion source



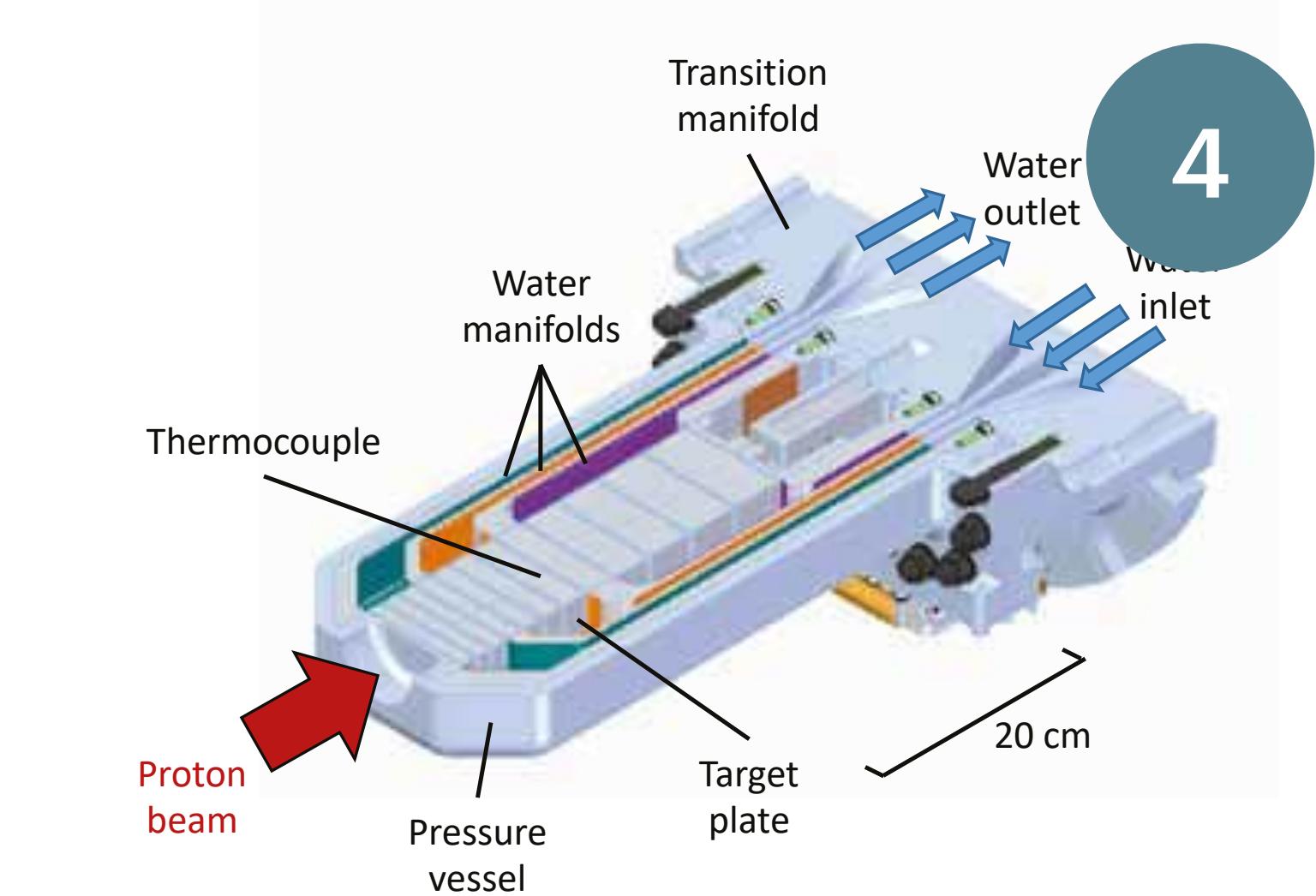
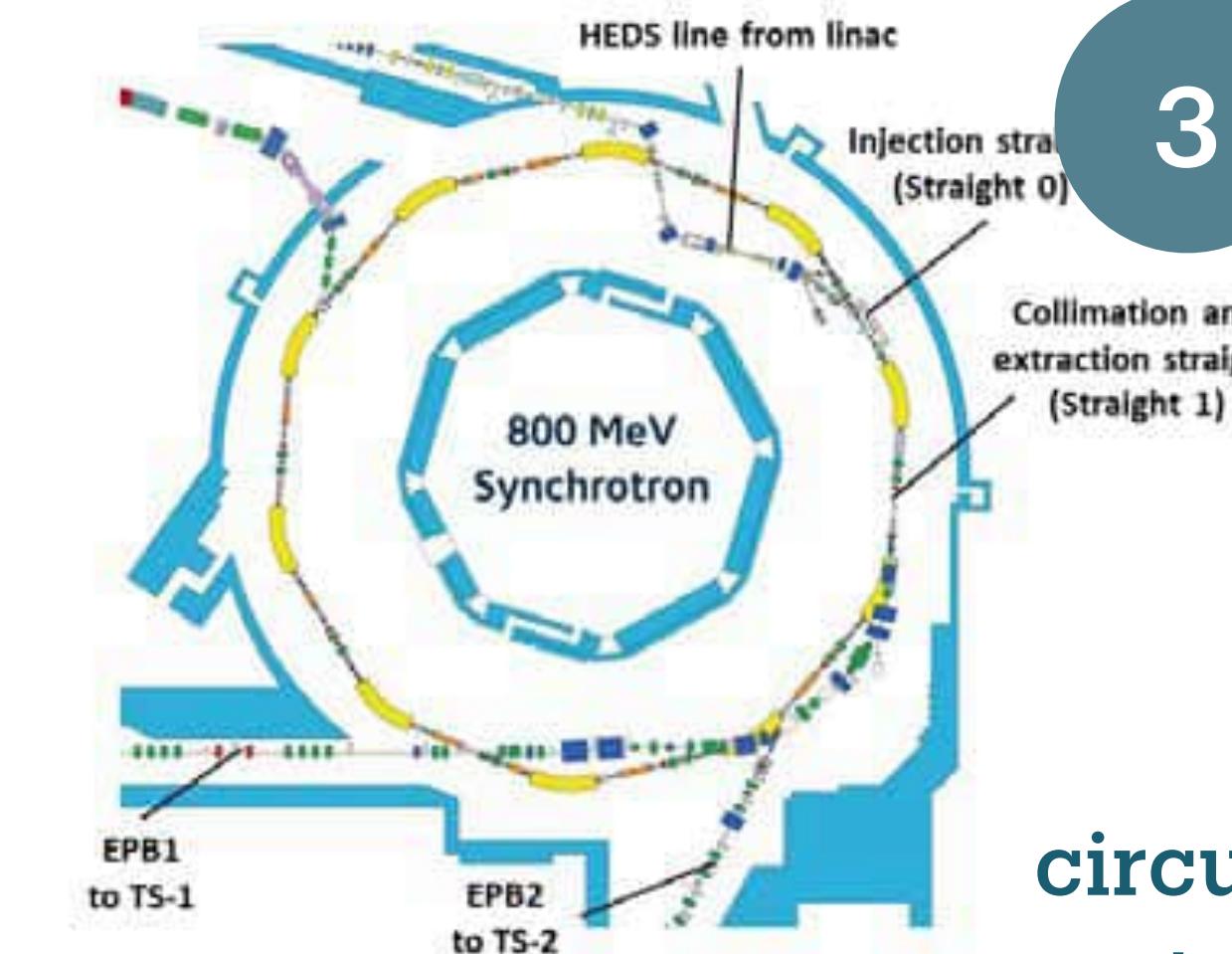
## RF Quadrupole



## The RF LINAC

Linac length 50 m—>  
proton up to 37% c

## The Syncrotron

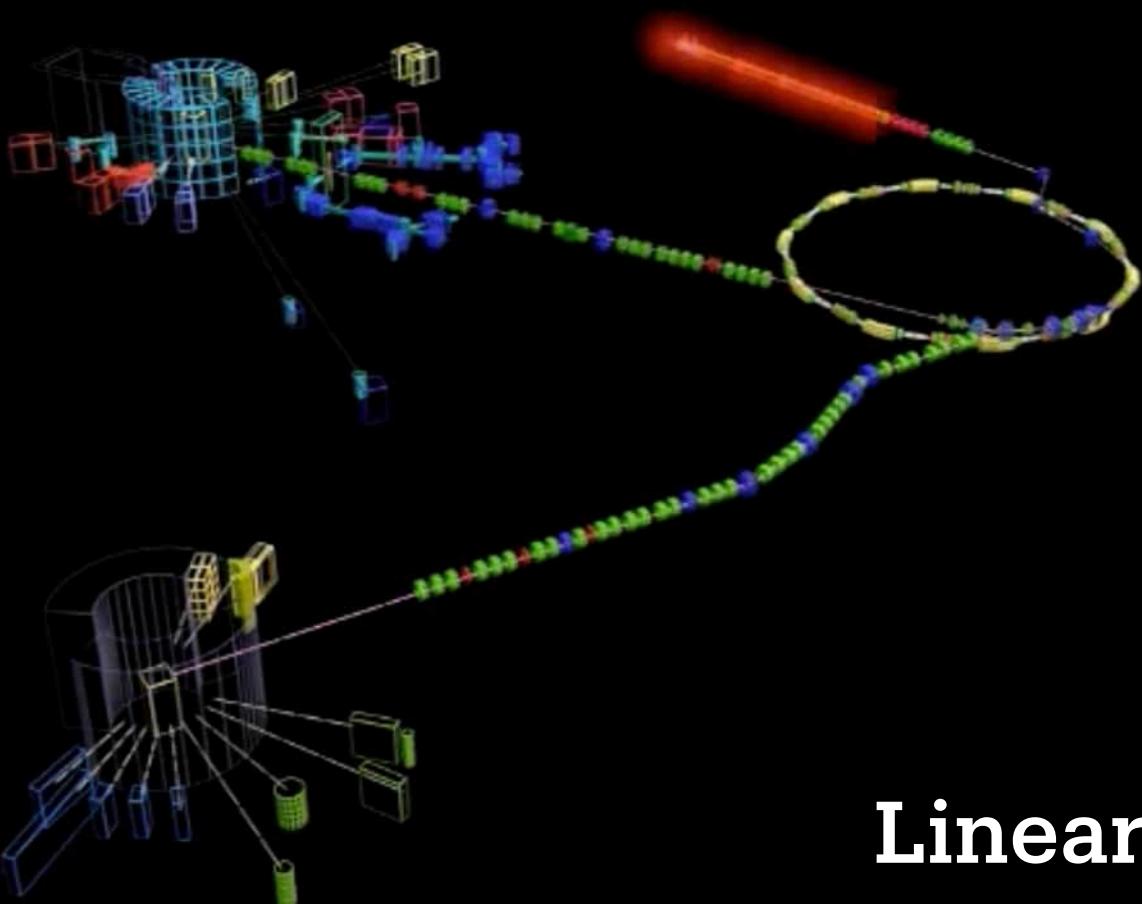


Schematic diagram of TS-1  
neutron-producing target.

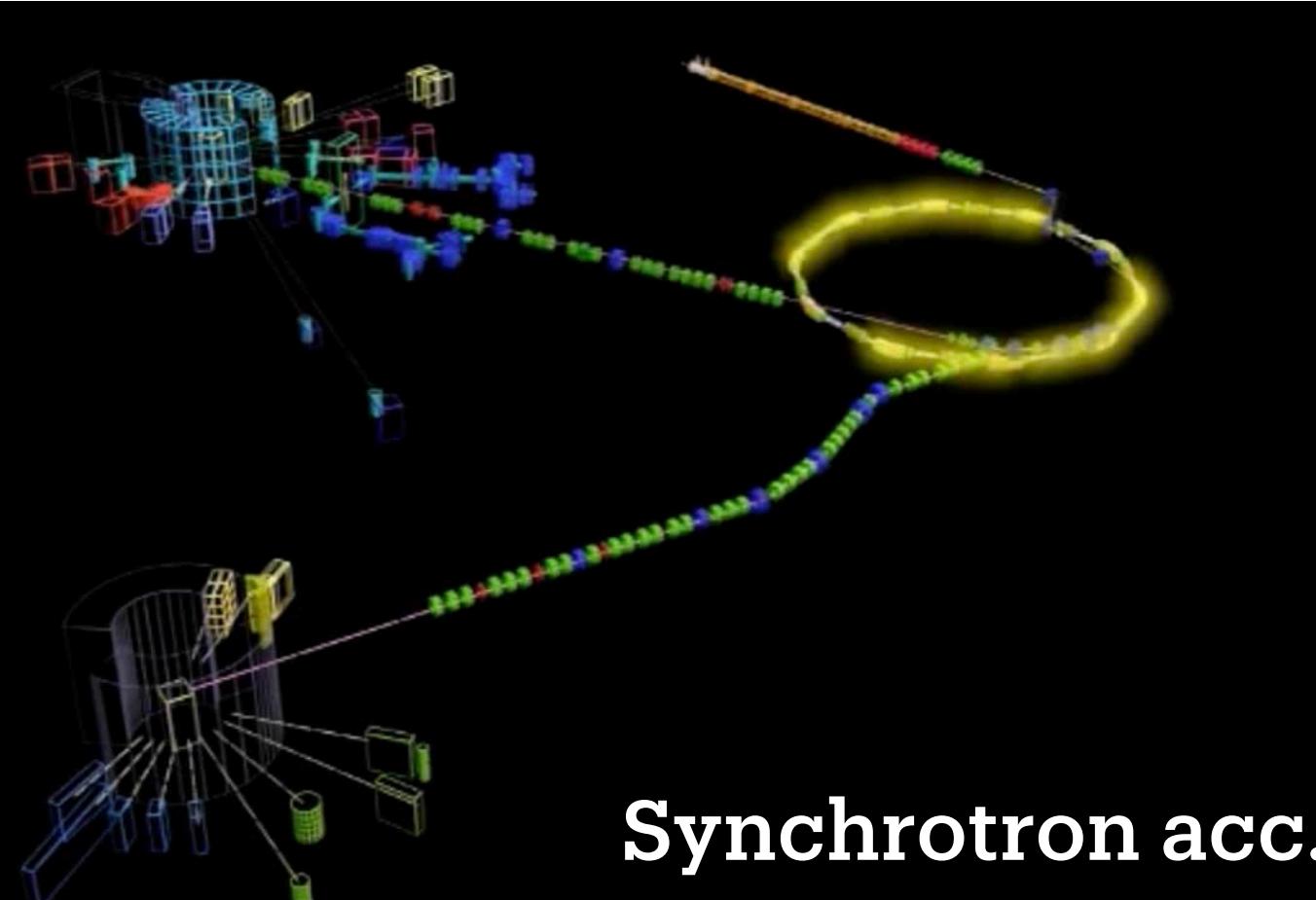
## Kicker magnet

circular orbit 163 m—>  
up tp 84% c

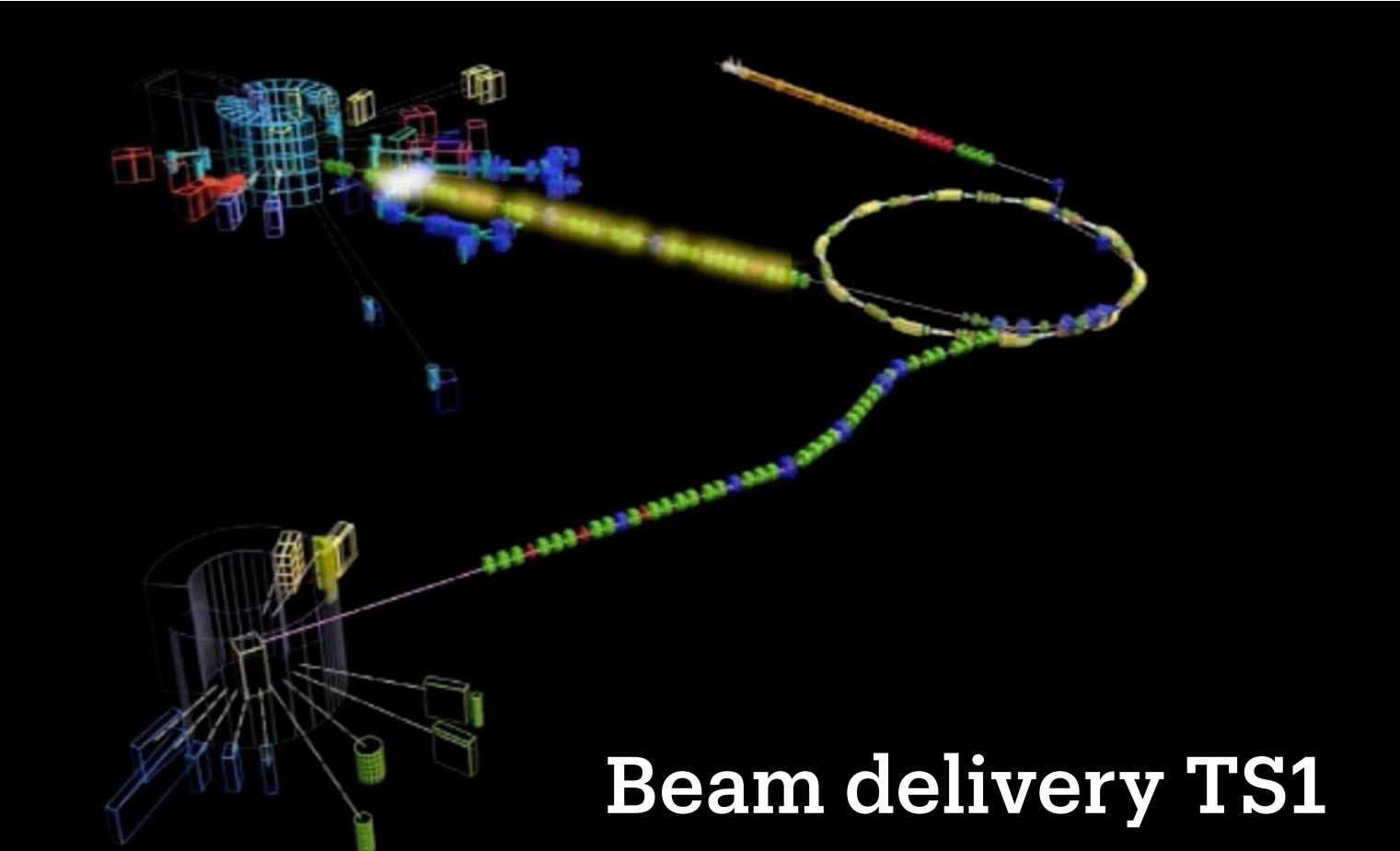
# BEAM Time Lapse in ISIS TS1



Linear acc.



Synchrotron acc.



Beam delivery TS1

ISIS ion source creates negatively charged hydrogen ions made up of two electrons and one proton. These are fed into two linear accelerator sections through which they are focused and accelerated up to 36% the speed of light

The third final accelerator is a circular synchrotron.

At the entrance the ions are stripped of their electrons by a thin alumina foil leaving the bare protons

The synchrotron is made up of ten sections each consisting of a bending magnet to keep the protons on their circular path and five focusing magnets

The beam is separated in 2 bunches

After 12000 revolutions of the synchrotron the protons are traveling at 84% of the speed of light at this speed they could travel six times around the earth in just one second. Through a kick magnet they are delivered to the TS1 or TS2 station

Every 5 pulses TS1 1 is delivered to TS2

# Design of Target for Spallation Sources

- The essential function of the targets in spallation sources is to convert the high-energy proton beam into as many neutrons as possible, whilst occupying as small a volume as possible.
- Minimising the volume within which protons are converted to neutrons results in the highest neutron fluxes.
- Within the target, most of the power in the proton beam appears as heat, which has to be carried away by cooling .
- For neutron production to occur, nuclear collisions must take place before the incident particle reaches the end of its range.

## Choosing the right thickness for a given $E_p$ :

$\lambda$  is almost constant above 100 MeV and approximately expressed by

$$\lambda = 33A^{1/3}, \text{ giving about } 200 \text{ g cm}^{-2} \text{ for heavy nuclei} \rightarrow \lambda = 10 \text{ cm}$$

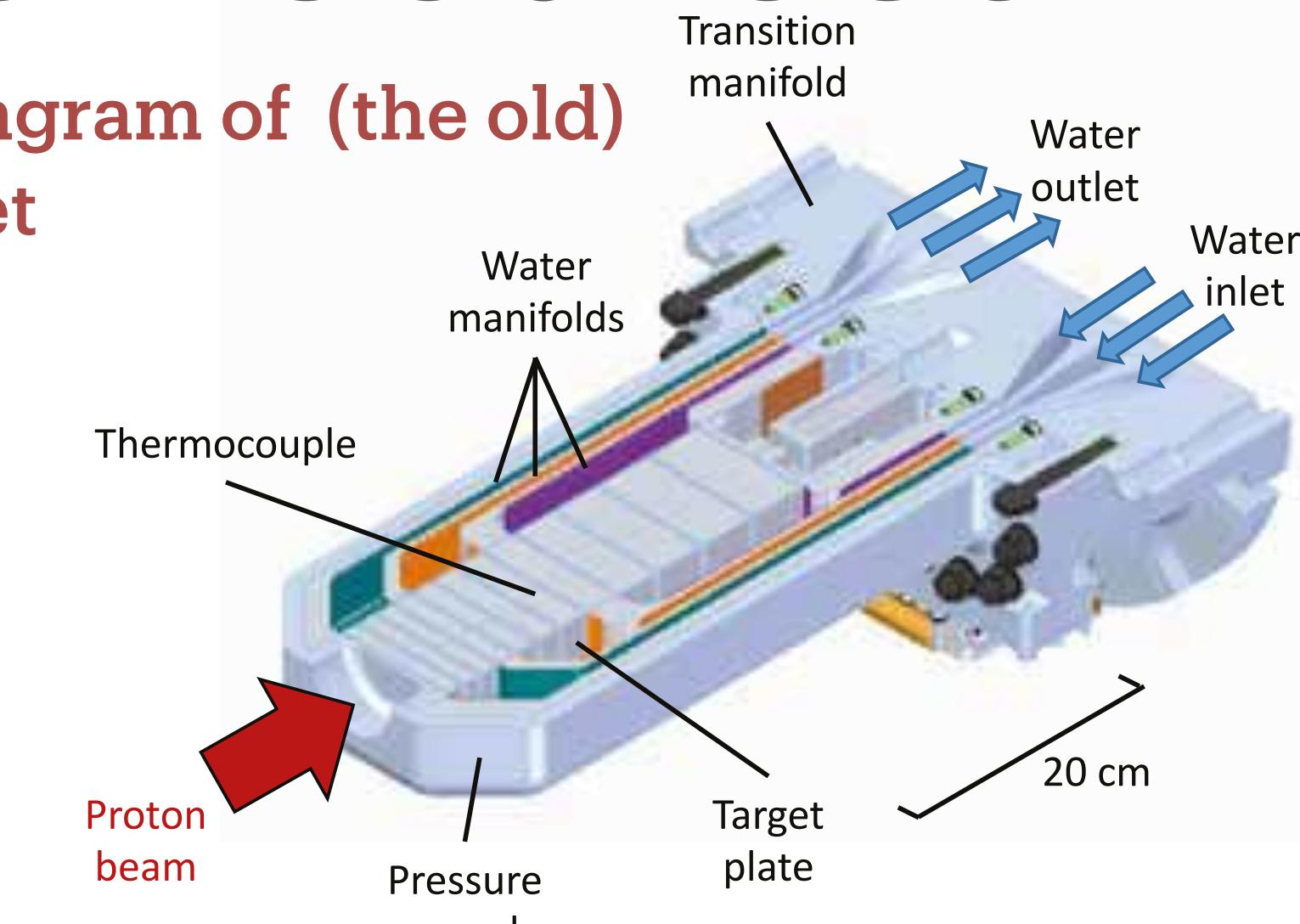
The probability for the nuclear collision,  $P_n$ :

$$P_n = 1 - \exp(-R/\lambda), \text{ where } R \text{ is the particle range}$$

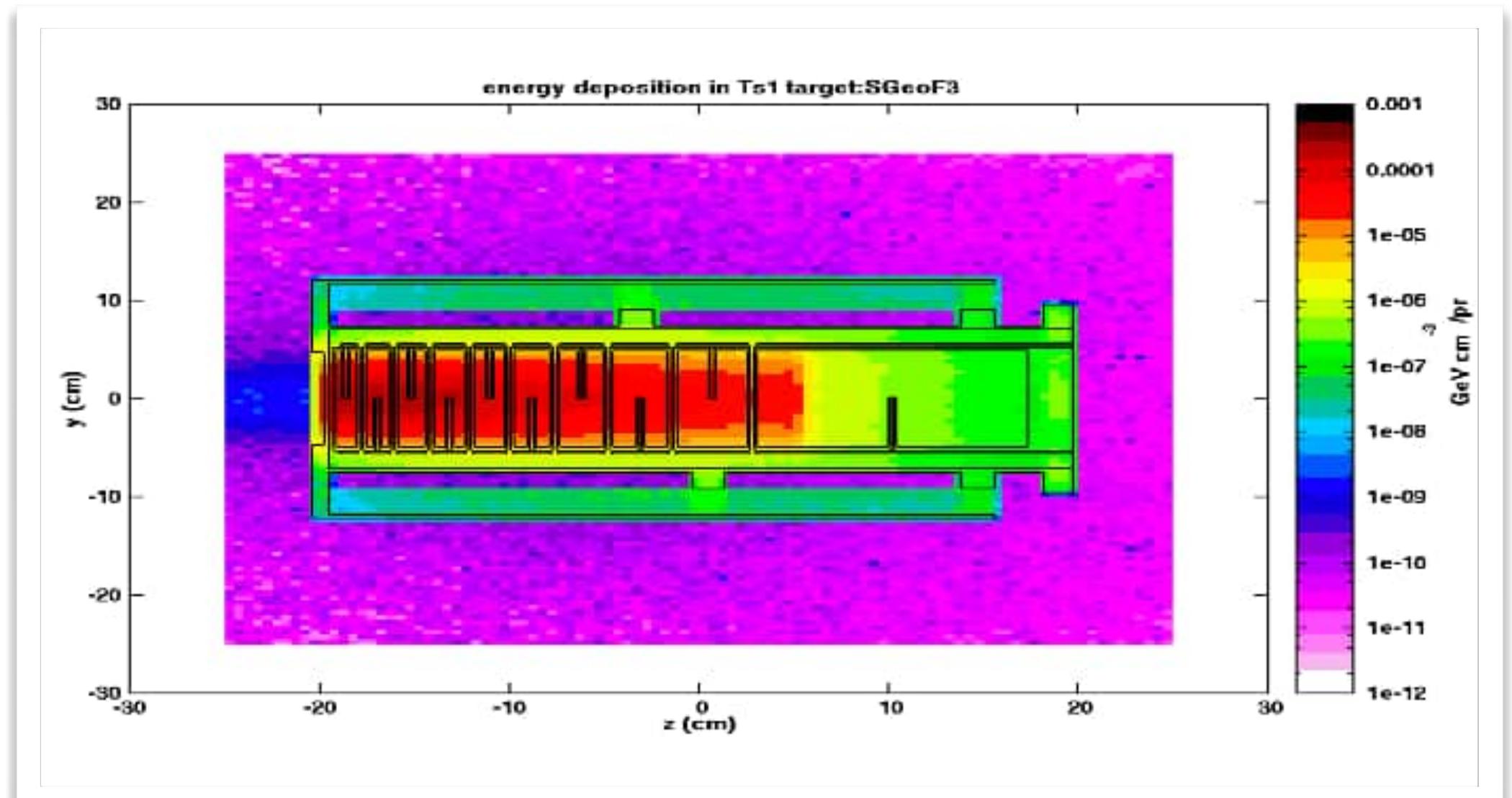
To get  $P_n > 90\% \rightarrow R(E) > 2.5 \lambda \rightarrow R > 25 \text{ cm}$   
proton of 800 MeV in W  $\rightarrow R(E_p) = 25 \text{ cm}$

A target thickness larger than 25 cm satisfies the above condition  $\rightarrow$   
The ISIS target  $L = 40 \text{ cm}$  (sum of plate thickness = 30 cm)

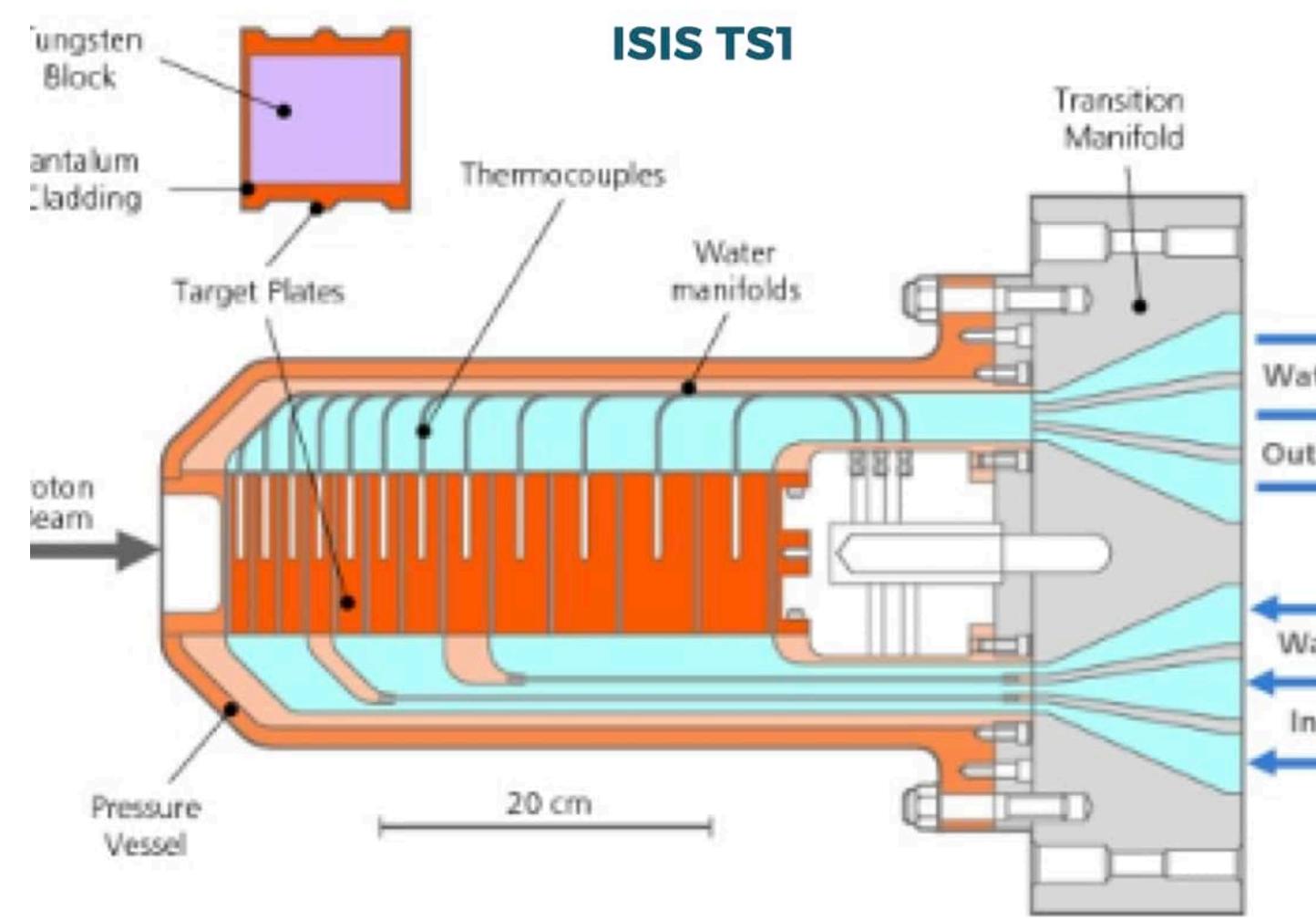
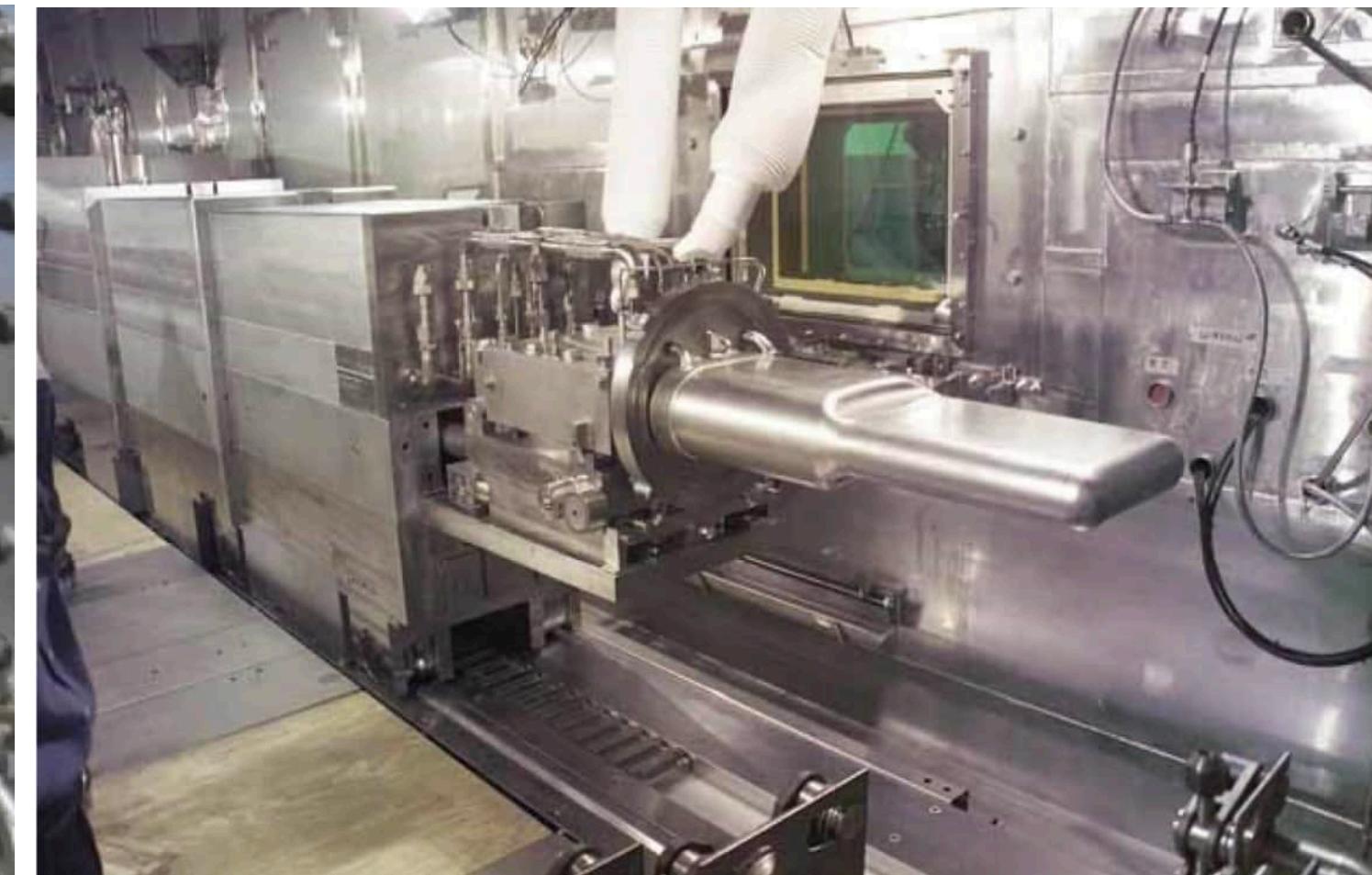
Schematic diagram of (the old)  
ISIS-TS1 target



Results of MC Simulation of Energy deposited in the target



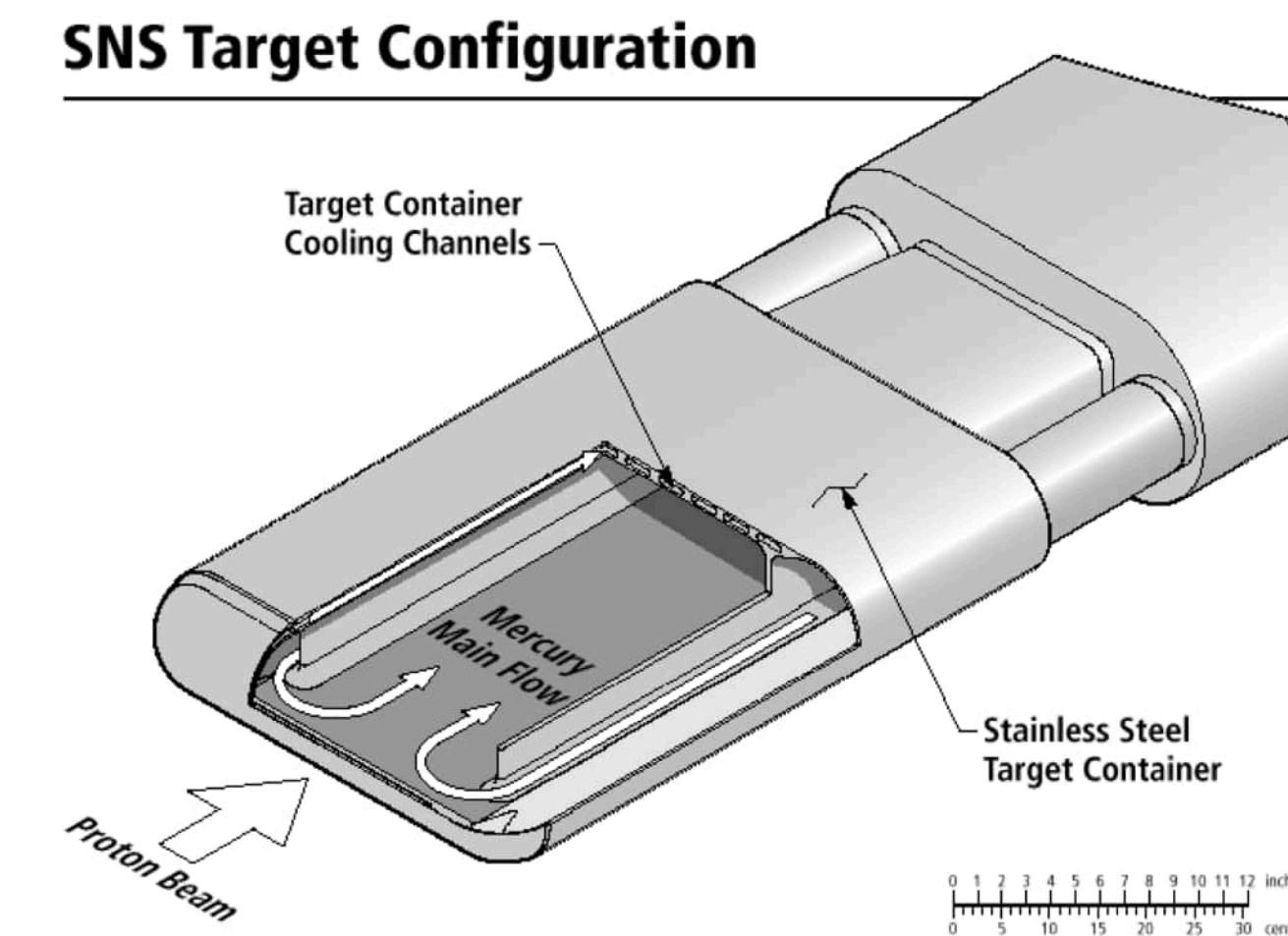
# Spallation Targets as a Function of Beam Power



160 kW in TS1

ISIS

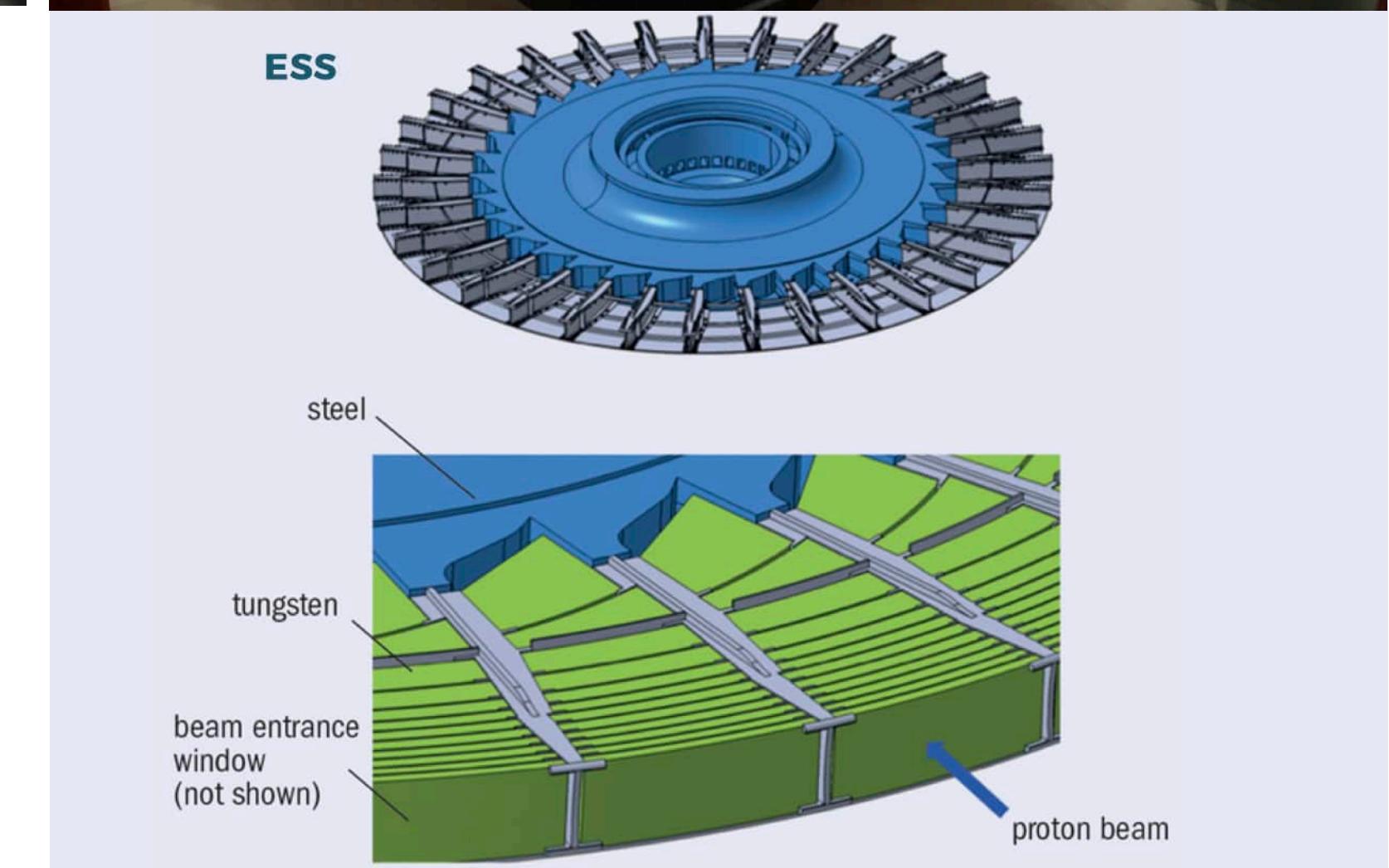
Segmented Target



1.4 MW

SNS

Liquid Circulating Target

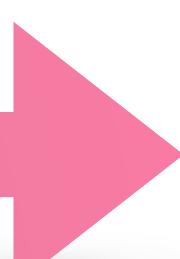


5 MW

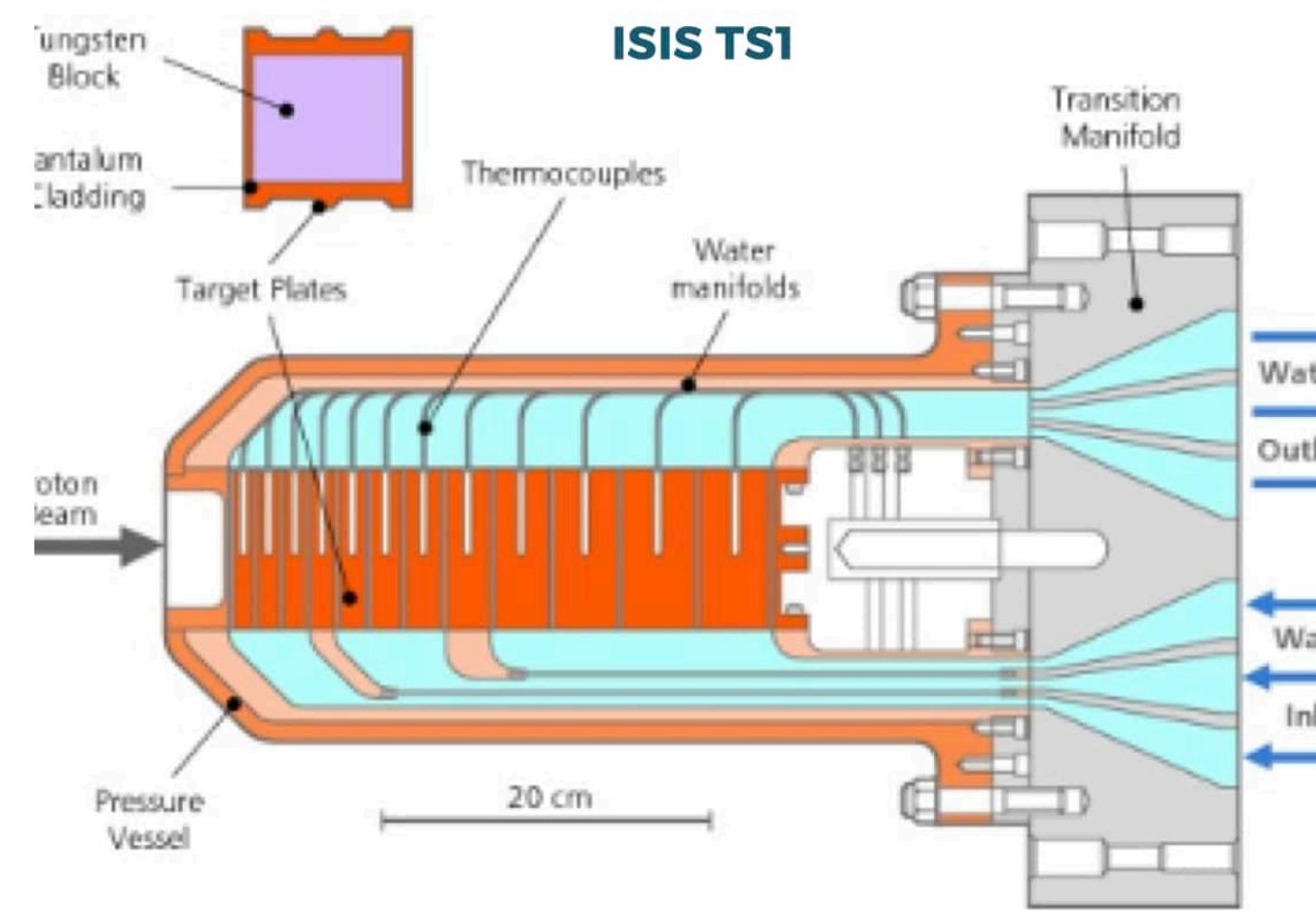
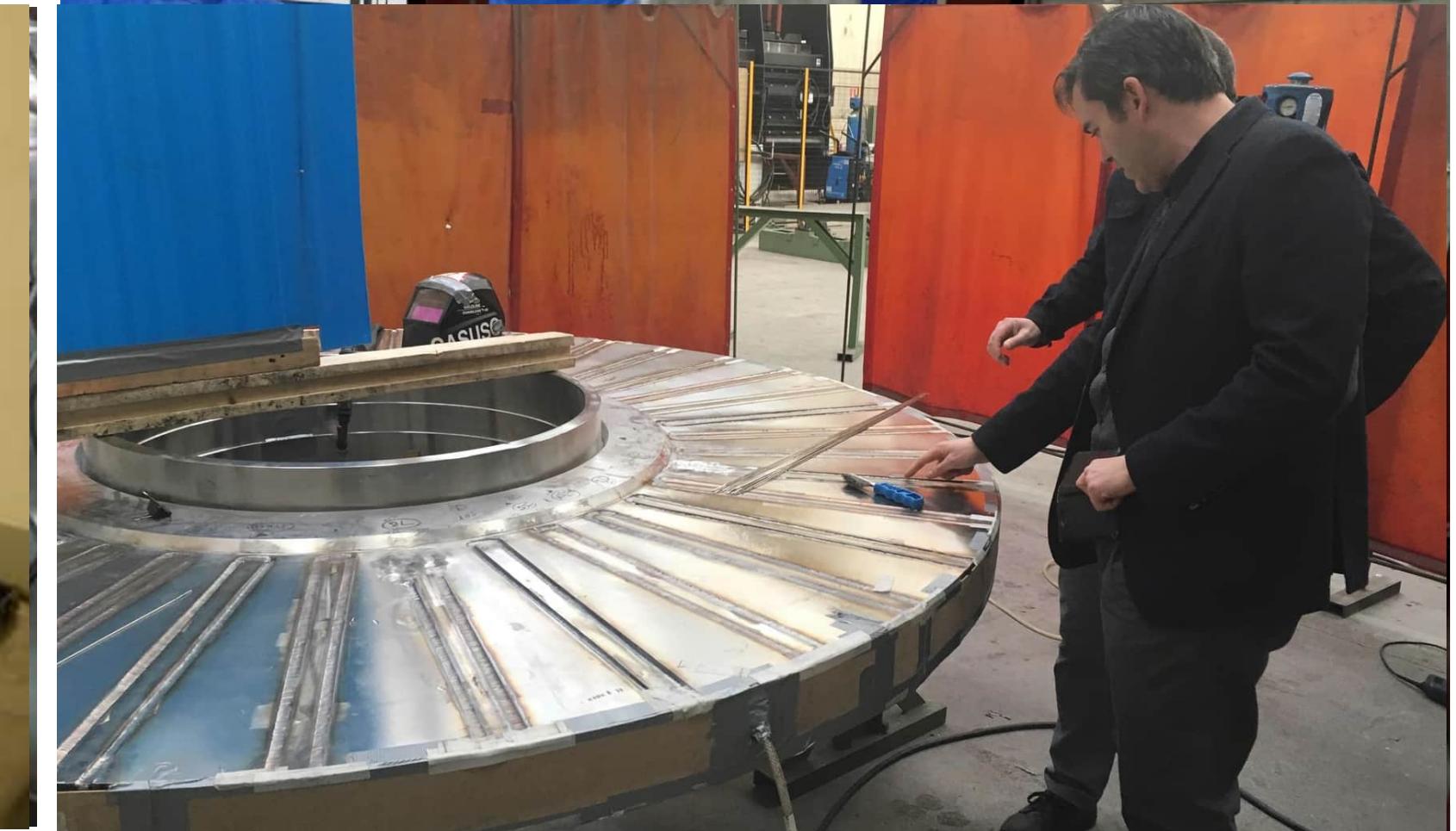
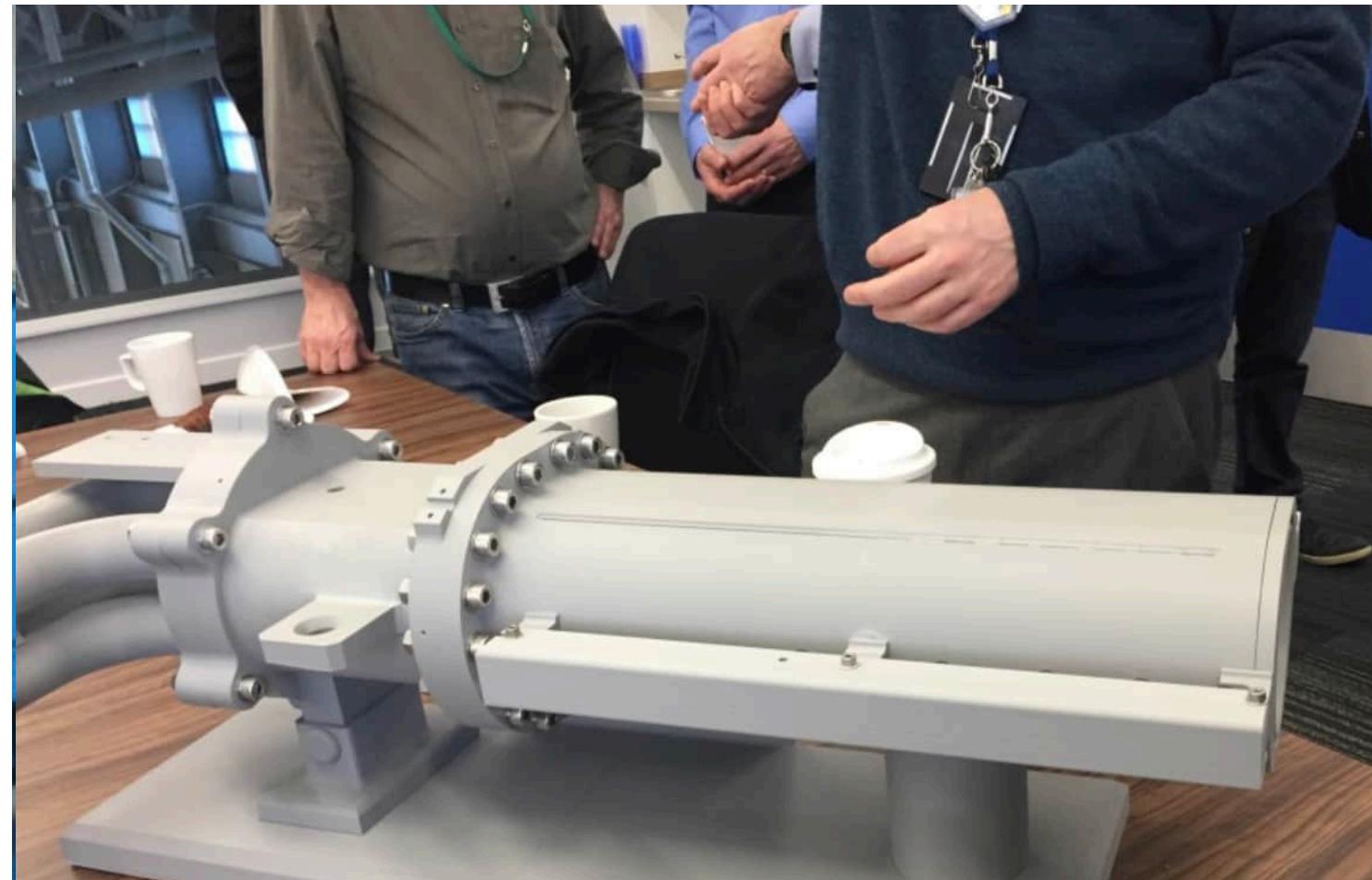
ESS

Rotating Target

Beam Power



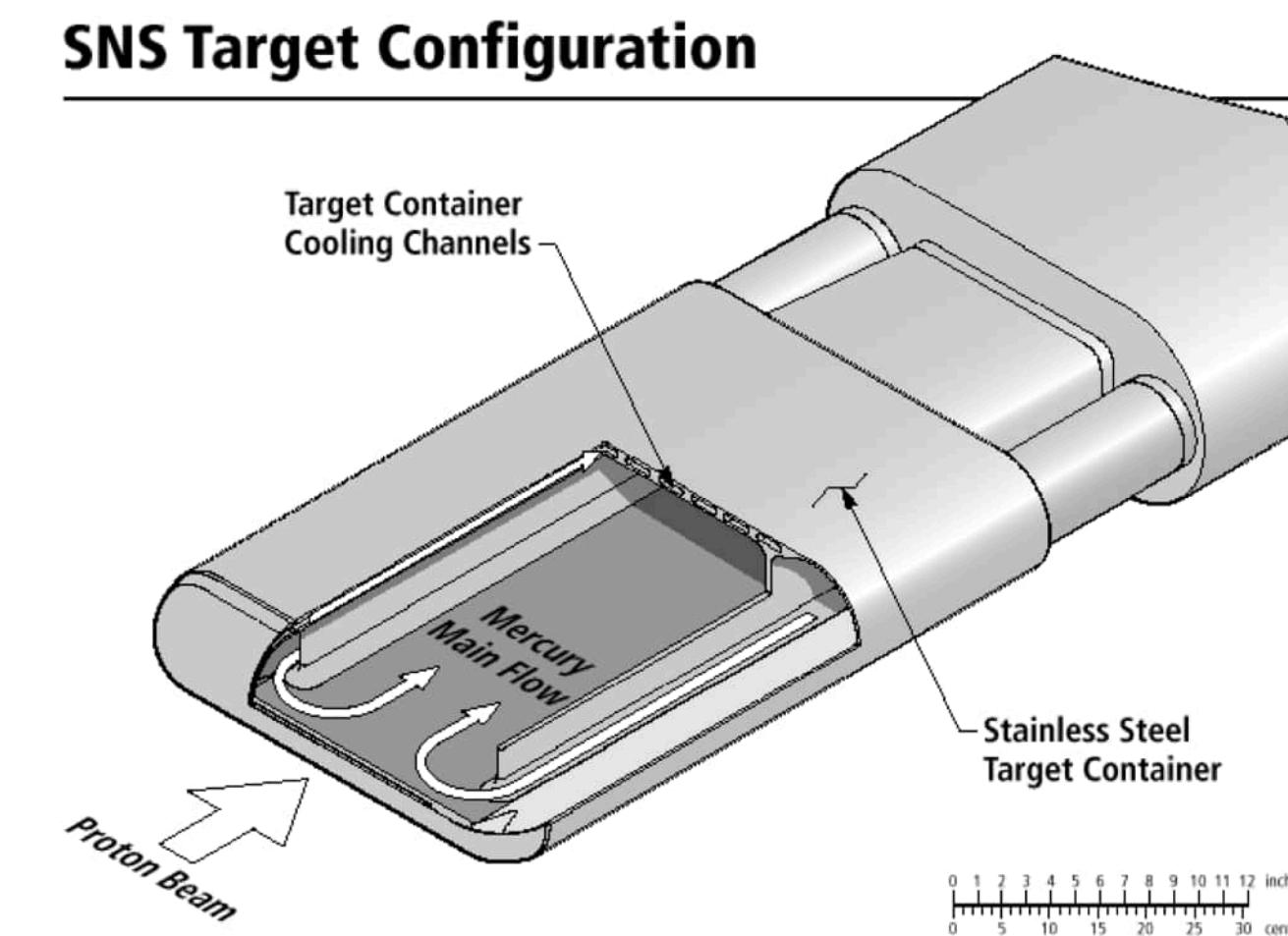
# Spallation Targets as a Function of Beam Power



160 kW in TS1

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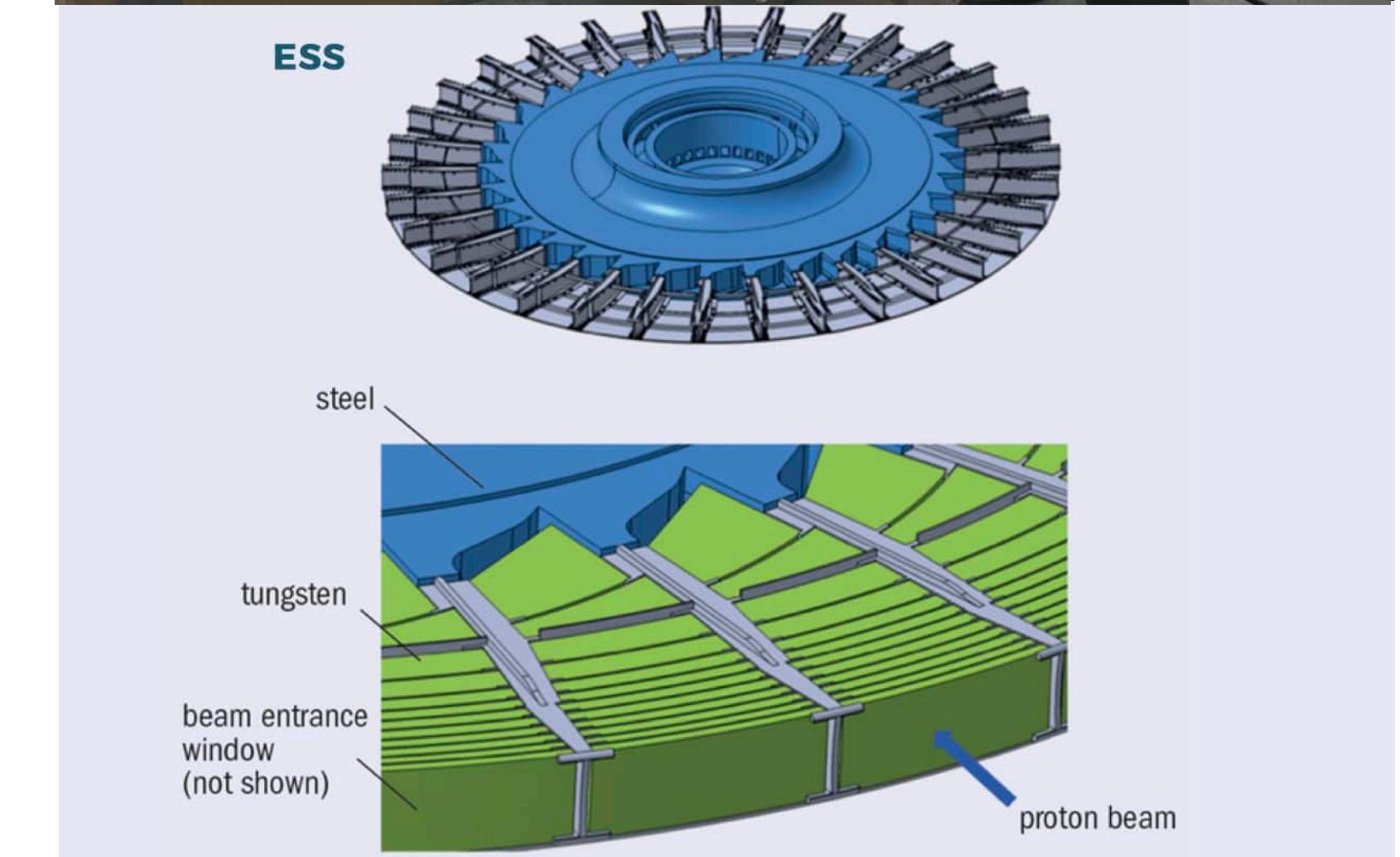
Segmented Target



1.4 MW

SNS

Liquid Circulating Target

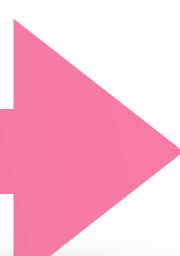


5 MW

ESS

Rotating Target

Beam Power

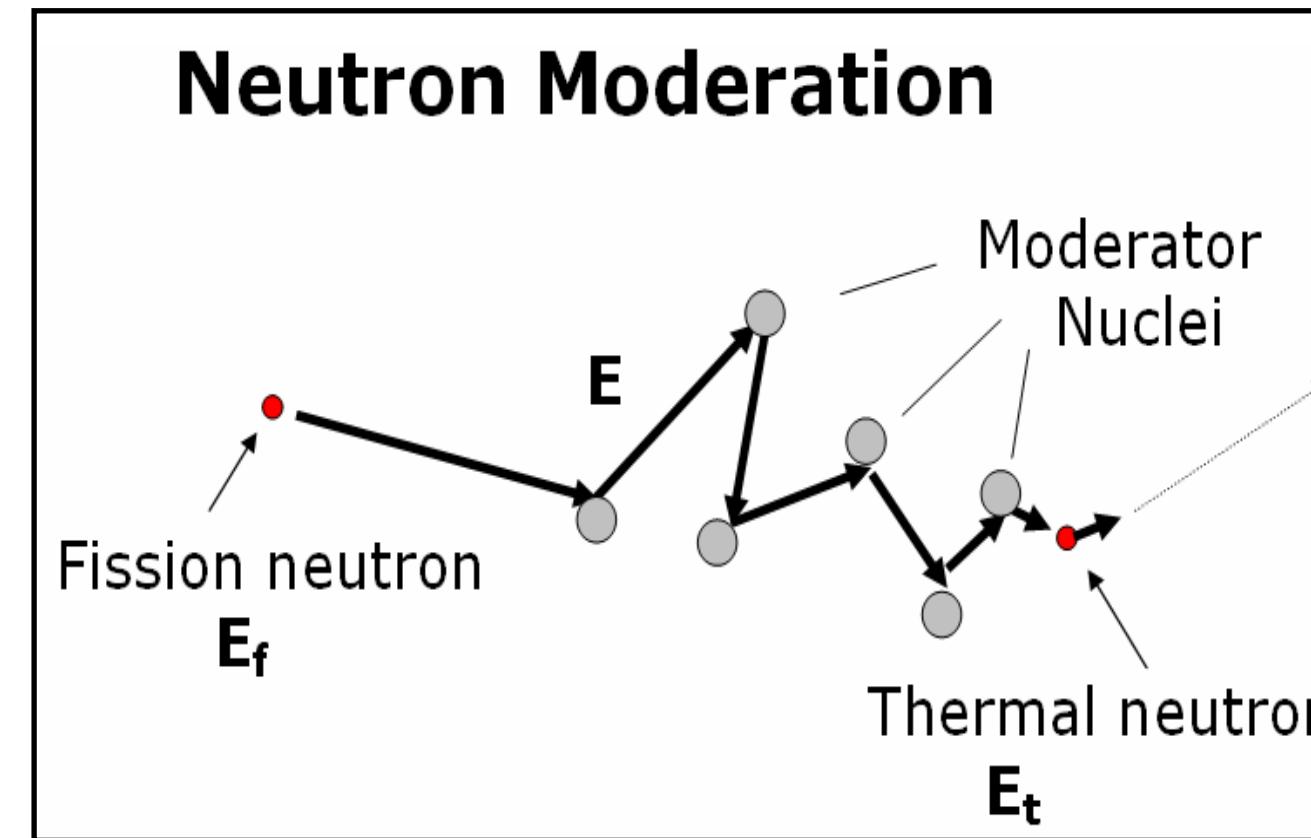
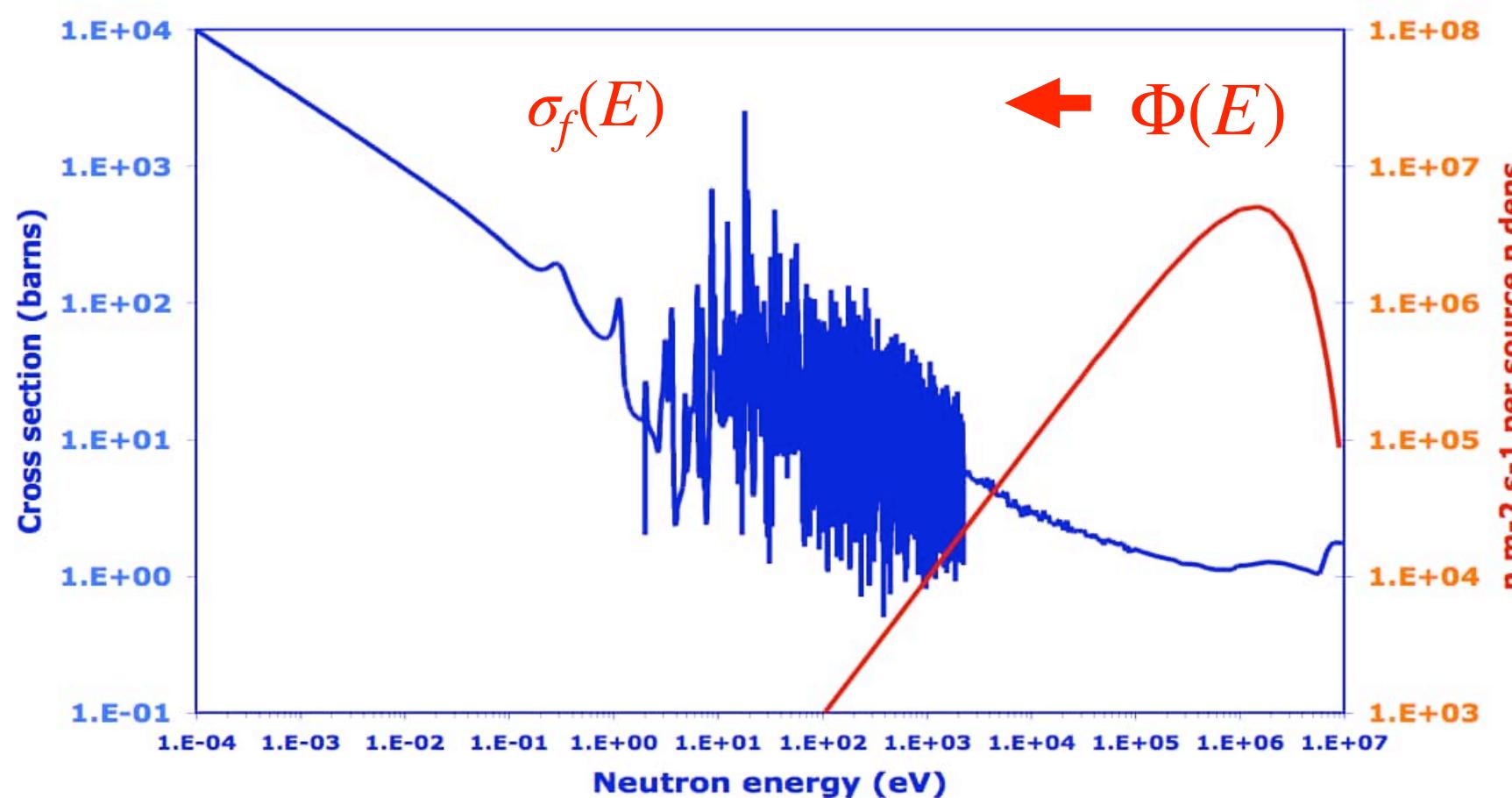


# **Design of a Nuclear Research Reactor Core**

**(Where Fission Neutrons Are Produced)**

# The Chain Reaction and the Critical Mass

Reactors are device used to initiate and control a chain of nuclear fission reaction.

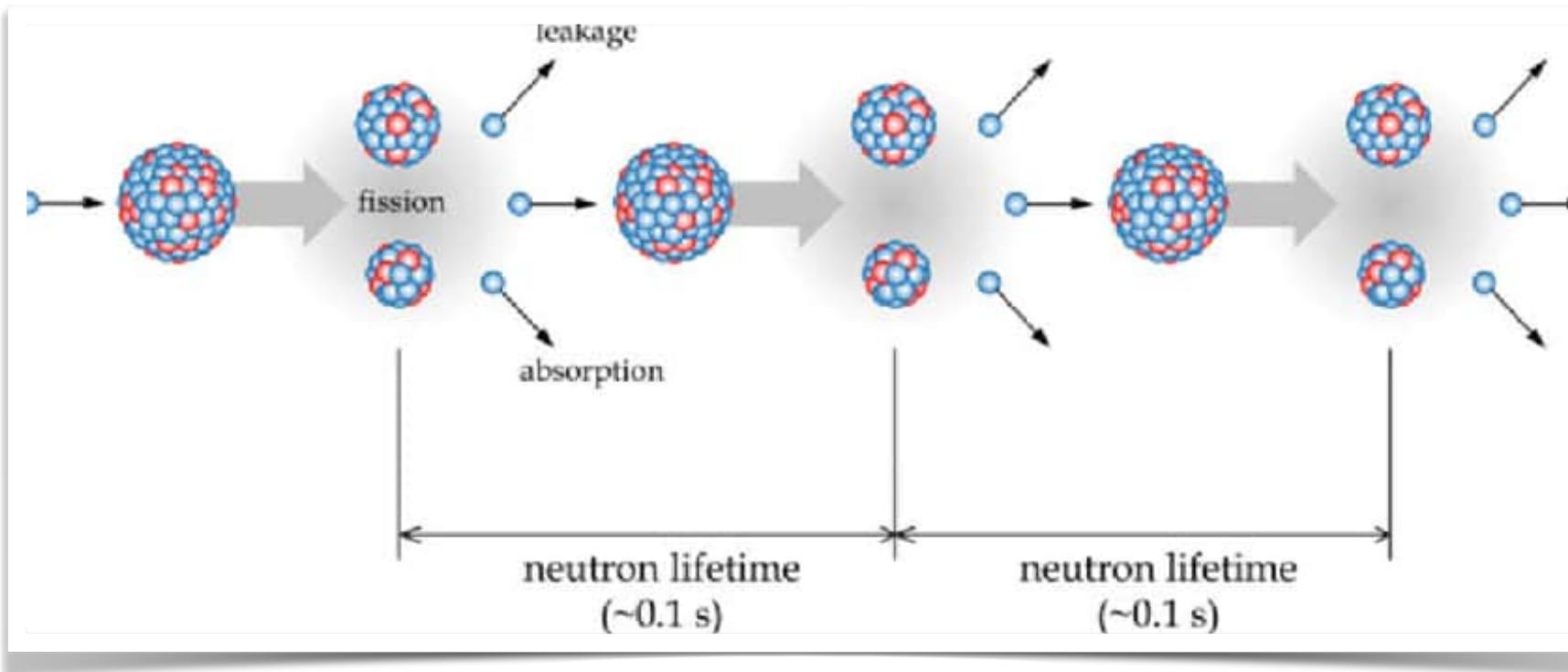


**Critical reactor**

$$k = \frac{\text{number of fissions in one generation}}{\text{number of fissions in preceding generation}} = 1$$

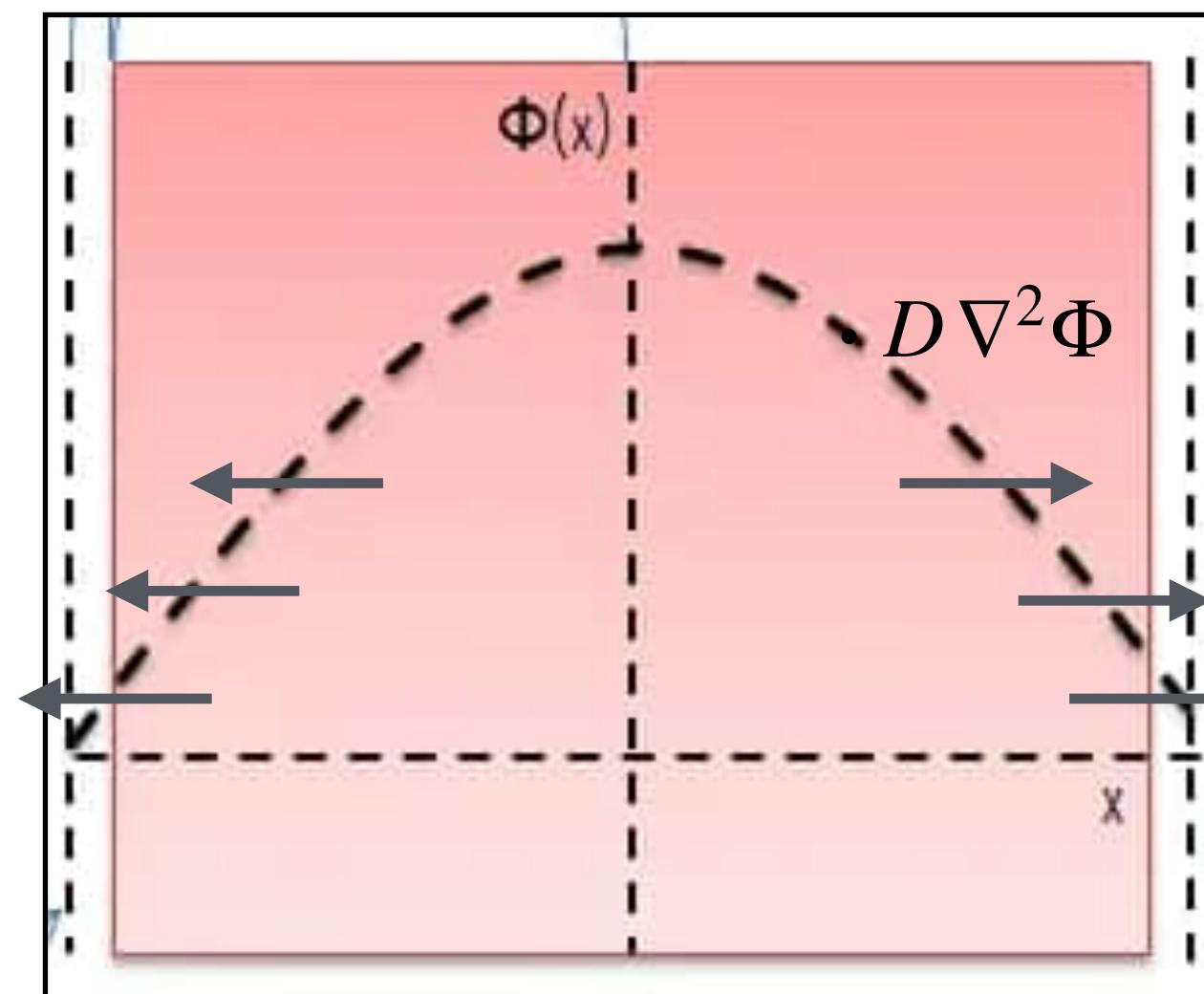
- Neutron population is not increasing neither decreasing → then reactor is called critical
- When production is less than losses due to leakage and absorption,  $k > 1 \rightarrow$  supercritical reactor
- if  $k < 1$  subcritical

Fission spectrum and  $^{235}\text{U}(n,f)$  cross section



Scheme of a self-sustaining fission chain

Criticality means sustaining fission chain w/o external source



$$RR = N\sigma\Phi = \Sigma\Phi$$

$$D \nabla^2 \Phi - \Sigma_a \Phi + S = 0$$

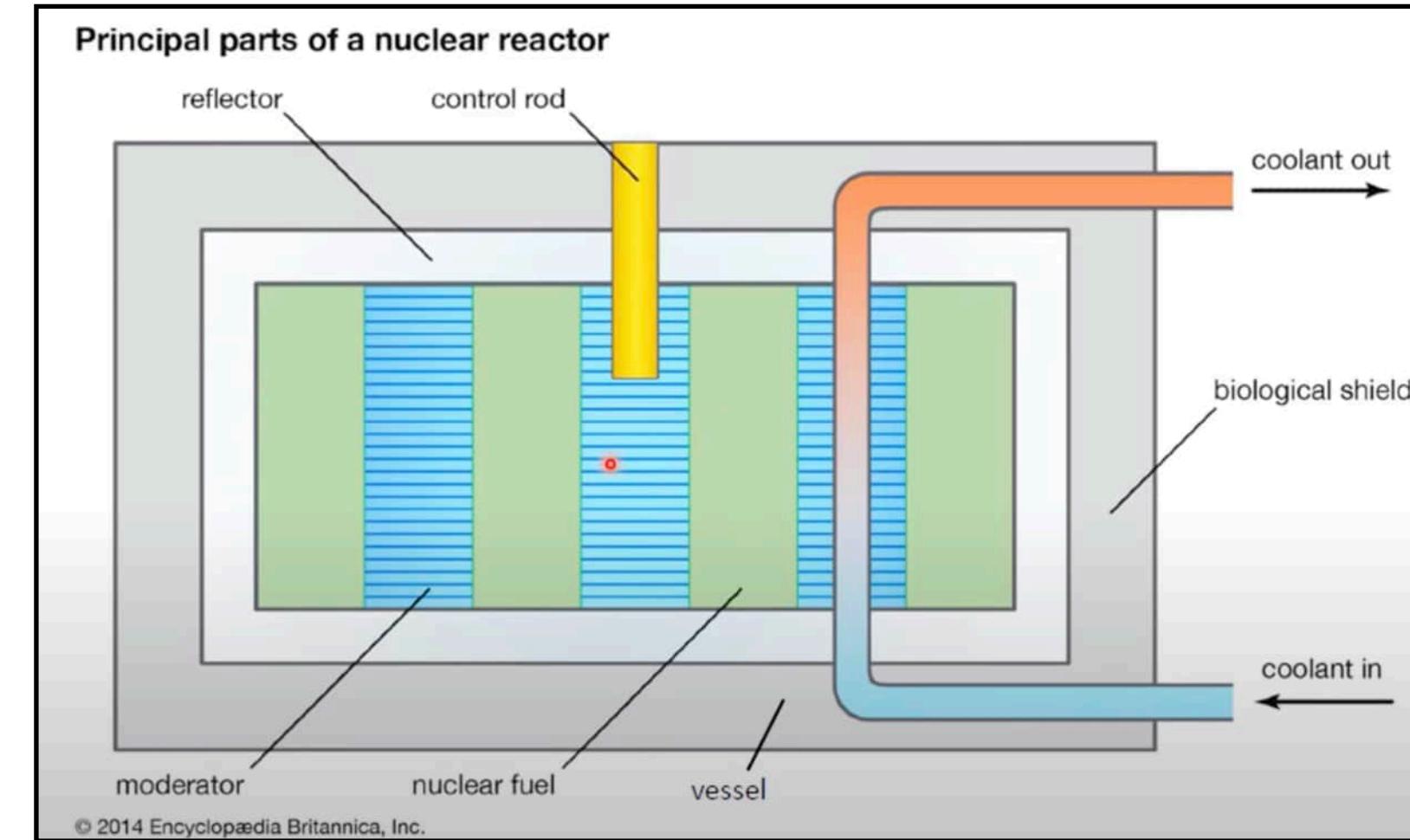
Leakage      Absorption capture      Source (fission)

$$S = \nu N_f \sigma_f \Phi \quad N_f = N_{235} \text{ per unit volume}$$

N.B.—>The spallation process, in contrast to fission, is not an exothermal process: energetic particles are required to drive it

# Reactor Core: NPP vs RR

Main components of a nuclear reactor core: fuel, moderator control rod, containment vessel



Parameter	NPP (isar2)	RR (FRM II)
Thermal power	3950 MW	20 MW
Electric power	1485 MW	0
Diameter	3.37 m	24 cm
Active length	3.66 m	70 cm
Enrichment	<5%	83%
N.fuel elem.	$^{193}$ (ca. 103 t Uranium)	1 (8,1 kg Uranium)
Coolant temperature	293-328 °C	39-52 °C
Coolant pressure	155 bar	Open Pool

## NPP

Energy produced from fission reaction is used to generate electricity

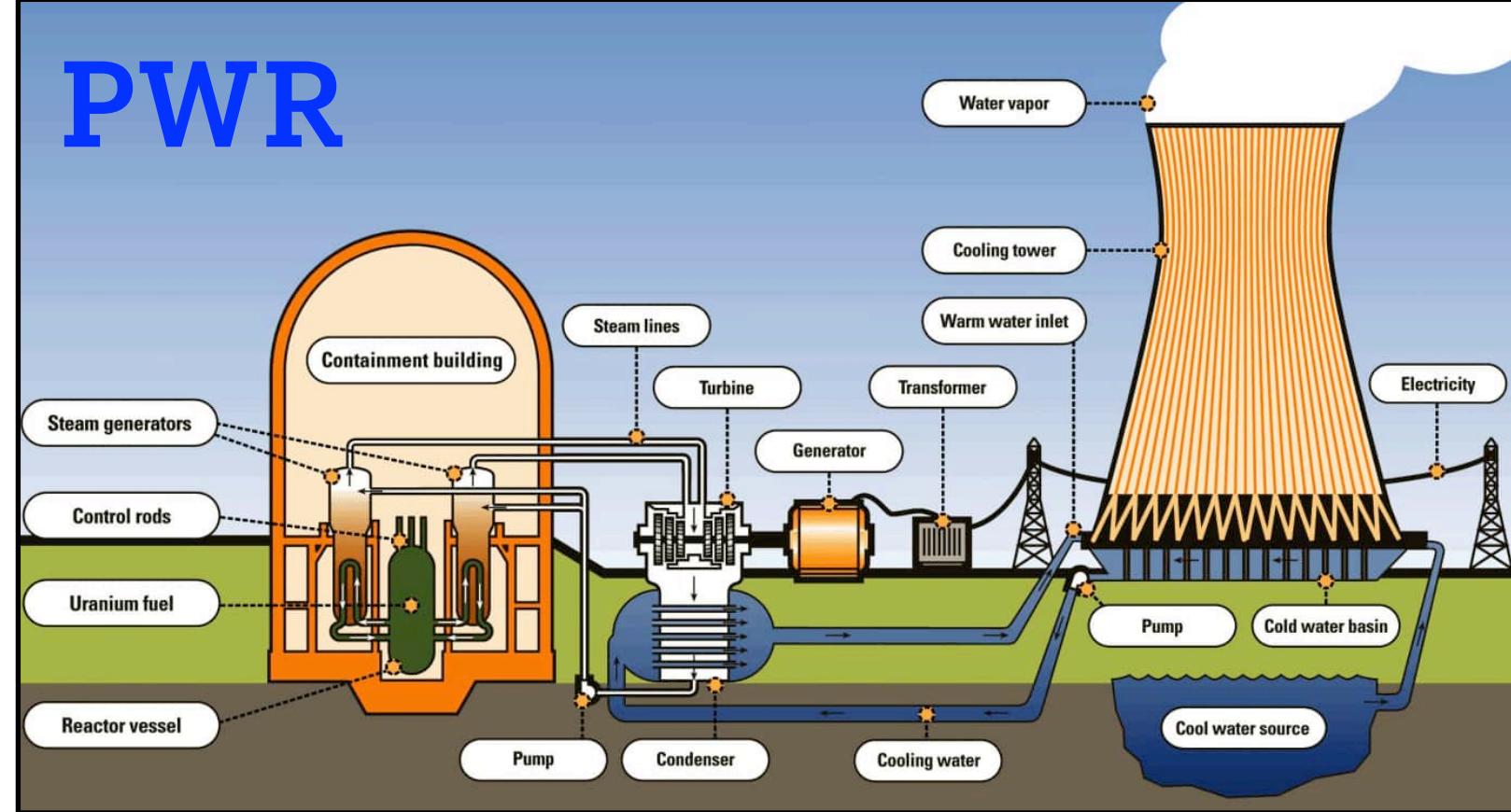
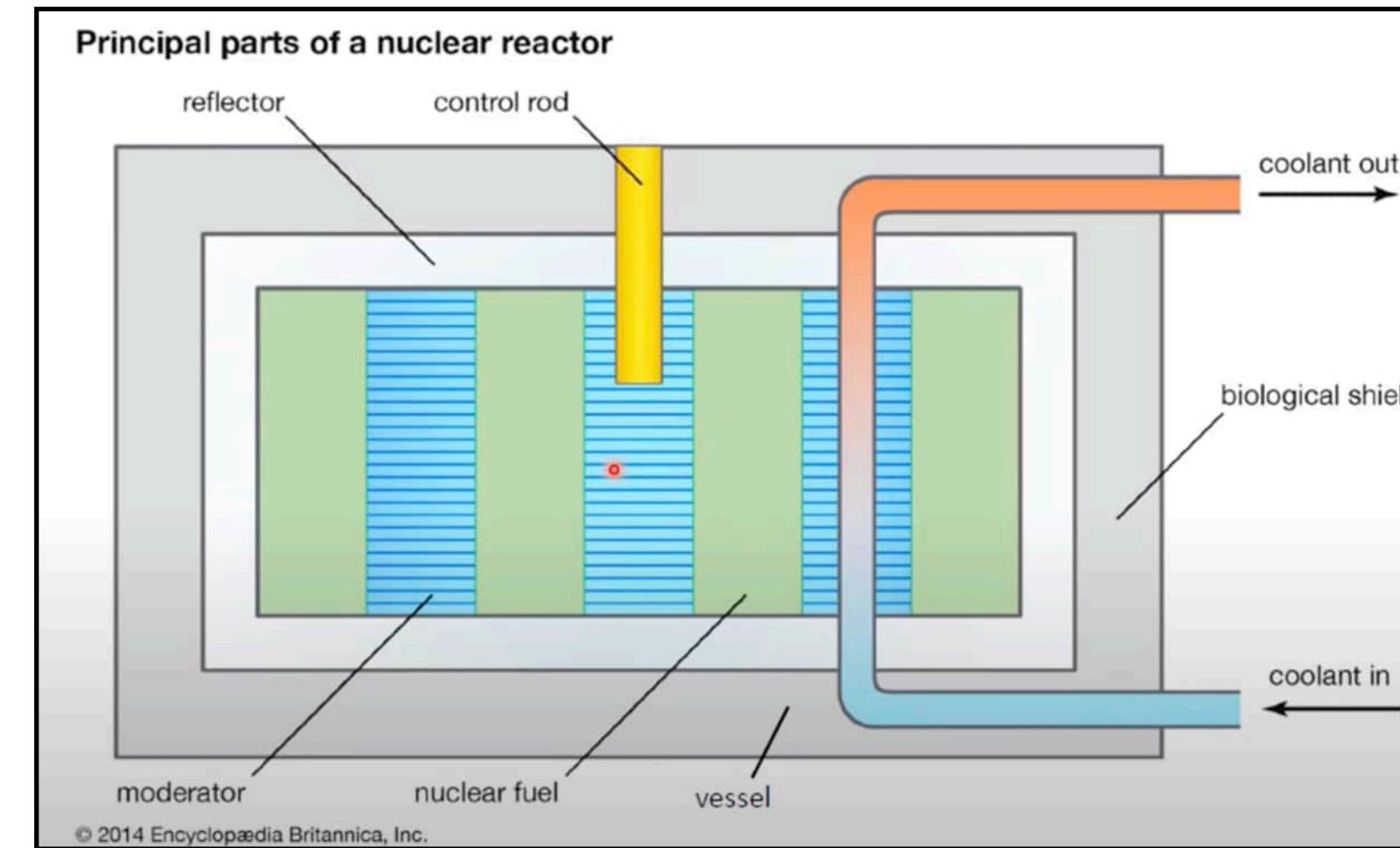


## Research Reactor

Neutrons and ionizing radiation are applied for research purposes: production of radionuclides, scattering experiments, medical applications

# Reactor Core: NPP vs RR

Main components of a nuclear reactor code: fuel, moderator control rod, containment vessel

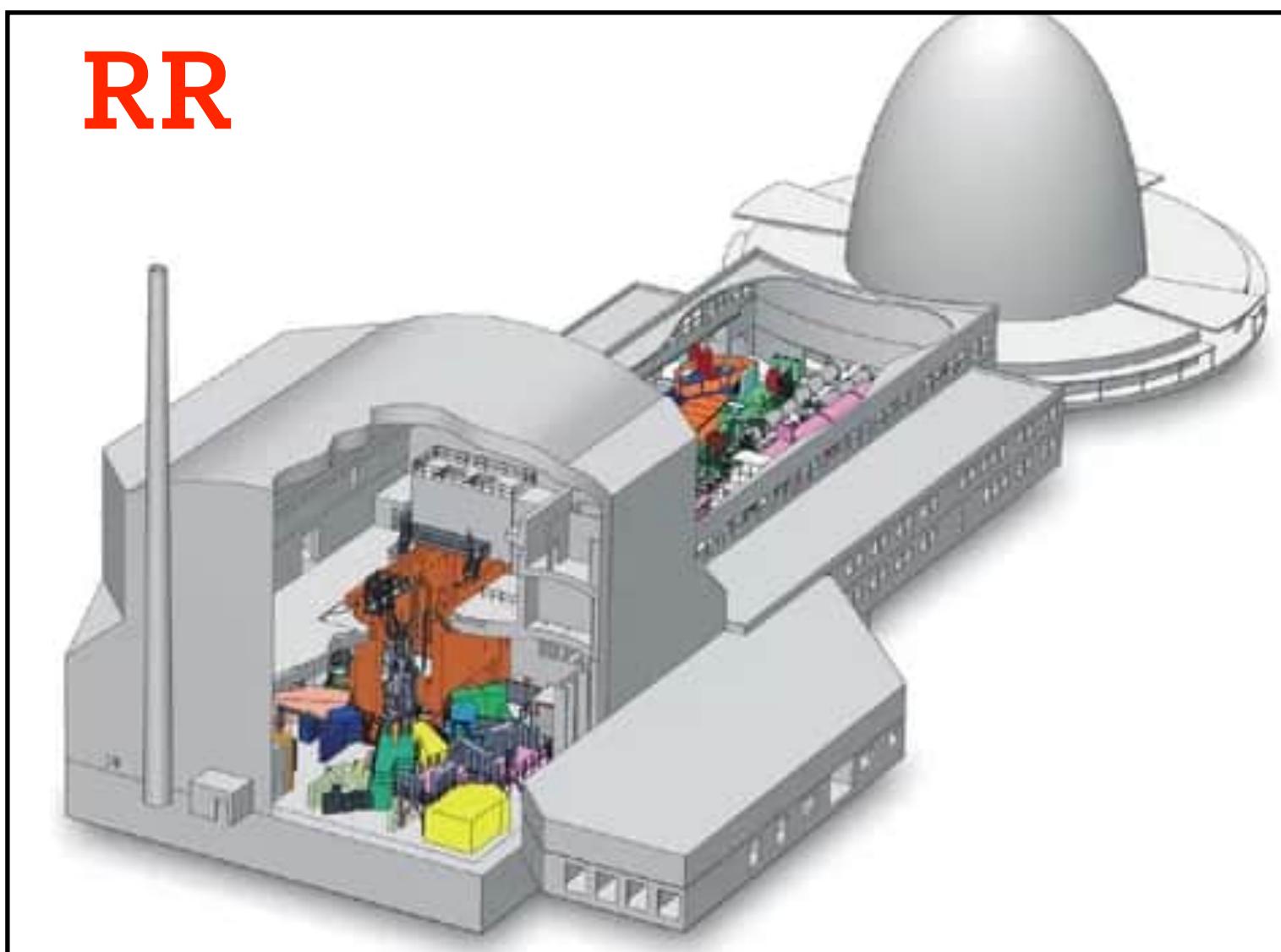


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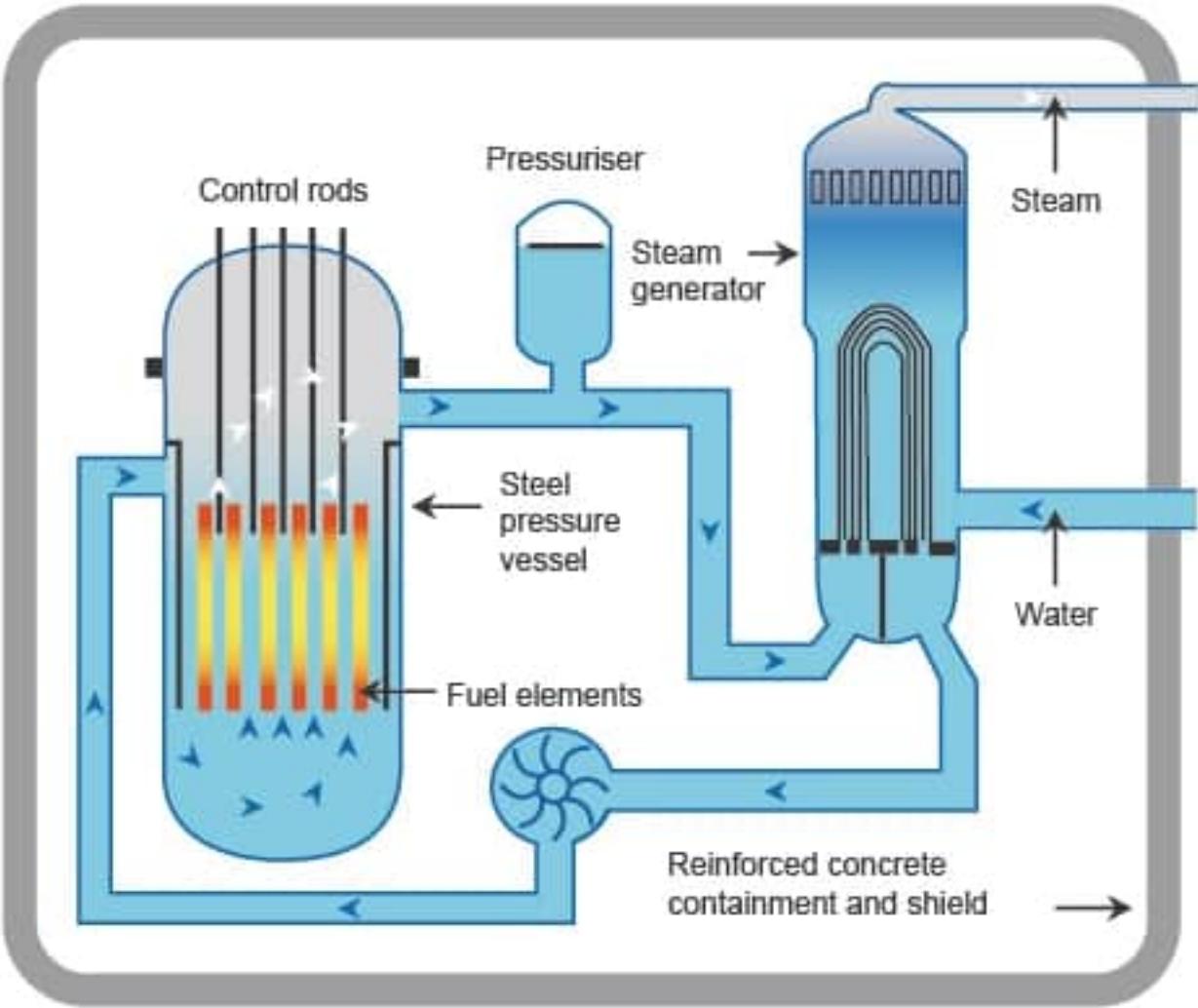
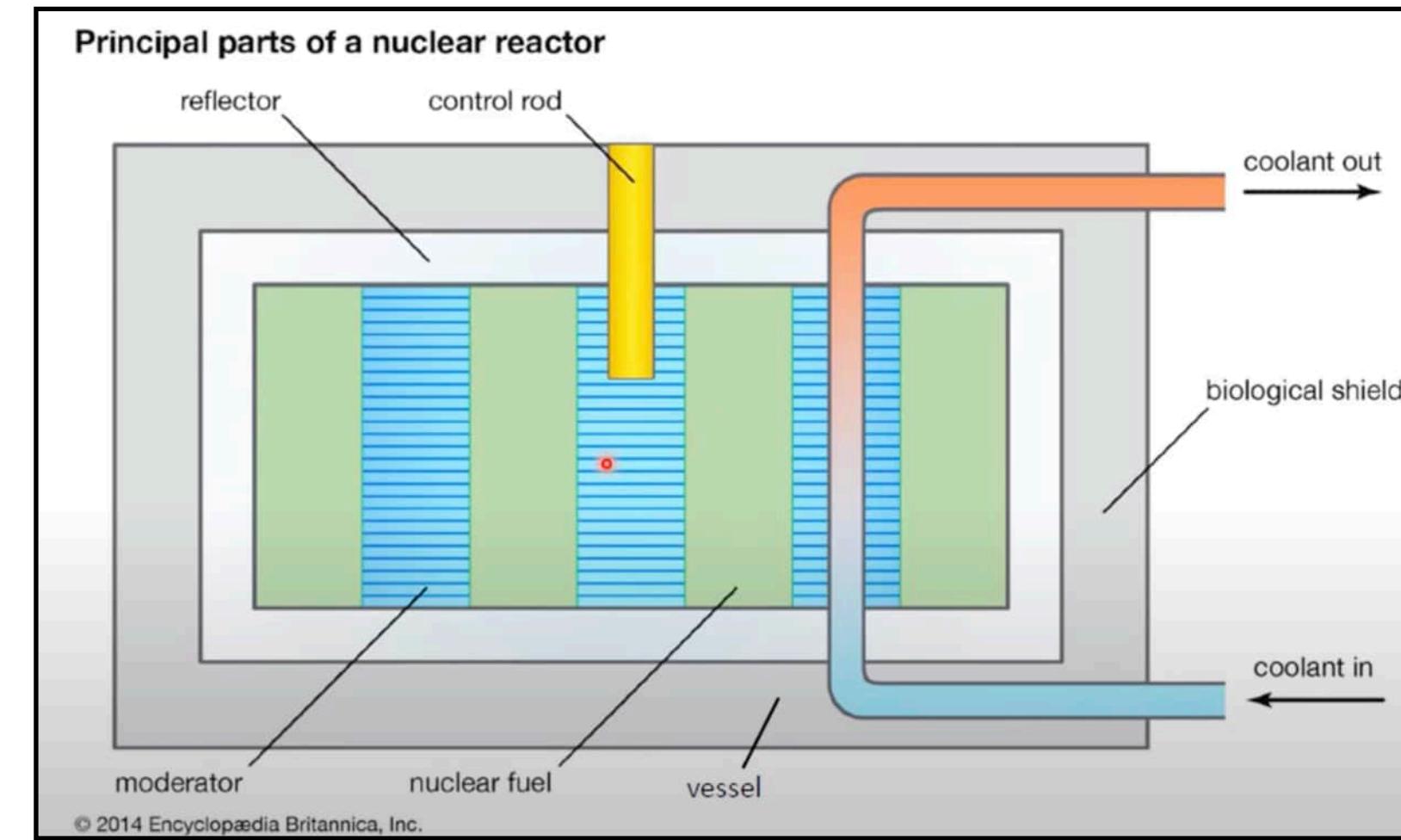
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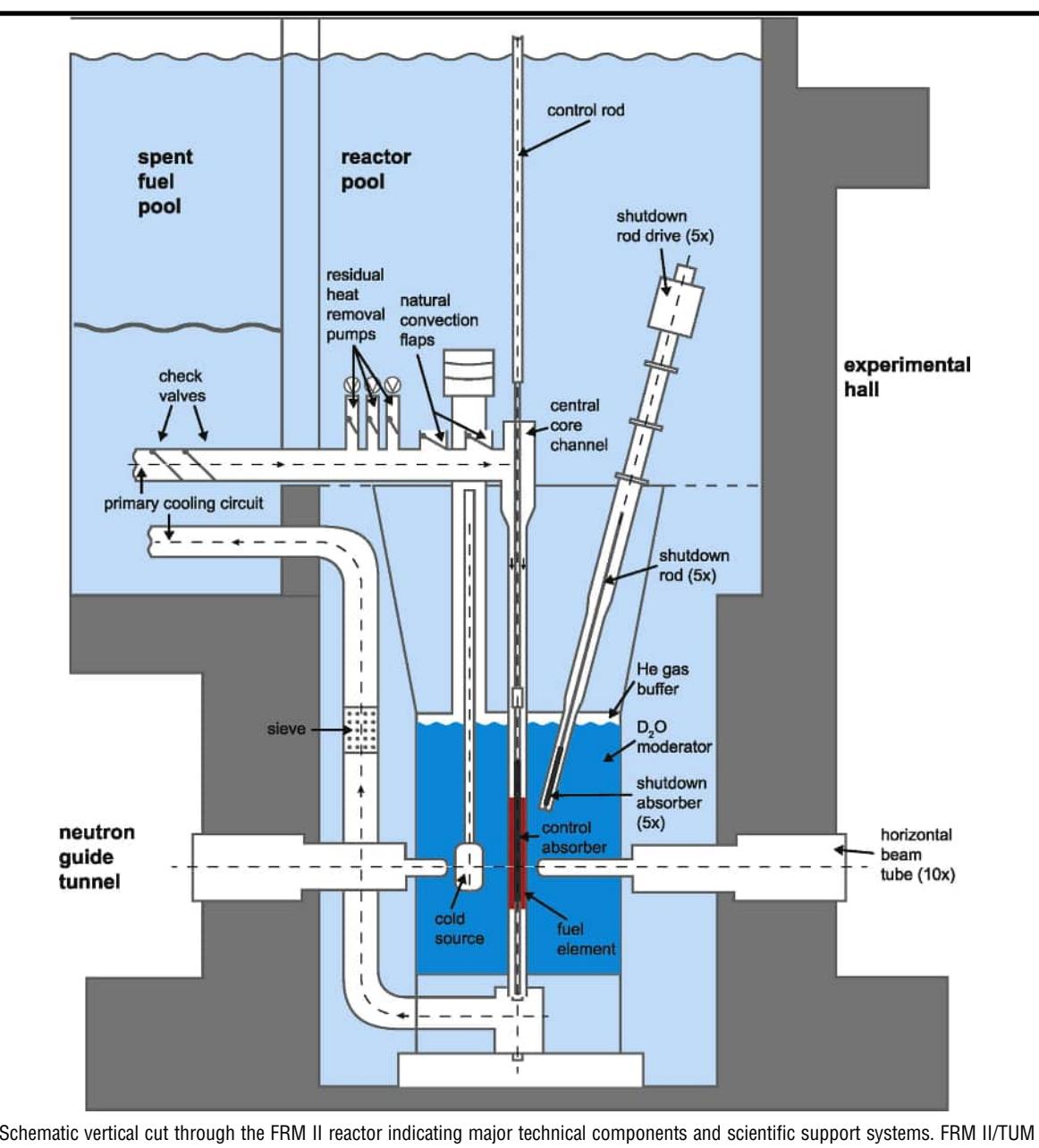


## NPP

Energy produced from fission reaction is used to generate electricity

## Research Reactor

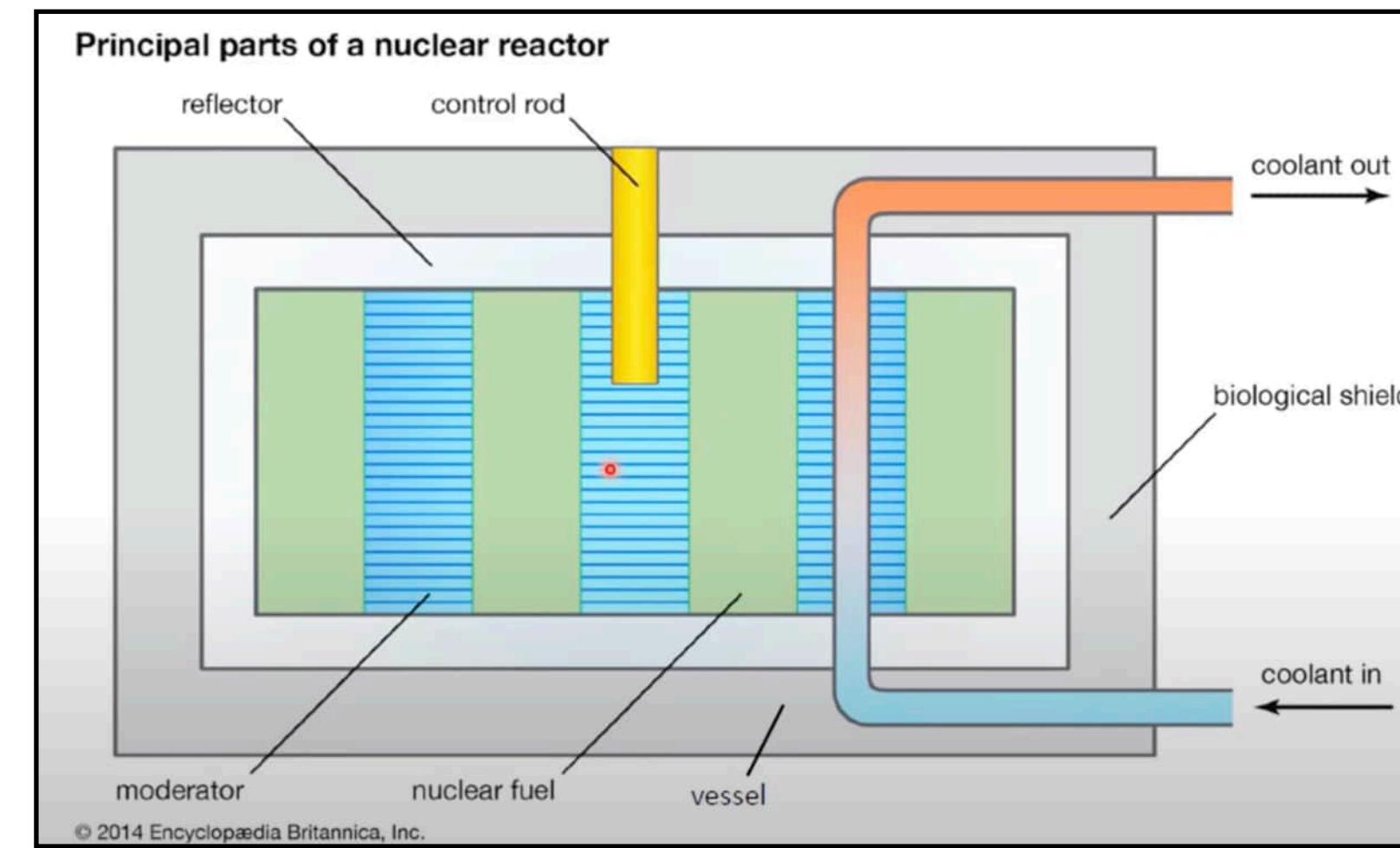
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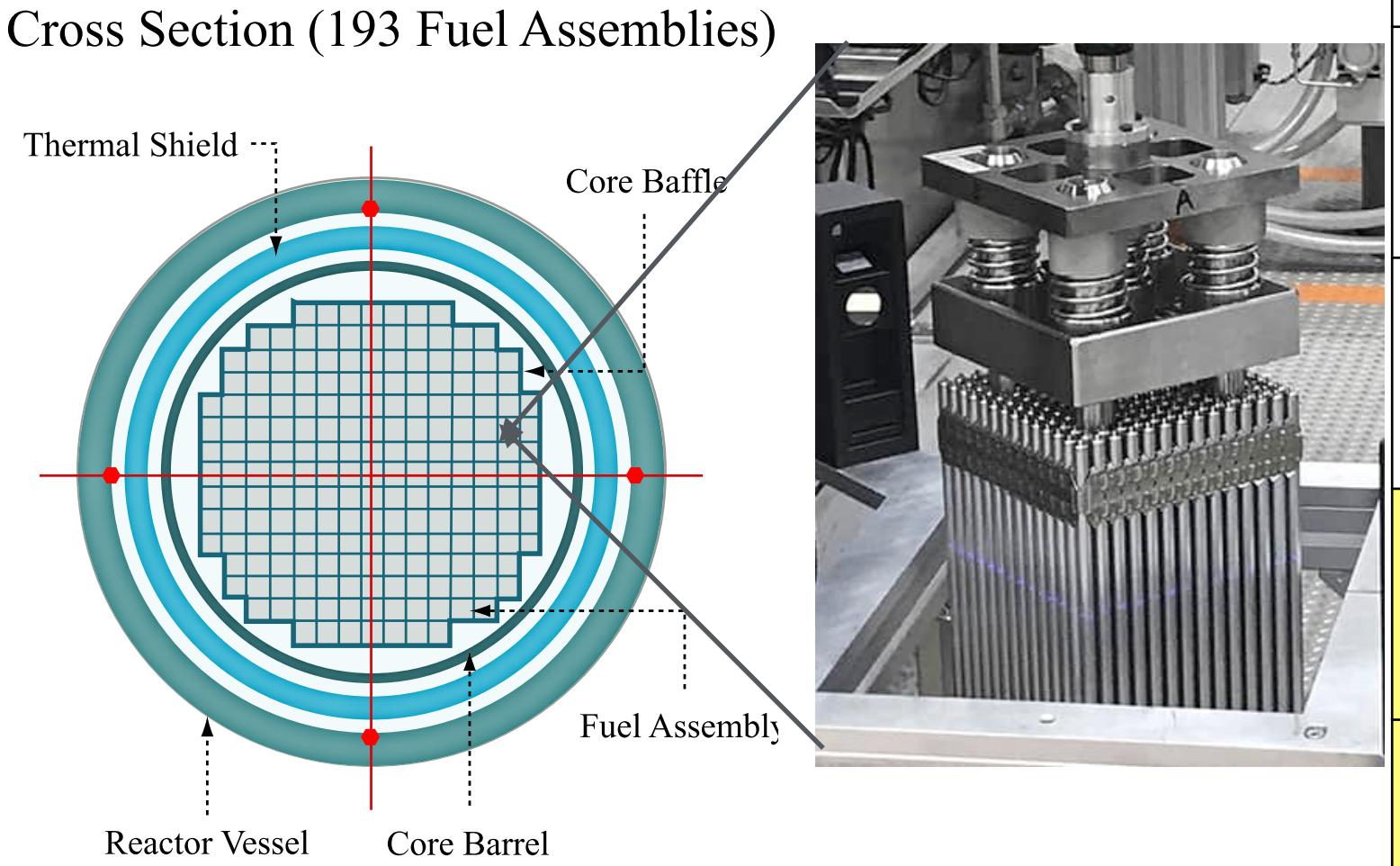
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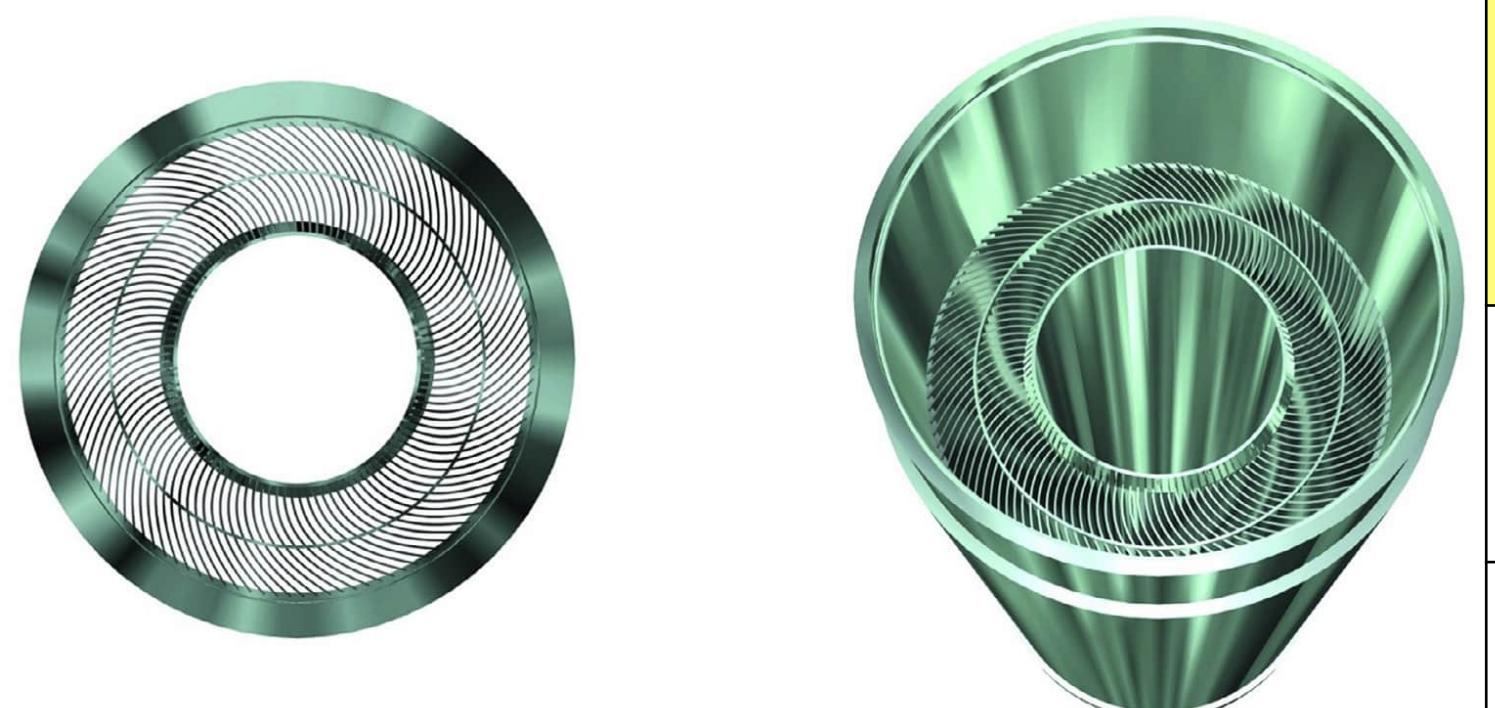
## Research Reactor

Neutrons and ionizing radiation are applied for research purposes: production of radionuclides, scattering experiments, medical applications

## PWR: fuel element



## RR: fuel element



View of the FRM II compact fuel element from different perspectives. FRM II/TUM

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# **Making “Good” Neutrons for Scattering Experiments**

**(Slowing Down Neutrons)**

# Doing Science with Low Energy Neutrons

- Neutrons always born fast. So they need to be moderated to be usable for scattering experiment

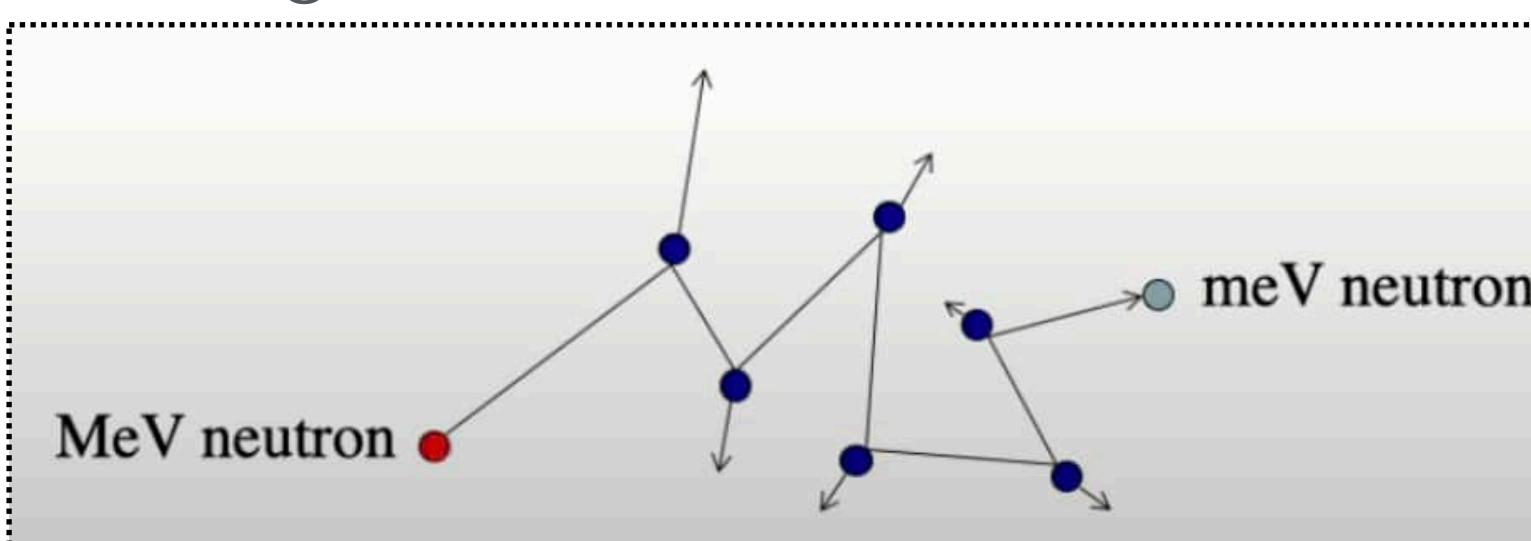
Some examples of applicative energy range for Bragg peaks

Bragg's Law  
 $2ds\sin\theta = \lambda$

Science	$d$ (Å)	$Q$ (Å $^{-1}$ )	$\lambda$ (Å)	$\theta$	$E$ (meV)
Polymers	10 000	0.001	20	0.1°	0.2
Micelles	1 000	0.01	10	0.6°	0.8
Metallurgy	100	0.1	5	3°	3.3
Crystal structures	10	1	2	11°	20.4
Glasses	1	10	1	52°	81
High precision crystallography	0.25	45	0.25	64°	1300

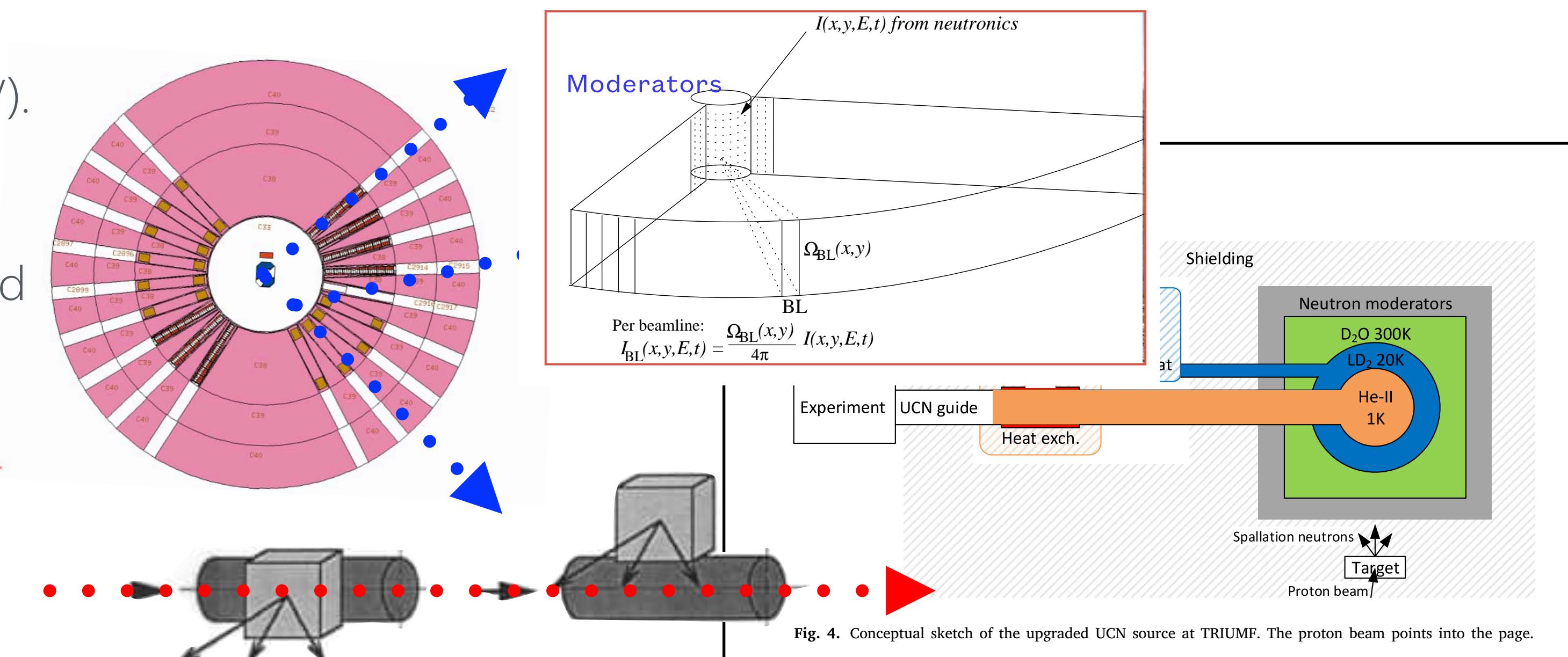
we cannot cope the whole range of science applications with only thermal neutrons (i.e 25 meV).

In order to be able to carry out different science experiment that involves different  $d$  spaces we need a broad range of lambda



In order to get an optimal design for moderators, it is needed to make an accurate choice of :

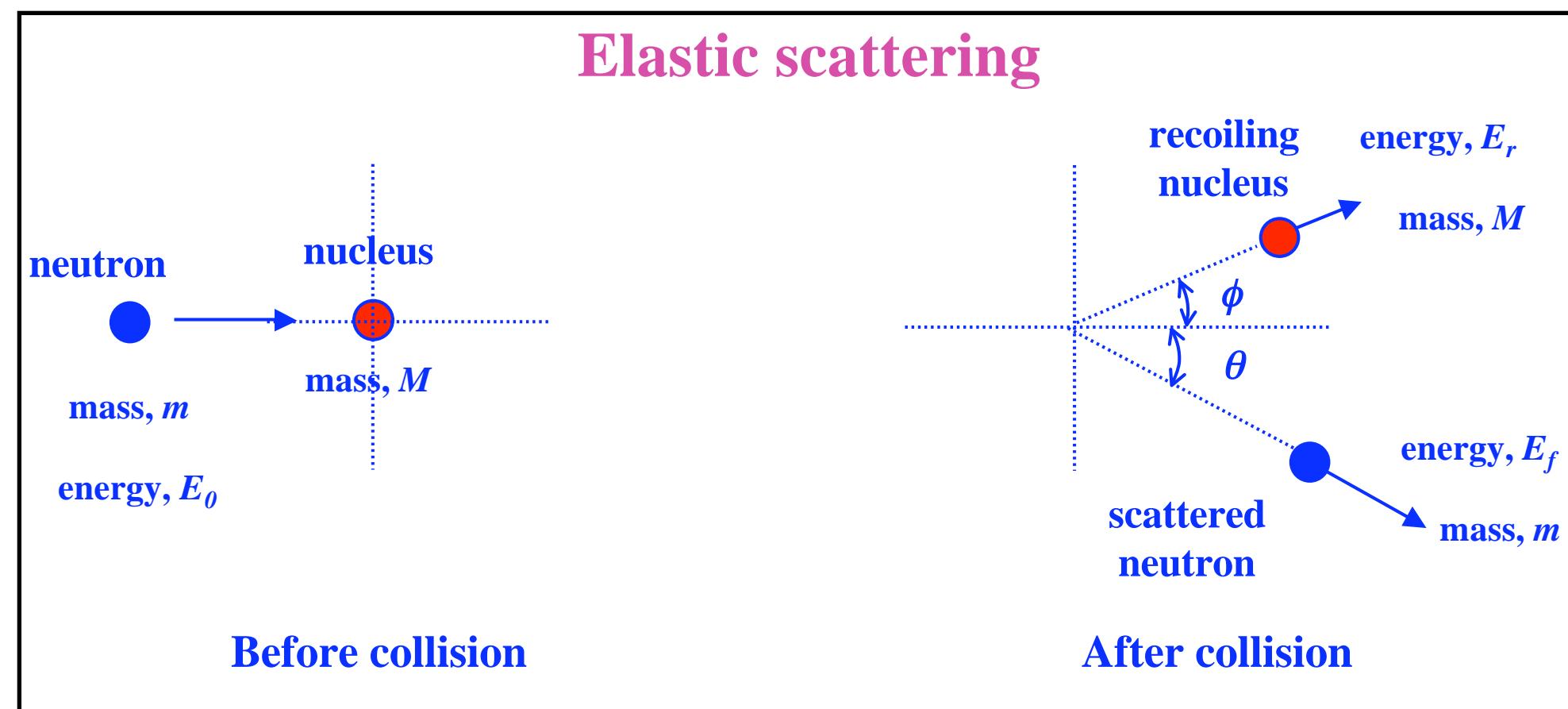
1. **material** (scattering xs)
2. **lay-out** (optimal arrangement between the source and the instrument beamlines, around the target, for SNS)
3. geometry shape and **thickness**
4. **coupling** with reflector (PSNS)
5. combination with pre moderating material for cold & ultra cold neutron



# **Moderator Material**

# Slowing Down Neutrons

Slowing-Down theory describes the loss of energy through repeated elastic collisions with effectively free nuclei. Neutron moderation occurs by “elastic” collision of neutrons with atoms of the moderator material

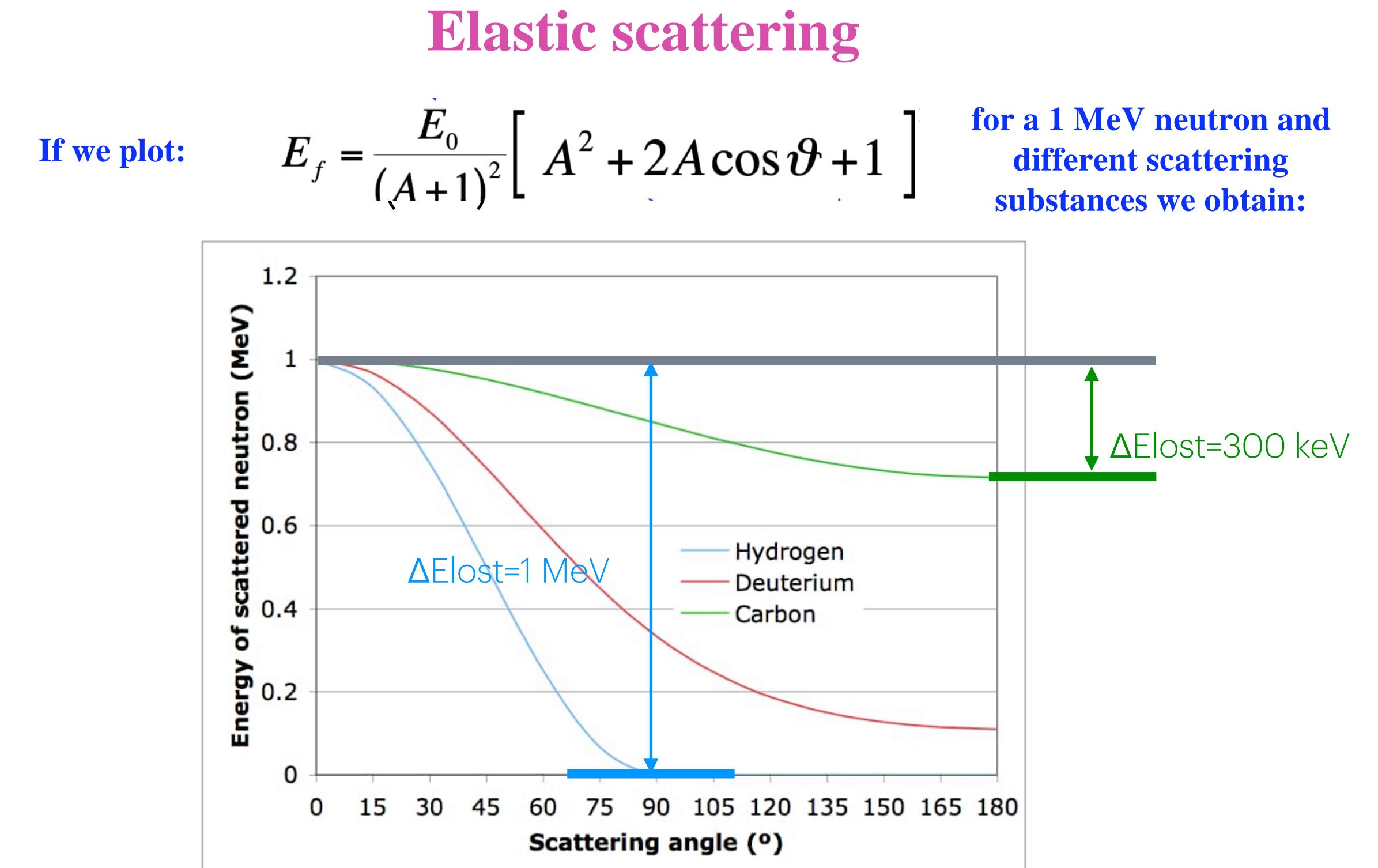


In any collision with a given scattering nucleus a neutron loses ON AVERAGE a constant fraction of the energy it had BEFORE THE COLLISION

$$\zeta = \overline{\ln\left(\frac{E_1}{E_2}\right)} = 1 + \frac{\alpha}{1 - \alpha} \ln(\alpha) \quad \alpha = \left(\frac{A - 1}{A + 1}\right)^2$$

- $\xi=1$  for  $A=1$  —> full energy loss in single collision

$$\#_{col} = \frac{\ln(E_0/E_1)}{\xi_s}$$



As you can see all the neutrons that would be scattered through  $90^\circ$  or more are basically stopped (and thermalised) in one collision with a hydrogen nucleus (proton). In the case of carbon the maximum energy loss is only about 300 keV even for a scattering angle of  $180^\circ$  ( $\pi$ ) so several collisions would be necessary to thermalise the neutrons.

# How to Choose a Good Moderator

Moderator	$\xi$ energy average loss per collision	Collisions 2 MeV to 1 eV	$\xi \Sigma s / \Sigma a$	$\Sigma s (1/cm)$ mean free path between collisions	Time to 1 eV (μs)
H <sub>2</sub> O	0.926	16	71	1.50	1.5
D <sub>2</sub> O	0.510	29	5670	0.35	9.7
Be	0.206	69	143	0.87	8.5
C (graphite)	0.158	91	192	0.38	25
Fe	0.035	411	35	0.96	43

$$\text{Moderation Ratio: } M = \zeta \Sigma_s / \Sigma_a$$

The moderating ratio or moderator quality is the complete measure of a moderator's effectiveness because it also takes into account the absorption effects. When absorption effects are high, the moderator will absorb most of the neutrons, leading to lower moderation or lower availability of thermal neutrons.

- Lighter nuclei require fewer collisions to slow the neutron and thermalise neutron faster
- **The quicker neutrons slow down, the less the initial narrow time distribution is corrupted**
- Anything with high hydrogen density is fast, and can make a slow neutron from a single collision
- Water, beryllium, and heavy water are all good moderators

# .....and Eventually Neutron Thermalisation

- Loss of energy – “slowing-down” only until up-scattering is relevant, that is up to thermalisation

$$N(E) = \frac{2\pi N}{(\pi kT)^{3/2}} E^{1/2} e^{-E/kT}.$$

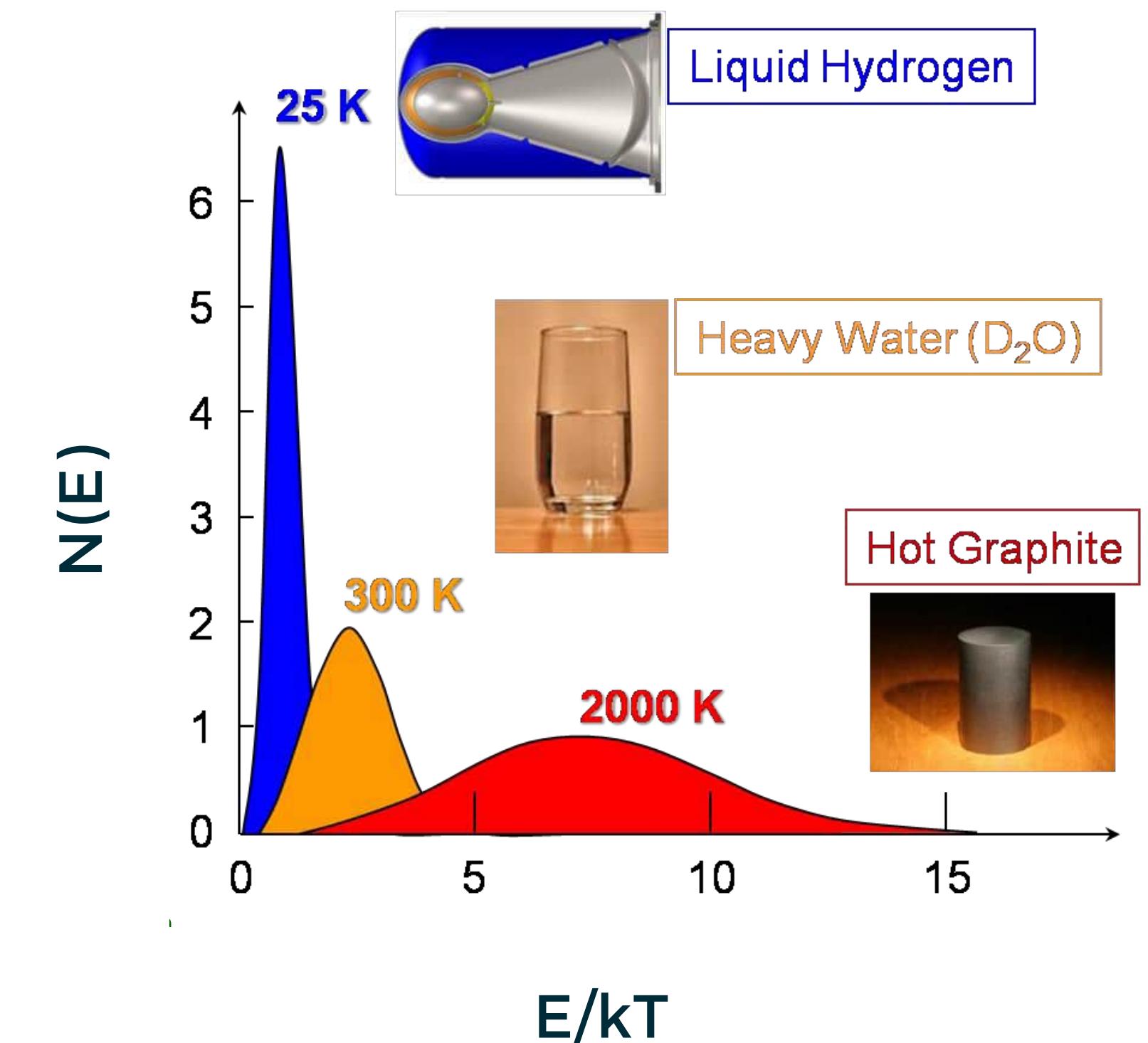
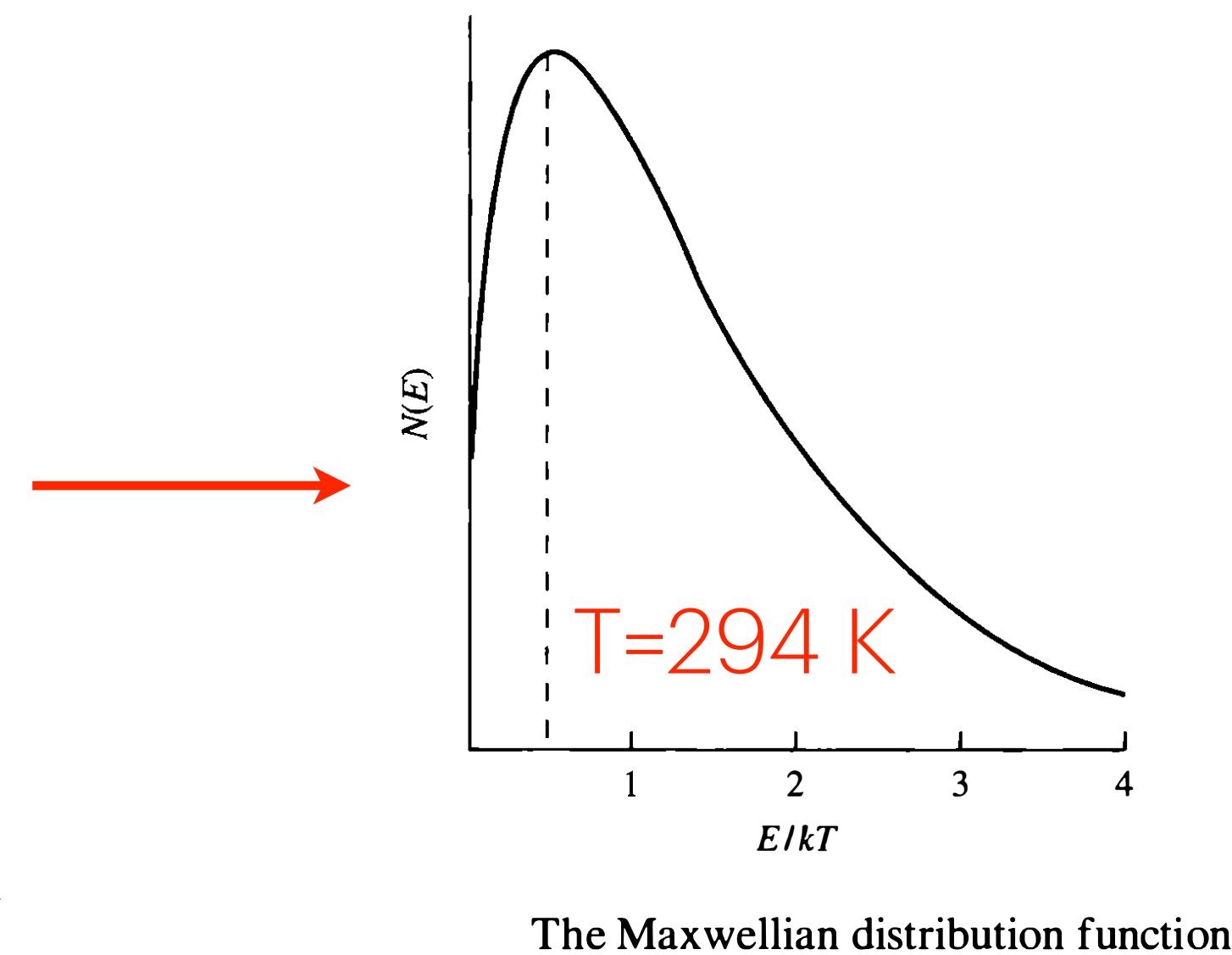
$$E = \frac{1}{2}kT$$

The most probable Energy

$$T_0 = 293.61^\circ\text{K}, \quad kT_0 = 0.0253 \text{ eV} \simeq \frac{1}{40} \text{ eV}$$

$$v(E) = \left(\frac{2E}{m}\right)^{1/2}$$

$$v(293.61 \text{ K}) = 2200 \text{ m/s} \quad E = 25 \text{ meV}$$



By changing the moderator temperature one can shift the thermal equilibrium (i.e most probable Energy in the spectrum) to a higher or lower energy

Whilst the atom free approx is considered reasonable in the range of energy  $> 1\text{eV}$ , below this energy accurate scattering kernel able to describe the mechanisms of energy loss by excitation of vibrational and rotational modes in molecules (nuclei have to be considered bound and no more free)

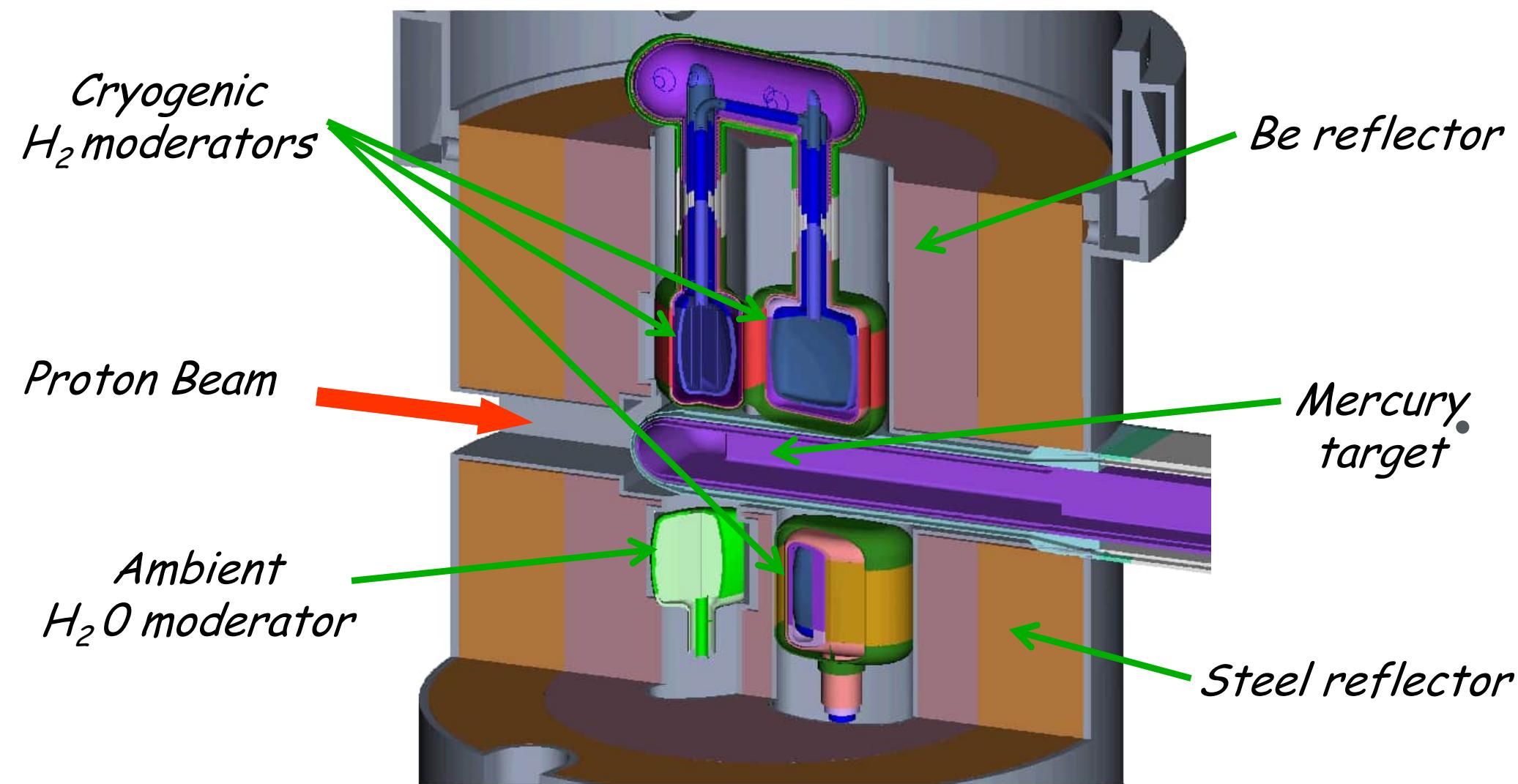
# **Moderators:**

## **Layout and Geometrical Shape**

# Moderators: Effective Source for Instruments

## Moderators layout in spallation sources and in reactors

### SNS vertical cross section

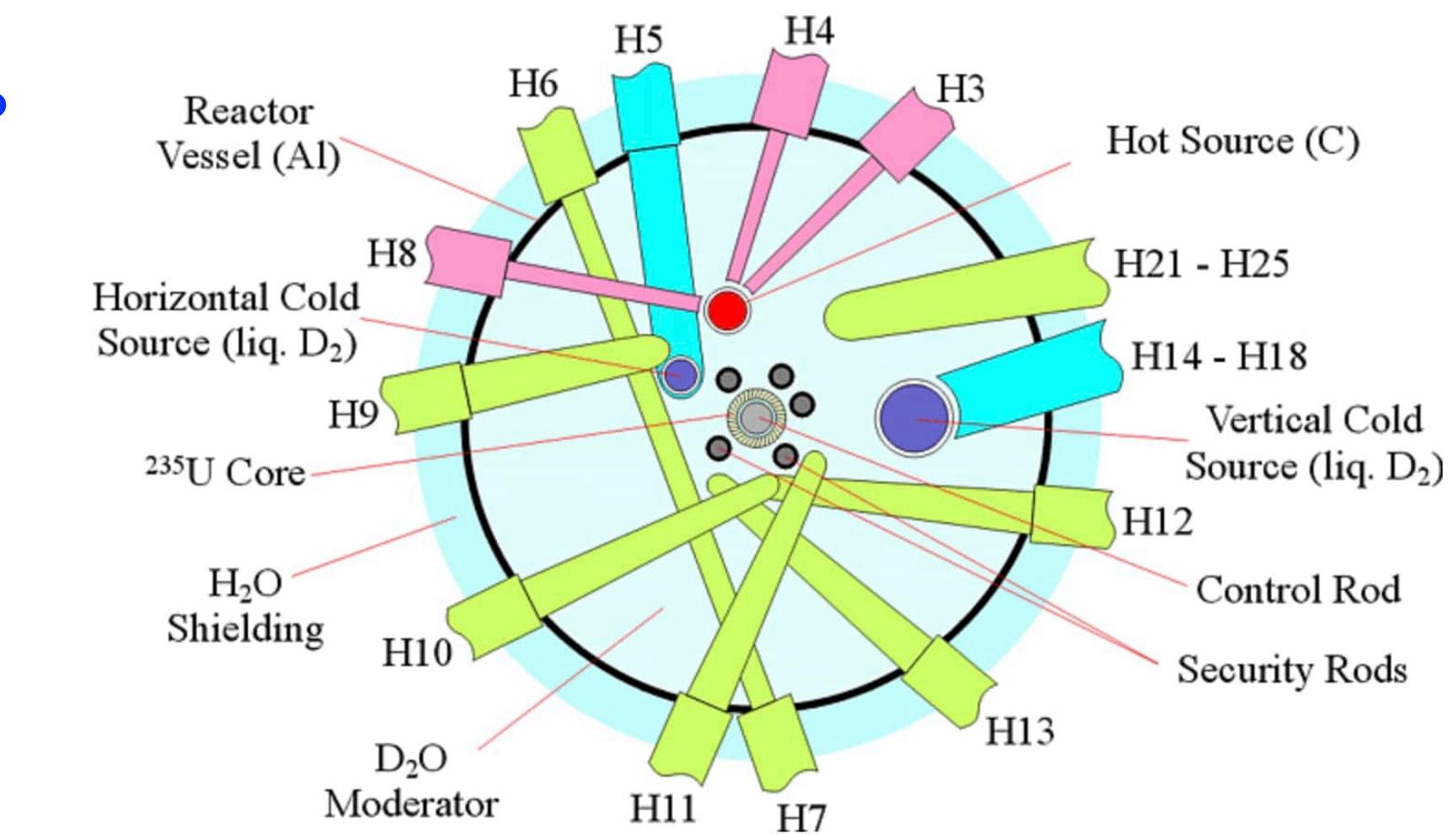


- Main goal of moderators: shift the pulse down in energy saving the intensity and reducing the broadening (for pulsed source)

BUT many neutrons are absorbed or lost on the road from the source to the sample and long tail could be added to the pulse!

- Usually target stations have a reflecting material (such as beryllium) to help the moderation process and improve the moderator efficiency.

### ILL horizontal cross section

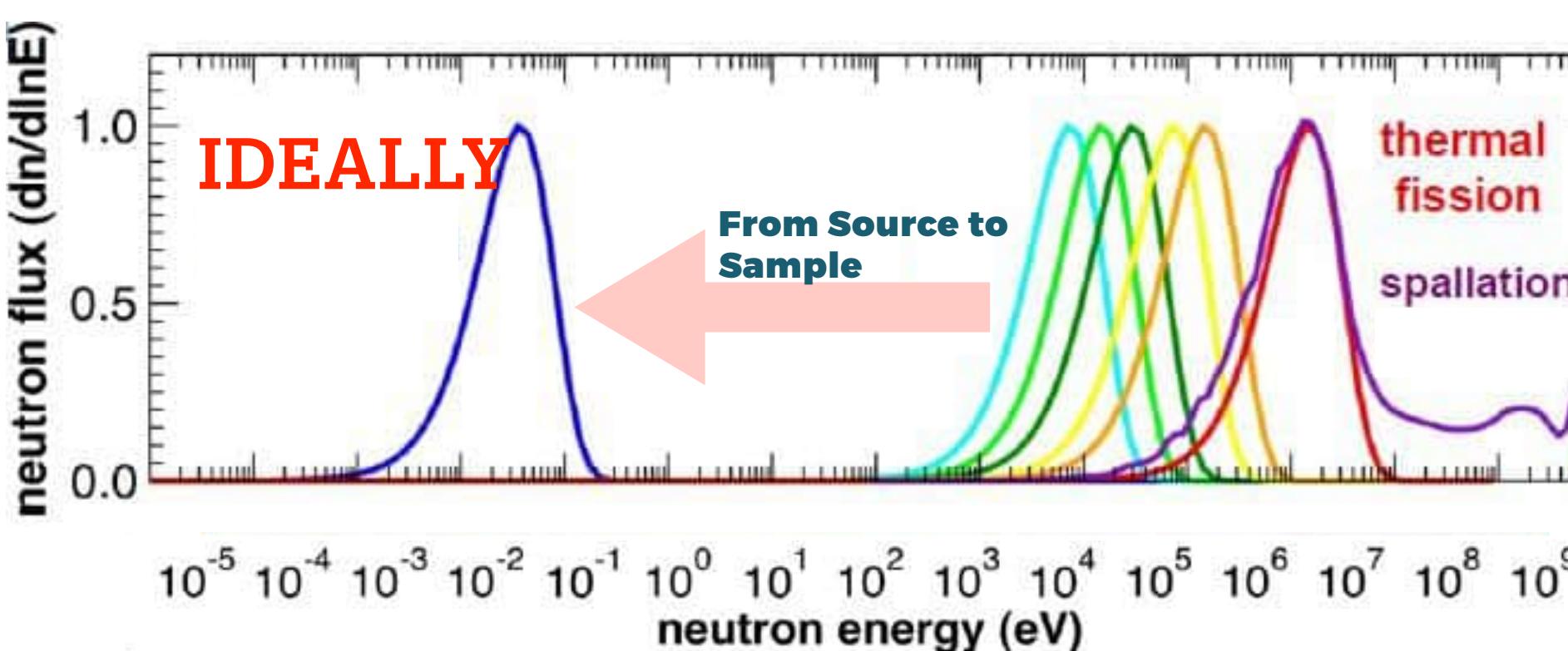


Thermal reactors already need to slow-down neutrons to keep the chain reaction working

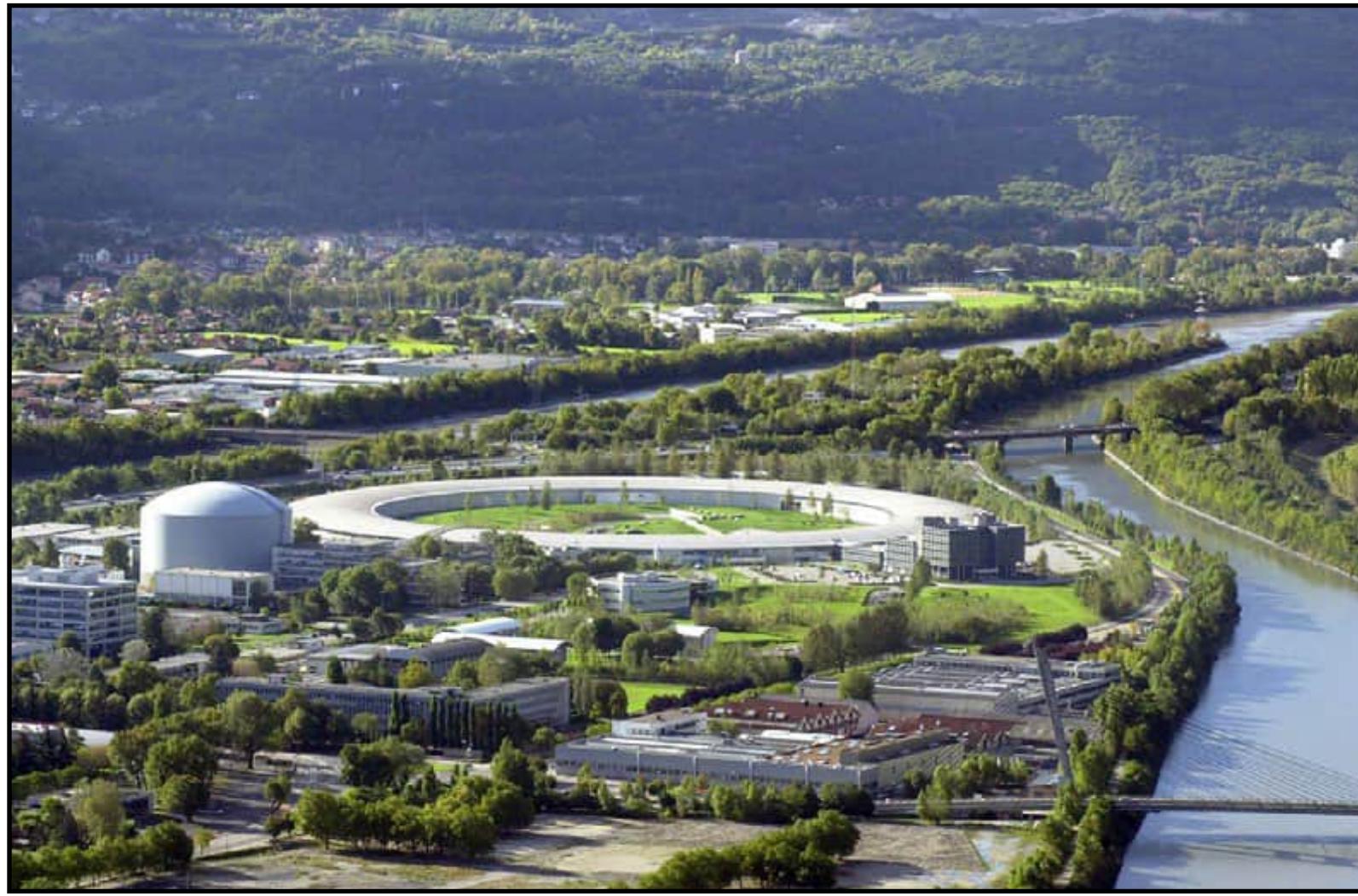
But additional moderators could be required to reach lower energy range.

Typically in research reactors, moderator for feeding the instruments are separated from moderator used for operational purposes

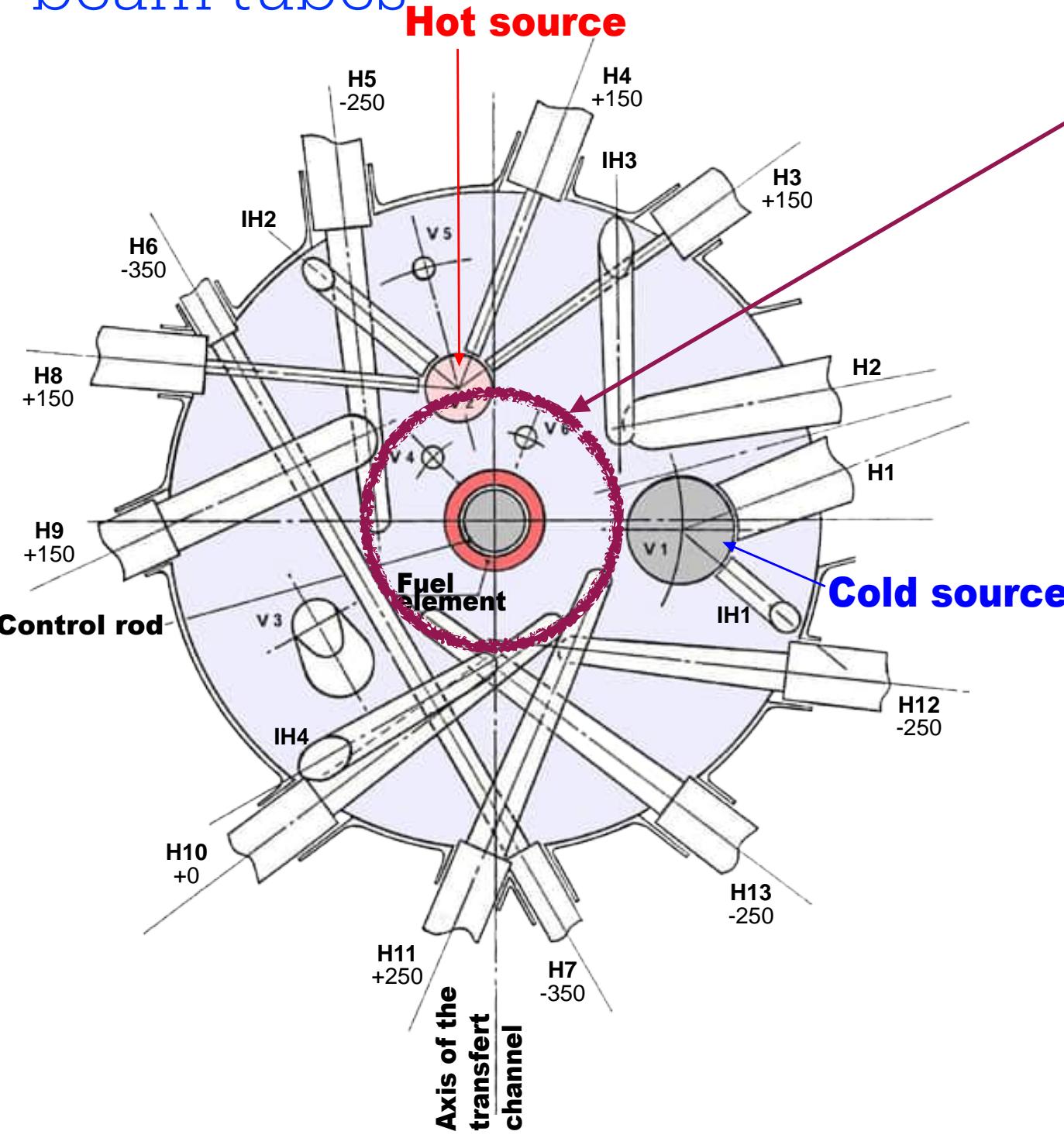
Produce and provide access to the neutrons from external beams



# RR thermal, Cold and Hot Moderators

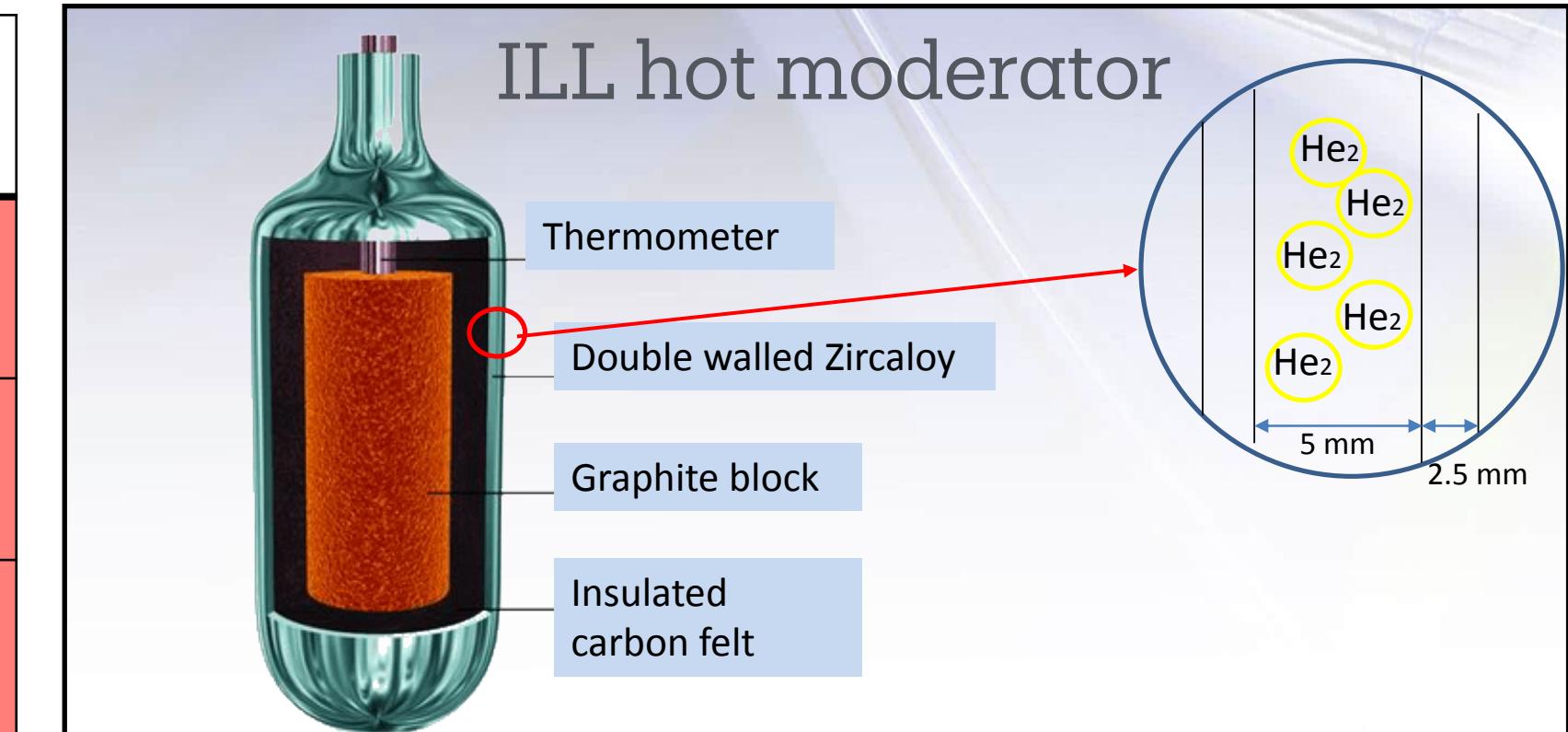
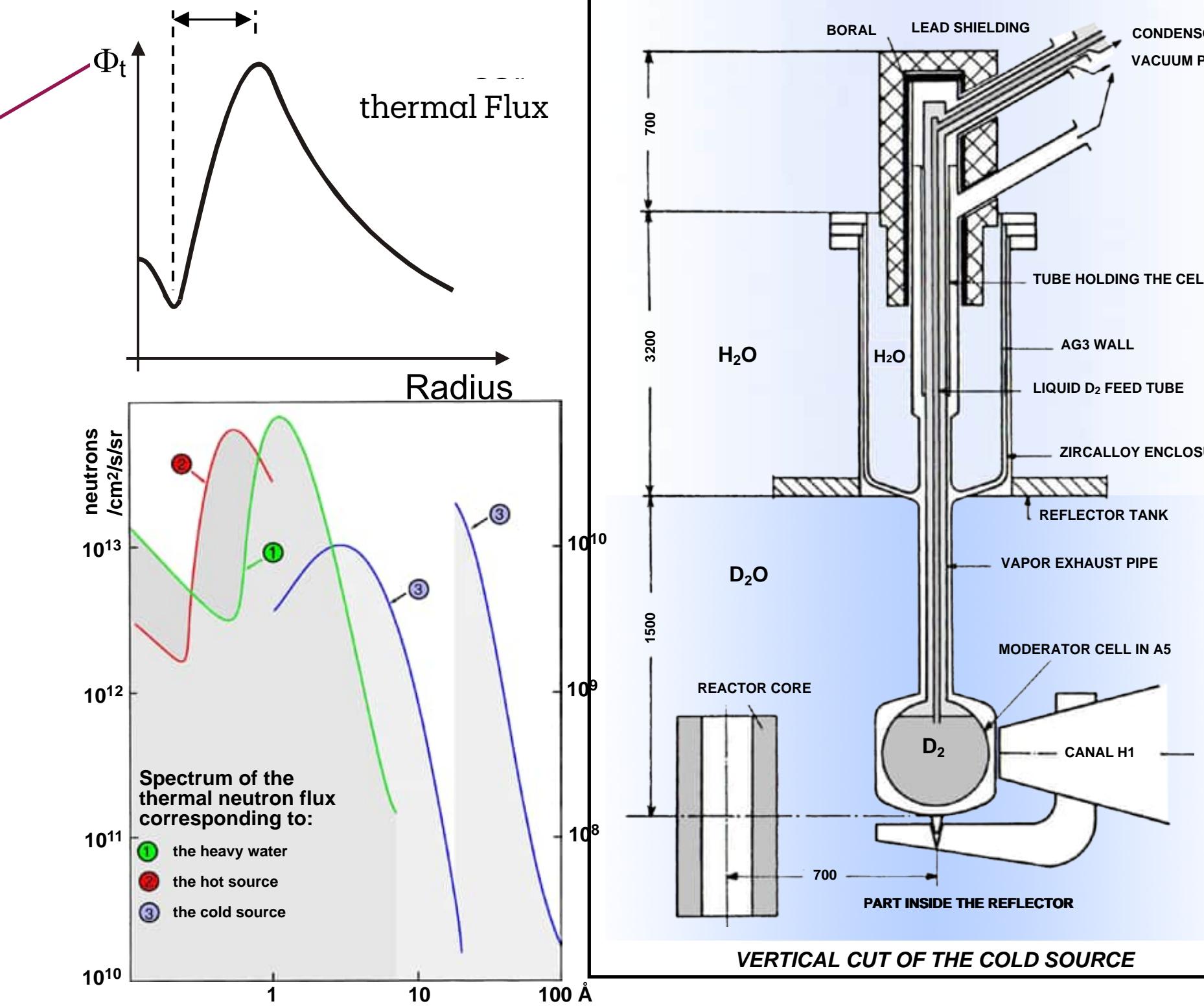


- Arrangement of the experimental beam tubes



- Highly enriched uranium (93% U235)
- Compact design for high brightness
- heavy water cooling
- heavy water reflector
- Single control road
- 57 MW thermal power
- Cold, Thermal and hot sources

ILL	cold	thermal	hot
moderator	liquid D <sub>2</sub>	liquid D <sub>2</sub> O	graphite
Mod. Temp.	20K	300K	2000K
neutron wavelenght	3→20 Å	1→ 3Å	0.2→1 Å

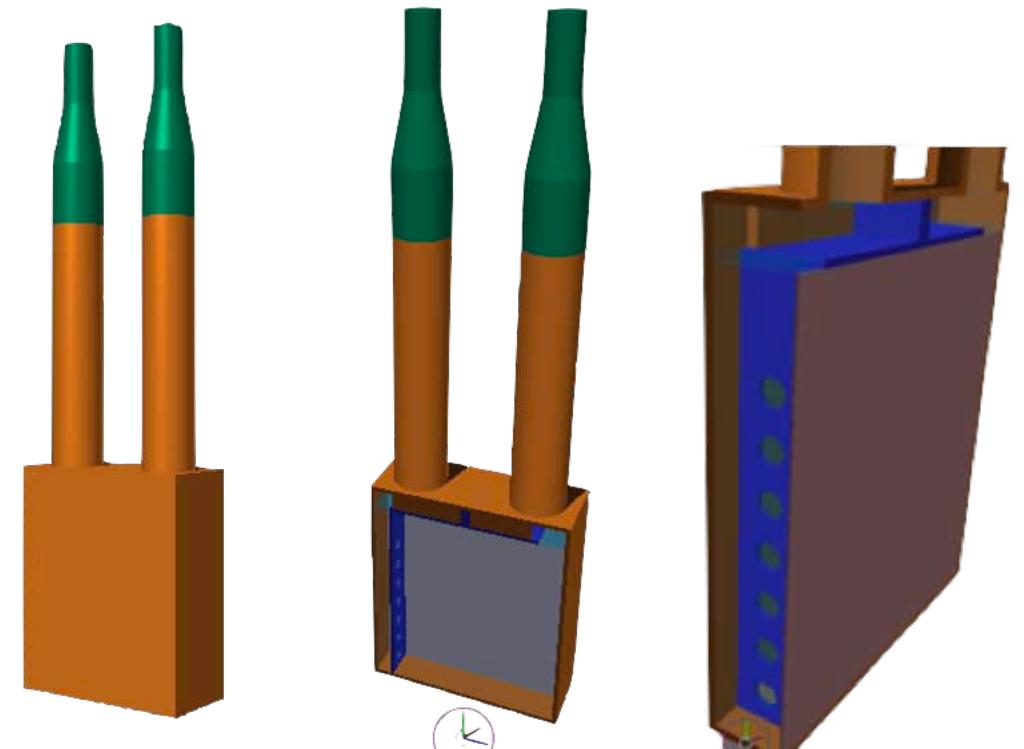


# Neutron Pulses from the Moderator Surface

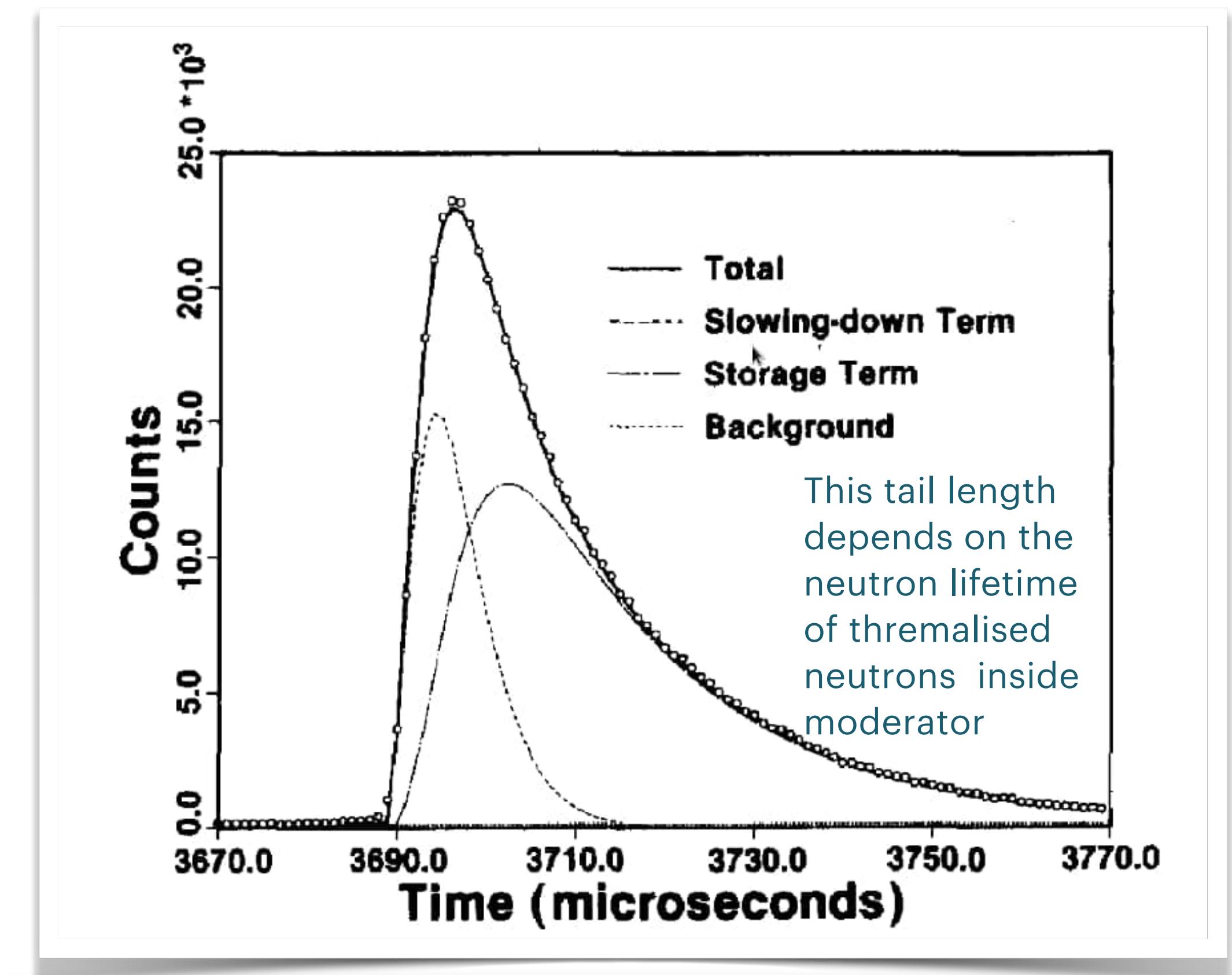
The function of a moderator is to convert leakage-neutrons (from a target) to slow-neutrons with energy spectrum and pulse characteristics as required for experiments.

- A relatively small hydrogenous moderator is essential for Pulsed SNS to obtain reasonably narrow pulses, but at a large penalty in time-integrated and peak intensities.
- Neutrons reach thermal equilibrium with the scattering material.
  1. Slowing-down neutrons come out quickly
  2. Thermalised population decays more slowly;
- Overall emission is the sum of slowing-down term and thermalised one
- Both the slowing-down and storage component will always be present in a moderator

Aluminum box dimension 10x10x5cm



ISIS water moderator



For short pulses: fast slowing down component and slow neutron lifetime in the moderator after the thermalisation are desirable (absorbing material inside the moderator)

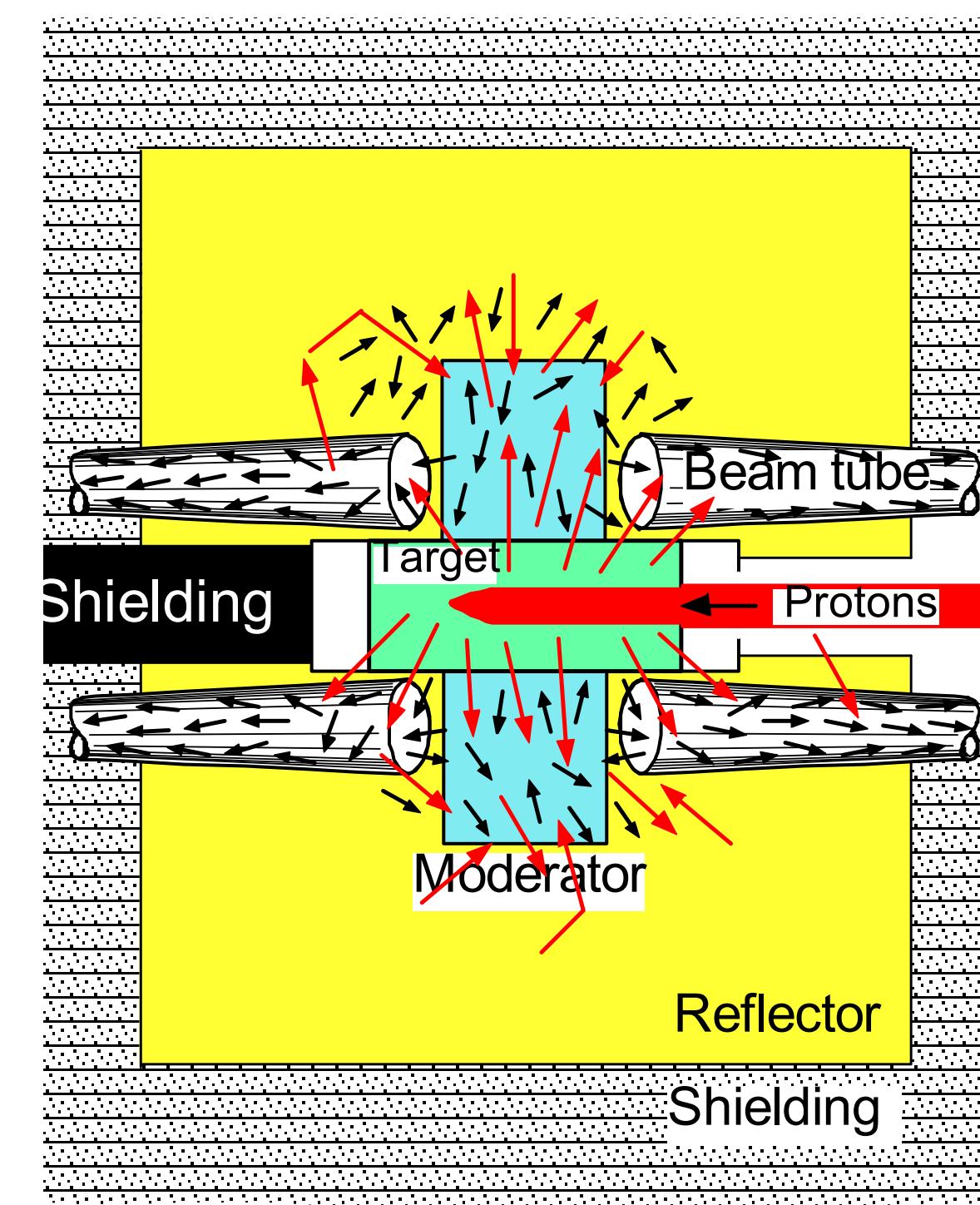
# Target Reflector And Moderator Assembly

Fast neutrons are generated by proton beam injected in a target

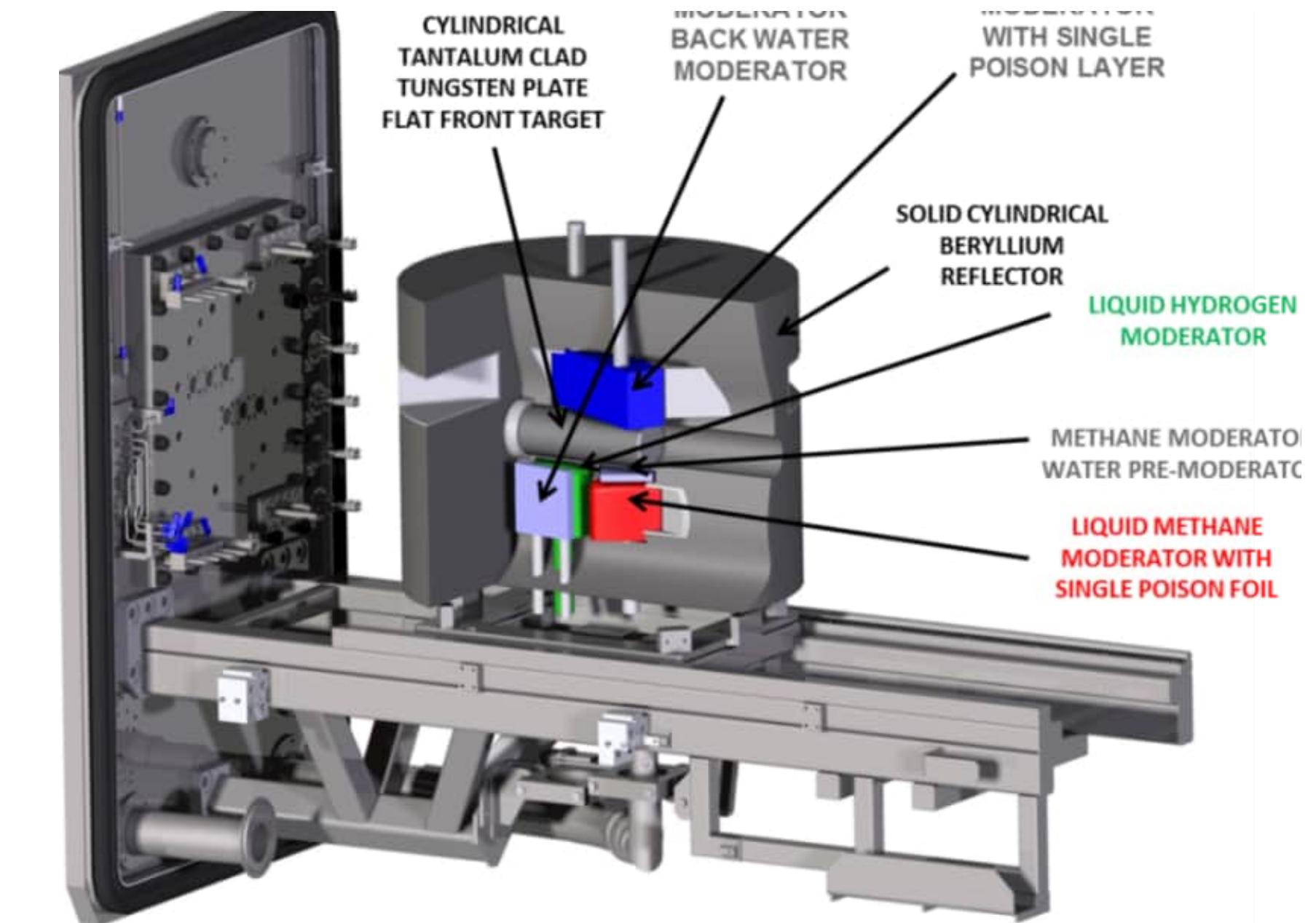
In order to be useful for scattering experiments neutrons must be slow-down to meV energy range by MODERATORS

Optimal coupling between moderator and target is required

In order to increase this coupling, moderator are surrounded by reflector whose task is to scatter those missing the moderator back into moderators



**ISIS TRAM module**



## The reflector role

The function of the reflector is to enhance the slow-neutron intensity by reflecting leakage-neutrons from the target that do not directly enter the moderators.

Typical reflector material: D<sub>2</sub>O, Be, Fe, Pb

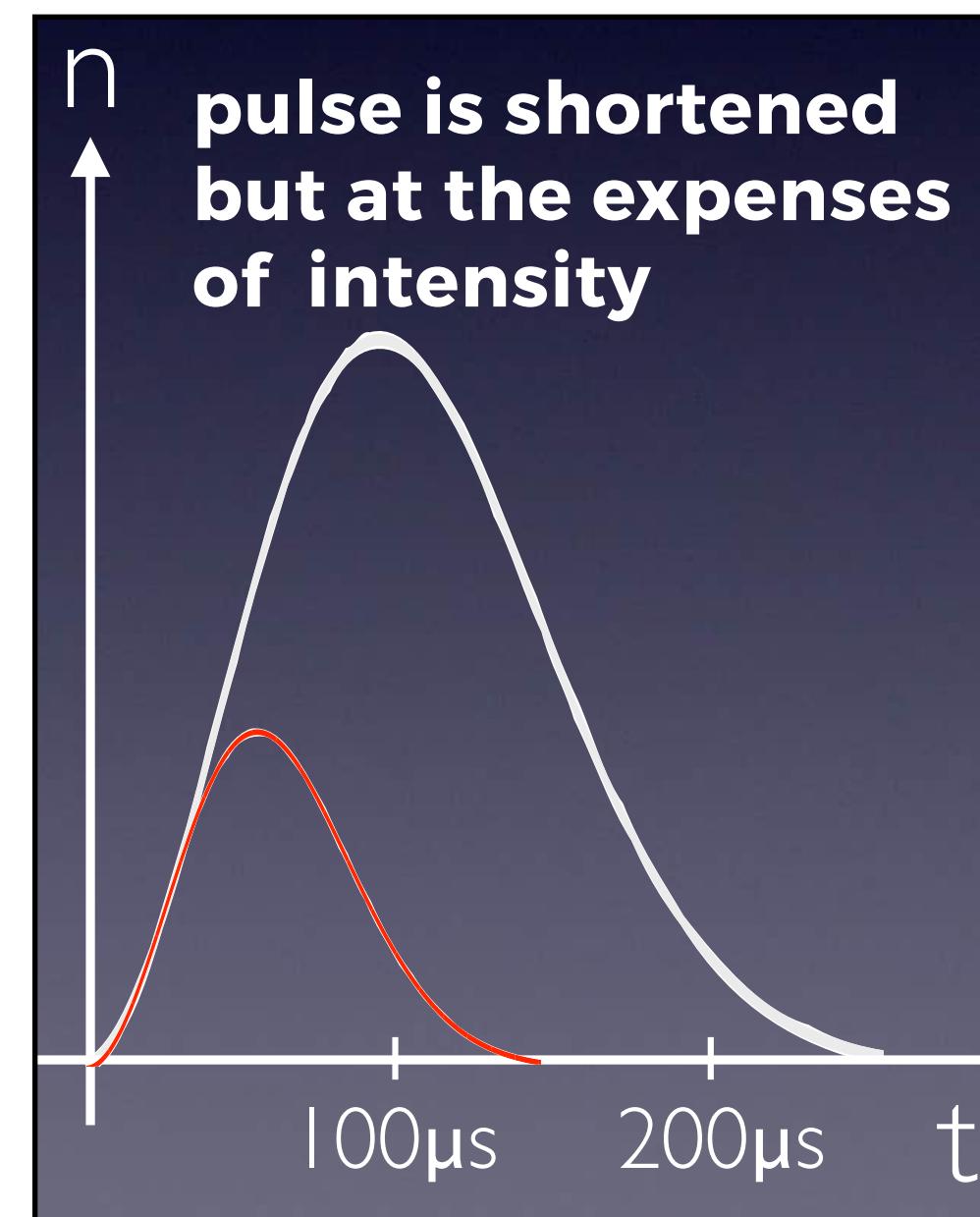
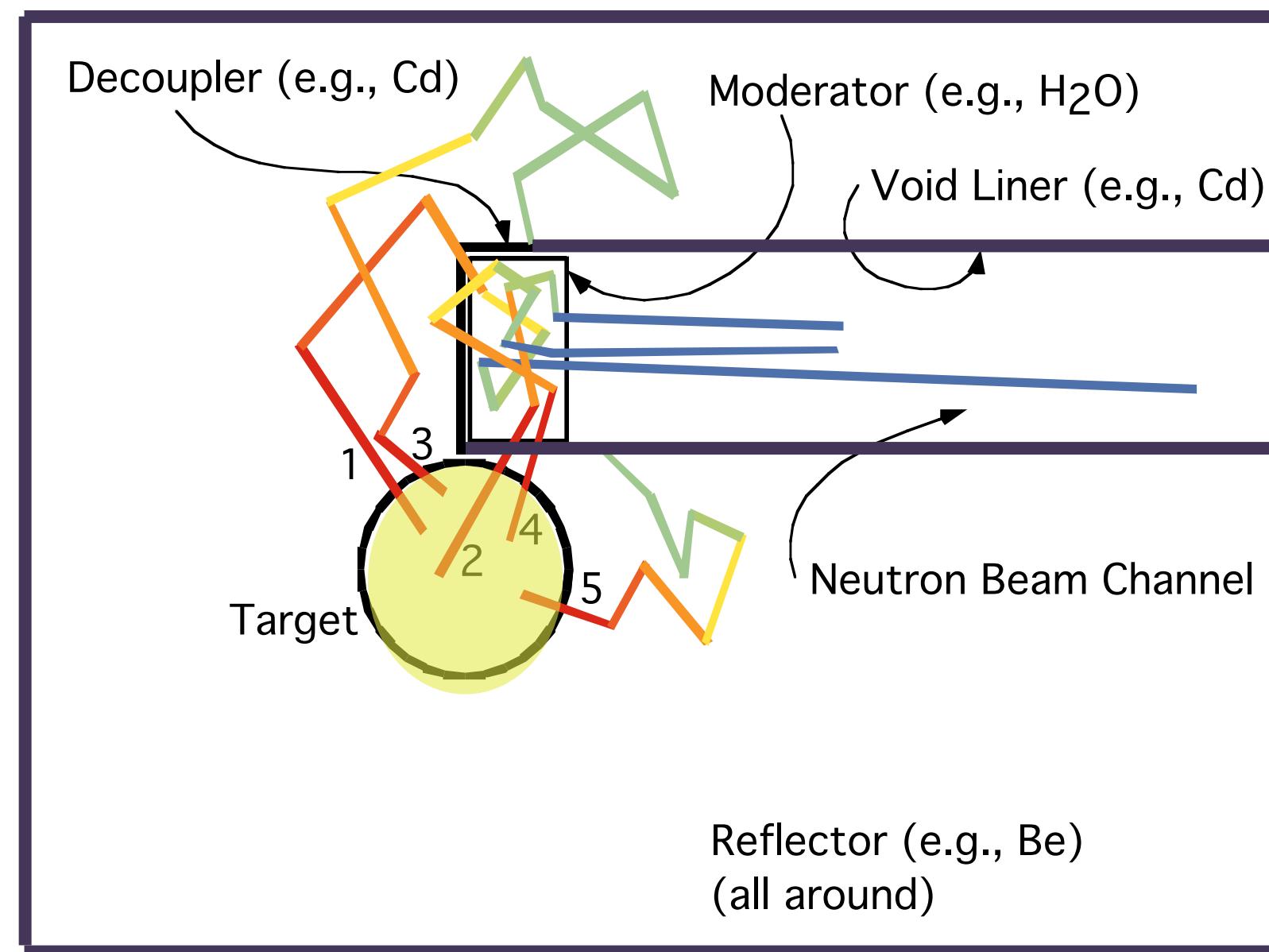
Neutrons can be also produced by inelastic interaction in reflector

Reflector are classified in 2 categories: moderating (Be) and not moderating (Pb)

# PSNS specific feature: Coupled and Decoupled Moderators

Spallation neutron sources offer a wide variety of choices for moderator and reflectors and require different technical solutions for different optimisation criteria.

This peculiarity makes spallation source more flexible than reactor ones



Purpose of a decoupled moderator: provide the narrowest possible neutron pulses with the highest possible peak intensities,

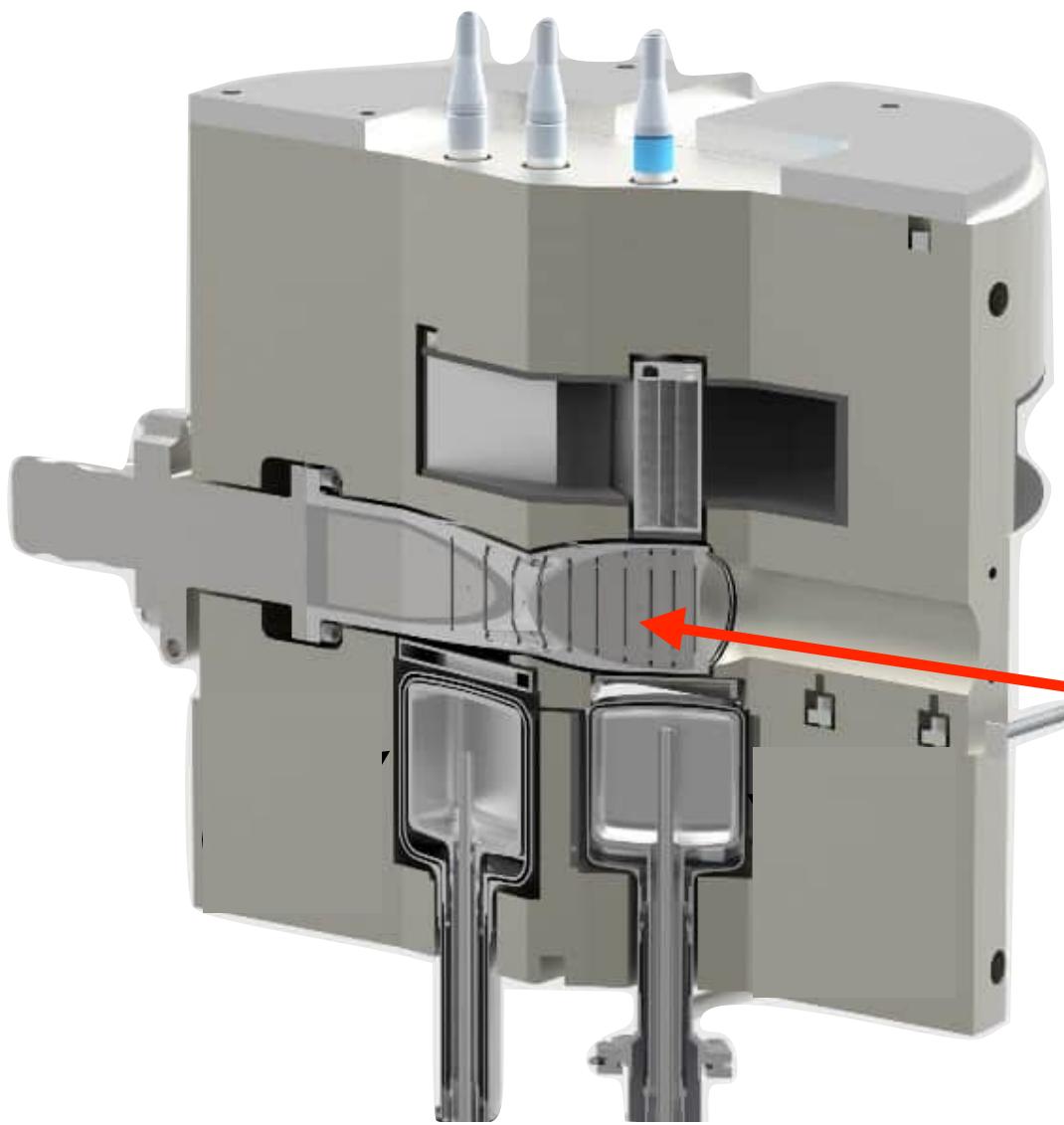
The use of a reflector is very important to enhance the slow-neutron intensity but it can cause broadening of pulses

Idea: decoupling the moderator from the reflector below the “decoupling Energy” by using a sheet of absorber material (B4C, Cd, Ag-In-Cd).  $E_d < 0.4$  eV for B4C,  $E_d > 0.5$  for Cd

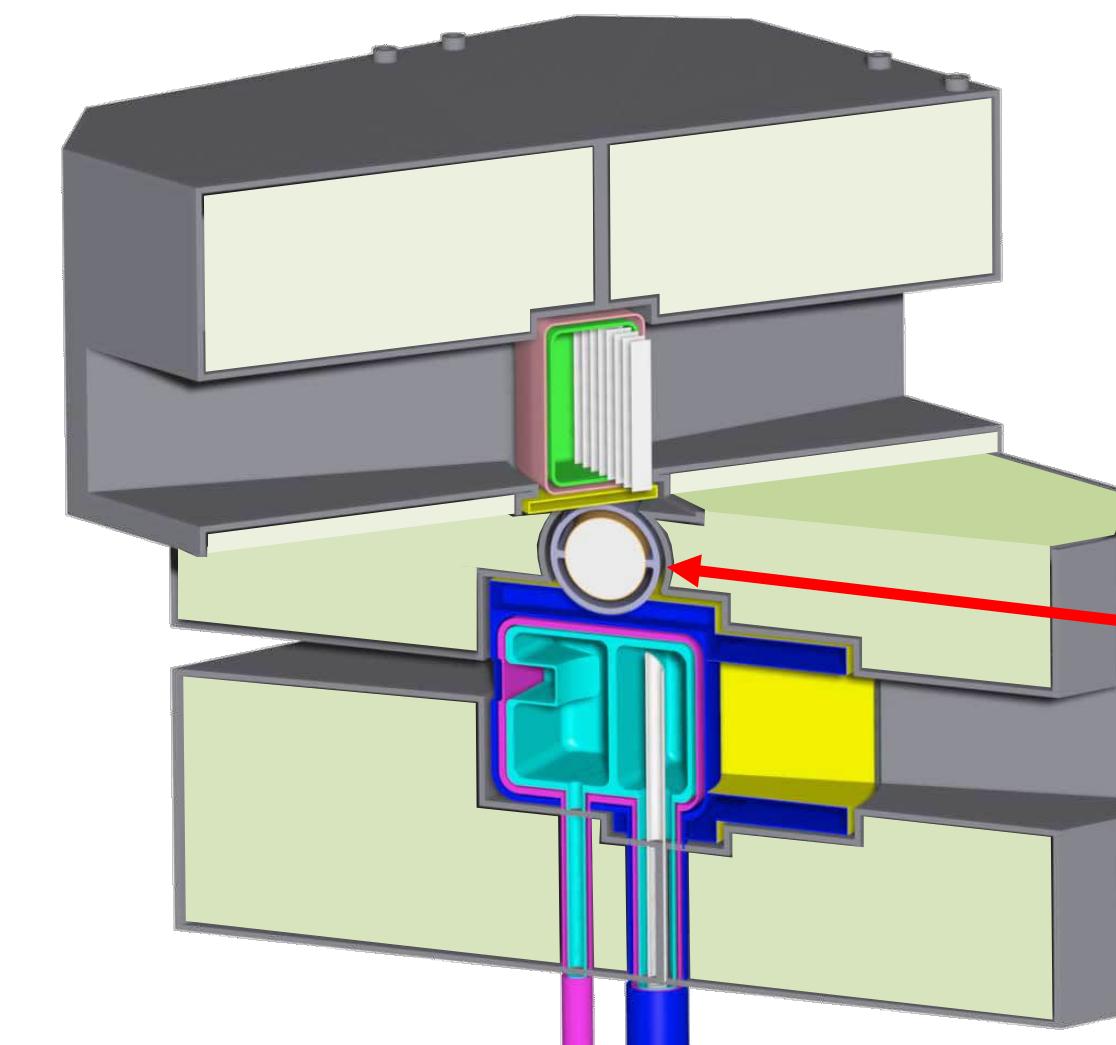
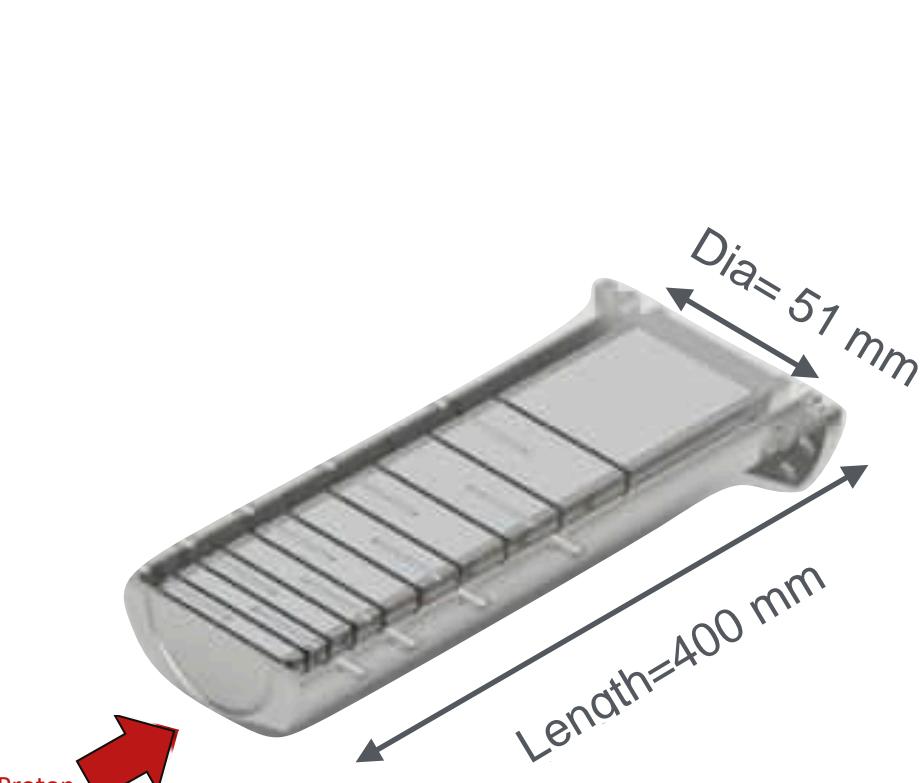
Only neutrons with  $E > E_d$  can pass through

On the contrary in case of maximum time average flux requirement (application for which intensity of the flux is more important than resolution) —> minimise neutron absorption in reflector and moderator (D2O preferred to H2O) and so moderator is fully coupled to reflector

# ISIS TS1 and TS2 TRAM



Target



TS2 TRAM (40 kW):

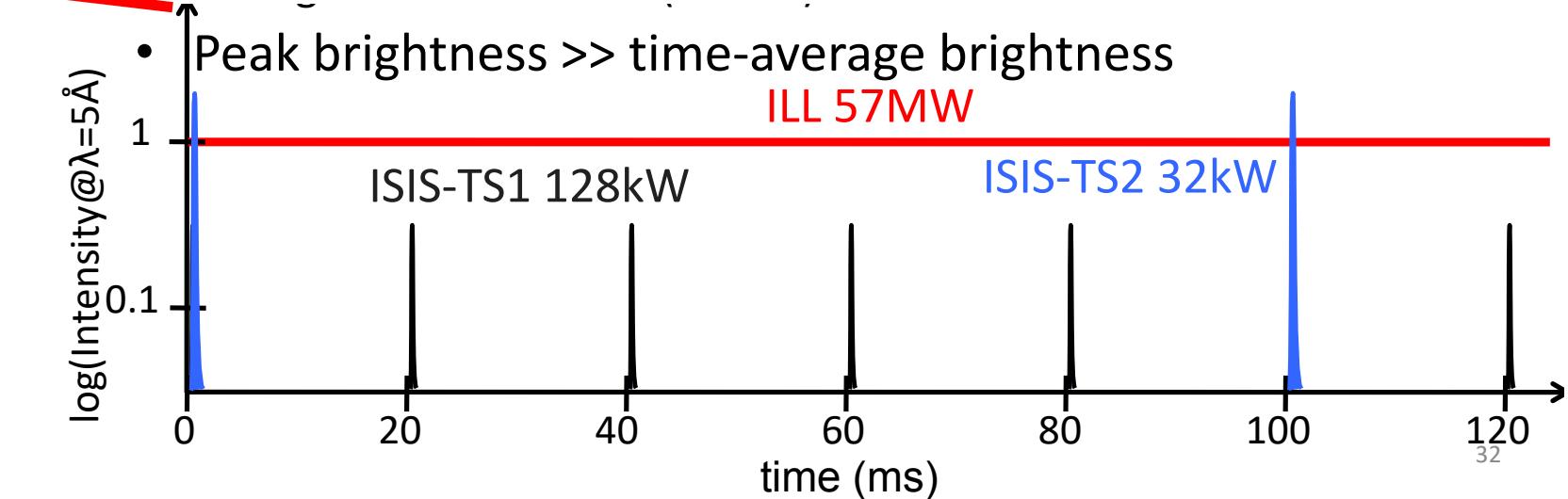
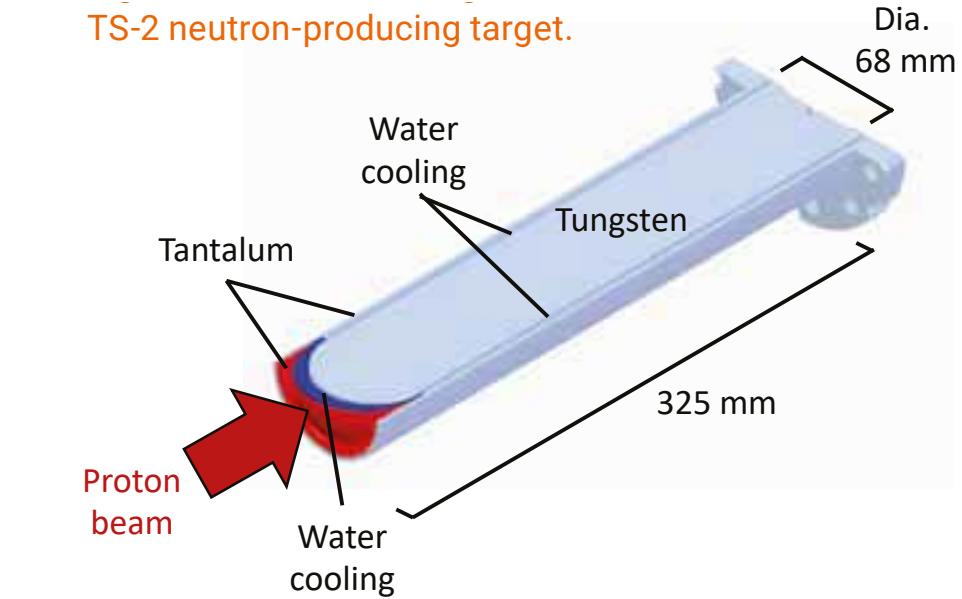
TS2 was developed to produce bright beams of low-energy neutrons enabling the ISIS science programme to expand in the key research areas of soft matter, advanced materials and bio-science

target → bulk W cylinder D= 68 mm, L=325 mm , Ta Clad.

above → coupled moderator with two viewed surfaces: LH<sub>2</sub> and grooved S-CH<sub>4</sub> (tapered hole to focus the beam and increase the brightness)

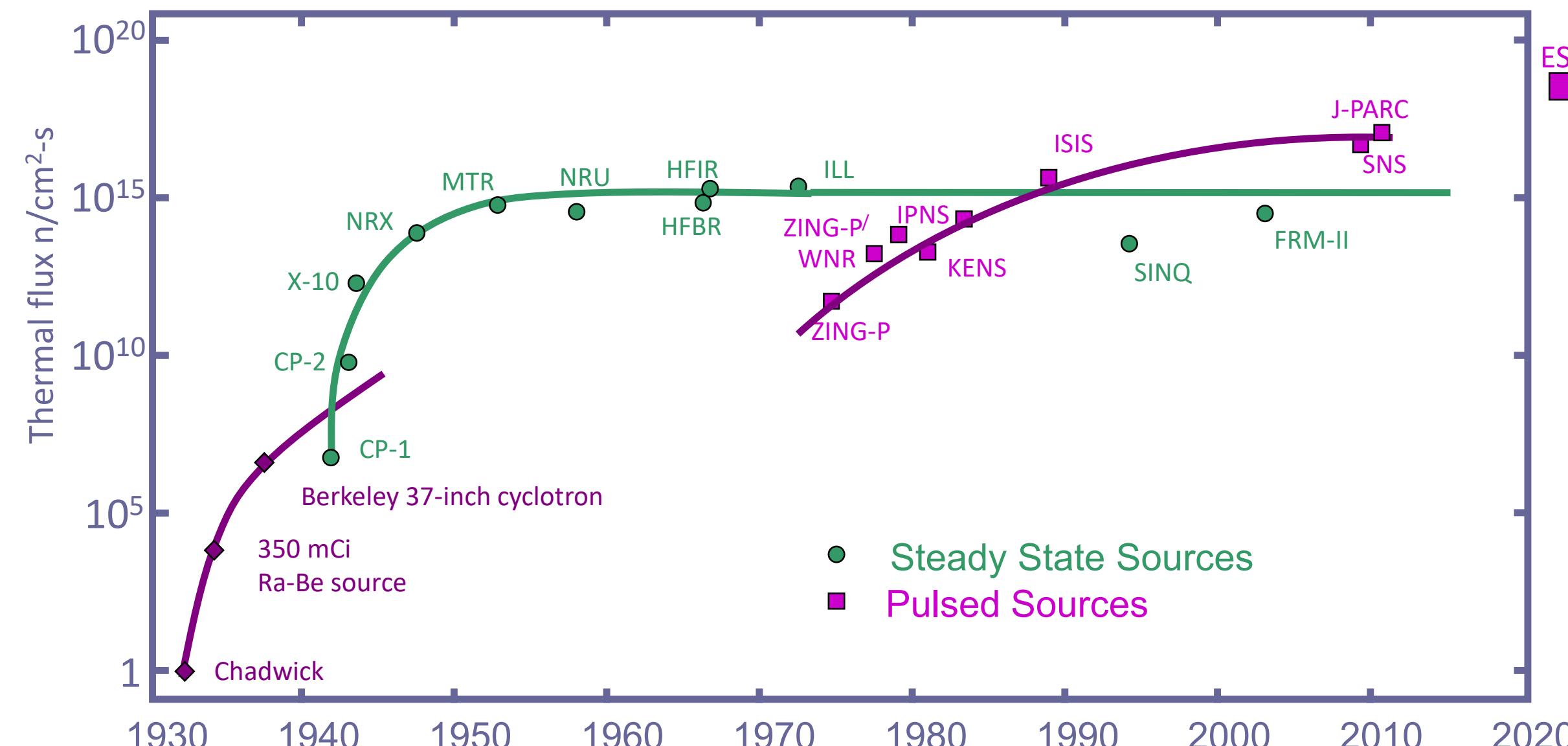
below → decoupled moderator LH<sub>2</sub> (for shorter pulses)

reflector → solid beryllium blocks plated in Nickel



# **Pulsed vs Continuous Neutron Source**

# Development of Neutron Science Facilities

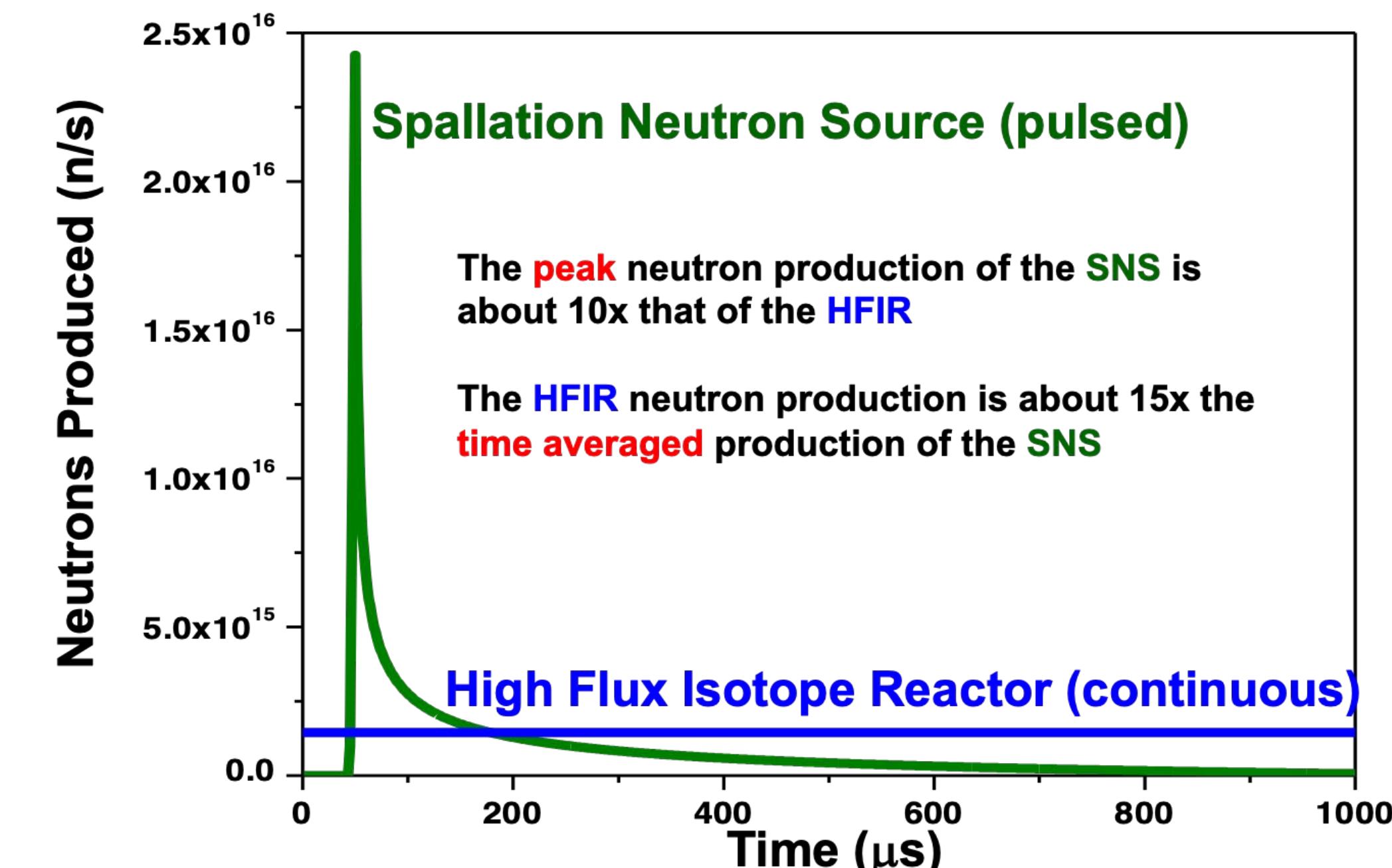
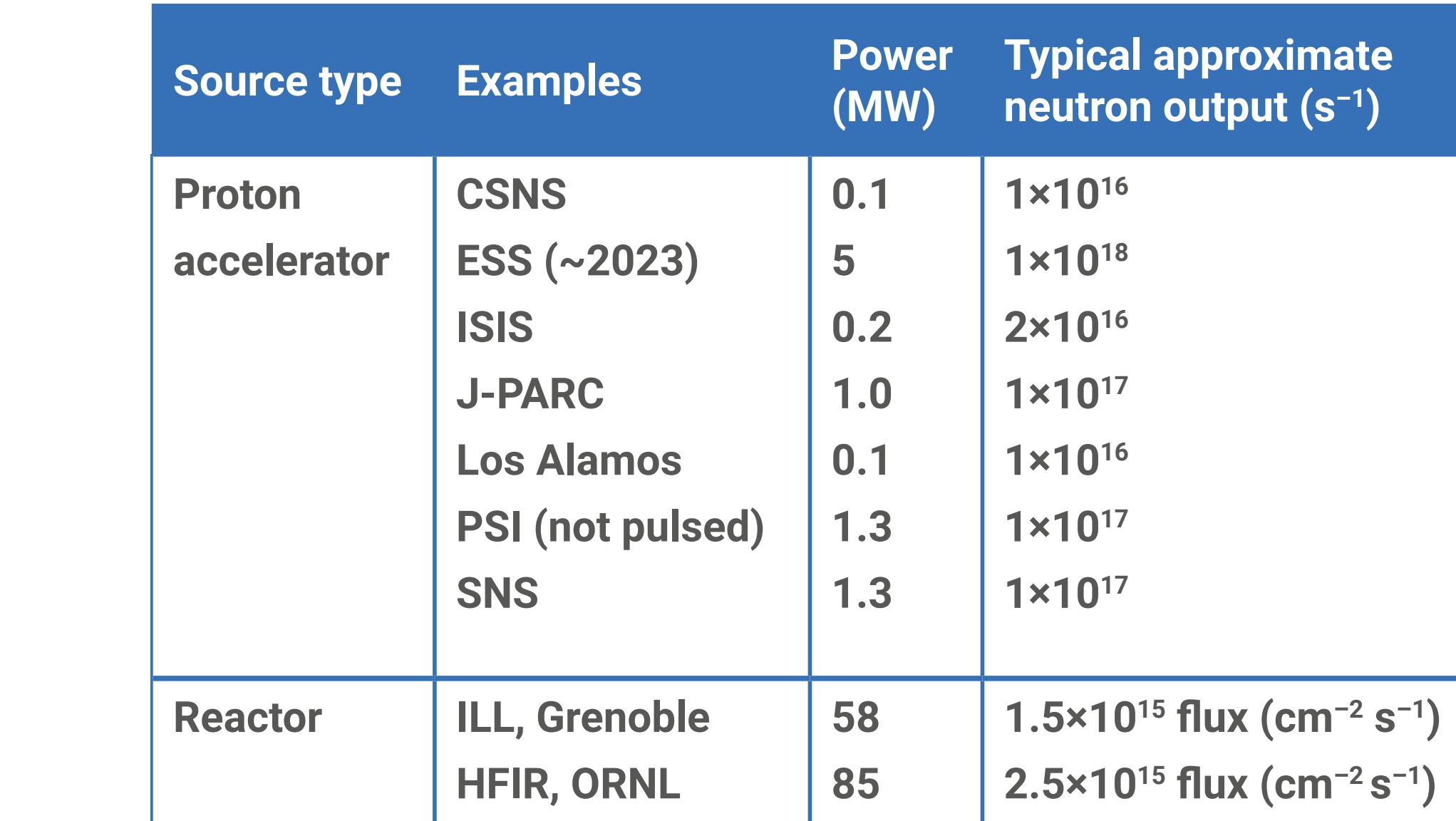


Reactors: great contribution to neutron scattering research, but the highest neutron flux available saturated since the early 1970s.

Pulsed spallation neutron sources (PSNS) came out in the late 1970s. The available flux is gradually getting higher with developments of high-intensity proton accelerators.

The neutron flux at the pulse-peak from PSNS could be 1–2 orders of magnitude higher than that from a high-flux reactor.

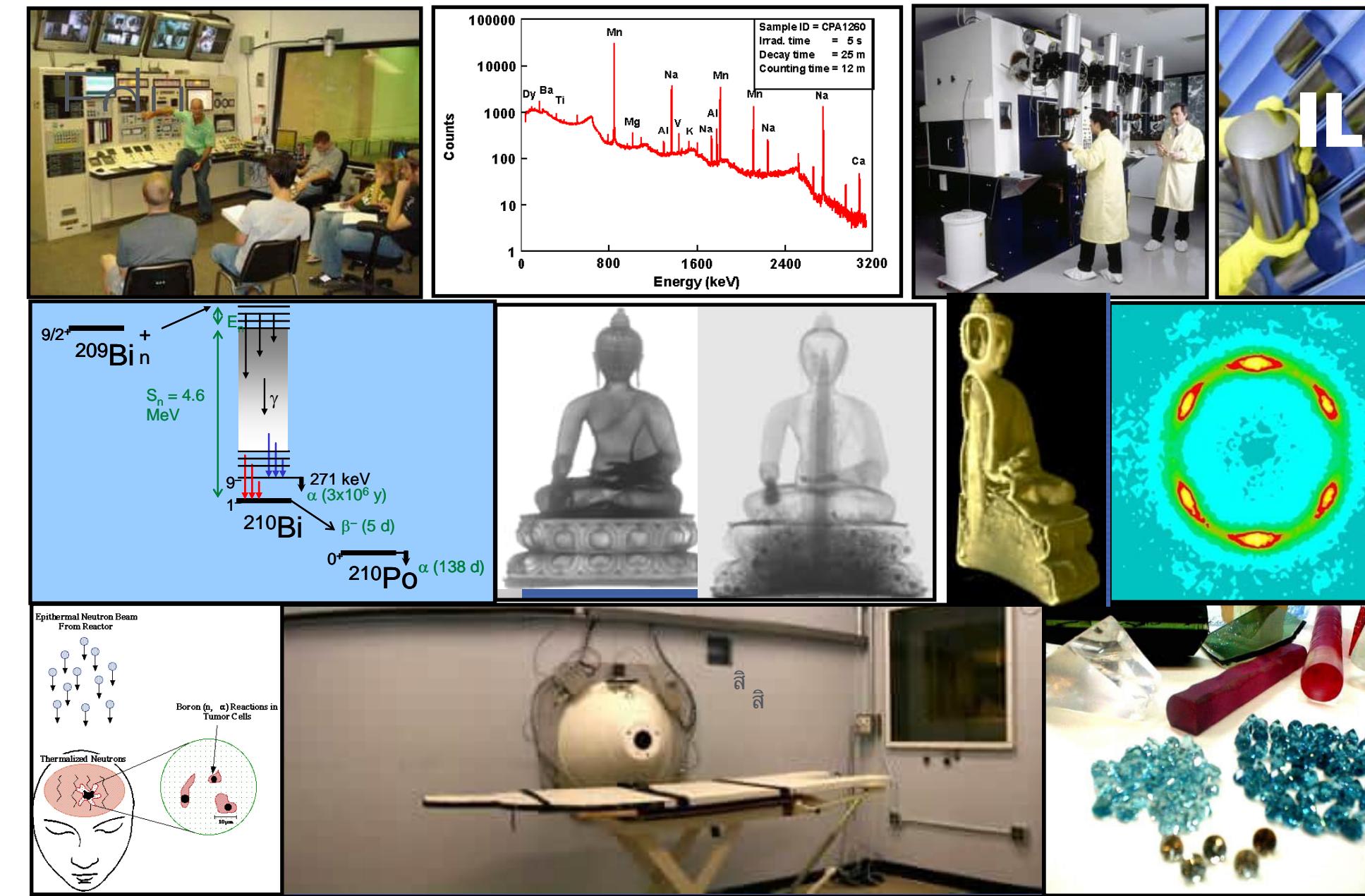
In new generation of long pulse sources (ESS 5MW), even the time-averaged flux is approaching a comparable level



# Large Scale Neutron Sources: Research Reactor Worldwide

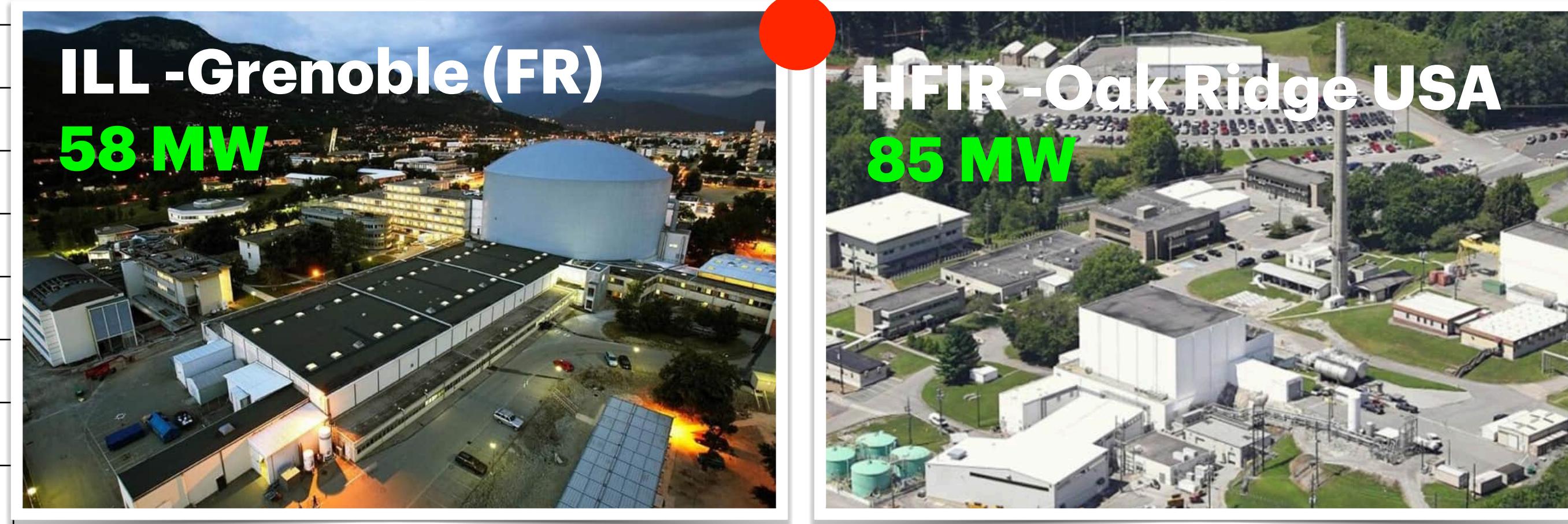
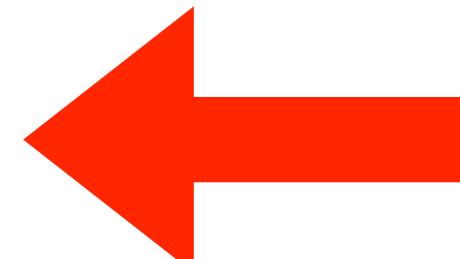
About 220 RRs operational in 53 countries . They are basically neutron factories

Country research	Operational research
Russia	52
USA	50
China	16
India	7
Argentina	5
Canada	5
Germany	5
Italy	5
Brazil	4
Iran	4
Kazakhstan	4
Belarus	3
Belgium	3
Czech Republic	3
France	3
Indonesia	3
Japan	3
Ukraine	3
Others	45
Total	223



## Main application fields

- Education & Training
- Neutron Activation Analysis & PGNAA
- Radioisotope production
- Neutron radiography
- Material/fuel testing/irradiations
- **Neutron scattering**
- Nuclear Data Measurements
- Si doping
- many others



Most powerful RRs for neutron scattering experiments:

ILL (France):  
 $1.5 \cdot 10^{15} \text{ n/cm}^2/\text{s}$

HIFR(USA):  
 $2.5 \cdot 10^{15} \text{ n/cm}^2/\text{s}$

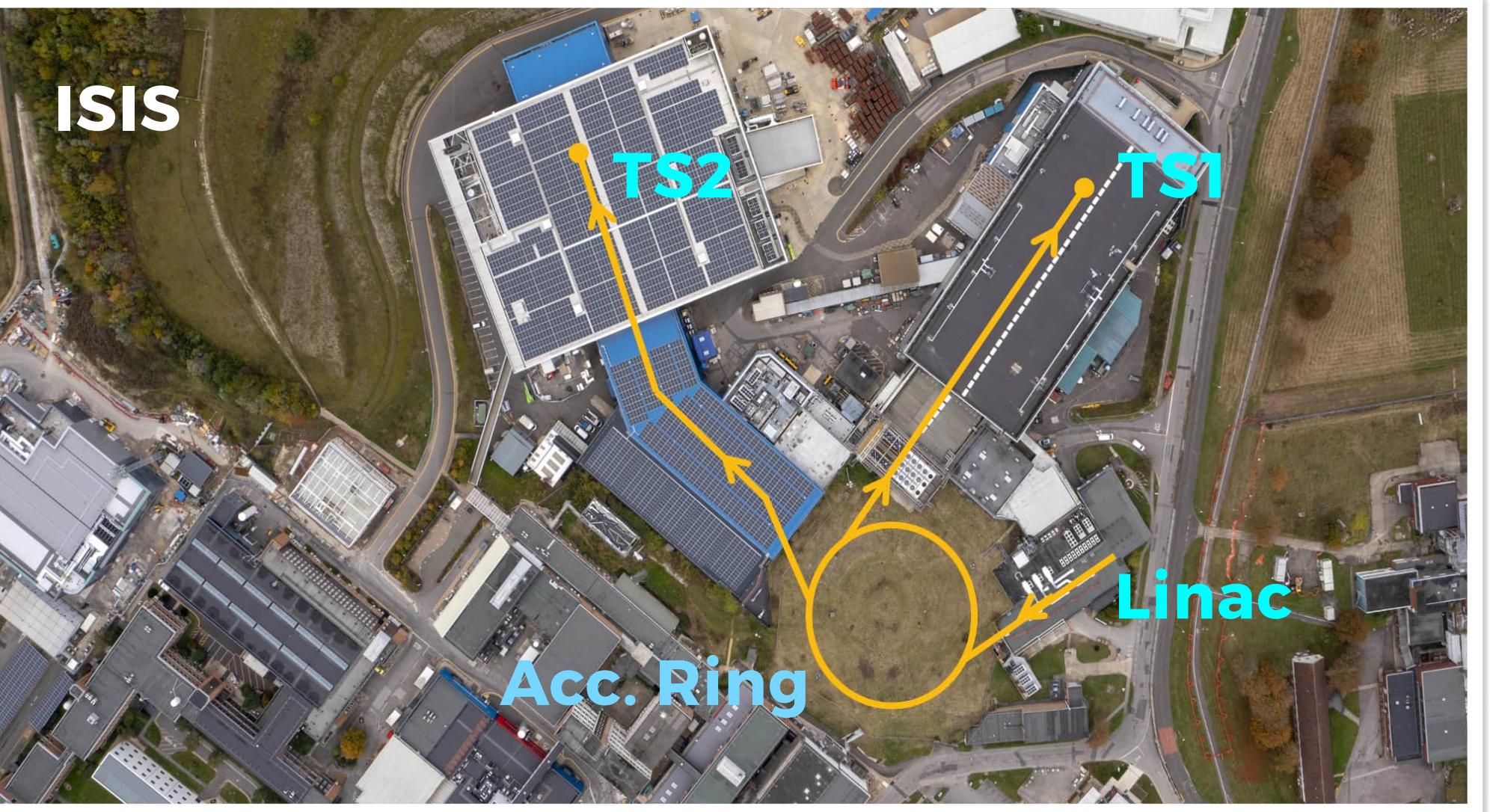
# Large Scale Accelerator-Driven Neutron Sources Worldwide

By Spallation

- **SNS** (Spallation Neutron Source) at **Oak Ridge, TN, USA**
- **ISIS** (Spallation Neutron Source) at **STFC, Harwell, UK**
- **n-TOF** (Spallation Neutron Source) at **CERN, Geneva**
- **IPNS** (Intense Pulsed Neutron Source) at **Argonne National Lab, IL, USA**
- **LANSCE** (Spallation Neutron Source) at **Los Alamos, NM, US**
- **LENS** (Low energy Neutron Source) at **Indiana University, IN, USA**
- **SINQ** Paul Scherrer Institut (PSI) in **Villigen PSI, Switzerland**
- **ESS** (European Spallation Source) in **Lund, Sweden**

By high energy e-

- **Gelina** (Geel Electron Linear Acceleration) in **Belgium**
- **n-ELBE** in **Dresden, Germany**
- **PNF** (Pohang Neutron Facility) in **Korea**
- **ORELA** at **Oak Ridge, TN, USA**



# Large Scale Accelerator-Driven Neutron Sources Worldwide

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## ISIS

**Accelerator:**  
Energy 800 MeV  
Rep.Rate =40 Hz  
Current =200  $\mu$ A  
Pulse length=120 ns

**TS1:**

Water (80%D<sub>2</sub>O-20%H<sub>2</sub>O) cooled W  
Avg. Power 160 kW  
Neutron beam ports=18 +5(mu)

**Accelerator:**  
Energy 800 MeV  
Rep.Rate =10 Hz  
Current =180  $\mu$ A  
Pulse length=120 ns

**TS2:**

Heavy Water cooled W  
Avg. Power 32kW  
Neutron beam ports=11

## ESS

**Accelerator:**  
Prot. Energy 2 GeV  
Rep.Rate =14 Hz  
Avg Current =62.5 mA  
Pulse length=2.86 ms

### Target station:

He-gas cooled rotating W-target  
Avg. Power 5 MW  
**Pulse power:** 125 MW  
Neutron beam ports=42

## SNS

**Accelerator:**  
Prot. Energy 1.0 GeV  
Rep.Rate =60 Hz  
Avg Current =26 mA  
Pulse length=0.7 ms

### Target station:

material Liquid Hg  
Avg. Power 1.4 MW  
Neutron beam ports=24

# Pulsed and Continuous Neutron Sources

**Pulsed sources-> neutrons are provided in bursts with a given repetition rate. These source are mainly accelerator driven sources**

- More High energy neutrons
- Measure when off (low backgrounds)
- Higher peak power possible
- Sharp pulses-high resolution
- Time of flight bandwidth controlled by source frequency
- Broader range of energy/wavelength
- Lower environmental footprint

**Continuous--> constant flux of neutrons over the running time. Mainly reactors**

- More Low energy neutrons (since the higher moderator efficiency)
- Easier to shield (low backgrounds)
- High average flux
- Sequence of choppers can be used to cut the continuous beam in pulses
- Experimental flexibility
- Environmental footprint

# C vs P: Technical and Applicative Differences

The attraction of pulsed source is obvious: intense burst of neutrons is produced in a time which is short compared with the period between pulses and so the heat generated during pulses can be efficiently removed.

## Accelerator main differences

- ✓ Synchrotrons or accumulator (compressor) rings provide short neutron pulses (ISIS,SNS)
- ✓ Linear accelerators provide long neutron pulse (ESS). Long pulses do not require pulse compressing ring
- ✓ Cyclotrons provide continuous beams of protons—> neutrons (PSI)

Spallation source	Short pulse	Long pulse	Continuous
Imp. Length	a few $\mu$ s	a few ms	-
Facilities	ISIS, SNS, JPARC	ESS	PSI

**Continuous source** favour those applications that require high time average flux (much lower fast and very high energy neutron background)—>**PSI**

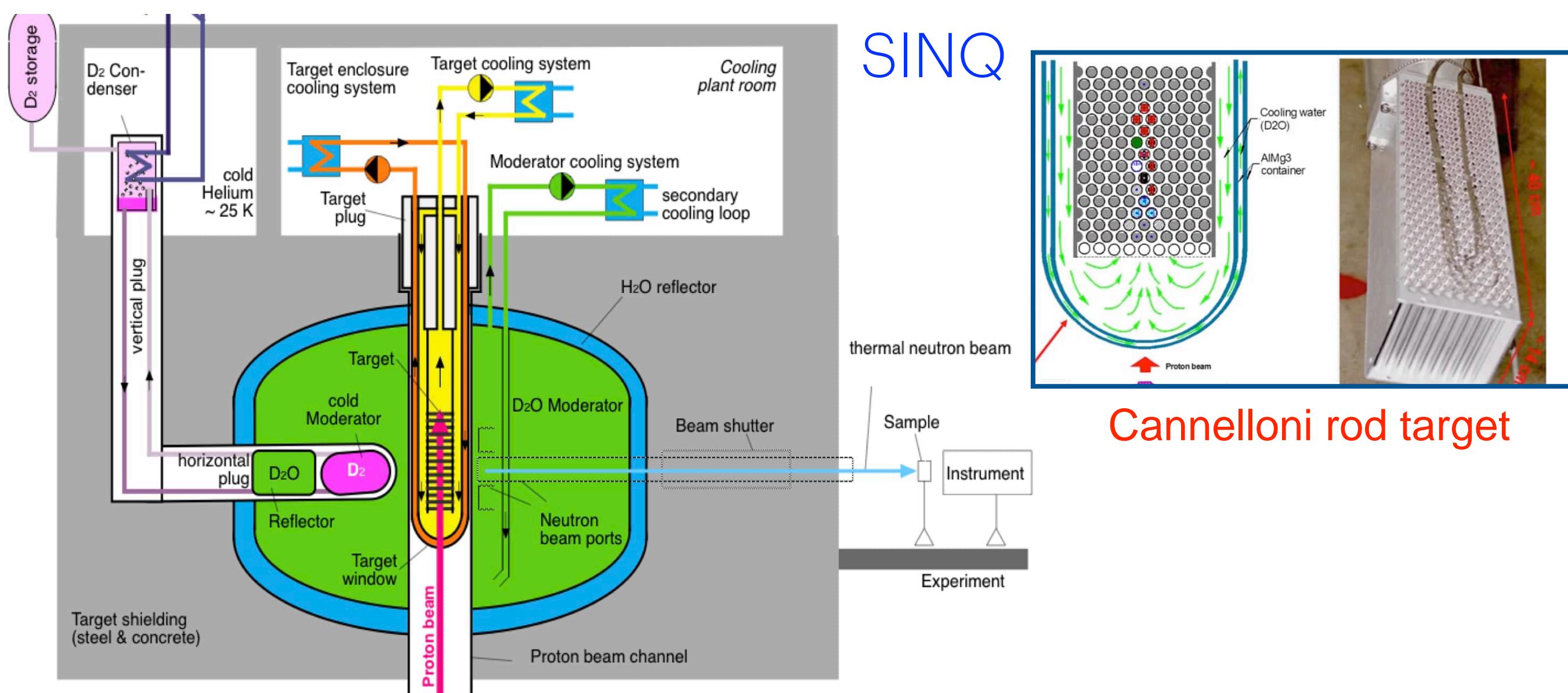
**Long Pulse** length of proton beam longer than time to slowdown the neutrons in the moderator. They are more suitable for high intense cold neutron beams or for application requiring high resolution with low wavelength, such as small-angle scattering —>**ESS**

**Short pulses** are required for epithermal neutron spectrum or for highest resolution in thermal region (serve better those applications with high peak flux at appropriate rate—> **ISIS**)

Reactor	Long pulse
Imp. Length	a few tens of ms
Facilities	IBR-2 DUBNA

Long pulse do not suffer from thermal shock one has to deal with in the targets of high power short pulse source

# Remarkable Exceptions: “Pulsed Reactor” and “Continuous Spallation Source”



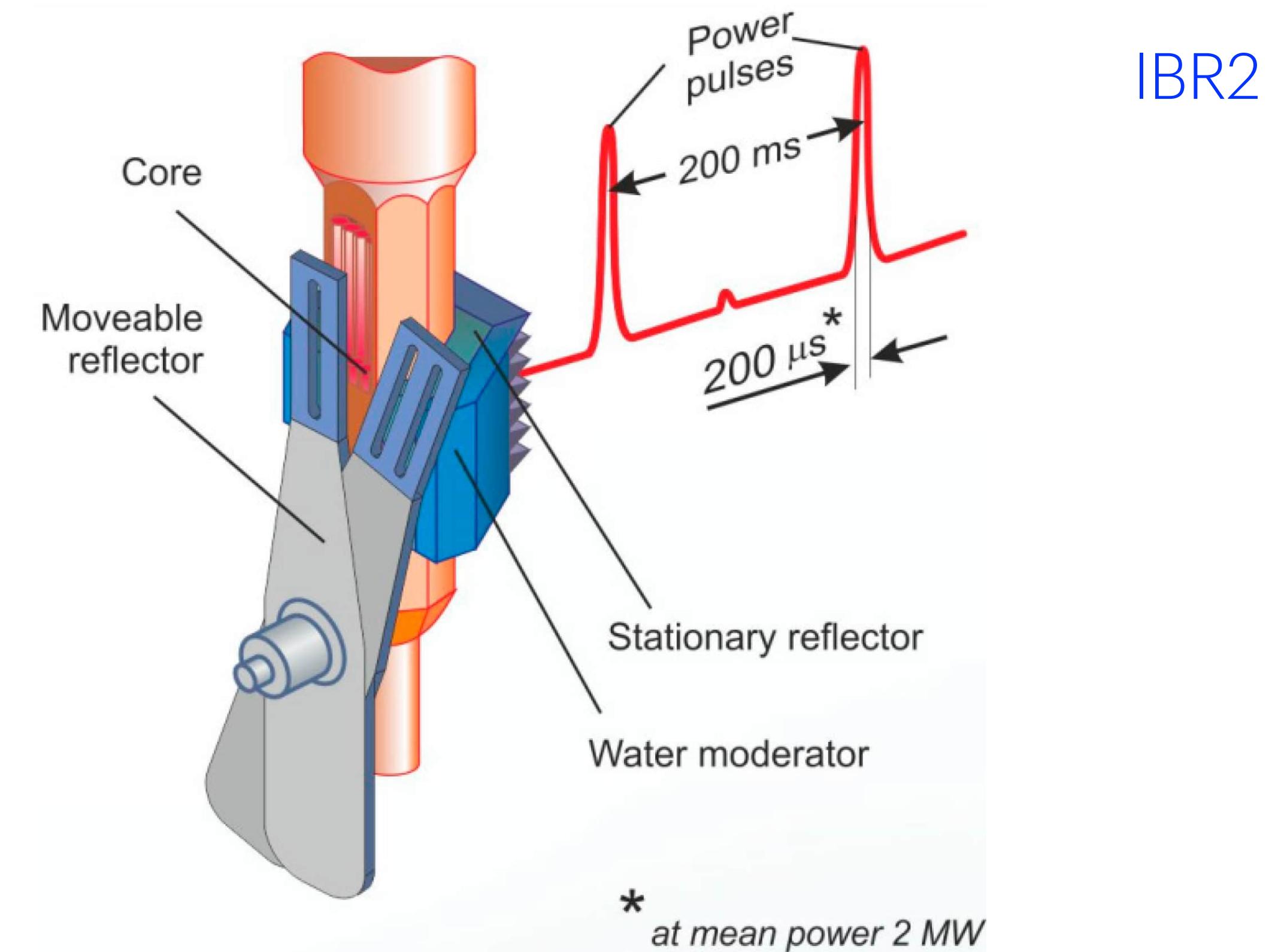
**Figure 1:** Schematic layout of the SINQ cooling and moderation systems

590 MeV proton and a current of approximately 1.2 mA. The beam power available for SINQ will thus be roughly 0.75 MW.

The proton beam is diverted by bending magnets vertically upwards from underneath into the heavy metal target (on Pb/Zr)

Spallation continuous source - the first and only one of its kind in the world - with a flux of about  $10^{14}$  n/cm<sup>2</sup>/s.

Beside thermal neutrons, a cold moderator of liquid deuterium (cold source) slows neutrons down and shifts their spectrum to lower energies.



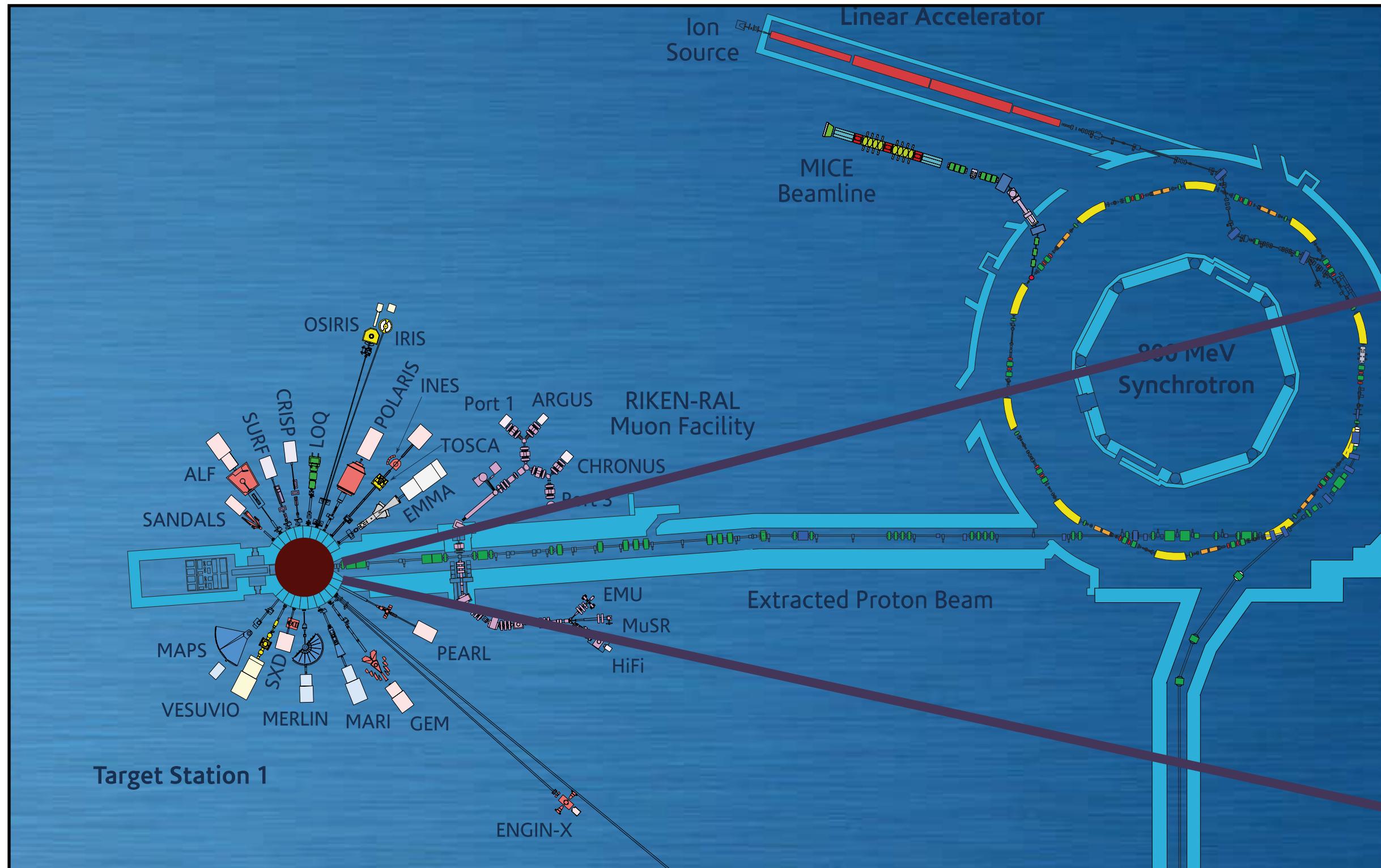
Pulsed fast reactor, IBR-2,

In this type of reactors neutron pulses with widths of several tenths of microseconds are generated periodically by mechanical modulation of reactivity.

The IBR-2 reactor with its unique technical approach produces one of the most intense neutron fluxes at the moderator surface among the world's reactors:  $\sim 10^{16}$  n/cm<sup>2</sup>/s, with a power of 1850 MW in pulse

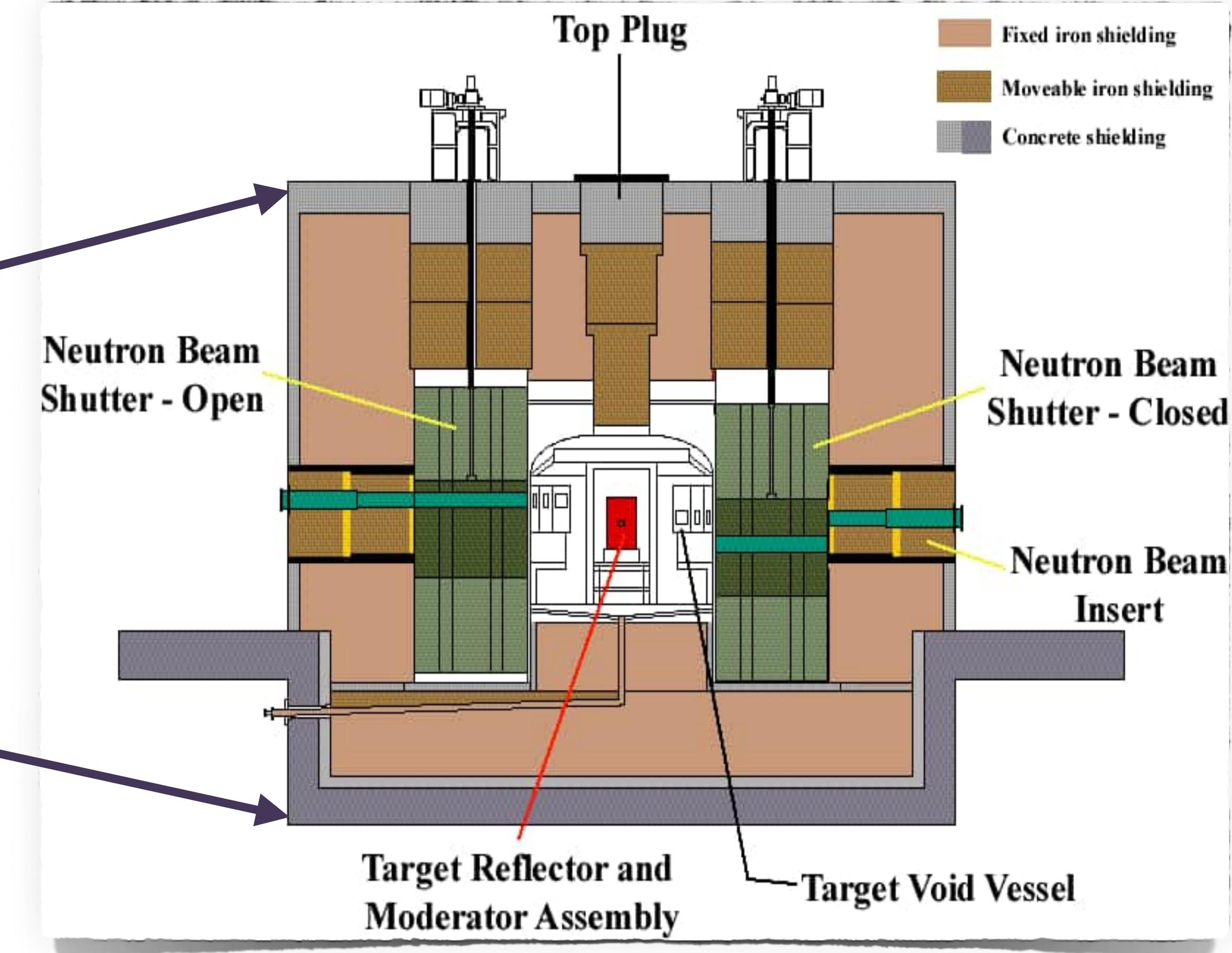
# **Neutron Beam Lines & Time of Flight Technique**

# Shield, Shutter and Instruments Layout (TS1 ISIS)



Instruments	Flight path to sample
Maps	12m
Vesuvio	11 m
Merlin	11.8 m
Mari	11.739 m
Pearl	12.8 m
Ines, Tosca	23.8 m
Polaris	14m

Instruments are located at several tens of meters from moderators. The detailed distances for ISIS TS1 is reported in the table



Shutter = 2 m of iron and concrete  
Dose rate at the shield surface < 0.75 mSV/h  
Take as reference that the dose of 1Ci AmBe at 10 cm from the center is about 0.35mSV/h

# Spallation Sources for Scattering: Main Energy and Time Requirements

In order to use neutron scattering for condensed matter study by TOF:

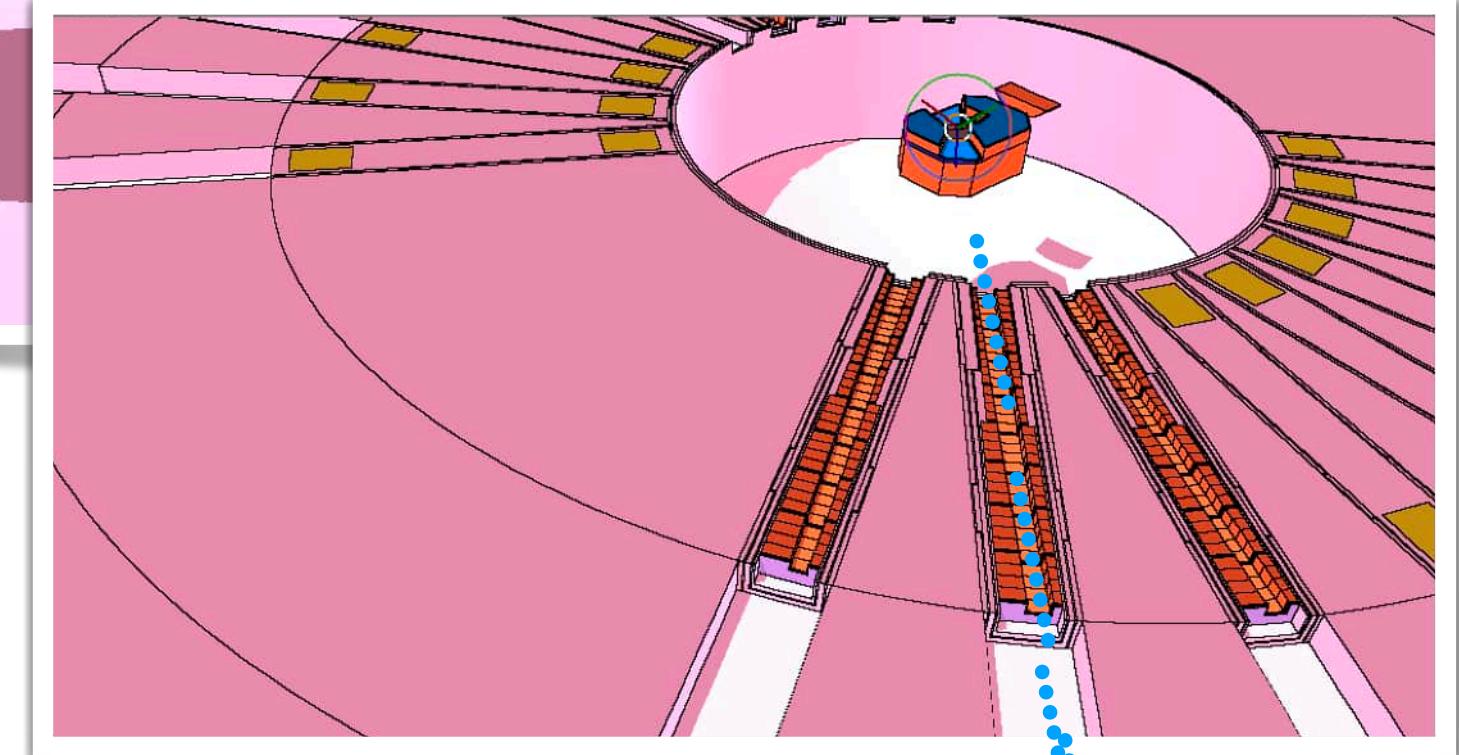
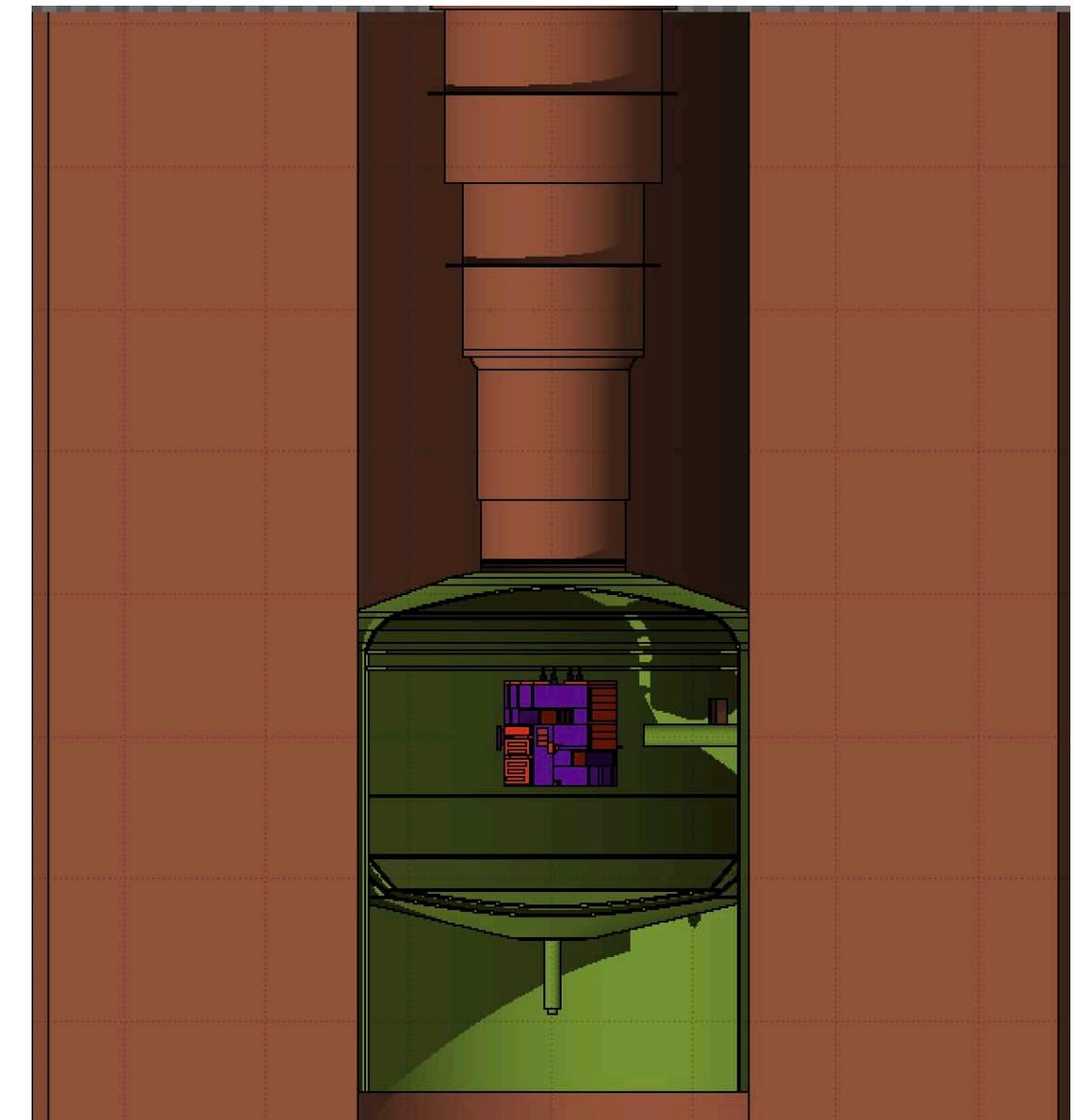
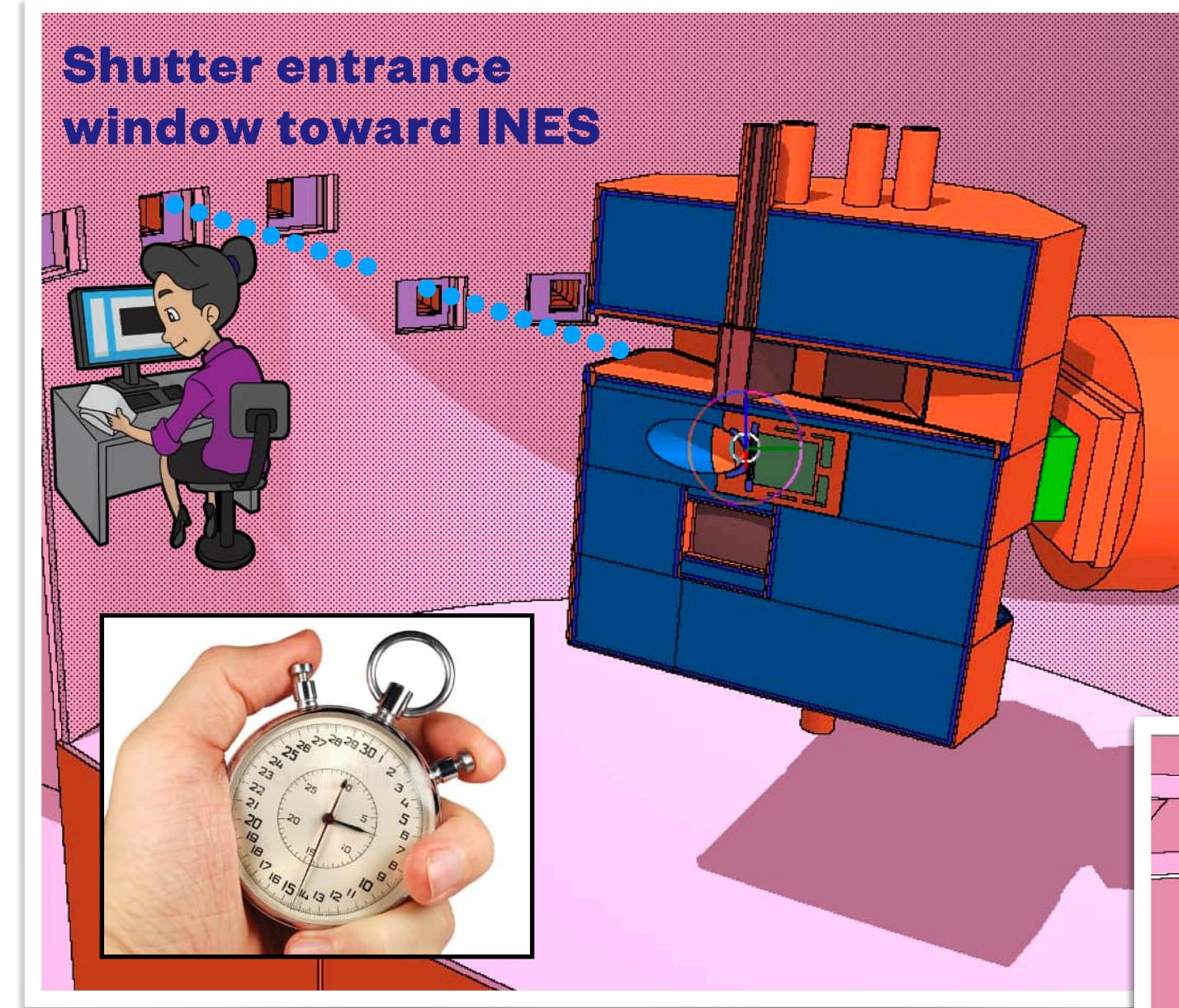
- Neutrons have to be between 10 eV and 0.1 meV (MODERATOR)
- Short beam pulses at the shutter entrance

**Time-of -flight techniques rely on approximation that all neutrons start out down a flight path at the same time.**

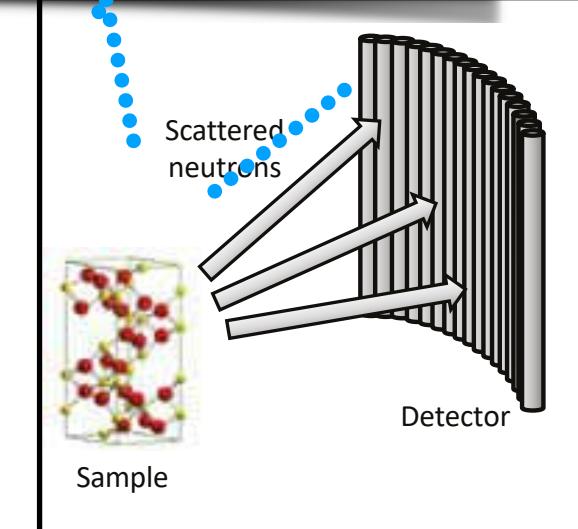
The distribution in emission time must be much smaller than the flight time: order of microseconds rather than tens of milliseconds.

The difference between arrival and departure is the time-of-flight, and knowing the distance (classical approach is applicable), the neutron velocity can be derived. It directly also gives energy and wavelength.

Time distribution of neutrons at the shutter inlet is of the order of a few microseconds



Thermal neutrons take about 10 ms to arrive at the sample, 20 m downstream from neutron moderators

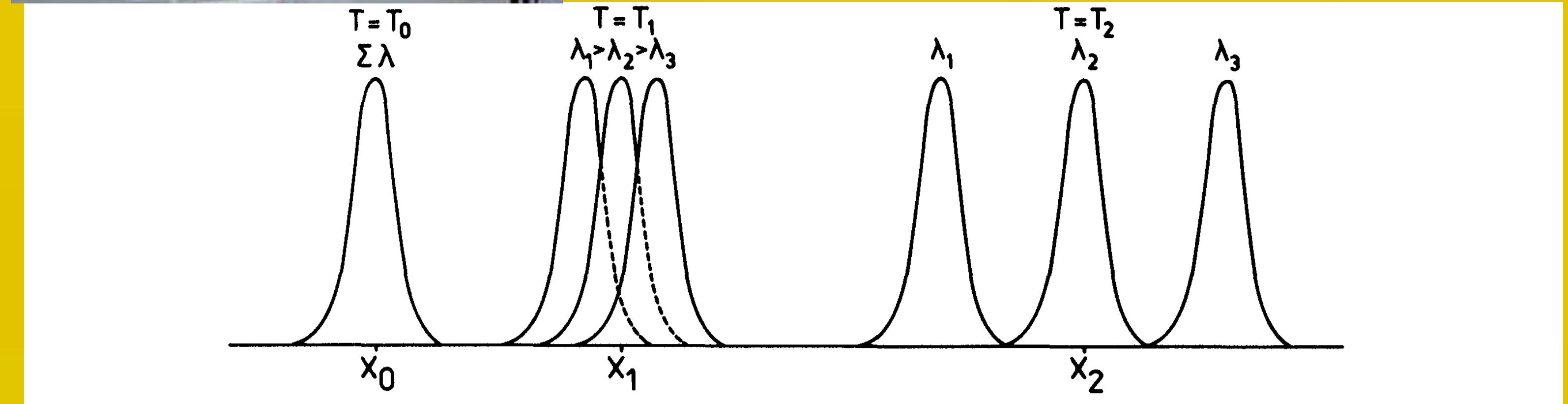


Time-of-flight technique relies on approximation that all neutrons start out down a flight path at the same time.

At time  $t_0$  the pulse leaves the moderator surface. It contains all wavelengths.

After a short time  $t$  the neutrons arrive at the point  $X_1$ . The  $\lambda_3$ -neutrons have passed it already, the  $\lambda_1$  neutrons are slower.

Passing a long flight path  $x_2$  (long time  $t_2$ ) the three different wavelength pulses are separated completely.



## TIME OF FLIGHT NEUTRON SPECTROMETRY

ToF=Difference between arrival and departure time-of-flight  
Knowing the distance, the neutron velocity can be derived.  
It directly also gives energy and wavelength.

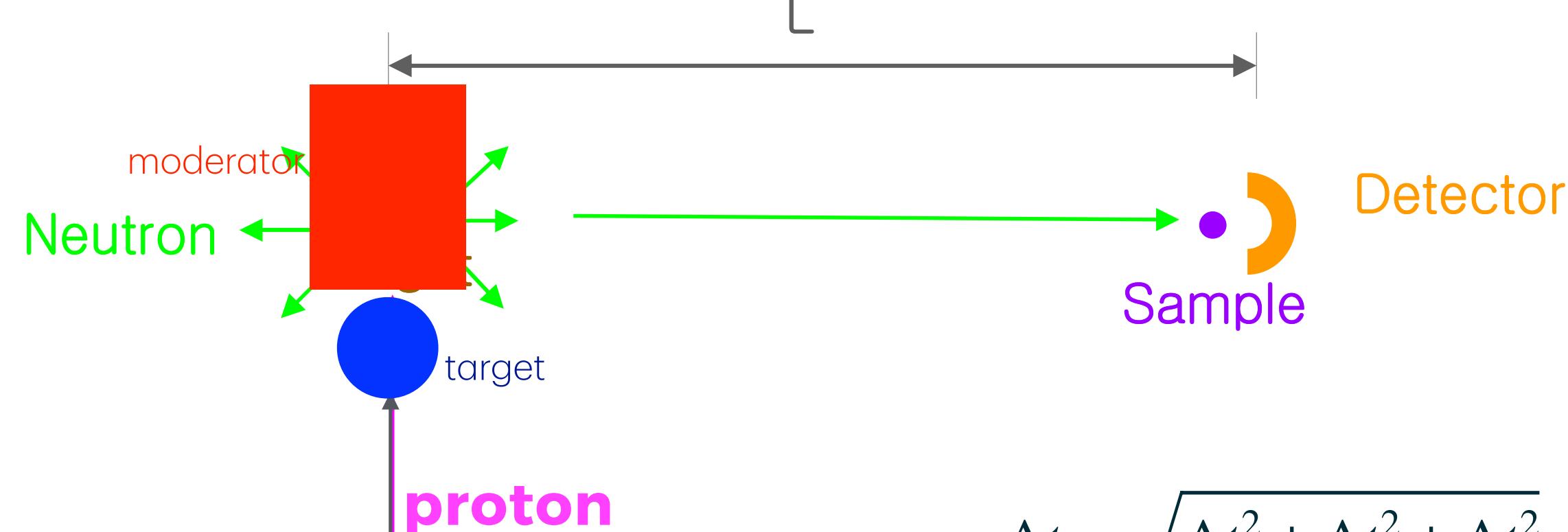
# Neutron Energy Measurements: TOF (Time of Flight)

- Time of flight method: to measure neutron energy of white spectrum
- Time t: Time for a neutron to reach the sample after pulse beam hit the neutron target
- Kinetic energy is determined by measuring t, knowing the fight length (L)

$$E = K \left( \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} - 1 \right)$$

E: Energy  
 K: Rest mass of neutron  
 v: Neutron velocity  
 c: Speed of light

$$\rightarrow E(eV) = \left( 72.3 \times \frac{L(m)}{t(\mu\text{sec})} \right)^2$$



- Energy resolution:

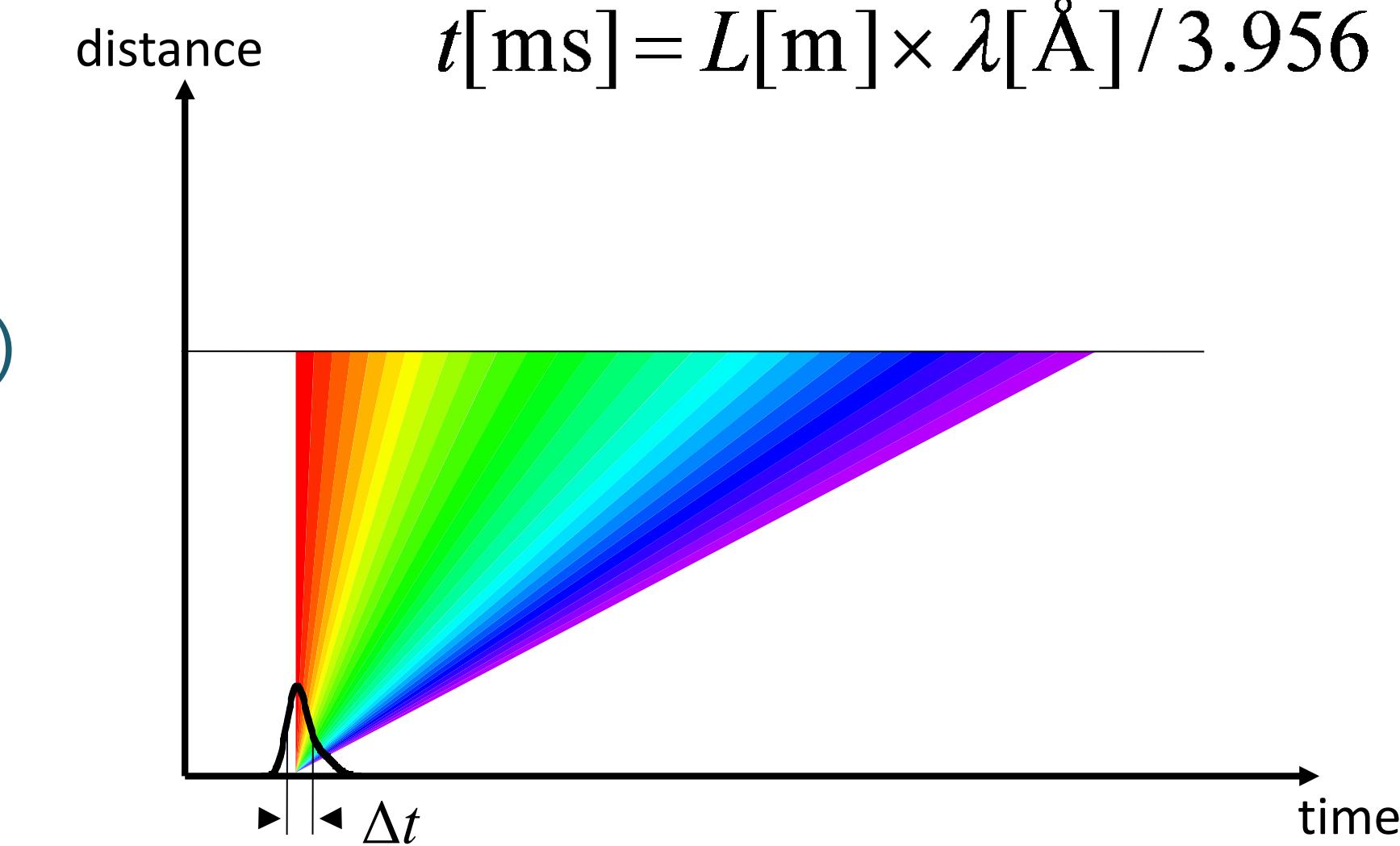
$$\frac{\Delta E}{E} = 2 \frac{\Delta t}{t} = 0.02766 \frac{\Delta t(\mu\text{sec})}{L(m)} \sqrt{E(eV)}$$

$$\Delta t = \sqrt{\Delta t_1^2 + \Delta t_2^2 + \Delta t_3^2}$$

$\Delta t_1 \rightarrow$  Proton Beam Pulse Width

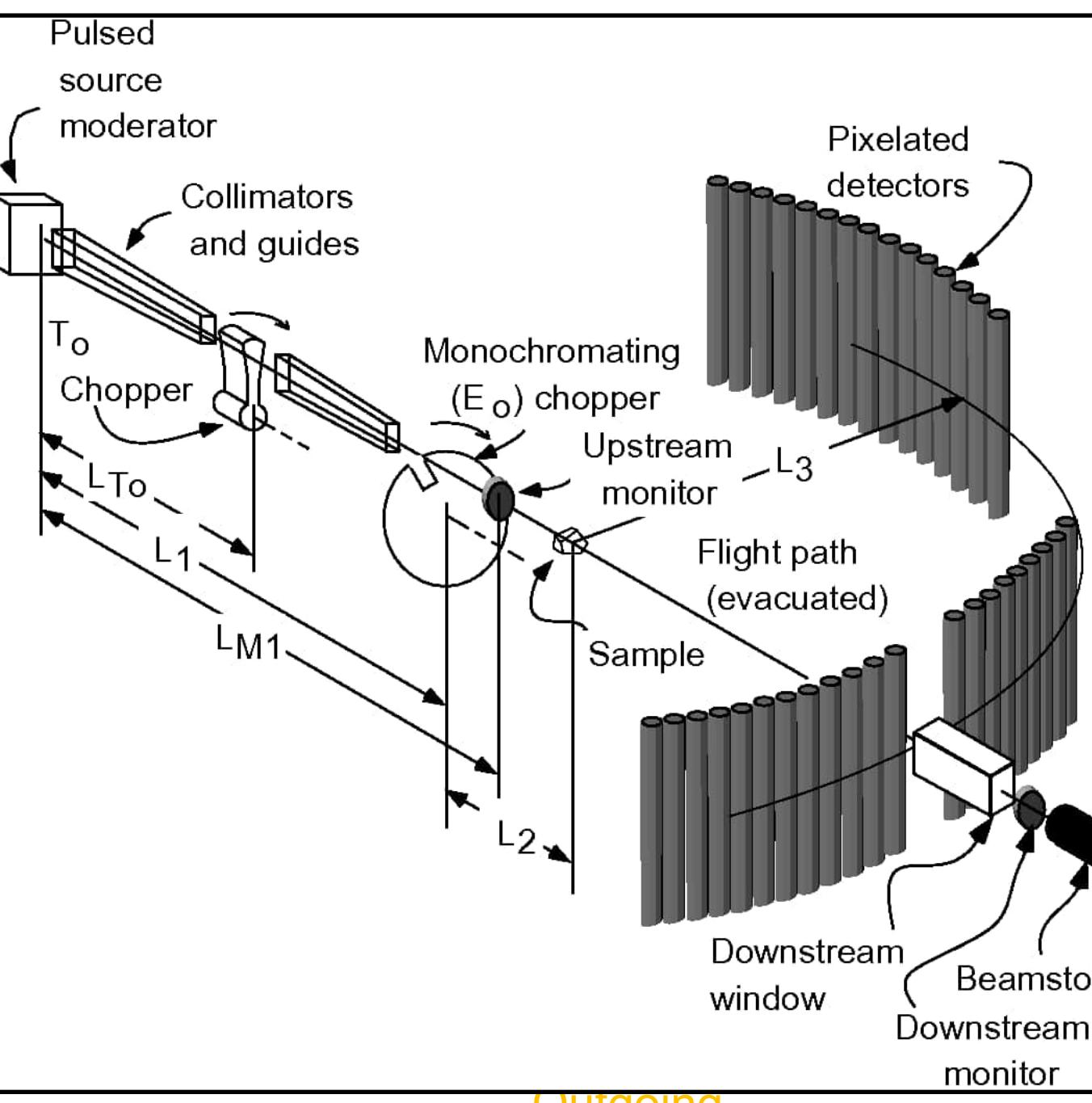
$\Delta t_2 \rightarrow$  Target Moderator Scattering time Width

$\Delta t_3 \rightarrow$  Detector & DAQ response Width

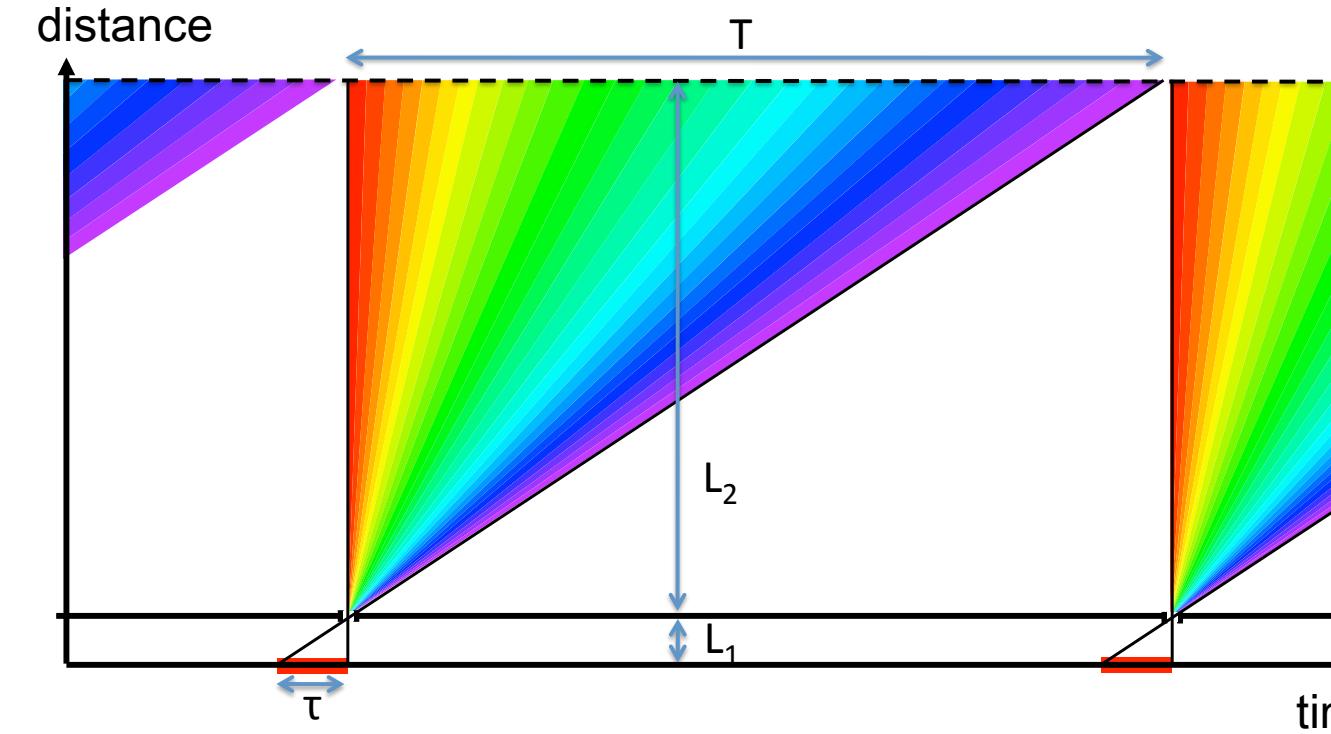


- Better energy resolution: Decreasing  $\Delta t$  or Increasing L
- Increasing L  $\rightarrow$  Decreasing flux, therefore try to decrease  $\Delta t$
- L should be decided by considering the energy resolution and neutron flux required for an instrument

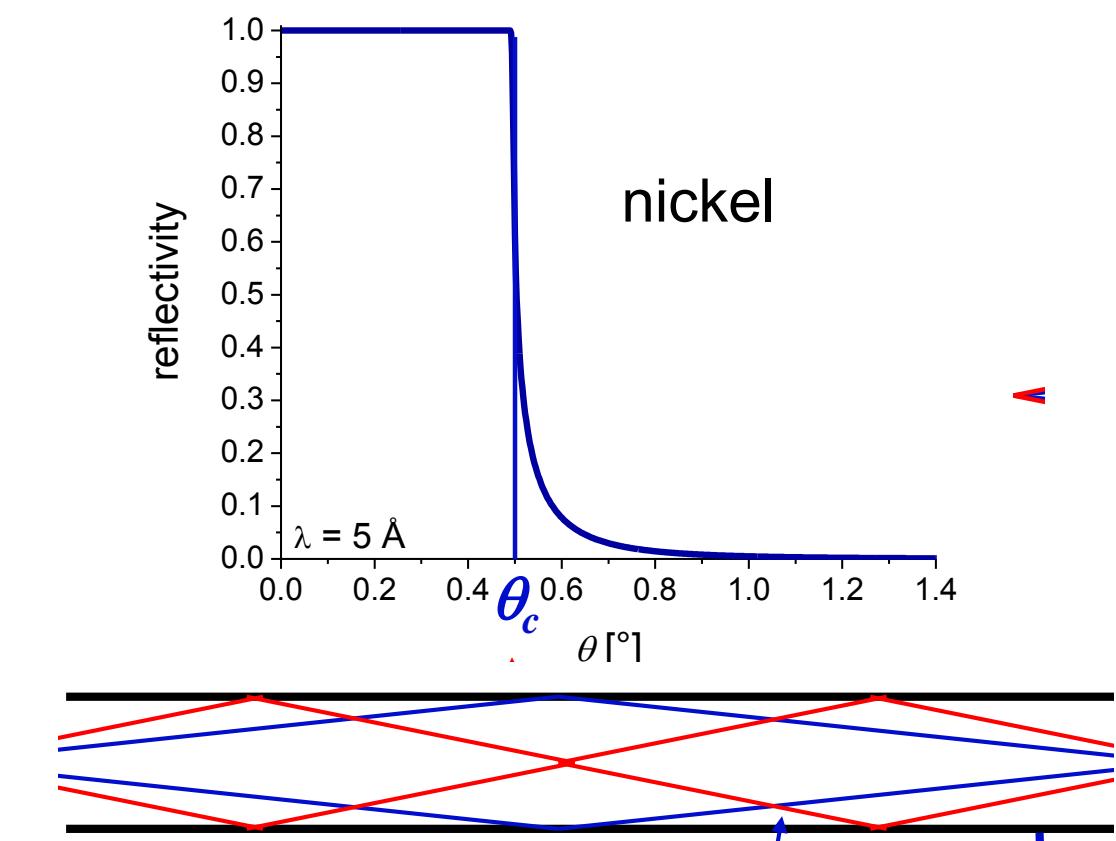
# Typical Components of an Instrument Beam Line



- **Collimators:** define the spatial and angular distributions of neutron beam on the sample, acting independently of wavelength. They are made of apertures trough materials opaque to neutrons.
- **Mirrors and waveguide.** Neutrons reflect perfectly from some surfaces which acts as mirrors provided that the grazing angle of incidence is smaller than the critical angle for the given wavelength (usually made of very flat Ni surface )
- **Chopper (spectrometers):** It is essentially a disc rotated at high speeds with some apertures through which neutrons may only pass during certain periods of the discs rotation. It consists of disk, motor, photosensor, and a housing. The chopper disk is generally made of two stainless steel slices with an absorber material (i.e cadmium) between them. Both steady and pulsed sources have several chopper spectrometers optimised for different applications (i.e high resolution, access the desired Q. etc)
- **Detector**—> neutron detection for getting sample properties measurements . Specific for each instruments. High efficiency and prompt answer are required



use of chopper at distance L1 for long pulse or reactor



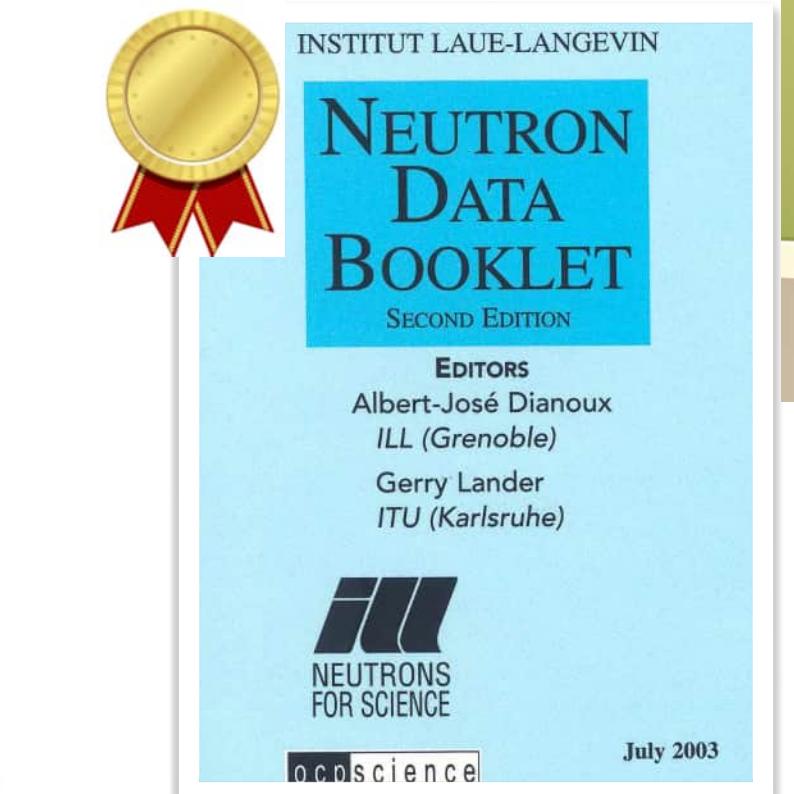
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downloadable for free from the web



## Neutronics of pulsed spallation neutron sources

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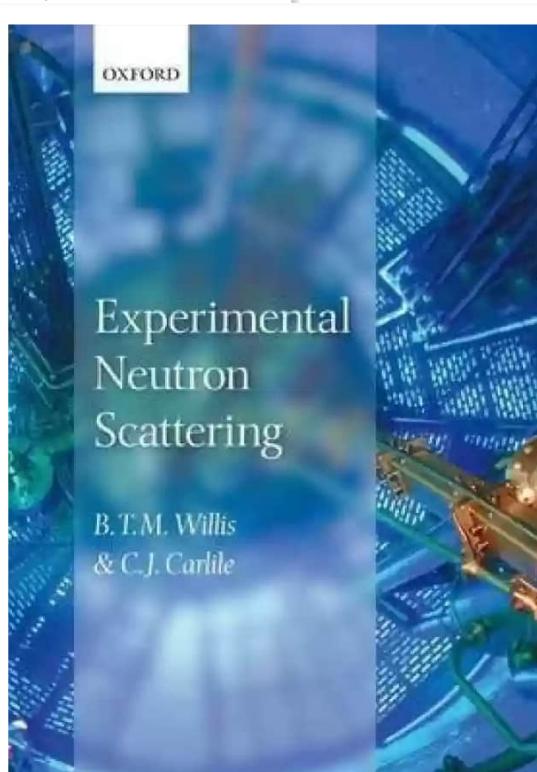
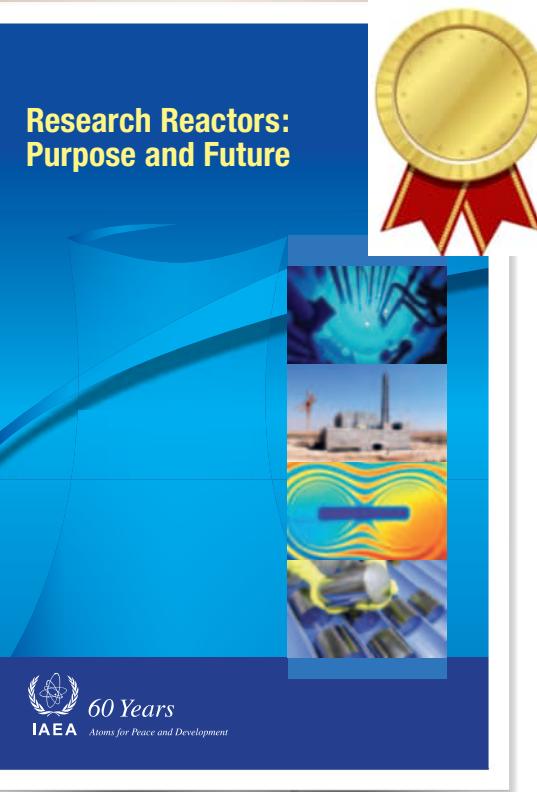
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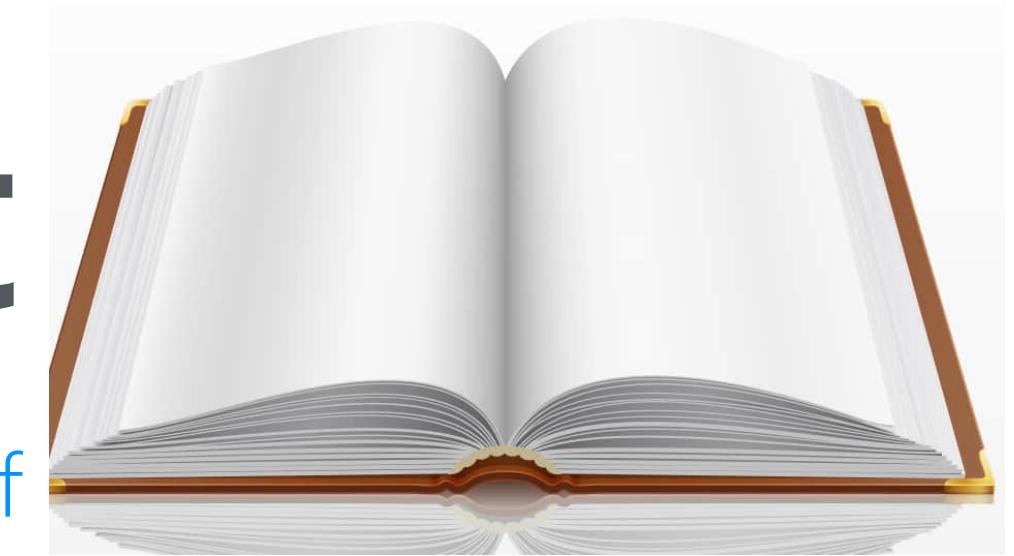
Physics and technology of spallation neutron sources\*

G.S. Bauer\*

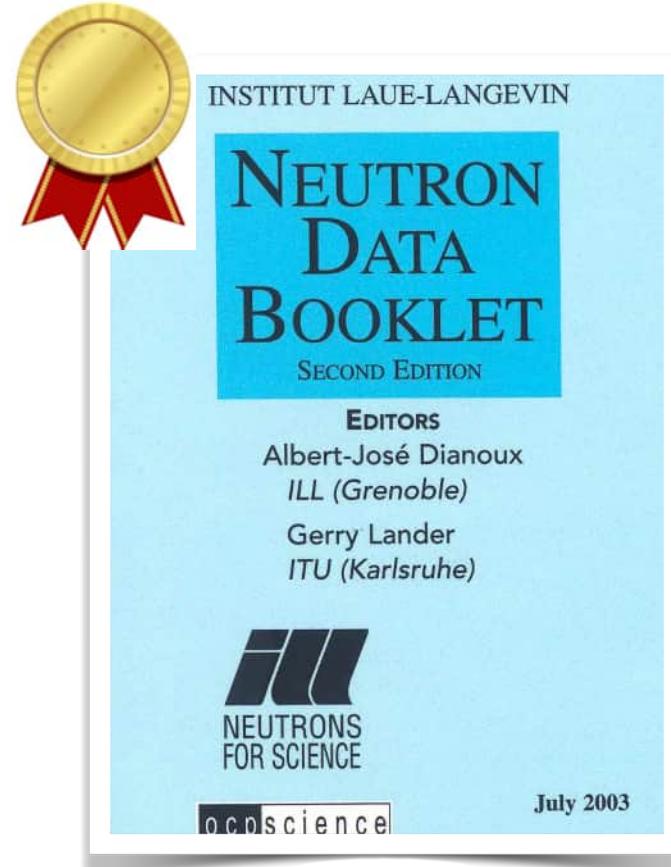
Spallation Neutron Source Division, WHG/A/252, Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland



# A Mention for Neutron D booklet



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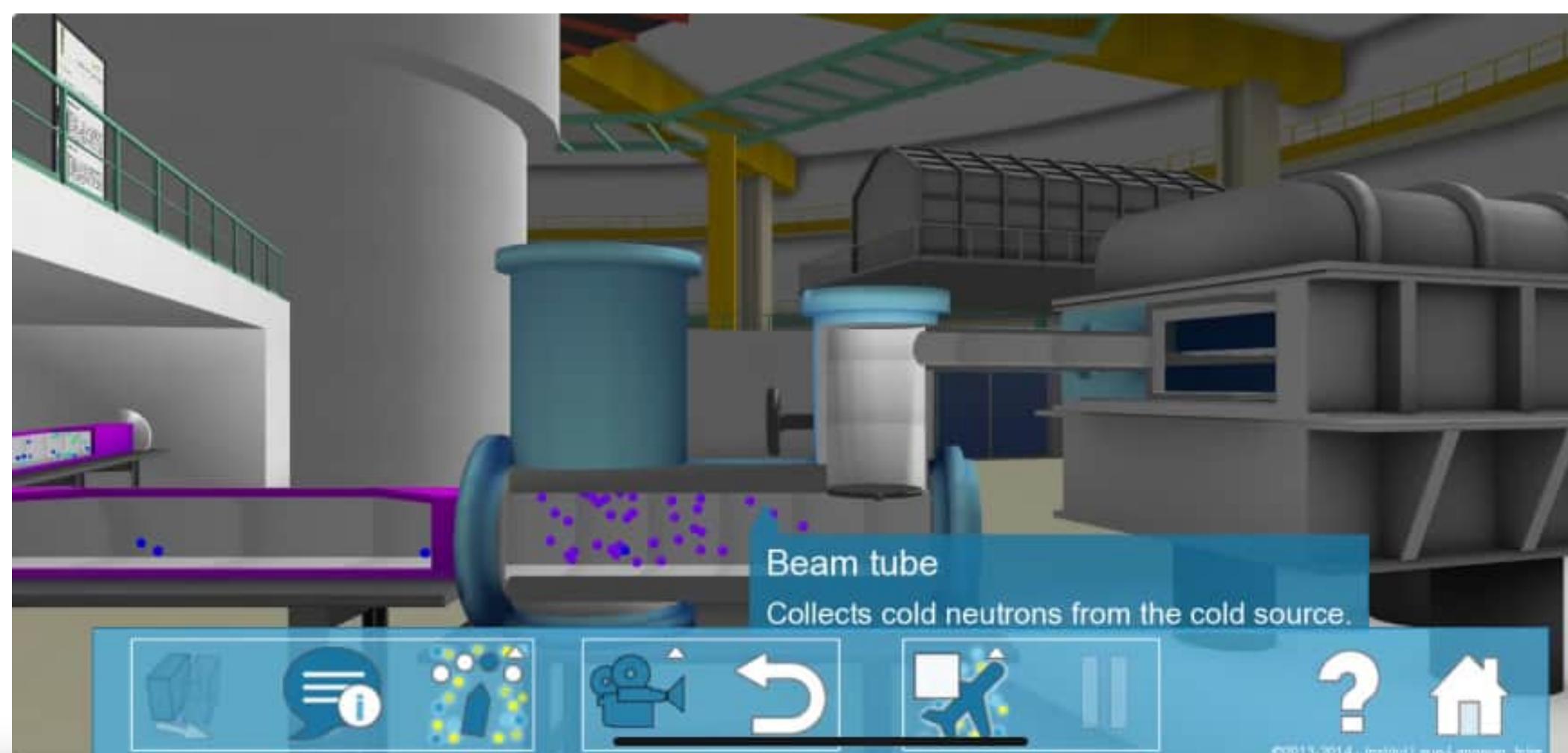
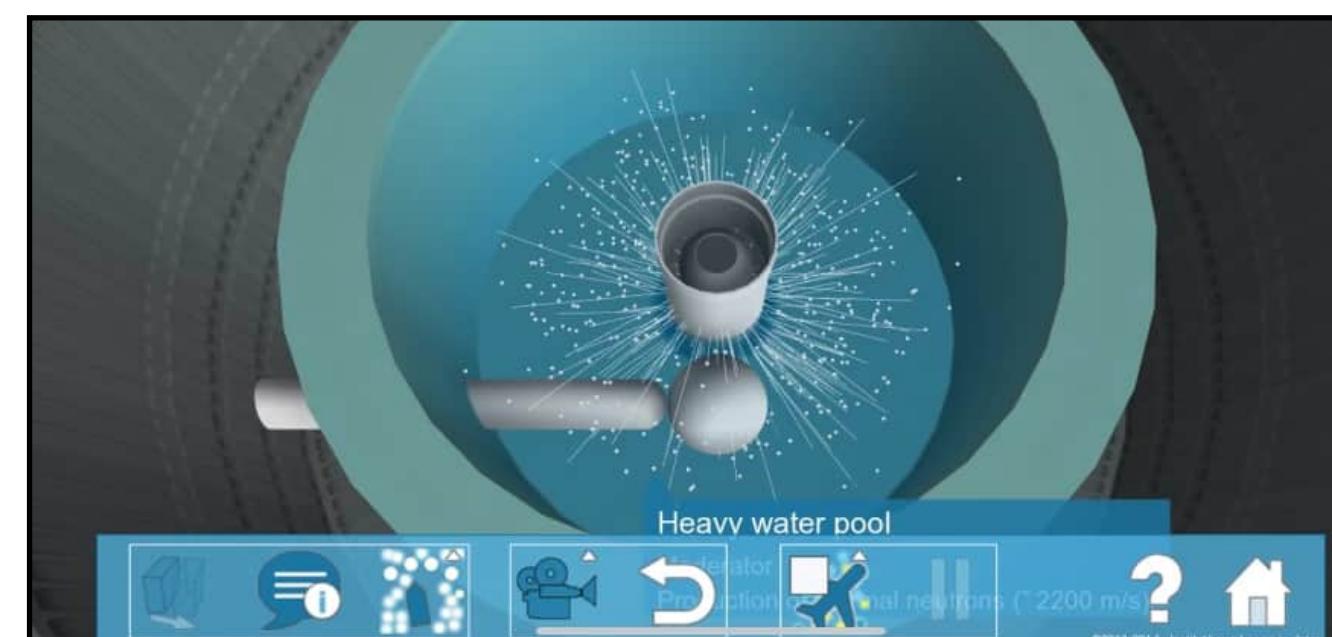
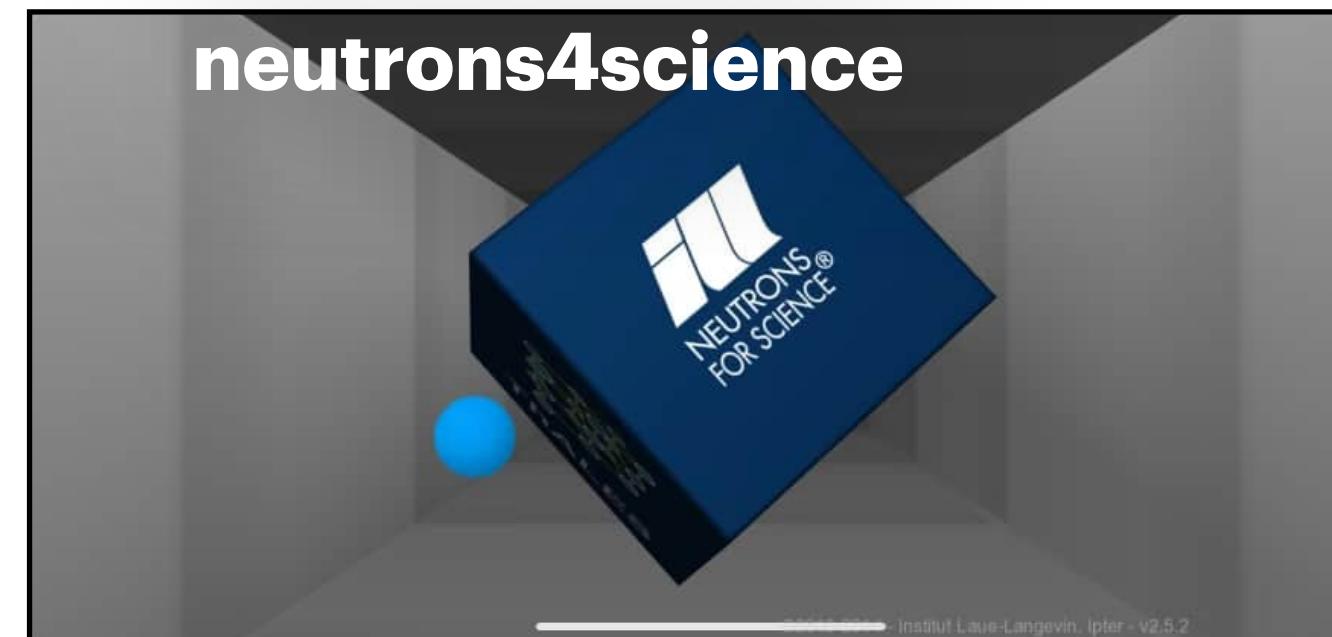
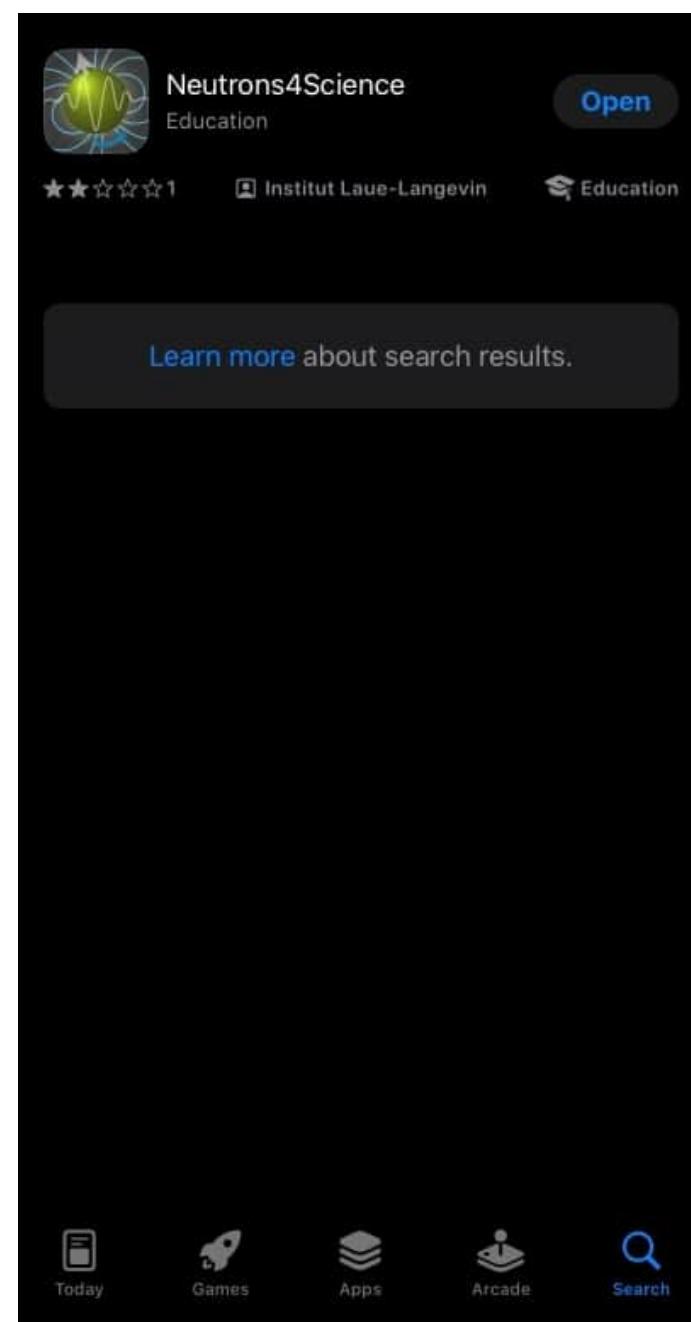
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# Useful Links and Funny Apps...on Neutrons



<https://www.ill.eu/users/support-labs-infrastructure/software-scientific-tools>

## Software & scientific tools

See also the [list of software supported](#) by the 'Computing of Science' group.

## WEB TOOL

## Activation table of elements

A table from LANSCE which allows you to calculate the activation of a sample after it has been in a neutron beam for one day and the amount of time for it to decay to 2nCi/g or less. Topics: Storage time, Prompt activation, Contact dose.

## Energy-dependent neutron cross-section

An interactive tool from the Nuclear Data Center (Japan Atomic Energy Agency) offering a compilation of many experimental data about nuclides.

**Neutron scattering length table** (from **NIST** Center for Neutron Research)

(Select the element, and you will get a list of scattering lengths and cross sections. All of this data was taken from the Special Feature section of neutron scattering lengths and cross sections of the elements and their isotopes in *Neutron News*, Vol. 3, No. 3, 1992, pp. 29-37.

Neutron scattering lengths (from Vienna Atominsti

It offers several different tables, including one in the form of a clickable periodic table.

Periodic table of elements ([www.ptable.com](http://www.ptable.com))

An interactive table compiling many data about elements.

# **BEST WISHES IN HONOUR OF JOHN M. CARPENTER**

**<<I wish you all good “neutron” sources>>**



**The father of the technique for utilising accelerator-induced intense pulses of neutrons for research  
and  
developer of the first spallation slow neutron source based on a proton synchrotron,  
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# **THANK YOU !**