EE2703: Applied Programming Lab Assignment No 5: The Resistor Problem

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1 Introduction

The aim of this assignment is to obtain the solution to the potential in a region subject to the given constraints by solving *Laplace's* equation in two-dimensions.

Laplace's equation in two-dimensions can be written as (in Cartesian coordinates):

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0$$

Using suitable approximations for a discrete point grid, we get

$$\phi_{i,j} = \frac{\phi_{i,j-1} + \phi_{i-1,j} + \phi_{i,j+1} + \phi_{i+1,j}}{4}$$

The physical significance of this would be that the potential at any point is the sum of the values at it's nearest neighbouring points. We have basically used this equation to solve the given problem by updating ϕ over many iterations till it converges within an acceptable error.

2 Solution

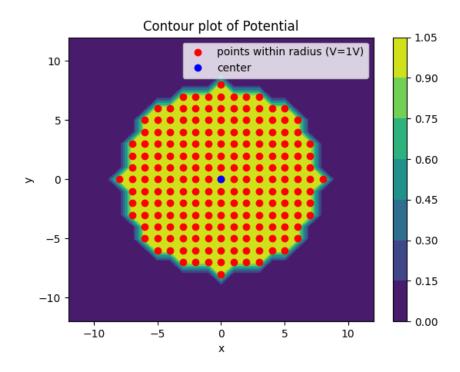
2.1 The potential solution using Laplace's equation

As mentioned earlier, we use the updation formula in multiple iterations and take care of the boundary conditions in each step as well. This is done using:

```
#creating the potential grid
phi = np.zeros((Ny,Nx))

#creating position vectors
x = np.linspace(-((Nx-1)//2), Nx//2, Nx) # since phi uses
integral indices, we cannot just do -Nx/2,Nx/2
```

```
6 y = np.linspace(-((Ny-1)//2), Ny//2, Ny)
8 X,Y = np.meshgrid(x,y)
10 #using numpy.where() to find the points within the radius
ii = np.where(X**2 + Y**2 <= radius**2)</pre>
13 #set potential at those points as one
14 phi[ii] = 1
15
^{16} #contour plot of phi versus x and y
plt.contourf(x, y, phi)
18 plt.colorbar()
19 plt.scatter(X[ii], Y[ii], c='r', marker = 'o', label = 'points
      within radius (V=1V)')
plt.plot(0,0, 'bo', label = 'center')
plt.xlabel('x')
plt.ylabel('y')
23 plt.title('Contour plot of Potential')
plt.legend()
plt.show()
26
27 def update_phi(phi, oldphi): #update phi using given formula
28 phi[1:-1, 1:-1] = (oldphi[0:-2, 1:-1] + oldphi[2:, 1:-1] +
      oldphi[1:-1, 0:-2] + oldphi[1:-1, 2:])/4
29
       return phi
  def boundary_conditions(phi): #boundary conditions
       phi[1:-1, 0] = phi[1:-1, 1]
phi[1:-1, -1] = phi[1:-1, -2]
32
33
       phi[0, 1:-1] = phi[1, 1:-1]
34
       phi[-1, 1:-1] = 0
35
       phi[ii] = 1
36
    return phi
```



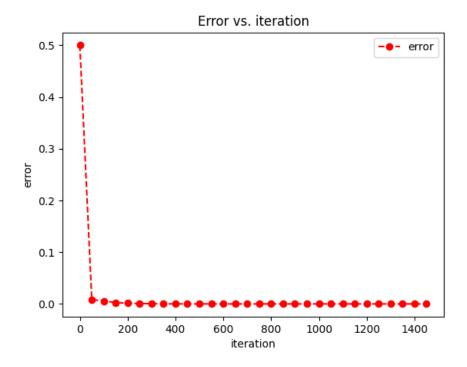
2.2 Error in estimations

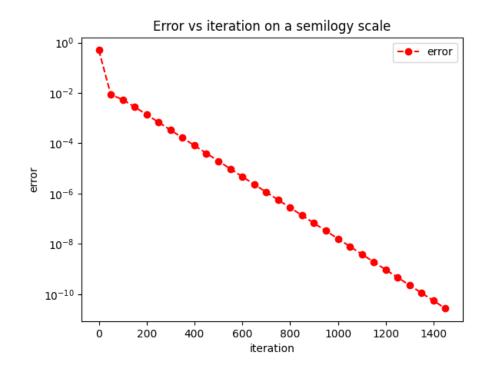
In every iteration, we keep track of the error by finding the maximum value of the error between the new phi matrix elements and the old phi matrix elements.

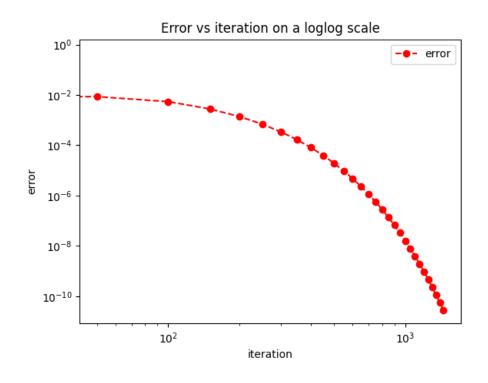
```
1 #keeping track of errors
2 errors = np.zeros(Niter)
4 for k in range(Niter):
      oldphi = np.copy(phi)
                              #making a copy of phi
      phi = update_phi(phi, oldphi) #updating phi
      phi = boundary_conditions(phi) #applying boundary
      conditions
      errors[k] = (np.abs(phi - oldphi).max()) #appending error
10 #plotting the errors
plt.plot(range(Niter)[::50], errors[::50], 'ro--', label = '
     error')
plt.xlabel('iteration')
plt.ylabel('error')
plt.title('Error vs. iteration')
plt.legend()
16 plt.show()
_{\rm 18} #plotting the errors of every 50th iteration on a semilogy plot
```

```
19 plt.semilogy(range(Niter)[::50], errors[::50], 'ro--', label =
      'error')
20 plt.xlabel('iteration')
21 plt.ylabel('error')
22 plt.title('Error vs iteration on a semilogy scale')
plt.legend()
24 plt.show()
25
_{\rm 26} #plotting the errors on a log log plot
27 plt.loglog(range(Niter)[::50], errors[::50], 'ro--', label = '
      error')
28 plt.xlabel('iteration')
29 plt.ylabel('error')
30 plt.title('Error vs iteration on a loglog scale')
31 plt.legend()
32 plt.show()
```

These errors are the plotted against the number of iterations on a linear, semilogy and a loglog scale.







We observe that the error vs number of iterations on a semilogy scale is nearly linear and thus we conclude that the error vs number of iterations might be a decaying exponential.

2.3 Least Squares Fit

Now, we attempt to extract the dependence of error on the number of iterations, by fitting an exponential to the plotted curve. This can be accomplished by creating a linear equation and then using np.linalg.lstsq().

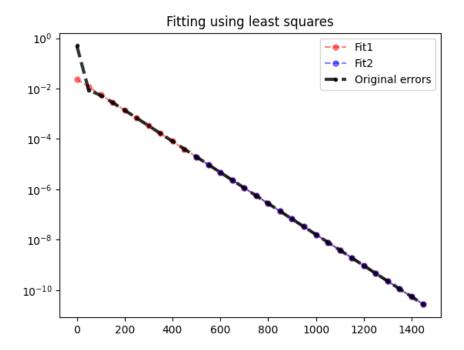
$$y = Ae^{Bx}$$
$$\log y = \log A + Bx$$

The above equation is linear, and thus np.linalg.lstsq() can be used to find the corresponding coefficients.

The corresponding code is:

```
1 #finding the least squares fit using lstsq
2 def error_fit(x, y):
      a = np.vstack([x, np.ones(len(x))]).T
      B, logA = np.linalg.lstsq(a, y)[0]
      return B, np.exp(logA)
7 #fitting an exponential
8 def exp_fit(x, A, B):
     return A*np.exp(B*x)
11 errors[errors == 0] = np.min(errors[errors != 0])*10**(-20) #
     ensuring there are no zero values before taking log
12
13 #finding fit 1
14 x1 = range(Niter)
15 y1 = np.log(errors)
16 B1, A1 = error_fit(x1, y1)
19 #finding fit 2
20 \times 2 = range(500, Niter)
21 y2 = np.log(errors[500:])
B2, A2 = error_fit(x2, y2)
23 plt.semilogy(x2[::50], exp_fit(x2[::50], A2, B2), 'bo--', label
      = 'Fit2', ms = 5, alpha = 0.5)
25 #original
26 plt.semilogy(range(Niter)[::50], errors[::50], 'ko--', label =
     'Original errors', ms = 3, linewidth = 3, alpha = 0.8)
27 plt.title('Fitting using least squares')
28 plt.legend()
29 plt.show()
```

After fitting, we plot the *fit1*, *fit2* approximations for every 50th element and the original error on the same plot.



As we see, both the fits are very close to the original graph, on zooming in, we notice that the *fit2* graph is much more closer to the original graph than *fit1*, this is because the magnitude of errors in the beginning is much larger, thus causing a larger error, whereas *fit2* only takes into account, the iterations after 500, where the errors are much smaller and thus the fit is more accurate.

2.4 Stopping condition

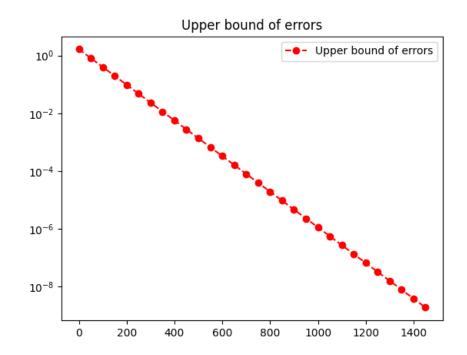
The upper bound for the error estimated with each iteration is given by:

$$Error = -\frac{A}{B} * \exp(B(N + 0.5))$$

We thus plot this upper bound of the errors vs the number of iterations on a semilogy plot.

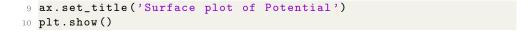
```
#upper bound of errors
def max_error(A, B, N):
    return -A*(np.exp(B*(N+0.5)))/B
```

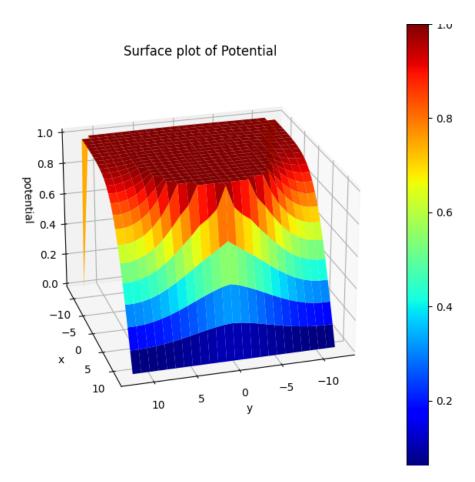
The corresponding plots are:



2.5 Surface Plot of potential

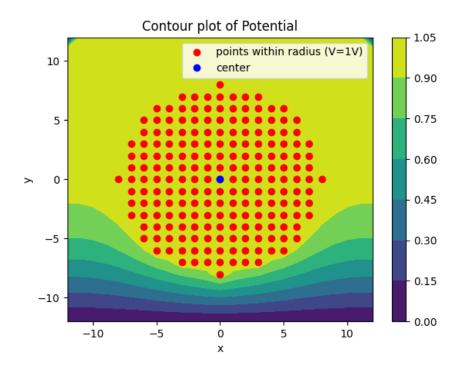
After running the algorithm using the suitable conditions, we plot our final potential grid on a 3-D surface plot for better visualisation using the plot_surface function.





2.6 Contour Plot of potential

We now plot the contourf plot for the potential after the algorithm is run using the plt.contourf function.



2.7 Vector Plot of Currents

Next, we try to find the currents in the system. We know that, the currents in the system are given as:

$$j_x = -\frac{\partial \phi}{\partial x}$$
$$j_y = -\frac{\partial \phi}{\partial y}$$

For our problem, this numerically translates to:

$$j_{x,ij} = \frac{1}{2}(\phi_{i,j-1} - \phi_{i,j+1})$$
$$j_{y,ij} = \frac{1}{2}(\phi_{i-1,j} - \phi_{i+1,j})$$

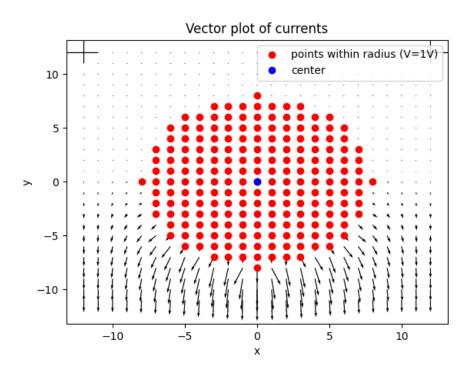
We then plot these current densities using the plt.quiver() function.

```
#vector plot of currents
#creating jx and jy of same dimensions as phi

jx = np.zeros(phi.shape)

jy = np.zeros(phi.shape)

jx [:, 1:-1] = (phi[:, 0:-2] - phi[:, 2:])/2
```



We thus notice that very little current flows through the top portion of the wire. This is because the lower surface is grounded (kept at zero potential), thus the easiest way for charge carriers to flow from the electrode would be directly through the lower half of the wire and not through the longer more resistance path through the upper half of the wire.

3 Conclusion

We thus have solved the *Laplace's* differential equation in a numerical fashion using gridpoints. The error is seen to be decaying almost exponentially and thus an exponential fit was found which fit the error curve almost exactly.

From the current density plots, we notice that most of the current only flows in the bottom half of the wire. Thus, we expect the bottom part of the wire to get heated the most due to Ohmic losses.