

# Time Domain Complex-valued ViscoElastic Wave Equation in an Isotropic, Homogenous Medium

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The idea presented here is based on paper:

Yang, Jidong, and Hejun Zhu. "A Time-Domain Complex Valued Wave Equation for Modeling Visco-Acoustic Wave Propagation." *Geophysical journal international* 215.2 (2018): 1064-1079.

The tensor of elastic moduli for an isotropic medium is given as:

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \sigma_{xy} \\ \sigma_{yz} \\ \sigma_{zx} \end{bmatrix} = \begin{bmatrix} \lambda + 2\mu & \lambda & \lambda & 0 & 0 & 0 \\ \lambda & \lambda + 2\mu & \lambda & 0 & 0 & 0 \\ \lambda & \lambda & \lambda + 2\mu & 0 & 0 & 0 \\ 0 & 0 & 0 & \mu & 0 & 0 \\ 0 & 0 & 0 & 0 & \mu & 0 \\ 0 & 0 & 0 & 0 & 0 & \mu \end{bmatrix} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \varepsilon_{zz} \\ 2\varepsilon_{xy} \\ 2\varepsilon_{yz} \\ 2\varepsilon_{zx} \end{bmatrix}$$

The 3D equations of motion for an isotropic linear-elastic medium can then be written as

$$\begin{aligned} \rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} \\ \rho \frac{\partial^2 u_y}{\partial t^2} &= \frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{yz}}{\partial z} \\ \rho \frac{\partial^2 u_z}{\partial t^2} &= \frac{\partial \sigma_{zx}}{\partial x} + \frac{\partial \sigma_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \\ \sigma_{xx} &= \lambda(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz}) + 2\mu\epsilon_{xx} \\ \sigma_{yy} &= \lambda(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz}) + 2\mu\epsilon_{yy} \\ \sigma_{zz} &= \lambda(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz}) + 2\mu\epsilon_{zz} \\ \sigma_{xy} &= 2\mu\epsilon_{xy} \\ \sigma_{yz} &= 2\mu\epsilon_{yz} \\ \sigma_{zx} &= 2\mu\epsilon_{zx} \end{aligned} \tag{1}$$

$$\begin{aligned} \mu &= v_s^2 \rho = \frac{v_p^2 \rho}{3} \\ \lambda &= v_p^2 \rho - 2\mu \end{aligned} \tag{2}$$

$$\rho \frac{\partial^2 u_x}{\partial t^2} = \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} \tag{3}$$

$$\begin{aligned}
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[\lambda(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz}) + 2\mu\epsilon_{xx}]}{\partial x} + \frac{\partial[2\mu\epsilon_{xy}]}{\partial y} + \frac{\partial[2\mu\epsilon_{xz}]}{\partial z} \\
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[\lambda(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})]}{\partial x} + \frac{\partial[2\mu\epsilon_{xx}]}{\partial x} + \frac{\partial[2\mu\epsilon_{xy}]}{\partial y} + \frac{\partial[2\mu\epsilon_{xz}]}{\partial z} \\
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[(v_p^2 \rho - 2\mu)(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})]}{\partial x} + \frac{\partial[2\mu\epsilon_{xx}]}{\partial x} + \frac{\partial[2\mu\epsilon_{xy}]}{\partial y} + \frac{\partial[2\mu\epsilon_{xz}]}{\partial z} \\
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[v_p^2 \rho \epsilon_{xx} + v_p^2 \rho \epsilon_{yy} + v_p^2 \rho \epsilon_{zz} - 2\mu\epsilon_{xx} - 2\mu\epsilon_{yy} - 2\mu\epsilon_{zz}]}{\partial x} + \frac{\partial[2\mu\epsilon_{xx}]}{\partial x} + \frac{\partial[2\mu\epsilon_{xy}]}{\partial y} + \frac{\partial[2\mu\epsilon_{xz}]}{\partial z} \\
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[v_p^2 \rho(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})]}{\partial x} + \frac{\partial(-2\mu\epsilon_{yy} - 2\mu\epsilon_{zz})}{\partial x} + \frac{\partial[2\mu\epsilon_{xy}]}{\partial y} + \frac{\partial[2\mu\epsilon_{xz}]}{\partial z} \\
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[v_p^2 \rho(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})]}{\partial x} + 2\mu \left( -\frac{\partial(\epsilon_{yy} + \epsilon_{zz})}{\partial x} + \frac{\partial\epsilon_{xy}}{\partial y} + \frac{\partial\epsilon_{xz}}{\partial z} \right) \\
\rho \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial[v_p^2 \rho(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})]}{\partial x} + 2\frac{v_p^2 \rho}{3} \left( -\frac{\partial(\epsilon_{yy} + \epsilon_{zz})}{\partial x} + \frac{\partial\epsilon_{xy}}{\partial y} + \frac{\partial\epsilon_{xz}}{\partial z} \right) \\
\frac{1}{v_p^2} \frac{\partial^2 u_x}{\partial t^2} &= \frac{\partial(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})}{\partial x} + \frac{2}{3} \left( -\frac{\partial(\epsilon_{yy} + \epsilon_{zz})}{\partial x} + \frac{\partial\epsilon_{xy}}{\partial y} + \frac{\partial\epsilon_{xz}}{\partial z} \right) \\
\frac{1}{v_p^2} \frac{\partial^2 u_x}{\partial t^2} &= \frac{1}{3} \left( \frac{\partial(\epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz})}{\partial x} + 2 \left( \frac{\partial\epsilon_{xx}}{\partial x} + \frac{\partial\epsilon_{xy}}{\partial y} + \frac{\partial\epsilon_{xz}}{\partial z} \right) \right) \\
\boxed{\frac{1}{v_p^2} \frac{\partial^2 u_x}{\partial t^2} &= \left[ \frac{\partial}{\partial x} \left( \epsilon_{xx} + \frac{1}{3}\epsilon_{yy} + \frac{1}{3}\epsilon_{zz} \right) \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{xy}}{\partial y} \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{xz}}{\partial z} \right]} & (4)
\end{aligned}$$

$$\boxed{\frac{1}{v_p^2} \frac{\partial^2 u_y}{\partial t^2} = \left[ \frac{\partial}{\partial y} \left( \frac{1}{3}\epsilon_{xx} + \epsilon_{yy} + \frac{1}{3}\epsilon_{zz} \right) \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{yx}}{\partial x} \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{yz}}{\partial z} \right]} \quad (5)$$

$$\boxed{\frac{1}{v_p^2} \frac{\partial^2 u_z}{\partial t^2} = \left[ \frac{\partial}{\partial z} \left( \frac{1}{3}\epsilon_{xx} + \frac{1}{3}\epsilon_{yy} + \epsilon_{zz} \right) \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{zy}}{\partial y} \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{zx}}{\partial x} \right]} \quad (6)$$

In frequency domain, the elastic equation in the  $x$  - direction (equation 4) can be written as:

$$\frac{-\omega^2 u_x}{c(\omega)^2} = \left[ \frac{\partial}{\partial x} \left( \epsilon_{xx} + \frac{1}{3}\epsilon_{yy} + \frac{1}{3}\epsilon_{zz} \right) \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{xy}}{\partial y} \right] + \left[ \frac{2}{3} \frac{\partial\epsilon_{xz}}{\partial z} \right]$$

$c(\omega)$  is a frequency-dependent phase velocity that can be modeled using the Kolsky-Futterman model [1]:

$$\begin{aligned}
c(\omega) &= c(\omega_0) \left[ 1 + \frac{1}{\pi Q} \ln \frac{\omega}{\omega_0} - \frac{i \operatorname{sgn}(\omega)}{2Q} \right] \\
\frac{1}{c(\omega)^2} &= \frac{1}{c(\omega_0)^2} \left[ 1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{i \operatorname{sgn}(\omega)}{Q} \right] \quad (7)
\end{aligned}$$

This equation describes how the velocity  $c(\omega)$  depends on frequency due to attenuation,  $\left[\frac{i \operatorname{sgn}(\omega)}{Q}\right]$ , and dispersion,  $\left[-\frac{2}{\pi Q} \ln \frac{\omega}{\omega_0}\right]$ .

$$\frac{-\omega^2}{c(\omega_0)^2} \left[1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{i \operatorname{sgn}(\omega)}{Q}\right] \hat{u}_x = \left[\frac{\partial}{\partial x} \left(\epsilon_{xx} + \frac{1}{3}\epsilon_{yy} + \frac{1}{3}\epsilon_{zz}\right)\right] + \left[\frac{2}{3} \frac{\partial \epsilon_{xy}}{\partial y}\right] + \left[\frac{2}{3} \frac{\partial \epsilon_{xz}}{\partial z}\right]$$

$$\begin{aligned} \epsilon_{xx} &= \frac{1}{2} \left( \frac{\partial u_x}{\partial x} + \frac{\partial u_x}{\partial x} \right), \epsilon_{yy} = \frac{1}{2} \left( \frac{\partial u_y}{\partial y} + \frac{\partial u_y}{\partial y} \right), \epsilon_{zz} = \frac{1}{2} \left( \frac{\partial u_z}{\partial z} + \frac{\partial u_z}{\partial z} \right) \\ \epsilon_{xy} &= \frac{1}{2} \left( \frac{\partial u_x}{\partial y} + \frac{\partial u_y}{\partial x} \right), \epsilon_{xz} = \frac{1}{2} \left( \frac{\partial u_x}{\partial z} + \frac{\partial u_z}{\partial x} \right) \end{aligned}$$

$$\frac{-\omega^2}{c(\omega_0)^2} \left[1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{i \operatorname{sgn}(\omega)}{Q}\right] \hat{u}_x = \frac{1}{3} \left( 2 \frac{\partial^2 u_x}{\partial x^2} + 2 \frac{\partial^2 u_y}{\partial x \partial y} + 2 \frac{\partial^2 u_z}{\partial x \partial z} + \frac{\partial^2 u_x}{\partial x^2} + \frac{\partial^2 u_x}{\partial y^2} + \frac{\partial^2 u_x}{\partial z^2} \right)$$

$$\frac{-\omega^2}{c(\omega_0)^2} \left[1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{i \operatorname{sgn}(\omega)}{Q}\right] \hat{u}_x = \frac{1}{3} \left( 3 \frac{\partial^2 u_x}{\partial x^2} + 2 \frac{\partial^2 u_y}{\partial x \partial y} + 2 \frac{\partial^2 u_z}{\partial x \partial z} + \frac{\partial^2 u_x}{\partial y^2} + \frac{\partial^2 u_x}{\partial z^2} \right)$$

Let  $U_x$  represent the analytical function of  $u_x$  which is the complex wavefield, encapsulating both the amplitude and phase information. So, function is  $U_x = u + iv$

we can define  $u$  and  $v$  as  $-i \operatorname{sgn}(\omega)v$  and  $i \operatorname{sgn}(\omega)u$ , respectively.  $\operatorname{sgn}(\omega)$  is the signum function.

$$\frac{-\omega^2}{c(\omega_0)^2} \left[1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{i \operatorname{sgn}(\omega)}{Q}\right] \hat{u}_x = \frac{1}{3} \left( 3 \frac{\partial^2 u_x}{\partial x^2} + 2 \frac{\partial^2 u_y}{\partial x \partial y} + 2 \frac{\partial^2 u_z}{\partial x \partial z} + \frac{\partial^2 u_x}{\partial y^2} + \frac{\partial^2 u_x}{\partial z^2} \right)$$

$$\frac{-\omega^2}{c(\omega_0)^2} \left[1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{i \operatorname{sgn}(\omega)}{Q}\right] \hat{v}_x = \frac{1}{3} \left( 3 \frac{\partial^2 v_x}{\partial x^2} + 2 \frac{\partial^2 v_y}{\partial x \partial y} + 2 \frac{\partial^2 v_z}{\partial x \partial z} + \frac{\partial^2 v_x}{\partial y^2} + \frac{\partial^2 v_x}{\partial z^2} \right)$$

this becomes;

$$\frac{-\omega^2}{c(\omega_0)^2} \left[ \hat{u}_x - \frac{2\hat{u}_x}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{\hat{v}_x}{Q} \right] = \frac{1}{3} \left( 3 \frac{\partial^2 u_x}{\partial x^2} + 2 \frac{\partial^2 u_y}{\partial x \partial y} + 2 \frac{\partial^2 u_z}{\partial x \partial z} + \frac{\partial^2 u_x}{\partial y^2} + \frac{\partial^2 u_x}{\partial z^2} \right) \quad (8)$$

$$\frac{-\omega^2}{c(\omega_0)^2} \left[ \hat{v}_x - \frac{2\hat{v}_x}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{-\hat{u}_x}{Q} \right] = \frac{1}{3} \left( 3 \frac{\partial^2 v_x}{\partial x^2} + 2 \frac{\partial^2 v_y}{\partial x \partial y} + 2 \frac{\partial^2 v_z}{\partial x \partial z} + \frac{\partial^2 v_x}{\partial y^2} + \frac{\partial^2 v_x}{\partial z^2} \right) \quad (9)$$

multiply equation 9 by  $i$ , then:

$$\frac{-\omega^2}{c(\omega_0)^2} \left[ i\hat{v}_x - \frac{2i\hat{v}_x}{\pi Q} \ln \frac{\omega}{\omega_0} + \frac{-i\hat{u}_x}{Q} \right] = \frac{1}{3} \left( 3i \frac{\partial^2 v_x}{\partial x^2} + 2i \frac{\partial^2 v_y}{\partial x \partial y} + 2i \frac{\partial^2 v_z}{\partial x \partial z} + i \frac{\partial^2 v_x}{\partial y^2} + i \frac{\partial^2 v_x}{\partial z^2} \right) \quad (10)$$

Equation 8 + Equation 10;  $U = u + iv$

$$\begin{aligned}
\frac{-\omega^2}{c(\omega_0)^2} \left[ \hat{U}_x - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} \hat{U}_x + \frac{\hat{v}_x - i\hat{u}_x}{Q} \right] &= \frac{1}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_y}{\partial x \partial y} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial y^2} + \frac{\partial^2 U_x}{\partial z^2} \right) \\
\frac{-\omega^2}{c(\omega_0)^2} \left[ \hat{U}_x - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} \hat{U}_x - \frac{i\hat{U}_x}{Q} \right] &= \frac{1}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_y}{\partial x \partial y} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial y^2} + \frac{\partial^2 U_x}{\partial z^2} \right) \\
\frac{-\omega^2}{c(\omega_0)^2} \left[ 1 - \frac{2}{\pi Q} \ln \frac{\omega}{\omega_0} - \frac{i}{Q} \right] \hat{U}_x &= \frac{1}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_y}{\partial x \partial y} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial y^2} + \frac{\partial^2 U_x}{\partial z^2} \right)
\end{aligned}$$

**Use approximation:** [2]

$$\omega^2 \ln \frac{\omega}{\omega_0} \approx a\omega^2 + b\omega + d$$

where

$$a = 5.7356, \quad b = -762.1606, \quad d = 4.6054 \times 10^4$$

$$\begin{aligned}
-\frac{1}{c(\omega_0)^2} \left[ \left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} \right) \omega^2 - \frac{2b}{\pi Q} \omega - \frac{2d}{\pi Q} \right] U(\omega) &= \frac{1}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_y}{\partial x \partial y} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial y^2} + \frac{\partial^2 U_x}{\partial z^2} \right) \\
\left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} \right) \ddot{U} - \frac{2bi}{\pi Q} \dot{U} + \frac{2d}{\pi Q} U &= \frac{c(\omega_0)^2}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_y}{\partial x \partial y} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial y^2} + \frac{\partial^2 U_x}{\partial z^2} \right)
\end{aligned}$$

**Discretization:**  $U(z, x, t) \rightarrow U_{ij}^n$ , where  $z = i\Delta z = ih$ ,  $x = jh$ ,  $t = n\Delta t$

**Approximations:**

$$\dot{U} \approx \frac{U^{n+1} - U^{n-1}}{2\Delta t}, \quad \ddot{U} \approx \frac{U^{n+1} - 2U^n + U^{n-1}}{\Delta t^2}$$

let:

$$\begin{aligned}
\ddot{U}_x &= \frac{1}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_y}{\partial x \partial y} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial y^2} + \frac{\partial^2 U_x}{\partial z^2} \right) \\
\ddot{U}_x &= \frac{1}{3} \left( 3 \frac{\partial^2 U_x}{\partial x^2} + 2 \frac{\partial^2 U_z}{\partial x \partial z} + \frac{\partial^2 U_x}{\partial z^2} \right) \\
\left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} \right) \frac{U^{n+1} - 2U^n + U^{n-1}}{\Delta t^2} - \frac{2bi}{\pi Q} \frac{U^{n+1} - U^{n-1}}{2\Delta t} + \frac{2d}{\pi Q} U^n &= c_0^2 \ddot{U}_x^n, \\
\left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} - \frac{bi}{\pi Q} \Delta t \right) U_x^{n+1} - \left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} - \frac{d}{\pi Q} \Delta t^2 \right) 2U_x^n + \left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} + \frac{bi}{\pi Q} \Delta t \right) U_x^{n-1} &= c_0^2 \Delta t^2 \ddot{U}_x^n \\
\boxed{\left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} - \frac{bi}{\pi Q} \Delta t \right) U_x^{n+1} = \left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} - \frac{d}{\pi Q} \Delta t^2 \right) 2U_x^n - \left( 1 - \frac{2a}{\pi Q} - \frac{i}{Q} + \frac{bi}{\pi Q} \Delta t \right) U_x^{n-1} + c_0^2 \Delta t^2 \ddot{U}_x^n}
\end{aligned}$$

$$\begin{aligned}
\ddot{U}_x(i, k) \approx \frac{1}{3} \left( \frac{3(U_{x(i+1,k)} - 2U_{x(i,k)} + U_{x(i-1,k)}))}{h_x^2} + \frac{2(U_{z(i+1,k+1)} - U_{z(i+1,k-1)} - U_{z(i-1,k+1)} + U_{z(i-1,k-1)})}{4h_x h_z} \right. \\
\left. + \frac{(U_{x(i,k+1)} - 2U_{x(i,k)} + U_{x(i,k-1)})}{h_z^2} \right)
\end{aligned}$$

## References

- [1] Walter I Futterman. “Dispersive body waves”. In: *Journal of Geophysical research* 67.13 (1962), pp. 5279–5291.
- [2] Jidong Yang and Hejun Zhu. “A time-domain complex-valued wave equation for modelling visco-acoustic wave propagation”. In: *Geophysical journal international* 215.2 (2018), pp. 1064–1079.