

# The Fermi Hole

Ivan Kulesh    Martijn Papendrecht

Delft University of Technology, The Netherlands

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## Some properties of Fermi gas

Particles dwell in a large but finite volume  $\mathcal{V}$ .

Periodical boundary conditions.

Restriction on values of the wave vector:

$$\mathbf{k} = 2\pi \left( \frac{n_x}{L_x}, \frac{n_y}{L_y}, \frac{n_z}{L_z} \right) \quad (1)$$

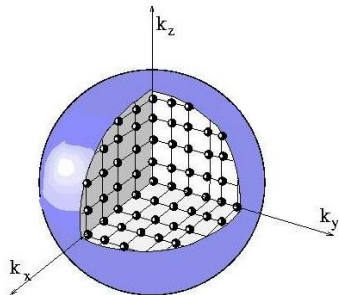
The volume of the unit cell in k-space:

$$dk_x dk_y dk_z = d^3k = \frac{(2\pi)^3}{\mathcal{V}} \quad (2)$$

## Some properties of Fermi gas

Particles occupy all energy levels below Fermi energy. Volume in k-space:

$$V_k = \frac{4}{3}\pi k_F^3 \quad (3)$$



## Some properties of Fermi gas

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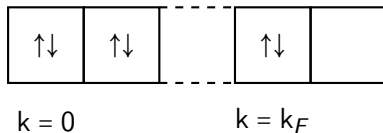
$$V_k = \frac{4}{3}\pi k_F^3 \quad (3)$$

Number of particles:

$$N = g_\sigma \frac{V_k}{d^3k} = \frac{2 \frac{4}{3}\pi k_F^3 \mathcal{V}}{8\pi^3} = \frac{k_F^3 \mathcal{V}}{3\pi^2} \implies \quad (4)$$
$$k_F^3 = \frac{3\pi^2 N}{\mathcal{V}}$$

## Fermionic ground state

In the ground state  $|g\rangle$  fermions occupy all available levels in the  $k$ -space up to  $k_F$ :



## Calculating sum

To continue we need to find

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{\mathcal{V}} \sum_{\mathbf{k}} e^{-i\mathbf{k}\mathbf{r}} \langle g | \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} | g \rangle$$

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Number of particles operator

$$\hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} = \hat{n}_{\mathbf{k}\sigma}$$

$$\langle g | \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} | g \rangle = \langle g | \hat{n}_{\mathbf{k}\sigma} | g \rangle = \begin{cases} 1 & k \leq k_F \\ 0 & k > k_F \end{cases}$$



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Heaviside step function

$$\langle g | \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} | g \rangle = H(k_F - |\mathbf{k}|)$$

## Calculating sum

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Change to summation inside the Fermi sphere

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{\mathcal{V}} \sum_{|\mathbf{k}| \leq k_F} e^{-i\mathbf{k}\mathbf{r}}$$

# From summation to integration

Multiply and divide by unit cell volume

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{\mathcal{V}} \sum_{|\mathbf{k}| \leq k_F} e^{-i\mathbf{k}\mathbf{r}} = \frac{1}{\mathcal{V}} \frac{\mathcal{V}}{8\pi^3} \sum_{|\mathbf{k}| \leq k_F} e^{-i\mathbf{k}\mathbf{r}} d^3\mathbf{k}$$

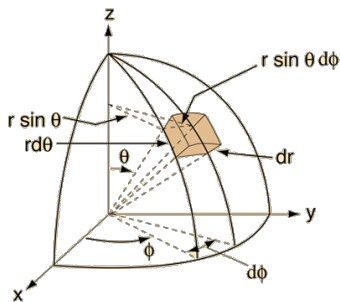
In the limit  $\mathcal{V} \rightarrow \infty$

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{8\pi^3} \sum_{|\mathbf{k}| \leq k_F} e^{-i\mathbf{k}\mathbf{r}} d^3\mathbf{k} \rightarrow \frac{1}{8\pi^3} \int_{|\mathbf{k}| \leq k_F} e^{-i\mathbf{k}\mathbf{r}} d^3\mathbf{k}$$

# Evaluation of the integral

We will integrate in spherical coordinates

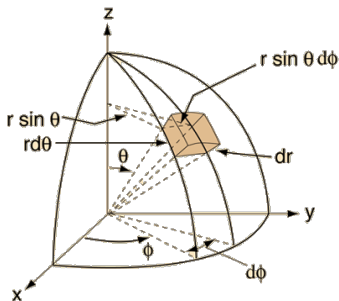
$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{8\pi^3} \int_{|\mathbf{k}| \leq k_F} e^{-i\mathbf{k}\mathbf{r}} d^3k$$



# Evaluation of the integral

We will integrate in spherical coordinates

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{8\pi^3} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} e^{-i\mathbf{k}\mathbf{r} \cos \theta} k^2 \sin \theta \, d\phi d\theta dk$$



# Evaluation of the integral

Integration over  $\phi$

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{8\pi^3} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} e^{-i\mathbf{k}\mathbf{r}\cos\theta} k^2 \sin\theta \, d\phi d\theta dk$$

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} e^{-i\mathbf{k}\mathbf{r}\cos\theta} k^2 \sin\theta \, d\theta dk$$

# Evaluation of the integral

Integration over  $\theta$

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} e^{-ikr \cos \theta} k^2 \sin \theta \, d\theta \, dk$$

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} e^{-ikr \cos \theta} k^2 -d \cos \theta \, d\theta \, dk$$

$$F(\mathbf{k}, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \int_{\eta=-1}^1 e^{ikr \eta} k^2 \, d\eta \, dk$$

# Evaluation of the integral

Integration over  $\theta$

$$F(k, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} e^{-ikr \cos \theta} k^2 \sin \theta \, d\theta \, dk$$

$$F(k, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \int_{\theta=0}^{\pi} e^{-ikr \cos \theta} k^2 -d \cos \theta \, d\theta \, dk$$

$$F(k, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \frac{2}{kr} \sin(kr) k^2 \, dk$$



# Evaluation of the integral

Integration over  $k$

$$F(k, \mathbf{r}) = \frac{1}{4\pi^2} \int_{k=0}^{k_F} \frac{2}{kr} \sin(kr) k^2 dk = \frac{1}{2\pi^2 r^3} \int_{k=0}^{k_F} \sin(kr) kr dk$$

$$F(k, \mathbf{r}) = \frac{1}{2\pi^2 r^3} \{ \sin(k_F r) - k_F r \cos(k_F r) \}$$

# Evaluation of the integral

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$$F(k, \mathbf{r}) = \frac{k_F^3}{2\pi^2 r^3 k_F^3} \{ \sin(k_F r) - k_F r \cos(k_F r) \}$$

Fermi wavenumber

$$k_F^3 = \frac{3\pi^2 N}{\mathcal{V}}$$

## Evaluation of the integral

$$\frac{1}{\mathcal{V}} \sum_{\mathbf{k}} e^{-i\mathbf{k}\mathbf{r}} \langle g | \hat{c}_{\mathbf{k}\sigma}^\dagger \hat{c}_{\mathbf{k}\sigma} | g \rangle = \left( \frac{N}{2\mathcal{V}} \right) \frac{3 \{ \sin(k_F r) - k_F r \cos(k_F r) \}}{(k_F r)^3}$$

No divergence at zero: at  $k_F r \rightarrow 0 \iff r \ll \lambda_F$

$$3 \frac{\sin(k_F r) - k_F r \cos(k_F r)}{(k_F r)^3} \approx 3 \frac{x - x^3/6 - x(1 - x^2/2)}{x^3} = 3 \frac{1}{3} = 1$$

Begin with the equations.

$$|\phi_\sigma(\mathbf{x})\rangle = \hat{\psi}_\sigma(\mathbf{x}) |g\rangle$$

$$\hat{\psi}_\sigma(\mathbf{x}) = \sum_k \hat{c}_{k,\sigma} \psi_{k,\sigma}(\mathbf{x})$$

$$\psi_k(\vec{x}) = \frac{1}{\sqrt{\mathcal{V}}} e^{-i\vec{k}\cdot\vec{x}}$$

$$\langle \phi_\sigma(\mathbf{x}) | \hat{\psi}_{\sigma'}^\dagger(\mathbf{x}') \hat{\psi}_{\sigma'}(\mathbf{x}') | \psi_\sigma(\mathbf{x}) \rangle$$

$$\langle g | \hat{\psi}_\sigma^\dagger(\mathbf{x}) \hat{\psi}_{\sigma'}^\dagger(\mathbf{x}') \hat{\psi}_{\sigma'}(\mathbf{x}') \hat{\psi}_\sigma(\mathbf{x}) | g \rangle$$

$$\langle g | \sum_k c_{k,\sigma}^\dagger \psi_{k,\sigma}^*(\mathbf{x}) \sum_l c_{l,\sigma'}^\dagger \psi_{l,\sigma'}^*(\mathbf{x}') \sum_m c_{m,\sigma'} \psi_{m,\sigma'}(\mathbf{x}') \sum_n c_{n,\sigma} \psi_{n,\sigma}(\mathbf{x}) | g \rangle$$

Now we see two creation and two annihilation operators, but N should be conserved, so there are two cases:

$$k, \sigma = m, \sigma'$$

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