

# § 8 变分法

## § 8.1 量子力学的变分原理

$$\text{能量 } E(\psi) = \frac{\langle \psi | \hat{H} | \psi \rangle}{\langle \psi | \psi \rangle}, \text{ 基态能量 } E_0$$

$$E(\psi) \geq E_0$$

$$E(\psi) = \frac{\sum_n E_n |\langle n | \psi \rangle|^2}{\sum_n |\langle n | \psi \rangle|^2} \geq \frac{\sum_n E_n |\langle n | \psi \rangle|^2}{\sum_n |\langle n | \psi \rangle|^2}$$

$$\text{Ritz法: } |\psi\rangle = \sum_i \alpha_i |i\rangle$$

$$\frac{\partial E}{\partial \alpha_i} = 0$$

例: 用变分法求解一维谐振子的基态.

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$$

$$\langle x | \psi \rangle = A e^{-bx^2}$$

$$E(b) = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} = \frac{\hbar^2}{2m} b + \frac{m\omega^2}{8b}$$

$$\frac{\partial E(b)}{\partial b} = 0 \rightarrow b = \frac{m\omega}{2\hbar}, \quad E_0 \approx \frac{1}{2} \hbar \omega$$

$$\langle x | \psi \rangle = \left(\frac{2b}{\pi}\right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2}$$

## § 8.2 氦原子

$$\hat{H} = -\frac{\hbar^2}{2m} (\nabla_1^2 + \nabla_2^2) - \frac{e^2}{4\pi\epsilon_0} \left( \frac{2}{r_1} + \frac{2}{r_2} - \frac{1}{|r_1 - r_2|} \right)$$

忽略电子的相互作用:

$$\text{基态: } \phi_0(r_1, r_2) = \psi_{100}(r_1) \psi_{100}(r_2)$$

代入计算能量:

$$\langle \hat{H} \rangle = 2Z^2 E_1 + \frac{e^2}{4\pi\epsilon_0} \int \phi_0^2(r_1, r_2) \frac{1}{|r_1 - r_2|} d^3r_1 d^3r_2$$

$$\frac{1}{|r_1 - r_2|} = \frac{1}{|\vec{r}_1^2 + \vec{r}_2^2 - 2\vec{r}_1 \cdot \vec{r}_2 \cos\theta|} = \begin{cases} \sum_l \frac{r_1^l}{r_2^{l+1}} P_l(\cos\theta) & 0 \leq r_1 \leq r_2 \\ \sum_l \frac{r_2^l}{r_1^{l+1}} P_l(\cos\theta) & r_1 > r_2 \end{cases}$$

$$\int_{-1}^1 P_m(x) P_n(x) dx = \frac{1}{2n+1} \delta_{mn}, \quad P_0(x) = 1$$

$$\int \phi_0^2(r_1, r_2) \frac{1}{|r_1 - r_2|} d^3r_1 d^3r_2$$

=

$$= \frac{32}{8a}$$

考虑屏蔽效应, 假设 1 个电子屏蔽后的核电荷为  $z - \eta$