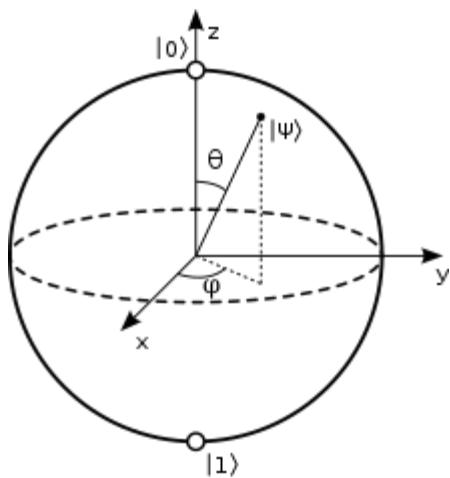


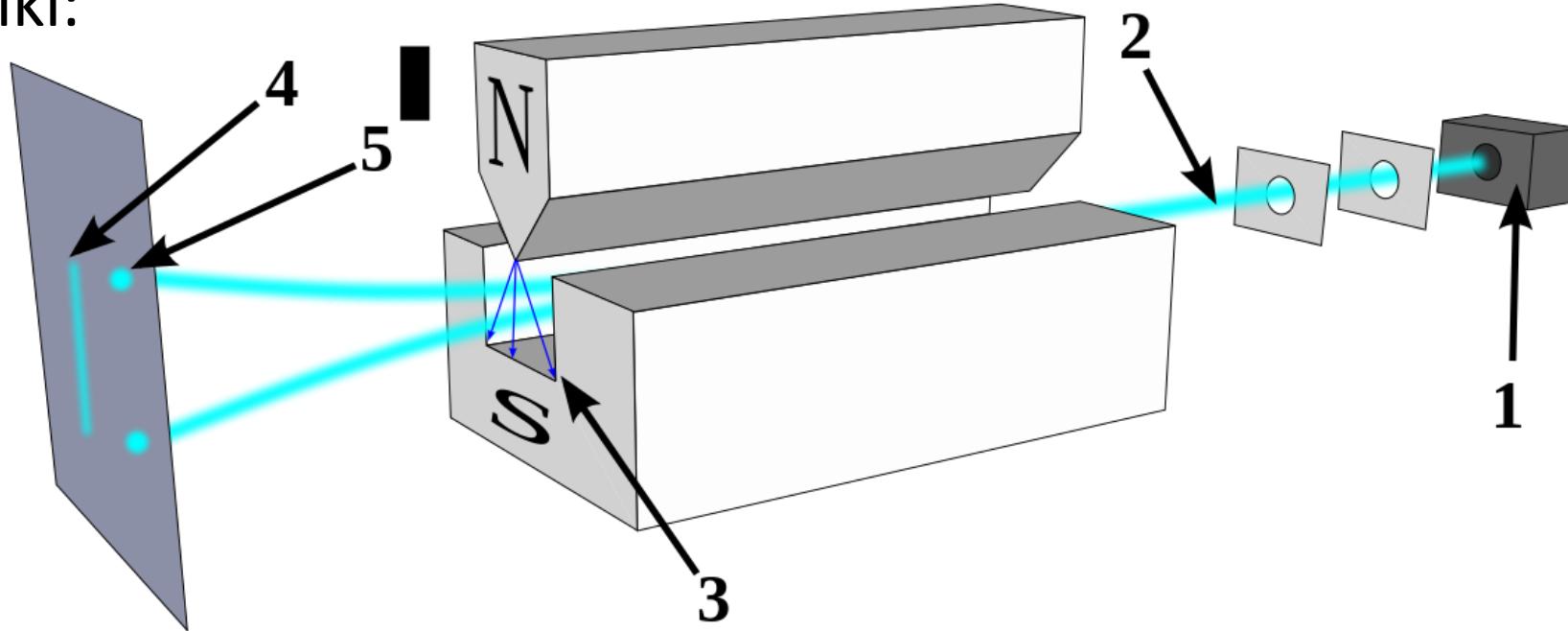
Intro to Quantum Computing

Maths of Qubits Part 1



Recall: The Stern-Gerlach Experiment

- From wiki:



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- Since we have rotated the magnet by 90° , the two new orientations of N-S pole are \rightarrow and \leftarrow .
- By rotating the apparatus we ***induce a new base of measurement*** and the spin of $|\psi\rangle$ (before measurement) is in superposition $c|\rightarrow\rangle + d|\leftarrow\rangle$, with $c^2 + d^2 = 1$.

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- We can also measure on any other direction we like!

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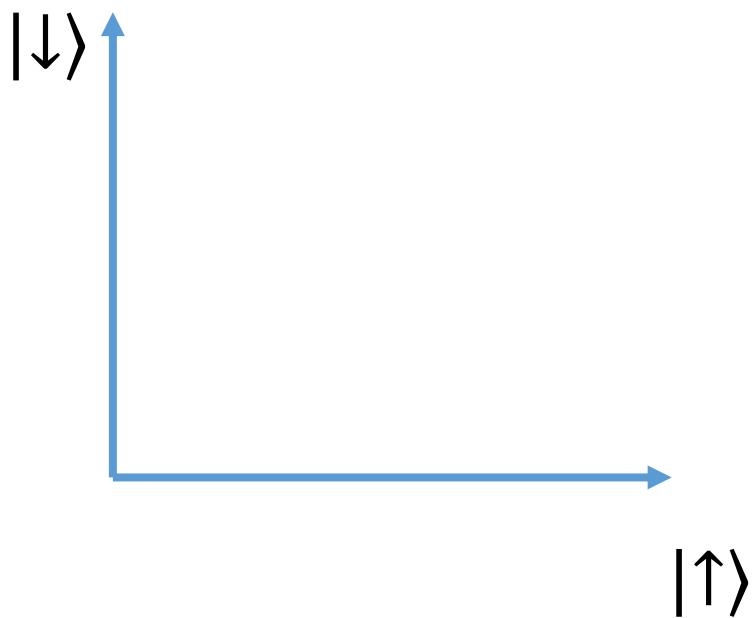
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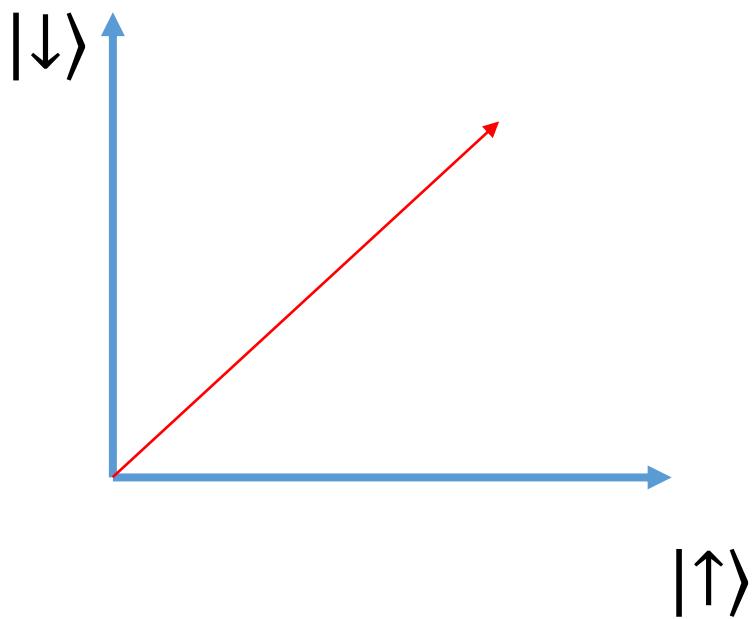
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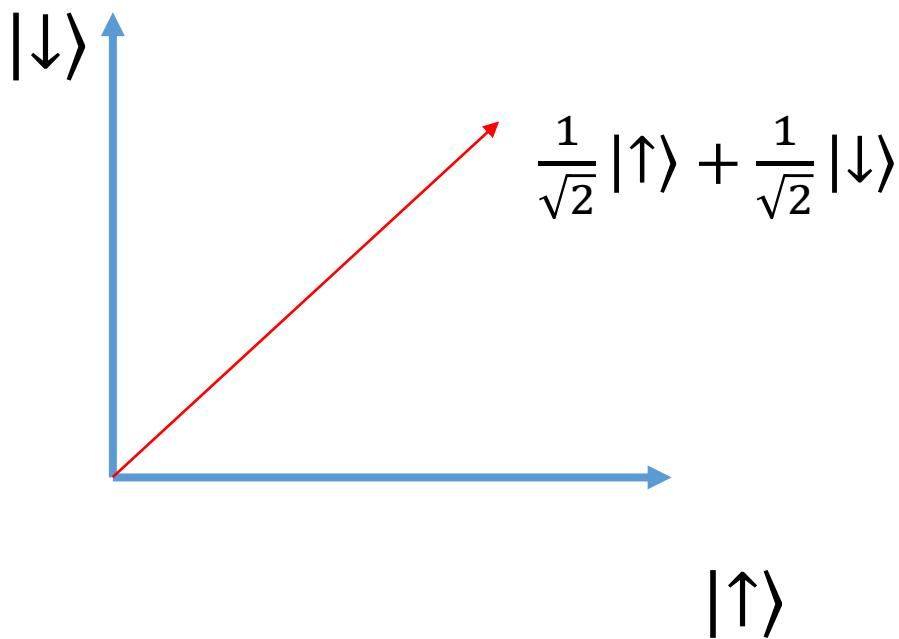
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Braket (Dirac's) notation

- Any quantum quantity or system ψ is described as (or by) $|\psi\rangle$.
- It is an element of a ***complex Hilbert Space***.
- We will not deal (as discussed) with complex number but for now we will present the ideas using complex numbers for full generality.

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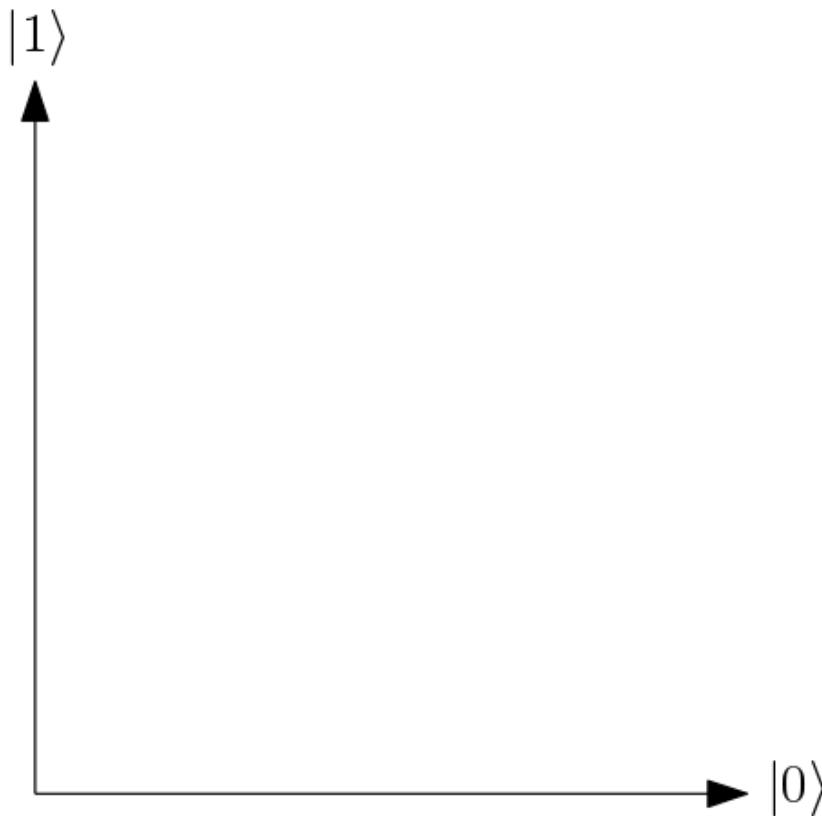
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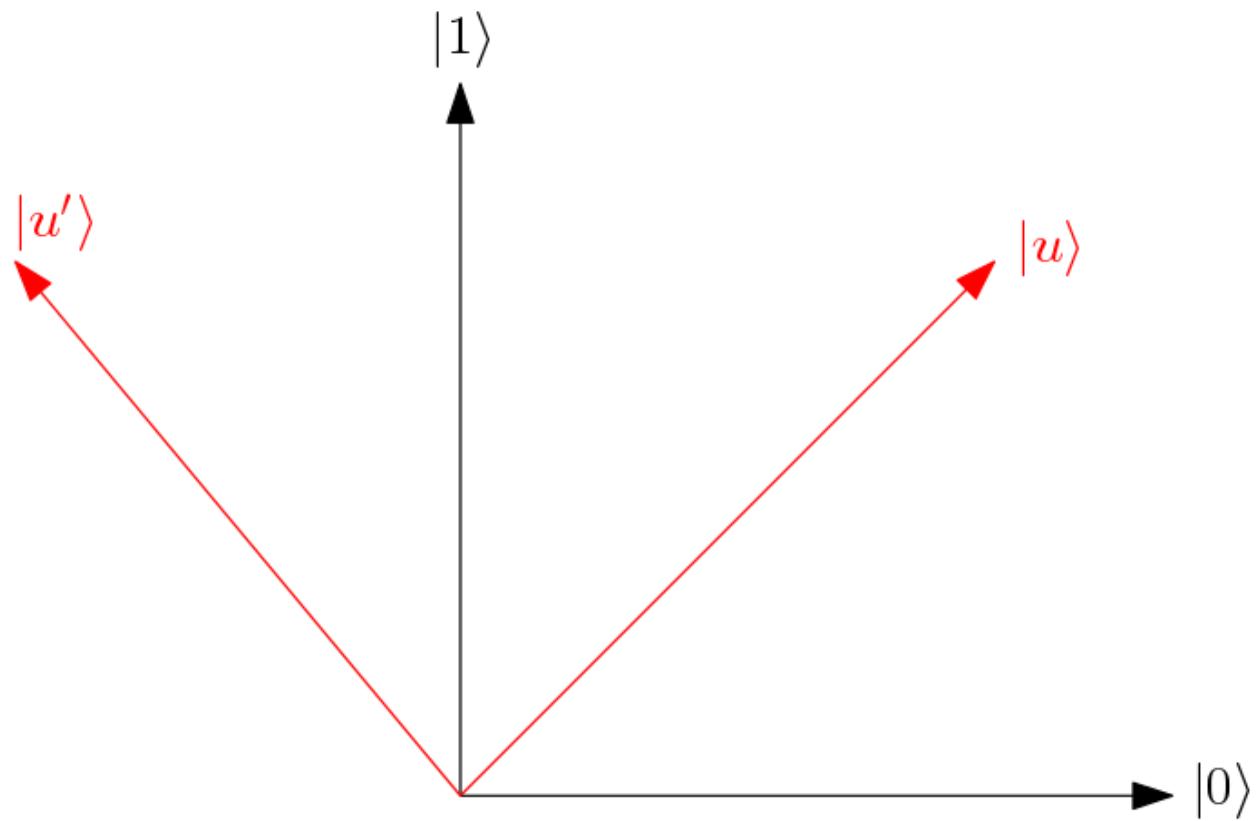
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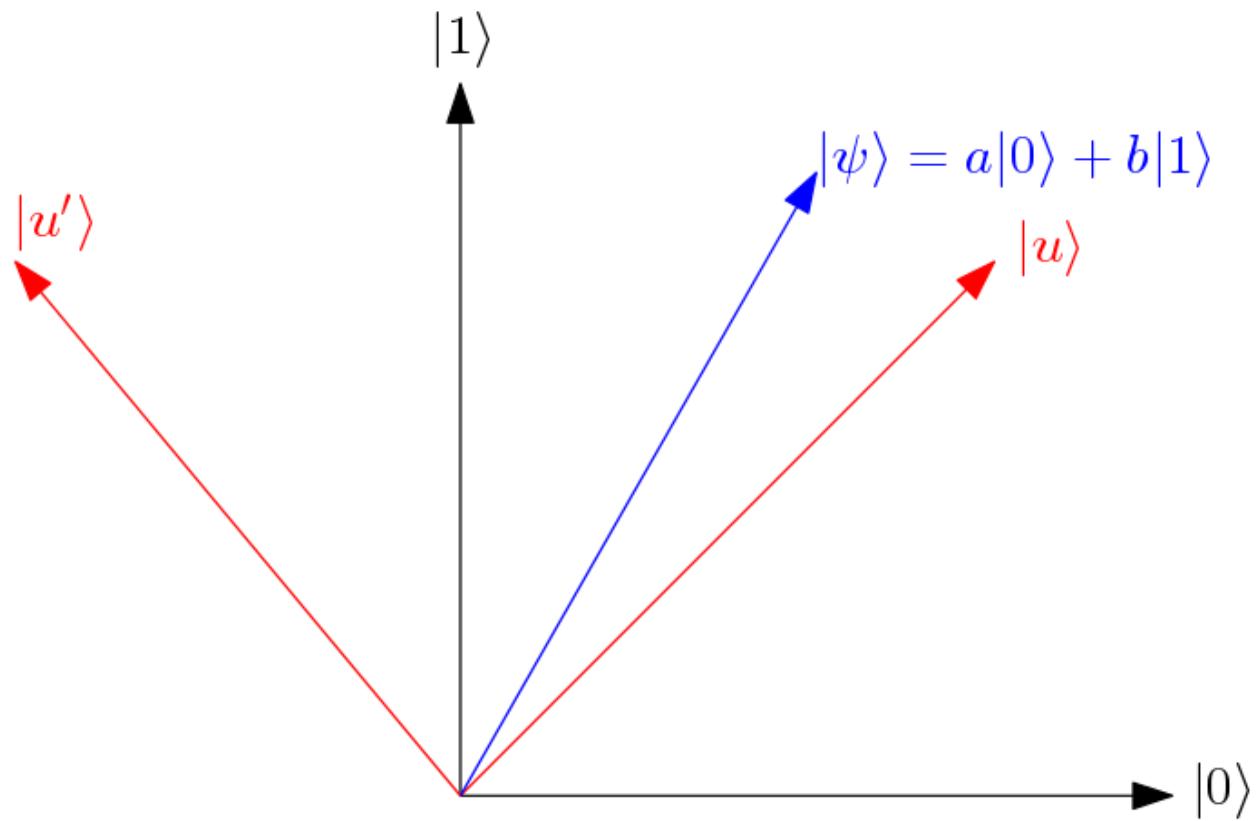
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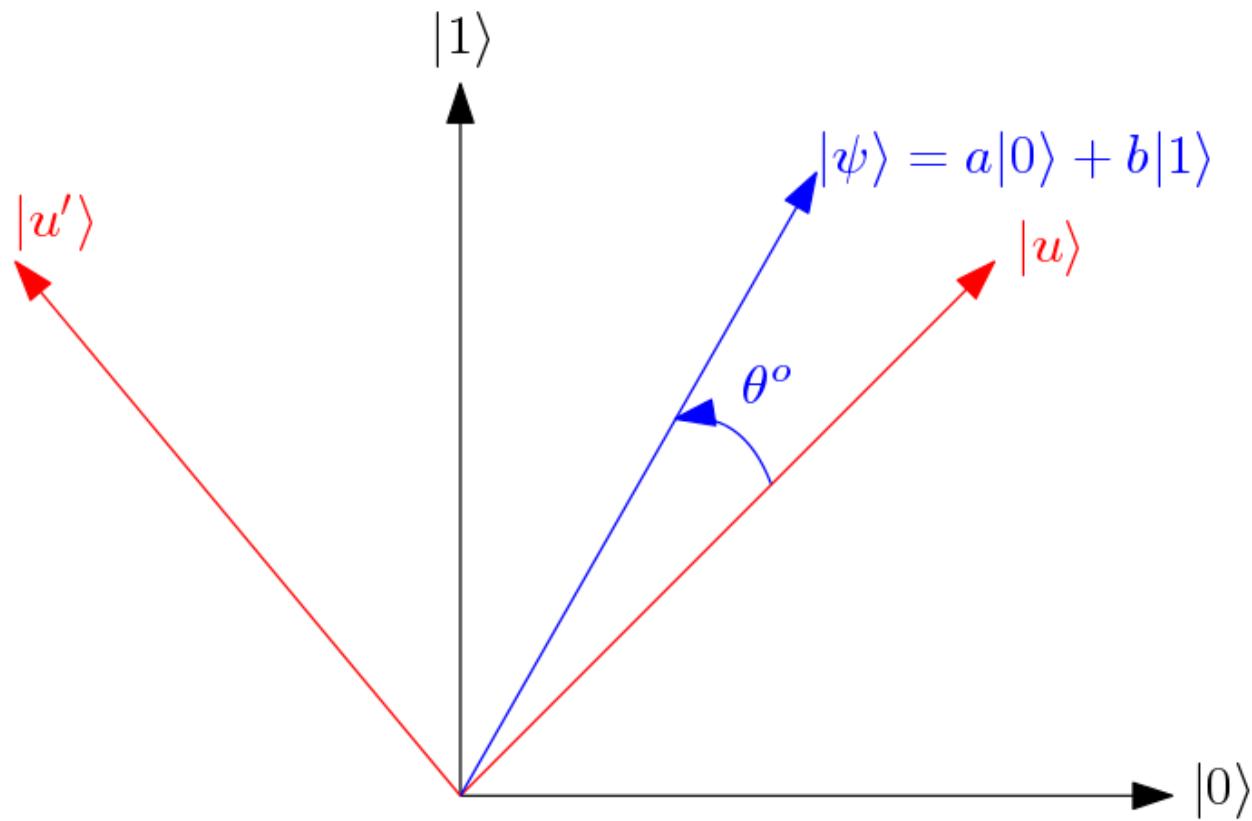
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- So we do not measure anymore if the qbit has spin **up** or **down** but if it is in $|u\rangle$ state or $|u'\rangle$ state.

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- We will give a short recap of the notions we will be using in the bra-ket vocabulary.

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- More compactly we denoted it as $\langle a | b \rangle$ (i.e., $a^T b$).

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Geometrically: $||\mathbf{a}\rangle + |\mathbf{b}\rangle|^2 = ||\mathbf{a}\rangle|^2 + ||\mathbf{b}\rangle|^2 + 2\langle \mathbf{a} | \mathbf{b} \rangle$.

If $\langle \mathbf{a} | \mathbf{b} \rangle = 0$ then the above gives Pythagorean Theorem (and vice versa).

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- $\begin{bmatrix} 1 \\ 0 \end{bmatrix}, \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ or $\begin{bmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{bmatrix}, \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix}$ or $\begin{bmatrix} \frac{1}{2} \\ \frac{\sqrt{3}}{2} \end{bmatrix}, \begin{bmatrix} \frac{-\sqrt{3}}{2} \\ \frac{1}{2} \end{bmatrix}$ are all orthonormal.

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- The next two correspond to $\{| \rightarrow \rangle, | \leftarrow \rangle\}$ and $\{| \nearrow \rangle, | \nwarrow \rangle\}$ respectively.

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- We can (easily) show that $\begin{bmatrix} a_1 \\ a_2 \end{bmatrix} = \frac{a_1 - a_2}{\sqrt{2}} |\rightarrow\rangle + \frac{a_1 + a_2}{\sqrt{2}} |\leftarrow\rangle$.

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- Form $A = [|b_1\rangle \dots |b_n\rangle]$ and take its transpose $A^T = \begin{bmatrix} \langle b_1 | \\ \vdots \\ \langle b_n | \end{bmatrix}$.

Linear combinations of basis vectors

- Let $|\psi\rangle$ be a n -dimensional ket and $\{|\mathbf{b}_1\rangle, |\mathbf{b}_2\rangle, \dots, |\mathbf{b}_n\rangle\}$ is a basis for \mathbb{R}^n .
- We want to write $|\psi\rangle = w_1|\mathbf{b}_1\rangle + \dots + w_n|\mathbf{b}_n\rangle$.
- We can (easily) show that $w_i = \langle \mathbf{b}_i | \psi \rangle$ and so
$$|\psi\rangle = \langle \mathbf{b}_1 | \psi \rangle |\mathbf{b}_1\rangle + \langle \mathbf{b}_2 | \psi \rangle |\mathbf{b}_2\rangle + \dots + \langle \mathbf{b}_n | \psi \rangle |\mathbf{b}_n\rangle.$$
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- Form $A = [|\mathbf{b}_1\rangle \dots |\mathbf{b}_n\rangle]$ and take its transpose $A^T = \begin{bmatrix} \langle \mathbf{b}_1 | \\ \vdots \\ \langle \mathbf{b}_n | \end{bmatrix}$.

- Then $\begin{bmatrix} w_1 \\ \vdots \\ w_n \end{bmatrix} = A^T |\psi\rangle = \begin{bmatrix} \langle \mathbf{b}_1 | \psi \rangle \\ \vdots \\ \langle \mathbf{b}_n | \psi \rangle \end{bmatrix}$.

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- Example: $\begin{bmatrix} 1 & 0 \\ 0 & 2 \end{bmatrix} \otimes [1 \ 2 \ 3] = ?$
- When tensor multiplying kets, the following are all equivalent: $|\psi\rangle \otimes |\phi\rangle = |\psi\rangle|\phi\rangle = |\psi\phi\rangle$. Tensor multiplication ***is associative***.