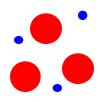




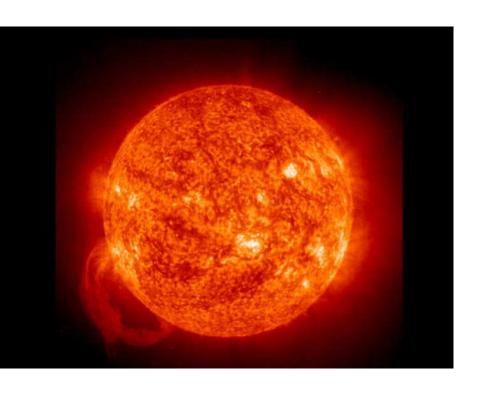


erc

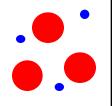




Ideal magnetohydrodynamics

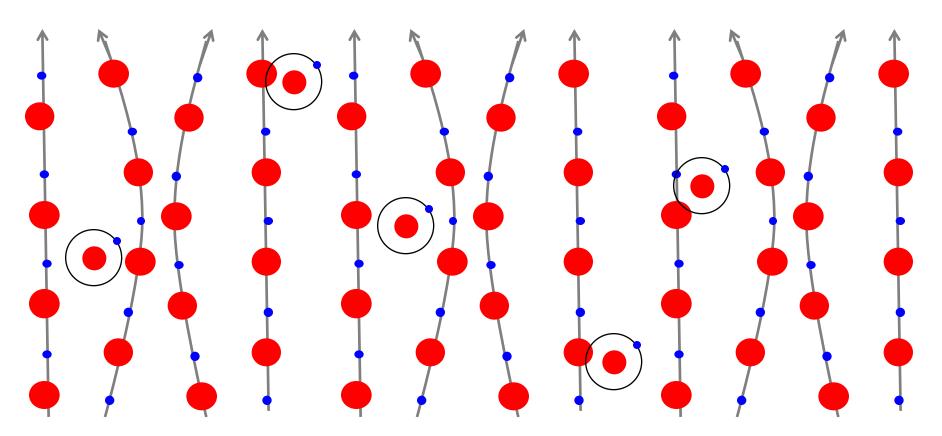


$$\frac{\mathrm{d}\boldsymbol{B}}{\mathrm{d}t} = (\boldsymbol{B} \cdot \boldsymbol{\nabla}) \boldsymbol{v} - \boldsymbol{B} (\boldsymbol{\nabla} \cdot \boldsymbol{v})$$

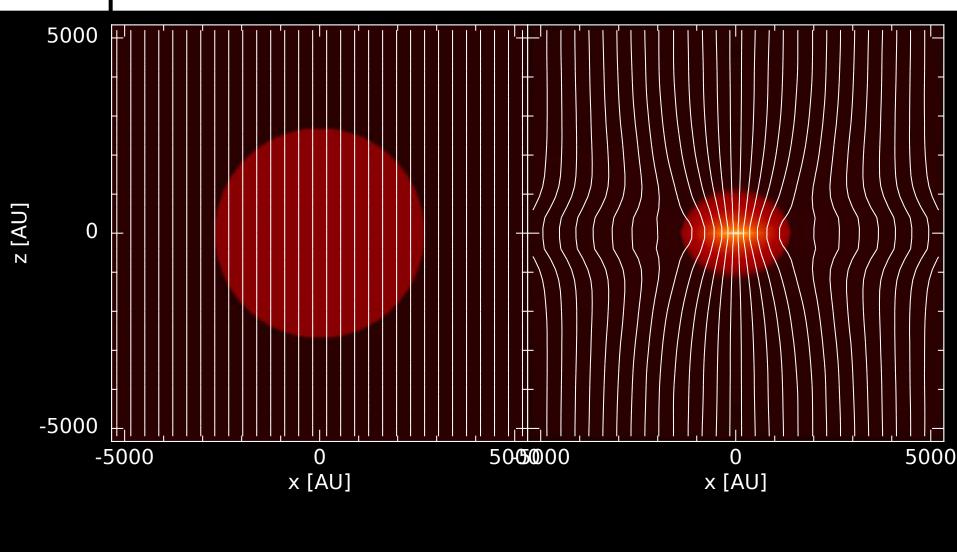


Ideal MHD

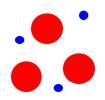
- ➤ Fully ionised plasma
- **•** + •
- ➤Zero resistivity & infinite conductivity
- ► Ions & electrons are tied to the magnetic field

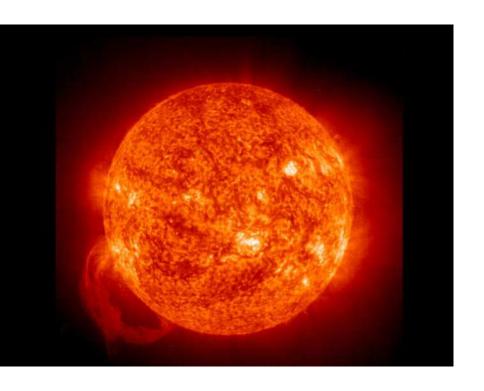






Density (rendered) + Magnetic field lines Ideal MHD. Left: Initial conditions. Right: at $\rho_{max} = 10^{-9} \text{g cm}^{-3}$

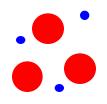




$$\frac{\mathrm{d}\boldsymbol{B}}{\mathrm{d}t} = (\boldsymbol{B} \cdot \boldsymbol{\nabla}) \boldsymbol{v} - \boldsymbol{B} (\boldsymbol{\nabla} \cdot \boldsymbol{v}) + \boldsymbol{\nabla} \times \eta_{\mathrm{art}} (\boldsymbol{\nabla} \times \boldsymbol{B})$$

where

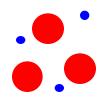
$$\eta_{
m art}pproxrac{1}{2}lpha_{
m B}v_{
m sig}h$$



➤ Artificial resistivity (Tricco & Price, 2013)

$$\begin{split} \frac{\mathrm{d}B_{a}^{i}}{\mathrm{d}t} &= -\frac{1}{\Omega_{a}\rho_{a}} \sum_{b} m_{b} \left[v_{ab}^{i} B_{a}^{j} \nabla_{a}^{j} W_{ab} \left(h_{a} \right) - B_{a}^{i} v_{ab}^{j} \nabla_{a}^{j} W_{ab} \left(h_{a} \right) \right] + \frac{\mathrm{d}B_{a}^{i}}{\mathrm{d}t} \bigg|_{\mathrm{art}} \\ \frac{\mathrm{d}B_{a}^{i}}{\mathrm{d}t} \bigg|_{\mathrm{art}} &= \frac{\rho_{a}}{2} \sum_{b} m_{b} B_{ab}^{i} \left[\frac{\alpha_{a}^{\mathrm{B}} v_{\mathrm{sig},a} \hat{r}_{ab}^{j} \nabla_{a}^{j} W_{ab} \left(h_{a} \right)}{\Omega_{a} \rho_{a}^{2}} + \frac{\alpha_{b}^{\mathrm{B}} v_{\mathrm{sig},b} \hat{r}_{ab}^{j} \nabla_{a}^{j} W_{ab} \left(h_{b} \right)}{\Omega_{b} \rho_{b}^{2}} \right] \\ v_{ab}^{i} &= v_{a}^{i} - v_{b}^{i} \\ B_{ab}^{i} &= B_{a}^{i} - B_{b}^{i} \\ v_{\mathrm{sig},a} &= \sqrt{c_{\mathrm{s},a}^{2} + v_{\mathrm{A},a}^{2}} \\ \alpha_{a}^{\mathrm{B}} &= \min \left(\frac{h_{a} \left| \nabla B_{a} \right|}{\left| B_{a} \right|}, 1 \right) \\ |\nabla B_{a}| &\equiv \sqrt{\sum_{i} \sum_{j} \left| \frac{\partial B_{a}^{i}}{\partial x_{a}^{j}} \right|^{2}} \end{split}$$

 \triangleright Always applied if there is a gradient in the magnetic field (i.e. $|\nabla \mathbf{B}| > 0$)

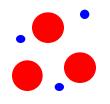


> Artificial resistivity (Price, et al, submitted)

$$\frac{\mathrm{d}B_a^i}{\mathrm{d}t} = -\frac{1}{\Omega_a \rho_a} \sum_b m_b \left[v_{ab}^i B_a^j \nabla_a^j W_{ab} \left(h_a \right) - B_a^i v_{ab}^j \nabla_a^j W_{ab} \left(h_a \right) \right] + \left. \frac{\mathrm{d}B_a^i}{\mathrm{d}t} \right|_{\mathrm{art}}$$

$$\frac{\mathrm{d}B_{a}^{i}}{\mathrm{d}t}\Big|_{\mathrm{art}} = \frac{\rho_{a}}{2} \sum_{b} m_{b} \alpha^{\mathrm{B}} v_{\mathrm{sig},ab} B_{ab}^{i} \left[\frac{\hat{r}_{ab}^{j} \nabla_{a}^{j} W_{ab}(h_{a})}{\Omega_{a} \rho_{a}^{2}} + \frac{\hat{r}_{ab}^{j} \nabla_{a}^{j} W_{ab}(h_{b})}{\Omega_{b} \rho_{b}^{2}} \right]
B_{ab}^{i} = B_{a}^{i} - B_{b}^{i}
v_{\mathrm{sig},ab} = |\mathbf{v}_{ab} \times \hat{\mathbf{r}}_{ab}|$$

- > Always applied for non-zero velocity
- Less resistive that that from Tricco & Price (2013)

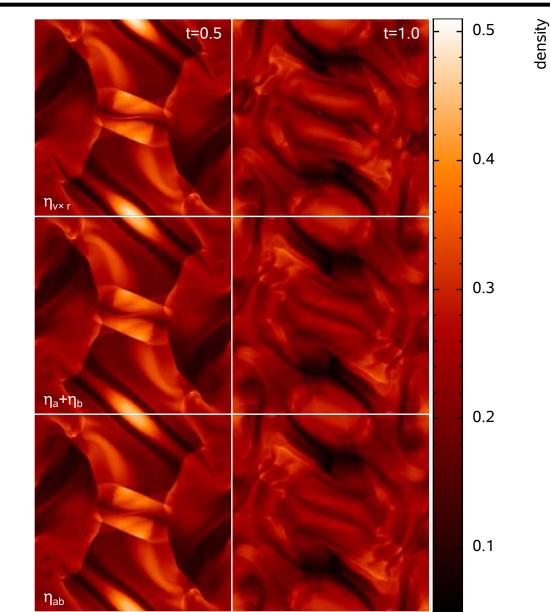


> Price et. al. (2017) artificial resistivity

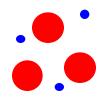
$$egin{array}{lll} v_{\mathrm{sig},ab} &=& |oldsymbol{v}_{ab} imes \hat{oldsymbol{r}}_{ab}| \ lpha^{\mathrm{B}} &\equiv& 1 \end{array}$$

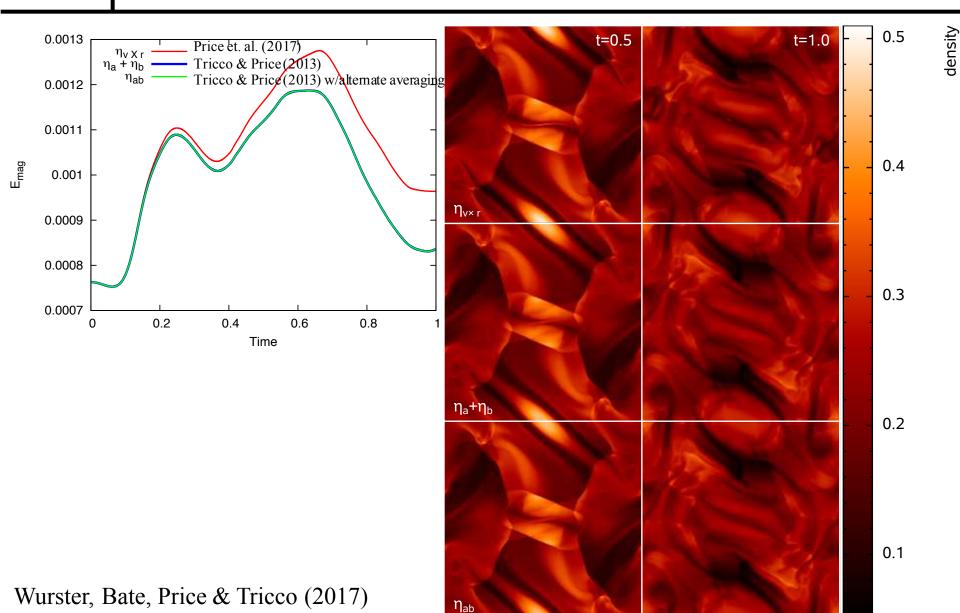
> Tricco & Price (2013)

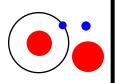
> Tricco & Price (2013) with alternate averaging



Wurster, Bate, Price & Tricco (2017)

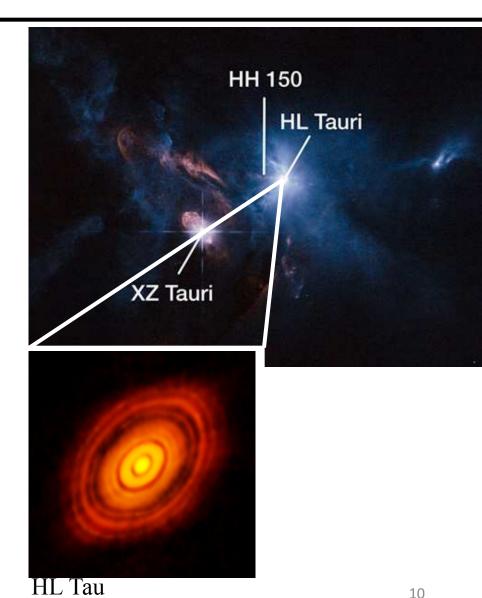






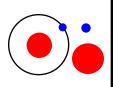
Motivation: Non-ideal MHD





 $\sim 10^{-12}$

Orion Molecular Cloud Ionisation fraction $\sim 10^{-14}$



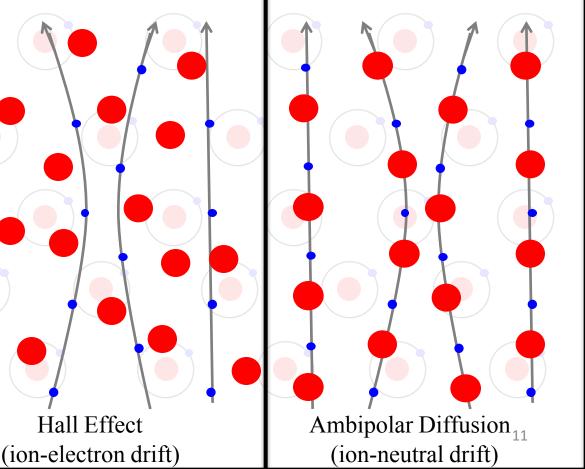
Non-ideal MHD

➤ Partially ionised plasma

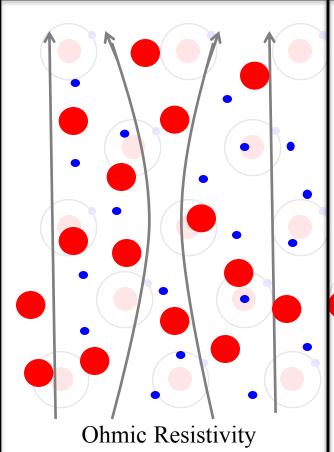


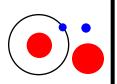
Hall Effect

- ➤ Non-zero resistivity & conductivity
- ➤ Ions, electrons & neutrals behaviour is environment-dependent

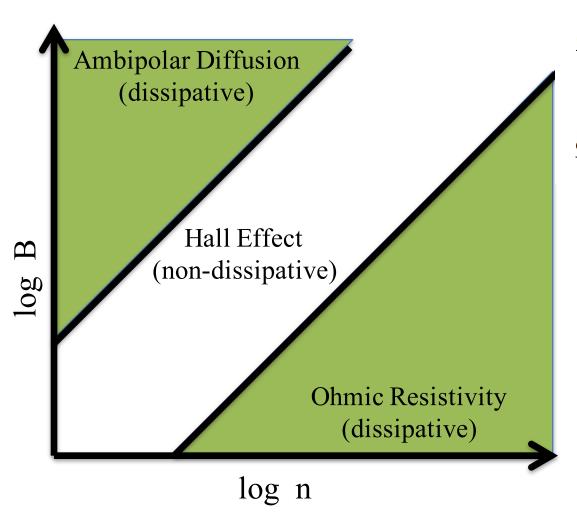


 \boldsymbol{B}

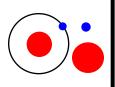




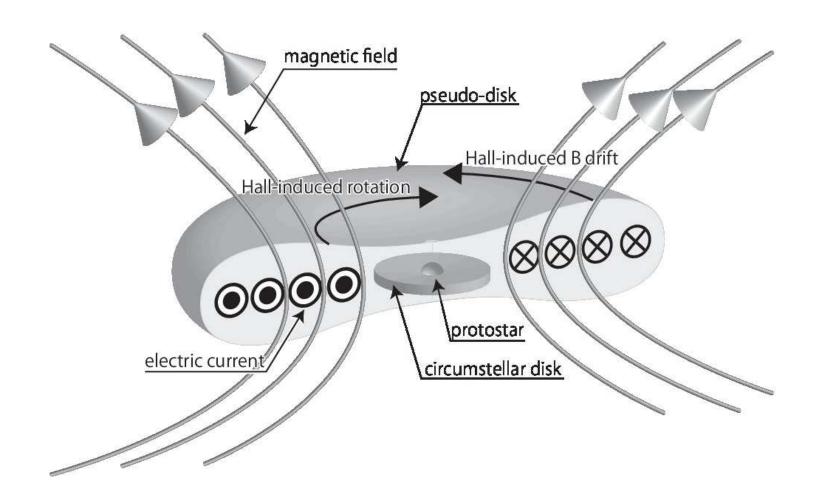
Non-ideal MHD

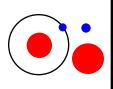


$$\begin{split} \frac{\mathrm{d}\boldsymbol{B}}{\mathrm{d}t}\bigg|_{\mathrm{OR}} &= -\boldsymbol{\nabla}\times\boldsymbol{\eta}_{\mathrm{OR}}\left(\boldsymbol{\nabla}\times\boldsymbol{B}\right),\\ \frac{\mathrm{d}\boldsymbol{B}}{\mathrm{d}t}\bigg|_{\mathrm{HE}} &= -\boldsymbol{\nabla}\times\boldsymbol{\eta}_{\mathrm{HE}}\left[\left(\boldsymbol{\nabla}\times\boldsymbol{B}\right)\times\hat{\boldsymbol{B}}\right],\\ \frac{\mathrm{d}\boldsymbol{B}}{\mathrm{d}t}\bigg|_{\mathrm{AD}} &= \boldsymbol{\nabla}\times\boldsymbol{\eta}_{\mathrm{AD}}\left\{\left[\left(\boldsymbol{\nabla}\times\boldsymbol{B}\right)\times\hat{\boldsymbol{B}}\right]\times\hat{\boldsymbol{B}}\right\}. \end{split}$$

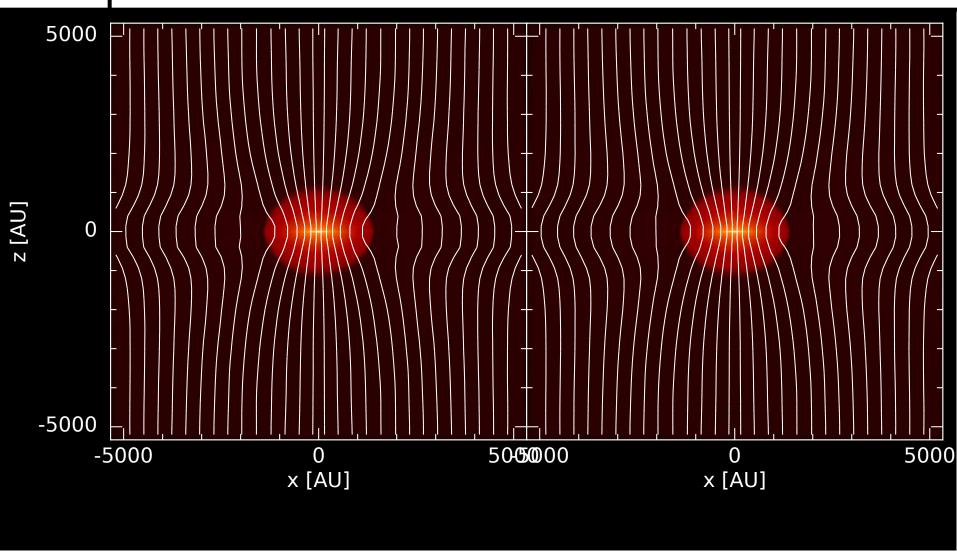


Non-ideal MHD: Hall effect

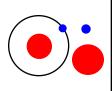




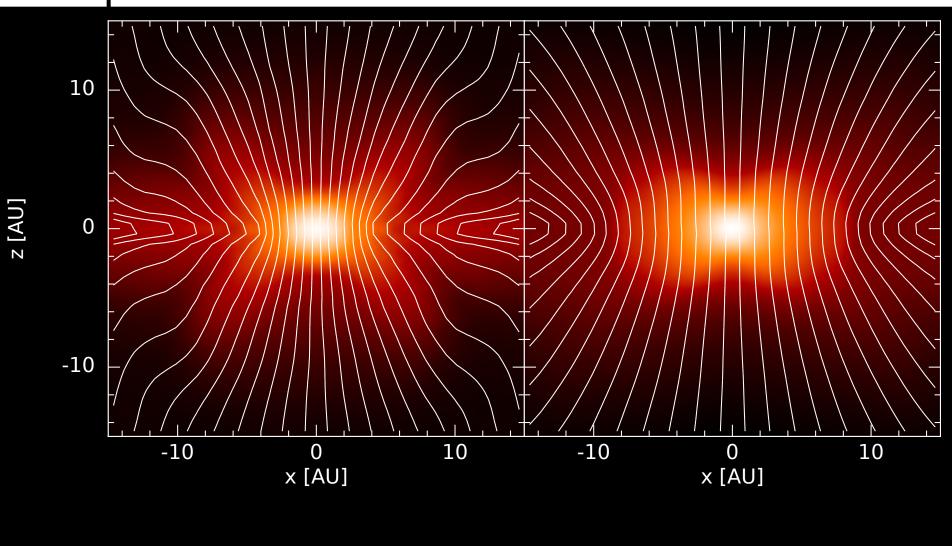
Ideal vs non-ideal MHD



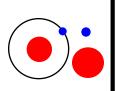
Density (rendered) + Magnetic field lines During first core phase. Left: ideal MHD. Right: non-ideal MHD



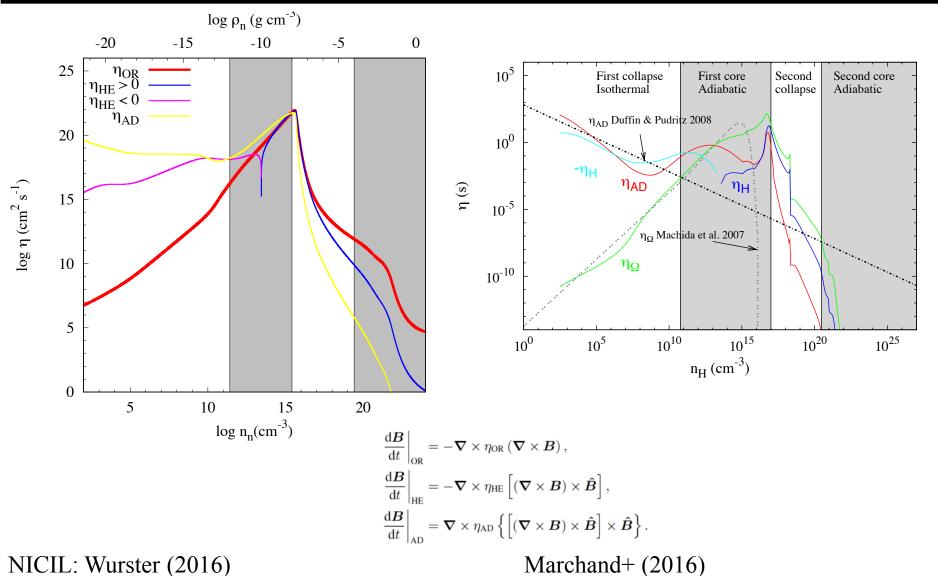
Ideal vs non-ideal MHD



Density (rendered) + Magnetic field lines During first core phase. Left: ideal MHD. Right: non-ideal MHD

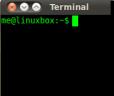


Non-ideal MHD



NICIL: Wurster (2016)

NICIL v1.2.3 is implemented in the current git version of Phantom



Non-ideal MHD in Phantom: the NICIL library

- > Phantom includes the NICIL code (Wurster 2016)
 - ➤ Publically available at https://bitbucket.org/jameswurster/nicil
- ➤ When compiling, set NONIDEALMHD=yes
- Realistic defaults are set; these will self-consistently calculate the non-ideal coefficients
- >Fully parameterisable
- ➤ Primary parameters are included in *Phantom*'s .in file
- ➤ All parameters are included at the top of nicil. F90
- ➤ Important parameters that can be modified
 - ➤ Included non-ideal MHD terms (default = ohmic + Hall + ambipolar)
 - ➤ Ionisation source (default = cosmic rays + thermal)
 - \triangleright Cosmic ray ionisation rate (default = 10^{-17} s⁻¹)
 - Elements that can be thermally ionised (cannot be modified through .in file)
 - For ain properties (default = fixed size of $0.1\mu m$; alternate is MRN, but is slow)
- >Important values are summarised in the dump files and the .ev file
- ➤ Can optionally preselect non-ideal MHD coefficients (preferably for tests only)
- ➤ All coefficients and required variables are calculated at runtime

mp(iH2) * massj_mp(in mp(iHe) * massj_mp(in

≻Continuum equations

$$\frac{\mathrm{d}\boldsymbol{B}}{\mathrm{d}t} = (\boldsymbol{B} \cdot \boldsymbol{\nabla}) \boldsymbol{v} - \boldsymbol{B} (\boldsymbol{\nabla} \cdot \boldsymbol{v})
+ \boldsymbol{\nabla} \times \eta_{\mathrm{art}} (\boldsymbol{\nabla} \times \boldsymbol{B})
+ \boldsymbol{\nabla} \times \eta_{\mathrm{OR}} (\boldsymbol{\nabla} \times \boldsymbol{B})
+ \boldsymbol{\nabla} \times \eta_{\mathrm{HE}} \left[(\boldsymbol{\nabla} \times \boldsymbol{B}) \times \hat{\boldsymbol{B}} \right]
+ \boldsymbol{\nabla} \times \eta_{\mathrm{AD}} \left\{ \left[(\boldsymbol{\nabla} \times \boldsymbol{B}) \times \hat{\boldsymbol{B}} \right] \times \hat{\boldsymbol{B}} \right\}$$

➤ SPMHD equations

$$\frac{\mathrm{d}B_{a}^{i}}{\mathrm{d}t} = -\frac{1}{\Omega_{a}\rho_{a}}\sum_{b}m_{b}\left[v_{ab}^{i}B_{a}^{j}\nabla_{a}^{j}W_{ab}\left(h_{a}\right)\right. - B_{a}^{i}v_{ab}^{j}\nabla_{a}^{j}W_{ab}\left(h_{a}\right)\right] + \left.\frac{\mathrm{d}B_{a}^{i}}{\mathrm{d}t}\right|_{\text{non-ideal}}$$

$$\frac{\mathrm{d}\boldsymbol{B}_{a}}{\mathrm{d}t}\Big|_{\text{non-ideal}} = -\rho_{a} \sum_{b} m_{b} \left[\frac{\boldsymbol{D}_{a}}{\Omega_{a} \rho_{a}^{2}} \times \nabla_{a} W_{ab}(h_{a}) + \frac{\boldsymbol{D}_{b}}{\Omega_{b} \rho_{b}^{2}} \times \nabla_{a} W_{ab}(h_{b}) \right],$$

$$\boldsymbol{D}_{a}^{\text{OR}} = -\eta_{\text{OR}} \boldsymbol{J}_{a}, \quad \boldsymbol{D}_{a}^{\text{HE}} = -\eta_{\text{HE}} \boldsymbol{J}_{a} \times \hat{\boldsymbol{B}}_{a}, \quad \boldsymbol{D}_{a}^{\text{AD}} = \eta_{\text{AD}} \left(\boldsymbol{J}_{a} \times \hat{\boldsymbol{B}}_{a} \right) \times \hat{\boldsymbol{B}}_{a}.$$

Wurster, Price & Ayliffe (2014)

```
ss (CGS; sigma units are contrates, and the maximum sigmaviRn is for cosmic ray (mj_mp(iH2) * massj_mp(in (mj_mp(iH2) * massj_mp(in (mj_mp(iH2) * massj_mp(in (mj_mp(iH2) * massj_mp(j-(mj_mp(iH) * massj_mp(j-(mj_mp(iH) * massj_mp(j-2.81d-9 * sq
```

do i = 1, N

enddo

Implementation

```
Density Loop:
do i = 1, N
    do j = 1, N_{\text{neigh}}
        Using j, calculate density of i
        Using j, calculate current density, oldsymbol{J} = oldsymbol{
abla} 	imes oldsymbol{B}, of i
    enddo
    Using new density of i_{\star} calculate \eta_{\text{nimhd}}
enddo
Force Loop:
do i = 1, N
    Calculate oldsymbol{J}_i 	imes oldsymbol{B}_i and (oldsymbol{J}_i 	imes oldsymbol{B}_i) 	imes oldsymbol{B}_i
    do j = 1, N_{\text{neigh}}
        Calculate J_i \times B_i and (J_i \times B_i) \times B_i
        Using j, calculate \mathrm{d}B/\mathrm{d}t_{\mathrm{non-ideal}} of i
    enddo
    Calculate non-ideal timesteps
enddo
Step Loop:
```

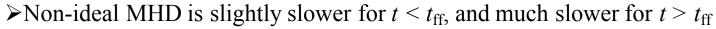
Updated magnetic field of i, using ideal, non-ideal and artificial terms

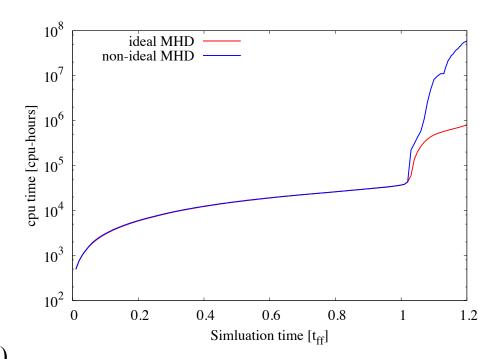
Implementation

➤Timestepping:

$$dt_{\text{Courant}} = C_{\text{c}} \frac{h}{v_{\text{sig}}}$$
$$dt_{\text{nimhd}} = C_{\text{ni}} \frac{h^2}{|\eta|}$$

- ➤ Phantom includes super-timestepping (Alexiades, Amiez & Gremaud 1996)
- ➤ Right: cpu-hours required for the 10^6 particle models with μ_0 =5 in Wurster, Price & Bate (2016)







Conclusions

- Artificial resistivity is required to stabilised magnetohydrodynamics equations
- ➤ Ideal MHD is a poor approximation for modelling molecular clouds or protoplanetary discs
- ➤ Non-ideal MHD requires an assumption of chemistry
- The non-ideal MHD coefficients are not dependent on neighbours
- The non-ideal MHD contribution to the magnetic field evolution is dependent on neighbours
- Non-ideal MHD introduces a diffusion timestep $\propto h^2$, hence can be computationally expensive

"Always code as if the guy who ends up maintaining your code will be a violent psychopath who knows where you live."

~ John Woods





j.wurster@exeter.ac.uk

http://www.astro.ex.ac.uk/people/wurster/

Presentation available at http://www.astro.ex.ac.uk/people/wurster/files/spmhd_resistivity.pdf

Nicil's git repository: https://bitbucket.org/jameswurster/nicil