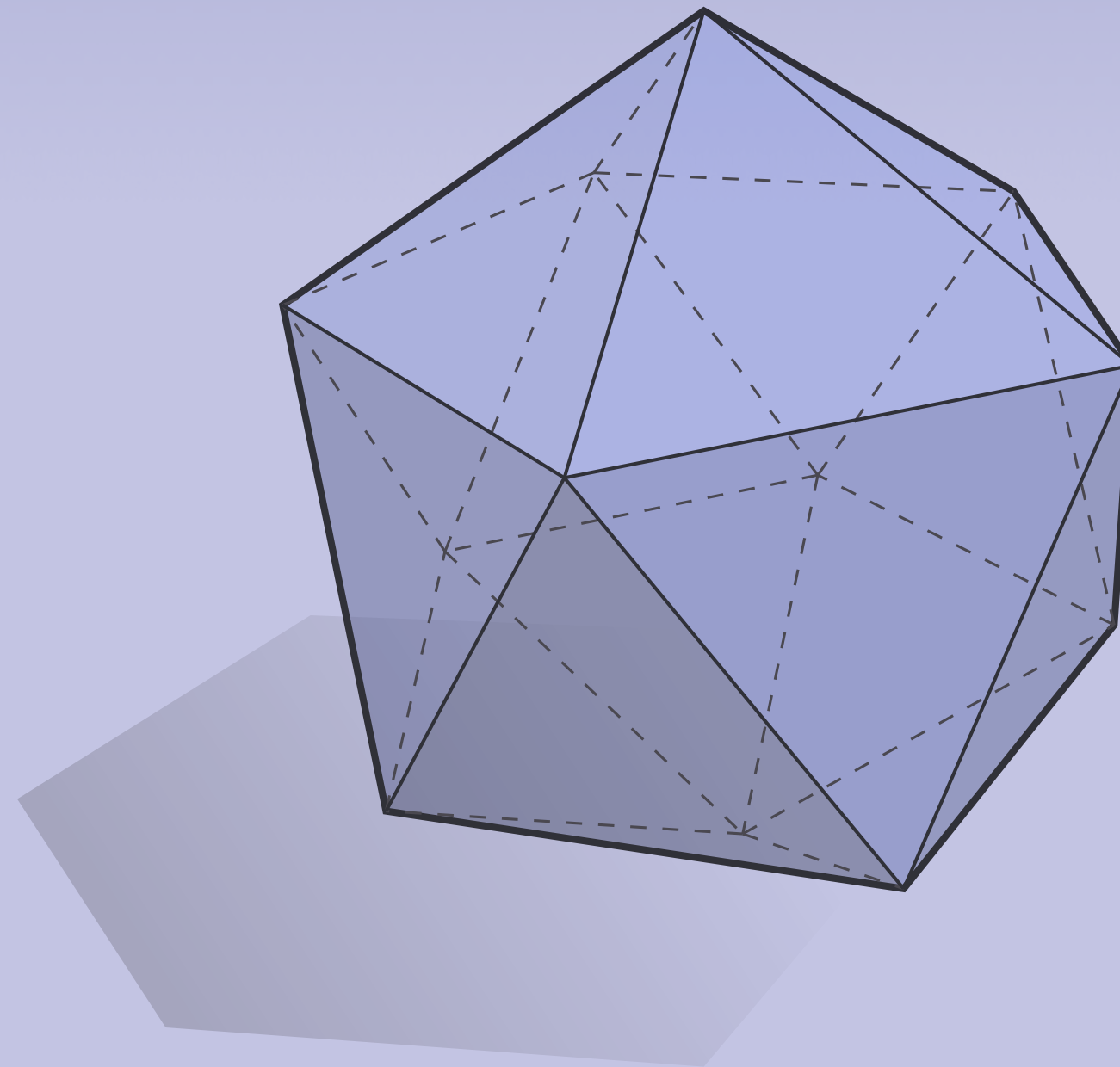


DISCRETE DIFFERENTIAL
GEOMETRY:
AN APPLIED INTRODUCTION

Keenan Crane • CMU 15-458/858B • Fall 2017

LECTURE 4:

DIFFERENTIAL FORMS IN R^n



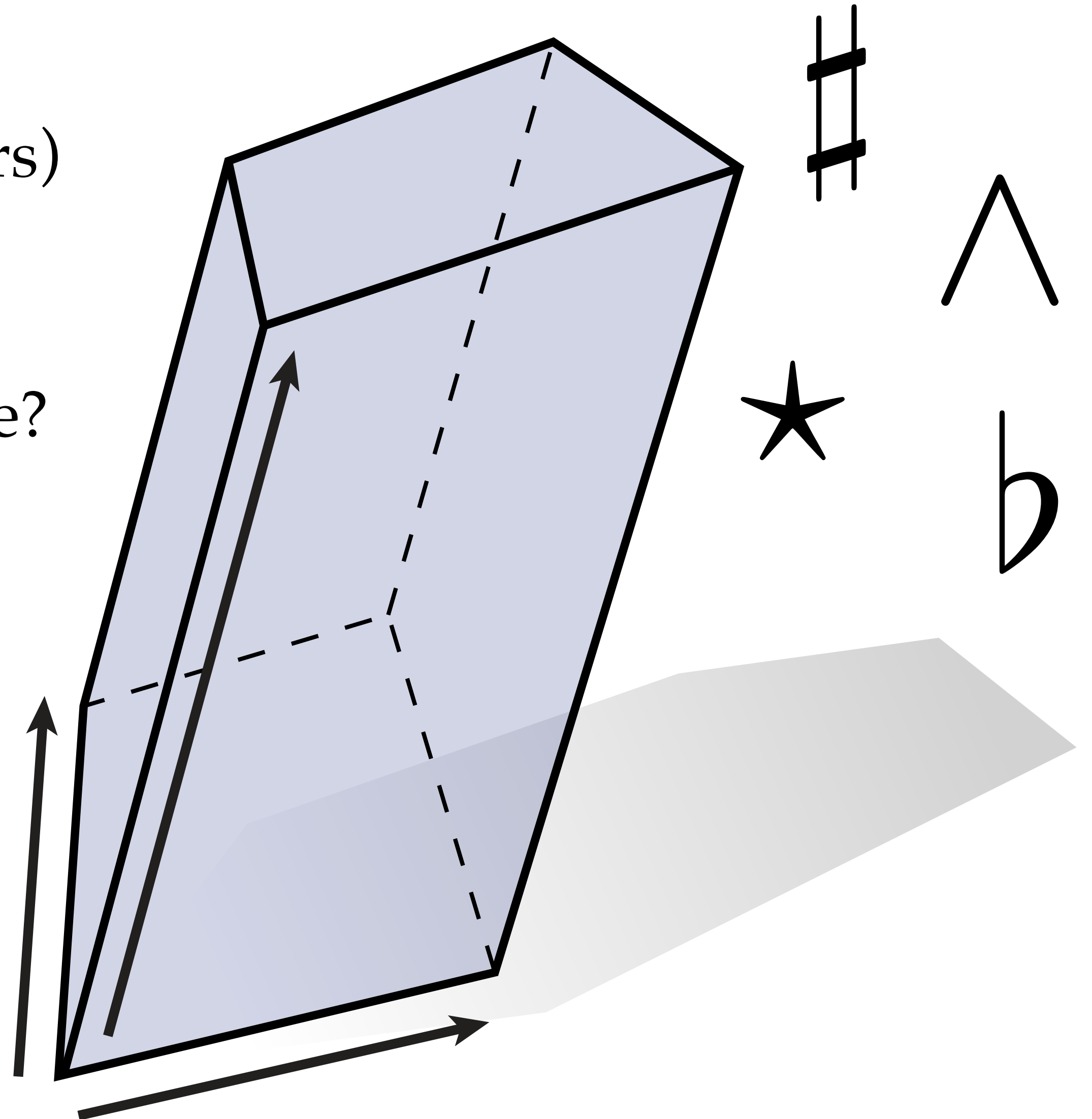
DISCRETE DIFFERENTIAL GEOMETRY:

AN APPLIED INTRODUCTION

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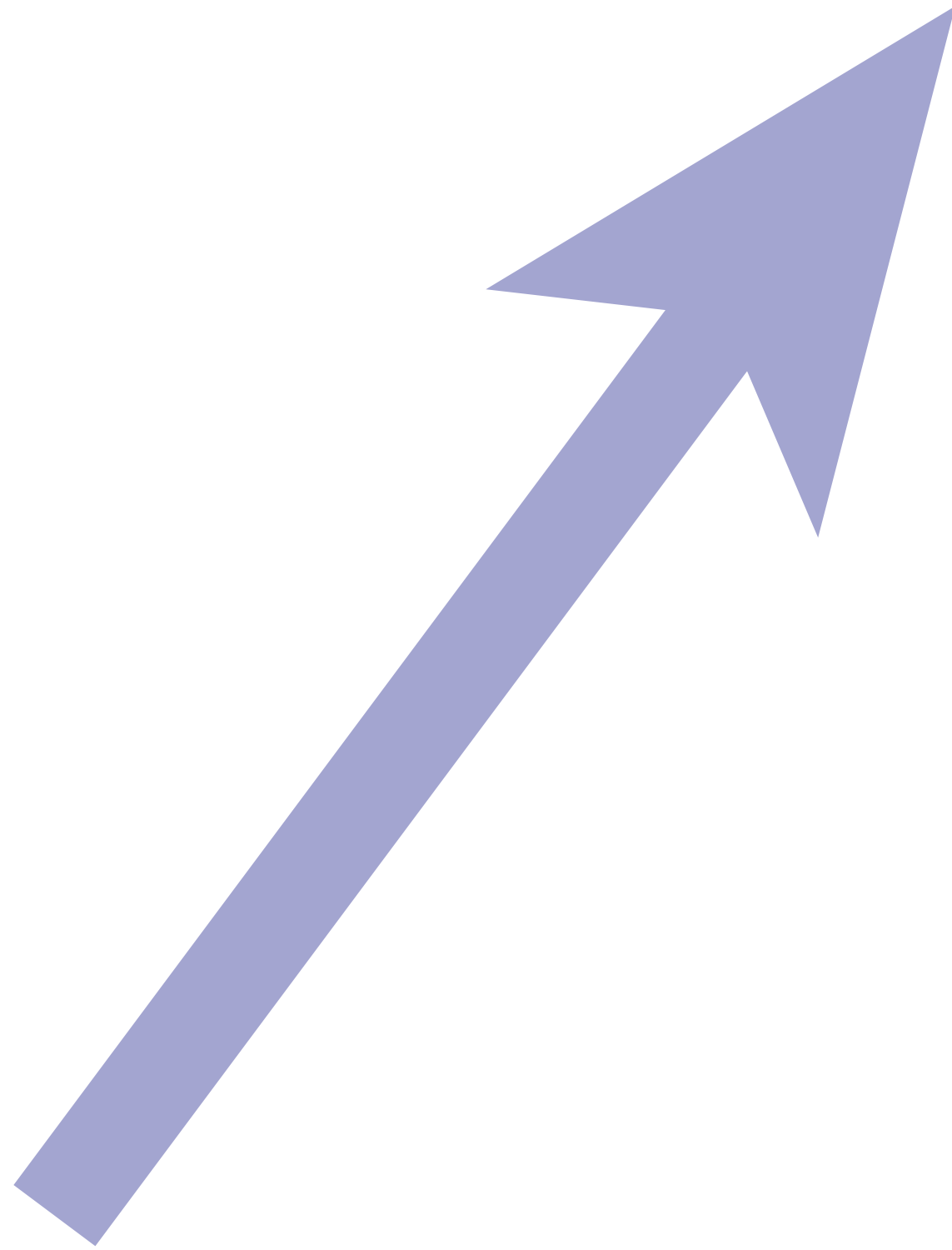
Exterior Calculus—Overview

- *Last time:*
 - **Exterior algebra**—“little volumes” (k -vectors)
- *This time:*
 - **Exterior calculus**—how do k -vectors change?
 - Before we considered just isolated k -vectors
 - Now we'll put a k -vector at each point
 - For now, stick to “flat” spaces (\mathbb{R}^n)
 - Later, consider “curved” spaces (manifolds)

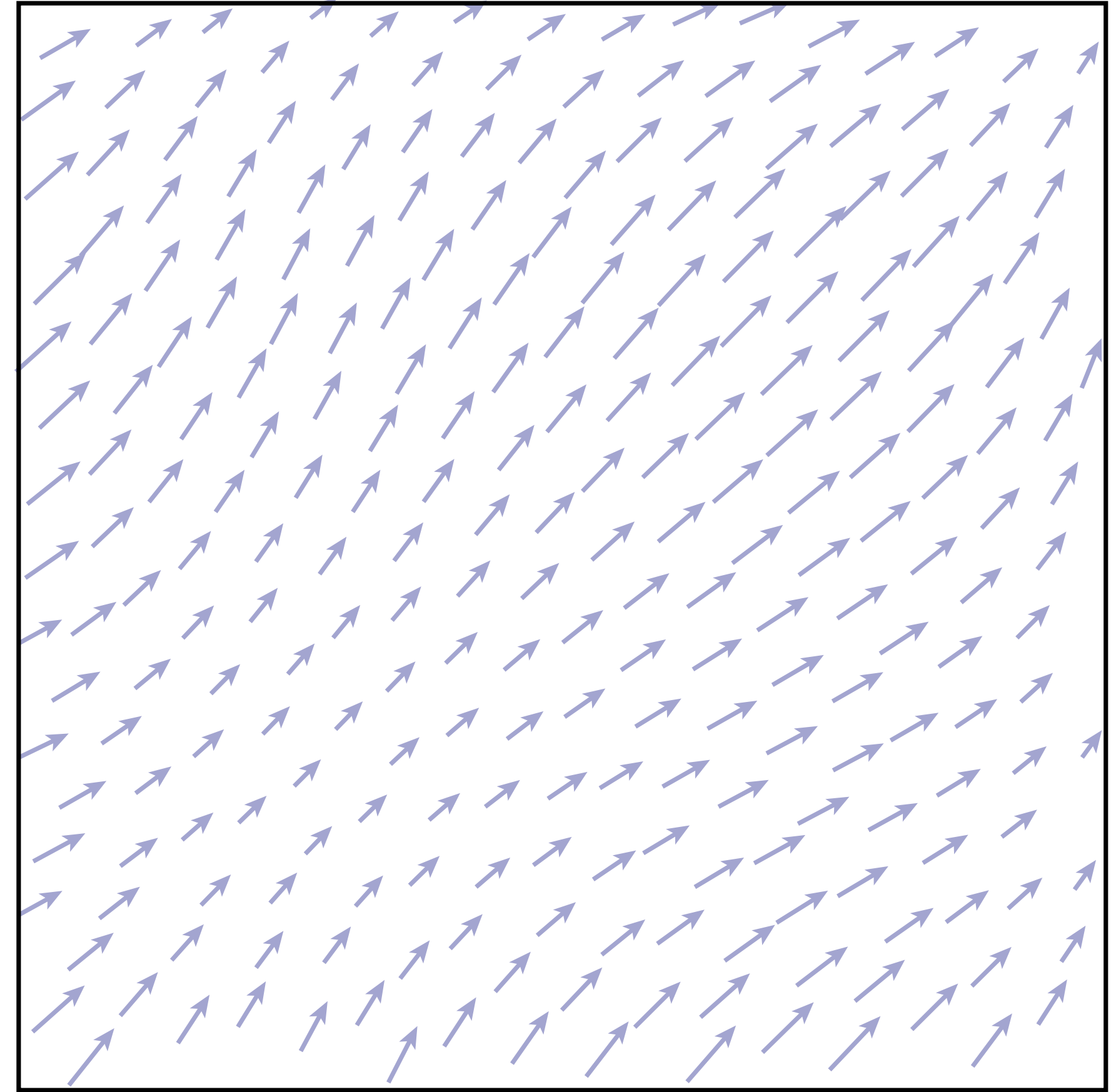


Review: Vector vs. Vector Field

- Recall that a vector *field* is an assignment of a vector to each point:



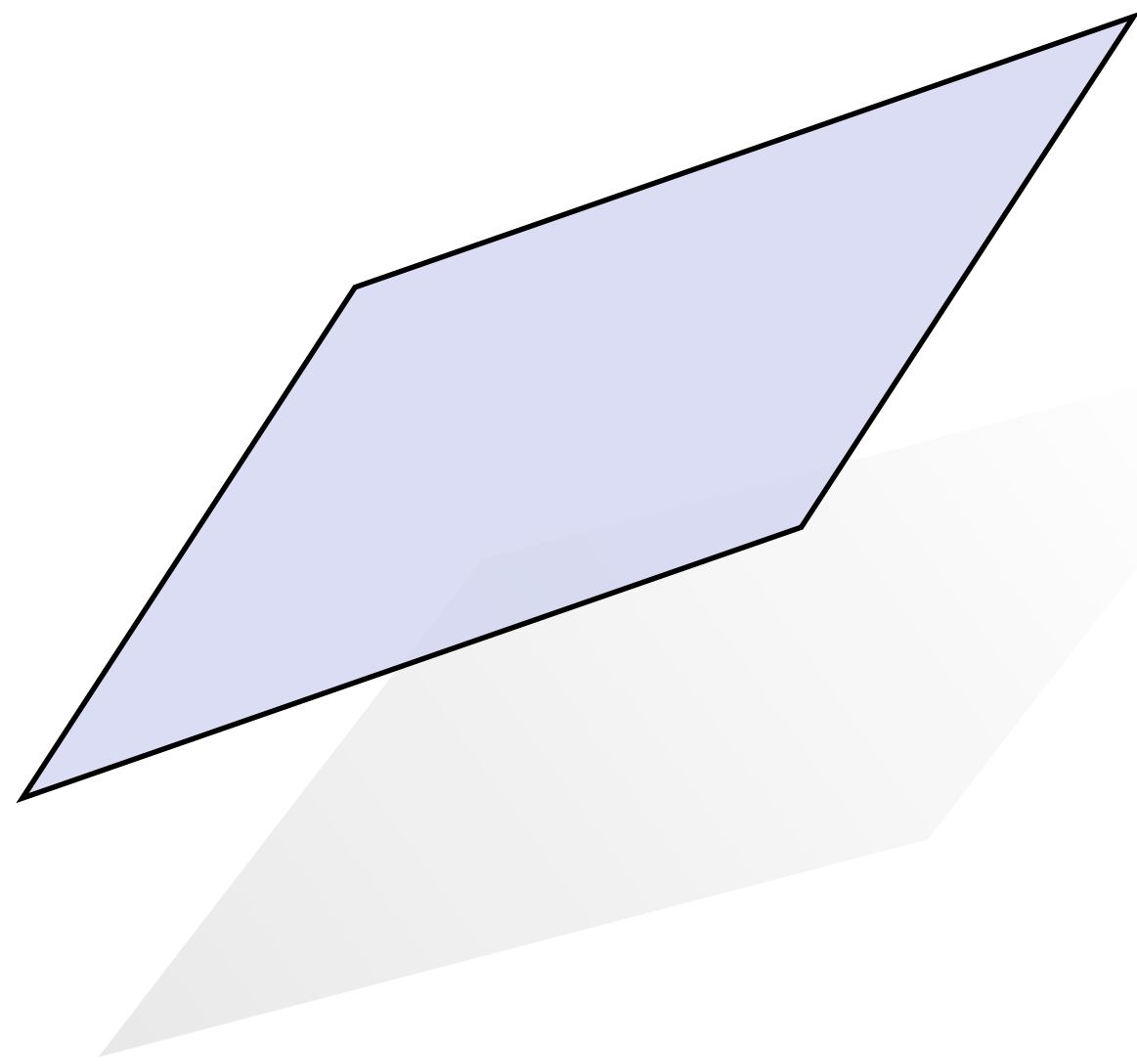
vector



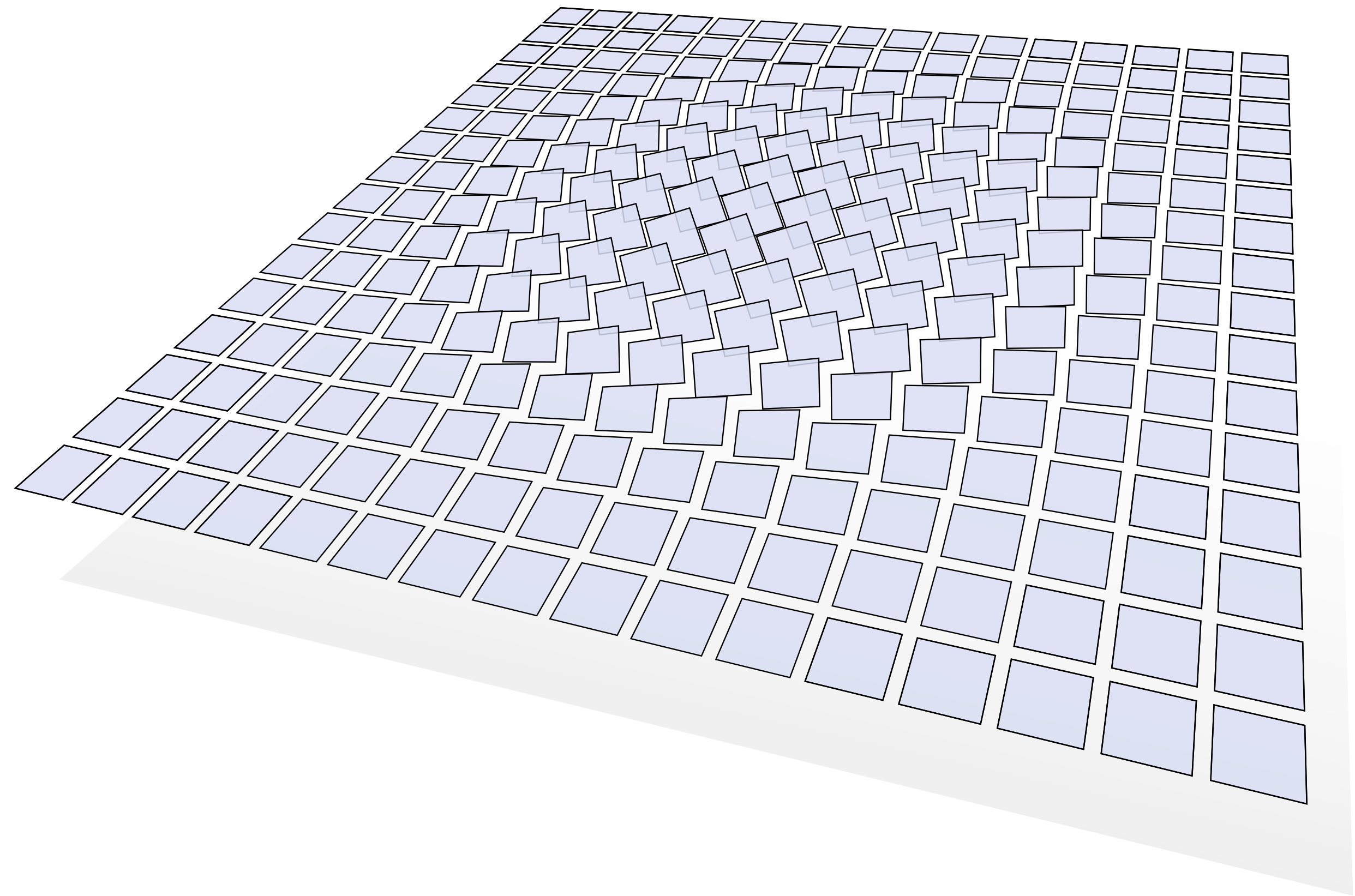
vector field

Preview: k -Vector vs. Differential Form

- A *differential k -form* will be an assignment of an object like a k -vector to each point:



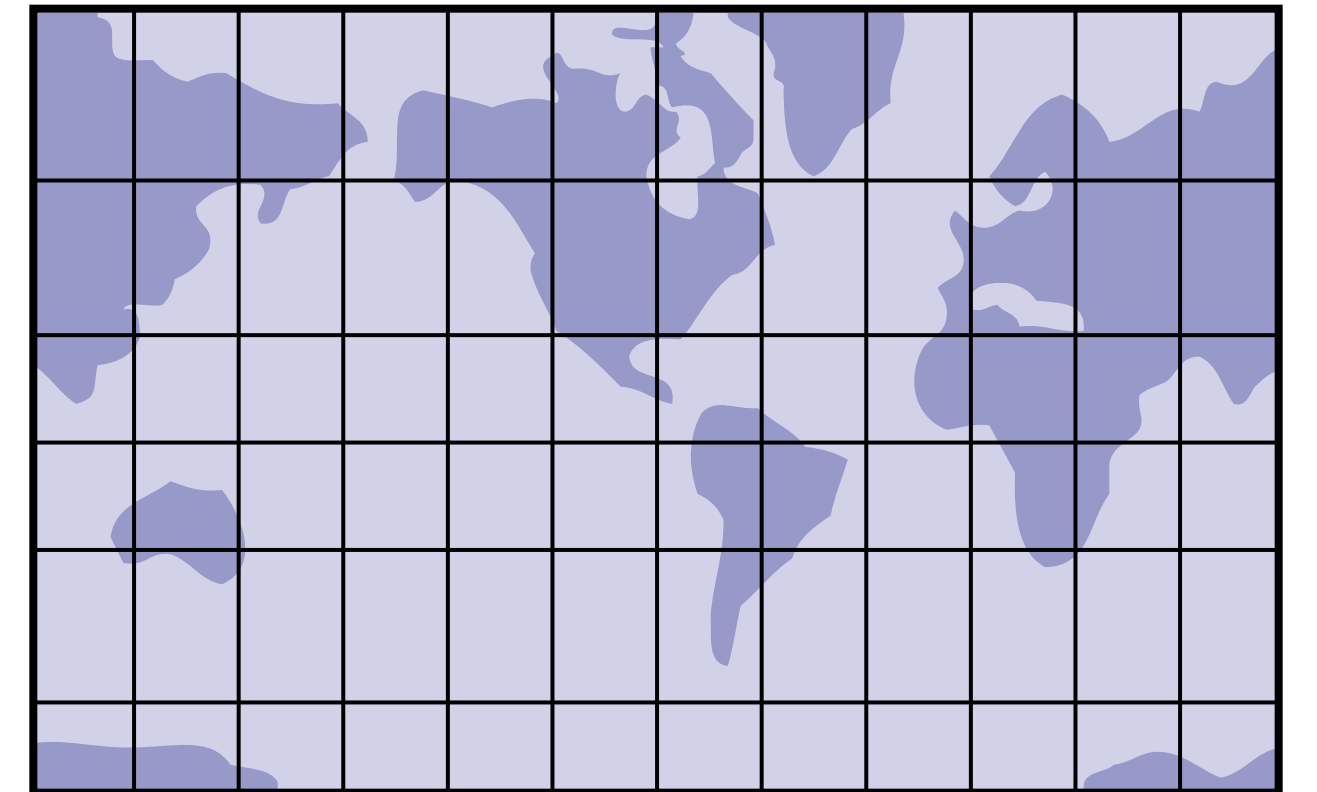
k -vector



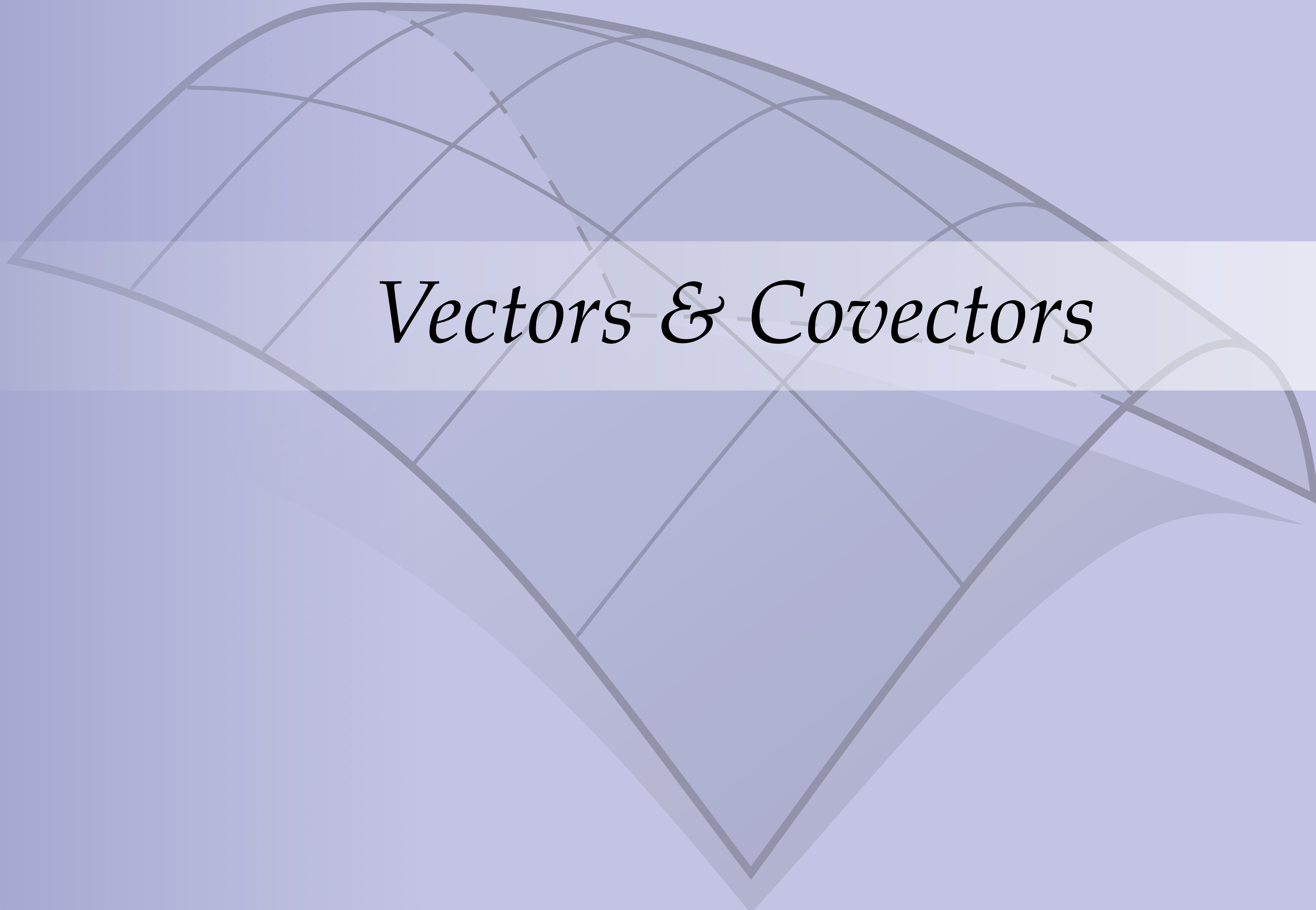
differential form

Exterior Calculus: Flat vs. Curved Spaces

- For now, we'll only consider *flat* spaces like the 2D plane
 - Keeps all our calculations simple
 - Don't have to define *manifolds* (yet!)
- True power of exterior calculus revealed on *curved* spaces
 - In flat spaces, vectors and forms look *very* similar*
 - Difference will be less superficial on curved spaces
 - Close relationship to *curvature* (geometry)
 - Also close relationship to *mass* (physics)



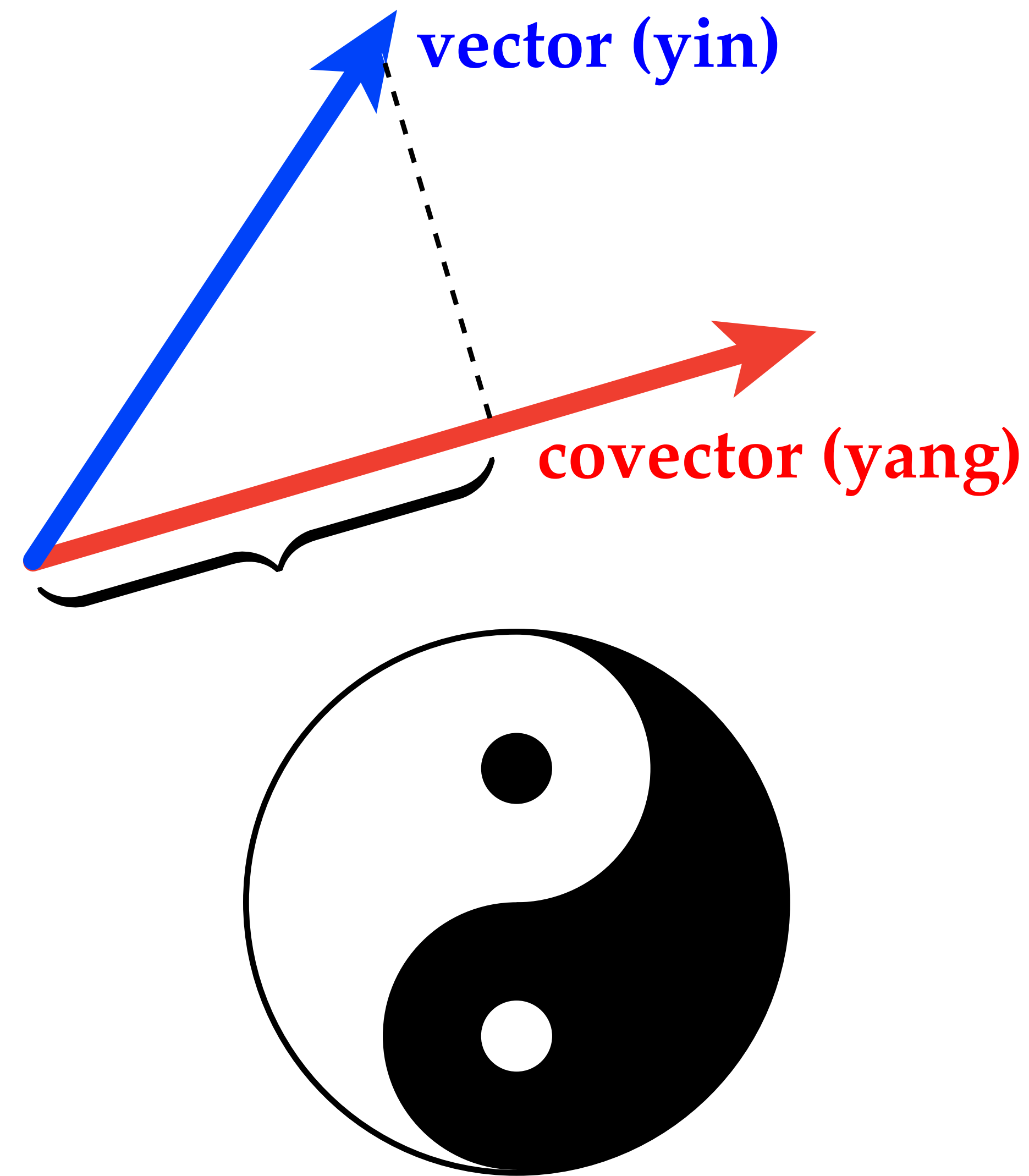
*So similar in fact you may wonder why we even bother to make a distinction! Well, stay tuned... :-)



Vectors & Covectors

Vector-Covector Duality

- Much of the expressivity of exterior calculus comes from the relationship between *vectors* and *covectors*.
- Loosely speaking:
 - **covectors** are objects that “*measure*”
 - **vectors** are objects that “*get measured*”
- We say that vectors and covectors are “dual”—roughly speaking, they are two sides of the same coin (but with important differences...).
- Ultimately, differential forms will be built up by taking wedges of covectors (rather than vectors).



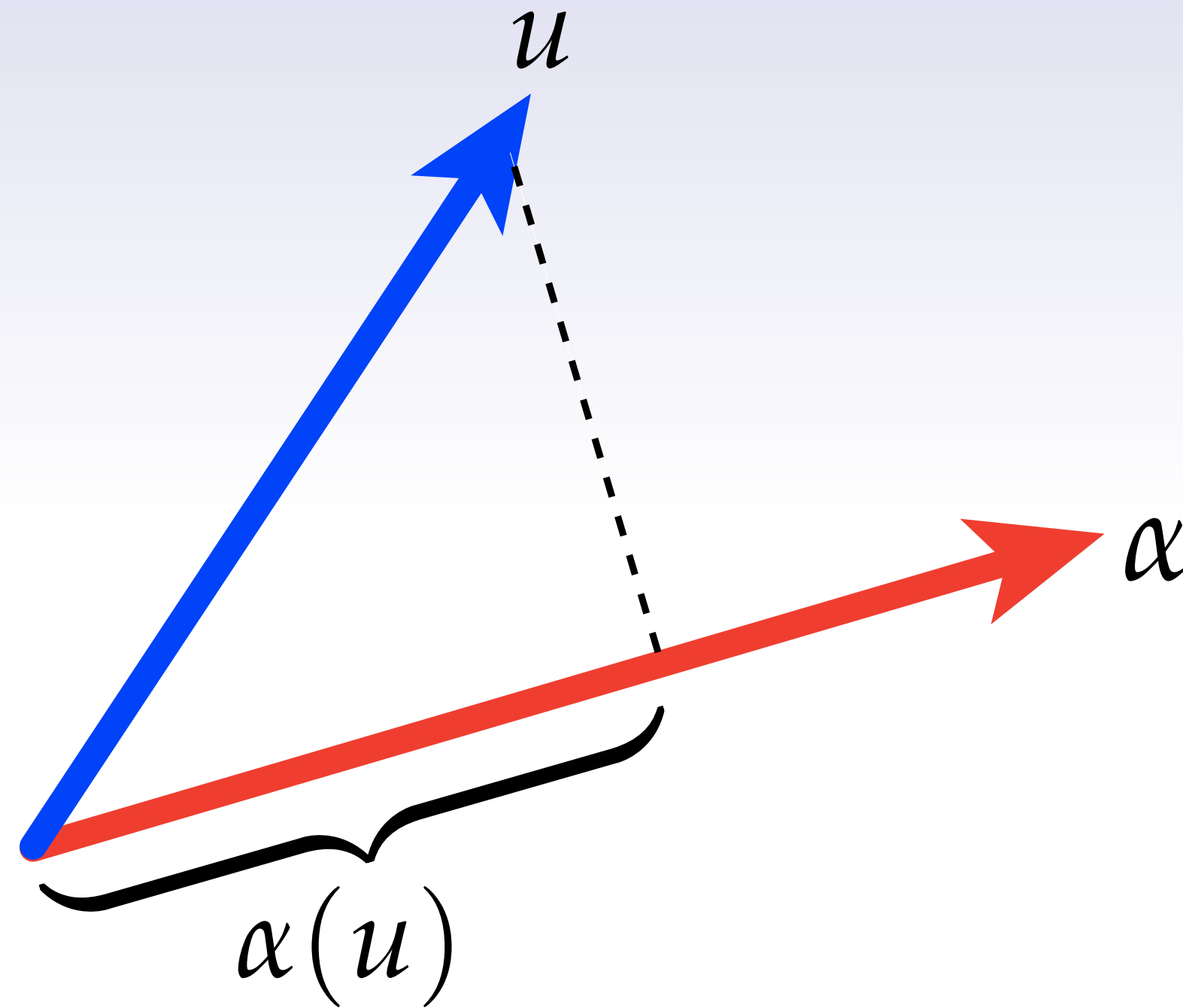
Analogy: Row & Column Vectors

In matrix algebra, we make a distinction between *row vectors* and *column vectors*:

$$\begin{bmatrix} \alpha_1 & \alpha_2 & \alpha_3 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \end{bmatrix}$$

Q: Why do we make the distinction? What does it mean geometrically?
What does it mean as a *linear map*? (Is this convention useful?)

Vectors and Covectors



Analogy: *row vs. column vector*

Key idea: a covector *measures* length of vector along a particular direction

Dual Space & Covectors

Definition. Let V be any real vector space. Its *dual space* V^* is the collection of all linear functions $\alpha : V \rightarrow \mathbb{R}$ together with the operations of *addition*

$$(\alpha + \beta)(u) := \alpha(u) + \beta(u)$$

and *scalar multiplication*

$$(c\alpha)(u) := c(\alpha(u))$$

for all $\alpha, \beta \in V^*$, $u \in V$, and $c \in \mathbb{R}$.

Definition. An element of a dual vector space is called a *dual vector* or a *covector*.

(Note: unrelated to *Hodge dual*!)

Covectors — Example (R^3)

- As a concrete example, let's consider the vector space $V = R^3$
- Recall that a map f is *linear* if for all vectors \mathbf{u} , \mathbf{v} and scalars a , we have

$$f(\mathbf{u} + \mathbf{v}) = f(\mathbf{u}) + f(\mathbf{v}) \quad \text{and} \quad f(a\mathbf{u}) = af(\mathbf{u})$$

- **Q:** What's an example of a *linear* map from R^3 to R ?

- Suppose we express our vectors in coordinates $\mathbf{u} = (x, y, z)$

$$\begin{bmatrix} x \\ y \\ z \end{bmatrix}$$

vector

- **A:** One of *many* possible examples: $f(x, y, z) = x + 2y + 3z$

- **Q:** What are *all* the possibilities?

$$\begin{bmatrix} a & b & c \end{bmatrix}$$

covector

- **A:** They all just look like $f(x, y, z) = ax + by + cz$ for constants a, b, c

- In other words in Euclidean R^3 , a *covector* looks like just another 3-vector!

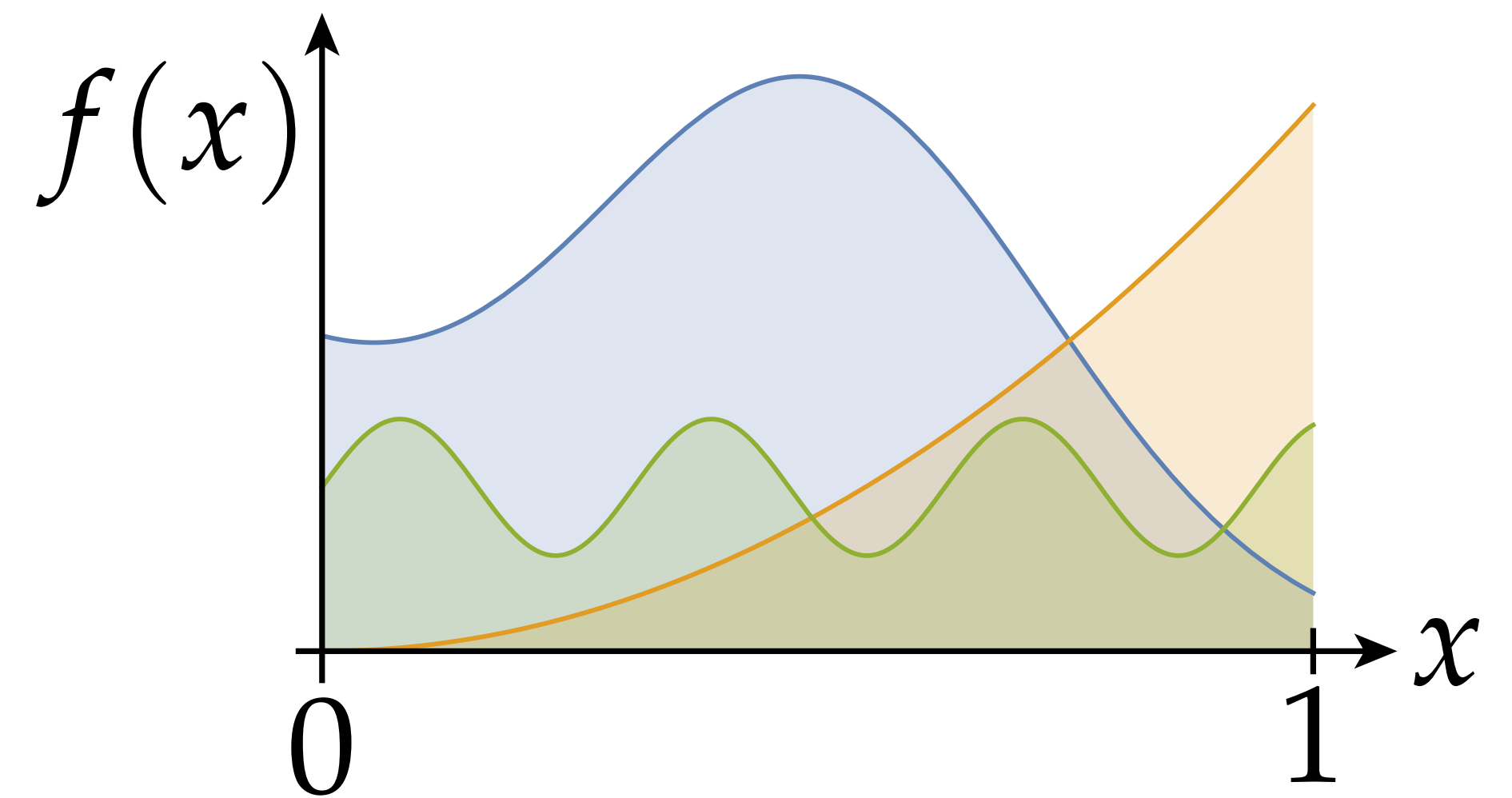
Covectors—Example (Functions)

- If covectors are just the same as vectors, why even bother?
- Here's a more interesting example:

Example. Let V be the set of integrable functions $f : [0, 1] \rightarrow \mathbb{R}$, and consider maps

- $\phi : V \rightarrow \mathbb{R}; f \mapsto \int_0^1 f(x) dx$
- $\delta : V \rightarrow \mathbb{R}; f \mapsto f(0)$

Is V a vector space? Are ϕ and δ covectors?



Key idea: the difference between primal & dual vectors is not merely superficial!

Sharp and Flat

$$\begin{bmatrix} u_1 \\ u_2 \\ u_3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix} \xrightarrow{\text{T}} \begin{bmatrix} u_1 & u_2 & u_3 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix}$$

$$u, v \xrightarrow{b} u^b(v)$$

$$\alpha, \beta \xrightarrow{\sharp} \alpha(\beta^\sharp)$$

Analogy: *transpose*

(What's up with the musical symbols? Will see in a bit...)

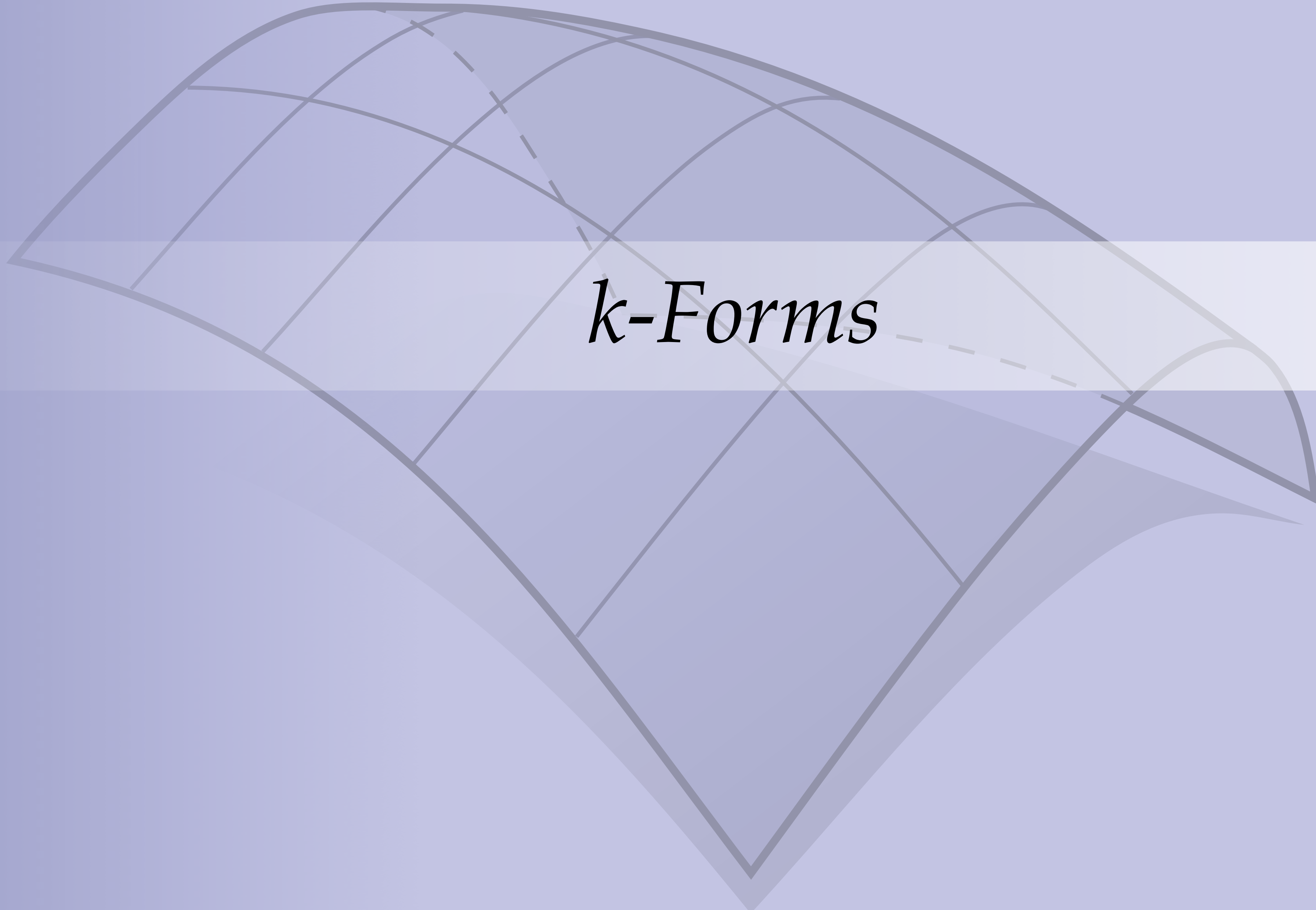
Sharp and Flat w/ Inner Product

$$\begin{bmatrix} u_1 & u_2 & u_3 \end{bmatrix} \begin{bmatrix} M_{11} & M_{12} & M_{13} \\ M_{12} & M_{22} & M_{23} \\ M_{13} & M_{23} & M_{33} \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix}$$

$$u^\flat(v) = u^\top M v \quad \Longleftrightarrow \quad u^\flat(\cdot) = \langle u, \cdot \rangle$$

$$\alpha(\beta^\sharp) = \alpha M^{-1} \beta^\top \quad \Longleftrightarrow \quad \langle \alpha^\sharp, \cdot \rangle = \alpha(\cdot)$$

Basic idea: applying the flat of a vector is the same as taking the inner product; taking the inner product w/ the sharp is same as applying the original 1-form.



k -Forms

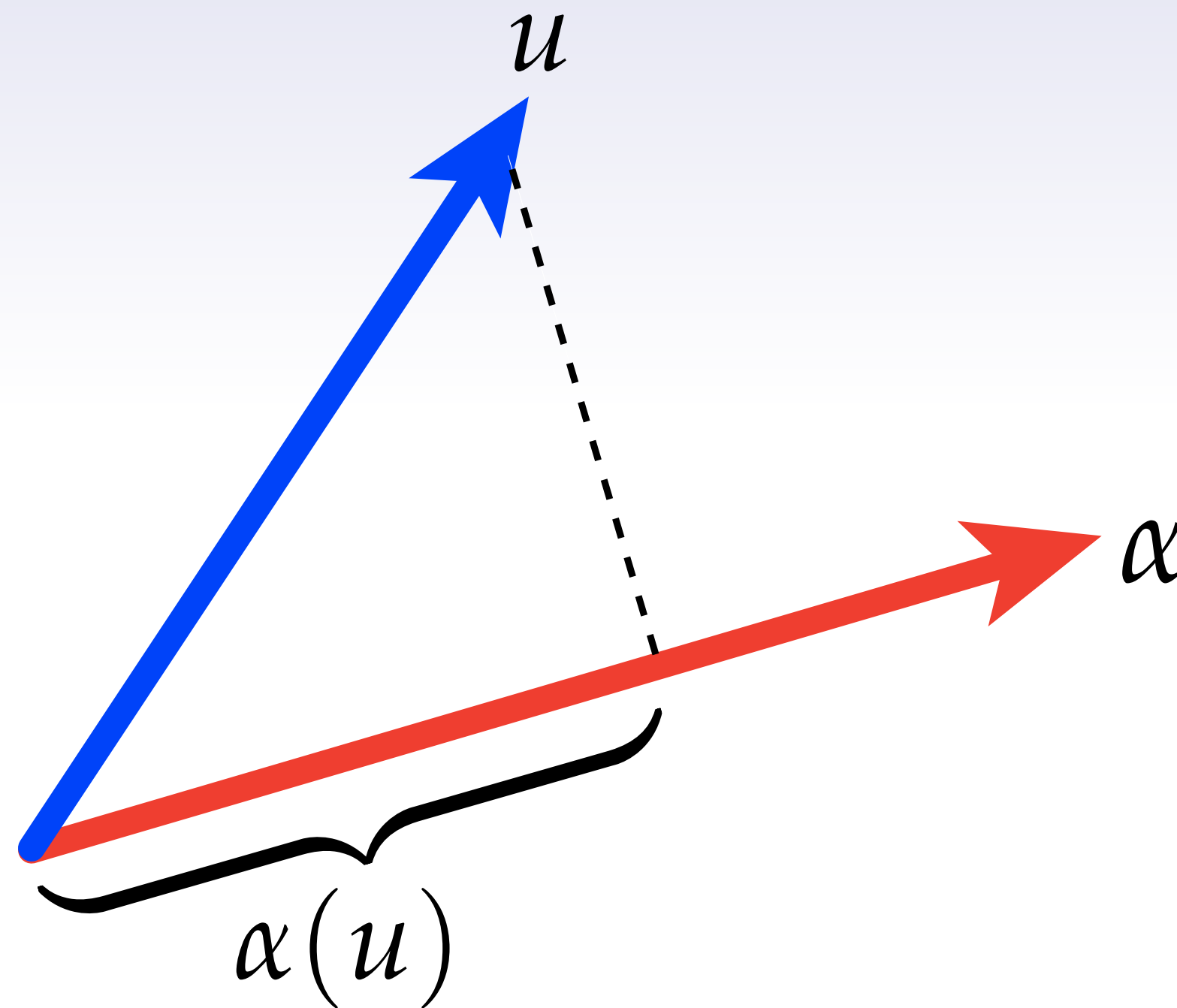
Covectors, Meet Exterior Algebra

- So far we've studied two distinct concepts:
 - **exterior algebra**—build up “volumes” from vectors
 - **covectors**—linear maps from vectors to scalars
- Combine to get an *exterior algebra of covectors*
 - Will call these objects *k-forms*
 - Just as a covector measures vectors...
 - ...a *k*-form will *measure* *k*-vectors.
 - In particular, measurements will be **multilinear**, *i.e.*, linear in each 1-vector.

	primal	dual
vector space	vectors	covectors
exterior algebra	<i>k</i> -vectors	<i>k</i> -forms

Measurement of Vectors

Geometrically, what does it mean to take a **multilinear** measurement of a 1-vector?

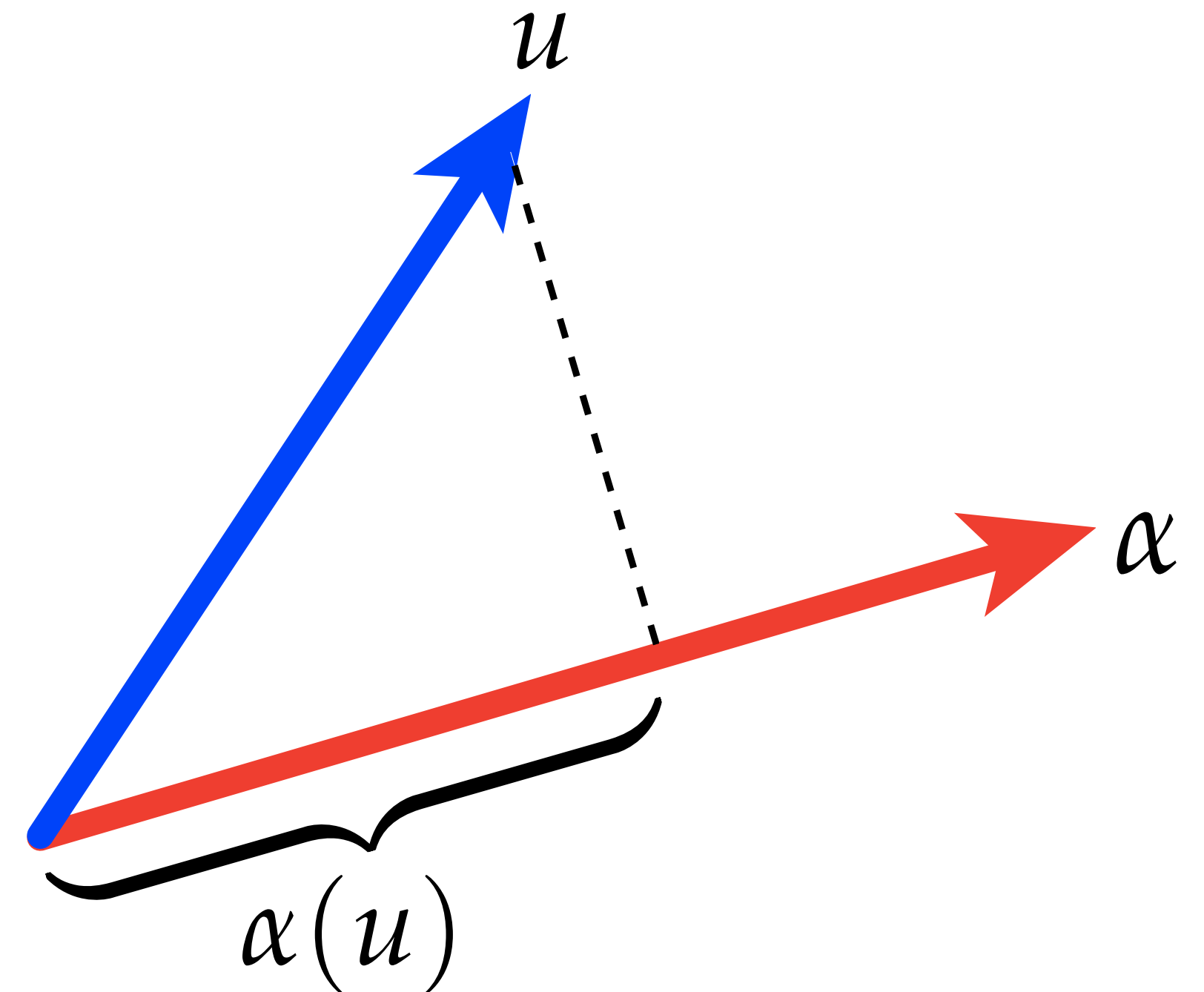


Observation: only thing we can do is measure extent along some other vector.

Computing the Projected Length

- Concretely, how do we compute projected length of one vector along another?
 - If α has unit norm then, we can just take the usual *dot product*
 - Since we think of u as the vector “*getting measured*” and α as the vector “*doing the measurement*”, we’ll write this as a function $\alpha(u)$:

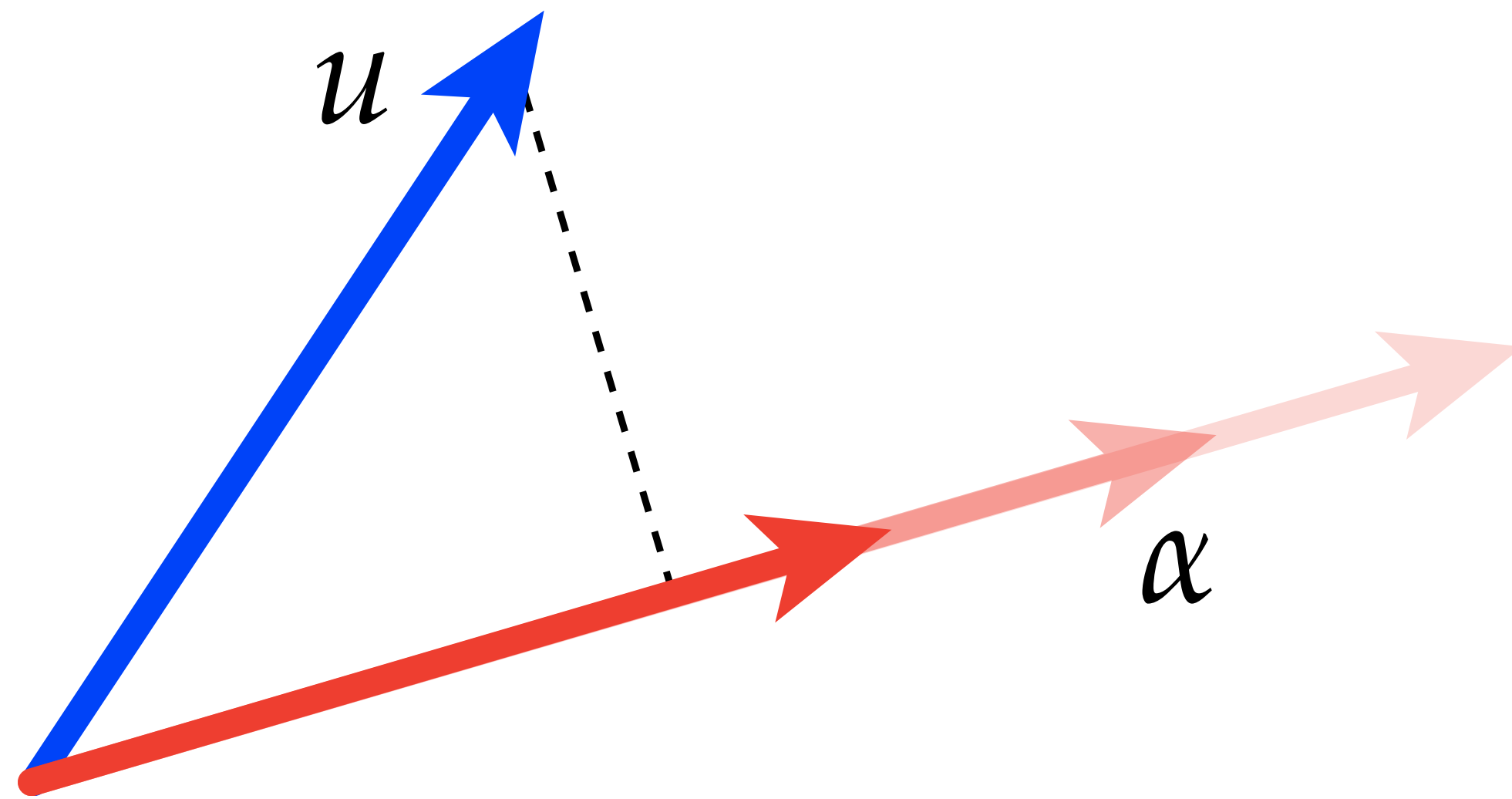
$$\alpha(u) = \sum_{i=1}^n \alpha_i u^i$$



1-form

We can of course apply this same expression when α does not have unit length:

$$\alpha(u) := \sum_i \alpha_i u^i$$

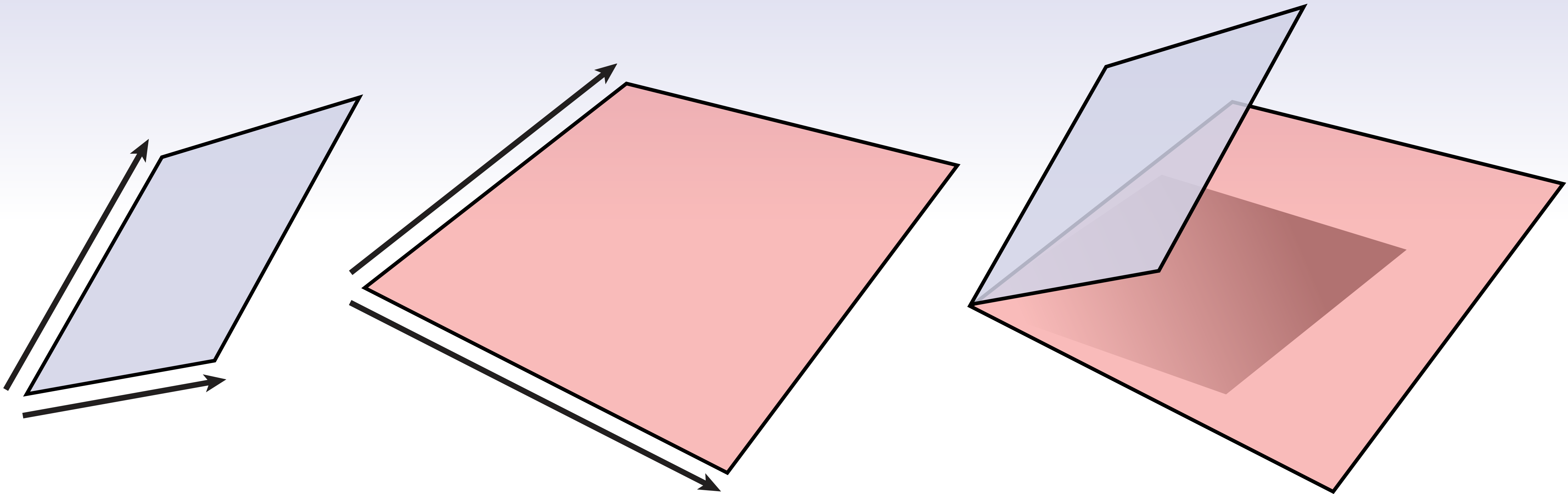


Interpretation?

Projected length gets scaled by magnitude of α .

Measurement of 2-Vectors

Geometrically, what does it mean to take a **multilinear** measurement of a 2-vector?



Intuition: size of “shadow” of one parallelogram on another.

Computing the Projected Area

- Concretely, how do we compute projected area of a parallelogram onto a plane?
 - First, project vectors defining parallelogram (u,v)
 - Then apply standard formula for area (cross product)
- Suppose for instance α, β are orthonormal:

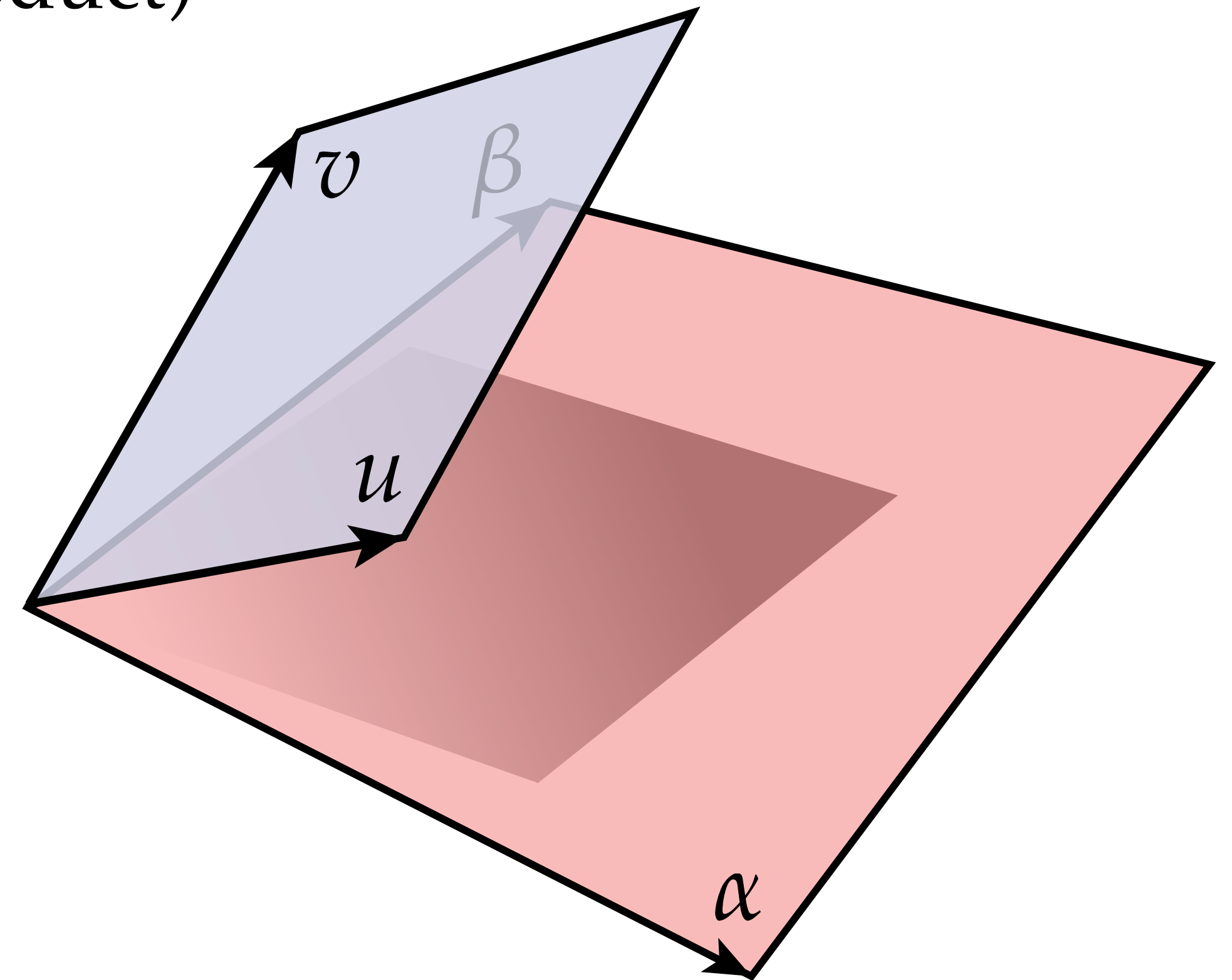
Projection

$$u \mapsto (\alpha(u), \beta(u))$$

$$v \mapsto (\alpha(v), \beta(v))$$

Area

$$\alpha(u)\beta(v) - \alpha(v)\beta(u)$$

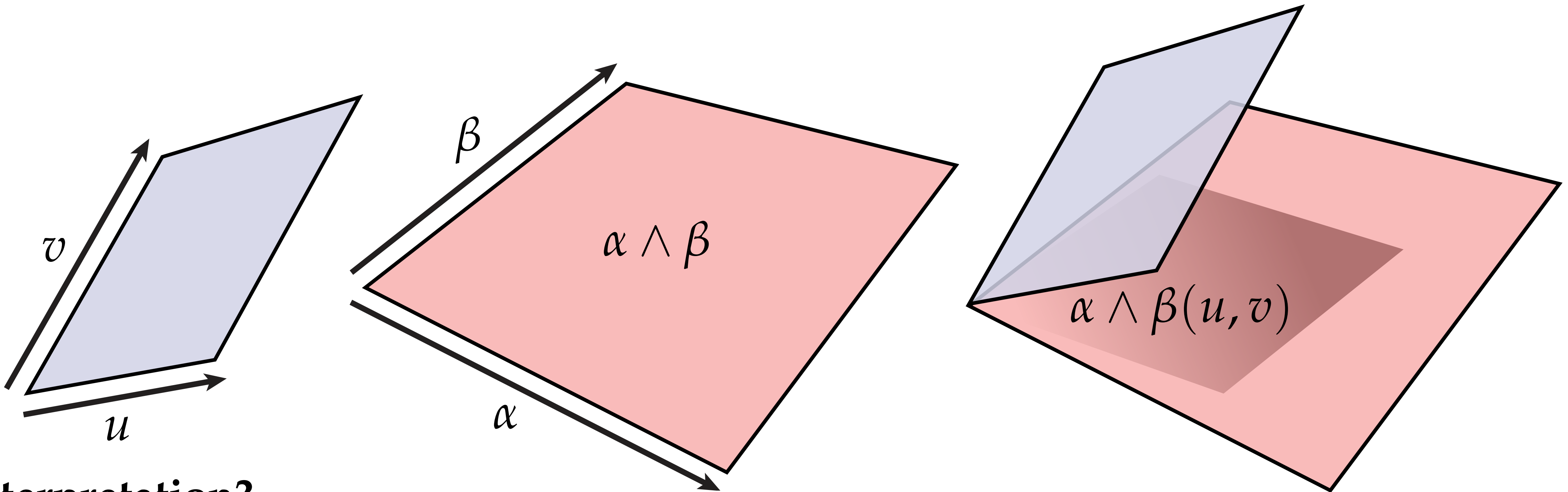


(Notice that in the projection we are treating α, β as 1-forms.)

2-form

We can of course apply this same expression when α, β are not orthonormal:

$$(\alpha \wedge \beta)(u, v) := \alpha(u)\beta(v) - \alpha(v)\beta(u)$$



Interpretation?

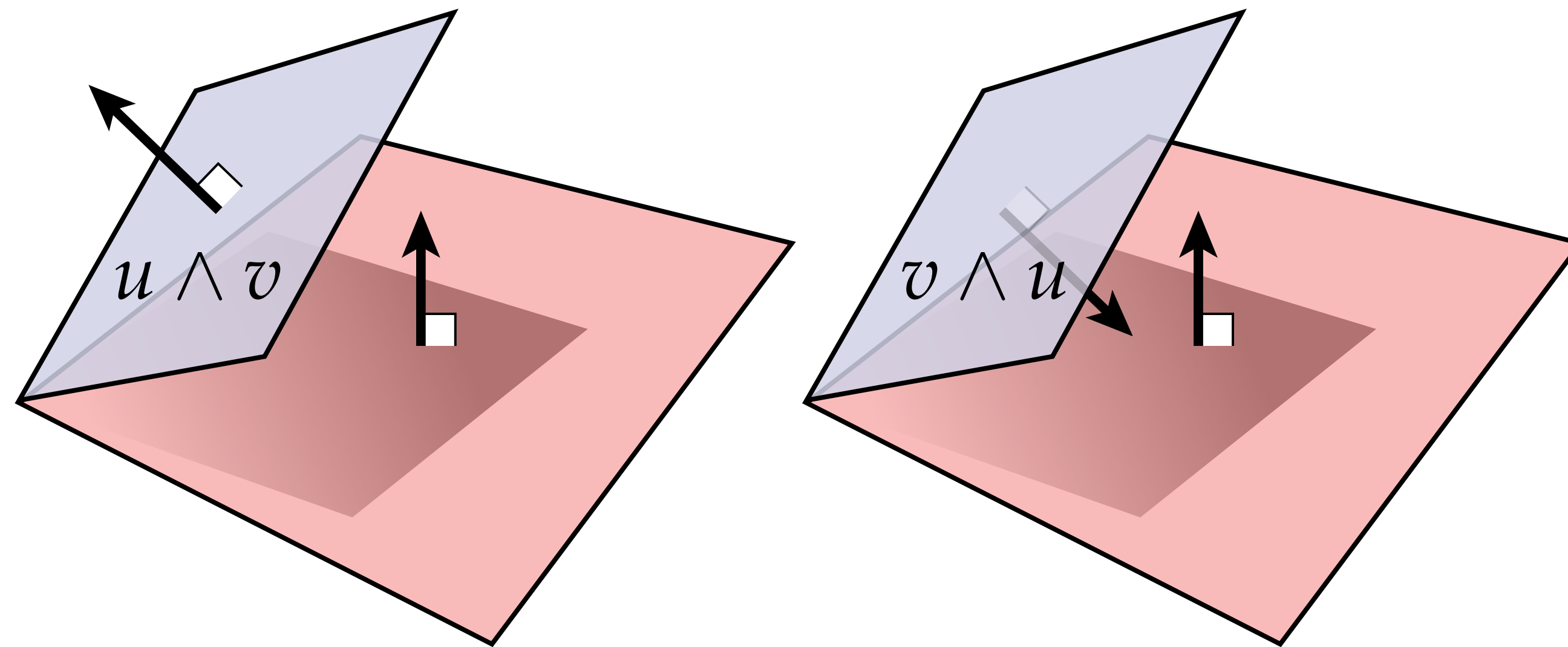
Projected area of u, v gets scaled by area of parallelogram with edges α, β .

Antisymmetry of 2-Forms

Notice that exchanging the arguments of a 2-form reverses sign:

$$\begin{aligned}(\alpha \wedge \beta)(v, u) &= \alpha(v)\beta(u) - \alpha(u)\beta(v) \\ &= -(\alpha(u)\beta(v) - \alpha(v)\beta(u)) \\ &= -(\alpha \wedge \beta)(u, v)\end{aligned}$$

Q: What does this *antisymmetry* mean geometrically?



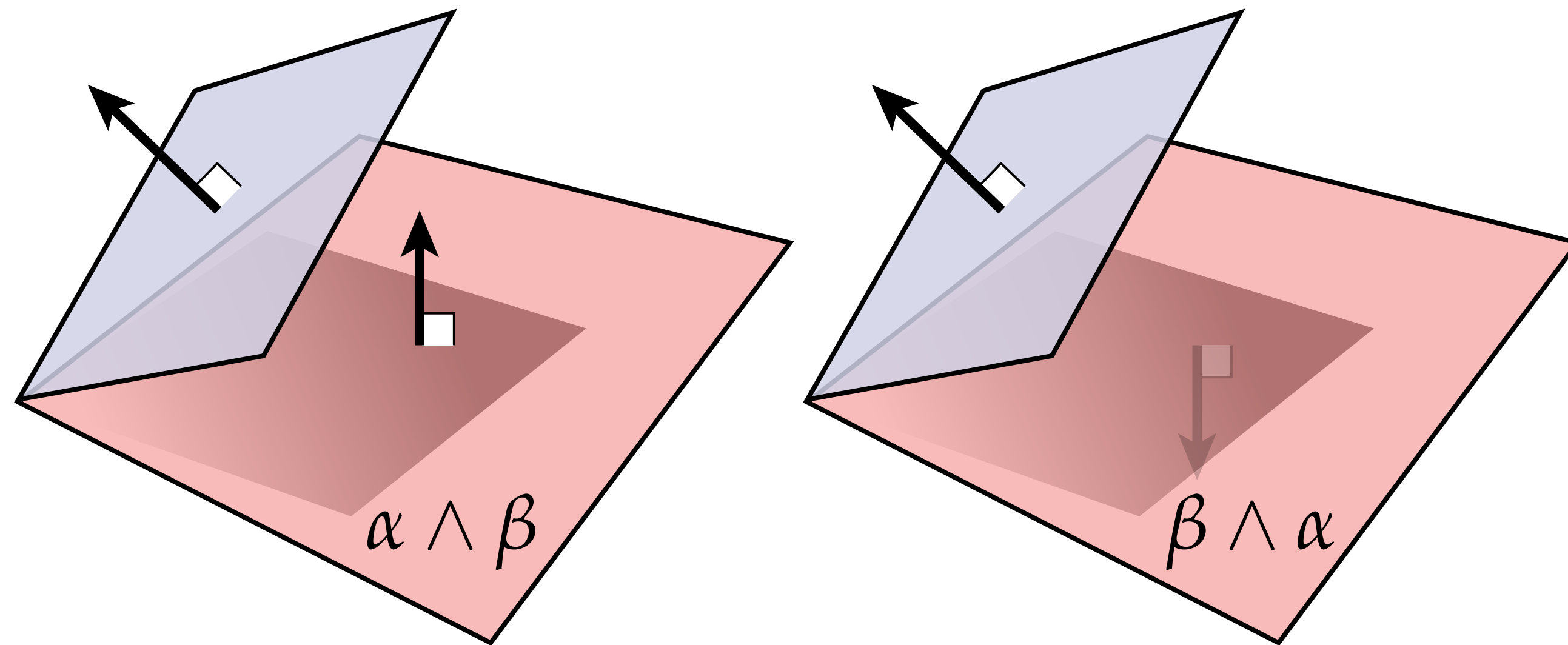
A: Opposite *orientations* of argument 2-vector.

Antisymmetry of 2-Forms

Remember that exchanging the arguments to a wedge product *also* reverses sign:

$$\begin{aligned}(\beta \wedge \alpha)(u, v) &= \beta(u)\alpha(v) - \beta(v)\alpha(u) \\ &= -(\alpha(u)\beta(v) - \alpha(v)\beta(u)) \\ &= -(\alpha \wedge \beta)(u, v)\end{aligned}$$

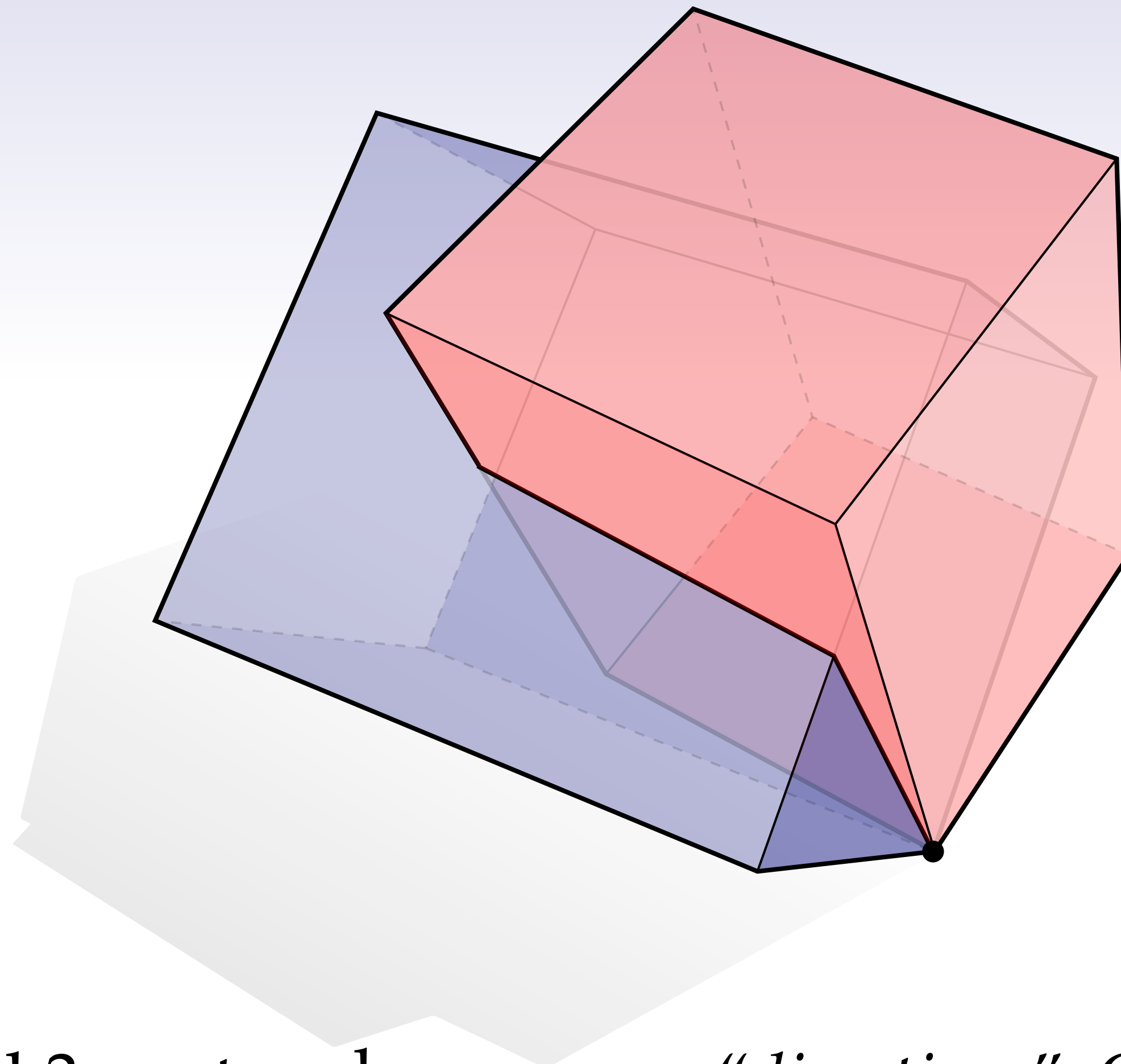
Q: What does this other kind of antisymmetry mean geometrically?



A: Opposite *orientations* of argument 2-vector.

Measurement of 3-Vectors

Geometrically, what does it mean to take a **multilinear** measurement of a 3-vector?



Observation: in R^3 , all 3-vectors have same “*direction*.” Can only measure *magnitude*.

Computing the Projected Volume

- Concretely, how do we compute the volume of a parallelepiped w/ edges u, v, w ?
 - Suppose (α, β, γ) is an orthonormal basis
 - Project vectors u, v, w onto this basis
 - Then apply standard formula for volume (determinant)

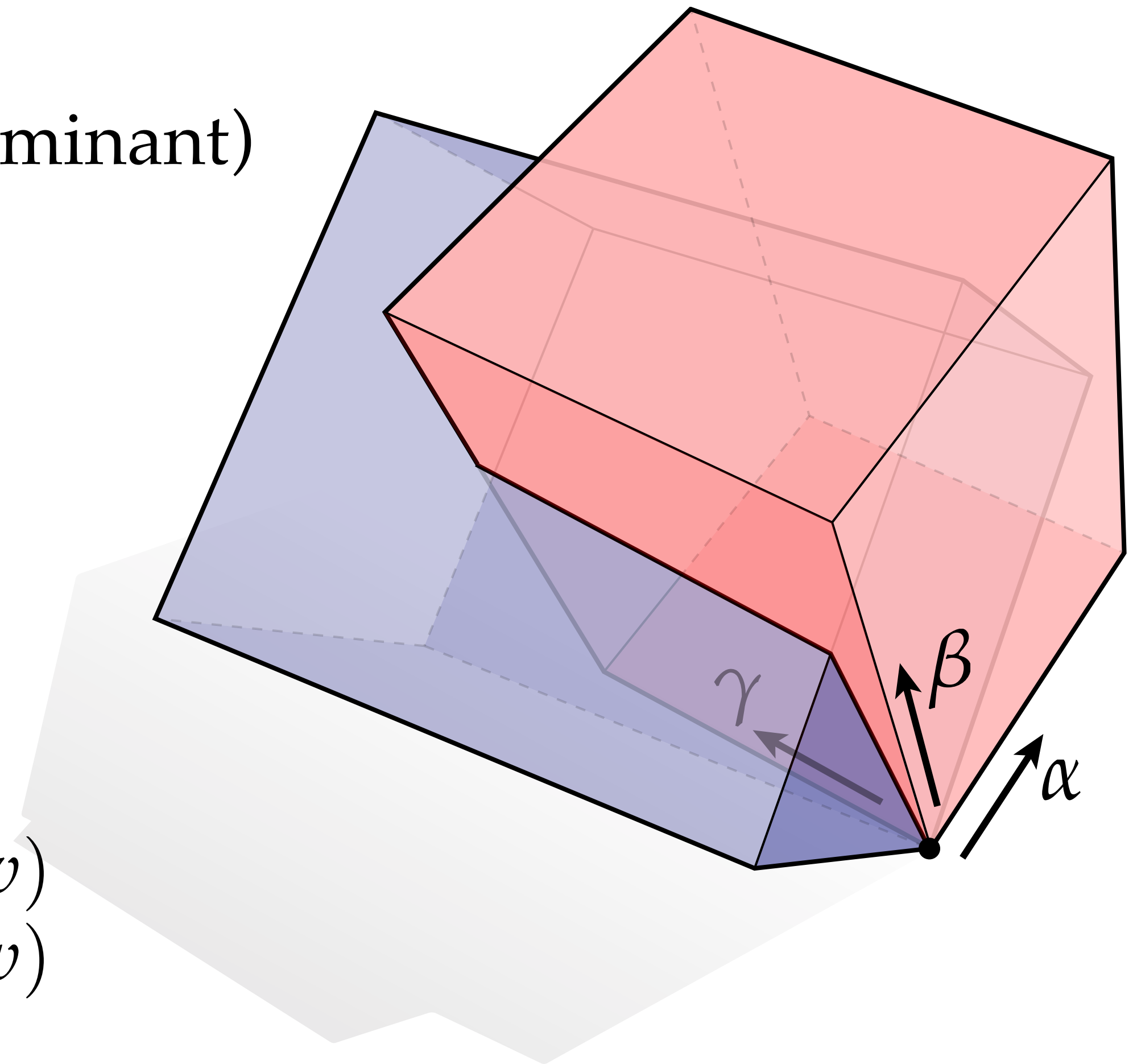
Projection

$$\begin{aligned} u &\mapsto (\alpha(u), \beta(u), \gamma(u)) \\ v &\mapsto (\alpha(v), \beta(v), \gamma(v)) \\ w &\mapsto (\alpha(w), \beta(w), \gamma(w)) \end{aligned}$$

Volume

$$\begin{vmatrix} \alpha(u) & \alpha(v) & \alpha(w) \\ \beta(u) & \beta(v) & \beta(w) \\ \gamma(u) & \gamma(v) & \gamma(w) \end{vmatrix}$$

$$\begin{aligned} &= \alpha(u)\beta(v)\gamma(w) + \alpha(v)\beta(w)\gamma(u) + \alpha(w)\beta(u)\gamma(v) \\ &\quad - \alpha(u)\beta(w)\gamma(v) - \alpha(w)\beta(v)\gamma(u) - \alpha(v)\beta(u)\gamma(w) \end{aligned}$$



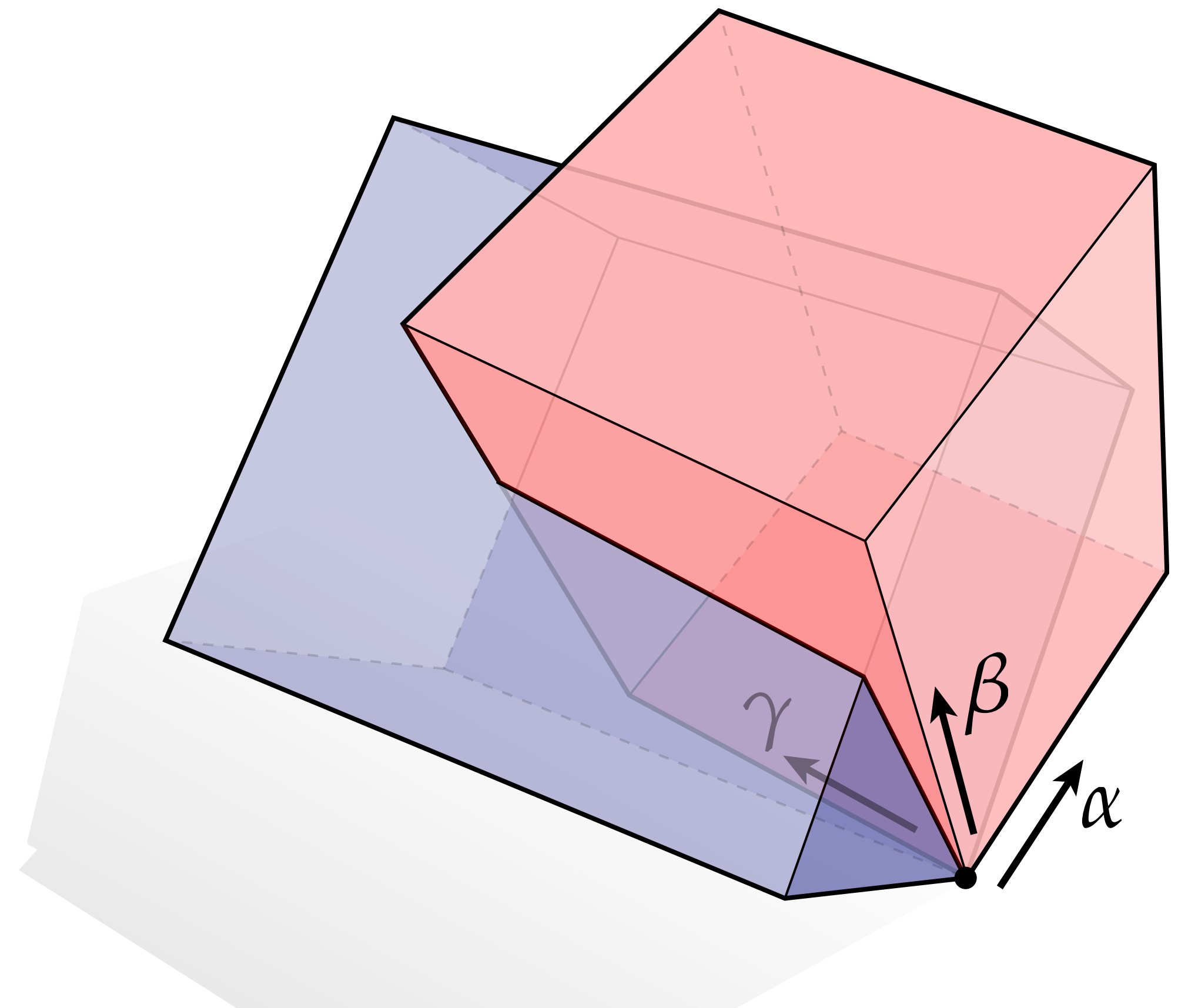
3-form

We can of course apply this same expression when α, β, γ are not orthonormal:

$$\begin{aligned} (\alpha \wedge \beta \wedge \gamma)(u, v, w) &:= \alpha(u)\beta(v)\gamma(w) + \alpha(v)\beta(w)\gamma(u) + \alpha(w)\beta(u)\gamma(v) \\ &\quad - \alpha(u)\beta(w)\gamma(v) - \alpha(w)\beta(v)\gamma(u) - \alpha(v)\beta(u)\gamma(w) \end{aligned}$$

Interpretation (in R^3)?

Volume of u, v, w gets scaled by volume of α, β, γ .



k -Form

- More generally, k -form is a *fully antisymmetric, multilinear* measurement of a k -vector.
- Typically think of this as a map from k vectors to a scalar:

$$\alpha : \underbrace{V \times \cdots V}_{k \text{ times}} \rightarrow \mathbb{R}$$

- *Multilinear* means “linear in each argument.” E.g., for a 2-form:

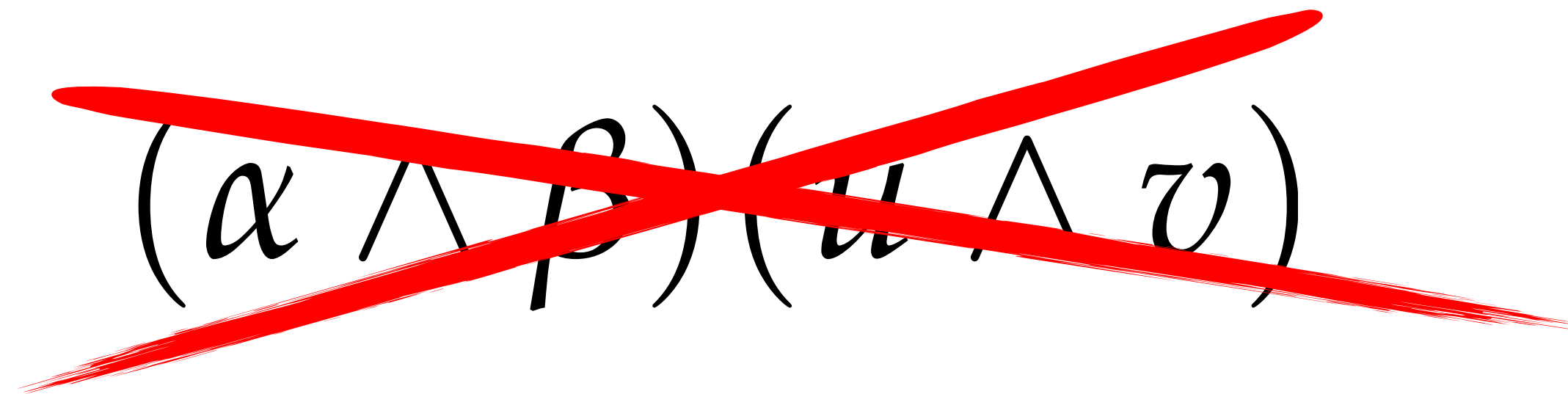
$$\begin{aligned} \alpha(au + bv, w) &= a\alpha(u, w) + b\alpha(v, w) \\ \alpha(u, av + bw) &= a\alpha(u, v) + b\alpha(u, w) \end{aligned} \quad , \quad \forall u, v, w \in V, a, b \in \mathbb{R}$$

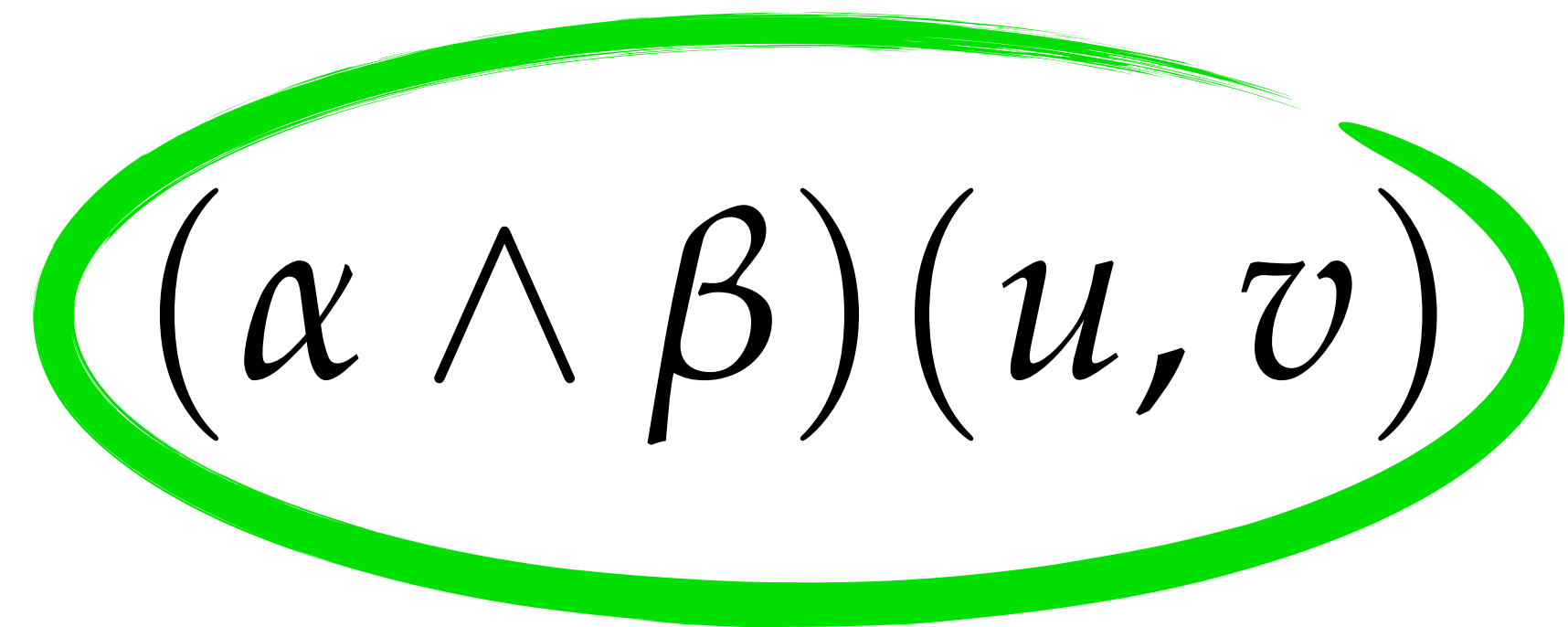
- *Fully antisymmetric* means exchanging two arguments reverses sign. E.g., 3-form:

$$\begin{aligned} \alpha(u, v, w) &= \alpha(v, w, u) = \alpha(w, u, v) = \\ &= -\alpha(u, w, v) = -\alpha(w, v, u) = -\alpha(v, u, w) \end{aligned}$$

A Note on Notation

- A k -form effectively measures a k -vector
- However, *nobody* writes the argument k -vector using a wedge*
- Instead, the convention is to write a list of vectors:


$$(\alpha \wedge \beta)(u \wedge v)$$


$$(\alpha \wedge \beta)(u, v)$$

*I have no idea why.

k-Forms and Determinants

- For 3-forms, saw that we could express a *k*-form via a *determinant*
- Captures the fact that *k*-forms are measurements of *volume*
- How does this work more generally?
 - **Conceptually:** “project” onto *k*-dimensional space and take determinant there

$$(\alpha_1 \wedge \cdots \wedge \alpha_k)(u_1, \dots, u_k) := \begin{vmatrix} \alpha_1(u_1) & \cdots & \alpha_1(u_k) \\ \vdots & \ddots & \vdots \\ \alpha_k(u_1) & \cdots & \alpha_k(u_k) \end{vmatrix}$$

k=1:

$$\det \left(\begin{bmatrix} \alpha_1(u_1) \end{bmatrix} \right) = \alpha_1(u_1)$$

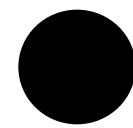
(Determinant of a 1x1 matrix is just the entry of that matrix!)

k=2:

$$\begin{aligned} & \det \left(\begin{bmatrix} \alpha_1(u_1) & \alpha_1(u_2) \\ \alpha_2(u_1) & \alpha_2(u_2) \end{bmatrix} \right) \\ &= \alpha_1(u_1)\alpha_2(u_2) - \alpha_1(u_2)\alpha_2(u_1) \end{aligned}$$

0-Forms

- What's a 0-form?
 - In general, a k -form takes k vectors and produces a scalar
 - So a 0-form must take 0 vectors and produce a scalar
 - I.e., *a 0-form is a scalar!*
- Basically looks like this:



Note: still has *magnitude*, even though it has only one possible “direction.”



k -Forms in Coordinates

Dual Basis

In an n -dimensional vector space V , can express vectors v in a basis e_1, \dots, e_n :

$$v = v^1 e_1 + \dots + v^n e_n$$

The scalar values v^i are the *coordinates* of v .

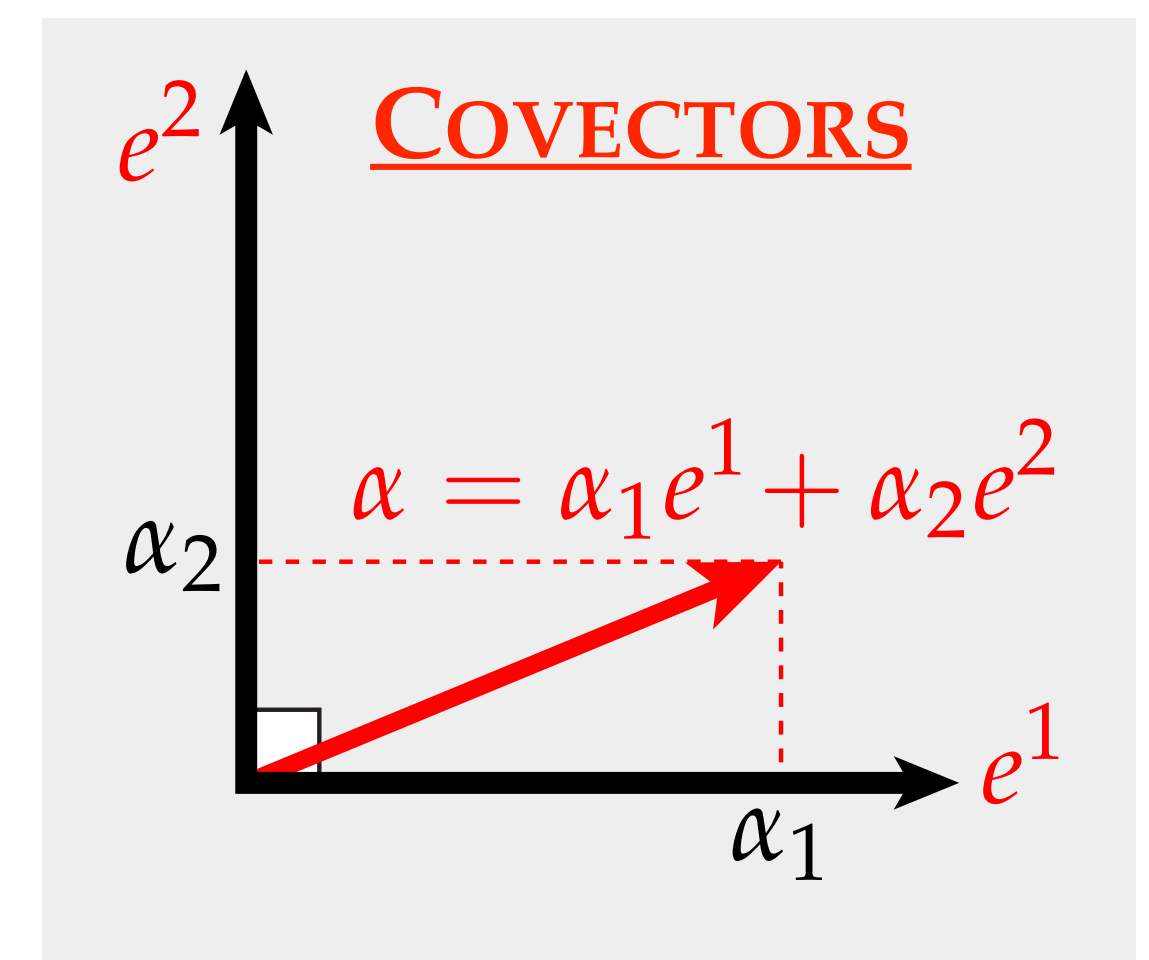
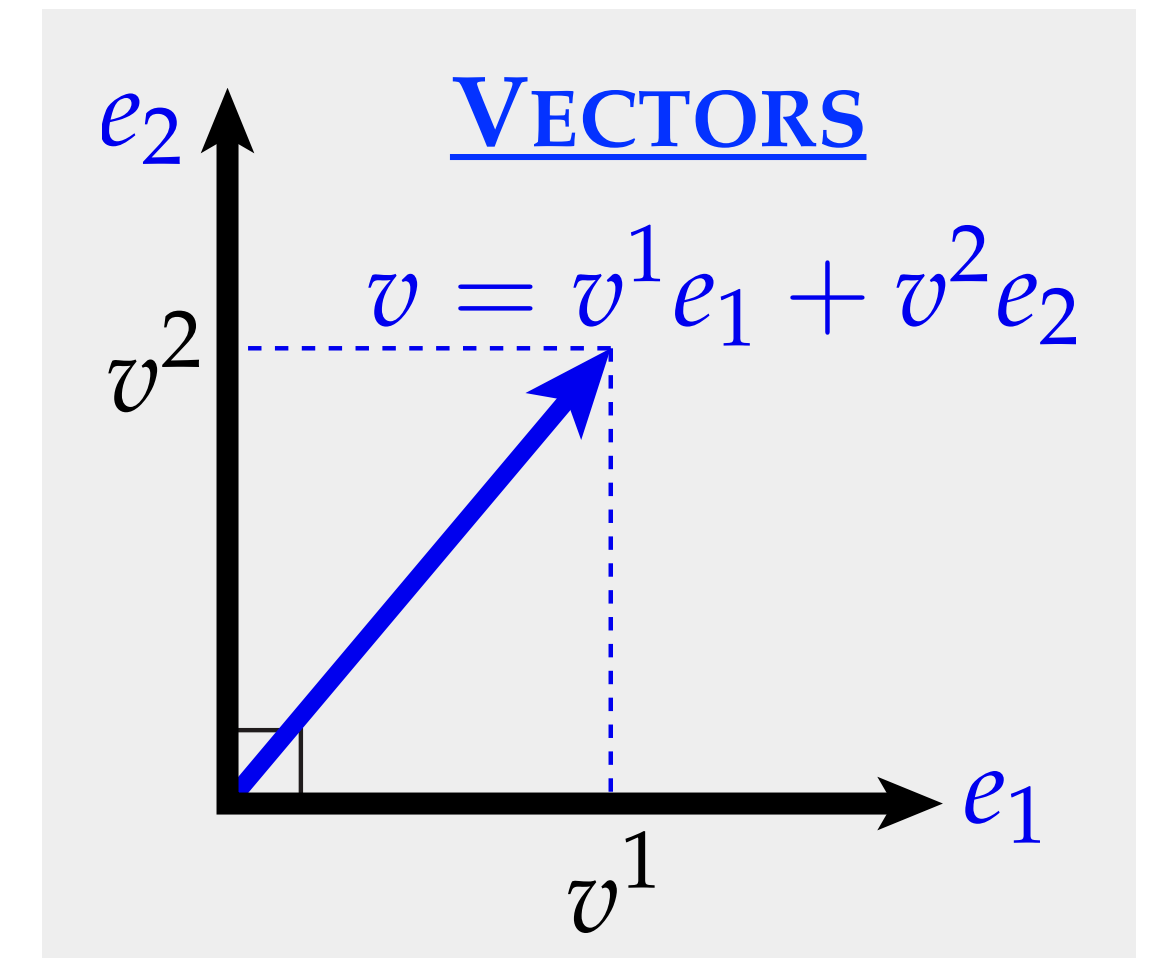
We can also write *covectors* α in a so-called *dual basis* e^1, \dots, e^n :

$$\alpha = \alpha_1 e^1 + \dots + \alpha_n e^n$$

These bases have a special relationship, namely:

$$e^i(e_j) = \begin{cases} 1, & i = j \\ 0, & \text{otherwise} \end{cases}$$

(**Q:** What does e^i mean, geometrically?)



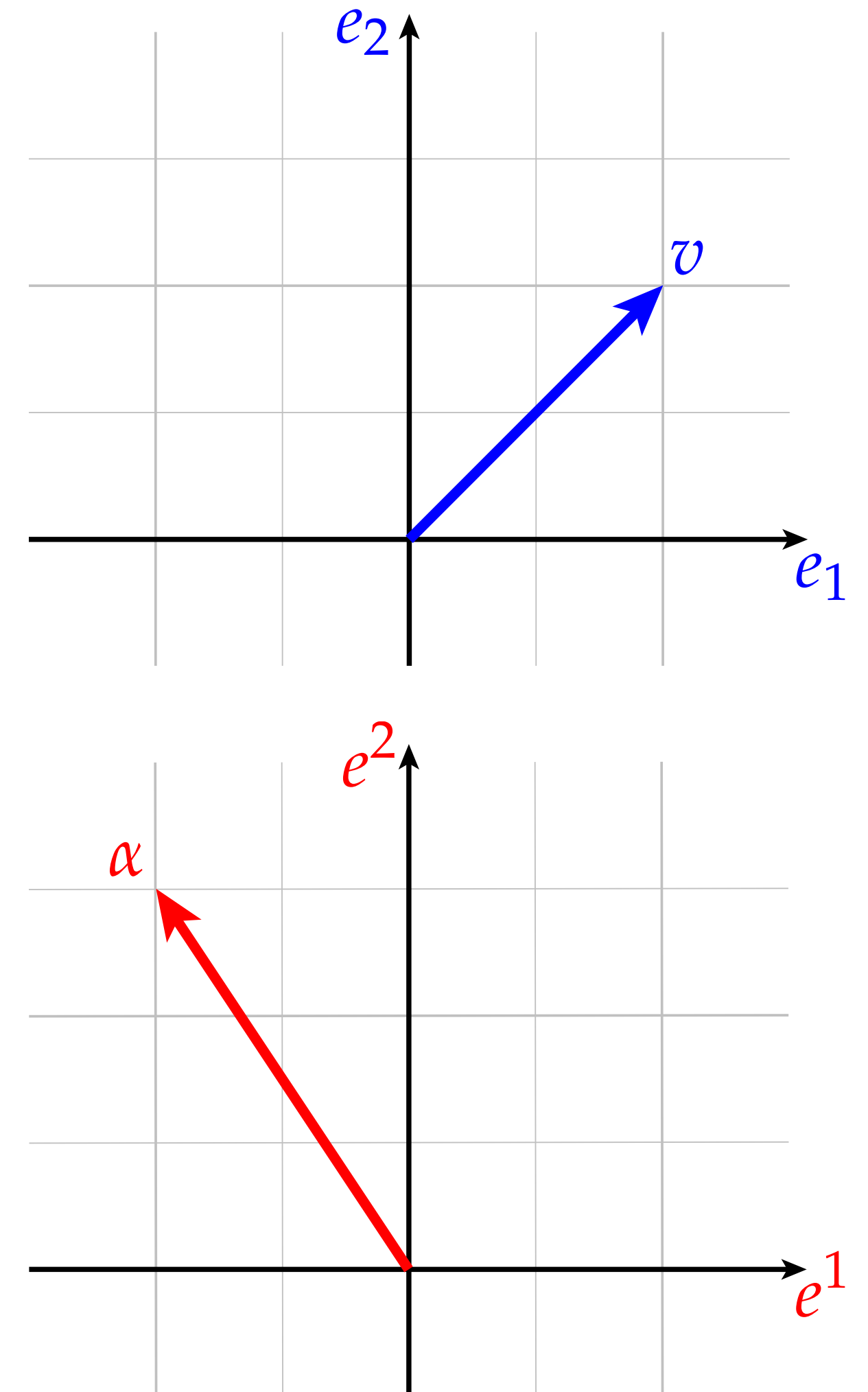
1-form — Example in Coordinates

- Some simple calculations in coordinates help to solidify understanding of k -forms.
- Let's start with a vector v and a 1-form α in the plane:

$$v = 2e_1 + 2e_2$$

$$\alpha = -2e^1 + 3e^2$$

$$\begin{aligned}\alpha(v) &= (-2e^1 + 3e^2)(2e_1 + 2e_2) \\ &= -2e^1(2e_1 + 2e_2) + 3e^2(2e_1 + 2e_2) \\ &= \cancel{-4e^1(e_1)}^1 \cancel{-4e^1(e_2)}^0 + \cancel{6e^2(e_1)}^0 \cancel{6e^2(e_2)}^1 \\ &= -4 + 6 \quad (\text{Just like a dot product!}) \\ &= 2.\end{aligned}$$



2-form—Example in Coordinates

Consider the following vectors and covectors:

$$\begin{aligned} u &= 2e_1 + 2e_2 & \alpha &= e^1 + 3e^2 \\ v &= -2e_1 + 2e_2 & \beta &= 2e^1 + e^2 \end{aligned}$$

We then have:

$$(\alpha \wedge \beta)(u, v) = \alpha(u)\beta(v) - \alpha(v)\beta(u)$$

$$\alpha(u) = 1 \cdot 2 + 3 \cdot 2 = 8$$

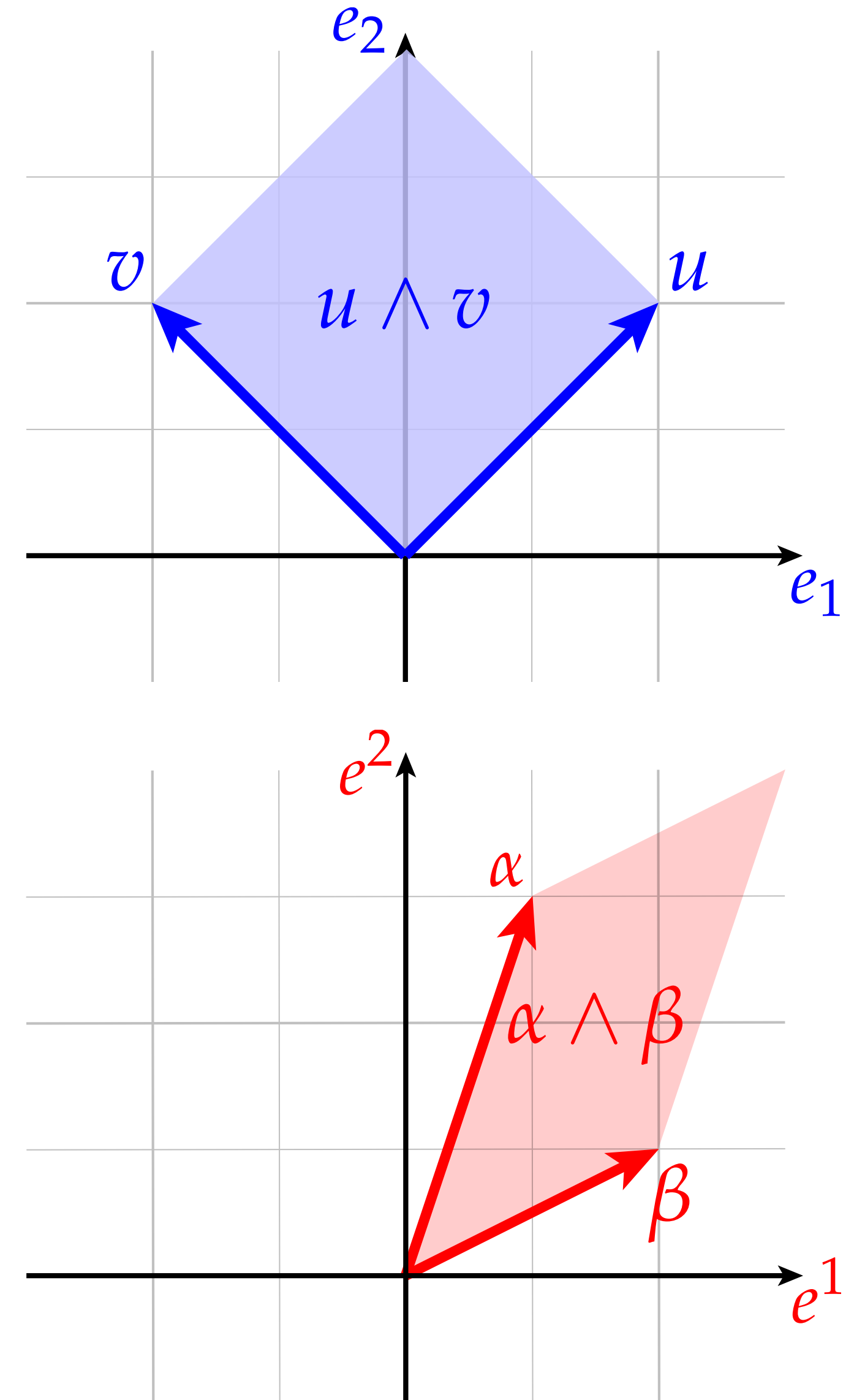
$$\beta(v) = \dots = -2$$

$$\alpha(v) = \dots = 4$$

$$\beta(u) = \dots = 6$$

$$\Rightarrow (\alpha \wedge \beta)(u, v) = 8 \cdot (-2) - 4 \cdot 6 = -40.$$

Q: What does this value mean, geometrically? Why is it *negative*?



Einstein Summation Notation

Why are some indices “up” and others “down”?

Bemerkung zur Vereinfachung der Schreibweise der Ausdrücke.

Ein Blick auf die Gleichungen dieses Paragraphen zeigt, daß über Indizes, die zweimal unter einem Summenzeichen auftreten [z. B. der Index ν in (5)], stets summiert wird, und zwar *nur* über zweimal auftretende Indizes. Es ist deshalb möglich, ohne die Klarheit zu beeinträchtigen, die Summenzeichen wegzulassen. Dafür führen wir die Vorschrift ein: Tritt ein Index in einem Term eines Ausdruckes zweimal auf, so ist über ihn stets zu summieren, wenn nicht ausdrücklich das Gegenteil bemerkt ist.

Neuere Berechnung des Elementartensors

$$\frac{1}{2} \left(\frac{\partial^2 g_{im}}{\partial x_k \partial x_l} + \frac{\partial^2 g_{kl}}{\partial x_i \partial x_m} - \frac{\partial^2 g_{il}}{\partial x_k \partial x_m} - \frac{\partial^2 g_{km}}{\partial x_i \partial x_l} \right) - \frac{1}{4} g_{ic} \left(\frac{\partial g_{ie}}{\partial x_l} + \frac{\partial g_{le}}{\partial x_i} - \frac{\partial g_{il}}{\partial x_e} \right) \left(\frac{\partial g_{ec}}{\partial x_m} + \frac{\partial g_{mc}}{\partial x_e} - \frac{\partial g_{me}}{\partial x_c} \right) \Big|_{x_{kl}}$$

bleibt stehen.

$$\frac{1}{2} g_{kl} \frac{\partial^2 g_{im}}{\partial x_k \partial x_l} \text{ bleibt stehen.}$$

$$g_{kl} \left[\frac{\partial^2 g_{im}}{\partial x_k \partial x_l} \right] = g_{kl} \left(2 \frac{\partial g_{il}}{\partial x_k} - \frac{\partial g_{kl}}{\partial x_i} \right) = 0 \quad \left| \frac{\partial}{\partial x_m} \right.$$

$$g_{kl} \left[\frac{\partial^2 g_{im}}{\partial x_k \partial x_l} \right] = g_{kl} \left(2 \frac{\partial g_{mk}}{\partial x_l} - \frac{\partial g_{kl}}{\partial x_m} \right) = 0 \quad \left| \frac{\partial}{\partial x_i} \right.$$

$$2 g_{kl} \left(\frac{\partial^2 g_{il}}{\partial x_k \partial x_m} + \frac{\partial^2 g_{mk}}{\partial x_i \partial x_l} - \frac{\partial^2 g_{kl}}{\partial x_i \partial x_m} \right) + \frac{\partial g_{kl}}{\partial x_m} \left(2 \frac{\partial g_{il}}{\partial x_k} - \frac{\partial g_{kl}}{\partial x_i} \right) + \frac{\partial g_{kl}}{\partial x_i} \left(2 \frac{\partial g_{mk}}{\partial x_l} - \frac{\partial g_{kl}}{\partial x_m} \right)$$

$$- \frac{1}{2} g_{kl} \left(\frac{\partial^2 g_{il}}{\partial x_k \partial x_m} + \frac{\partial^2 g_{mk}}{\partial x_i \partial x_l} - \frac{\partial^2 g_{kl}}{\partial x_i \partial x_m} \right) = \frac{1}{4} \left[\frac{\partial g_{kl}}{\partial x_m} \left(2 \frac{\partial g_{il}}{\partial x_k} - \frac{\partial g_{kl}}{\partial x_i} \right) + \frac{\partial g_{kl}}{\partial x_i} \left(2 \frac{\partial g_{mk}}{\partial x_l} - \frac{\partial g_{kl}}{\partial x_m} \right) \right]$$

zweites Glied:

$$- \frac{1}{4} g_{ic} \frac{\partial g_{ie}}{\partial x_l} \frac{\partial g_{ec}}{\partial x_m} g_{kl} \quad \begin{matrix} + \frac{1}{4} \frac{\partial g_{ic}}{\partial x_i} \frac{\partial g_{ec}}{\partial x_m} g_{kl} \\ + \frac{1}{4} \frac{\partial g_{ic}}{\partial x_l} \frac{\partial g_{ec}}{\partial x_m} g_{kl} \end{matrix}$$

$$- \frac{1}{4} g_{ic} \left(\frac{\partial g_{ie}}{\partial x_l} - \frac{\partial g_{il}}{\partial x_e} \right) \left(\frac{\partial g_{ec}}{\partial x_m} - \frac{\partial g_{mc}}{\partial x_e} \right) g_{kl}$$

$$= - \frac{1}{2} g_{ic} g_{kl} \frac{\partial g_{ie}}{\partial x_l} \frac{\partial g_{mc}}{\partial x_m} + \frac{1}{2} g_{ic} g_{kl} \frac{\partial g_{il}}{\partial x_e} \frac{\partial g_{mc}}{\partial x_m}$$

Die mit 2 multiplizierte Elementartensor erhält also die Form

$$g_{kl} \frac{\partial^2 g_{im}}{\partial x_k \partial x_l} - \frac{1}{2} \frac{\partial g_{kl}}{\partial x_m} \frac{\partial g_{il}}{\partial x_k} + \frac{\partial g_{kl}}{\partial x_m} \frac{\partial g_{il}}{\partial x_k} + \frac{\partial g_{kl}}{\partial x_i} \frac{\partial g_{mk}}{\partial x_l}$$

$$- g_{ic} g_{kl} \frac{\partial g_{ie}}{\partial x_l} \frac{\partial g_{mc}}{\partial x_m} + g_{ic} g_{kl} \frac{\partial g_{il}}{\partial x_e} \frac{\partial g_{mc}}{\partial x_m}$$

Resultat sicher. Gilt für Koordinaten, die der Gl. $\Delta \varphi = 0$ genügen.

— Einstein, “Die Grundlage der allgemeinen Relativitätstheorie” (1916)

Einstein Summation Notation

Key idea: sum over repeated indices.

$$x^i y_i := \sum_{i=1}^n x^i y_i$$

NOTE ON A SIMPLIFIED WAY OF WRITING EXPRESSIONS

A look at the equations of this paragraph show that there is always a summation over indices which occur twice, and only for twice-repeated indices. It is therefore possible, without detracting from clarity, to omit the sum sign. For this we introduce a rule: if an index in an expression appears twice, then a sum is implicitly taken over this index, unless specifically noted to the contrary.

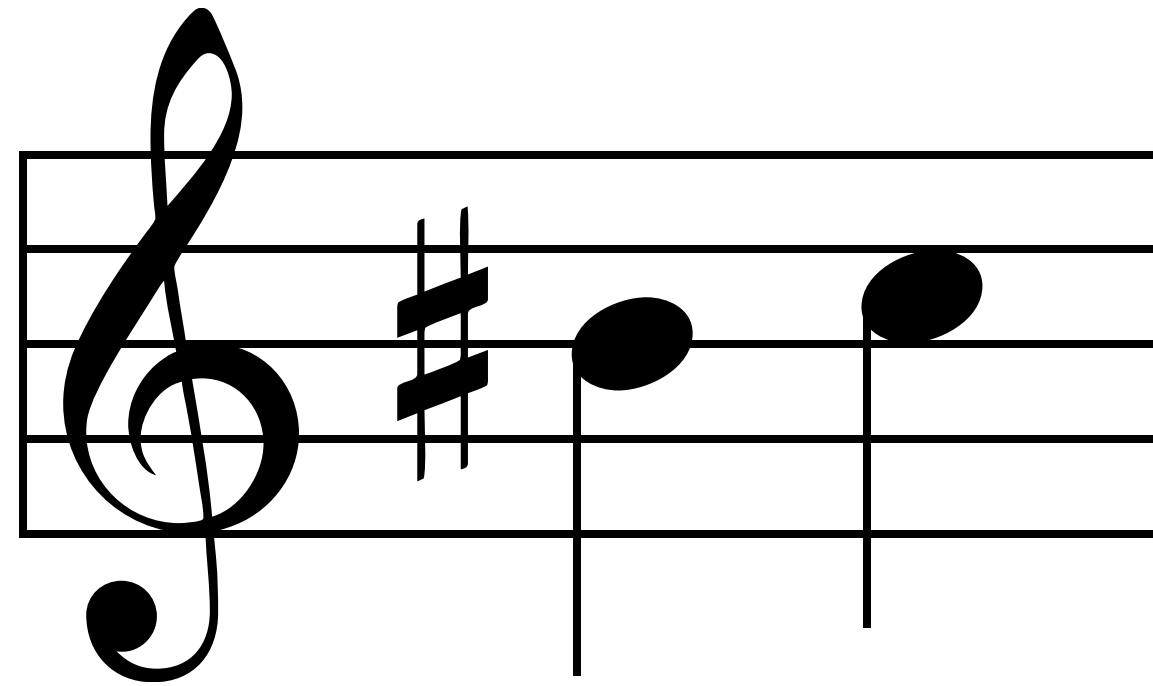
Handwritten manuscript page showing the derivation of the Ricci tensor from the Riemann tensor using Einstein summation notation. The text is in German and includes the following steps:

- Top line: $\frac{1}{2} \left(\frac{\partial^2 g_{im}}{\partial x_k \partial x_l} + \frac{\partial^2 g_{kl}}{\partial x_i \partial x_m} - \frac{\partial^2 g_{il}}{\partial x_k \partial x_m} - \frac{\partial^2 g_{km}}{\partial x_i \partial x_l} \right) - \frac{1}{4} g_{ic} \left(\frac{\partial g_{ie}}{\partial x_k} + \frac{\partial g_{ek}}{\partial x_i} - \frac{\partial g_{il}}{\partial x_e} \right) \left(\frac{\partial g_{me}}{\partial x_k} + \frac{\partial g_{mk}}{\partial x_m} - \frac{\partial g_{mk}}{\partial x_e} \right) g_{kl}$
- Second line: $\frac{1}{2} g_{kl} \frac{\partial^2 g_{im}}{\partial x_k \partial x_l}$ bleibt stehen.
- Third line: $g_{kl} \left[\begin{smallmatrix} k & l \\ i & i \end{smallmatrix} \right] = g_{kl} \left(2 \frac{\partial g_{il}}{\partial x_k} - \frac{\partial g_{kl}}{\partial x_i} \right) = 0 \quad \left| \frac{\partial}{\partial x_m} \right.$
- Fourth line: $g_{kl} \left[\begin{smallmatrix} k & l \\ m & m \end{smallmatrix} \right] = g_{kl} \left(2 \frac{\partial g_{mk}}{\partial x_l} - \frac{\partial g_{kl}}{\partial x_m} \right) = 0 \quad \left| \frac{\partial}{\partial x_i} \right.$
- Fifth line: $2 g_{kl} \left(\frac{\partial^2 g_{il}}{\partial x_k \partial x_m} + \frac{\partial^2 g_{mk}}{\partial x_i \partial x_l} - \frac{\partial^2 g_{kl}}{\partial x_i \partial x_m} \right) + \frac{\partial g_{kl}}{\partial x_m} \left(2 \frac{\partial g_{il}}{\partial x_k} - \frac{\partial g_{kl}}{\partial x_i} \right) + \frac{\partial g_{kl}}{\partial x_i} \left(2 \frac{\partial g_{mk}}{\partial x_l} - \frac{\partial g_{kl}}{\partial x_m} \right)$
- Sixth line: $-\frac{1}{2} g_{kl} \left(\frac{\partial g_{ie}}{\partial x_k} - \frac{\partial g_{il}}{\partial x_e} \right) \left(\frac{\partial g_{me}}{\partial x_k} - \frac{\partial g_{mk}}{\partial x_e} \right) g_{kl}$
- Seventh line: $-\frac{1}{4} g_{ic} \frac{\partial g_{ie}}{\partial x_k} \frac{\partial g_{kc}}{\partial x_m} g_{kl} - \frac{1}{4} g_{ic} \frac{\partial g_{ie}}{\partial x_l} \frac{\partial g_{lc}}{\partial x_m} g_{kl}$
- Eighth line: $-\frac{1}{4} g_{ic} \left(\frac{\partial g_{ie}}{\partial x_l} - \frac{\partial g_{il}}{\partial x_e} \right) \left(\frac{\partial g_{me}}{\partial x_k} - \frac{\partial g_{mk}}{\partial x_e} \right) g_{kl}$
- Ninth line: $= -\frac{1}{2} g_{ic} g_{kl} \frac{\partial g_{ie}}{\partial x_l} \frac{\partial g_{me}}{\partial x_k} + \frac{1}{2} g_{ic} g_{kl} \frac{\partial g_{il}}{\partial x_k} \frac{\partial g_{me}}{\partial x_l}$
- Tenth line: $\frac{\partial g_{im}}{\partial x_k \partial x_l} - \frac{1}{2} \frac{\partial g_{kl}}{\partial x_m \partial x_i} + \frac{\partial g_{kl}}{\partial x_m \partial x_k} + \frac{\partial g_{kl}}{\partial x_i \partial x_l} - \frac{\partial g_{kl}}{\partial x_i \partial x_m}$
- Eleventh line: $-g_{ic} g_{kl} \frac{\partial g_{ie}}{\partial x_l} \frac{\partial g_{me}}{\partial x_k} + g_{ic} g_{kl} \frac{\partial g_{il}}{\partial x_k} \frac{\partial g_{me}}{\partial x_l}$
- Bottom line: $\text{Resultat sicher. Gilt für Koordinaten, die der Gl. } \Delta \varphi = 0 \text{ genügen.}$

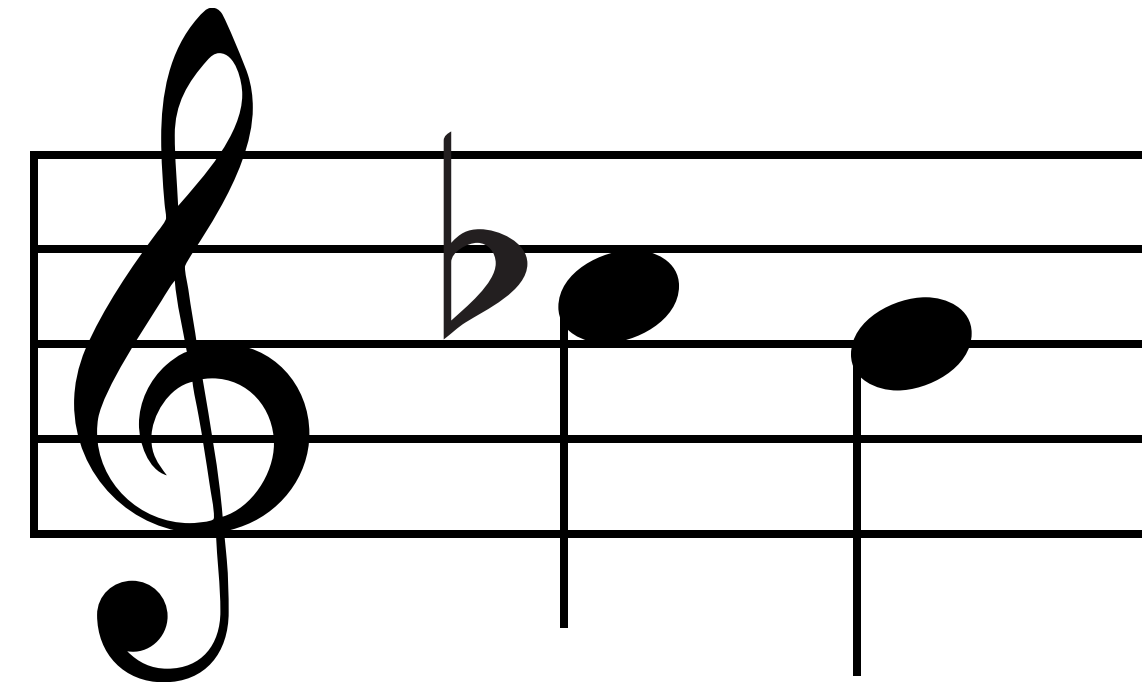
— Einstein, “Die Grundlage der allgemeinen Relativitätstheorie” (1916)

Sharp and Flat in Coordinates

Q: What do sharp and flat do on a musical staff?



(raise pitch)



(lower pitch)

Likewise, sharp and flat *raise* and *lower* indices of coefficients for 1-forms/vectors.

Suppose for instance that $\alpha(v) = \langle u, v \rangle$ for all $v \in V$. Then

$$\alpha = \alpha_1 e^1 + \cdots + \alpha_n e^n \quad \begin{array}{c} \xrightarrow{\#} \\ \xleftarrow{b} \end{array} \quad u = u^1 e_1 + \cdots + u^n e_n$$

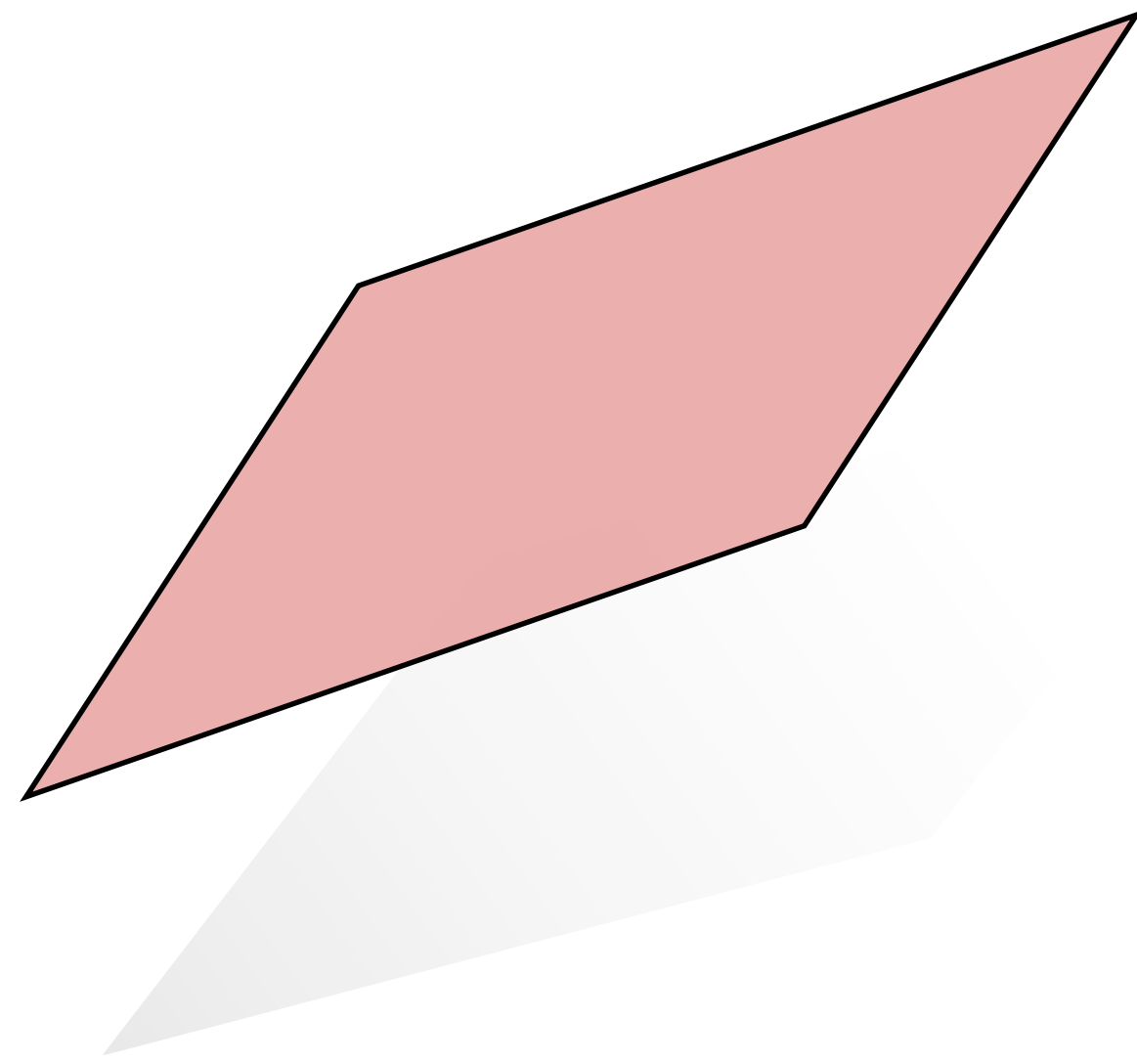
(Sometimes called the *musical isomorphisms*.)



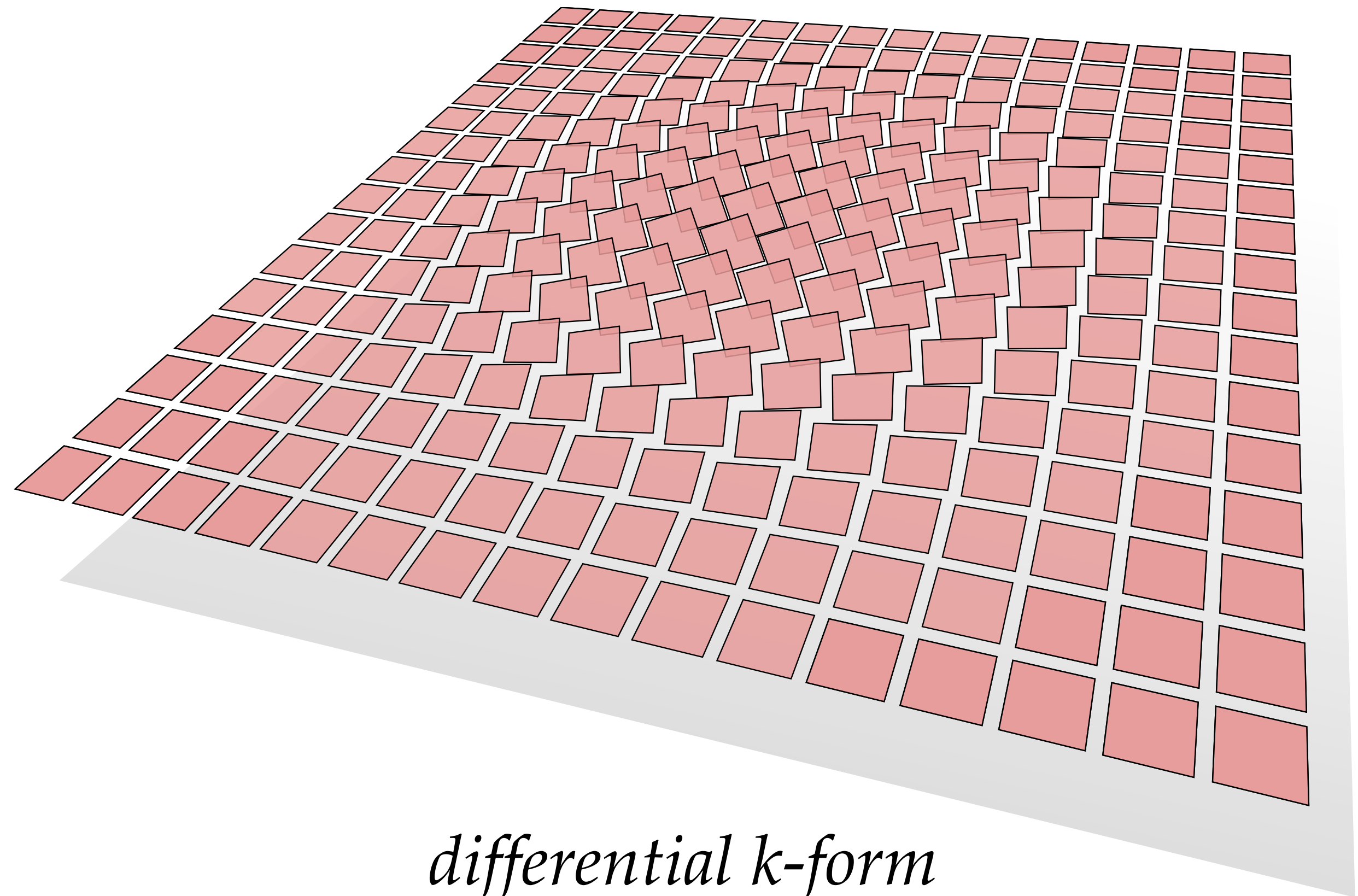
Differential k -Forms

Differential Form

- A *differential k -form* is an assignment of a k -form to each point*:



k-form

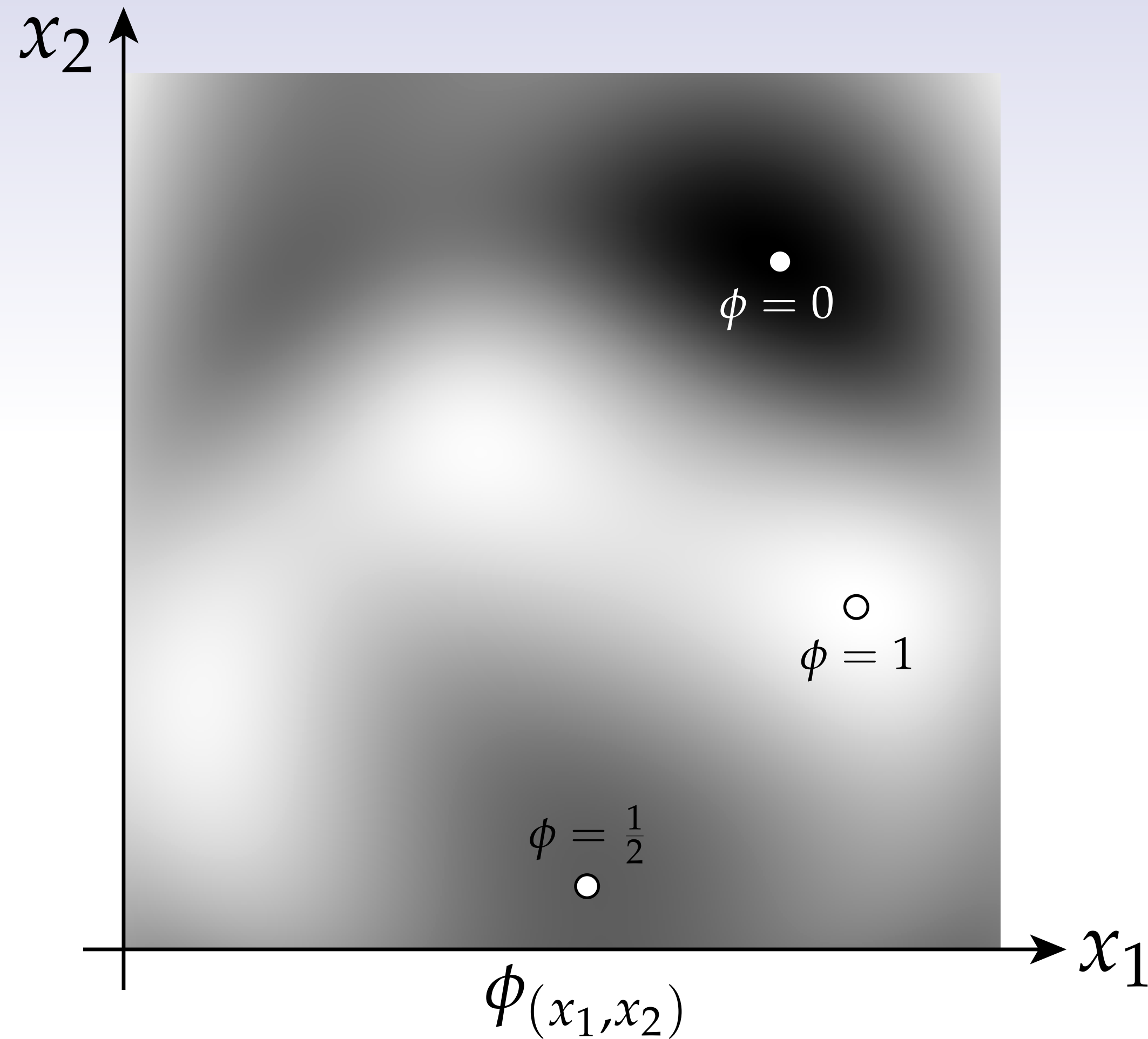


differential k-form

*Common (and confusing!) to abbreviate “differential k -form” as just “ k -form”!

Differential 0-Form

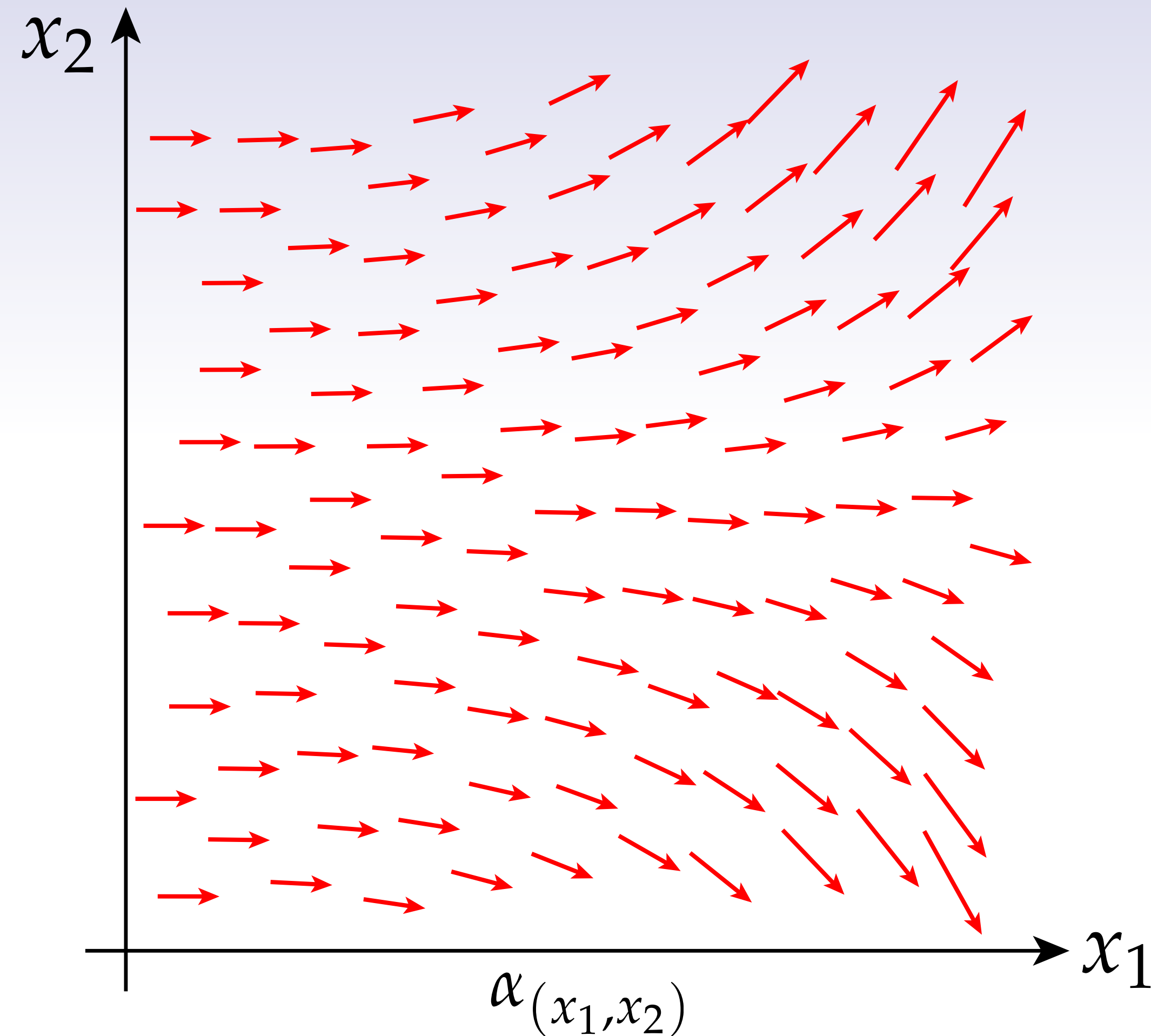
Assigns a scalar to each point. E.g., in 2D we have a value at each point (x_1, x_2) :



Note: exactly the same thing as a *scalar function*!

Differential 1-Form

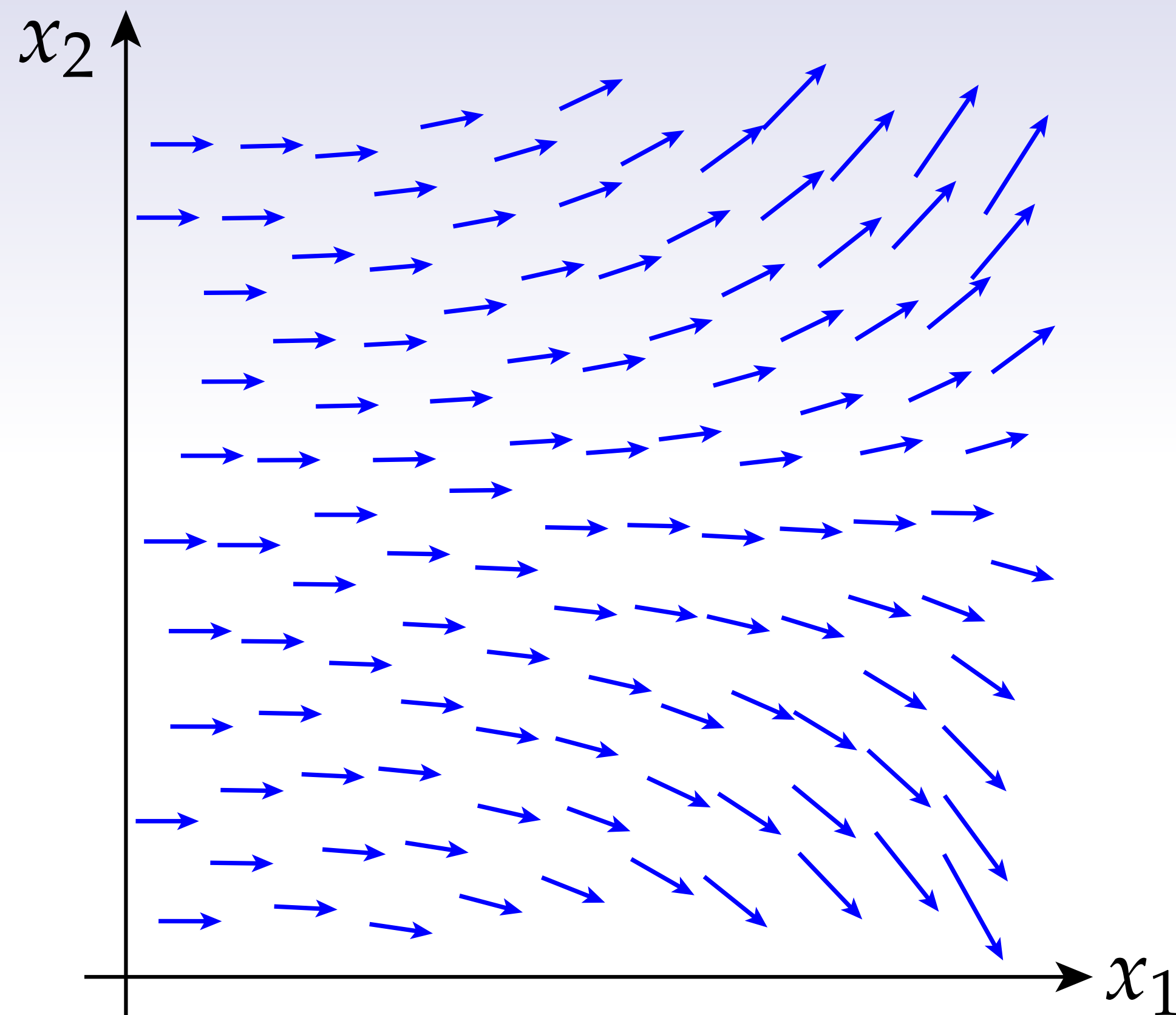
Assigns a 1-form each point. E.g., in 2D we have a 1-form at each point (x_1, x_2) :



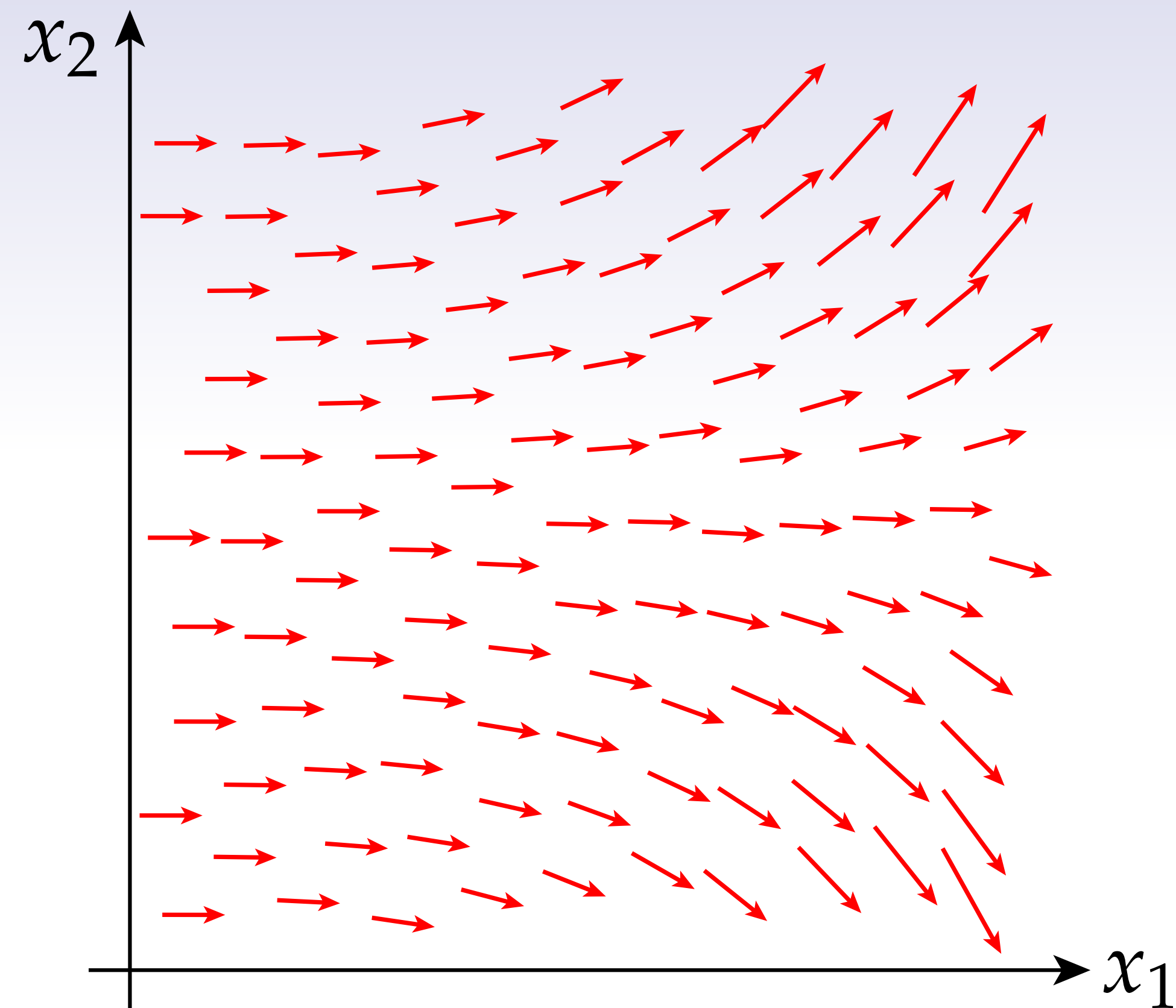
Note: NOT the same thing as a vector field!

Vector Field vs. Differential 1-Form

Superficially, vector fields and differential 1-forms look the same (in R^n):



vector field

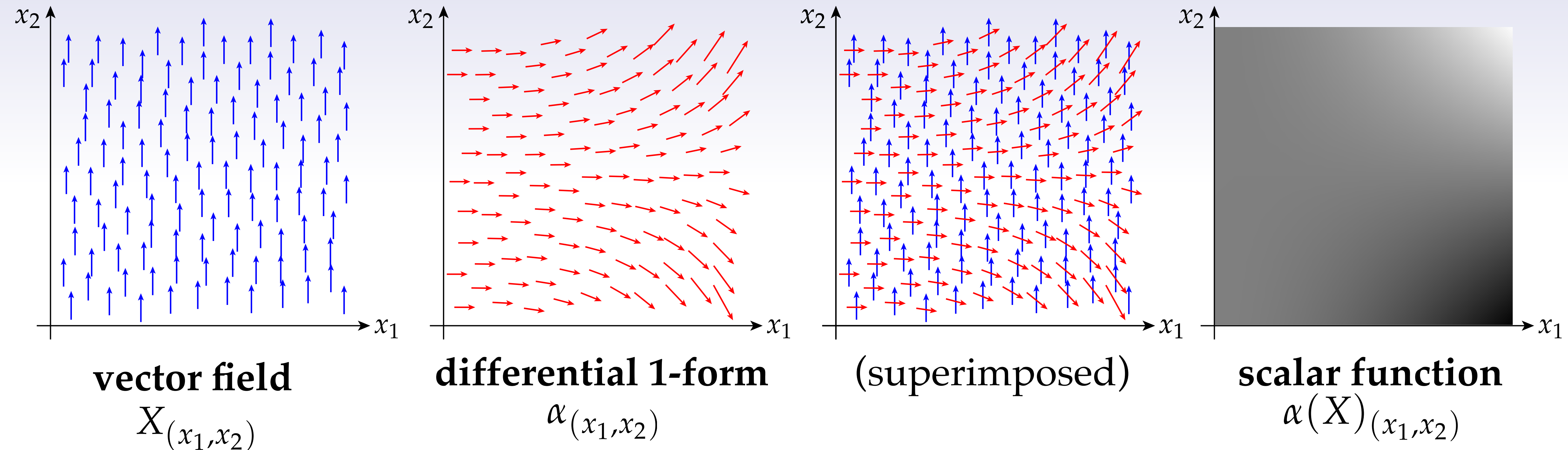


1-form

But recall that a 1-form is a *linear function* from a vector to a scalar (here, at each point.)

Applying a Differential 1-Form to a Vector Field

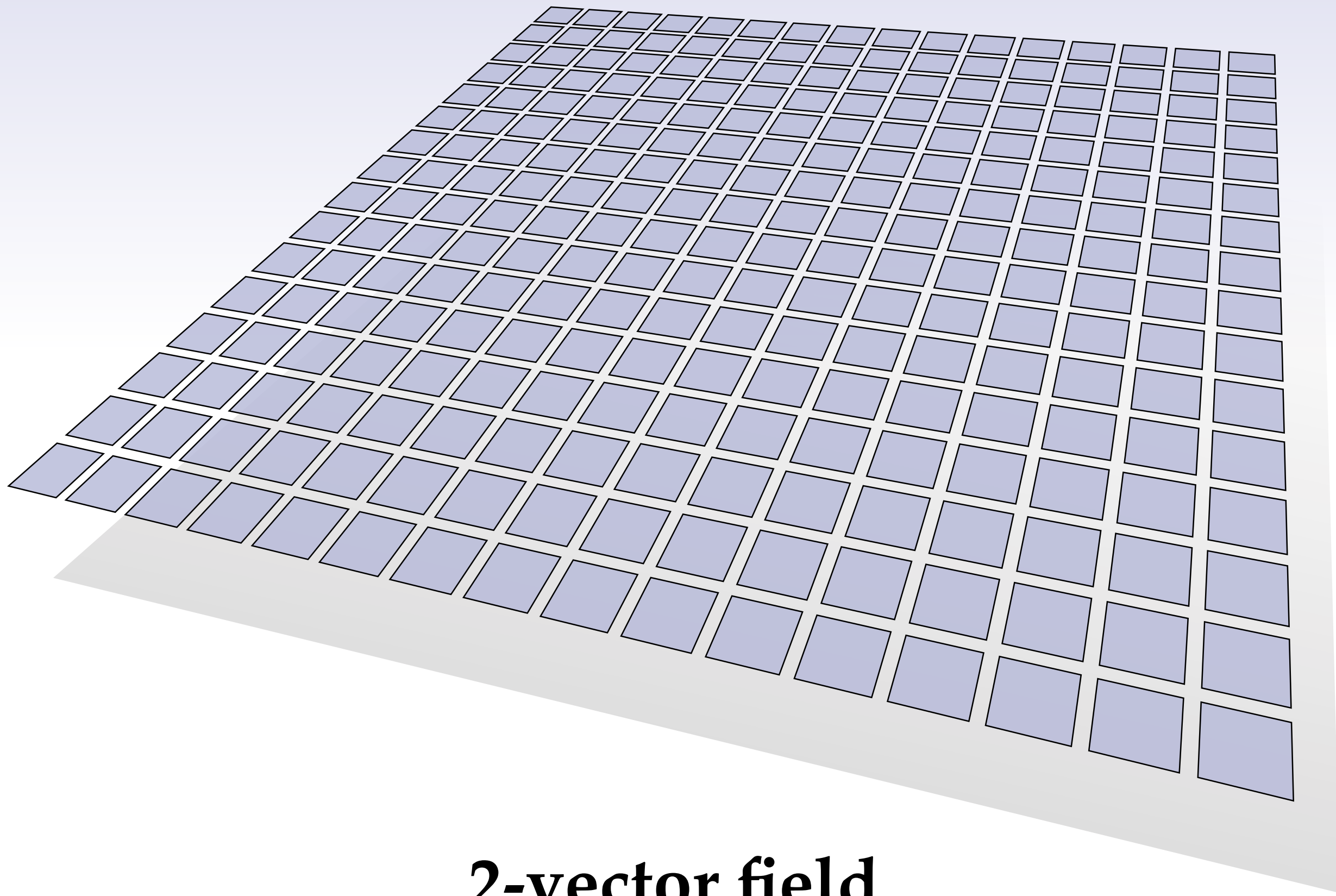
At each point (x_1, x_2) , we can therefore use a 1-form to measure a vector field:



Intuition: resulting function indicates “how strong” X is along α .

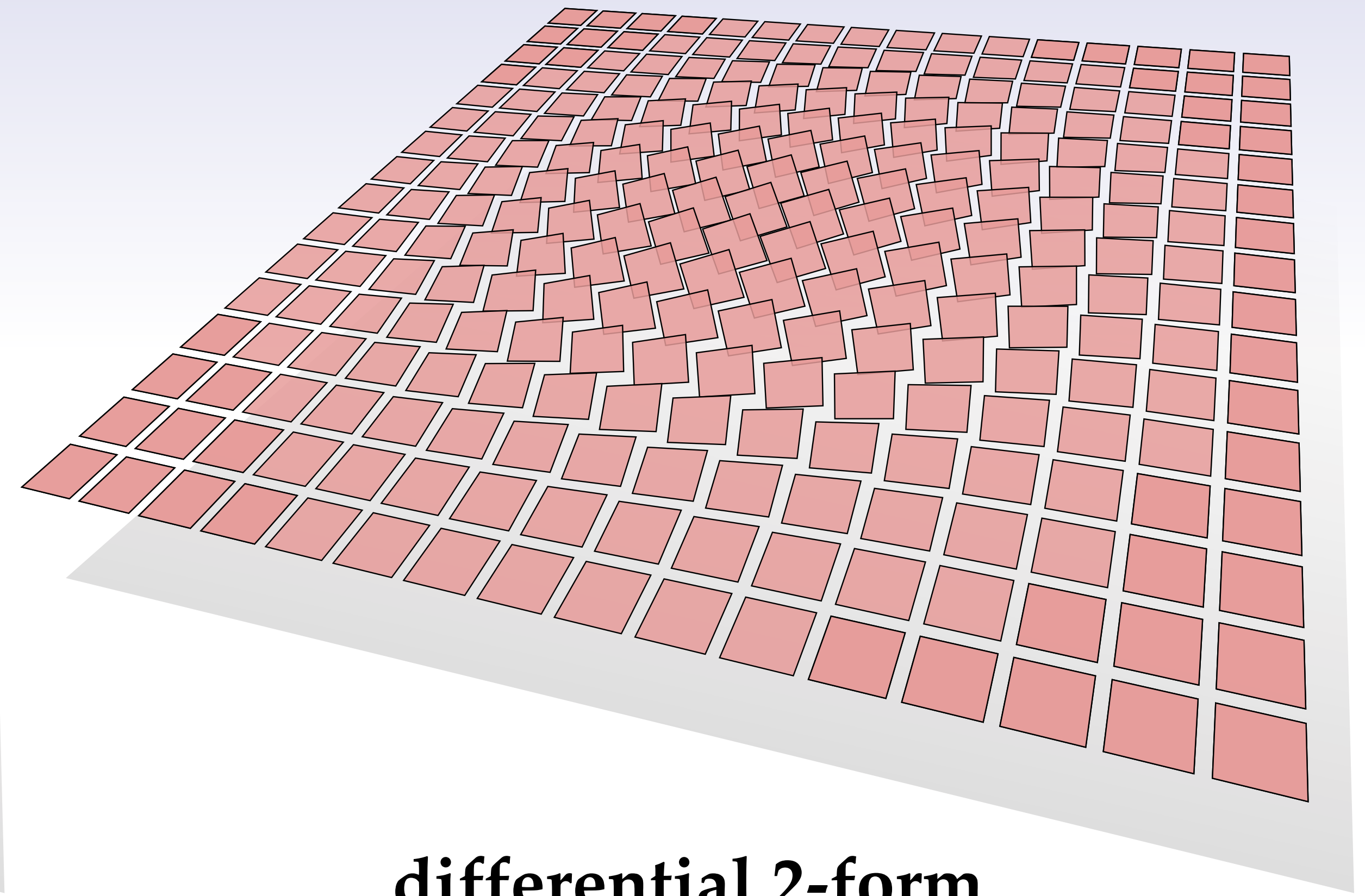
Differential 2-Forms

Likewise, a differential 2-form is an area measurement at each point (x_1, x_2, x_3) :



2-vector field

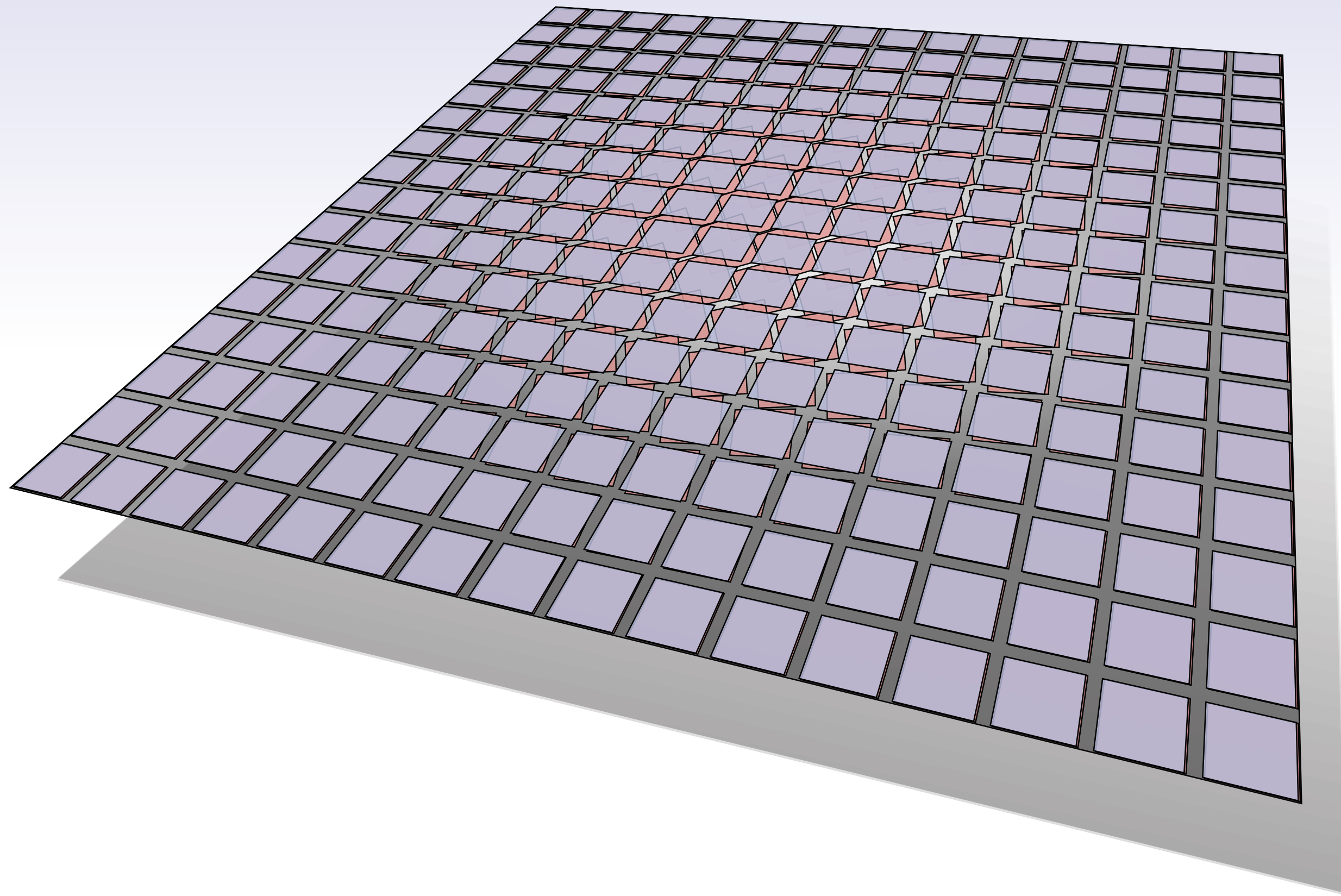
$$(X \wedge Y)_{(x_1, x_2, x_3)}$$



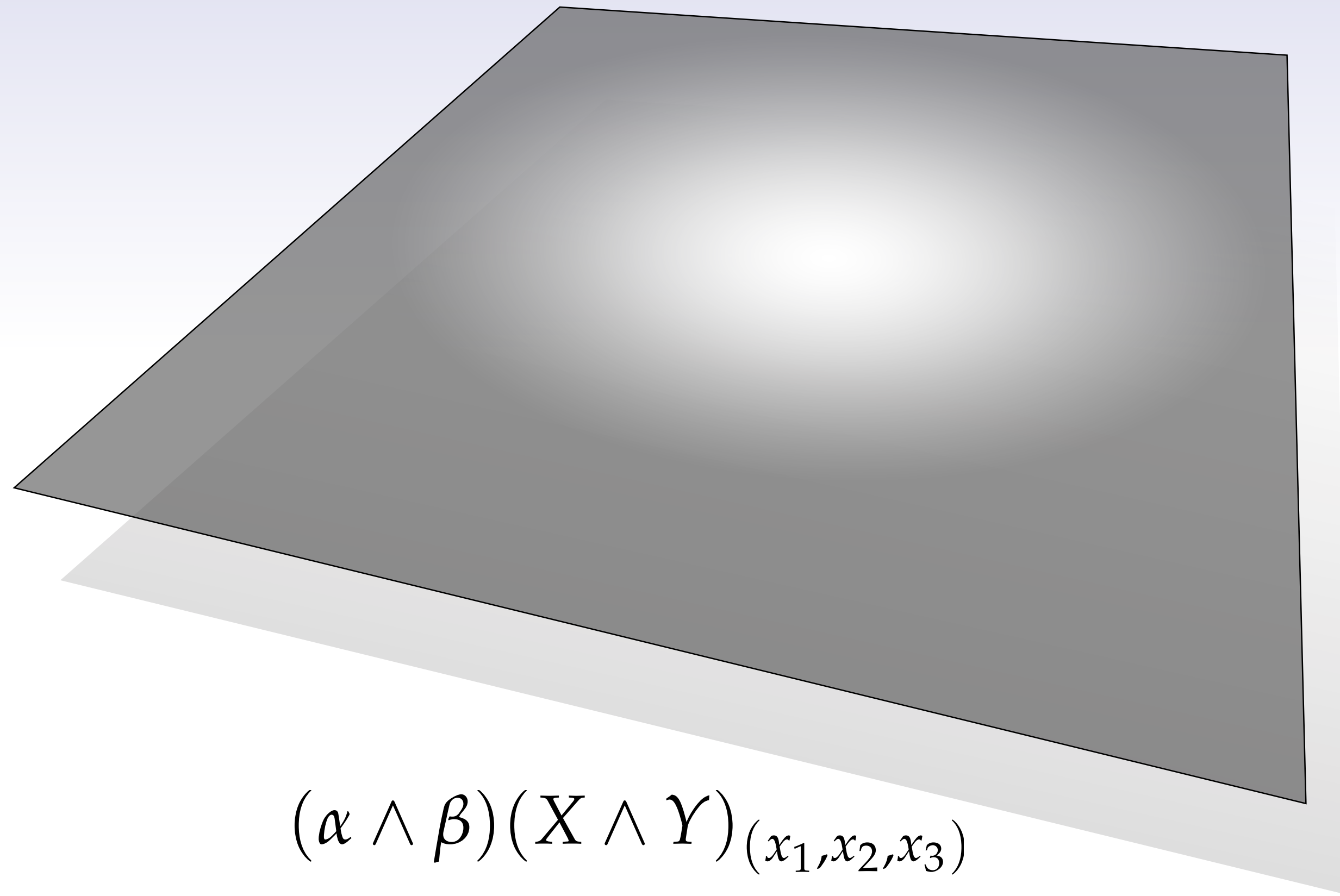
differential 2-form

$$(\alpha \wedge \beta)_{(x_1, x_2, x_3)}$$

Differential 2-Forms



Differential 2-Forms



$$(\alpha \wedge \beta)(X \wedge Y)_{(x_1, x_2, x_3)}$$

Resulting function says how much a 2-vector field “lines up” with a given 2-form.

Pointwise Operations on Differential k -Forms

- Most operations on differential k -forms simply apply that operation at each point.
- E.g., consider two differential forms α, β on R^n . At each point $p := (x_1, \dots, x_n)$,

$$(\star\alpha)_p := \star(\alpha_p)$$

$$(\alpha \wedge \beta)_p := (\alpha_p) \wedge (\beta_p)$$

- In other words, to get the Hodge star of the *differential k -form*, we just apply the Hodge star to the individual k forms at each point p ; to take the wedge of two differential k -forms we just wedge their values at each point.
- Likewise, if X_1, \dots, X_k are vector fields on all of R^n , then

$$\alpha(X_1, \dots, X_k)_p := (\alpha_p)((X_1)_p, \dots, (X_k)_p)$$

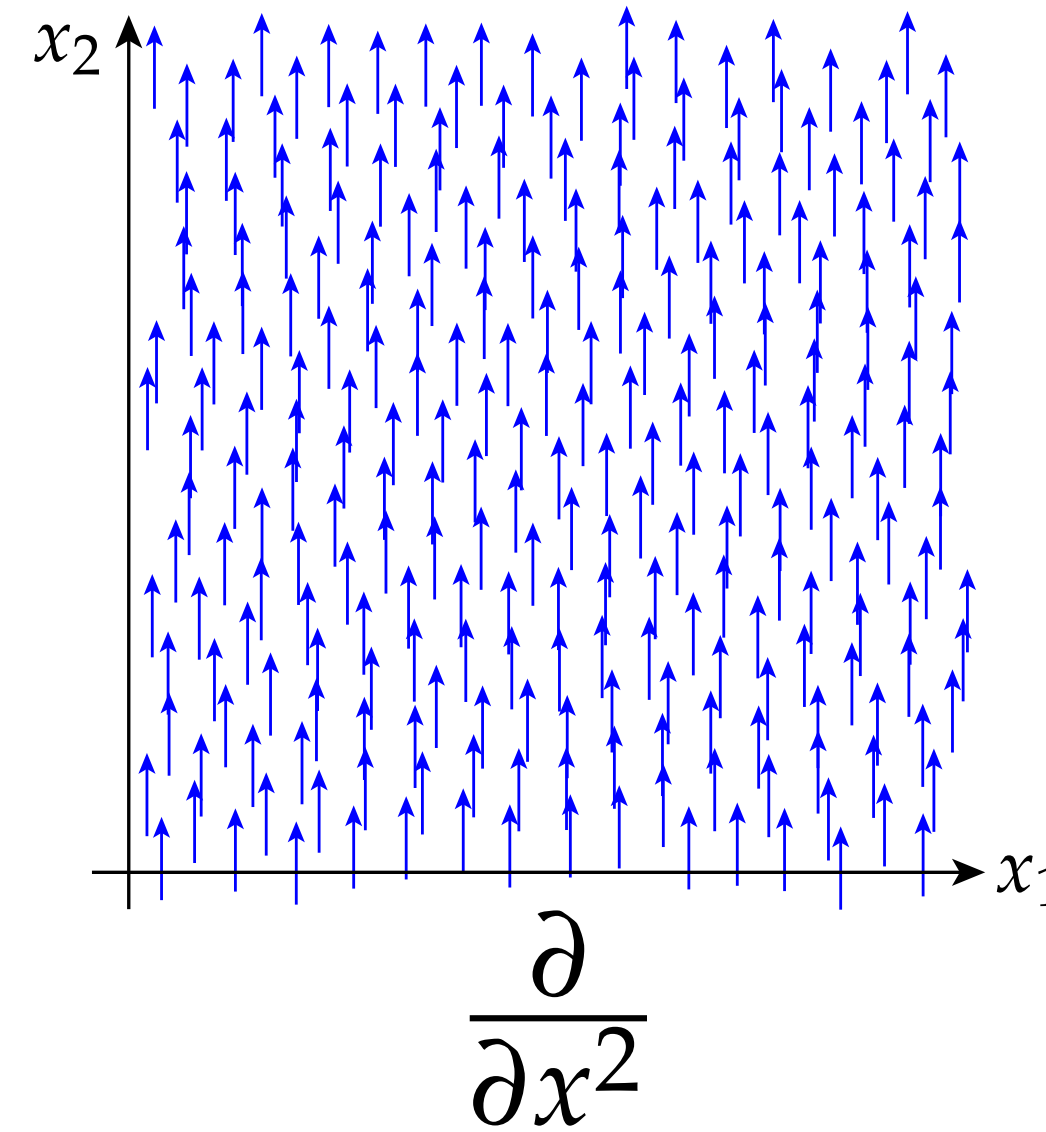
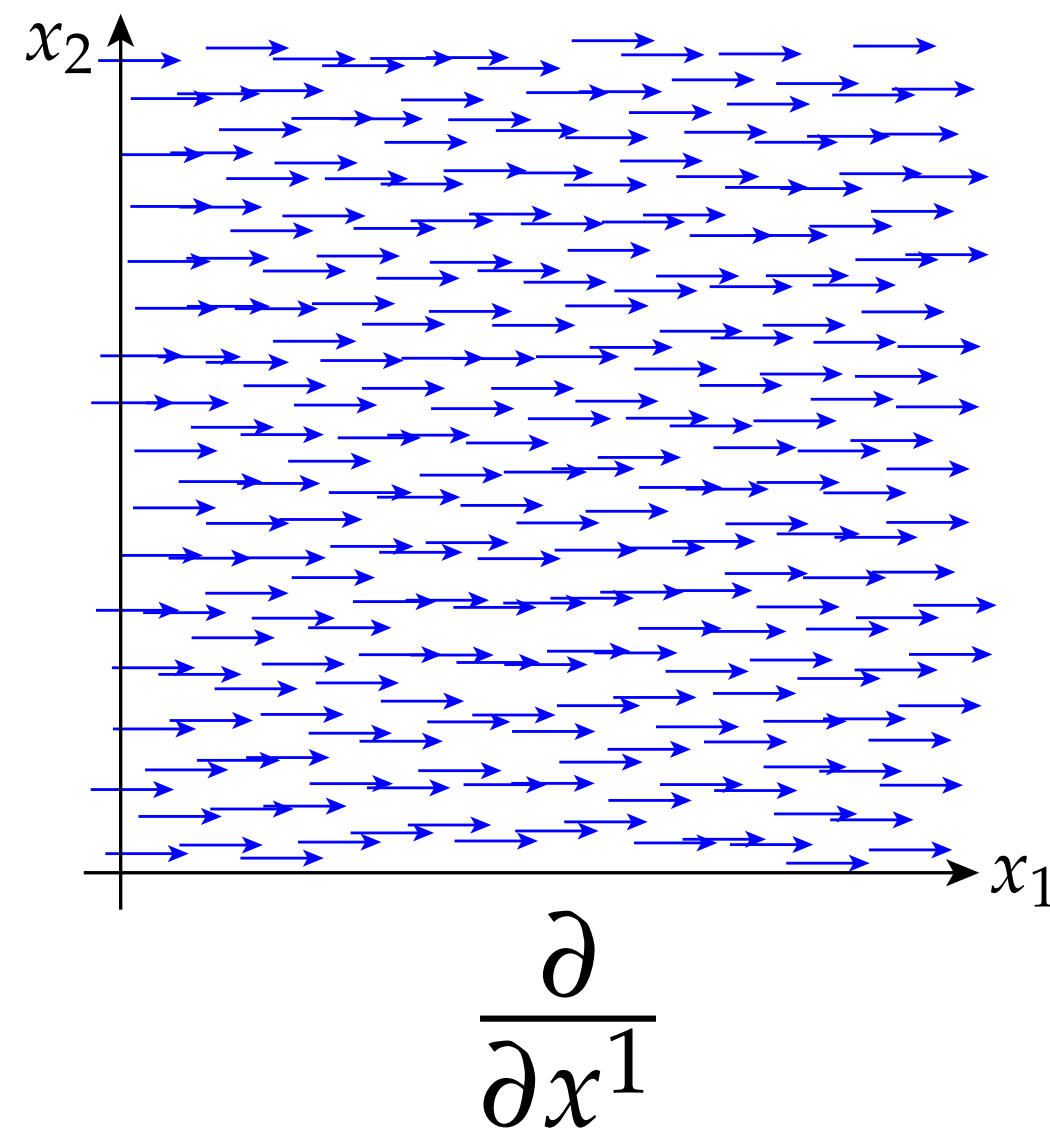
Typically we just drop the p entirely and write $\star\alpha, \alpha \wedge \beta, \alpha(X, Y), etc.$



Differential k -Forms in Coordinates

Basis Vector Fields

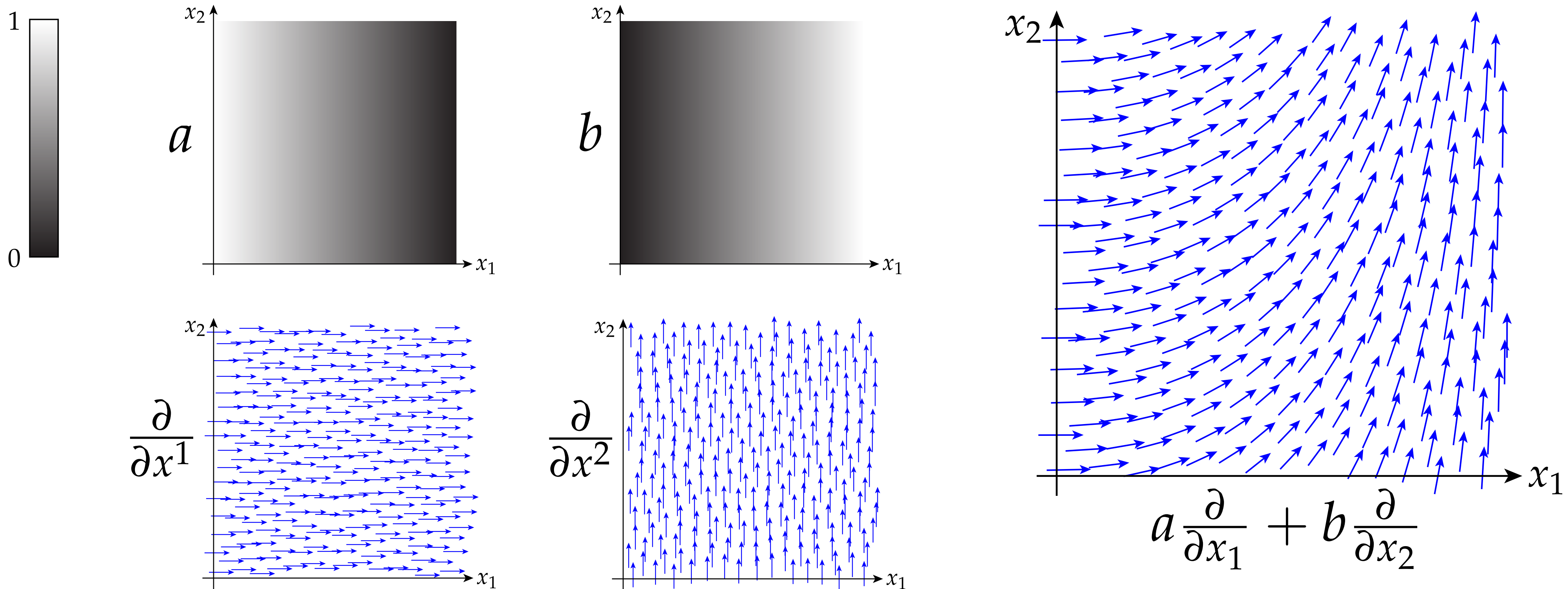
- Just as we can pick a basis for *vectors*, we can also pick a basis for *vector fields*
- The standard basis for vector fields on R^n are just **constant** vector fields of unit magnitude pointing along each of the coordinate axes:



- For historical reasons, these fields have funny-looking names that look like partial derivatives. But you will do yourself a *huge* favor by **forgetting that they have anything at all to do with derivatives!** (For now...)

Basis Expansion of Vector Fields

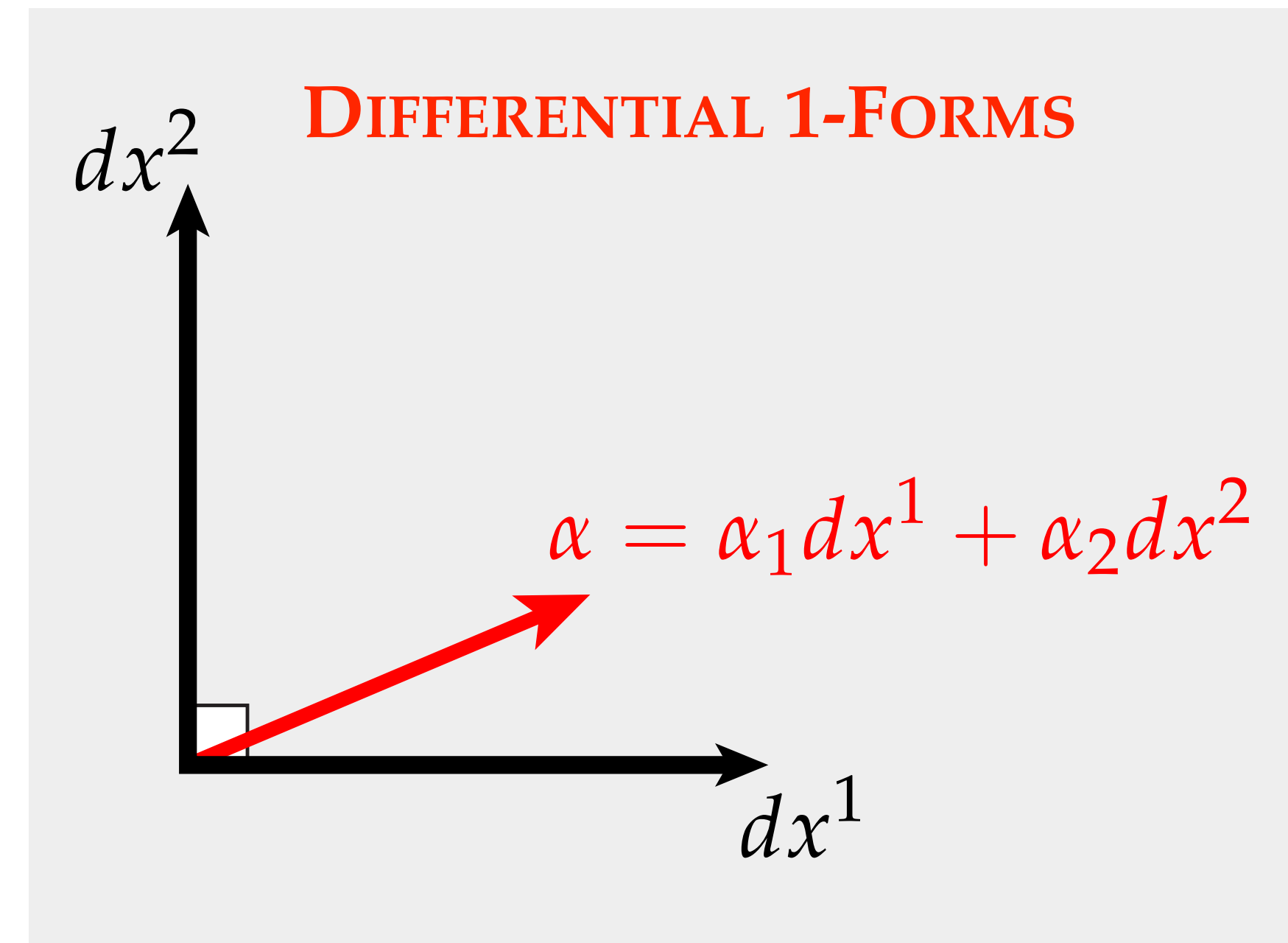
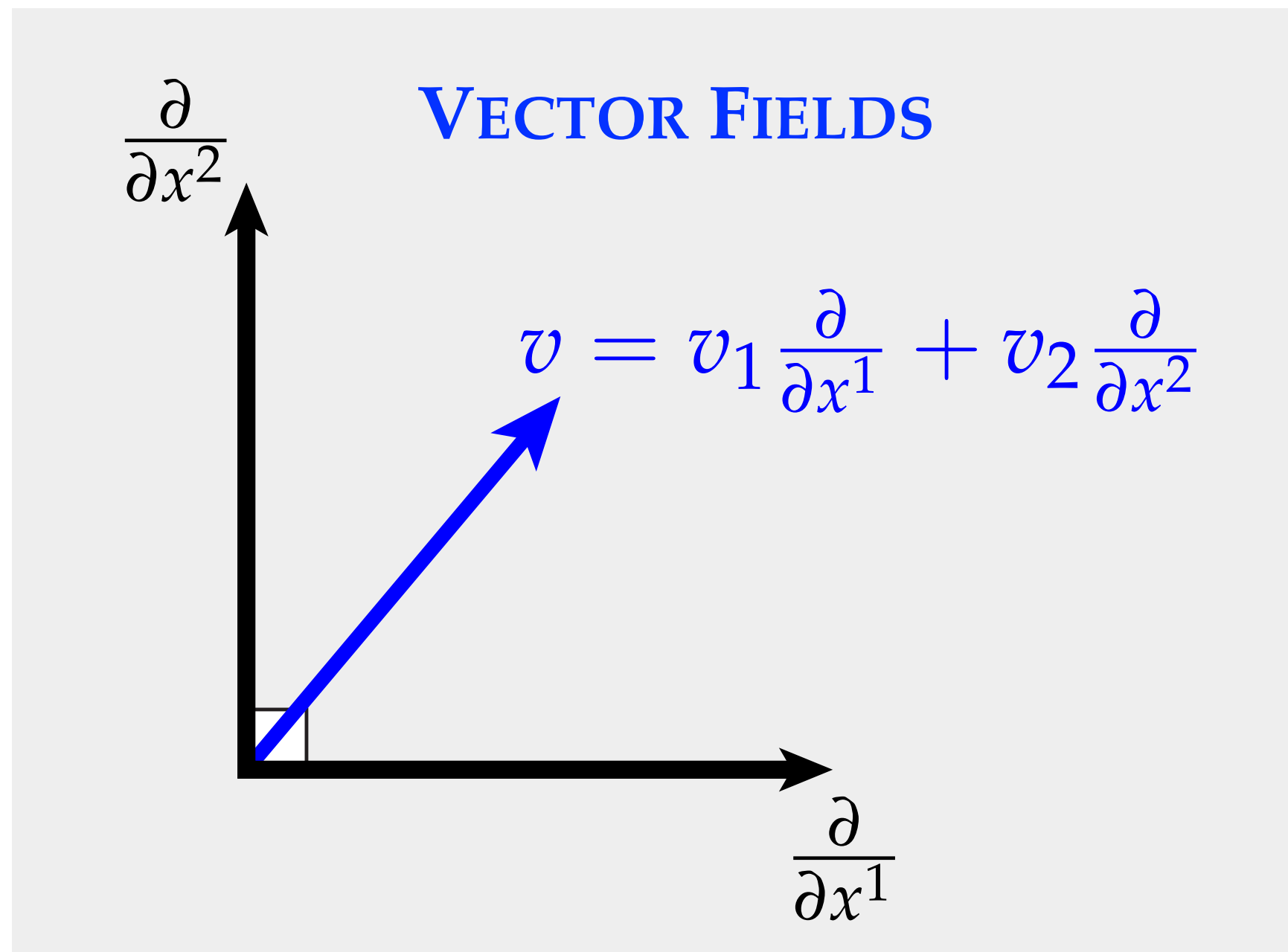
- Any other vector field is then a linear combination of the basis vector fields...
- ...but, coefficients of linear combination vary across the domain:



Q: What would happen if we didn't allow coefficients to vary?

Bases for Vector Fields and Differential 1-forms

The story is nearly identical for differential 1-forms, but with different bases:

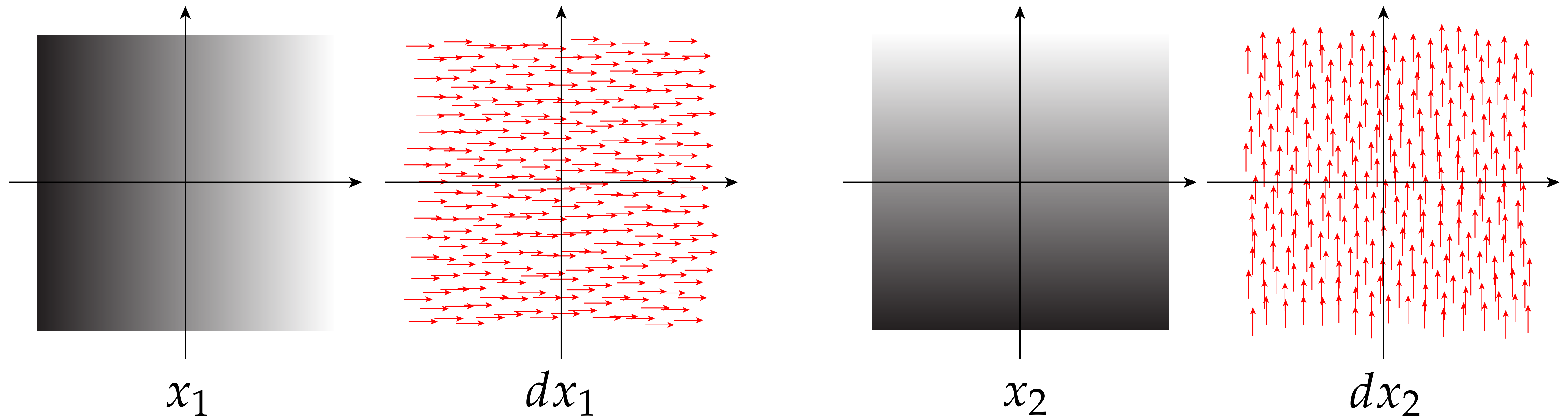


$$dx^i \left(\frac{\partial}{\partial x^j} \right) = \delta_j^i := \begin{cases} 1, & i = j \\ 0, & \text{otherwise} \end{cases}$$

Stay sane: think of these symbols as *bases*; forget they look like *derivatives*!

Coordinate Bases as Derivatives

Q: That being said, why the heck do we use symbols that look like *derivatives*?

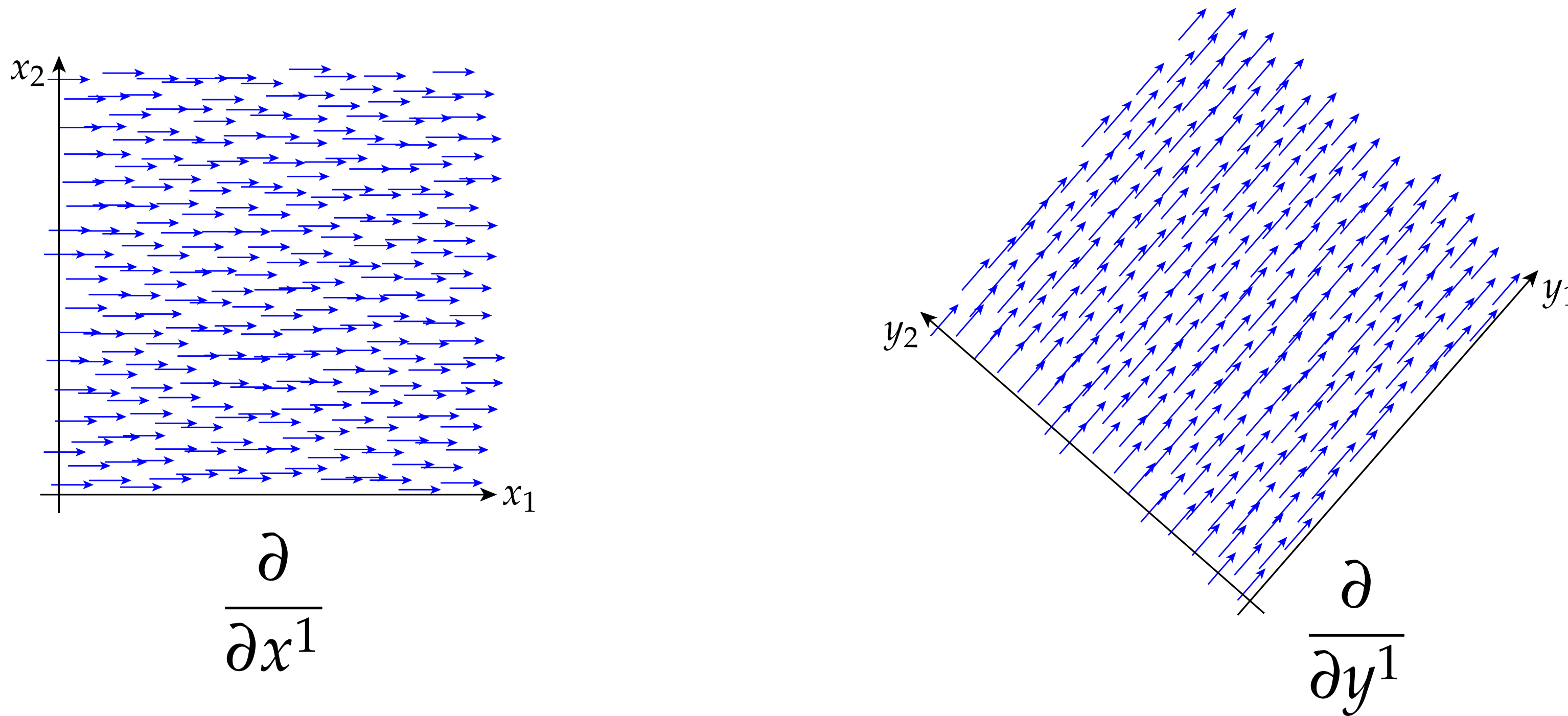


Key idea: derivative of each coordinate function yields a constant basis field.

*We'll give a more precise meaning to "d" in a little bit.

Coordinate Notation—Further Apologies

- There is at least one good reason for using this notation for basis fields
- Imagine a situation where we're working with two different coordinate systems:



- Including the name of the coordinates in our name for the basis vector field (or basis differential 1-form) makes it clear which one we mean. Not true with e_i , X_i , etc.

Example: Hodge Star of Differential 1-form

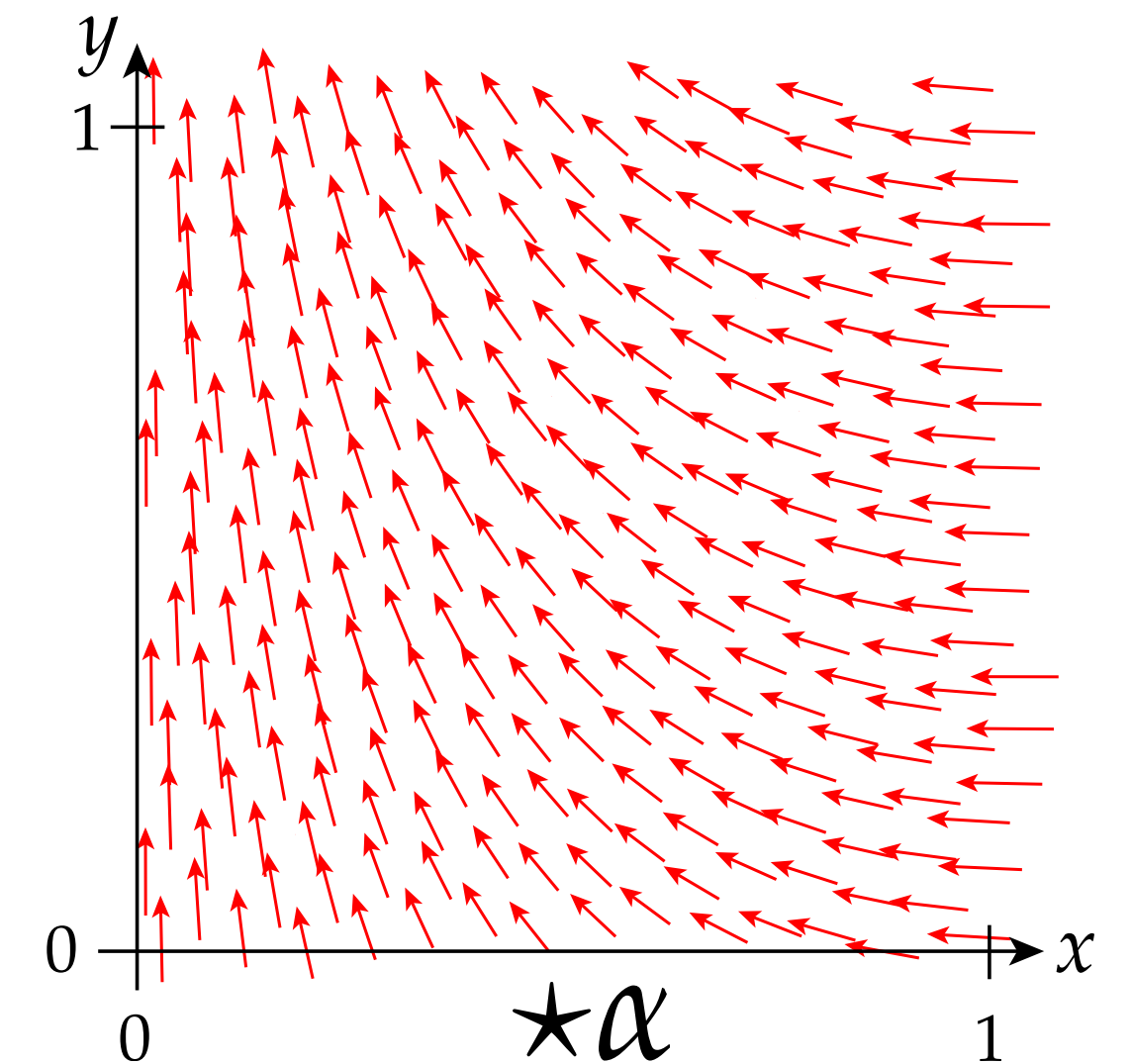
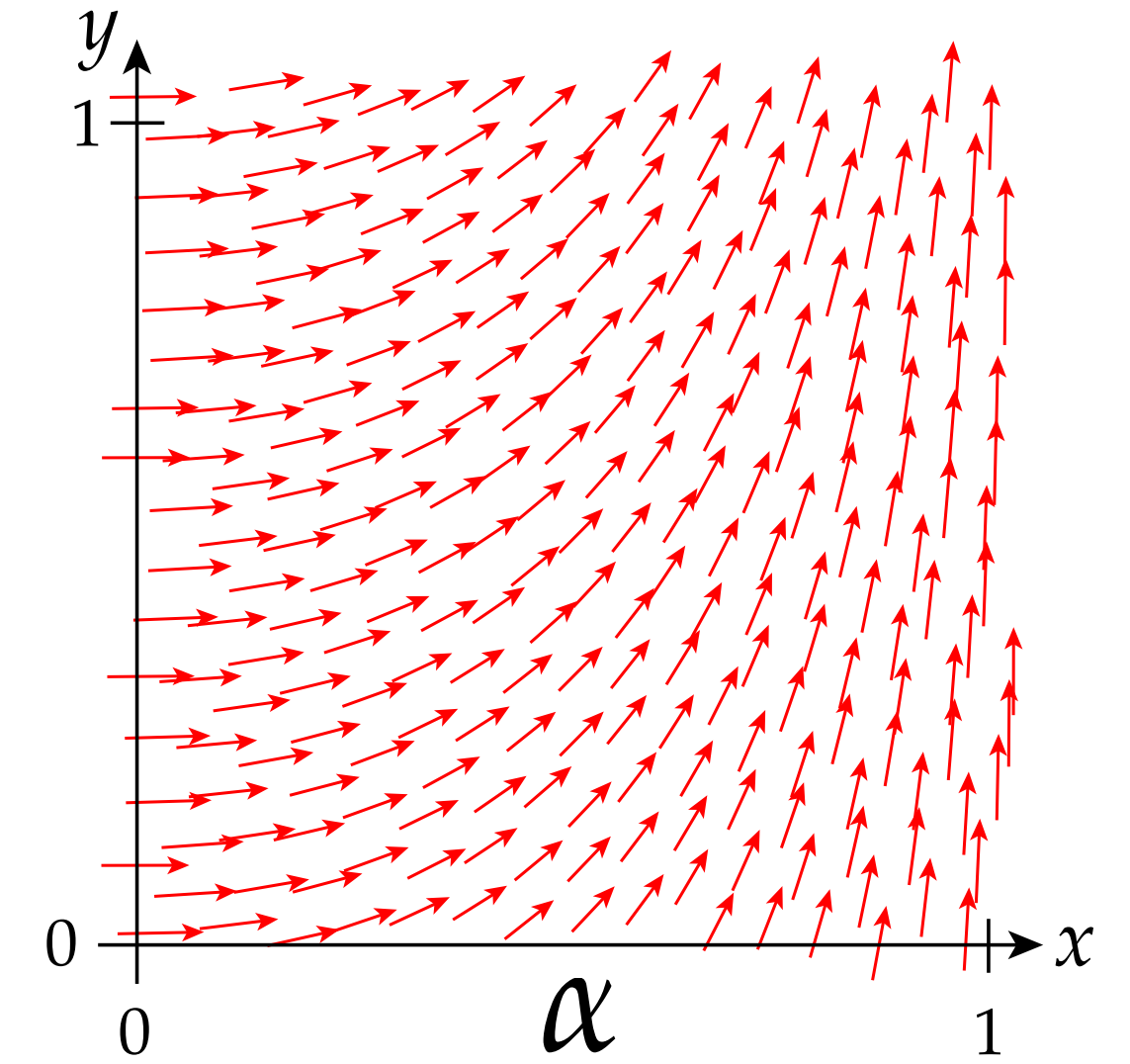
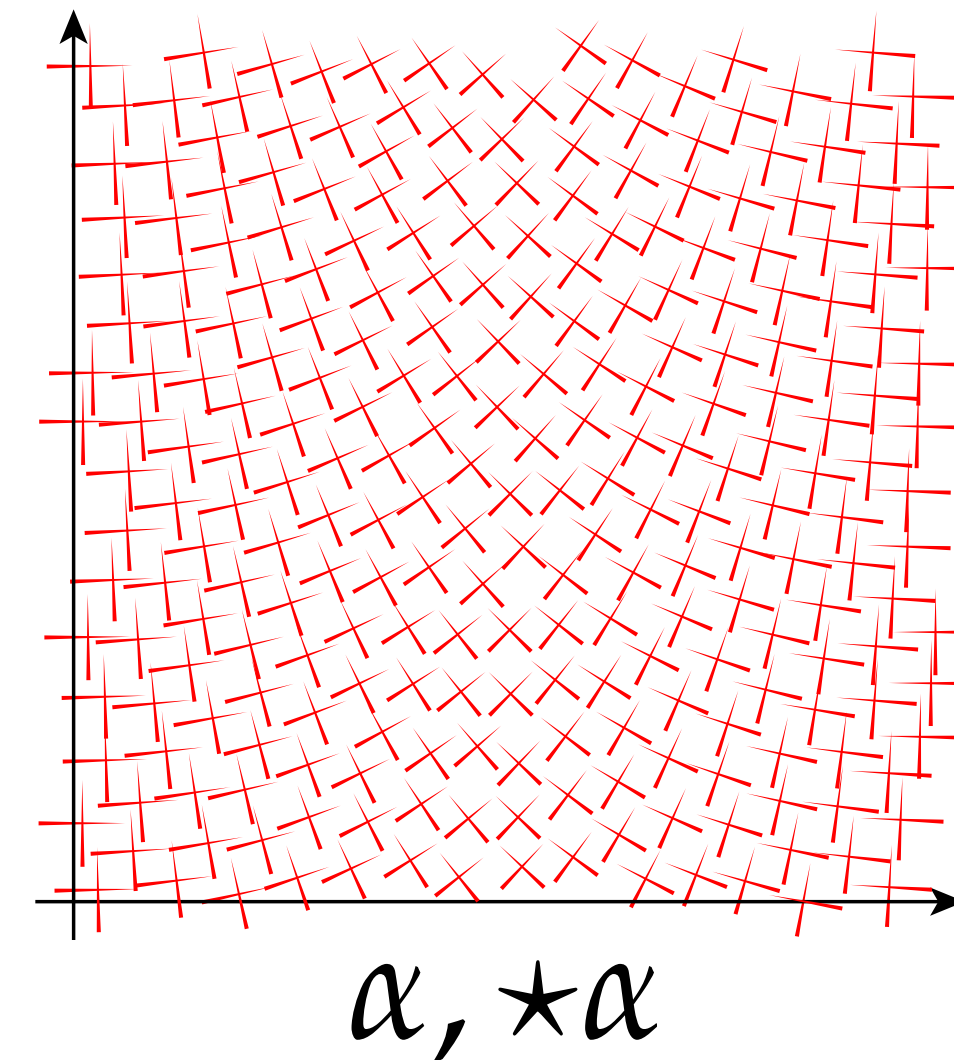
- Consider the differential 1-form $\alpha := (1 - x)dx + xdy$
 - Use coordinates (x, y) instead of (x_1, x_2)
 - Notice this expression varies over space

Q: What's its Hodge star?

$$\begin{aligned}\star\alpha &= \star((1 - x)dx) + \star(xdy) \\ &= (1 - x)(\star dx) + x(\star dy) \\ &= (1 - x)dy + -xdx\end{aligned}$$

Recall that in 2D, 1-form Hodge star is quarter-turn.

So, when we overlay the two we get little crosses...



Example: Wedge of Differential 1-Forms

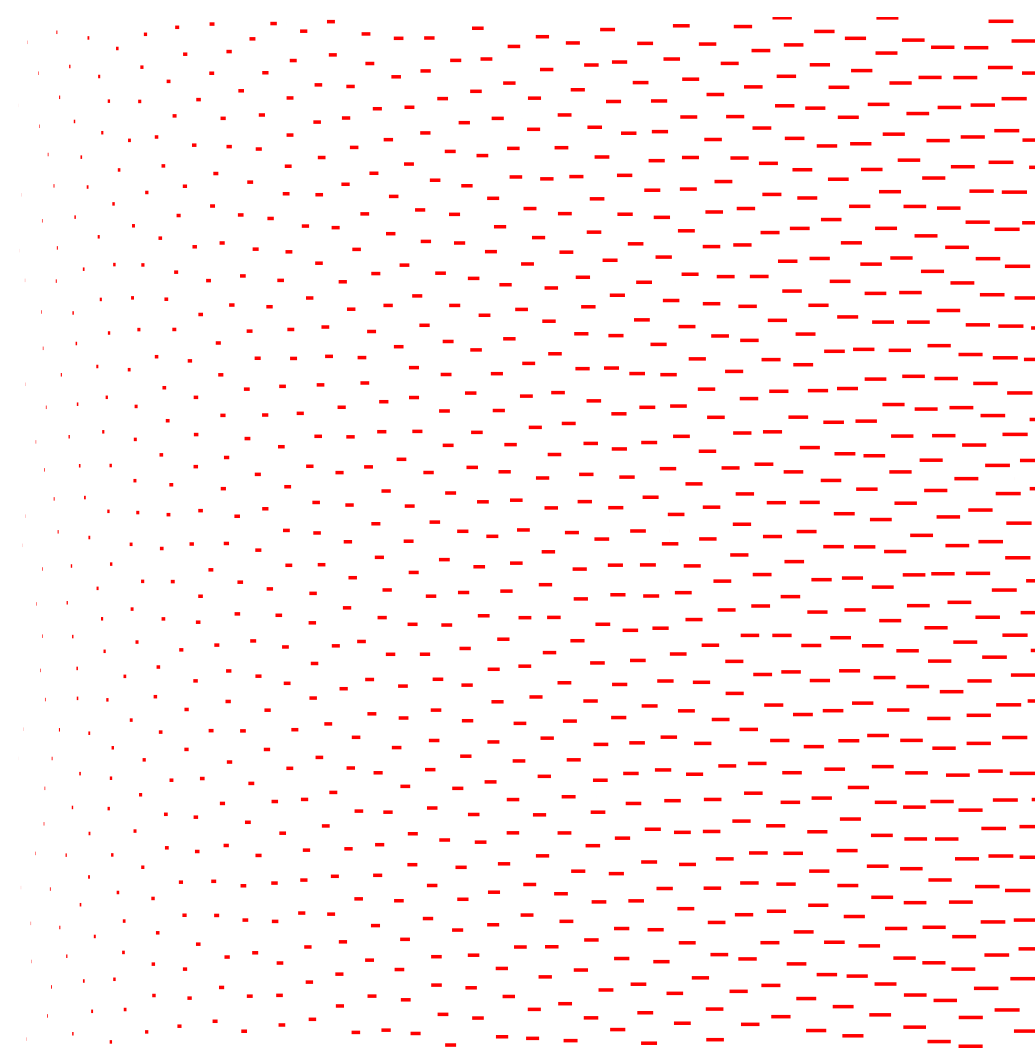
Consider the differential 1-forms*

$$\alpha := xdx, \quad \beta := (1-x)dx + (1-y)dy$$

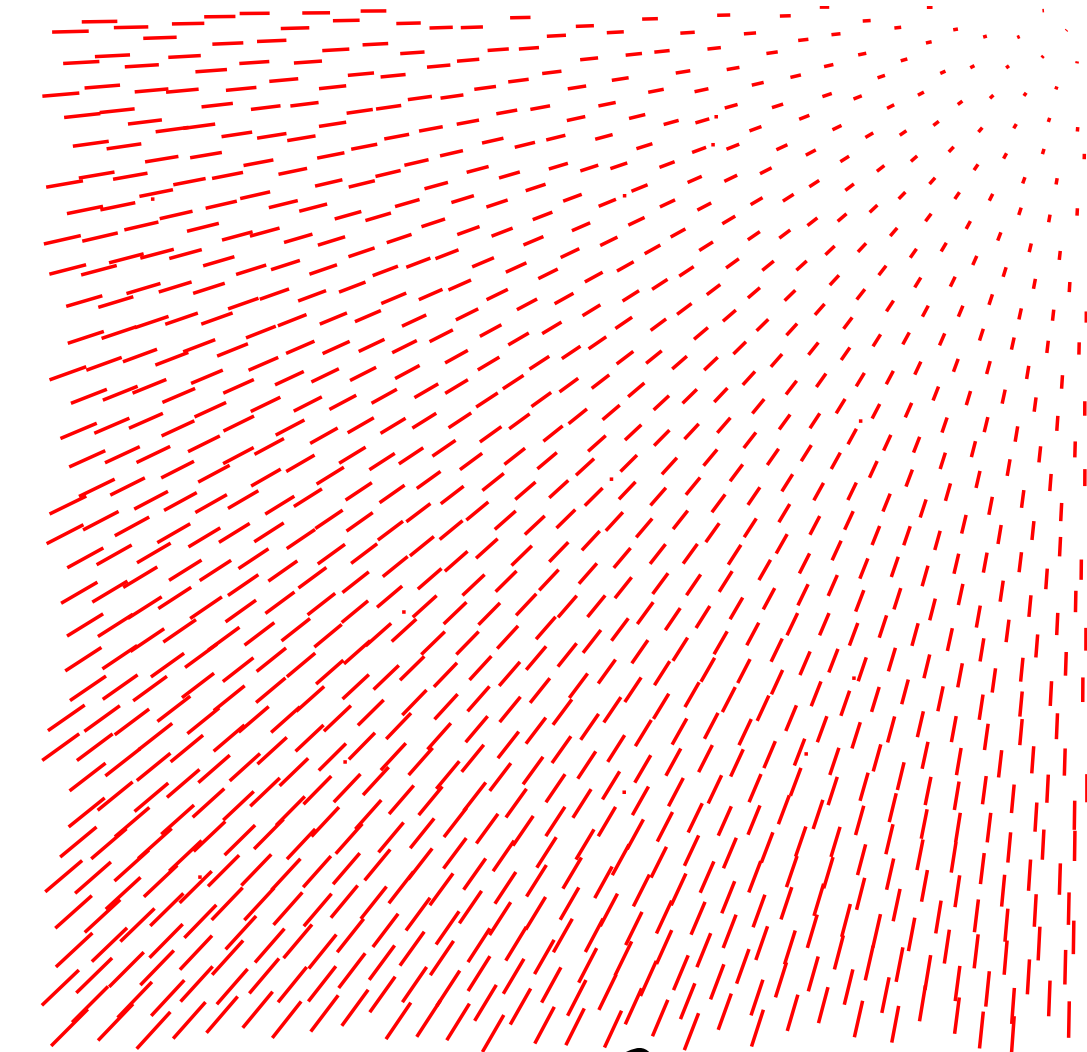
Q: What's their wedge product?

$$\begin{aligned}\alpha \wedge \beta &= (xdx) \wedge ((1-x)dx + (1-y)dy) \\ &= (xdx) \wedge ((1-x)dx) + (xdx) \wedge ((1-y)dy) \\ &= x(1-x)\cancel{dx \wedge dx}^0 + x(1-y)dx \wedge dy \\ &= (x - xy)dx \wedge dy\end{aligned}$$

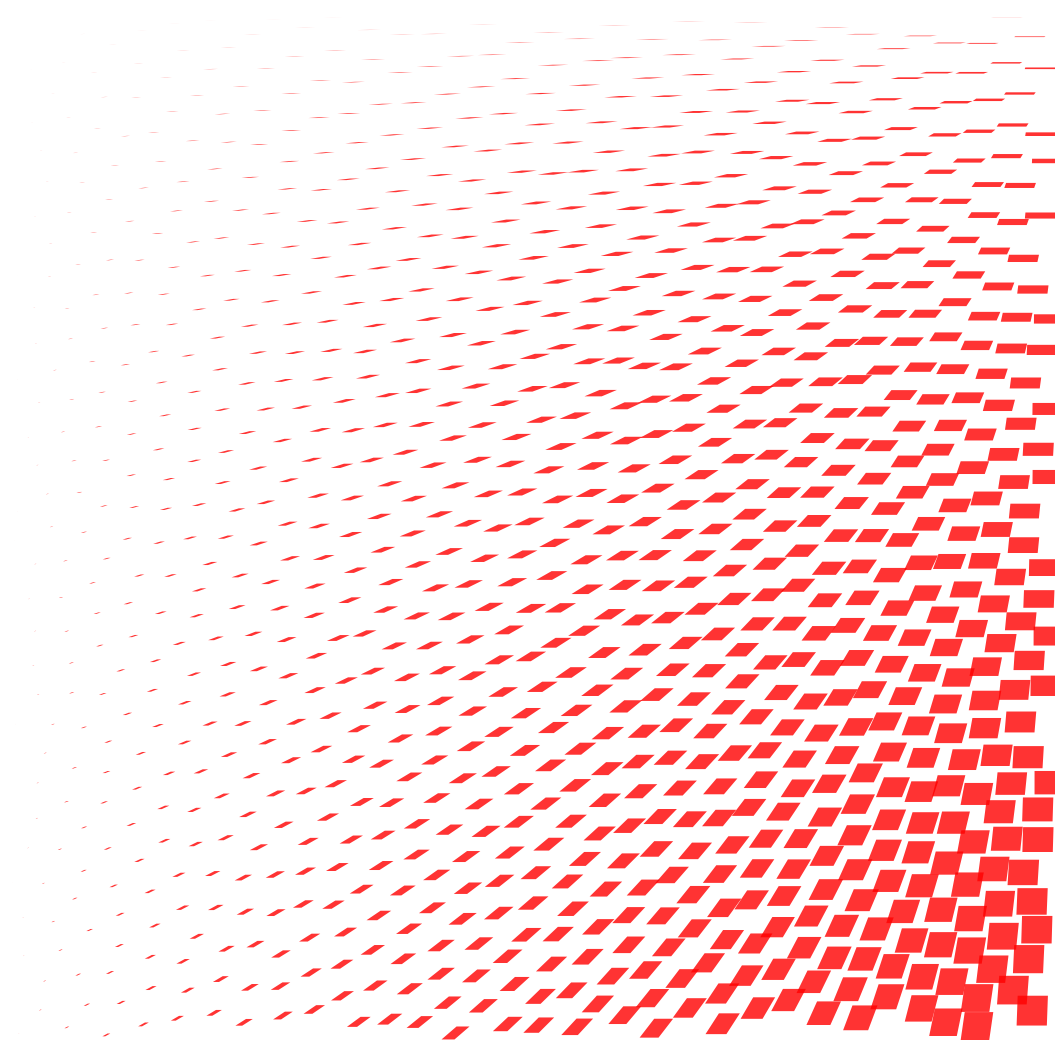
(What does the result **look** like?)



α



β



$\alpha \wedge \beta$

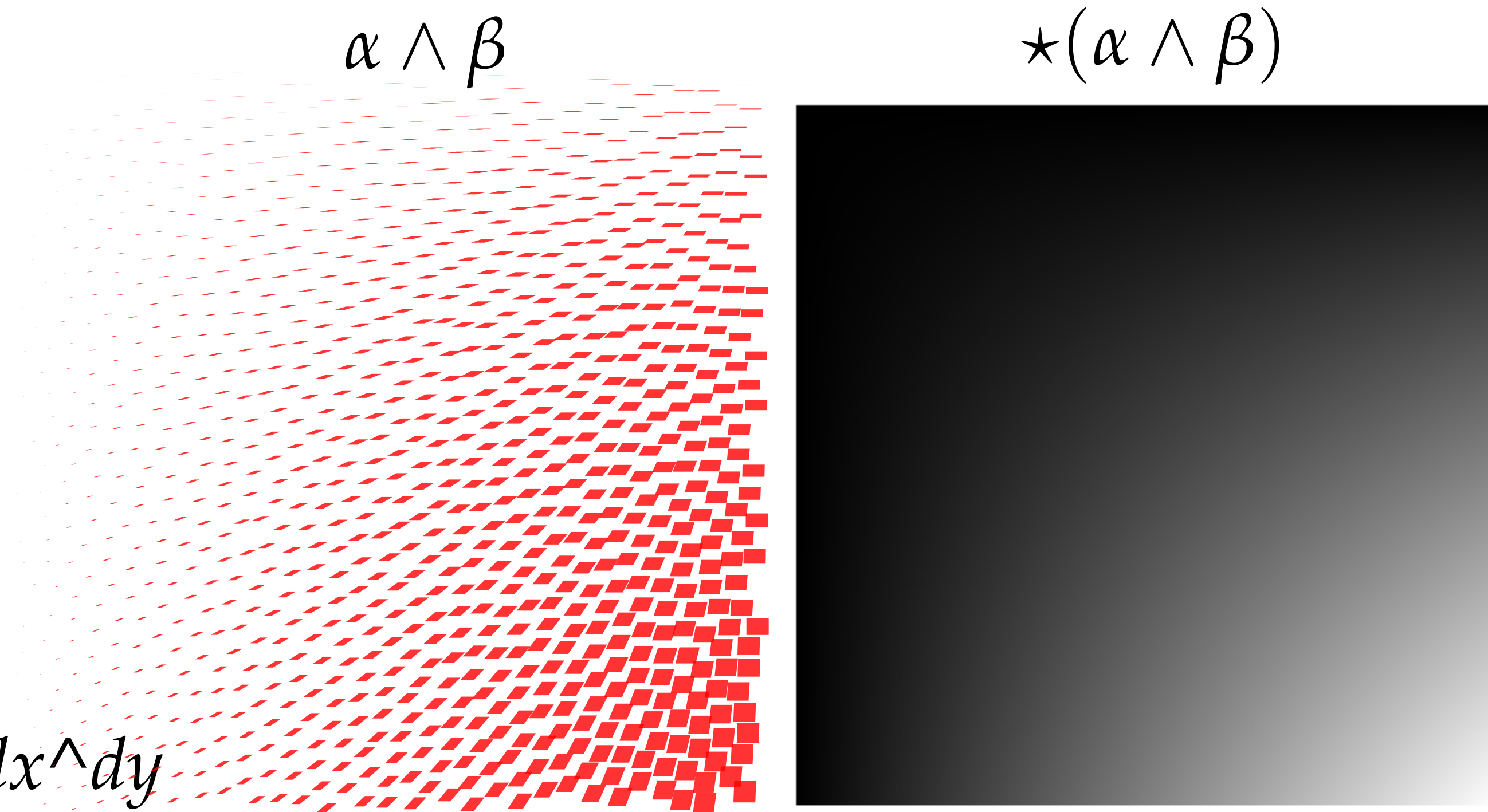
*All plots in this slide (and the next few slides) are over the unit square $[0,1] \times [0,1]$.

Volume Form / Differential n -form

- Our picture has little parallelograms
- But what information does our differential 2-form actually encode?

$$\alpha \wedge \beta = (x - xy)dx \wedge dy$$

- Has magnitude $(x-xy)$, and “direction” $dx \wedge dy$
- But in the plane, *every* differential 2-form will be a multiple of $dx \wedge dy$!
 - More precisely, some scalar function times $dx \wedge dy$, which measures *unit area*
- In n -dimensions, any *positive* multiple of $dx^1 \wedge dx^2 \wedge \cdots \wedge dx^n$ is called a *volume form*.
 - Provides some meaningful (i.e., nonzero, nonnegative) notion of volume.



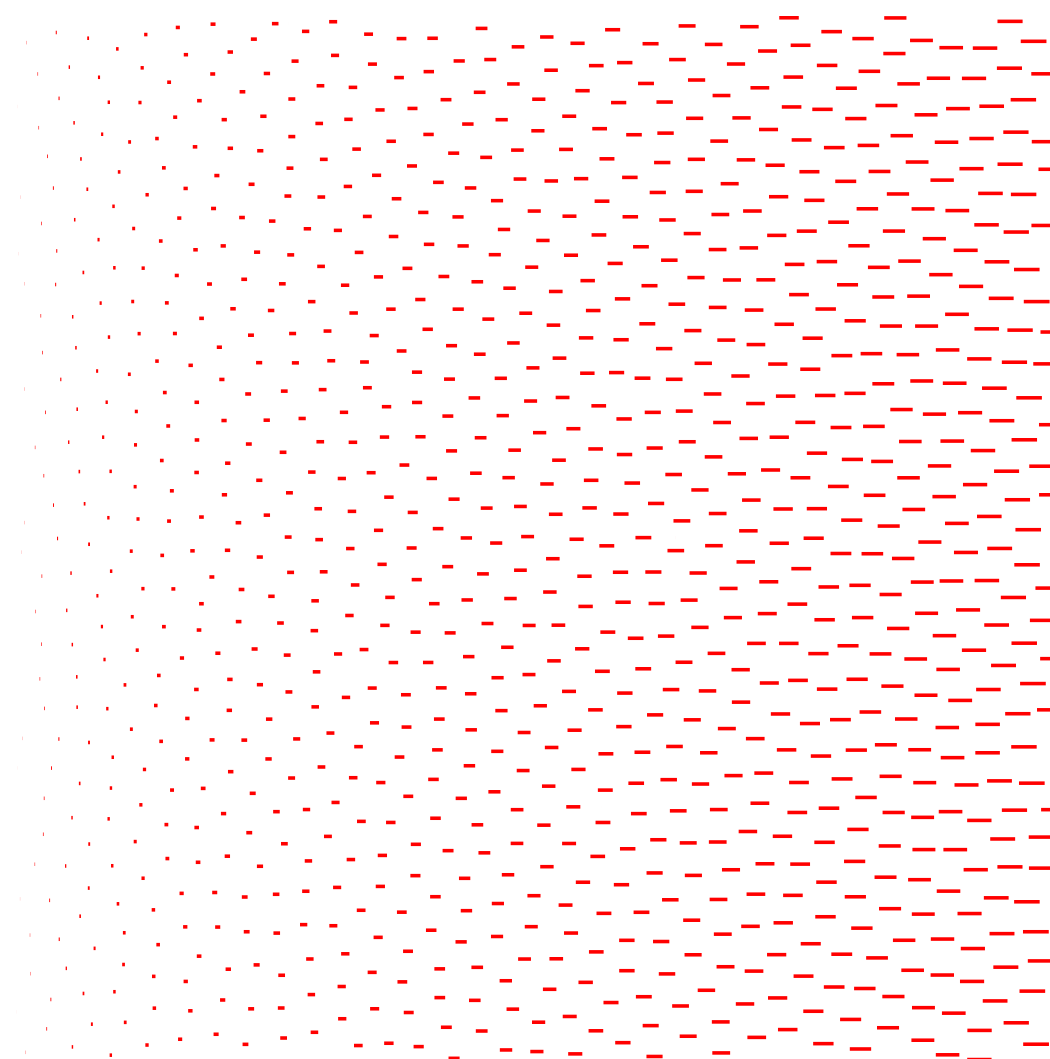
Applying a Differential 1-Form to a Vector Field

- The whole point of a differential 1-form is to measure vector fields. So let's do it!

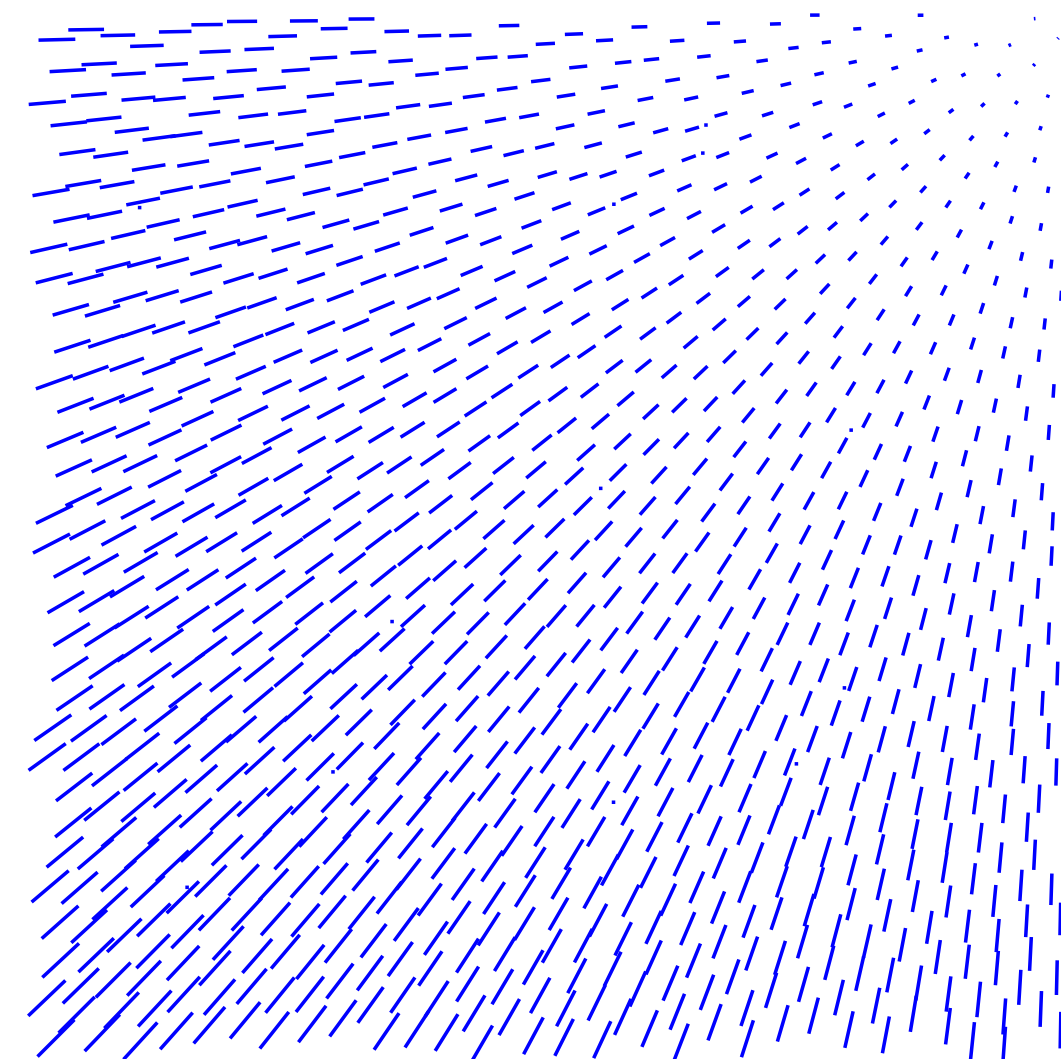
$$\begin{aligned}\alpha(X) &= (xdx) \left((1-x) \frac{\partial}{\partial x} + (1-y) \frac{\partial}{\partial y} \right) \\ &= (xdx) \left((1-x) \frac{\partial}{\partial x} \right) + (xdx) \left((1-y) \frac{\partial}{\partial y} \right) \\ &= (x - x^2) \cancel{dx \left(\frac{\partial}{\partial x} \right)}^1 + (x - xy) \cancel{dx \left(\frac{\partial}{\partial y} \right)}^0 \\ &= x - x^2\end{aligned}$$

$$\alpha := xdx$$

$$X := (1-x) \frac{\partial}{\partial x} + (1-y) \frac{\partial}{\partial y}$$



α



X

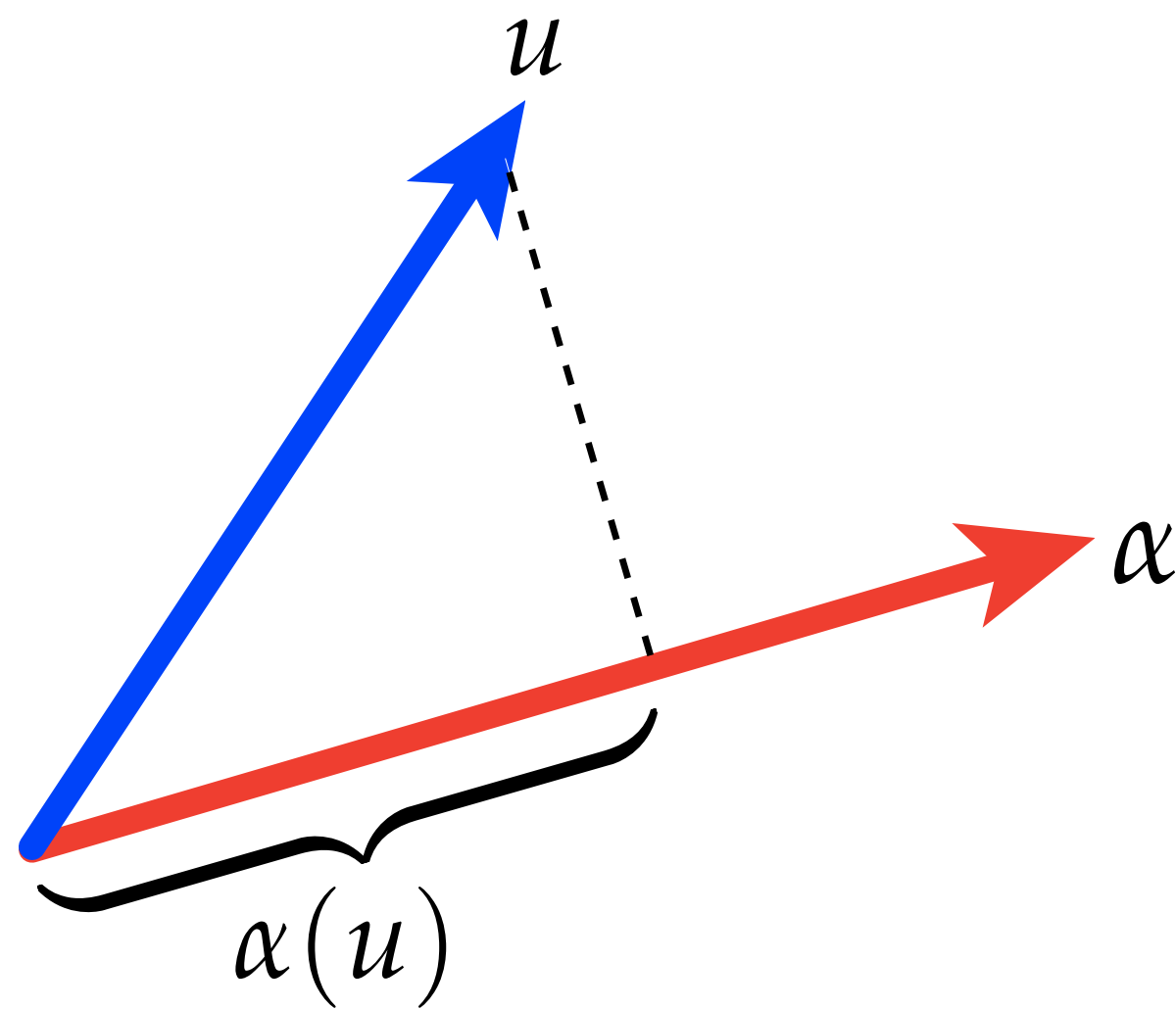


$\alpha(X)$

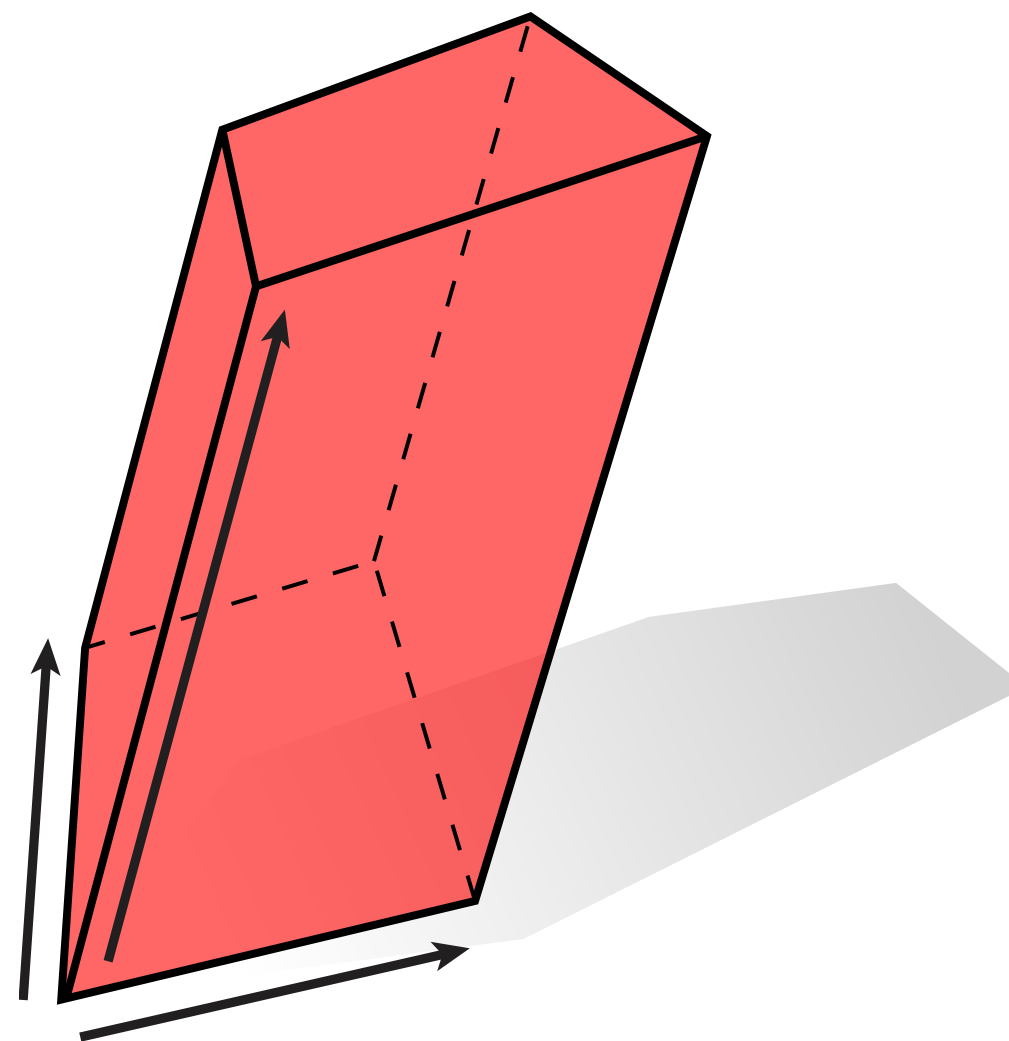
(Kind of like a dot product...)

Differential Forms in R^n - Summary

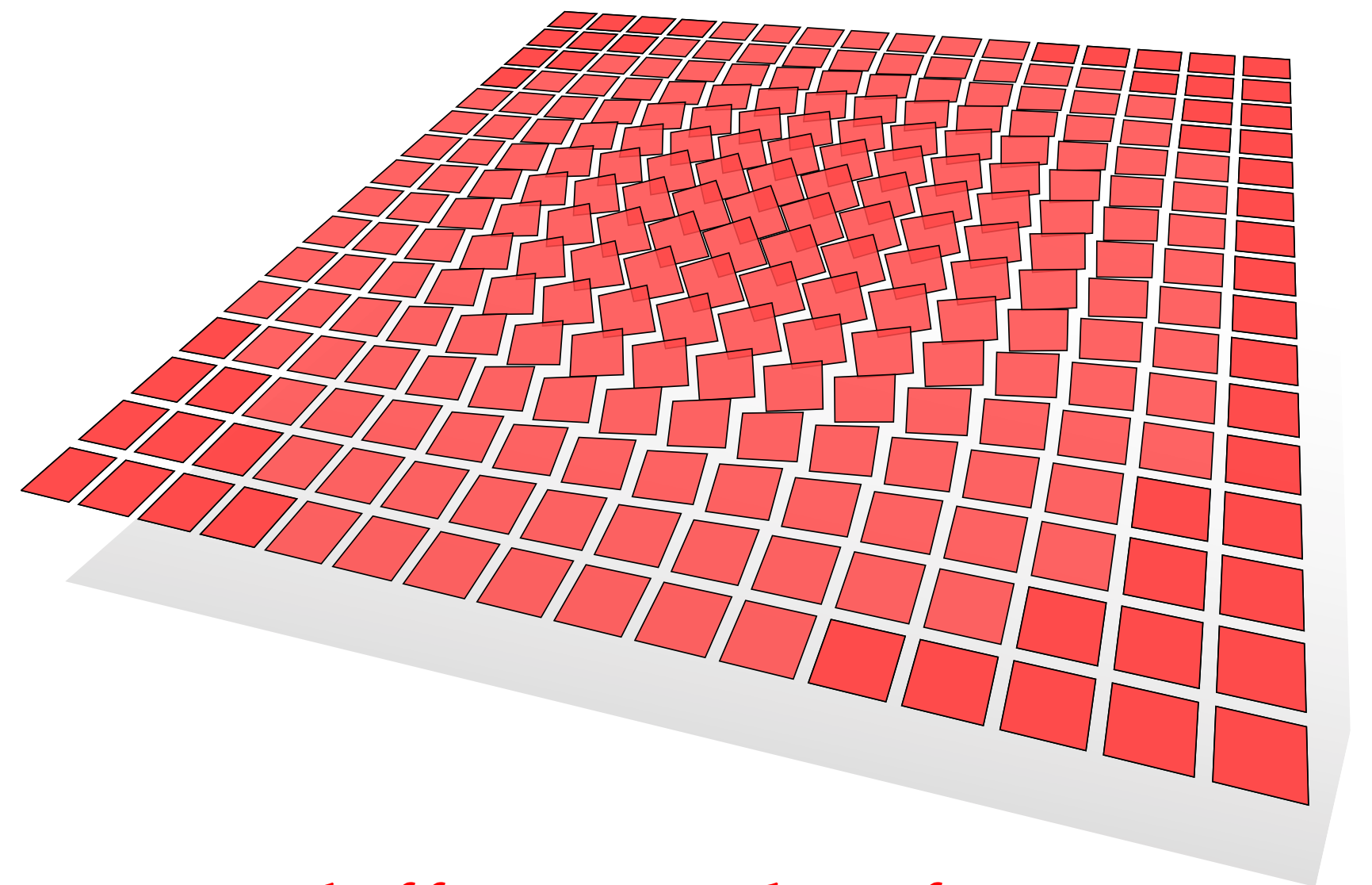
- Started with a vector space V (e.g., R^n)
 - *(1-forms)* Dual space V^* of covectors, i.e., linear measurements of vectors
 - *(k-forms)* Wedge together k covectors to get a measurement of k -dim. volumes
 - *(differential k-forms)* Put a k -form at each point of space



1-form



3-form



differential 2-form

Exterior Algebra & Differential Forms—Summary

	primal	dual
vector space	vectors	covectors
exterior algebra	k -vectors	k -forms
spatially-varying	k -vector fields	differential k -forms

Where Are We Going Next?

GOAL: develop *discrete exterior calculus* (DEC)

Prerequisites:

Linear algebra: “little arrows” (vectors)

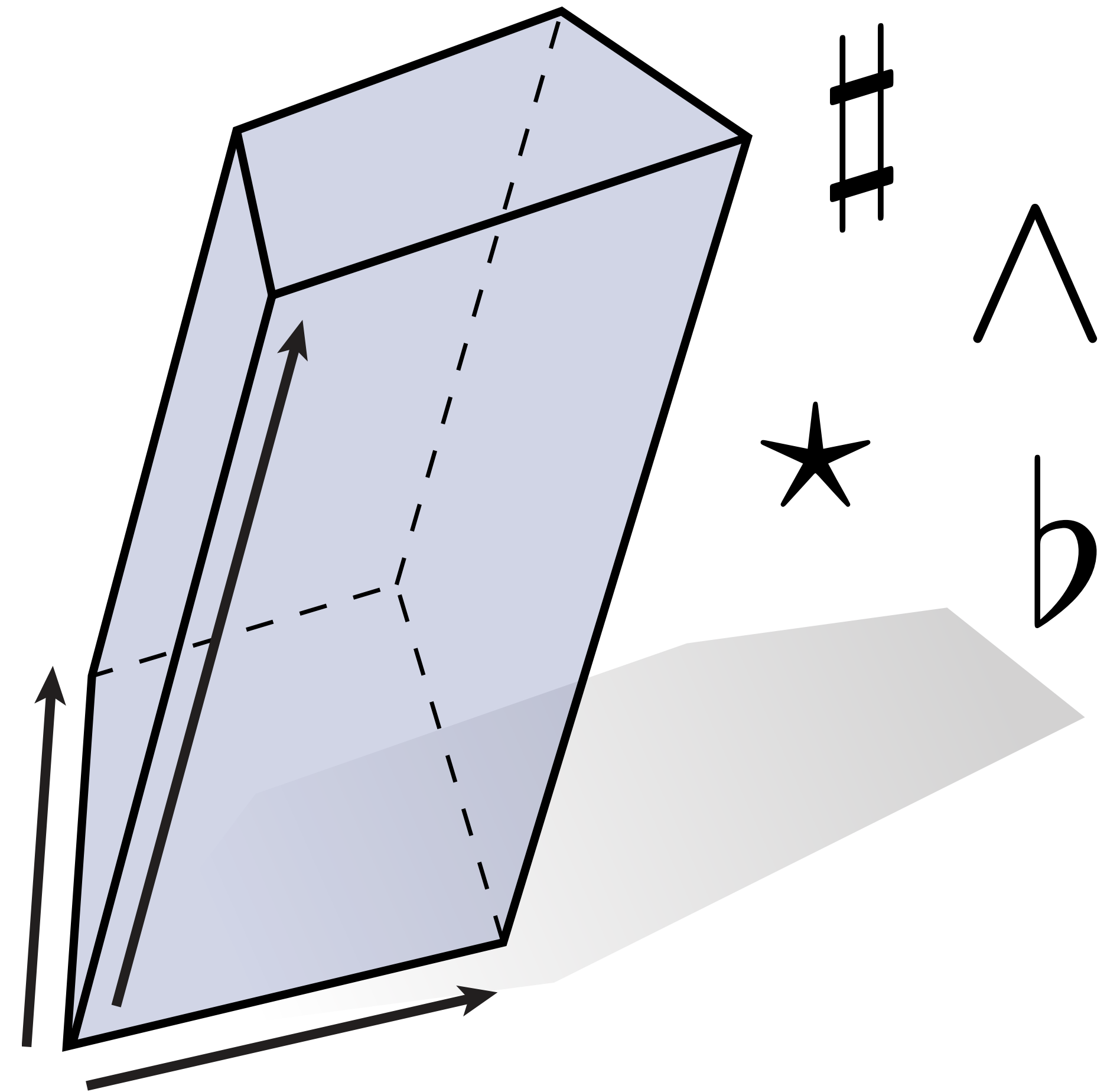
Vector Calculus: how do vectors *change*?

Next few lectures:

Exterior algebra: “little volumes” (k -vectors)

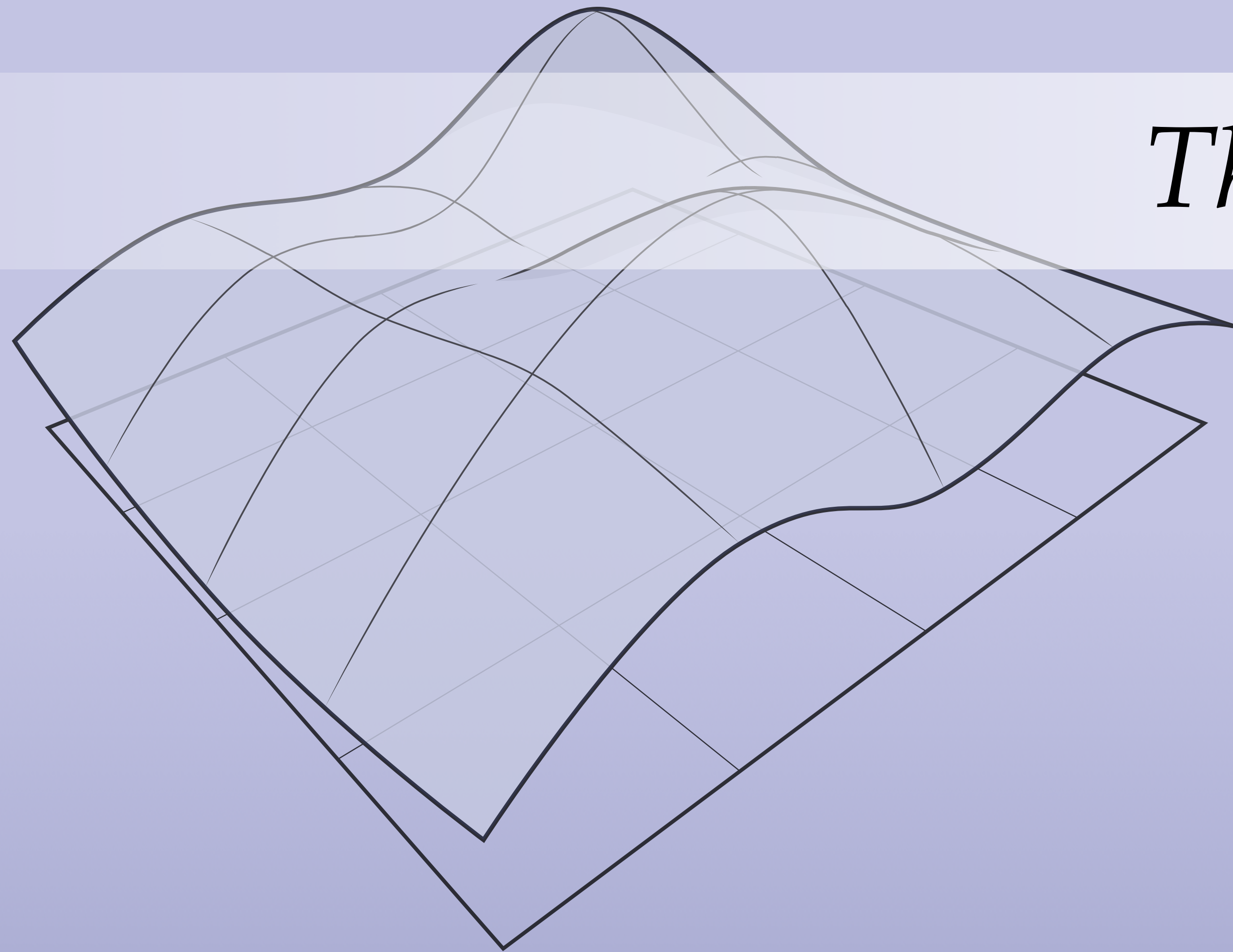
Exterior calculus: how do k -vectors change?

DEC: how do we do all of this on meshes?



Basic idea: replace vector calculus with computation on meshes.

Thanks!



DISCRETE DIFFERENTIAL GEOMETRY

AN APPLIED INTRODUCTION