

# QUANTUM BUILDING BLOCKS

## Recap

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# Single-Qubit Quantum Systems

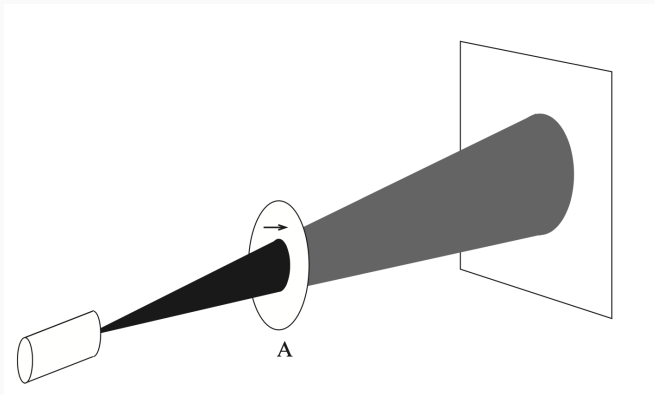
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# The Quantum Mechanics of Photon Polarization

There is a simple experiment with light polarization illustrates some of the **nonintuitive** behavior of quantum systems, behavior that is **exploited** to good effect in quantum algorithms and protocols.

# The Quantum Mechanics of Photon Polarization

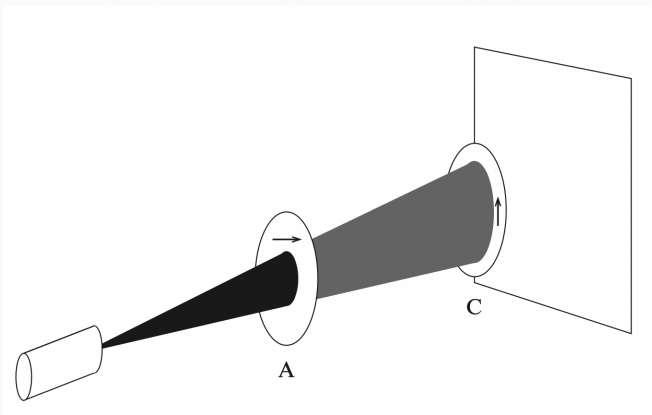
Shine a beam of light on a projection screen. When polaroid **A** is placed between the light source and the screen, the intensity of the light reaching the screen is **reduced**.



**Figure 1:** Single polaroid weaken unpolarized light by 50 percent.

# The Quantum Mechanics of Photon Polarization

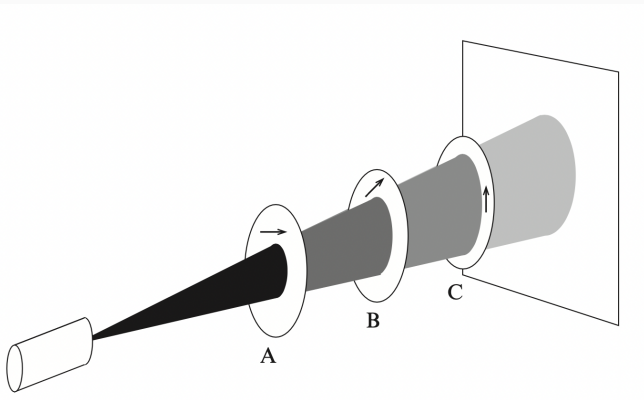
Next, if we place polaroid **C** and if polaroid C is rotated so that its polarization is orthogonal (vertical) to the polarization of **A**, no light reaches the screen.



**Figure 2:** Two orthogonal polaroids block all photons.

# The Quantum Mechanics of Photon Polarization

Finally, place polaroid **B** between polaroids **A** and **C**. Surprisingly, at most polarization angles of **B**, light shines on the screen. The intensity of this light will be maximal if the polarization of **B** is at  $45^\circ$  to both **A** and **C**.



**Figure 3:** Inserting a third polaroid allows photons to pass.

# The Quantum Mechanics of Photon Polarization

Clearly the polaroids cannot be acting as **simple sieves**; otherwise, inserting polaroid B could not increase the number of photons reaching the screen.

## Some notation before an explanation

Quantum mechanics models a photon's polarization state by a unit vector, a vector of length 1, pointing in the appropriate direction.

We write  $|\uparrow\rangle$  and  $|\rightarrow\rangle$  for the unit vectors that represent vertical and horizontal polarization respectively.

In quantum mechanics, the standard notation for a vector representing a quantum state is  $|\nu\rangle$ , just as  $\vec{v}$  or  $\mathbf{v}$  are notations used for vectors in other settings.

An arbitrary polarization can be expressed as a linear combination

$$|\nu\rangle = a|\uparrow\rangle + b|\rightarrow\rangle$$

of the two basis vectors  $|\uparrow\rangle$  and  $|\rightarrow\rangle$ .



## Polaroid interacts with a photon

When a photon with polarization  $|\nu\rangle = a|\uparrow\rangle + b|\rightarrow\rangle$  meets a polaroid with preferred axis  $|\uparrow\rangle$ , the photon will get through with probability  $|a|^2$  and will be absorbed with probability  $|b|^2$ .

Any photon that passes through the polaroid will now be polarized in the direction of the polaroid's preferred axis.

## A Quantum Explanation

In the experiment, any photons that pass through polaroid **A**, will leave as a  $|\rightarrow\rangle$ .

So it has no chance of passing through polaroid **C**, which was vertical.

To understand what will happen, when we insert **B**, let's rewrite

$$|\rightarrow\rangle = \frac{1}{\sqrt{2}} |\nearrow\rangle - \frac{1}{\sqrt{2}} |\nwarrow\rangle.$$

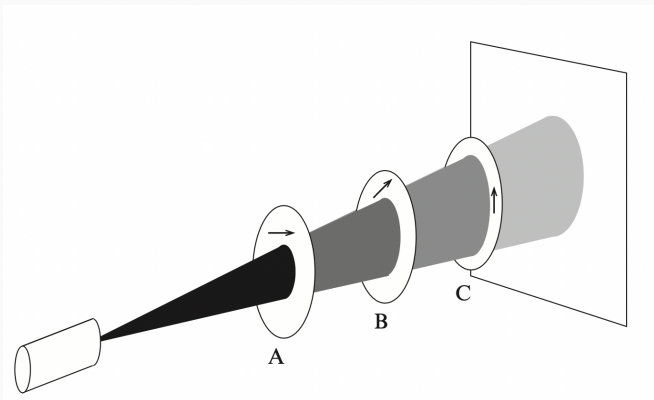
So, any photon in  $|\rightarrow\rangle$  state will have a  $\frac{1}{\sqrt{2}}$  amplitude in  $|\nearrow\rangle$  direction. Thus, it has a probability of  $\frac{1}{2}$  passing through **B**.

The same happens with **C** as

$$|\nearrow\rangle = \frac{1}{\sqrt{2}} |\uparrow\rangle + \frac{1}{\sqrt{2}} |\rightarrow\rangle.$$

# The Quantum Mechanics of Photon Polarization

We see how the model explains that  $\frac{1}{8}$  of light pass through **A**, **B**, **C**.



# Single Quantum Bits

The space of possible polarization states of a photon is an example of a quantum bit, or **qubit**. A qubit has a continuum of possible values: any state represented by a unit vector  $a|\uparrow\rangle + b|\rightarrow\rangle$  is a legitimate qubit value.

Any quantum mechanical system that can be modeled by a two-dimensional complex vector space can be viewed as a qubit.

It includes photon polarization, electron spin, and the ground state together with an excited state of an atom.

It is as yet unclear which two-state systems will be most suitable for physical realizations of quantum computers; it is likely that a **variety of physical representation** of qubits will be used.

## $|0\rangle$ and $|1\rangle$

We just need to specify any two orthonormal (perpendicular) states and call them  $|0\rangle$  and  $|1\rangle$ . It can be any states. For example in the photon polarization example we could equally set  $\{|0\rangle, |1\rangle\} = \{|\rightarrow\rangle, |\uparrow\rangle\}$  or  $\{|0\rangle, |1\rangle\} = \{|\nearrow\rangle, |\nwarrow\rangle\}$ .

We will use the convention that  $|0\rangle = |\uparrow\rangle$ ,  $|1\rangle = |\rightarrow\rangle$ , which means  $|\nearrow\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$  and  $|\nwarrow\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$ .

We will call  $\{|0\rangle, |1\rangle\}$  as standard basis.

And write any qubit in the state  $a|0\rangle + b|1\rangle$  as  $\begin{pmatrix} a \\ b \end{pmatrix}$ .

# Single-Qubit Measurement

Quantum theory postulates that any device that measures a two-state quantum system must have two preferred states whose representative vectors,  $\{|u\rangle, |u^\perp\rangle\}$ , form an **orthonormal** basis for the associated vector space.

The probability of measuring  $|u\rangle$  or  $|u^\perp\rangle$  is equal to the square of the magnitude of  $|u\rangle$  and  $|u^\perp\rangle$ , respectively.

And measured qubit has to gain amplitude of 1 in the measured state, so the measurement outcome is always one of the two basis vectors.

This behavior of measurement is an **axiom** of quantum mechanics.

For this reason, whenever we say “*measure a qubit*,” we must specify with respect to which basis the measurement takes place.

## Pop-science misconception: $|\nu\rangle$ is not prob. mix of $|0\rangle, |1\rangle$

Rather,  $|\nu\rangle$  is a definite state, which, when measured in certain bases, gives deterministic results, but for some other bases gives random results.

For example a state

$$|\nearrow\rangle = \frac{1}{\sqrt{2}} (|\rightarrow\rangle + |\uparrow\rangle)$$

is deterministic in Hadamard basis  $\{|\nearrow\rangle, |\nwarrow\rangle\}$ , however random in  $\{|\rightarrow\rangle, |\uparrow\rangle\}$ .

It is OK to think that  $a|0\rangle + b|1\rangle$  is at the same time  $|0\rangle, |1\rangle$ , however we have to be careful to distinguish between

$$|+\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle), |-\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \text{ and } |i\rangle = \frac{1}{\sqrt{2}} (|0\rangle + i|1\rangle)$$

that behave differently in many cases, yet have the same proportion of  $|0\rangle$  and  $|1\rangle$ .

## A global phase vs relative phase

Let's consider a qubit state  $|1\rangle$ ,  $-|1\rangle$ , and  $i|1\rangle$ . Can we differentiate them?

Our measuring devices can not. This is what we call a global phase. We can not tell a difference between states  $|\phi\rangle$  and  $c \cdot |\phi\rangle$ , when  $|c| = 1$  in other symbols, ( $c = e^{i\theta}$  for some  $\theta$ ).

We only care the difference between the given states, their relative phases.

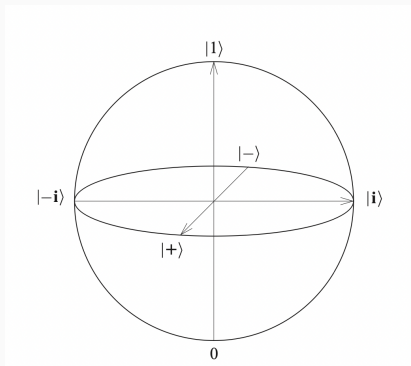
While multiplication with a unit constant does not change a quantum state vector, relative phases in a superposition do represent distinct quantum states: eventhough  $|v_1\rangle \sim e^{i\theta} |v_1\rangle$ ,

$$\frac{1}{\sqrt{2}} (|v_1\rangle + |v_2\rangle) \not\sim \frac{1}{\sqrt{2}} (e^{i\theta} |v_1\rangle + |v_2\rangle).$$



# Bloch Sphere

It turns out it is also convenient to represent a qubit as a point in a 3 dimensional space, where each diameter is a orthonormal basis of a qubit. Note that the angles are doubled in this representation.



**Figure 4:** Location of certain single-qubit states on the surface of the Bloch sphere.