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Physics 441: Assignment #4 - Potentials

Due on Friday, May 31, 2013

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Problem 3.3

Find the general solution to Laplace's equation in spherical coordinates, for the case where V depends only on r . Do the same for cylindrical coordinates, assuming V depends only on s

Laplace's equation in spherical coordinates takes the following form:

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} = 0$$

The problem says that V only depends on r , so we can simplify this to parts that only contain derivatives in r . In this case, Laplace's equation becomes:

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) = 0$$

We can cancel out the leading $\frac{1}{r^2}$ (because the left hand side is 0) and get

$$\frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) = 0$$

which simplifies down to

$$\left(r^2 \frac{\partial V}{\partial r} \right) = c \rightarrow \frac{\partial V}{\partial r} = \frac{c}{r^2} \rightarrow V = -\frac{c}{r} + k$$

Laplace's equation in cylindrical coordinates is

$$\nabla^2 V = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial V}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

Doing the same and keeping only terms with derivatives in s we simplify to

$$\nabla^2 V = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial V}{\partial s} \right) = 0$$

I again multiply both sides by s to arrive at the following expression, which I simplify to get the final result:

$$\frac{\partial}{\partial s} \left(s \frac{\partial V}{\partial s} \right) = 0 \rightarrow s \frac{\partial V}{\partial s} = c \rightarrow \frac{\partial V}{\partial s} = \frac{c}{s} \rightarrow V = c \ln s + k$$

□

Problem 3.10

A uniform line charge λ is placed on an infinite straight wire, a distance d above a grounded conducting plane. (Let's say the wire runs parallel to the x -axis and directly above it, and the conducting plane is the xy plane.)

1. Find the potential in the region above the plane [HINT: refer to problem 2.52]
2. Find the charge density σ induced on the conducting plane

1. We know that energy (potential in parenthesis) is a function of $\hat{r}r^2$ ($-\hat{r}r$) in 3d and \hat{r}/r ($\ln r \hat{r}$) in 3d. Our problem is a 2-d problem (we are dealing with a line). Using that, and remembering we pick up factor of $2\pi\epsilon_0$ from integrating $E \rightarrow V$ we know that the general formula for potential in 2d is the following.

$$\begin{aligned} V(\mathbf{s}) &= \frac{\lambda}{2\pi\epsilon_0} \ln(\mathbf{a}/\mathbf{s}) \\ &= \frac{\lambda}{4\pi\epsilon_0} \ln(\mathbf{a}^2/\mathbf{s}^2) \end{aligned}$$

In this case we say that $\mathbf{a} = -\mathbf{s}$ (it is an image problem!). We also say that \mathbf{a} is the distance away from the x -axis in the xy plane. An expression for this is $\mathbf{a} = y + (z + d)$. We can now apply this to our expression for the potential energy to get the answer:

$$V(y, z) = \frac{\lambda}{4\pi\epsilon_0} \ln \left(\frac{y^2 + (z + d)^2}{y^2 + (z - d)^2} \right)$$

2. Now we need to find σ on the plane ($z = 0$). We will use the equation at the top of page 126:

$$\sigma = -\epsilon_0 \left. \frac{\partial V}{\partial z} \right|_{z=0}$$

We apply this equation and simplify to get $\sigma(y)$ (Note I let the computer do the algebra for me and I have included the code below)

$$\begin{aligned}\sigma(y) &= -\epsilon_0 \left. \frac{\partial V}{\partial z} \right|_{z=0} \\ &= \left. \frac{\lambda (y^2 + (-d+z)^2) \left(\frac{(2d-2z)(y^2+(d+z)^2)}{(y^2+(-d+z)^2)^2} + \frac{2d+2z}{y^2+(-d+z)^2} \right)}{4e_0\pi (y^2 + (d+z)^2)} \right|_{z=0} \\ &= -\frac{d\lambda}{\pi(d^2 + y^2)}\end{aligned}$$

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1  import sympy as sym
   y, z, d, lamb, pi, e0 = sym.symbols('y, z, d, lamb, pi, e0 ')
   expr = lamb / (4 * pi * e0) * sym.log((y**2 + (z + d) **2) / (y**2 + (z-d)**2))
6
   print('\ve_0 * dv/dz: ')
   sym.pprint(-e0 * expr.diff(z))
   print(sym.latex(-e0 * expr.diff(z)))
11
   print('\n\n\n\n \ve_0 * dv/dz at z=0: ')
   sym.pprint(-e0 * expr.diff(z).subs({z:0}))
   print(sym.latex(-e0 * expr.diff(z).subs({z:0})))

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□

Problem 3.14

For the infinite slot (see example 3.3), determine the charge density $\sigma(y)$ on the strip at $x = 0$, assuming it is a conductor at constant potential V_0

Problem 3.19

The potential at the surface of a sphere (radius R) is given by

$$V_0 = k \cos 3\theta$$

where k is a constant. Find the potential inside and outside the sphere, as well as the surface charge density $\sigma(\theta)$ on the sphere. (assume there is not charge inside or outside the sphere.)

Problem 3.24

Solve Laplace's equation by separation of variables in cylindrical coordinates, assuming there is no dependence on z (cylindrical symmetry). [Make sure you find all solutions to the radial equation; in particular, your results must accommodate the case of an infinite line charge, for which (of course) we already know the answer.]

Problem 3.29

For particles (one of charge q , one of charge $3q$, and two of charge $-2q$) are placed at the following points:

- $q \rightarrow (0, 0, -1)$
- $3q \rightarrow (0, 0, 1)$
- $-2q \rightarrow (0, -1, 0)$
- $-2q \rightarrow (0, 1, 0)$

Find a simple approximate formula for the potential, valid at points far from the origin (Express the answer in spherical coordinates)

Problem 3.31

For the dipole in example 3.10, expand $1/r_{\pm}$ to order $(d/r)^3$, and use this to determine the quadropole and octopole terms in the potential.
