A Note on Elliptic Coordinates

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Explicit Equations to Transform from Cartesian to Elliptic Coordinates

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Abstract: Explicit equations are obtained to convert Cartesian coordinates to elliptic coordinates, based on which a function in elliptic coordinates can be readily mapped in physical space. Application to Kirchhoff vortex shows that its elliptical core induces two counter-rotating irrotational eddies.

Keywords: Elliptic Coordinates, Cartesian Coordinates, Kirchhoff Vortex

1. Introduction

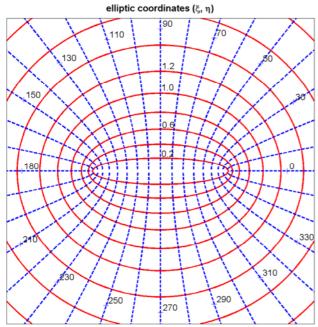


Figure 1. Coordinate lines for an elliptic coordinate system with a = 1, b = 0.5.

The elliptical coordinate system (ξ, η) as a two-dimensional orthogonal coordinate system has many dynamical and engineering applications, such as Kirchhoff vortex [1], insect aerodynamics [2], hydrodynamic wave

diffraction [3], and theoretical physics [4]. Its coordinate lines are confocal ellipses and hyperbolae and the transformation from elliptic to Cartesian coordinates is given by

$$x = c \cosh(\xi) \cos(\eta)$$

$$y = c \sinh(\xi) \sin(\eta)$$

$$c^{2} = a^{2} - b^{2}$$

$$\xi \ge 0, \quad 0 \le \eta < 2\pi$$
(1)

where a and b denote the semi-major and semi-minor axes of the ellipse and c is the elliptical eccentricity (Figure 1).

In the meantime, explicit equations to transform from Cartesian to elliptic coordinates have not been found in the existing literature [5, 6, 7]. Such a conversion relation would be useful in mapping an elliptic-coordinate solution, for example the flow field of Kirchhoff vortex, in physical space.

2. Cartesian to Elliptic Coordinates

In order to invert the functional relation (1), we first eliminate ξ and have

$$\frac{x^2}{\cos^2(\eta)} - \frac{y^2}{\sin^2(\eta)} = c^2$$

which means curves of constant η are hyperbolae. The focus distance is c and the eccentricity is $e = \sec(\eta)$.

Let
$$p = \sin^2(\eta)$$
, we have

$$\frac{x^2}{1-p} - \frac{y^2}{p} = c^2$$

which becomes

$$c^{2}p^{2} + (x^{2} + y^{2} - c^{2})p - y^{2} = 0$$
 (2)

Then eliminating η from (1) we have

$$\frac{x^2}{\cosh^2(\xi)} + \frac{y^2}{\sinh^2(\xi)} = c^2$$

It shows that curves of constant ξ are ellipses. The focus distance is c and the eccentricity is $e = \cosh^{-1}(\xi)$.

Let $q = -\sinh^2(\xi)$, we have

$$\frac{x^2}{1-q} - \frac{y^2}{q} = c^2$$

which leads to

$$c^{2}q^{2} + (x^{2} + y^{2} - c^{2})q - y^{2} = 0$$
 (3)

It is essentially the same as (2). Therefore (p,q) constitute the two roots of a quadratic equation. Since $0 \le p \le 1$, $q \le 0$, we have $p \ge q$, and the two roots are

$$p = \frac{-B + \sqrt{B^2 + 4c^2y^2}}{2c^2}, \quad q = \frac{-B - \sqrt{B^2 + 4c^2y^2}}{2c^2} \quad (4)$$

in which $B = x^2 + y^2 - c^2$.

From the definition of p we obtain

$$\eta_0 = \arcsin(\sqrt{p})$$
 (5)

It has four cases depending on which quadrant the Cartesian point (x, y) is located, i.e.,

$$\eta = \eta_0, x \ge 0, y \ge 0
\eta = \pi - \eta_0, x < 0, y \ge 0
\eta = \pi + \eta_0, x \le 0, y < 0
\eta = 2\pi - \eta_0, x > 0, y < 0$$
(6)

Based on the definition of q, we can solve ξ from quadratic equation

$$e^{4\xi} + (4q-2)e^{2\xi} + 1 = 0$$

which has two roots

$$e^{2\xi} = 1 - 2q \pm 2\sqrt{q^2 - q}$$

Since $q \le 0$, both roots are real and denoted as (ξ_1, ξ_2) . They clearly satisfy $e^{2\xi_1} \cdot e^{2\xi_2} = 1$, which leads to $\xi_2 = -\xi_1 < 0$. Because in elliptical coordinates only non-negative ξ value is considered, we obtain

$$\xi = \frac{1}{2} \ln \left(1 - 2q + 2\sqrt{q^2 - q} \right) \tag{7}$$

Eqs. (4-7) are explicit equations to derive elliptic coordinates from Cartesian grid. They can easily be realized via computation software such as Matlab.

3. Application to Kirchhoff Vortex

Kirchhoff vortex is a rotating elliptical region of uniform vorticity ω embedded in an irrotational ideal fluid [8]. It is the simplest example of non-smooth weak solutions to the Euler equations and has wide application in vortex dynamics [9-11]. It has a discontinuity of vorticity across its elliptical boundary

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

where a and b are semi-major and semi-minor axes, respectively. The boundary corresponds to contour $\xi = \operatorname{arcsinh}(b/c)$, and the vortex rotates with constant angular velocity

$$\Omega = \frac{ab}{(a+b)^2}\omega$$

around the origin. From Act. 159 of Lamb [1], the streamfunction outside the core is

$$\psi = \frac{1}{4}\Omega(a+b)^2 e^{-2\xi} \cos(2\eta) + \frac{1}{2}\omega ab\xi$$
 (8)

The instantaneous streamlines in a unsteady flow are given by the curves $\psi = const$. In a rotating frame with angular velocity Ω , Kirchhoff vortex looks steady and its streamfunction outside the elliptic core is related to the inertial-frame streamfunction (8) by

$$\psi_{R} = \psi - \frac{1}{2}\Omega[x(\xi,\eta)^{2} + y(\xi,\eta)^{2}]$$
 (9)

From (9) it is straightforward to make a conformal mapping of steady streamfunction on a uniform mesh of elliptic coordinates (Figure 2).

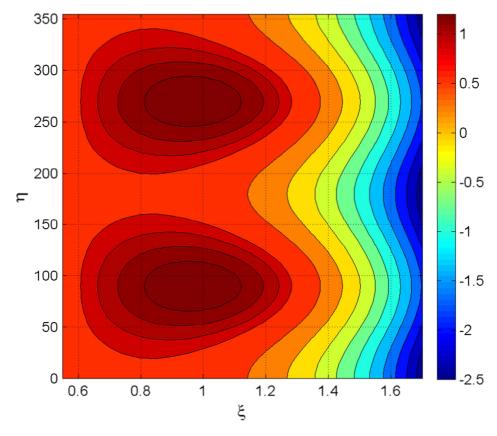


Figure 2. Mapping of streamfunction (9) on a uniform mesh of (ξ, η) , with $a = 1, b = 0.5, \omega = 10$.

In order to view the flow field in physical space, we use the conversion equations (4-7) to plot streamfunction (9) on a uniform Cartesian mesh in Figure 3, which shows the elliptical core of Kirchhoff vortex induces two irrotational eddies that rotate in opposite directions.

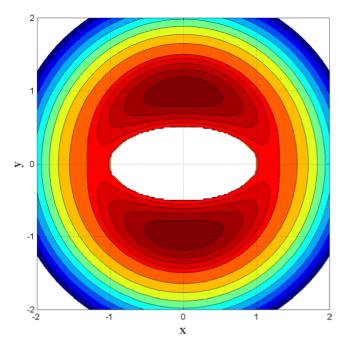


Figure 3. Mapping of streamfunction (9) on a uniform mesh of xy-plane, with $a=1,b=0.5,\omega=10$. The boundary of vortex core corresponds to $\xi=0.55$.

4. Conclusion

This study obtains explicit equations that convert Cartesian coordinates to elliptic coordinates. The conversion relation can be easily realized with computer software and used to map a known function of elliptic-coordinates, such as the streamfunction of Kirchhoff vortex, on a uniform Cartesian mesh.

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Matlab mfile scripts

```
function [zeta,eta]=xy2el(x,y,a,b)
%-----
% Conversion from Cartesian mesh (x,y) to elliptic coordinates (zeta,eta)
% zeta (>=0), eta [0 2pi)
% a -- semi-major axis
% b -- semi-minor axis
% c -- elliptic eccentricity
% (x,y,zeta,eta) same-size matrix
% Author: Che Sun (CAS, 2016)
% Source doi:10.11648/j.mma.20170204.12
% Note 1: An early code error was pointed out by Dr. Vinod Varghese.
% Note 2: A potential round-off error at x=0, as pointed out by Dr. Peter Gangl,
        is avoided by requiring root p<=1.
%
%-----
c2=(a^2-b^2); c=2*c2;
x2=x.^2; y2=y.^2;
B=x2+y2-c2;
del2=B.^2+2*c*y2; del=sqrt(del2);
p=(-B+del)/c; iv=find(p>1); p(iv)=1;
p=sqrt(p); et0=asin(p); eta=x;
i=find(x>=0&y>=0); eta(i)=et0(i);
i=find(x<0&y>=0); eta(i)=pi-et0(i);
i=find(x<=0&y<0); eta(i)=pi+et0(i);
i=find(x>0&y<0); eta(i)=2*pi-et0(i);
q=(-B-del)/c;
del2=q.^2-q; del=sqrt(del2);
zeta=log(1-2*q+2*del)/2;% only keep the positive root
function [x,y]=el2xy(zeta,eta,a,b)
%-----
% Conversion from elliptic coordinates (zeta,eta) to Cartesian coordinates
% zeta (>=0), eta [0 2pi)
% a -- semi-major axis
% b -- semi-minor axis
% c -- elliptic eccentricity
% (x,y,zeta,eta) same-size matrix
% Author: Che Sun (CAS, 2016)
% Source doi:10.11648/j.mma.20170204.12
%-----
c2=a^2-b^2; c=sqrt(c2);
x=c*cosh(zeta).*cos(eta);
y=c*sinh(zeta).*sin(eta);
```