

## **DSP & Digital Filters**

#### **Lectures 2-3 Three Different Fourier Transforms**

#### DR TANIA STATHAKI

READER (ASSOCIATE PROFESSOR) IN SIGNAL PROCESSING IMPERIAL COLLEGE LONDON

#### **Three different Fourier Transforms**

There are three useful representations of signals in frequency domain.

- Continuous Time Fourier Transform (CTFT)
  - Continuous aperiodic signals. Continuous time and continuous frequency.
- Discrete Time Fourier Transform (DTFT)
  - Discrete aperiodic signals. Discrete time and continuous frequency.
- Discrete Fourier Transform (DFT)
  - Discrete periodic signals. Discrete Time and discrete frequency.

	Forward Transform	Inverse Transform
CTFT	$X(j\Omega) = \int_{-\infty}^{\infty} x(t)e^{-j\Omega t}dt$ $\Omega$ : "real" frequency	$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\Omega) e^{j\Omega t} d\Omega$
DTFT	$X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}$ $\omega = \Omega T \text{: "normalised" angular frequency}$	$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} X(e^{j\omega}) e^{j\omega n} d\omega$
DFT	$X[k] = \sum_{n=0}^{N-1} x[n]e^{-j2\pi \frac{kn}{N}}$	$x[n] = \frac{1}{N} \sum_{n=0}^{N-1} X[k] e^{j2\pi \frac{kn}{N}}$

#### **Discrete Time Fourier Transform**

• The discrete-time Fourier transform (DTFT)  $X(e^{j\omega})$  of a sequence x[n] is given by

$$X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}$$

• In general  $X(e^{j\omega})$  is a complex function of the real variable  $\omega$  and can be written as

$$X(e^{j\omega}) = X_{\rm re}(e^{j\omega}) + jX_{\rm im}(e^{j\omega})$$

where  $X_{\rm re}(e^{j\omega})$  and  $X_{\rm im}(e^{j\omega})$  are the real and imaginary parts of  $X(e^{j\omega})$  and are real functions of  $\omega$ .

•  $X(e^{j\omega})$  can alternatively be expressed as

$$X(e^{j\omega}) = |X(e^{j\omega})|e^{j\theta(\omega)}$$

where  $|X(e^{j\omega})|$  and  $\theta(\omega)$  are the amplitude and phase of  $X(e^{j\omega})$  and are real functions of  $\omega$  as well.

#### **Discrete Time Fourier Transform**

- For a real sequence x[n],  $|X(e^{j\omega})|$  and  $X_{re}(e^{j\omega})$  are even functions of  $\omega$ , whereas,  $\theta(\omega)$  and  $X_{im}(e^{j\omega})$  are odd functions of  $\omega$ .
- Note that for any integer k

$$X(e^{j\omega}) = |X(e^{j\omega})|e^{j[\theta(\omega) + 2\pi k]} = |X(e^{j\omega})|e^{j\theta(\omega)}$$

- The above property indicates that the phase function  $\theta(\omega)$  cannot be uniquely specified for the DTFT. Recall that the same observation holds for the CTFT.
- Unless otherwise stated, we shall assume that the phase function  $\theta(\omega)$  is restricted to the following range of values:

$$-\pi \Theta \theta(\omega) < \pi$$

called the *principal value*.

#### **Discrete Time Fourier Transform**

- The phase response of DTFT might exhibit discontinuities of  $2\pi$  radians in the plot.
  - [In numerical computations, when the computed phase function is outside the range  $[-\pi,\pi]$ , the phase is computed modulo  $2\pi$  to bring the computed value to the above range.]
- An alternate type of phase function that is a continuous function of  $\omega$  is often used in that case.
- It is derived from the original phase function by removing the discontinuities of  $2\pi$ .
- The process of removing the discontinuities is called *phase* unwrapping.
- Sometimes the continuous phase function generated by unwrapping is denoted as  $\theta_c(\omega)$

#### **Discrete Time Fourier Transform Periodicity**

• Unlike the Continuous Time Fourier Transform, the DTFT is a periodic function in  $\omega$  with period  $2\pi$ .

$$X\left(e^{j(\omega_0+2\pi k)}\right) = \sum_{n=-\infty}^{\infty} x[n]e^{-j(\omega_0+2\pi k)n} = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega_0n} e^{-j2\pi kn} = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega_0n} = X(e^{j\omega_0}), \text{ for any integer } k.$$

- Therefore,  $X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}$  imitates a Fourier Series representation of the periodic function  $X(e^{j\omega})$ .
- As a result, the Fourier Series coefficients x[n] can be derived from  $X(e^{j\omega})$  using the Fourier integral

$$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} X(e^{j\omega})e^{j\omega n} d\omega$$
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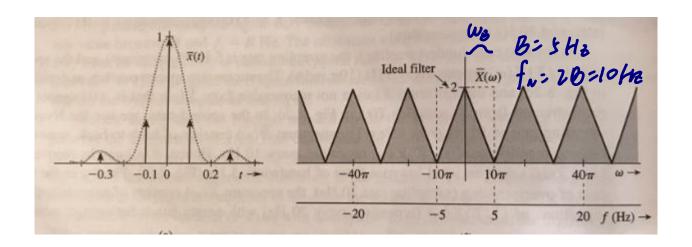
called the Inverse DTFT (IDTFT).

 Periodicity of DTFT is not a new concept; we know from sampling theory, that sampling a continuous signal results in a periodic repetition of its CTFT. Imperial College London

#### **Revision**

## **Nyquist sampling: Just about the correct sampling rate**

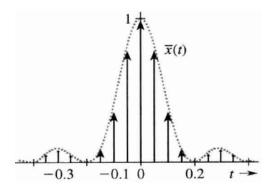
- In that case we use the Nyquist sampling rate of 10Hz.
- The spectrum X̄(ω) consists of back-to-back, non-overlapping repetitions of <sup>1</sup>/<sub>T<sub>s</sub></sub> X(ω) repeating every 10Hz.
   In order to recover X(ω) from X̄(ω) we must use an ideal lowpass filter
- In order to recover  $X(\omega)$  from  $X(\omega)$  we must use an ideal lowpass filter of bandwidth 5Hz. This is shown in the right figure below with the dotted line.

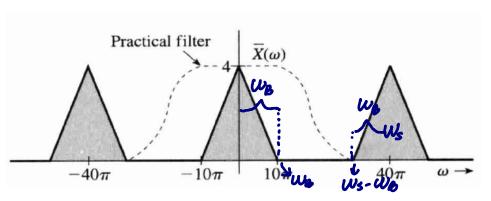


#### **Revision**

## **Oversampling: What happens if we sample too quickly?**

- Sampling at higher than the Nyquist rate (in this case 20Hz) makes reconstruction easier.
- The spectrum  $\bar{X}(\omega)$  consists of non-overlapping repetitions of  $\frac{1}{T_s}X(\omega)$ , repeating every 20Hz with empty bands between successive cycles.
- In order to recover  $X(\omega)$  from  $\bar{X}(\omega)$  we can use a practical lowpass filter and not necessarily an ideal one. This is shown in the right figure below with the dotted line.

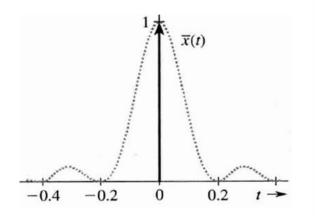


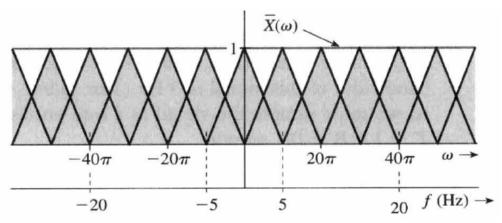


• The filter we use for reconstruction must have gain  $T_s$  and bandwidth of any value between B and  $(f_s - B)Hz$ .

# Revision Undersampling: What happens if we sample too slowly?

- Sampling at lower than the Nyquist rate (in this case 5Hz) makes reconstruction impossible.
- The spectrum  $\bar{X}(\omega)$  consists of overlapping repetitions of  $\frac{1}{T_s}X(\omega)$  repeating every 5Hz.
- $X(\omega)$  is not recoverable from  $\bar{X}(\omega)$ .
- Sampling below the Nyquist rate corrupts the signal. This type of distortion is called <u>aliasing</u>.





## **More DTFT Properties**

- The DTFT is the z –transform evaluated at  $z = e^{j\omega}$ .
  - [Recall that  $X(z) = \sum_{-\infty}^{\infty} x[n] z^{-n}$ ].
  - Therefore, the DTFT converges if the ROC includes |z| = 1 ( $z = e^{j\omega}$ ).
- The DTFT is the same as the CTFT of a signal comprising impulses of appropriate heights at the sample instances.

$$x_{\delta}(t) = \sum_{n} x[n]\delta(t - nT) = x(t)\sum_{-\infty}^{\infty} \delta(t - nT)$$

• Recall that x[n] = x(nT)

$$X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n} = \sum_{n=-\infty}^{\infty} x[n] \int_{-\infty}^{\infty} \delta(t-nT)e^{-j\omega \frac{t}{T}} dt$$

$$= \int_{-\infty}^{\infty} \left[\sum_{n=-\infty}^{\infty} x[n] \, \delta(t-nT)\right] e^{-j\omega \frac{t}{T}} dt = \int_{-\infty}^{\infty} x_{\delta}(t) e^{-j\Omega t} dt$$

- For the above the condition  $\sum_{n=-\infty}^{\infty} |x[n]| < \infty$  must hold.
- $\omega = \Omega T$

#### **Examples**

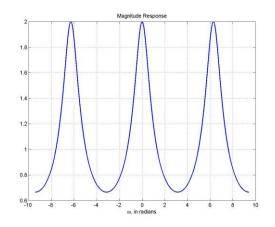
• The DTFT of a shifted discrete Dirac function  $\delta[n-k]$  is given by:

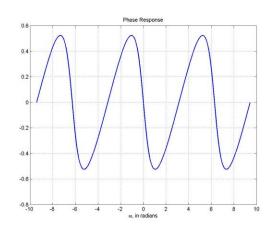
$$\Delta(\omega) = \sum_{n=-\infty}^{\infty} \delta[n-k]e^{-j\omega n} = e^{-j\omega k}$$

• The DTFT of the causal sequence  $x[n] = \alpha^n u[n]$ ,  $|\alpha| < 1$  is given by:

$$X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} \alpha^n u[n] e^{-j\omega n} = \sum_{n=0}^{\infty} \alpha^n e^{-j\omega n} = \sum_{n=0}^{\infty} (\alpha e^{-j\omega})^n = \frac{1}{1-\alpha e^{-j\omega}}$$
 if  $|\alpha e^{-j\omega}| = |\alpha| < 1$ 

• For  $\alpha=0.5$ , the magnitude and phase of  $X(e^{j\omega})=1/(1-0.5e^{-j\omega})$  are shown below.





Imperial College DT1: 
$$\chi(e^{jw}) = \sum_{k=-\infty}^{\infty} \chi(k)e^{jwk}$$
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IDF1:  $\chi(k) = \sum_{k=-\infty}^{\infty} \chi(e^{jw})e^{jwk} dw$ 

#### **Inverse Discrete Time Fourier Transform**

$$X[n] = \frac{1}{2\pi} \int_{-\pi}^{4\pi} \chi(e^{jw}) e^{jwn} dw = \frac{1}{2\pi} \int_{-\pi}^{\pi} \sum_{l=-\infty}^{\infty} X[l] e^{-jwl} e^{jwn} dw = \sum_{l=-\infty}^{\infty} \chi[l] \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{jw(n-l)} dw$$

$$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} \left( \sum_{\ell=-\infty}^{\infty} x[\ell] e^{j\omega n} d\omega \right) = \sum_{\ell=-\infty}^{\infty} x[\ell] \int_{-\pi}^{\pi} e^{j\omega(n-\ell)} d\omega$$

$$= \sum_{\ell=-\infty}^{\infty} x[\ell] \left( \sum_{\ell=-\infty}^{\infty} x[\ell] e^{j\omega(n-\ell)} d\omega \right) = \sum_{\ell=-\infty}^{\infty} x[\ell] \frac{\sin \pi (n-\ell)}{\pi (n-\ell)}$$

$$= \sum_{\ell=-\infty}^{\infty} x[\ell] \left( \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{j\omega(n-\ell)} d\omega \right) = \sum_{\ell=-\infty}^{\infty} x[\ell] \frac{\sin \pi (n-\ell)}{\pi (n-\ell)}$$

(Note that the order of integration and summation can be interchanged if the summation inside the top brackets converges uniformly, i.e., if  $X(e^{j\omega})$ exists.)

#### **Inverse Discrete Time Fourier Transform cont.**

$$x[\ell] \frac{\sin \pi (n-\ell)}{\pi (n-\ell)} = \begin{cases} 1 & n=\ell\\ 0 & n \neq \ell \end{cases}$$

Hence,

$$\sum_{\ell=-\infty}^{\infty} x[\ell] \frac{\sin \pi (n-\ell)}{\pi (n-\ell)} = \sum_{\ell=-\infty}^{\infty} x[\ell] \delta[n-\ell] = x[n]$$

## Discrete Time Fourier Transform: uniform convergence

- An infinite series of the form  $X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}$  may or may not converge.
- Let  $X_K(e^{j\omega}) = \sum_{n=-K}^K x[n]e^{-j\omega n}$
- For uniform convergence (strong convergence) of  $X(e^{j\omega})$  we require:  $\lim_{K\to\infty} X_K(e^{j\omega}) = X(e^{j\omega})$
- If x[n] is an **absolutely summable** sequence, i.e., if  $\sum_{n=-\infty}^{\infty} |x[n]| < \infty$ , then

$$\left|X(e^{j\omega})\right| = \left|\sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}\right| \le \sum_{n=-\infty}^{\infty} |x[n]| \left|e^{-j\omega n}\right| \le \sum_{n=-\infty}^{\infty} |x[n]| < \infty$$

for all values of  $\omega$ 

• Thus, the absolute summability of x[n] is a sufficient condition for the existence of the DTFT  $X(e^{j\omega})$ .

## **Examples**

The sequence  $x[n] = \alpha^n u[n]$  is absolutely summable for  $|\alpha| < 1$  since

$$\sum_{n=-\infty}^{\infty} |\alpha^n| u[n] = \sum_{n=0}^{\infty} |\alpha^n| = \frac{1}{1-|\alpha|} < \infty$$

- - sequence has always finite energy.
  - However, a finite energy sequence is not necessarily absolutely summable.
    - The sequence  $x[n] = \begin{cases} 1/n & n \ge 1 \\ 0 & n < 0 \end{cases}$

has finite energy equal to  $\sum_{n=1}^{\infty} (\frac{1}{n})^2 = \pi^2/6$  but is not absolutely summable.

#### Discrete Time Fourier Transform: mean square convergence

• To represent a finite energy sequence x[n] that is not absolutely summable by DTFT, it is necessary to consider the so called *mean-square* **convergence** (weak convergence) of  $X(e^{j\omega})$ :

$$\lim_{K \to \infty} \int_{-\pi}^{n} \left| X(e^{j\omega}) - X_K(e^{j\omega}) \right|^2 d\omega = 0$$

where  $X_K(e^{j\omega}) = \sum_{n=-K}^K x[n]e^{-j\omega n}$ .

- Here, the total energy of the error  $X(e^{j\omega}) X_K(e^{j\omega})$  must approach zero at each value of  $\omega$  as K goes to  $\infty$ .
- In such a case, the absolute value of the error may not go to zero as

K goes to 
$$\infty$$
 and the DTFT is no longer bounded.

Uniform convergence

 $\lim_{k\to\infty} |\chi(e^{i\omega}) - \chi_k(e^{i\omega})| = 0$  (error  $\to 0$ )

mean-square convergence

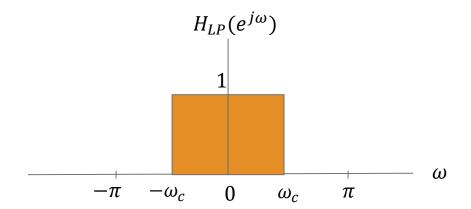
 $\lim_{k\to\infty} |\chi(e^{i\omega}) - \chi_k(e^{i\omega})|^2 d\omega = 0$  (error energy  $\to 0$ )

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#### **Example**

Consider the DTFT:

$$H_{LP}(e^{j\omega}) = \begin{cases} 1, & 0 \le |\omega| \le \omega_c \\ 0, & \omega_c < |\omega| \le \pi \end{cases}$$



The inverse DTFT is given by

$$h_{LP}[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} H_{LP}(e^{j\omega}) e^{j\omega n} d\omega = \frac{1}{2\pi} \int_{-\omega_c}^{\omega_c} e^{j\omega n} d\omega$$

$$=\frac{1}{2\pi}\left(\frac{e^{j\omega_{c}n}}{jn}-\frac{e^{-j\omega_{c}n}}{jn}\right)=\frac{\sin\omega_{c}n}{\pi n},\,-\infty< n<\infty$$

- The energy of  $h_{LP}[n]$  is given by  $E_h = \frac{1}{2\pi} \int_{-\pi}^{\pi} |H_{LP}(e^{j\omega})|^2 d\omega = \frac{\omega_c}{\pi}$ .
- $h_{LP}[n]$  is a finite-energy sequence, but it is not absolutely summable.

## **Example cont.**

As a result

$$\sum_{n=-K}^{K} h_{LP}[n]e^{-j\omega n} = \sum_{n=-K}^{K} \frac{\sin \omega_c n}{\pi n} e^{-j\omega n}$$

does not uniformly converge to 
$$H_{LP}\left(e^{j\omega}\right) = \begin{cases} 1, & \text{converge} \\ 1, & 0 \leq |\omega| \leq \omega_c \\ 0, & \omega_c < |\omega| \leq \pi \end{cases}$$

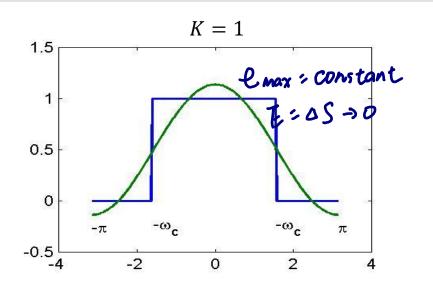
for all values of  $\omega$ , but converges to  $H_{LP}(e^{j\omega})$  in the mean-square sense.

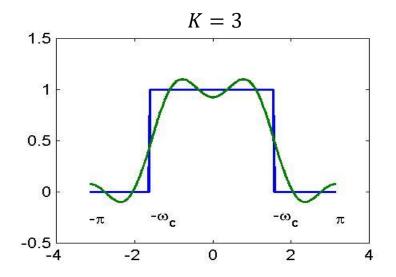
The mean-square convergence property of the sequence  $h_{LP}[n]$  can be further illustrated by examining the plot of the function

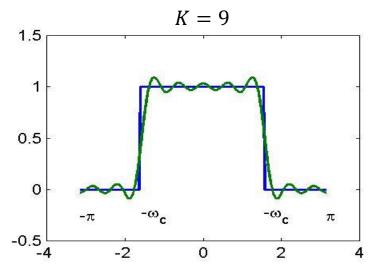
$$H_{LP,K}(e^{j\omega}) = \sum_{n=-K}^{K} \frac{\sin \omega_c n}{\pi n} e^{-j\omega n}$$

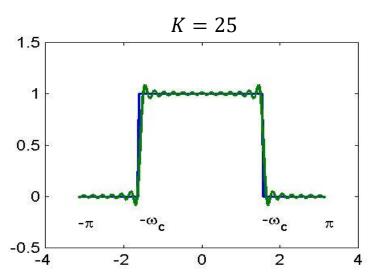
for various values of K as shown next.

## **Example cont.**









## **Example cont.**

- As it can be seen from these plots, independent of the value of K there are ripples in the plot of  $H_{LP,K}(e^{j\omega})$  around both sides of the point  $\omega = \omega_c$ .
- The number of ripples increases as *K* increases with the height of the largest ripple remaining the same for all values of *K*.
- As *K* goes to infinity, the condition

$$\lim_{K \to \infty} \int_{-\pi}^{K} \left| H_{LP}(e^{j\omega}) - H_{LP,K}(e^{j\omega}) \right|^{2} d\omega = 0$$

holds, indicating the convergence of  $H_{LP,K}(e^{j\omega})$  to  $H_{LP}(e^{j\omega})$ .

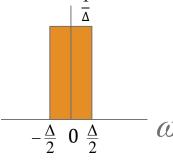
• The oscillatory behavior observed in  $H_{LP,K}(e^{j\omega})$  is known as the **Gibbs phenomenon**.

#### **Neither absolutely- nor square- summable**

- The DTFT can also be defined for a certain class of sequences which are neither absolutely summable nor square summable.
- Examples of such sequences are the unit step sequence u[n], the sinusoidal sequence  $\cos(\omega_o n + \varphi)$  and the complex exponential sequence  $A\alpha^n$ . These are neither absolutely summable nor square summable.
- For this type of sequences, a DTFT representation is possible using Dirac delta functions.
- A *Dirac delta function*  $\delta(\omega)$  is a "function" of  $\omega$  with infinite height, zero width, and unit area.
- It is the limiting form of a unit area pulse function  $p_{\Delta}(\omega)$  as  $\Delta$  goes to zero

$$\delta(\omega) = \lim_{\Delta \to 0} p_{\Delta}(\omega) \qquad \qquad p_{\Delta}(\omega)$$
 satisfying

$$\int_{-\infty}^{\infty} p_{\Delta}(\omega) d\omega = 1, \, p_{\Delta}(\omega) = 0, \, \omega \neq 0$$



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$$\chi[n] := \frac{\sum_{k=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} \sum_{k=$$

• Consider the complex exponential sequence  $x[n] = e^{j\omega_0 n}$ ,  $\omega_0$  real. Its DTFT is given by

$$X(e^{j\omega}) = \sum_{k=-\infty}^{\infty} 2\pi\delta(\omega - \omega_o + 2\pi k)$$

where  $\delta(\omega)$  is an impulse function of  $\omega$  and  $-\pi \leq \omega_o \leq \pi$ .

• To verify the above we can take the IDTFT of  $X(e^{j\omega})$  above:

$$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} \sum_{k=-\infty}^{\infty} 2\pi \delta(\omega - \omega_o + 2\pi k) e^{j\omega n} d\omega = \int_{-\pi}^{\pi} \delta(\omega - \omega_o) e^{j\omega n} d\omega = e^{j\omega_o n}$$

## **DTFT** properties (listed without proof)

<b>Type of Property</b>	Sequence	Discrete-Time Fourier Transform
	g[n] $h[n]$	$G(e^{j\omega}) \ H(e^{j\omega})$
Linearity	$\alpha g[n] + \beta h[n]$	$\alpha G(e^{j\omega}) + \beta H(e^{j\omega})$
Time-shifting	$g[n-n_o]$	$e^{-j\omega n_o}G(e^{j\omega})$
Frequency-shifting	$e^{j\omega_o n}g[n]$	$G\left(e^{j(\omega-\omega_o)}\right)$
Differentiation in frequency	ng[n]	$j\frac{dG(e^{j\omega})}{d\omega}$
Convolution	$g[n] \circledast h[n]$	$G(e^{j\omega})H(e^{j\omega})$
Modulation	g[n]h[n]	$\frac{1}{2\pi} \int_{-\pi}^{\pi} G(e^{j\theta}) H(e^{j(\omega-\theta)}) d\theta$
Parseval's relation	$\sum_{n=-\infty}^{\infty} g[n]h^*[$	$[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} G(e^{j\omega}) H^*(e^{j\omega}) d\omega$

#### **DTFT** properties (listed without proof)

Sequence	Discrete-Time Fourier Transform	
x[n]	$X(e^{j\omega})$ $x[n]: A$	– complex sequence
x[-n]	$X(e^{-j\omega})$	
$x^*[-n]$	$X^*(e^{j\omega})$	
$Re\{x[n]\}$	$X_{\rm cs}(e^{j\omega}) = \frac{1}{2} \{ X(e^{j\omega}) + X^*(e^{-j\omega}) \}$	
$j\operatorname{Im}\{x[n]\}$	$X_{\mathrm{ca}}(e^{j\omega}) = \frac{1}{2} \{ X(e^{j\omega}) - X^*(e^{-j\omega}) \}$	
$x_{cs}[n]$	$X_{\mathrm{re}}(e^{j\omega})$	
$x_{ca}[n]$	$jX_{\mathrm{im}}(e^{j\omega})$	

Note:  $X_{cs}(e^{j\omega})$  and  $X_{ca}(e^{j\omega})$  are the conjugate-symmetric and conjugate-antisymmetric parts of  $X(e^{j\omega})$ , respectively. Likewise,  $x_{cs}[n]$  and  $x_{ca}[n]$  are the conjugate-symmetric and conjugate-antisymmetric parts of x[n], respectively.



## **DTFT** properties (listed without proof)

		_
Sequence	Discrete-Time Fourier Transform	
x[n]	$X(e^{j\omega}) = X_{\rm re}(e^{j\omega}) + jX_{\rm im}(e^{j\omega})$	x[n]: A real sequence
$x_{\text{ev}}[n]$ $x_{\text{od}}[n]$	$X_{\rm re}(e^{j\omega})$ $jX_{\rm im}(e^{j\omega})$	
		-
	$X(e^{j\omega}) = X^*(e^{-j\omega})$ $X_{re}(e^{j\omega}) = X_{re}(e^{-j\omega})$	
Symmetry relations	$X_{\rm im}(e^{j\omega}) = -X_{\rm im}(e^{-j\omega})$	
	$ X(e^{j\omega})  =  X(e^{-j\omega}) $ $\arg\{X(e^{j\omega})\} = -\arg\{X(e^{-j\omega})\}$	
	$m_{\theta}(x_{i}(c_{i})) = m_{\theta}(x_{i}(c_{i}))$	

Note:  $x_{ev}[n]$  and  $x_{od}[n]$  denote the even and odd parts of x[n], respectively.

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#### **Common DTFT pairs**

$$\delta[n] \leftrightarrow 1$$

$$1 \leftrightarrow \sum_{k=-\infty}^{\infty} 2\pi\delta(\omega + 2\pi k)$$

$$u[n] \leftrightarrow \frac{1}{1 - e^{-j\omega}} + \sum_{k=-\infty}^{\infty} \pi\delta(\omega + 2\pi k)$$

$$e^{j\omega_{o}n} \leftrightarrow \sum_{k=-\infty}^{\infty} 2\pi\delta(\omega - \omega_{o} + 2\pi k)$$

$$\alpha^{n}u[n], (|\alpha| < 1) \leftrightarrow \frac{1}{1 - \alpha e^{-j\omega}}$$

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$$hu[n] \longleftrightarrow j \frac{d}{dw} \chi(e^{jw}) = \frac{1}{1-\alpha e^{jw}}$$

$$j \frac{d}{dw} \chi(e^{jw}) = j \frac{-\alpha j e^{jw}}{(1-\alpha e^{jw})^2} = \frac{\alpha e^{-jw}}{(1-\alpha e^{jw})^2}$$
Example

Determine the DTFT of the sequence

$$y[n] = (n+1)\alpha^n u[n], |\alpha| < 1$$

- Let  $x[n] = \alpha^n u[n]$ ,  $|\alpha| < 1$ . We can, therefore, write y[n] = nx[n] + x[n]
- From tables, the DTFT of x[n] is given by

$$X(e^{j\omega}) = \frac{1}{1 - \alpha e^{-j\omega}}$$

• Using the differentiation property of the DTFT given in previous tables, we observe that the DTFT of nx[n] is given by

$$j\frac{dX(e^{j\omega})}{d\omega} = j\frac{d}{d\omega}\left(\frac{1}{1 - \alpha e^{-j\omega}}\right) = \frac{\alpha e^{-j\omega}}{(1 - \alpha e^{-j\omega})^2}$$

 Next using the linearity property of the DTFT given in previous tables we arrive at

$$Y(e^{j\omega}) = \frac{\alpha e^{-j\omega}}{(1 - \alpha e^{-j\omega})^2} + \frac{1}{1 - \alpha e^{-j\omega}} = \frac{1}{(1 - \alpha e^{-j\omega})^2}$$

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$$V[n-1] \leftrightarrow e^{\int_{0}^{w} V(e^{jw})}$$

$$d_{0}V(e^{jw}) + \alpha_{1}e^{-jw}V(e^{jw}) = P_{0} + P_{1}e^{-jw}$$

$$V(e^{jw}) = \frac{P_{0} + P_{1}e^{-jw}}{\alpha_{0} + \alpha_{1}e^{-jw}}$$

## **Example**

• Determine the DTFT of the sequence v[n] defined by

$$d_0v[n] + d_1v[n-1] = p_0\delta[n] + p_1\delta[n-1], |d_1/d_0| < 1$$

- From previous tables, we see that the DTFT of  $\delta[n]$  is 1.
- Using the time-shifting property of the DTFT given in previous tables, we observe that the DTFT of  $\delta[n-1]$  is  $e^{-j\omega}V(e^{j\omega})$ .
- Using the linearity property of previous tables we then obtain the frequency-domain representation of  $d_0v[n]+d_1v[n-1]=p_0\delta[n]+p_1\delta[n-1]$  as

$$d_0V(e^{j\omega})+d_1e^{-j\omega}V(e^{j\omega})=p_0+p_1e^{-j\omega}$$

Solving the above equation we get

$$V(e^{j\omega}) = \frac{p_0 + p_1 e^{-j\omega}}{d_0 + d_1 e^{-j\omega}}$$

#### **Energy Density Spectrum**

• The total energy of a finite-energy sequence g[n] is given by

$$\varepsilon_g = \sum_{n=-\infty}^{\infty} |g[n]|^2$$

From Parseval's Theorem we know that

$$\varepsilon_g = \sum_{n=-\infty}^{\infty} |g[n]|^2 = \frac{1}{2\pi} \int_{-\pi}^{\pi} |G(e^{j\omega})|^2 d\omega$$

- The quantity  $S_{gg}(\omega) = \left| G(e^{j\omega}) \right|^2$  is called the **energy density spectrum**.
- The area under this curve in the range  $-\pi \le \omega \le \pi$  divided by  $2\pi$  is the energy of the sequence.

#### **Example**

Compute the energy of the sequence

$$h_{LP}[n] = \frac{\sin \omega_c n}{\pi n}, -\infty < n < \infty$$

Here,

$$\sum_{n=-\infty}^{\infty} |h_{LP}[n]|^2 = \frac{1}{2\pi} \int_{-\pi}^{\pi} |H_{LP}(e^{j\omega})|^2 d\omega$$

$$H_{LP}(e^{j\omega}) = \begin{cases} 1, & 0 \le |\omega| \le \omega_c \\ 0, & \omega_c < |\omega| \le \pi \end{cases}$$

Therefore,

$$\sum_{n=-\infty}^{\infty} |h_{LP}[n]|^2 = \frac{1}{2\pi} \int_{-\omega_c}^{\omega_c} d\omega = \frac{\omega_c}{\pi} < \infty$$

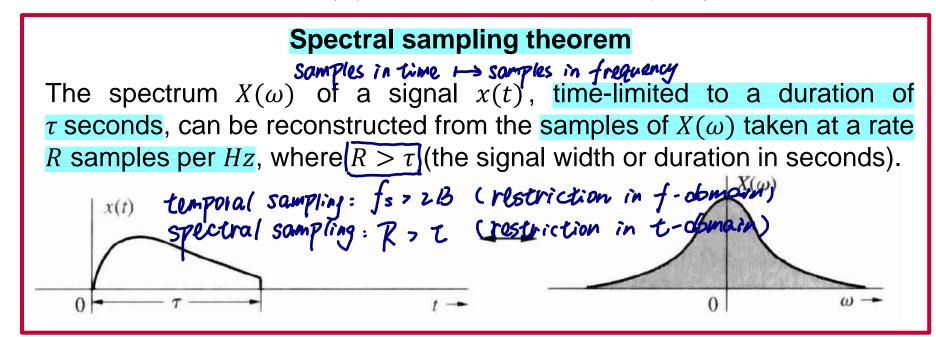
• Hence,  $h_{LP}[n]$  is a finite energy sequence.

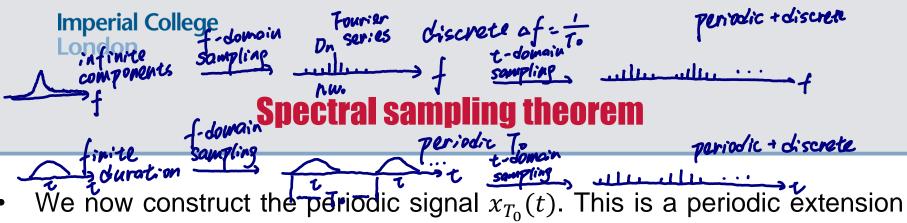
## **Introduction. Time sampling theorem resume.**

- We wish to perform spectral analysis using digital computers.
- Therefore, we must somehow sample the Discrete Time Fourier Transform of the signal!
- We will compute a discrete version of the DTFT of a <u>sampled</u>, <u>finite-duration</u> signal. This transform is known as the Discrete Fourier Transform (DFT).
- The goal is to understand how DFT is related to the original Fourier transform.
- We showed that a signal bandlimited to BHz can be reconstructed from signal samples if they are obtained at a rate of  $f_s > 2B$  samples per second.
- Not that the signal spectrum exists over the frequency range (in Hz) from -B to B.
- The interval 2B is called **spectral width**. Note the difference between spectral width (2B) and bandwidth (B).
- In time sampling theorem:  $f_s > 2B$  or  $f_s >$  (spectral width).

## Time sampling theorem has a dual: Spectral sampling theorem

- Consider a time-limited signal x(t) with a spectrum  $X(\omega)$ .
- In general, a time-limited signal is 0 for  $t < T_1$  and  $t > T_2$ . The duration of the signal is  $\tau = T_2 T_1$ . Below we assume that  $T_1 = 0$ .
- Recall that  $X(\omega) = \int_{-\infty}^{\infty} x(t)e^{-j\omega t}dt = \int_{0}^{\tau} x(t)e^{-j\omega t}dt$ .
- The Fourier transform  $X(\omega)$  is assumed real for simplicity.





- of x(t) with period  $T_0 > \tau$ .
- This periodic signal can be expressed using Fourier series.

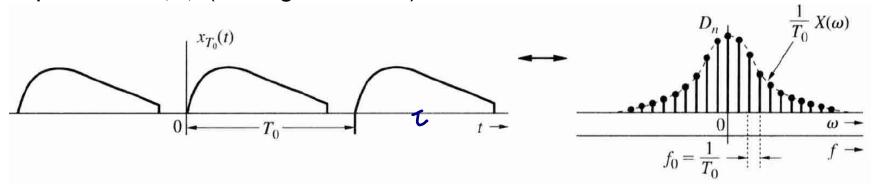
$$x_{T_0}(t) = \sum_{n=-\infty}^{n=\infty} D_n e^{jn\omega_0 t}, \ \omega_0 = \frac{2\pi}{T_0}$$

$$D_n = \frac{1}{T_0} \int_0^{T_0} x(t) \ e^{-jn\omega_0 t} dt = \frac{1}{T_0} \int_0^{\tau} x(t) \ e^{-jn\omega_0 t} dt = \frac{1}{T_0} X(n\omega_0)$$

- The result indicates that the coefficients of the Fourier series for  $x_{T_0}(t)$ are the values of  $X(\omega)$  taken at integer multiples of  $\omega_0$  and scaled by
- We call **spectrum of a periodic signal** the weights of the exponential terms in its Fourier series representation.
- The above implies that the spectrum of the periodic signal  $x_{T_0}(t)$  is the sampled version of spectrum  $X(\omega)$ .

## **Spectral sampling theorem cont.**

The spectrum of the periodic signal  $x_{T_0}(t)$  is the sampled version of spectrum  $X(\omega)$  (see figure below).



- If successive cycles of  $x_{T_0}(t)$  do not overlap, x(t) can be recovered from  $x_{T_0}(t)$ .
- If we know x(t) we can find  $X(\omega)$ .
- The above imply that  $X(\omega)$  can be reconstructed from its samples.
- These samples are separated by the so called fundamental frequency  $f_0 = \frac{1}{T_0} Hz$  or  $\omega_0 = 2\pi f_0 rads/s$  of the periodic signal  $x_{T_0}(t)$ . Therefore, the condition for recovery is  $T_0 > \tau \Rightarrow f_0 < \frac{1}{\tau} Hz$ .

#### **Spectral interpolation formula**

• To reconstruct the spectrum  $X(\omega)$  from the samples of  $X(\omega)$ , the samples should be taken at frequency intervals  $f_0 < \frac{1}{\tau}Hz$ . If the sampling rate is R frequency samples/Hz we have:

$$R = \frac{1}{f_0} > \tau$$
 samples/ $Hz$ 

 We know that the continuous version of a signal can be recovered from its sampled version through the so called signal interpolation formula: (refer to a Signals and Systems book for the proof of it)

$$x(t) = \sum_{n} x(nT_s)h(t - nT_s) = \sum_{n} x(nT_s)\operatorname{sinc}\left(\frac{\pi t}{T_s} - n\pi\right)$$

We use the dual of the approach employed to derive the signal interpolation formula above, to obtain the **spectral interpolation formula** as follows. We assume that x(t) is time-limited to  $\tau$  and centred at  $T_c$ . We can prove that:

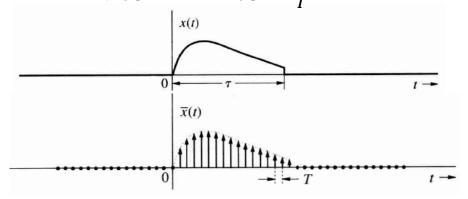
$$X(\omega) = \sum_{n=-\infty} X(n\omega_0) \operatorname{sinc}\left(\frac{\omega T_0}{2} - n\pi\right) e^{-j(\omega - n\omega_0)T_C}, \ \omega_0 = \frac{2\pi}{T_0}, \ T_0 > \tau$$

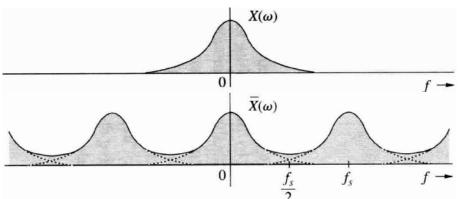
## **Spectral interpolation formula: Proof.**

- We know that  $x_{T_0}(t) = \sum_{n=-\infty}^{n=\infty} D_n e^{jn\omega_0 t}$ ,  $\omega_0 = \frac{2\pi}{T_0}$
- Therefore,  $\mathcal{F}\{x_{T_0}(t)\} = 2\pi \sum_{n=-\infty}^{n=\infty} D_n \, \delta(\omega n\omega_0)$ [It is easier to prove that  $\mathcal{F}^{-1}\{\sum_{n=-\infty}^{n=\infty} D_n \, \delta(\omega - n\omega_0)\} = x_{T_0}(t)$ ]
- We can write  $x(t) = x_{T_0}(t) \cdot \text{rect}\left(\frac{t T_c}{T_0}\right)$  (1) [We were given that x(t) is centred at  $T_c$ ]
- We know that  $\mathcal{F}\left\{\operatorname{rect}\left(\frac{t}{T_0}\right)\right\} = T_0\operatorname{sinc}\left(\frac{\omega T_0}{2}\right)$ .
- Therefore,  $\mathcal{F}\left\{\operatorname{rect}\left(\frac{t-T_c}{T_0}\right)\right\} = T_0\operatorname{sinc}\left(\frac{\omega T_0}{2}\right)e^{-j\omega T_c}$ .
- From (1) we see that  $X(\omega) = \frac{1}{2\pi} \mathcal{F}\{x_{T_0}(t)\} * \mathcal{F}\left\{\operatorname{rect}\left(\frac{t-T_c}{T_0}\right)\right\}$
- $X(\omega) = \frac{1}{2\pi} 2\pi \left[\sum_{n=-\infty}^{n=\infty} D_n \,\delta(\omega n\omega_0)\right] * T_0 \operatorname{sinc}\left(\frac{\omega T_0}{2}\right) e^{-j\omega T_C}$   $X(\omega) = \sum_{n=-\infty} D_n T_0 \operatorname{sinc}\left[\frac{(\omega n\omega_0)T_0}{2}\right] e^{-j(\omega n\omega_0)T_C}, \,\omega_0 = \frac{2\pi}{T_0}, \,T_0 > \tau$   $X(\omega) = \sum_{n=-\infty} X(n\omega_0) \operatorname{sinc}\left(\frac{\omega T_0}{2} n\pi\right) e^{-j(\omega n\omega_0)T_C}$

#### **Discrete Fourier Transform DFT**

- The numerical computation of the Fourier transform requires samples of x(t) since computers can work only with discrete values.
- Furthermore, the Fourier transform can only be computed at some discrete values of  $\omega$ .
- The goal of what follows is to relate the samples of  $X(\omega)$  with the samples of x(t).
- Consider a time-limited signal x(t). Its spectrum  $X(\omega)$  will not be bandlimited (try to think why). In other words aliasing after sampling cannot be avoided. (Loil effect)
- The spectrum  $\bar{X}(\omega)$  of the sampled signal  $\bar{x}(t)$  consist of  $X(\omega)$  repeating every  $f_SHz$  with  $f_S=\frac{1}{T}$ .

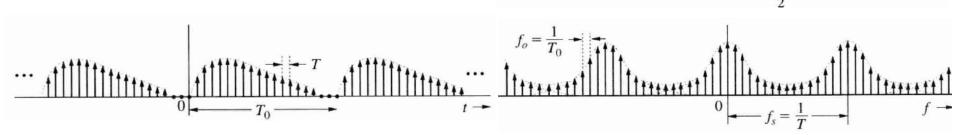




#### **Discrete Fourier Transform DFT cont.**

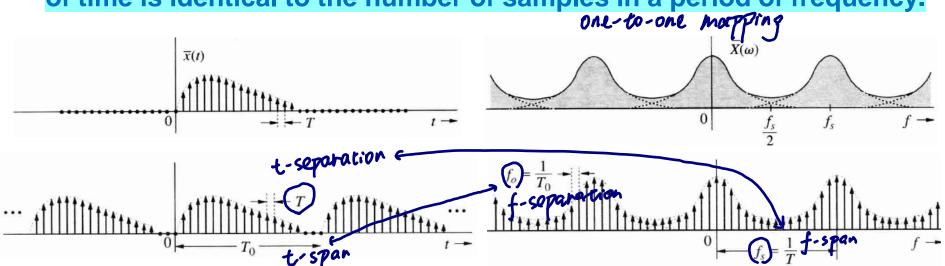
- Suppose now that the sampled signal  $\bar{x}(t)$  is repeated periodically every  $T_0$  seconds.
- According to the spectral sampling theorem, this operation results in sampling the spectrum at a rate of  $T_0$  samples/Hz. This means that the samples are spaced at  $f_0 = \frac{1}{T_0}Hz$ .
- Therefore, when a signal is sampled and periodically repeated, its spectrum is also sampled and periodically repeated.

• The goal of what follows is to relate the samples of  $X(\omega)$  with the samples of x(t).



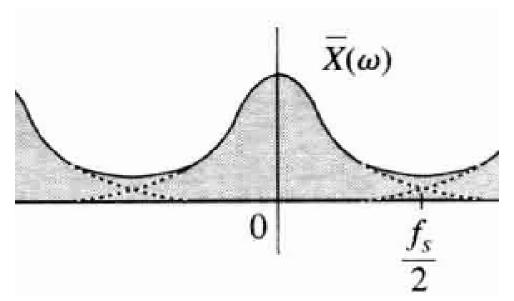
#### **Discrete Fourier Transform DFT cont.**

- The number of samples of the discrete signal in one period  $T_0$  is  $N_0 = \frac{T_0}{T}$  (figure below left).
- The number of samples of the discrete spectrum in one period is  $N_0' = \frac{f_s}{f_0}$ .
- We see that  $N_0' = \frac{f_s}{f_0} = \frac{\frac{1}{T}}{\frac{1}{T_0}} = \frac{T_0}{T} = N_0$ .
- This is an interesting observation: the number of samples in a period of time is identical to the number of samples in a period of frequency.



# **Aliasing and leakage effects**

• Since  $X(\omega)$  is not bandlimited, we will get some aliasing effect:



• Furthermore, if x(t) is not time limited, we need to truncate x(t) with a window function. This leads to a "leakage" effect (refer to a Signals and Systems book for the demonstration of it).

#### **Formal definition of DFT**

• If x(nT) and  $X(k\omega_0)$  are the  $n^{\text{th}}$  and  $k^{\text{th}}$  samples of x(t) and  $X(\omega)$  respectively, we define:

$$x[n] = \widehat{T}x(nT) = \frac{T_0}{N_0}x(nT)$$

$$X[k] = X(k\omega_0), \, \omega_0 = 2\pi f_0 = \frac{2\pi}{T_0}$$

It can be shown that x[n] and X[k] are related by the following equations:

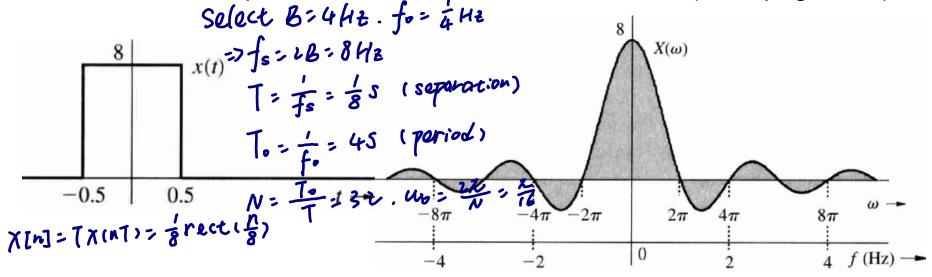
$$X[k] = \sum_{n=0}^{N_0 - 1} x[n] e^{-jnk\Omega_0}$$
(1)

$$x[n] = \frac{1}{N_0} \sum_{k=0}^{N_0 - 1} X[k] e^{jkn\Omega_0} , \Omega_0 = \omega_0 T = \frac{2\pi}{N_0}$$
 (2)

- The equations (1) and (2) above are the direct and inverse Discrete Fourier Transforms respectively, known as DFT and IDFT.
- In the above equations, the summation is performed from 0 to  $N_0 1$ . It can be shown that the summation can be performed over any successive  $N_0$  values of n or k.

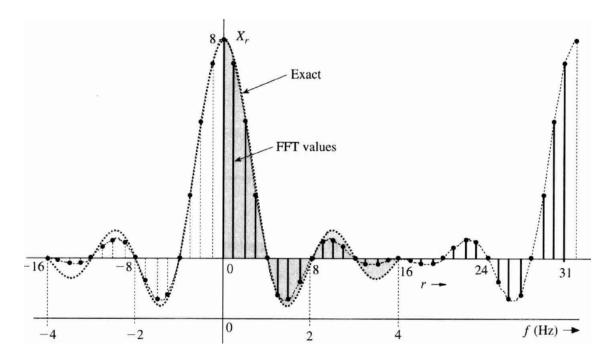
# **Example**

• Use DFT to compute the Fourier transform of 8rect(t) (Lathi page 808.)

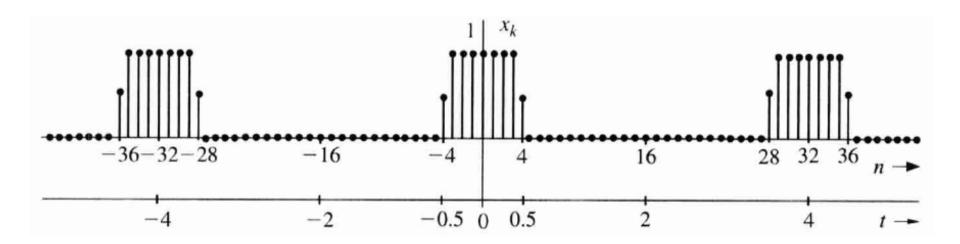


- The essential bandwidth B (calculated by finding where the amplitude response drops to 1% of its peak value) is well above 16Hz. However, we select B = 4Hz:
  - To observe the effects of aliasing.
  - In order not to end up with a huge number of samples in time.

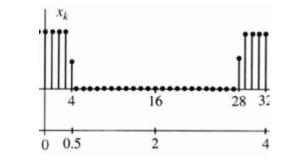
- B = 4Hz,  $f_S = 8Hz$ ,  $T = \frac{1}{f_S} = \frac{1}{8}$ .
- For the frequency resolution we choose  $f_0 = \frac{1}{4}Hz$ . This choice gives us 4 samples in each lobe of  $X(\omega)$  and  $T_0 = \frac{1}{f_0} = 4s$ .



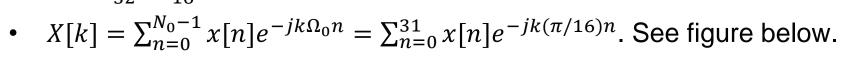
- $N_0 = \frac{T_0}{T} = \frac{4}{1/8} = 32$ . Therefore, we must repeat x(t) every 4s and take samples every  $\frac{1}{8}s$ . This yields 32 samples in a period.
- $x[n] = Tx(nT) = \frac{1}{8}x(\frac{n}{8})$  with x(t) = 8rect(t).
- The DFT of the signal x[n] is obtained by taking any full period of x[n] (i.e.,  $N_0$  samples) and not necessarily  $N_0$  over the interval  $(0, T_0)$  as we assumed in the theoretical analysis of DFT.

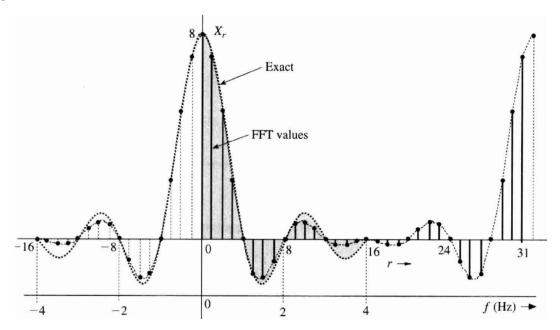


• 
$$x[n] = \begin{cases} 1 & 0 \le n \le 3 & \text{and} & 29 \le n \le 31 \\ 0 & 5 \le n \le 27 \\ 0.5 & n = 4,28 \end{cases}$$

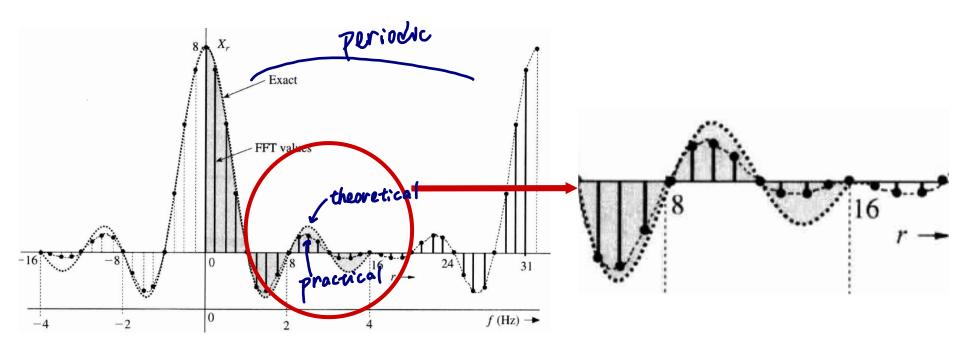


$$\bullet \quad \Omega_0 = \frac{2\pi}{32} = \frac{\pi}{16}$$





- Observe that X[k] is periodic.
- The dotted curve depicts the Fourier transform of x(t) = 8rect(t).
- The aliasing error is quite visible when we use a single graph to compare the superimposed plots. The error increases rapidly with k.



#### **Appendix: Proof of DFT relationships**

For the sampled signal we have:

$$\overline{x(t)} = \sum_{n=0}^{N_0 - 1} x(nT) \delta(t - nT).$$

• Since  $\delta(t - nT) \Leftrightarrow e^{-jn\omega T}$ 

$$\overline{X(\omega)} = \sum_{n=0}^{N_0 - 1} x(nT)e^{-jn\omega T}$$

• For  $|\omega| \leq \frac{\omega_s}{2}$ ,  $\overline{X(\omega)}$  the Fourier transform of  $\overline{x(t)}$  is  $\frac{X(\omega)}{T}$ , i.e.,

$$X(\omega) = T\overline{X(\omega)} = T\sum_{n=0}^{N_0 - 1} x(nT)e^{-jn\omega T}, |\omega| \le \frac{\omega_s}{2}$$
$$X[k] = X(k\omega_0) = T\sum_{n=0}^{N_0 - 1} x(nT)e^{-jnk\omega_0 T}$$

- If we let  $\omega_0 T = \Omega_0$  then  $\Omega_0 = \omega_0 T = 2\pi f_0 T = \frac{2\pi}{N_0}$  and also Tx(nT) = x[n].
- Therefore,  $X[k] = \sum_{n=0}^{N_0-1} x[n]e^{-jnk\Omega_0}$

# **Appendix: Proof of DFT relationships**

To prove the inverse relationship write:

$$\sum_{k=0}^{N_0-1} X[k] e^{jkm\Omega_0} = \sum_{k=0}^{N_0-1} \left[ \sum_{n=0}^{N_0-1} x[n] e^{-jnk\Omega_0} \right] e^{jkm\Omega_0} \Rightarrow$$

$$\sum_{k=0}^{N_0-1} X[k] e^{jkm\Omega_0} = \sum_{n=0}^{N_0-1} x[n] \left[ \sum_{k=0}^{N_0-1} e^{-jk(n-m)\Omega_0} \right]$$

- $\sum_{k=0}^{N_0-1} e^{-jk(n-m)\Omega_0} = \sum_{k=0}^{N_0-1} e^{-jk(n-m)\frac{2\pi}{N_0}} = \begin{cases} N_0 & n-m=rN_0, r \in \mathbb{Z} \\ 0 & \text{otherwise} \end{cases}$
- Since  $0 \le m, n \le N_0 1$  the only multiple of  $N_0$  that the term (n m) can be is 0. Therefore:

$$\sum_{k=0}^{N_0-1} e^{-jk(n-m)\frac{2\pi}{N_0}} = \begin{cases} N_0 & n-m=0 \Rightarrow n=m\\ 0 & \text{otherwise} \end{cases}$$

Therefore,

$$x_m = \frac{1}{N_0} \sum_{k=0}^{N_0 - 1} X[k] e^{jkm\Omega_0}, \ \Omega_0 = \frac{2\pi}{N_0}$$



#### **Continue with Dr Mike Brookes's notes**

 For the rest of the material related to DFT refer to Dr Mike Brookes's notes Three Different Fourier Transforms, from section Symmetries to the end.