

# Chapter 3.

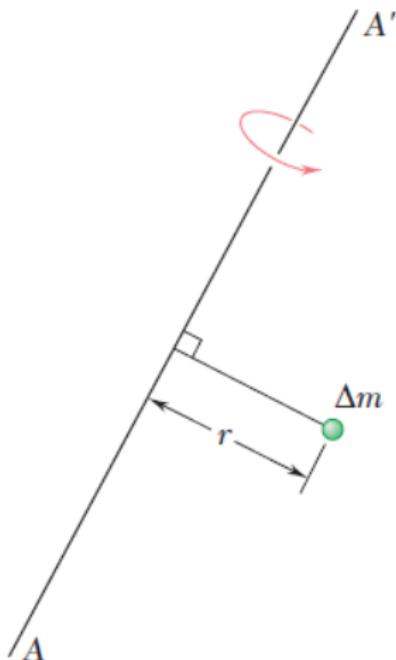
## Geometry of mass

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# Moment of Inertia of a Simple Mass

Consider a small mass  $\Delta m$  mounted on a rod of negligible mass that can rotate freely about an axis  $AA'$ . If we apply a couple to the system, the rod and mass (assumed initially at rest) rotate about  $AA'$ .

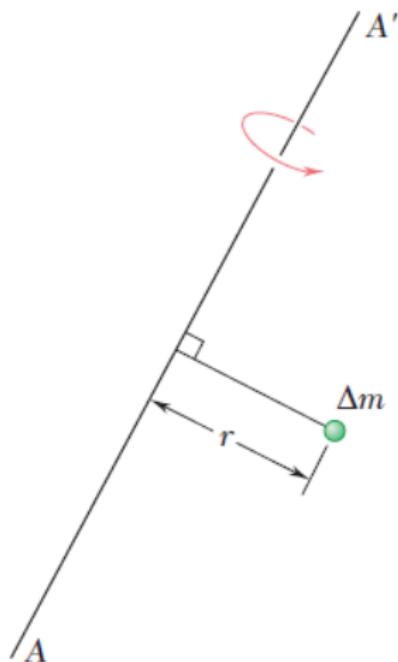


# Moment of Inertia of a Simple Mass

The time required for the system to reach a given rotation speed is proportional to the mass  $\Delta m$  and the square of the distance  $r$ .

The product  $r^2\Delta m$  is called the moment of inertia of the mass  $\Delta m$  to axis  $AA'$ .

It provides a measure of the inertia of the system (i.e. a measure of the resistance the system offers when we try to set it in motion).



# Moment of Inertia of a Simple Mass

Let's consider a body of mass  $m$  rotating about an axis  $AA'$ .

The body is composed of elements of masses  $\Delta m_i$ .

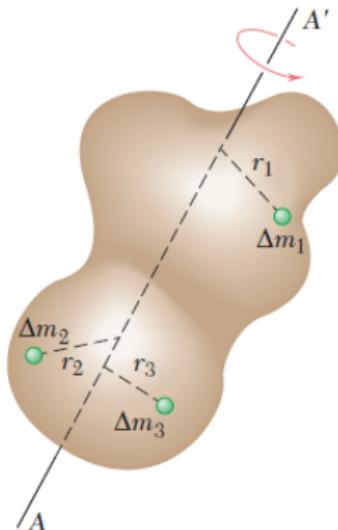
The body's resistance to being

rotated is measured by the sum (the moment of inertia of the body to axis  $AA'$ ):

$$\sum r_i^2 \cdot \Delta m_i$$

Increasing the number of elements, we find that the moment of inertia is equal, in the limit, to the integral:

$$I = \int r^2 \cdot dm$$

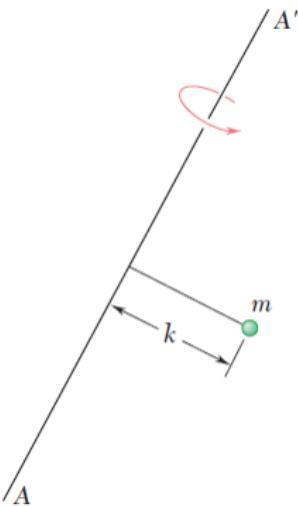


# Radius of gyration of a mass

We define the radius of gyration  $k$  of the body with respect to axis  $AA'$  by the relation:

$$I = k^2 \cdot m \text{ or } k = \sqrt{\frac{I}{m}}$$

The radius of gyration  $k$  represents the distance at which the entire mass of the body should be concentrated if its moment of inertia to  $AA'$  is to remain unchanged.



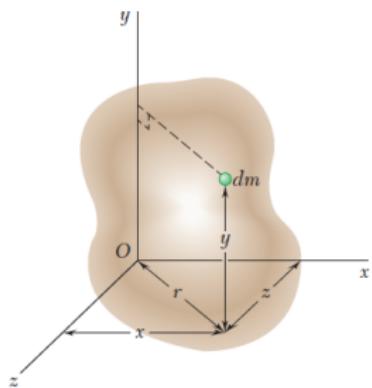
In SI units: the radius of gyration  $k$  is expressed in meters and the mass  $m$  in kilograms, so the unit for the moment of inertia of a mass is  $\text{kg} \cdot \text{m}^2$ .

# Moments of inertia to coordinate axes

We can express the moment of inertia of a body to a coordinate axis in terms of the coordinates  $x$ ,  $y$ , and  $z$  of the element of mass  $dm$ .

For example, the moment of inertia of the body with respect to the  $y$ -axis is:

$$I_y = \int r^2 dm = \int (z^2 + x^2) dm$$



We obtain similar expressions for the moments of inertia to the  $x$  and  $z$  axes.

$$I_z = \int (y^2 + x^2) dm \text{ and } I_x = \int (y^2 + z^2) dm$$

# Mass Products of Inertia

The moment of inertia  $I_{OL}$  of the body to  $OL$  is equal to  $\int p^2 dm$ , where  $p$  denotes the perpendicular distance from the element of mass  $dm$  to the axis  $OL$ .

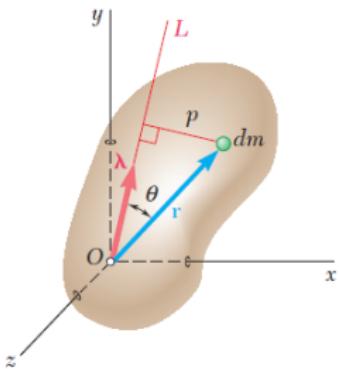
Let's consider:

$\vec{\lambda}$  is the unit vector along  $OL$ .

$\vec{r}$  is the position vector of the element  $dm$

Then, the perpendicular distance  $p$  is equal to:

$$p = r \cdot \sin(\theta) = \|\vec{\lambda} \wedge \vec{r}\|$$



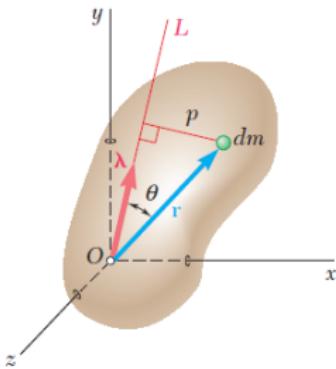
# Mass Products of Inertia

We therefore have:  $I_{OL} = \int p^2 dm = \int \|\vec{\lambda} \wedge \vec{r}\|^2 dm$

If we take:

$$\vec{\lambda}(\lambda_x, \lambda_y, \lambda_z)$$

$$\vec{r}(x, y, z)$$



we find:

$$I_{OL} = \lambda_x^2 \int (y^2 + z^2) dm + \lambda_y^2 \int (z^2 + x^2) dm + \lambda_z^2 \int (x^2 + y^2) dm$$

$$-2\lambda_x\lambda_y \int xy dm - 2\lambda_y\lambda_z \int yz dm - 2\lambda_z\lambda_x \int zx dm$$

# Mass Products of Inertia

The last three integrals which involve products of coordinates, are called the products of inertia of the body:

$$I_{xy} = \int xydm ; I_{yz} = \int yzdm ; I_{zx} = \int zx dm$$

Then:

$$I_{OL} = \lambda_x^2 I_x + \lambda_y^2 I_y + \lambda_z^2 I_z - 2\lambda_x\lambda_y I_{xy} - 2\lambda_y\lambda_z I_{yz} - 2\lambda_z\lambda_x I_{zx}$$

Or in matrix form:

$$I_{OL} = \begin{Bmatrix} \lambda_x & \lambda_y & \lambda_z \end{Bmatrix} \times \begin{bmatrix} I_x & -I_{xy} & -I_{xz} \\ -I_{yx} & I_y & -I_{yz} \\ -I_{zx} & -I_{zy} & I_z \end{bmatrix} \times \begin{Bmatrix} \lambda_x \\ \lambda_y \\ \lambda_z \end{Bmatrix}$$

$$I_{OL} = \vec{\lambda}^T \times I_{Oxyz} \times \vec{\lambda}$$

# Product of inertia with respect to two perpendicular lines

Let  $\Delta$  and  $\Delta'$  be two straight lines with directions  $\vec{n}$  and  $\vec{t}$  (unit vectors).

$$I_{nt} = \int x_n x_t dm$$

$x_n$  and  $x_t$  are the coordinates of  $P$  in the reference frame  $(O, \vec{n}, \vec{t})$ .

$$I_{nt} = -\vec{n}^T \cdot I_{(Oxyz)} \cdot \vec{t}$$

# Change of reference frame in rotation

$I_{(Oxyz)}$  is the inertia tensor with respect to a reference frame  $R(Oxyz)$ .

We can calculate the tensor in another frame  $R'(Ox'y'z')$  defined by the rotation of  $R$ , using the following formulas:

$$I_{(Ox'y'z')} = P_{(R' \rightarrow R)} I_{(Oxyz)} P_{(R \rightarrow R')}$$

# Parallel-Axis Theorem for Mass Moments of Inertia

Consider a body of mass  $m$  and let  $(Oxyz)$  be a system of rectangular coordinates whose origin is at the arbitrary point  $O$ . Let  $(Gx'y'z')$  be a system of parallel axes (the origin is at the centre of gravity  $G$ ).

We denote the coordinates :

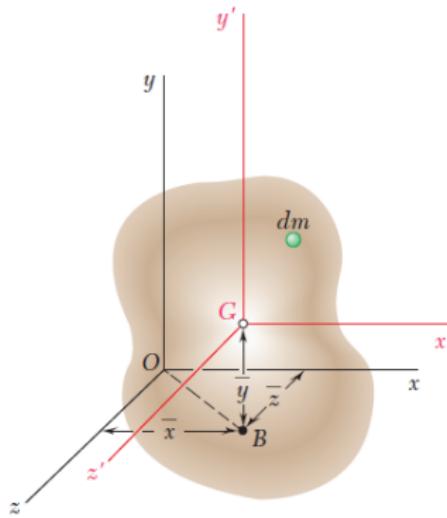
$G(\bar{x}, \bar{y}, \bar{z})$  in  $(Oxyz)$

$dm(x, y, z)$  in  $(Oxyz)$

$dm(x', y', z')$  in  $(Gx'y'z')$

Then :

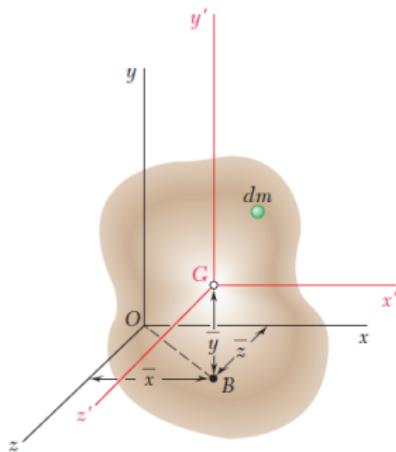
$$x = x' + \bar{x}; y = y' + \bar{y}; z = z' + \bar{z}$$



# Parallel-Axis Theorem for Mass Moments of Inertia

We can express the moment of inertia of the body to the x-axis  $I_x = \int (y^2 + z^2) dm$  as:

$$I_x = \int [(y' + \bar{y})^2 + (z' + \bar{z})^2] dm$$



$$I_x = \int (y'^2 + z'^2) dm + 2\bar{y} \int y' dm + 2\bar{z} \int z' dm + (\bar{y}^2 + \bar{z}^2) \int dm$$

$$I_x = \bar{I}_{x'} + m(\bar{y}^2 + \bar{z}^2)$$

# Parallel-Axis Theorem for Mass Moments of Inertia

Similarly,

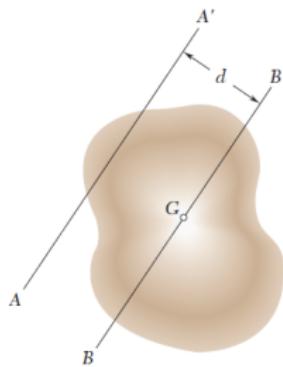
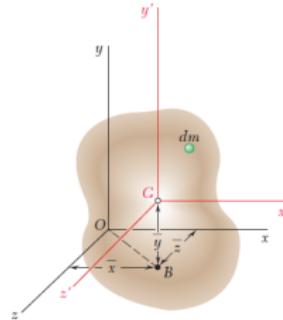
$$I_x = \bar{I}_{x'} + m(\bar{y}^2 + \bar{z}^2)$$

$$I_y = \bar{I}_{y'} + m(\bar{z}^2 + \bar{x}^2)$$

$$I_z = \bar{I}_{z'} + m(\bar{x}^2 + \bar{y}^2)$$

Parallel-axis theorem for mass moments of inertia

$$I = \bar{I} + md^2$$



# Particular cases

## Case 1. Plane solids:

In the case of plane solids, one of the element's coordinates is zero. If the solid is in the plane ( $Oxy$ ) then  $z = 0$ .

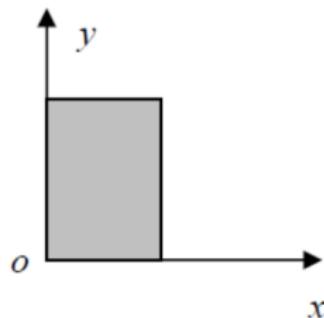
We immediately deduce that:

$$I_x = \int y^2 dm; I_y = \int x^2 dm$$

Hence:

$$I_z = \int (x^2 + y^2) dm = I_x + I_y$$

and  $I_{xz} = I_{yz} = 0$ ;  $I_{xy} = \int xy dm$



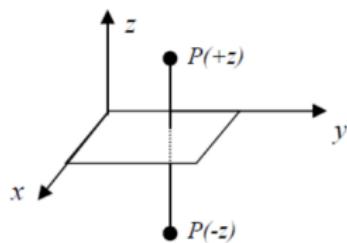
# Particular cases

## Case 2: plane of symmetry

If  $oxy$  is a plane of symmetry, then  $oz$  is a principal axis of inertia

$$\int xz dm = \int yz dm = 0 \rightarrow I_{xz} = I_{yz} = 0$$

$$I_{Oxyz}(S) = \begin{bmatrix} I_x & -I_{xy} & 0 \\ -I_{yx} & I_y & 0 \\ 0 & 0 & I_z \end{bmatrix}$$



If  $oyz$  is a plane of symmetry then  $I_{xy} = I_{xz} = 0$  ( $ox$  is a principal axis of inertia)

If  $oxz$  is a plane of symmetry then  $I_{yz} = I_{xy} = 0$  ( $oy$  is a principal axis of inertia)

## Particular cases

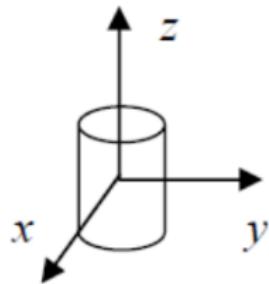
### Case 3: Axis of symmetry (revolution around an axis)

If the system is a body of revolution around  $oz$ , then any plane passing through  $oz$  is a plane of symmetry. Therefore:

$$I_{xy} = I_{xz} = I_{yz} = 0$$

and

$$I_x = I_y = \frac{I_z}{2} + \int z^2 dm$$



# Particular cases

## Case 4: Point of symmetry (spherical revolution)

If point O is a point of revolution then:

$$I_{xy} = I_{xz} = I_{yz} = 0$$

and

$$I_x = I_y = I_z = \frac{2}{3} \int (x^2 + y^2 + z^2) dm$$

