

# Chapter 21 Electric Potential



## § 1 Electric Potential Energy

- The similarity of **electrostatic** and **gravitational** force

$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} \hat{r} \quad \text{electrostatic}$$

$$\vec{F} = -G \frac{Mm}{r^2} \hat{r} \quad \text{gravitational}$$

- ➡ Both forces depend on the inverse square of the separation distance between the two objects.

# Electrostatic vs. gravitational



$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} \hat{r} \quad \text{electrostatic}$$

$$\vec{F} = -G \frac{Mm}{r^2} \hat{r} \quad \text{gravitational}$$

- The work done by the gravitational force on the object  $m$  depends only on the starting and finishing points and does not depend on the path taken between the points — **gravitational** force is a **conservative** force.

$$\Delta U = U_f - U_i = -W_{if} = -\int_i^f \vec{F}_c \cdot d\vec{s}$$

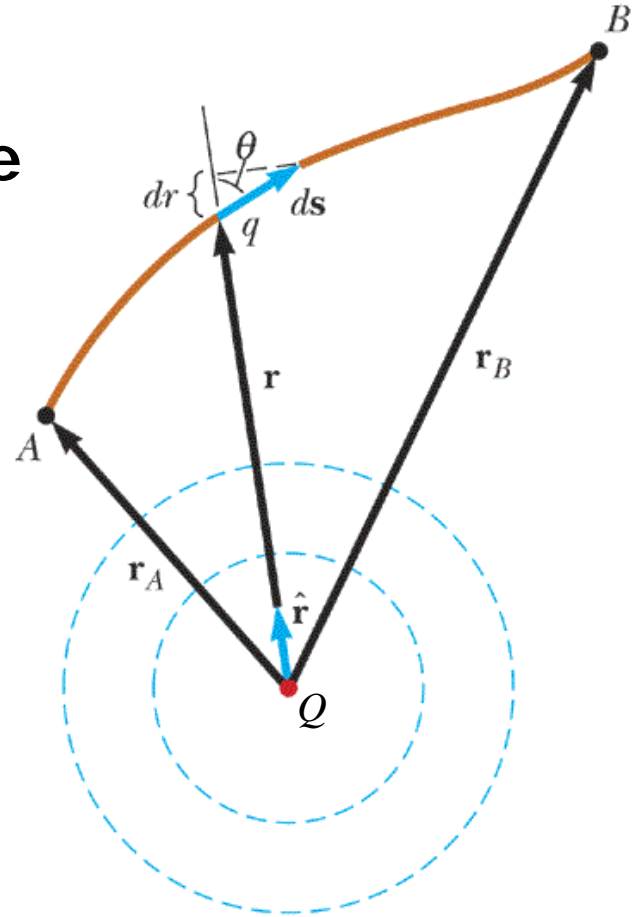
the gravitational potential energy difference

$$\Delta U = \left( -G \frac{Mm}{r_f} \right) - \left( -G \frac{Mm}{r_i} \right)$$

# The electric potential energy



- The electric potential energy
  - Because of the similarity of the electrostatic and gravitational force laws, the electrostatic force is also **conservative**, and there is a **potential energy** associated with the configuration of a system (the relative locations of the charges).



$$\vec{F} \cdot d\vec{s} = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} \hat{r} \cdot d\vec{s} = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} ds \cos \theta = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} dr$$

## The electric potential energy

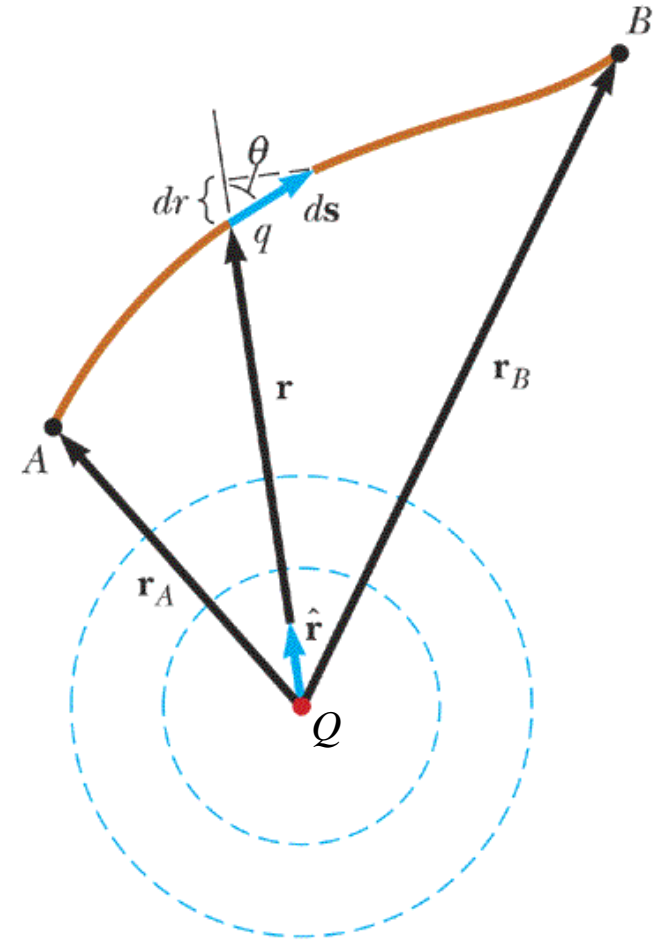


The electric potential energy difference between  $A$  and  $B$

$$\begin{aligned}\Delta U &= -\int_A^B \vec{F} \cdot d\vec{s} = -\int_{r_A}^{r_B} \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} dr \\ &= \left( \frac{1}{4\pi\epsilon_0} \frac{Qq}{r_B} \right) - \left( \frac{1}{4\pi\epsilon_0} \frac{Qq}{r_A} \right)\end{aligned}$$

If we set  $U_A(\infty) = 0$  as our reference potential energy, the potential energy at any point in space is

$$U(r) = \frac{1}{4\pi\epsilon_0} \frac{Qq}{r}$$





## ■ Electric potential

- ➡ Consider a **test** charge  $q_0$  in the **field** of charge  $Q$ . The potential energy  $U$  associates with both the test charge  $q_0$  and the source charge  $q$ , which means that the  $U$  characterizes the interaction of two charges with one another.

$$\Delta U_{BA} = \left( \frac{1}{4\pi\epsilon_0} \frac{q_0 Q}{r_B} \right) - \left( \frac{1}{4\pi\epsilon_0} \frac{q_0 Q}{r_A} \right) = - \int_A^B \vec{F} \cdot d\vec{s} = -q_0 \int_A^B \vec{E} \cdot d\vec{s}$$

- ➡ The potential energy per unit test charge, which is symbolized as  $\Delta U/q_0$ , is independent of the test charge  $q_0$ , and is characteristic only of the field of due to source charge  $q$  which we are investigating — we define the **electric potential difference**  $\Delta V$  to be the electric potential energy difference per unit test charge.

$$\Delta V_{BA} = V_B - V_A = \frac{\Delta U_{BA}}{q_0} = - \frac{1}{q_0} \int_A^B \vec{F} \cdot d\vec{s} = - \int_A^B \vec{E} \cdot d\vec{s}$$

$$\Delta V_{BA} = V_B - V_A = \frac{\Delta U_{BA}}{q_0} = -\frac{1}{q_0} \int_A^B \vec{F} \cdot d\vec{s} = -\int_A^B \vec{E} \cdot d\vec{s}$$

➡ If we set  $U(\infty) = 0$  as our reference potential

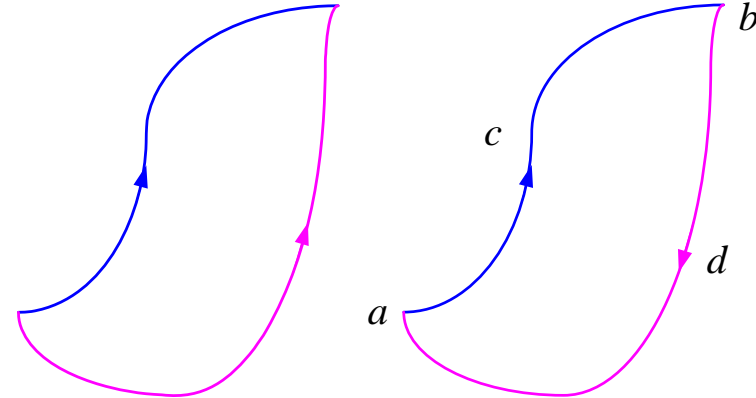
$$V_B = \int_B^\infty \vec{E} \cdot d\vec{s},$$

$$V_P = \int_P^\infty \vec{E} \cdot d\vec{s}$$

➡ SI unit:  $1V = 1 \text{ J/C}$

## ■ The Circulation law of electric potential

$$\int_{acb} \vec{E} \cdot d\vec{s} = \int_{adb} \vec{E} \cdot d\vec{s}$$



$$\int_{acb} \vec{E} \cdot d\vec{s} - \int_{adb} \vec{E} \cdot d\vec{s} = \int_{acb} \vec{E} \cdot d\vec{s} + \int_{bda} \vec{E} \cdot d\vec{s} = 0$$

$$\oint_L \vec{E} \cdot d\vec{s} = 0$$

**The circulation law of electric potential**

This law means that the electrostatic field is a **conservative** field !

# The Circulation law of electric potential

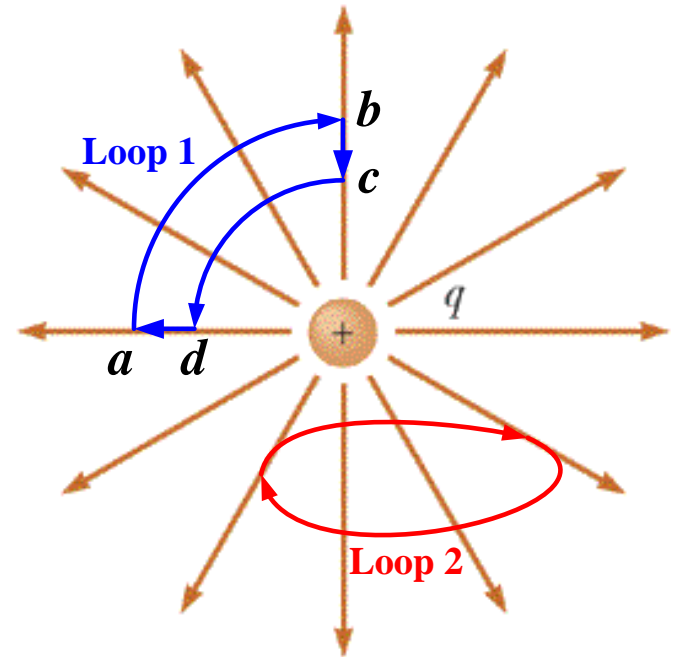


$$\oint_L \vec{E} \cdot d\vec{s} = 0$$

**Example: For an electric field of a point charge  $q$ .**

$$\begin{aligned} \oint_{\text{Loop 1}} \vec{E} \cdot d\vec{s} &= \int_a^b \vec{E} \cdot d\vec{s} + \int_b^c \vec{E} \cdot d\vec{s} \\ &+ \int_c^d \vec{E} \cdot d\vec{s} + \int_d^a \vec{E} \cdot d\vec{s} \\ &= \int_{r_b}^{r_c} k_e \frac{q}{r^2} dr + \int_{r_d}^{r_a} k_e \frac{q}{r^2} dr = 0 \end{aligned}$$

$$\oint_{\text{Loop 2}} \vec{E} \cdot d\vec{s} = 0$$







### ■ Summary of the laws for electrostatic field

#### ➤ Gauss' Law:

The electrostatic charge is the **source** of the electrostatic field.

$$\oiint_S \vec{E} \cdot d\vec{A} = \frac{q_{inside}}{\epsilon_0}$$

#### ➤ Circulation Law:

The electrostatic field is a **conservative** field. Therefore we can introduce a scalar field (electric potential) correlated to the electrostatic field.

$$\oint_L \vec{E} \cdot d\vec{s} = 0$$

## § 2 Calculating the Electric Potential



- If the electric field is known

$$V_P = \int_P^{\infty} \vec{E} \cdot d\vec{s}$$

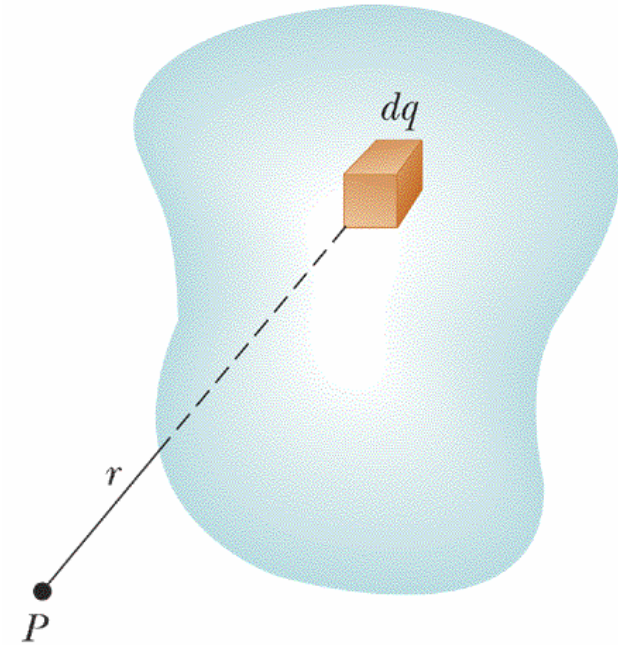
- If the charge distribution is known

➡ The electric potential due to individual charge particles

$$V = \frac{1}{4\pi\epsilon_0} \sum_i \frac{q_i}{r_i}$$

➡ The electric potential due to continuous charge distributions

$$V = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r}$$

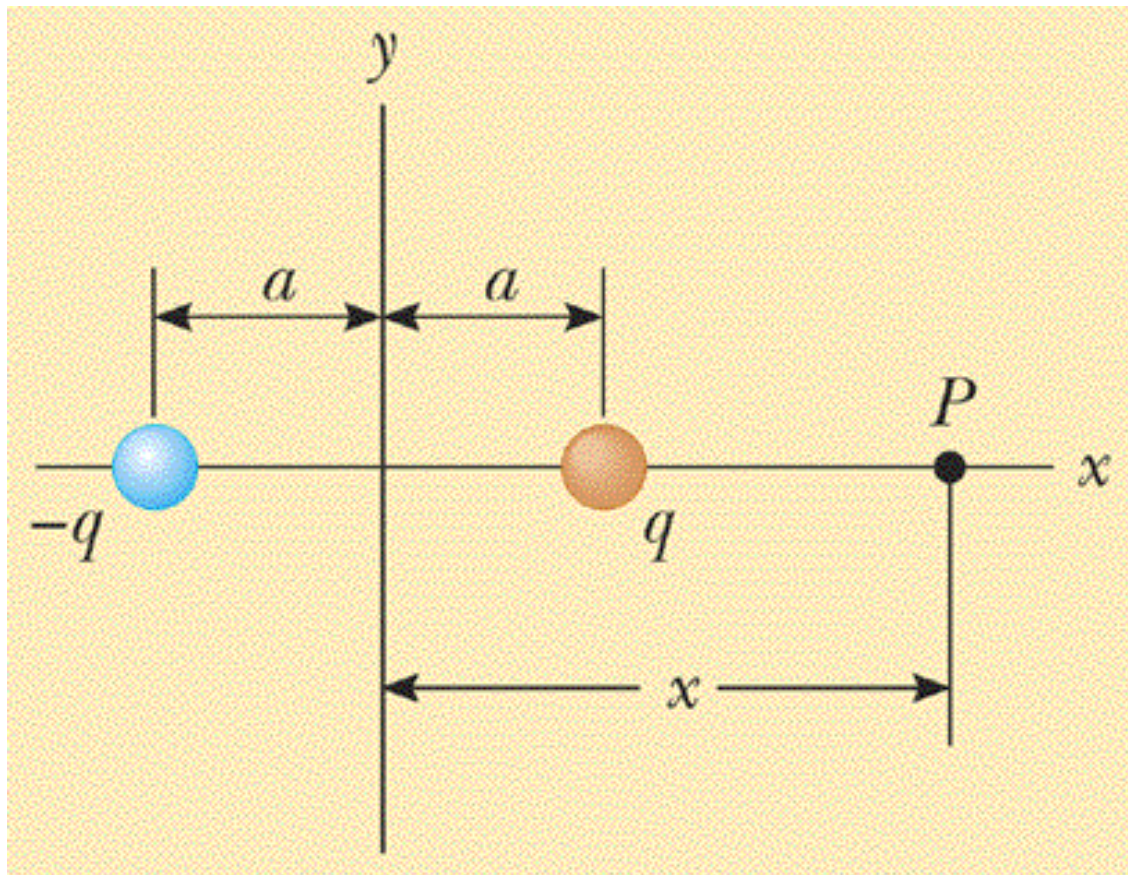


## Example — The Electric Dipole



### The electric potential of a dipole

**Example:** The dipole is along the  $x$  axis and is centered at the origin. Calculating the electric potential at any point  $P$  along the  $x$  axis.



## Example — The Electric Dipole



### The electric potential of a dipole

Example: The dipole is along the  $x$  axis and is centered at the origin.  
Calculating the electric potential at any point  $P$  along the  $x$  axis.

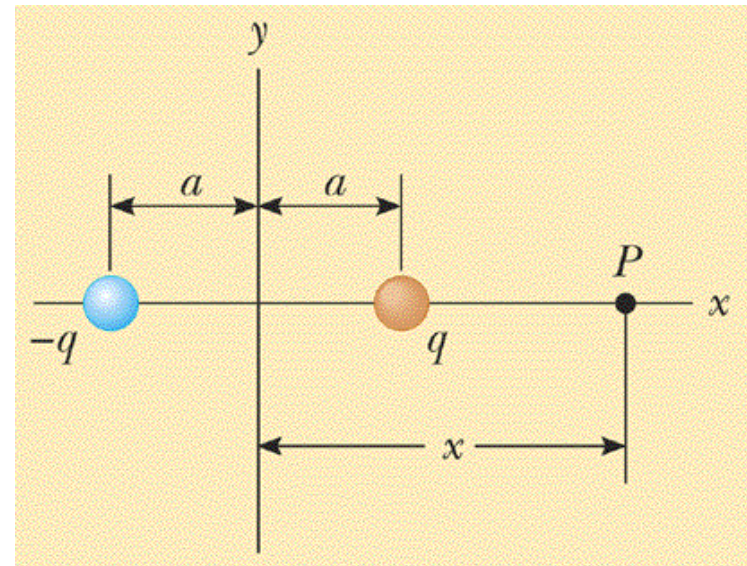
**Solution:**

$$V = \frac{1}{4\pi\epsilon_0} \left( \frac{q}{x-a} + \frac{-q}{x+a} \right)$$

$$= \frac{1}{4\pi\epsilon_0} \frac{2aq}{x^2 - a^2}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{p}{x^2 - a^2}$$

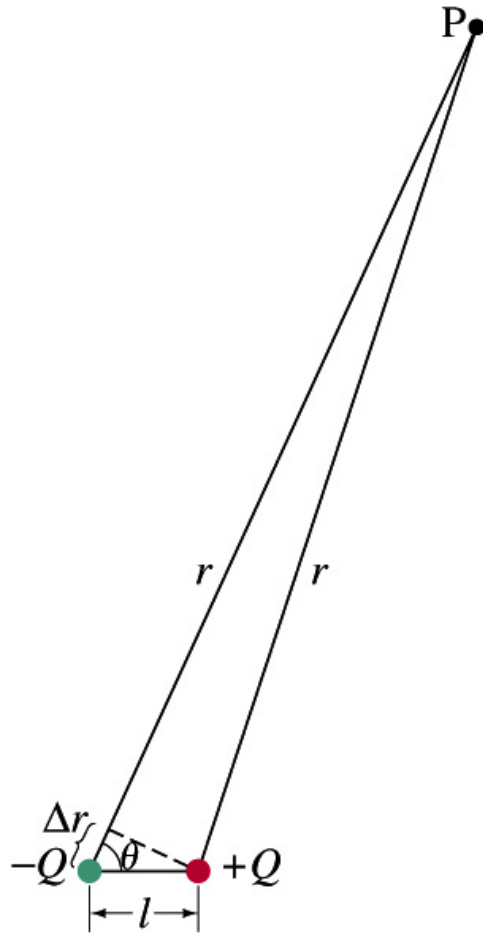
$$= \frac{1}{4\pi\epsilon_0} \frac{p}{x^2} \quad x \gg a$$



## Example — The Electric Dipole



### The electric potential of a dipole

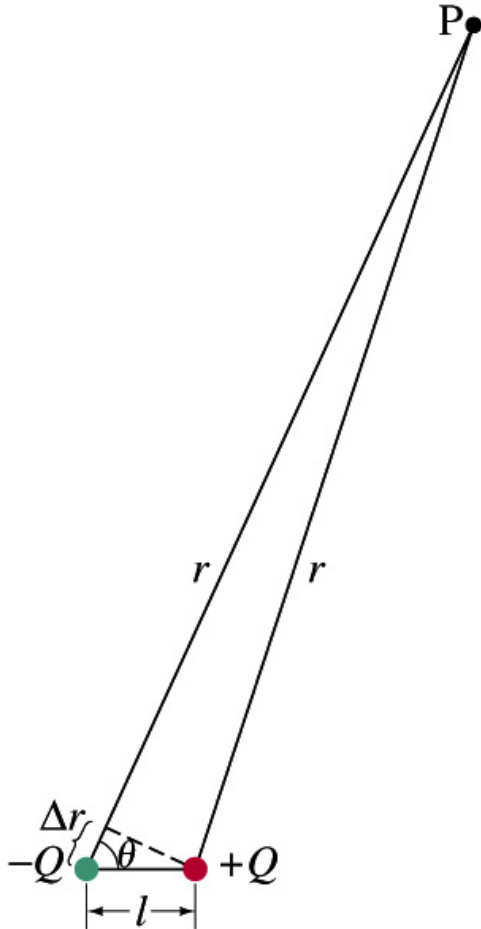


**Fig. 23-16**  
**(P512 § 21-6)**

## Example — The Electric Dipole



### The electric potential of a dipole



$$V = \frac{1}{4\pi\epsilon_0} \left( \frac{Q}{r} + \frac{-Q}{r + \Delta r} \right)$$
$$= \frac{Q}{4\pi\epsilon_0} \frac{\Delta r}{r(r + \Delta r)}$$

$$\Delta r \approx l \cos \theta, \quad r \gg \Delta r$$

$$V = \frac{Q}{4\pi\epsilon_0} \frac{l \cos \theta}{r^2}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{p \cos \theta}{r^2}$$

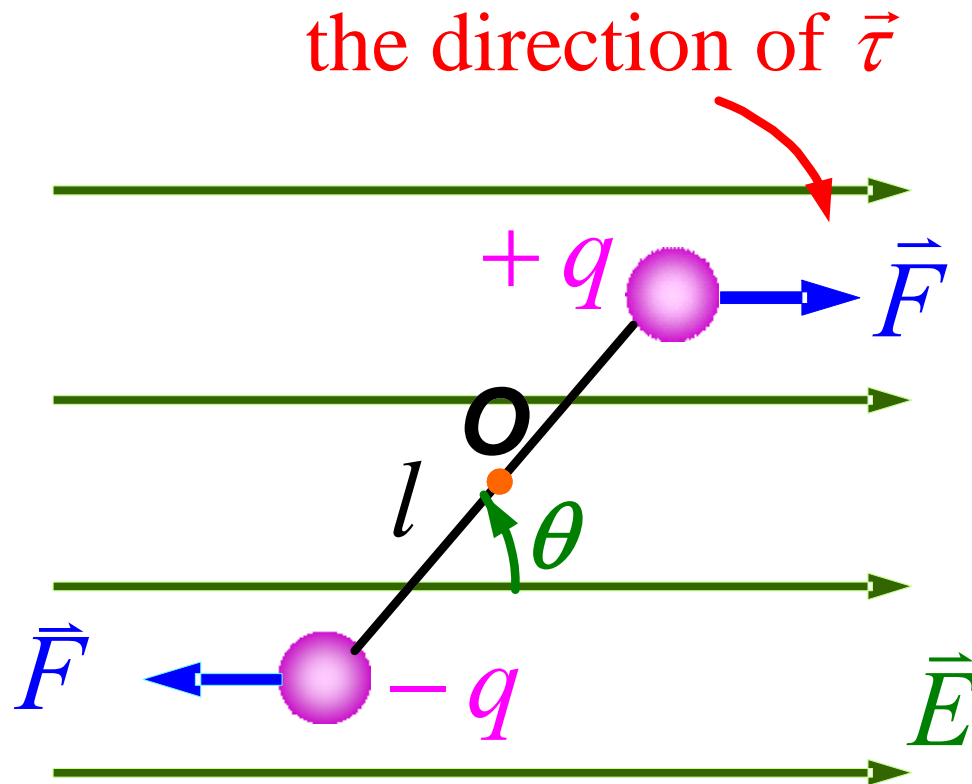
Fig. 23-16



## The Potential Energy of a Dipole in an External Field

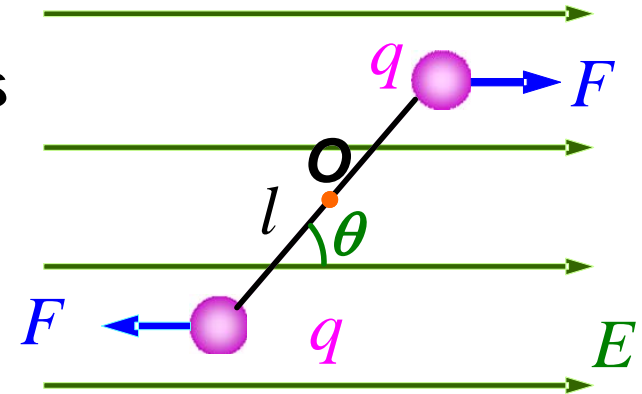


**Example: Find the potential energy of an electric dipole in a uniform external field.**



## Solution I:

The potential energy of a dipole is the sum of the potential energies of positive and negative charges in the field.



$$U = U_+ + U_- = qV(P_+) - qV(P_-)$$

$$= q[V(P_+) - V(P_-)]$$

$$= q \int_{P_+}^{P_-} \vec{E} \cdot d\vec{s} = q(-El \cos \theta) = -\vec{p} \cdot \vec{E}$$



# The Potential Energy of a Dipole in an External Field



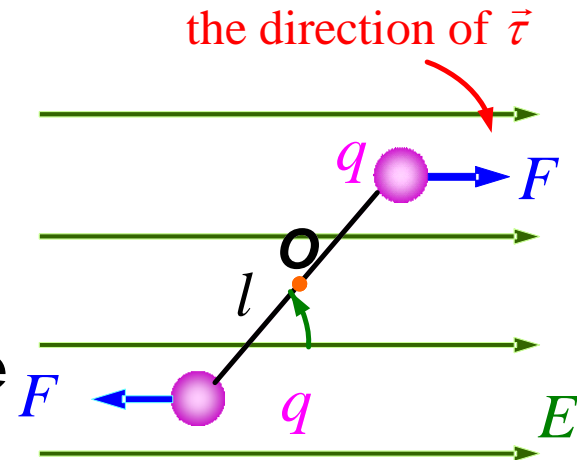
## Solution II: (P476)

$$\Delta U = U_f - U_i = -W_{if}$$

The work done on the dipole by the electric field to change the angle  $\theta$  from  $\theta_1$  to  $\theta_2$ :

$$\begin{aligned} W &= \int_{\theta_1}^{\theta_2} \tau d\theta = -pE \int_{\theta_1}^{\theta_2} \sin \theta d\theta \\ &= pE \cos \theta_2 - pE \cos \theta_1 \end{aligned}$$

The work done by a conservative force decreases the potential energy.



$$U_2 - U_1 = -W = (-pE \cos \theta_2) - (-pE \cos \theta_1)$$

$$U=0 \text{ when } \vec{p} \perp \vec{E} \quad \theta_1 = 90^\circ, \quad \cos \theta_1 = 0, \quad U = -pE \cos \theta$$

The vector description:

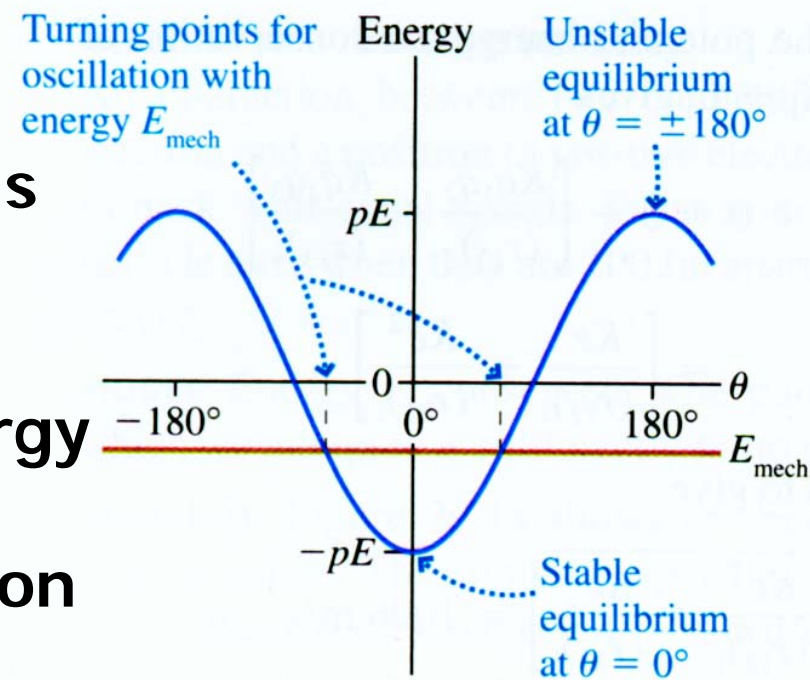
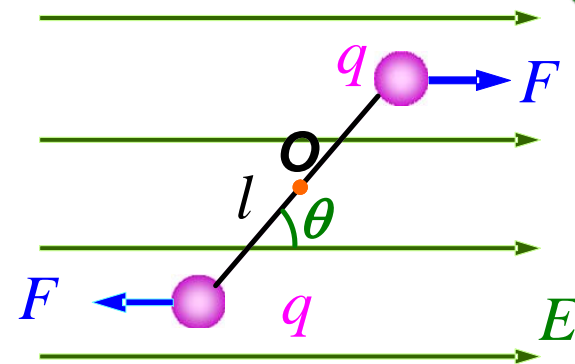
$$U = -\vec{p} \cdot \vec{E}$$

## The Potential Energy of a Dipole in an External Field



$$U = -\vec{p} \cdot \vec{E}$$

- The potential energy is minimum at  $\theta=0^\circ$ . This is the a point of **stable** equilibrium.
- The potential energy is maximum at  $\theta=\pm 180^\circ$ , which is at the point of **unstable** equilibrium.
- A dipole with mechanical energy  $E_{\text{mech}}$  will **oscillates** back and forth between turning points on either side of  $\theta=0^\circ$ .

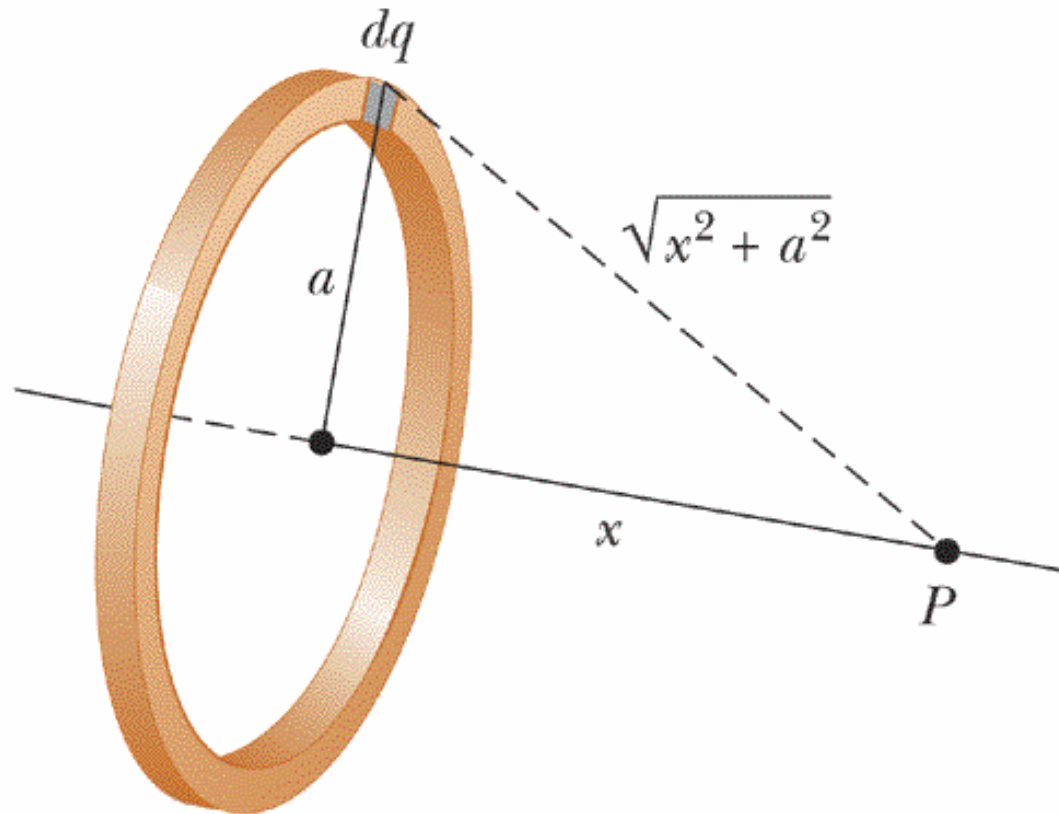


### Example (P510 Ex. 21-8)



The electric potential due to a uniformly charged ring

**Example:** Find the electric potential at a point **P** located on the axis of a uniformly charged ring of radius ***a*** and total charge ***Q***.



## Example



### The electric potential due to a uniformly charged ring

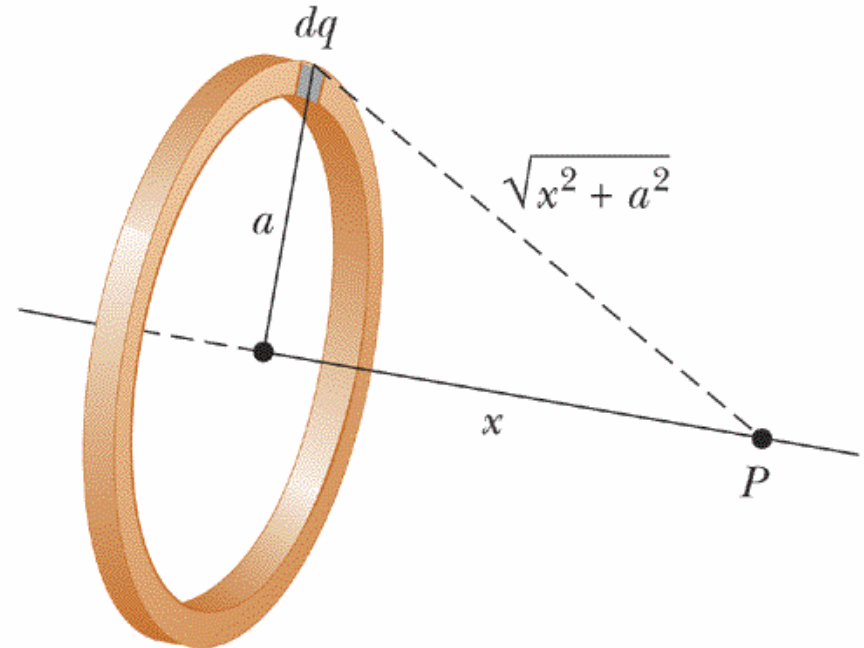
**Example:** Find the electric potential at a point P located on the axis of a uniformly charged ring of radius  $a$  and total charge  $Q$ .

**Solution:**

$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{\sqrt{x^2 + a^2}}$$

$$V = \frac{1}{4\pi\epsilon_0} \frac{1}{\sqrt{x^2 + a^2}} \int dq$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{\sqrt{x^2 + a^2}}$$

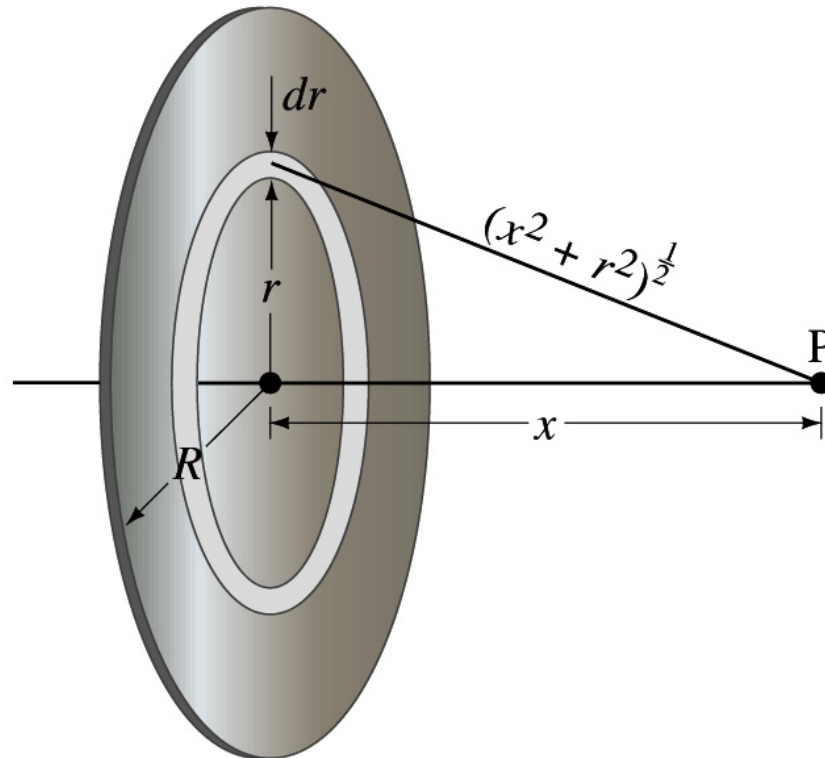


## Example (P510 Ex. 21-9)



The electric potential due to a uniformly charged disk

A thin flat disk, of radius  $R$ , carries a uniformly distributed charge  $Q$ . Determine the potential at a point P on the axis of the disk, a distance  $x$  from its center.



### Example (P510 Ex. 21-9)

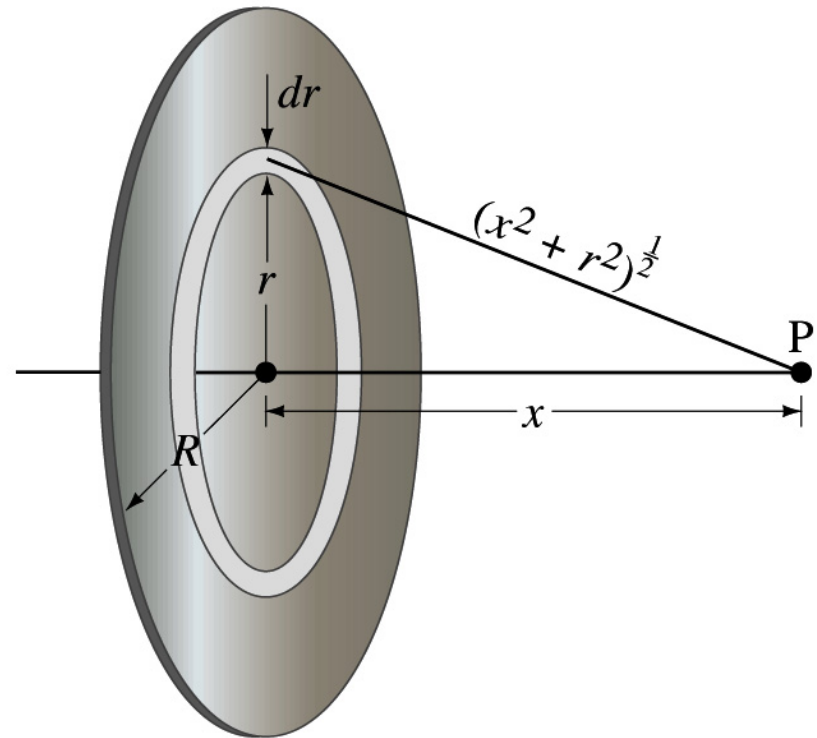


$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{\sqrt{x^2 + r^2}}$$

$$dA = (2\pi r)(dr)$$

$$dq = \frac{Q}{\pi R^2} dA = \frac{2Q}{R^2} r dr$$

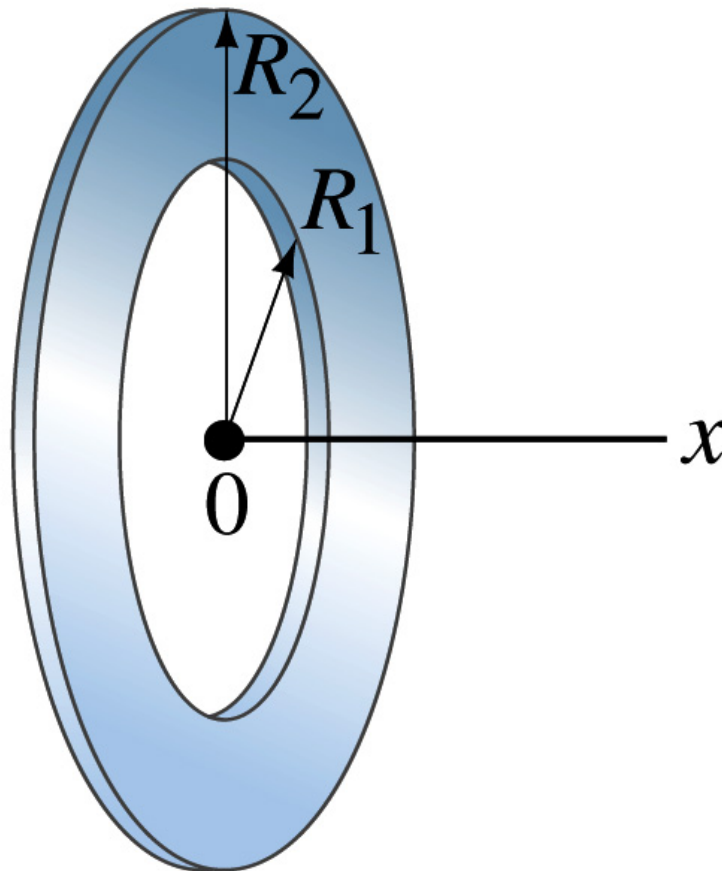
$$\begin{aligned} V &= \frac{1}{4\pi\epsilon_0} \int \frac{dq}{\sqrt{x^2 + r^2}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{2Q}{R^2} \int_0^R \frac{r dr}{\sqrt{x^2 + r^2}} \\ &= \frac{1}{2\pi\epsilon_0} \frac{Q}{R^2} (\sqrt{x^2 + R^2} - x) \end{aligned}$$



**Prob. 31** (Ch21 P520)



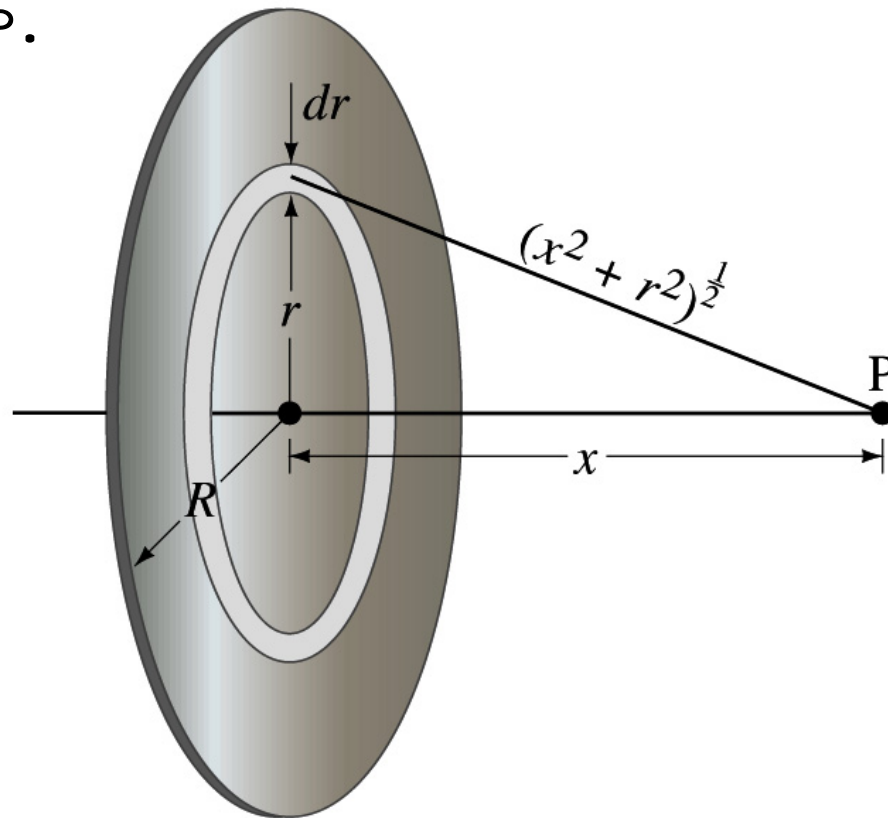
A flat ring of inner radius  $R_1$  and outer radius  $R_2$  carries a uniform surface charge density  $\sigma$ . Determine the electric potential at points along the  $x$  axis.



**Prob. 35 (Ch21 P521)**



Suppose the flat circular disk has a nonuniform surface charge density  $\sigma = ar^2$ , where  $r$  is measured from the center of the disk. Find the potential at points along the  $x$  axis, relative to  $V=0$  at  $x=\infty$ .





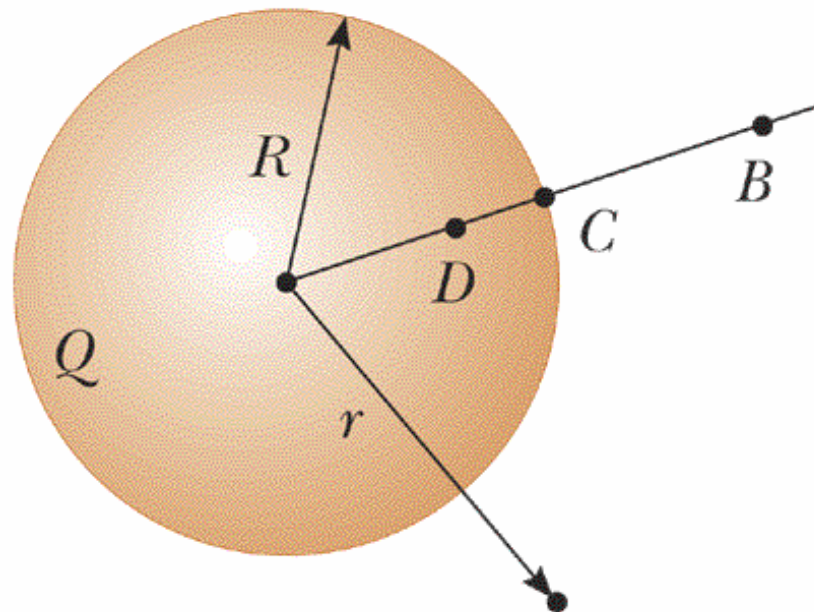
## Example



### The electric potential of a uniformly charged sphere

**Example:** An insulating solid sphere of radius  $R$  has a total charge  $Q$ , which is distributed uniformly throughout the volume of the sphere.

- (1) Find the electric potential at a point for  $r > R$ .
- (2) Find the electric potential at a point for  $r < R$ .



## Example

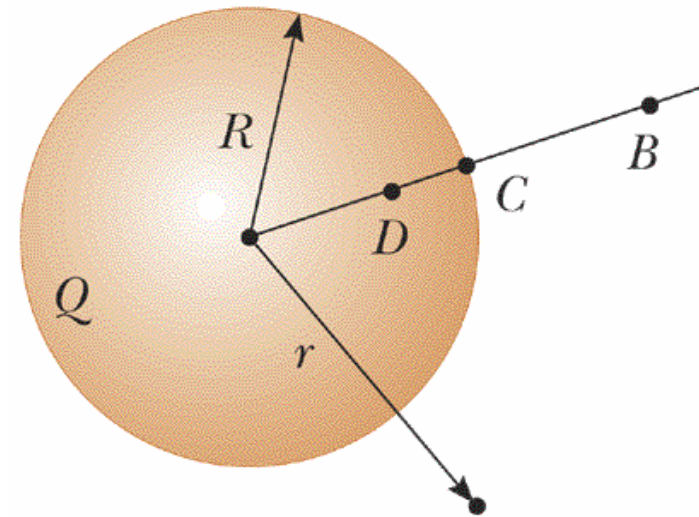


### The electric potential of a uniformly charged sphere

**Solution 1:** 
$$V = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r}$$

**Solution 2:** 
$$V_P = \int_P^\infty \vec{E} \cdot d\vec{s}$$

$$E = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} & \text{for } r > R \\ \frac{1}{4\pi\epsilon_0} \frac{r}{R^3} Q & \text{for } r < R \end{cases}$$



## Example – cont'd



**For  $r > R$**   $V_B = \int_r^\infty \vec{E} \cdot d\vec{s}$

$$= \frac{Q}{4\pi\epsilon_0} \int_r^\infty \frac{dr}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$$

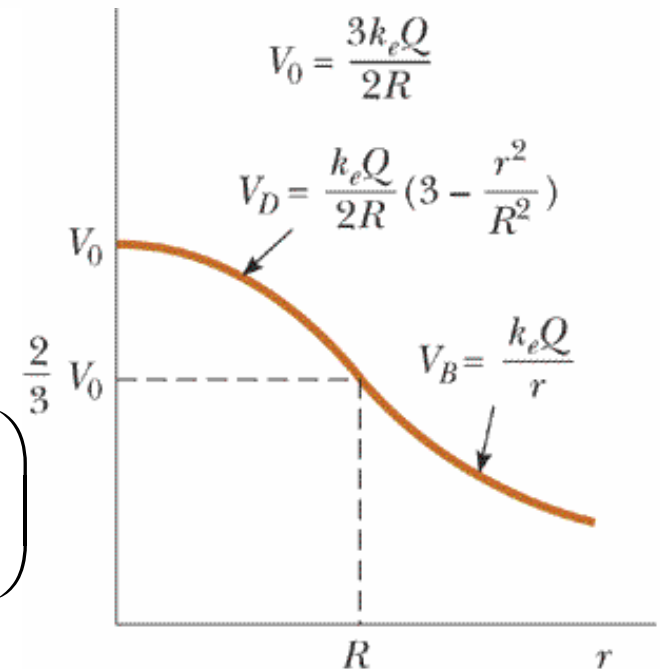
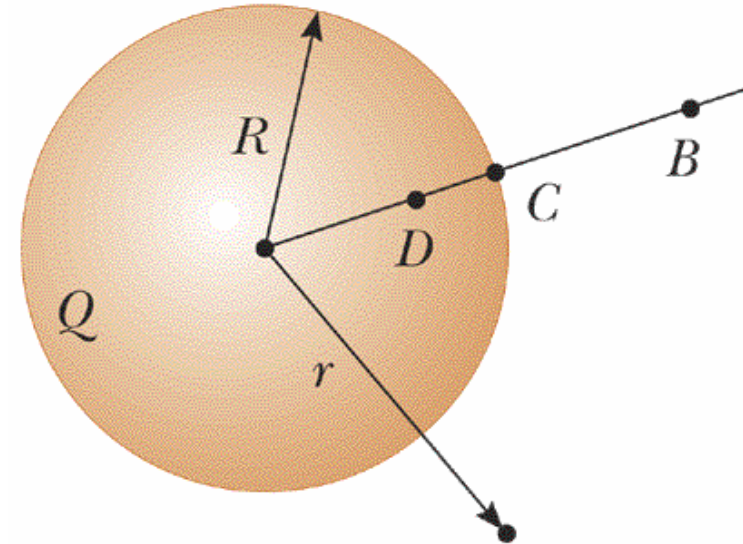
**For  $r < R$**

$$V_D = \int_r^R \vec{E} \cdot d\vec{s} + \int_R^\infty \vec{E} \cdot d\vec{s}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{R^3} \int_r^R r dr + \frac{Q}{4\pi\epsilon_0} \int_R^\infty \frac{dr}{r^2}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{2R^3} (R^2 - r^2) + \frac{1}{4\pi\epsilon_0} \frac{Q}{R}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q}{2R} \left( 3 - \frac{r^2}{R^2} \right) = \frac{Q}{8\pi\epsilon_0 R} \left( 3 - \frac{r^2}{R^2} \right)$$



**Ch21 Prob. 18, 34, 35, 43 (P520)**

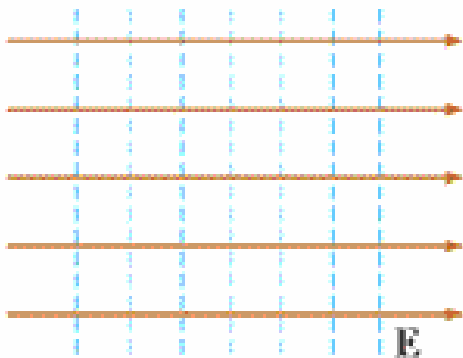
## § 3 Equipotential Surfaces

(P511, § 21-5)

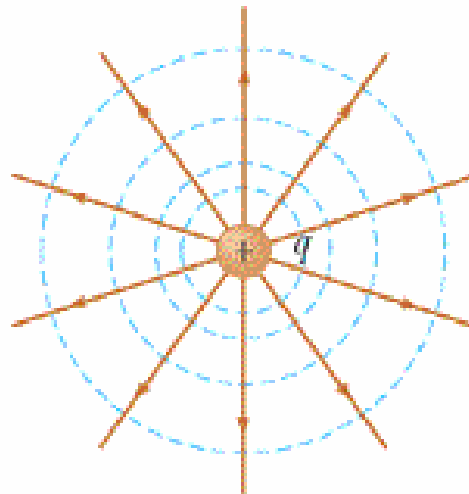


### ■ The equipotential surface

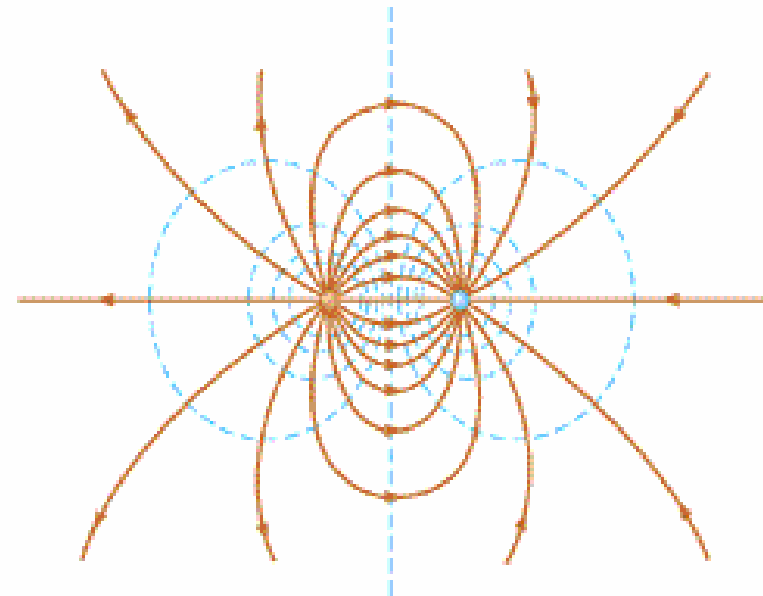
- ➡ An equipotential surface is a three-dimensional surface on which the electric potential  $V$  is the same at every point.



(a)



(b)



(c)



- The properties of the equipotential surface
  - If a test charge moves over an equipotential surface, the electric field can do no work on such a charge.

$$W_{ab} = -q_0 \Delta U = q_0 (U_a - U_b) = 0$$

- Field lines and equipotential surface are always mutually **perpendicular**.

A test charge  $q_0$  moves a distance  $d\vec{l}$  on an equipotential surface

$$dW = q_0 \vec{E} \cdot d\vec{l} = q_0 E \cos \theta dl = 0 \Rightarrow \vec{E} \perp d\vec{l}$$

- In regions where the magnitude of  $\vec{E}$  is **large**, the equipotential surface are **close** together.

## § 4 Potential Gradient



(P513 § 21-7)

$$\begin{aligned} -dV &= -\left(\frac{\partial V}{\partial x}dx + \frac{\partial V}{\partial y}dy + \frac{\partial V}{\partial z}dz\right) \\ &= \vec{E} \cdot d\vec{s} = E_x dx + E_y dy + E_z dz \end{aligned}$$

$$E_x = -\frac{\partial V}{\partial x}, \quad E_y = -\frac{\partial V}{\partial y}, \quad E_z = -\frac{\partial V}{\partial z}$$

$$\vec{E} = -\left(\hat{i}\frac{\partial V}{\partial x} + \hat{j}\frac{\partial V}{\partial y} + \hat{k}\frac{\partial V}{\partial z}\right) = -\left(\hat{i}\frac{\partial}{\partial x} + \hat{j}\frac{\partial}{\partial y} + \hat{k}\frac{\partial}{\partial z}\right)V$$

$$\boxed{\vec{E} = -\vec{\nabla}V}$$

$\vec{E}$  is the negative of the gradient of  $V$ .

## The meaning of the gradient



- ➡ Make a displacement  $d\vec{s}$  from one equipotential surface to the adjacent surface

$$-dV = \vec{E} \cdot d\vec{s} = E \cos \theta ds$$

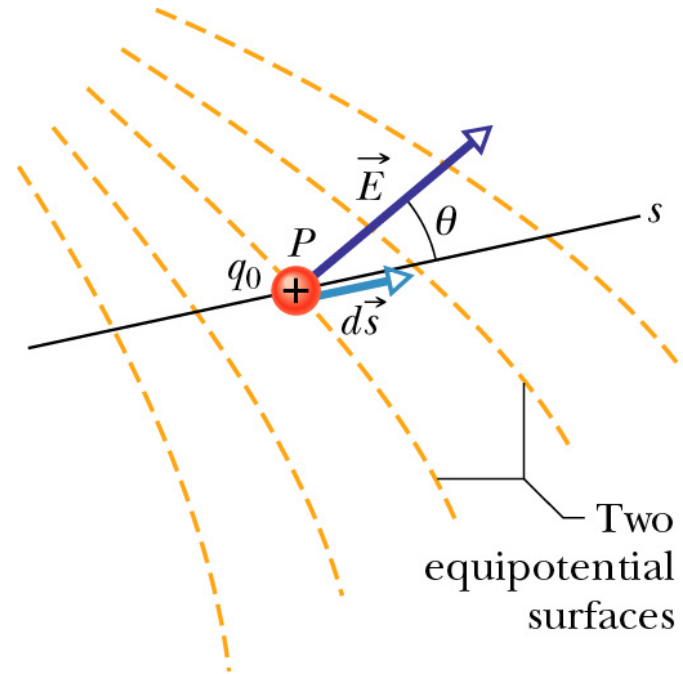
$$E \cos \theta = -\frac{dV}{ds}$$

$$E_s = -\frac{\partial V}{\partial s}$$

- ➡ The component of  $\vec{E}$  in any direction is the negative of the rate of change of the electric potential with distance in that direction.

- ➡ Take the  $s$  axis to be, in turn,  $x$ ,  $y$ , and  $z$  axis, we get the  $x$ ,  $y$ ,  $z$  components of  $\vec{E}$  at any point are

$$E_x = -\frac{\partial V}{\partial x}, \quad E_y = -\frac{\partial V}{\partial y}, \quad E_z = -\frac{\partial V}{\partial z}$$







- By Coulomb's law:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \hat{r}$$

➡ The most general method.

- By Gauss' law:

$$\oiint_S \vec{E} \cdot d\vec{A} = \frac{q_{inside}}{\epsilon_0}$$

➡ If charge distribution possesses a high degree of symmetry

- By gradient of V:

$$\vec{E} = -\vec{\nabla}V$$

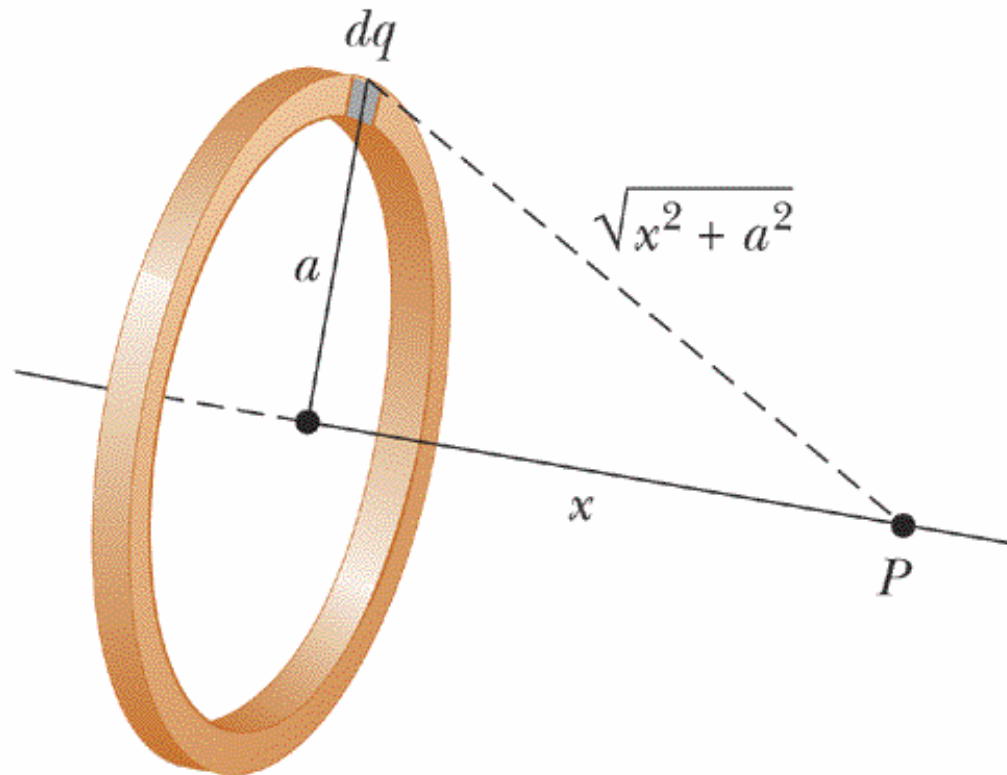
➡ If the potential is easy to obtain.

## Example



### A uniformly charged ring (P514 Ex. 21-11)

Find the electric field at a point  $P$  located on the axis of a uniformly charged ring of radius  $a$  and total charge  $Q$ .



## Example



**Solution: based on the electric potential:**

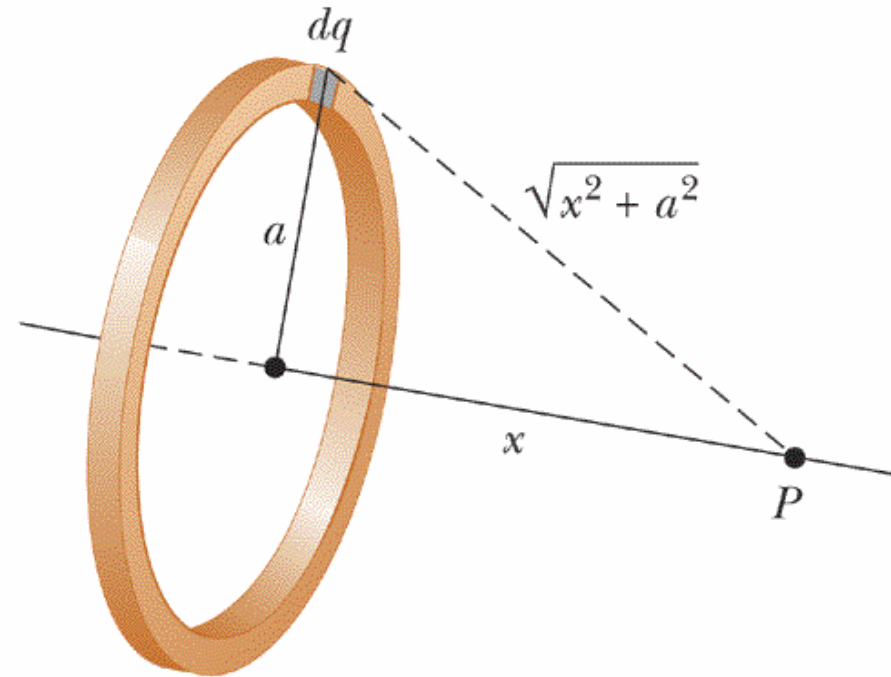
$$V = \frac{1}{4\pi\epsilon_0} \frac{Q}{\sqrt{x^2 + a^2}}$$

$$E = -\frac{\partial V}{\partial x}$$

$$= -\frac{Q}{4\pi\epsilon_0} \frac{d}{dx} (x^2 + a^2)^{-1/2}$$

$$= -\frac{Q}{4\pi\epsilon_0} \left( -\frac{1}{2} \right) (x^2 + a^2)^{-3/2} (2x)$$

$$= \frac{1}{4\pi\epsilon_0} \frac{xQ}{(x^2 + a^2)^{3/2}}$$

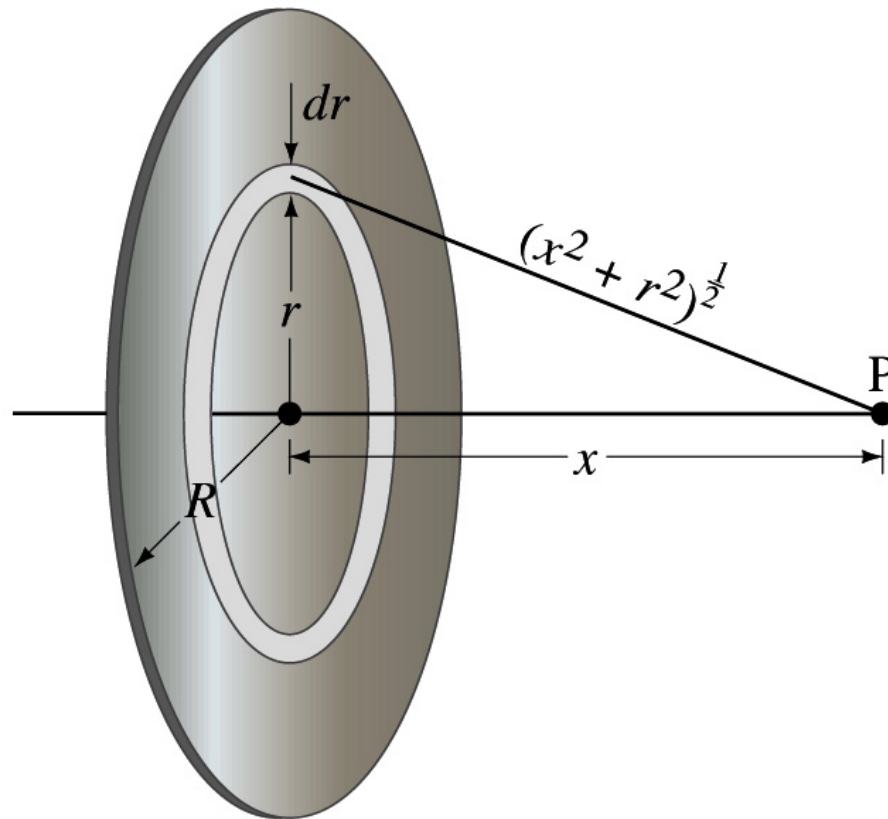


## Example



### A uniformly charged disk (P514 Ex. 21-11)

Find the electric field at a point  $P$  located on the axis of a uniformly charged disk of radius  $R$  and total charge  $Q$ .



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### A uniformly charged disk (P514 Ex. 21-11)

$$V = \frac{1}{2\pi\epsilon_0} \frac{Q}{R^2} (\sqrt{x^2 + R^2} - x)$$

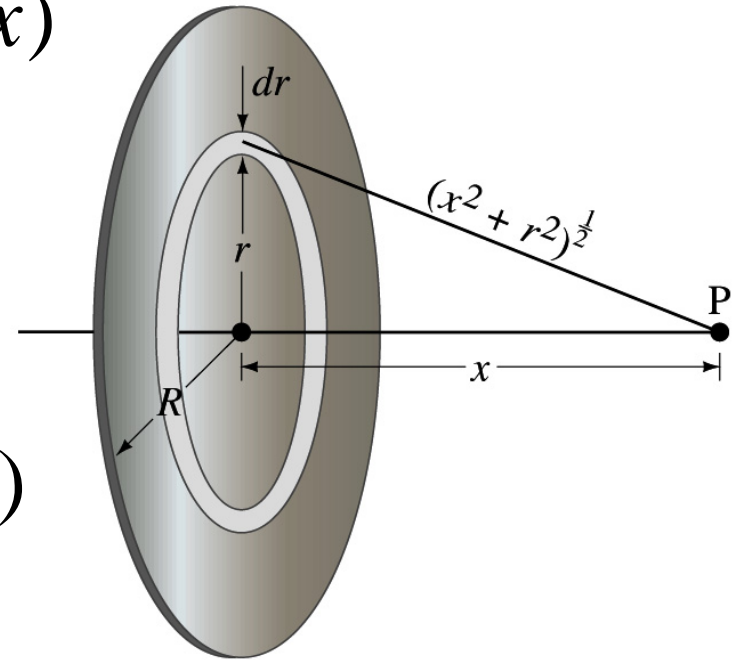
$$E_x = -\frac{\partial V}{\partial x}$$

$$= \frac{Q}{2\pi\epsilon_0 R^2} \left(1 - \frac{x}{\sqrt{x^2 + R^2}}\right)$$

$$E_y = E_z = 0$$

$$x \ll R,$$

$$E_x \approx \frac{Q}{2\pi\epsilon_0 R^2} = \frac{\sigma}{2\epsilon_0}$$



The electric potential in a region of space varies as

$$V = \frac{ay}{b^2 + y^2}$$

Determine  $\vec{E}$

From the spatial dependence of the electric potential,  $V(x, y, z) = ay/(b^2 + y^2)$ , we find the components of the electric field from the partial derivatives of  $V$ :

$$E_x = -\partial V / \partial x = 0;$$

$$E_y = -\partial V / \partial y = -a/(b^2 + y^2) - ay(-2y)/(b^2 + y^2)^2 = a(y^2 - b^2)/(b^2 + y^2)^2.$$

$$E_z = -\partial V / \partial z = 0.$$

We can write the electric field:

$$\vec{E} = a(y^2 - b^2)/(b^2 + y^2)^2 \hat{j}.$$

## **Ch21 Prob. 38, 47 (P521)**