# Computational Methods for Astrophysical Applications

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#### References

#### Books

- Riemann solvers and numerical methods for fluid dynamics A practical introduction, by Toro
- Numerical methods for conservation laws, by Leveque
- Finite volume methods for hyperbolic problems, by Leveque

#### Courses

- Rony Keppens & Jon Sundqvist 2016-2017
- Numerical PDE Techniques for Scientists and Engineers, by Dinshaw Balsara

#### Schools

- Les Houches
- Numerical techniques in MHD simulations, Köln University

**Lesson 1: Hyperbolic Partial Differential Equations** 

- Conservation laws
   Integral and differential forms
   Examples
  - Hyperbolic PDE

    Matrix formulation of conservation laws
    Time advance
    Linear advective equations
    The Riemann problem

- Conservation laws
  - Integral and differential forms
  - Examples
- 2 Hyperbolic PDE
  - System of (1D) 1st order linear PDE
  - Matrix formulation of conservation laws
  - Time advance
  - Linear advection equations
  - The Riemann problem

## Integral form

- Continuous medium hypothesis: from actual particles to particles of fluid
  - $\Rightarrow$  distribution function and integration of Boltzmann equation : the closure problem
  - ⇒ analogy with moments of the specific photon intensity (see Prof. Sundqvist's course)
- Consider any volume element dV of a fluid
  - $\Rightarrow$  XXX bilan XXX over dt of scalar quantity : inflow/outflow, sinks/sources
    - ⇒ for the momentum : tensor formulation and dyadic product
    - $\Rightarrow$  for energy
- Shocks and jump conditions : Rankine-Hugoniot

### Conservative form

• Provided the variables  $X(\mathbf{r},t)$  are XXX differentiable XXX

$$\partial_t U + \nabla \cdot \underbrace{[F(U)]}_{\text{fluxes}} = \underbrace{S(\mathbf{r}, \mathbf{v}, t)}_{\text{sources/sinks}}$$

- ⇒ Matrix formulation
- Properties of this differential form
  - ⇒ Coordinate / dimension independent formulation
  - $\Rightarrow$  Given inital and boundary conditions, can provide a general solution
  - $\Rightarrow$  Conservative form : Green-Ostrogradsky (or Gauss) law and Eulerian approach
    - $\Rightarrow$  Conservative variables :  $\rho$ ,  $\rho$ **v** and e

#### Primitive form

Lagrangian approach XXX Picture of a parachute XXX

$$D_t\left(\cdot\right) = \partial_t\left(\cdot\right) + \mathbf{v}\nabla\left(\cdot\right)$$

Continuity equation

$$D_t(\rho) = -\rho \nabla \mathbf{v}$$

- $\Rightarrow$  incompressible fluid  $\neq$  flow
- Navier-Stokes equation

$$\rho D_t(\mathbf{v}) = -\nabla P$$

- Energy equation
  - ⇒ internal and mechanical energy
  - ⇒ link w/ 1<sup>st</sup> principle of Thermodynamics
  - ⇒ entropy formulation

### Closure relation

Equation-of-state of an ideal gas

$$u=\frac{P}{\gamma-1}$$

- $\Rightarrow$  adiabatic index  $\gamma$
- A classic way-around : the polytropic assumption

$$P \propto \rho^{\alpha}$$

- $\Rightarrow$  polytropic index  $\alpha$
- ⇒ the isentropic case

## The linear advection equation

Consider the 1D continuity equation w/ constant & uniform speed v

$$\partial_t \rho + \mathbf{v} \partial_{\mathbf{x}} \rho = \mathbf{0}$$

- $\Rightarrow$  initial conditions  $\rho(x, t = 0) = \rho_0(x)$
- Analytic solution

$$\rho(\mathbf{x},t)=\rho_0(\mathbf{x}-\mathbf{v}t)$$

Fourier analysis

$$\rho\left(x,t\right)=r\left(k,t\right)e^{ikx}$$

#### XXX

- ⇒ non-dissipative
- ⇒ non-dispersive
- $\Rightarrow$  see also DB1#32-33
- Analytic preliminary analysis => mathematical properties of the PDE => appropriate numerical scheme
- Numerical solution El Mellah & Sundavist (KU Leuven)

## **Euler equations**

• Consider the viscous tensor  $\tau$ 

$$au = \mu \left[ \left( \mathbf{v} + \mathbf{v}^t 
ight) - rac{2}{3} \left( 
abla \cdot \mathbf{v} 
ot\!\!\!\!/ 
ight) 
ight]$$

- ⇒ see Landau Lifschitz section XXX for hypothesis & shape
- $\Rightarrow \mu$  dynamic viscosity in XXX
- ⇒ conservative formulation
- $\Rightarrow$  Dimensionless form, Reynolds number and singular limit (Lagrange paradox? XXX)

## Ideal Magneto-Hydrodynamics equations

Ideal MHD fundamental hypothesis

$$\mathbf{E} + \mathbf{j} \wedge \mathbf{B} = \mathbf{0}$$

- ⇒ **E** secondary variable
- Maxwell-Faraday and Maxwell-Ampère give the induction equation XXX :

$$\nabla \wedge \mathbf{B} - \partial_t ((\nabla \wedge \mathbf{B}) \wedge \mathbf{B}) = \mathbf{0}$$

- ⇒ magnetic pressure and XXX Prandtl XXX number
- ⇒ magnetic tension
- Frozen-flux theorem

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## Linear, non-linear, quasi-linear

Homogeneous

$$\partial_t \mathbf{U} + \mathbf{A} (\mathbf{U}, x, t) \, \partial_x \mathbf{U} = \mathbf{0} \tag{1}$$

XXX Choice diagram XXX

- Non-homogeneous : ... +  $\mathbf{B}(\mathbf{U}, x, t) = \mathbf{0}$ 
  - $\Rightarrow$  still linear if  $\mathbf{B} \propto \mathbf{U}$

Time and space dependent linear advection equation? Burgers equation? Conservation equations?

Toro 2.1 p.42

Systematic way to solve quasi-linear PDE?

## Eigenvalues & eigenvectors

Consider the following system of linear constant coeff. PDE:

$$\begin{aligned}
\partial_t \rho_1 + \partial \left( \rho_1 a_{11} + \rho_2 a_{12} \right) &= 0 \\
\partial_t \rho_2 + \partial \left( \rho_1 a_{21} + \rho_2 a_{22} \right) &= 0
\end{aligned} \tag{2}$$

- ⇒ Linear? Homogeneous?
- $\Rightarrow$  Matrix form  $\partial_t \mathbf{U} + A \partial_x \mathbf{U}$ ?
- ⇒ A diagonalisable?
- $\Rightarrow$  Diagonal matrix *D*?
- $\Rightarrow$  Eigenvector matrix K?

## Hyperbolic, parabolic and elliptic PDE

Elliptic <=> complex eigenvalues

⇒ ex: Cauchy-Riemann equations (2D steady)

$$\partial_x u - \partial_y v = 0$$

$$\partial_x v + \partial_y u = 0$$

Hyperbolic <=> real eigenvalues and linearly independent eigenvectors

- $\Rightarrow$  eigenvalues  $\lambda_i$  as wave speeds
- ⇒ Was (2) hyperbolic?

Systematic way to solve hyperbolic quasi-linear PDE?

## General solution of quasi-linear hyperbolic PDE

Toro 2.3.1

 $K^{-1} \times (1) =$  canonical form :

$$\partial_t \mathbf{V} + D\partial_x \mathbf{V} = 0$$

with the characteristic variables or waves:

$$V = K^{-1}\mathbf{U} \tag{3}$$

Given initial conditions  $\mathbf{U}(x,0)$ ,  $\mathbf{V}$  basis of decoupled linear advection solutions :

$$V_i(x,t) = V_i(x - \lambda_i t, 0)$$

whose linear superposition via (3) provides, the general solution

- $\Rightarrow$  In (2), what if  $a_{11}$  and  $a_{22} = 0$ ? What if  $a_{12} = a_{21}$ ?
- ⇒ Exo: section 6.3 of Leveque

## Linearization of non-linear PDE

Leveque 6.4

### Flux Jacobian matrix

 In matrix form, the 1D conservation laws without source terms are written:

$$\partial_t \mathbf{U} + \partial_x \left[ XXX\mathbf{F} \left( \mathbf{U} \right) \right] = 0$$

where

$$\mathbf{U} = (\rho, \rho \mathbf{v}, \mathbf{e})$$

and the vector of fluxes is:

$$F(\mathbf{U}) = \left(egin{array}{c} 
ho \mathbf{v} & \mathbf{v} + P \mathbb{W} \ (e + P) \mathbf{v} \end{array}
ight)$$

#### Flux Jacobian matrix

Chain rule => quasi-linear form

$$\partial_t \mathbf{U} + \mathbf{AXXX}(\mathbf{U}) \, \partial_x \mathbf{U} = 0 \tag{4}$$

with **AXXX** (**U**) the flux Jacobian matrix :

$$AXXX = \partial F/\partial U$$

⇒ conservation equations are quasi-linear

# **Explicit**

# **Implicit**

## **Properties**

## Characteristics

## Stencils, domain of dependence, range of influence

# The Riemann problem