

# Holographic Complexity in Non-Commutative Gauge Theory

Work with Stefan Eccles, Willy Fischler, and Ming-Lei Xiao  
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# Introduction

Given recent interest in the holographic complexity conjectures, it is interesting to test these conjectures in new contexts, to see if they hold up under closer scrutiny. With this in mind, we will examine holographic complexity in a geometry dual to non-commutative super Yang-Mills at finite temperature. In particular, we will be interested in the behaviour of the complexity as we vary the Moyal scale of the boundary theory. In the limit where the Moyal scale is sent to zero, we will recover the usual result for a planar black hole in  $\text{AdS}_5$ . On the other hand, as we make the Moyal scale large, we find that we get an enhancement of the late time complexification rate, which asymptotically approaches an enhancement.

# Holographic Complexity Recap

Recap of a few points from the last talk:

- Holographic complexity is motivated by the growth of the behind the horizon geometry
- Originally it was conjectured that the volume of a maximal spatial slice is dual to the circuit complexity of the boundary state (relative to some reference state)
- Relative circuit complexity of a state  $|\psi\rangle$  is defined as the minimum number of unitary gates  $g_i$  from some gate set  $G$  needed to build a quantum circuit which when applied to a reference state  $|\psi_0\rangle$  will produce  $\psi$  to within a tolerance  $\epsilon$
- The 'complexity = volume' conjecture is supported by the switchback effect, where we reproduce the expected behavior holographically

# Complexity = Action

Complexity = volume has a few unpleasant features

- For example, in order to reproduce the correct boundary behaviour, the volume must be multiplied by a non-universal length scale

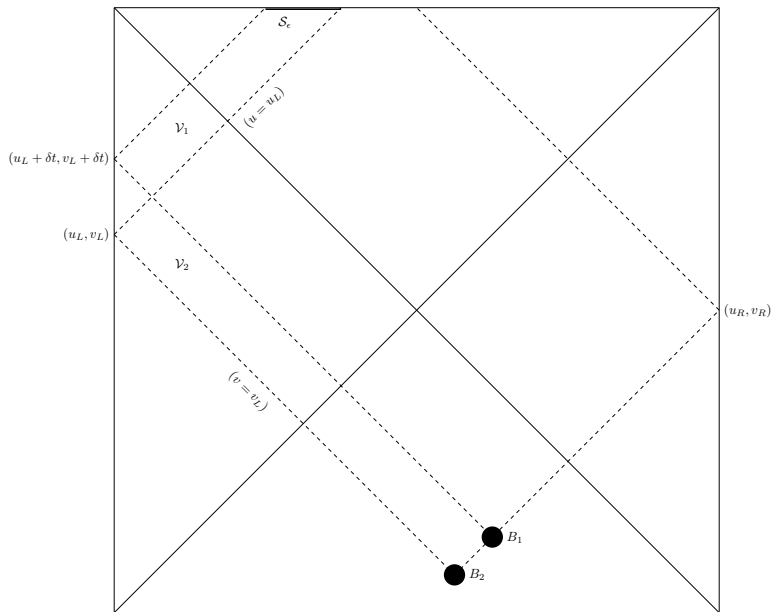
One might seek an alternative proposal, which still captures something about the behind the horizon geometry, but which does not have these features. In fact, such an alternative has been proposed by Susskind et al., and it goes by 'complexity = action.'

- According to complexity = action, the complexity is dual to the action of a so called 'Wheeler-DeWitt' (WDW) patch.
- Because this is an action, it can be undimensionalized with some multiple of  $\hbar$ .
- A universal choice for this coefficient is consistent with the expected large temperature behaviour.
- The action of a WDW patch also behaves in the appropriate way in the presence of shockwaves, so it still reproduces the switchback effect.

# The WDW patch

(insert figure of WDW patch here)

# Time derivative of WDW patch action



# Non-Commutative SYM

- We would like to test complexity = action in a new context. However, computing the complexity of a state in a strongly coupled field theory is not something we know how to do. Instead, we will rely on our intuition about the qualitative behavior of complexity in some novel situation.
- A good candidate is SYM living on a non-commutative geometry, i.e., a geometry where two of the spatial directions don't commute.
- A gravity dual to such a theory was derived in the late 90's by a number of authors [1, 2].
- Generalizations of this system to other numbers of dimensions have also been considered in, e.g., [3, 4].
- These solutions are obtained by considering a stack of Dp-branes in type II string theory and exciting a component of the NS-NS 2-form B-field along two of the worldsheet directions.
- Additional non-commutativities can be introduced by turning on additional components of the B-field.

# The gravity dual to NCSYM

The gravity dual to NCSYM is obtained in a manner similar to the standard commutative story.

- Consider a stack of D3-branes
- we turn on a 2-form  $B$  parallel to the two spatial dimensions of the branes
- This can be achieved by applying a combination of T-dualities and Gauge transformations
- The  $B$ -field has the effect of introducing non-commutativity in the worldbrane theory
- The near horizon geometry is a IIB SUGRA solution.
- This procedure can be generalized to  $D_p$ -branes, with  $p = 2, 4$ , or  $5$ .



## D3-brane results: Finite time behavior

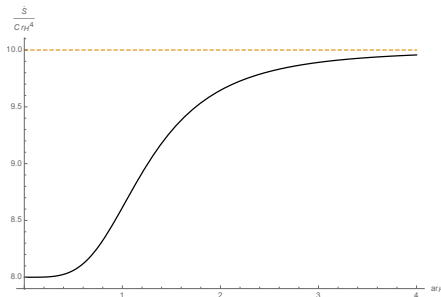
Similar behaviour to (cite paper by Rob, Shira, etc.), (of course as  $a \rightarrow 0$  must reproduce planar black hole)

- Logarithmic divergence at critical time (smooth out with thermal average?)
- Reaches true global maximum in order one time (in thermal units)
- approaches asymptotic value from above
- This behavior persists for all values of the Moyal scale ...
- ... but as the Moyal scale becomes very large, the true maximum shrinks, and the asymptotic value increases, so that the gap between them decreases

This behaviour calls into question the normalization to complexity = action as set by (cite Brown et al.), though the logic that to that normalization would seem already to be contradicted by (cite Cotrell 'complexity is simple')

## D3-brane results: Late time limit

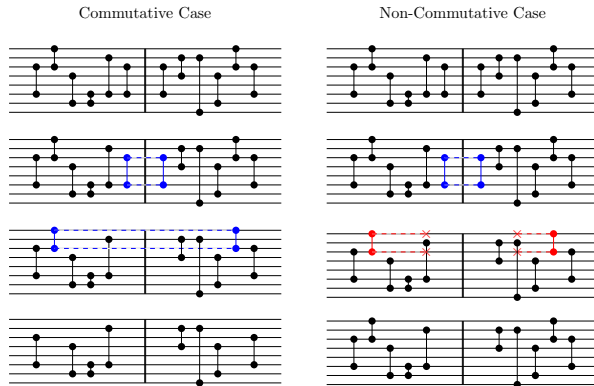
We see a 25% enhancement for large Moyal scale!



**Figure :** Late time action growth rate normalized by  $C = \frac{\alpha^4 \Omega_5 V_3}{g_s^2}$  and extra  $r_H$  dependence, versus  $a r_H$ , which is the Moyal scale measured in units of thermal length.

At first this seems a bit mysterious, but perhaps we should have expected it.

# Non-Commutativity and Complexity: A heuristic argument



**Figure :** Consider an optimal circuit that evolves our state by a small time  $\delta t$ . As we compose this circuit many times, we expect some cancellations between gates in adjacent copies of the circuit. Non-commutativity acts as an impediment to such cancellation by making fewer gates commute (due to non-locality), and so the final circuit after cancellation is more complex.

## Results for other values of $p$

$p$	$m = 0$	$m = 1$	$m = 2$
2	12	12	-
3	8	10	-
4	5	5	8
5	4	5	6

# Conclusions

- For  $p = 3, 5$  we do see an increase, as expected.
- Though we did not see an increase for  $p = 2$  or for  $p = 4$  with a single non-trivial commutator, we did not see a decrease either.
- Overall, the results are consistent with the heuristic argument above.
- This result is in tension with the idea that commutative black holes are the fastest possible computers
- In future work we plan to repeat our calculations for complexity = volume.

# References



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## Backup Slides

# Holographic Complexity

- Consider a two sided black hole in asymptotically AdS space.
- A holographic puzzle: what could be dual to the geometry behind the horizon?
  - ▶ AdS/CFT  $\Rightarrow$  a two sided black hole is dual to the thermofield double (TFD) state.
  - ▶ The volume of the wormhole keeps growing long after the thermalization time .
  - ▶ This would seem to indicate that most observables we consider for a thermal state will fail.
  - ▶ By contrast, the circuit complexity of a TFD state continues to grow long past the thermalization time
- (Possible) resolution: the behind the horizon geometry encodes the circuit complexity of the quantum state.



# Quantum Circuit Complexity

What is 'circuit complexity'?

- Consider a Hilbert space  $\mathcal{H}$ , e.g., the Hilbert space for  $N$  quantum bits.
- A universal gate set  $\{g_i\}$  for  $\mathcal{H}$  is a set of unitary operators on the Hilbert space such that any unitary  $U$  acting on  $\mathcal{H}$  can be approximated by some product  $\prod_i g_{\alpha_i}$  to within a small tolerance  $\epsilon$ .
- Such a product of gates is referred to as a quantum circuit.
- The quantum circuit complexity of a unitary  $U$  is then the minimum number of gates needed to approximate  $U$  to within the tolerance.
- In the example of qubits, one typically considers gates that act on a single qubit or pairs of qubits at a time.
- Given some reference state  $|\psi_R\rangle$ , one may define the complexity of a state  $|\psi\rangle$  as the minimum of complexity  $C(U)$  over all unitaries  $U$  such that  $|\psi\rangle = U |\psi_R\rangle$ .