

A Search for New Exotic Particles with Jets in Proton-Proton Collisions at the CMS Experiment

Dissertation

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2 Abstract

3 The Standard Model of particle physics has been very successful. But there are problems. They
4 can be explained by new exotic particles.

5 A search is presented for a new exotic particle. Events with jets are considered. A novel
6 technique is used. The search is based on Run-2 proton-proton collision data at a center-of-mass
7 energy of 13 TeV recorded with the CMS detector corresponding to an integrated luminosity of
8 137.1 fb^{-1} .

9 Upper limits are set on the production cross section of a new exotic particle as a function of
10 mass. Results are compared with theoretical predictions to obtain lower limits on the mass. At
11 95% confidence level, new exotic particles are excluded for masses below X GeV.

¹³ **Acknowledgements**

¹⁴ I want to thank the great people in the UZH group and at CMS Collaboration that have helped
¹⁵ me along the way, and with whom it was a great privilege and pleasure to work on diverse
¹⁶ projects and push the limits of our understanding of Nature.

¹⁷ Special thanks go to my supervisor, Prof. Dr. ...; my supervisor Dr. ...; ...

¹⁸ I would also like to thank my parents and family, for supporting me along my studies.

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60 1 Introduction

61 Over the last 100 years, physicists have made large strides in understanding the natural laws
62 that govern the universe as we know it. The Standard Model (SM) of particle physics describes
63 matter and its interactions.

64 In 2012, the discovery of the Higgs boson was announced by the ATLAS and CMS experiments [1–3] at CERN’s LHC. The Higgs was 47 years before that [4, 5]. Despite its many
65 successes, there are both theoretical and experimental reasons that imply the SM is not the final
66 theory of Nature.

67 This dissertation has the following structure: Chapter 2 discusses the SM, after which Chapter
68 3 motivates the search for new hypothetical particles. Chapter 4 describes the LHC accelerator,
69 CMS detector and the data collection. Chapter 5 describes the physics objects that are
70 relevant for this thesis. After these, Chapter 7 discusses the selections.

72 2 The Standard Model

73 This chapter introduces the Standard Model (SM) of particle physics. I left a bit of text with
74 equations, citations, figures, and tables as an example of this template.

75 2.1 The Standard Model

76 A rough timeline of the experimental discovery of particles is given by Fig. 2.1. More detailed
77 historical discussions can be found in Refs. [6–8].

78 With a mass of $m_t \approx 172.8 \text{ GeV}$ [10, p. 32], the top quark is by far the heaviest fermion.

79 The discovery of the Higgs was announced in 2012 [1–3, 11]. A single Higgs boson is the
80 simplest and most minimal solution, although it is possible to achieve SBB with more than one
81 Higgs boson. No evidence of additional Higgs bosons have been found so far [12].

82 2.1.1 The Standard Model Lagrangian

83 The SM is encoded in the *SM Lagrangian density*. The terms of the Lagrangian density can be
84 grouped into three separate *sectors*:

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}}. \quad (2.1)$$

85 Secondly, the SM is a *gauge theory*: It has an internal symmetry that corresponds to the
86 local $\text{SU}(3)_C \times \text{SU}(2)_L \times \text{U}(1)_Y$ symmetry group. The SM gauge group is spontaneously broken
87 to

$$\text{SU}(3)_C \times \text{SU}(2)_L \times \text{U}(1)_Y \xrightarrow{\text{SSB}} \text{SU}(3)_C \times \text{U}(1)_{\text{EM}} \quad (2.2)$$

88 by the nonzero vev of the Higgs field, where $\text{U}(1)_{\text{EM}}$ is the gauge symmetry group for QED.

89 2.1.2 The gauge sector

The gauge sector can be written as

$$\mathcal{L}_{\text{gauge}} = \sum_f \sum_{\psi_f} \bar{\psi}_f i\gamma^\mu D_\mu \psi_f - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu} - \frac{1}{4} W_{\mu\nu}^i W_i^{\mu\nu} - \frac{1}{4} B^{\mu\nu} B^{\mu\nu}, \quad (2.3)$$

90 in natural units $c = \hbar = 1$ [14], and where $a = 1, \dots, 8$ and $i = 1, 2, 3$. In the first term, matter
91 is given by five fermion fields ψ_f in three generations $f = 1, 2, 3$;

$$Q_L^f = \begin{pmatrix} u_L^f \\ d_L^f \end{pmatrix}, \quad u_R^f, \quad d_R^f, \quad L_L^f = \begin{pmatrix} \nu_L^f \\ e_L^f \end{pmatrix}, \quad e_R^f. \quad (2.4)$$

92 Their Dirac adjoints are given by $\bar{\psi}_f = \psi_f^\dagger \gamma^0$.

93 The covariant derivative

$$D_\mu = \partial_\mu + ig_s t_a G_\mu^a + ig T_i W_\mu^i + ig' \frac{Y}{2} B_\mu \quad (2.5)$$

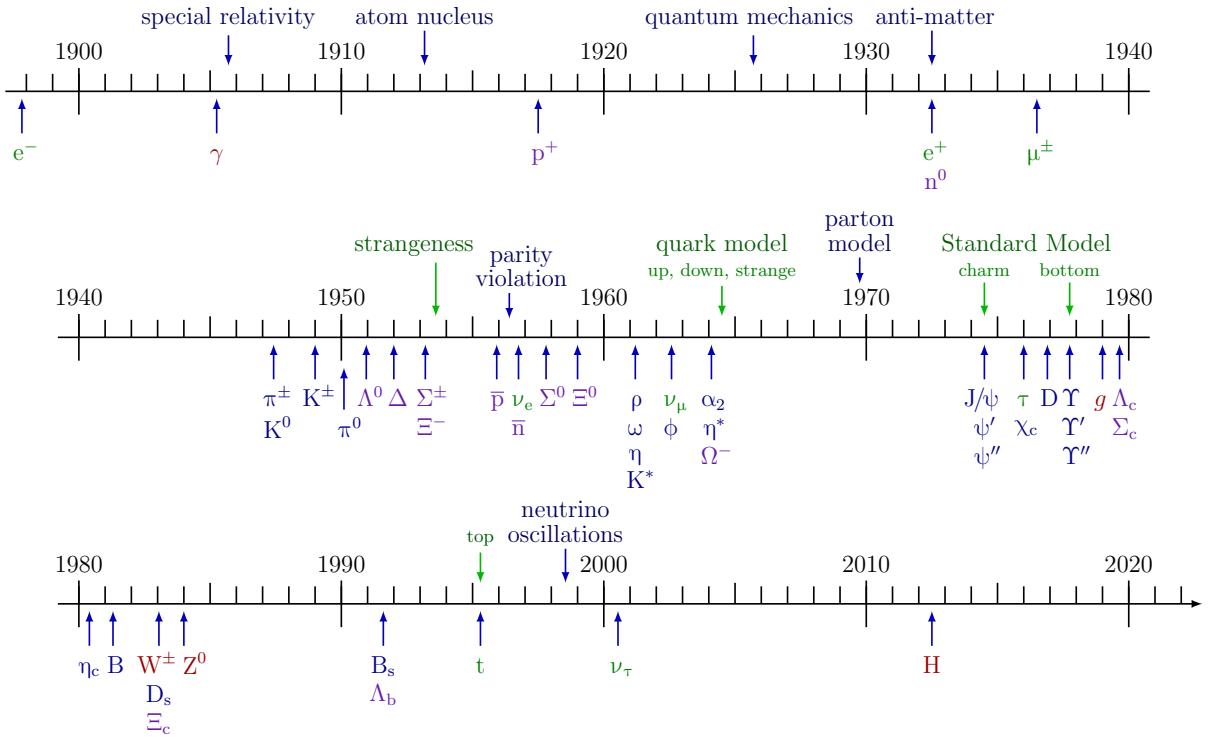


Figure 2.1: A timeline of particle physics in the last 130 years. Adapted from Ref. [9].

three generations of matter (fermions)			interactions / force carriers (bosons)	
I	II	III	g gluon	H higgs
mass charge spin $=2.2 \text{ MeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$ u up	mass charge spin $=1.28 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$ c charm	mass charge spin $=173.1 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$ t top	0 0 1 0 0 1 g gluon	$=124.97 \text{ GeV}/c^2$ 0 0 0 0 0 0 H higgs
mass charge spin $=4.7 \text{ MeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$ d down	mass charge spin $=96 \text{ MeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$ s strange	mass charge spin $=4.18 \text{ GeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$ b bottom	0 0 1 0 0 1 γ photon	$=105.66 \text{ MeV}/c^2$ -1 $\frac{1}{2}$ μ muon
mass charge spin $=0.511 \text{ MeV}/c^2$ -1 $\frac{1}{2}$ e electron	mass charge spin $=1.7768 \text{ GeV}/c^2$ -1 $\frac{1}{2}$ τ tau	mass charge spin $=91.19 \text{ GeV}/c^2$ -1 $\frac{1}{2}$ Z Z boson	$=80.433 \text{ GeV}/c^2$ ± 1 1 W W boson	$=80.433 \text{ GeV}/c^2$ ± 1 1 W W boson
mass charge spin $<1.0 \text{ eV}/c^2$ 0 $\frac{1}{2}$ ν _e electron neutrino	mass charge spin $<0.17 \text{ MeV}/c^2$ 0 $\frac{1}{2}$ ν _μ muon neutrino	mass charge spin $<18.2 \text{ MeV}/c^2$ 0 $\frac{1}{2}$ ν _τ tau neutrino		

Figure 2.2: A table of the particles in Standard Model of particle physics and their properties. Figure taken from Ref. [13].

Table 2.1: Summary of the representation and quantum numbers SM fields. Bold numbers indicate the dimension of the representation under the respective gauge group.

Field name	Symbol	Representations		Quantum numbers		
		SU(3) _C	SU(2) _L	T_3	$Y/2$	$Q = T_3 + Y/2$
Quark doublet	$Q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	3	2	$\begin{pmatrix} +\frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$	$+\frac{1}{6}$	$\begin{pmatrix} +\frac{2}{3} \\ -\frac{1}{3} \end{pmatrix}$
Up-quark singlet	u_R	3	1	$+\frac{2}{3}$	$+\frac{2}{3}$	$+\frac{2}{3}$
Down-quark singlet	d_R	3	1	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$
Lepton doublet	$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$	1	2	$\begin{pmatrix} +\frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$	$-\frac{1}{2}$	$\begin{pmatrix} 0 \\ -1 \end{pmatrix}$
Lepton singlet	e_R	1	1	0	-1	-1
Gluon field	G_μ^a	8	1	0	0	0
Weak gauge field	$W_\mu^i = \begin{pmatrix} W_\mu^+ \\ W_\mu^- \\ W_\mu^3 \end{pmatrix}$	1	3	$\begin{pmatrix} +1 \\ -1 \\ 0 \end{pmatrix}$	0	$\begin{pmatrix} +1 \\ -1 \\ 0 \end{pmatrix}$
Hypercharge field	B_μ	1	1	0	0	0
Higgs doublet	$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$	1	2	$\begin{pmatrix} +\frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$	$+\frac{1}{2}$	$\begin{pmatrix} +1 \\ 0 \end{pmatrix}$
Conjugate Higgs doublet	$\Phi^c = \begin{pmatrix} \phi^{0*} \\ \phi^- \end{pmatrix}$	1	2	$\begin{pmatrix} +\frac{1}{2} \\ -\frac{1}{2} \end{pmatrix}$	$-\frac{1}{2}$	$\begin{pmatrix} 0 \\ -1 \end{pmatrix}$

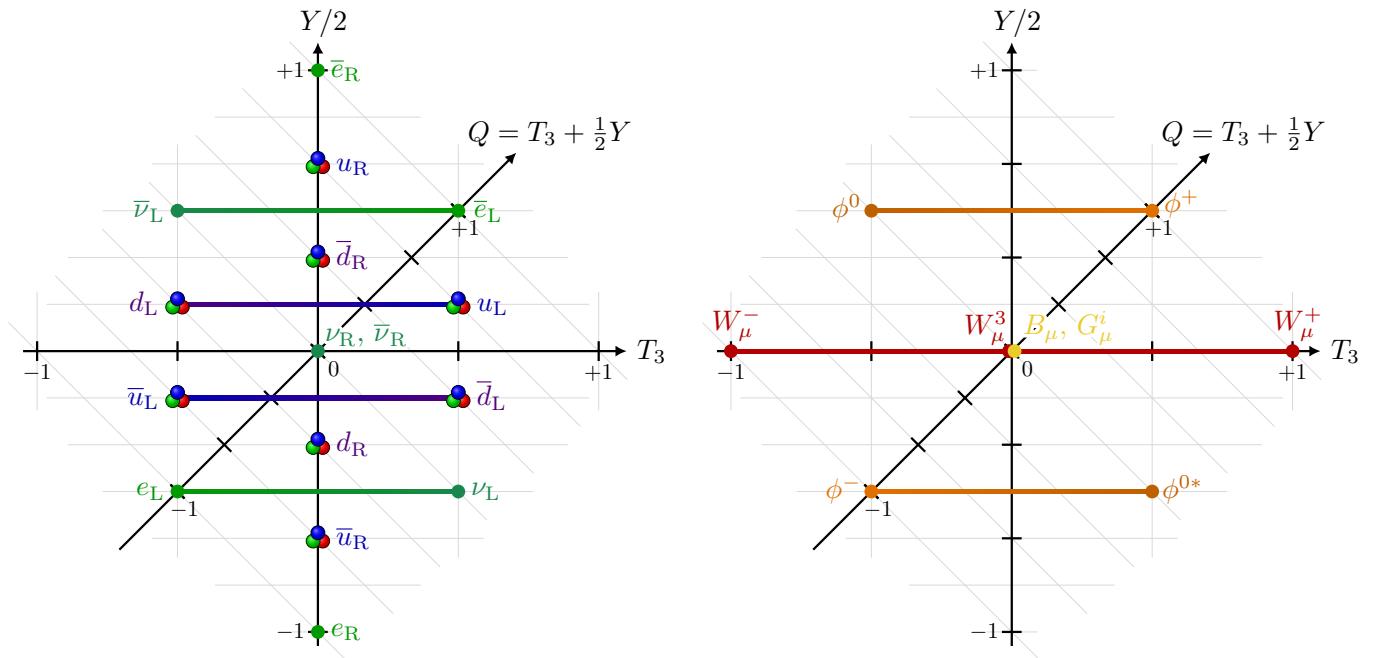


Figure 2.3: Graph of the $(T_3, Y/2)$ quantum numbers.

where T_i and t_a are the generators for $SU(2)_L$ and $SU(3)_C$, respectively, and Y is the weak hypercharge of the field that D_μ acts on.

Through the Higgs mechanism, the gauge bosons acquire mass and mix into the mass eigenstates

$$W_\mu^\pm = \frac{W_\mu^1 \mp iW_\mu^2}{\sqrt{2}}, \quad Z_\mu = W_\mu^3 \cos \theta_W + B_\mu^3 \sin \theta_W, \quad A_\mu = W_\mu^3 \sin \theta_W - B_\mu^3 \cos \theta_W, \quad (2.6)$$

where W_μ^\pm are the charged W boson fields, Z_μ is the neutral Z boson field, A_μ is the photon field, and θ_W is the Weinberg angle given by $\sin \theta_W = g'/\sqrt{g^2 + g'^2}$. After mixing, the new covariant derivative becomes

$$D_\mu = \partial_\mu + ig_s t_a G_\mu^a + i \frac{g}{\sqrt{2}} T^+ W_\mu^+ + i \frac{g}{\sqrt{2}} T^- W_\mu^- + i \frac{g}{\cos \theta_W} (T_3 - \sin^2 \theta_W Q) Z_\mu + ieQ A_\mu, \quad (2.7)$$

with ladder operators $T^\pm = T_1 \pm iT_2$.

2.1.3 The Yukawa sector

In order to explain mass of the fermions,

$$\mathcal{L}_{\text{Yukawa}} = -Y_d^{fg} \bar{Q}_L^f \Phi d_R^g - Y_u^{fg} \bar{Q}_L^f \Phi^c u_R^g - Y_e^{fg} \bar{L}_L^f \Phi e_R^g + \text{h.c.}, \quad (2.8)$$

with the conjugate Higgs doublet $\Phi^c = i\sigma_2 \Phi^\dagger$, and unitary Yukawa matrices Y_d^{fg} , Y_u^{fg} , $Y_e^{fg} \in U(3)$ that mixes the fermion generations. As Φ is a scalar field that forms a weak isospin doublet,

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}. \quad (2.9)$$

Because the Higgs vev is nonzero, we can fix the $SU(2)_L$ gauge. It is convenient to choose the unitary gauge

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h \end{pmatrix}, \quad (2.10)$$

with the vev $v \approx 246$ GeV, a real constant, and the neutral Higgs field h , a (real) scalar. Therefore, after SSB, we are left with

$$\mathcal{L}_{\text{Yukawa}} = -\frac{Y_d^{fg}}{\sqrt{2}} \bar{d}_L^f (v + h) d_R^g - \frac{Y_u^{fg}}{\sqrt{2}} \bar{u}_L^f (v + h) u_R^g - \frac{Y_e^{fg}}{\sqrt{2}} \bar{e}_L^f (v + h) e_R^g + \text{h.c.} \quad (2.11)$$

2.1.4 CKM matrix

To explain the suppressed decay rate of $K^0 \rightarrow \mu^+ \mu^-$, Glashow, Iliopoulos, and Maiani used a mixing matrix and proposed the existence of the charm quark [15]. This is referred to as the *Glashow–Iliopoulos–Maiani (GIM) mechanism*. The mixing of two quark generations is given by the *Cabibbo matrix*:

$$\begin{pmatrix} d' \\ s' \end{pmatrix} = \begin{pmatrix} \cos \theta_C & \sin \theta_C \\ -\sin \theta_C & \cos \theta_C \end{pmatrix} \begin{pmatrix} d \\ s \end{pmatrix}, \quad (2.12)$$

where $|d'\rangle$ and $|s'\rangle$ represent the weak eigenstates, which are linear combinations of the mass eigenstates $|d\rangle$ and $|s\rangle$. The mixing is quantified with the *Cabibbo angle* $\theta_C \approx 13.04^\circ$.

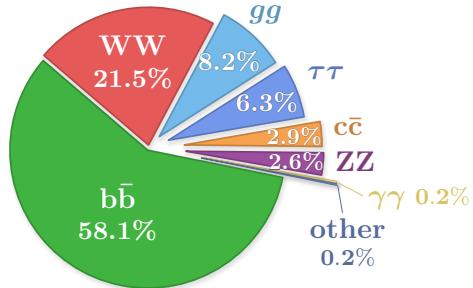


Figure 2.4: Pie chart of Higgs branching fractions.

117 The mixing of three generations is given by

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}, \quad (2.13)$$

118 where the matrix is the so-called *Cabbibo-Kobayashi-Maskawa (CKM) matrix*, denoted by V_{CKM} .

119 **2.2 Higgs**

120 A pie chart of Higgs decay can be found in Fig. 2.4.

121 **2.3 The physics of proton-proton collisions**

122 The substructure of the proton has been carefully studied in deep inelastic scattering (DIS)
123 experiments at SLAC [16, 17] and DESY [18], which collided electrons or positrons with protons.

124 In the parton model, *parton distribution functions* (PDFs) describe the probability of finding
125 a parton of type p with momentum fraction x in a collision at some momentum scale Q as a
126 function of the form $f_p(x, Q^2)$ [19–23]. The proton has *valence quarks* and *sea quarks*. To reflect
127 the proton quantum numbers, the PDFs are normalized:

$$\int_0^1 [f_u(x, Q^2) - f_{\bar{u}}(x, Q^2)] dx = 2, \quad \int_0^1 [f_d(x, Q^2) - f_{\bar{d}}(x, Q^2)] dx = 1. \quad (2.14)$$

128 This follows from the factorization theorem [25, 26], which provides a formula to calculate
129 the cross sections of a hard process with an integral of the form

$$\sigma(\text{pp} \rightarrow X + Y) = \sum_{i,j} \int_0^1 \int_0^1 f_i(x_1, Q^2) f_j(x_2, Q^2) \hat{\sigma}_{ij \rightarrow X}(x_1, x_2, Q^2) dx_1 dx_2, \quad (2.15)$$

130 where $\hat{\sigma}_{ij \rightarrow X}(x_1, x_2, Q^2)$ is the parton-level cross section of the hard process, $i + j \rightarrow X$, and
131 parton i (j) carries a fraction x_1 (x_2) of the momentum of its mother proton.

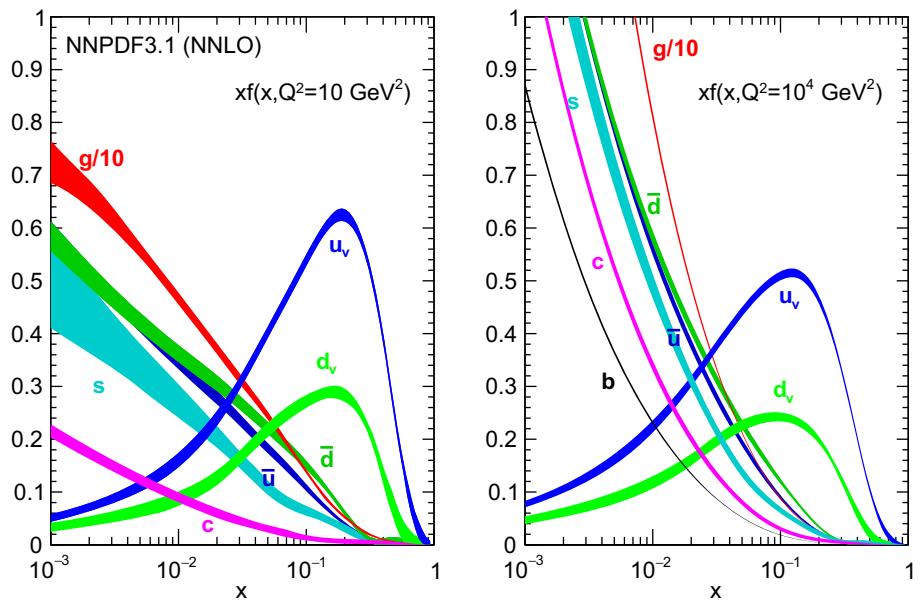


Figure 2.5: Plot of the proton PDF $f_p(x, Q^2)$ times the momentum fraction x as calculated by NNPDF3.1 at NNLO accuracy in perturbation theory for $Q^2 = 10 \text{ GeV}^2$ (left) and $Q^2 = 10^4 \text{ GeV}^2$ (right). Adapted from [24].

¹³² **3 Beyond the Standard Model**

¹³³ The SM of particle physics is successful. Fine-structure constant $\alpha = e^2/4\pi$ has been precisely
¹³⁴ calculated and measured [27–30]. In 2012, the discovery of the Higgs boson was announced by
¹³⁵ the ATLAS and CMS experiments at the CERN LHC [1–3, 11], more than 47 years after its
¹³⁶ prediction [4, 5]. But there are both problems. The following section highlights several examples
¹³⁷ of open questions in physics.

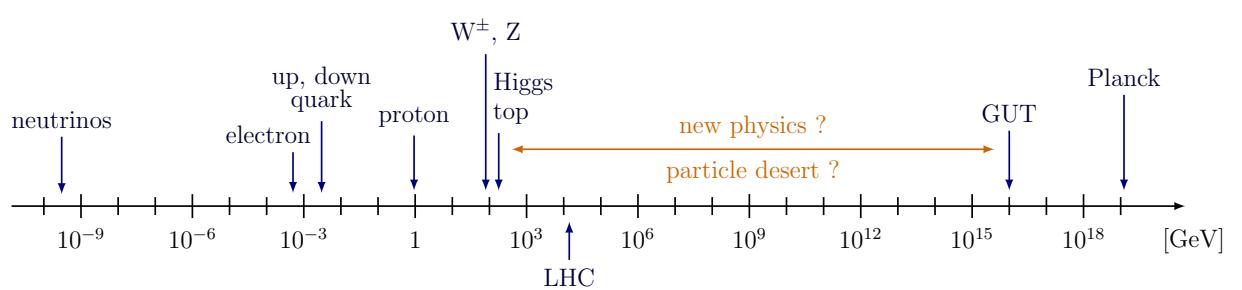


Figure 3.1: Logarithmic scale of different energy scales in particle physics. Adapted from Ref. [9].

¹³⁸ 4 The CMS Experiment

¹³⁹ This chapter will discuss the experimental setup.

¹⁴⁰ 4.1 The Large Hadron Collider

¹⁴¹ The Large Hadron Collider (LHC) is the world's largest and most powerful particle collider
¹⁴² with a 27 km circumference and a record collision energy of 13.6 TeV. It surpasses the Tevatron
¹⁴³ accelerator at Fermilab in the United States [32] (6.3 km, 1.96 TeV), and LEP (209 GeV) [10,
¹⁴⁴ Section 32].

¹⁴⁵ During the data-taking period of 2015 to 2018, referred to as "Run 2", the energy of each
¹⁴⁶ proton beam was 6.5 TeV, which means the center-of-mass energy was $\sqrt{s} = 13$ TeV. The LHC
¹⁴⁷ machine is described in more technical detail in Ref. [33].

¹⁴⁸ 4.1.1 Luminosity

¹⁴⁹ It depends on several beam parameters. The simplified expression is [10, p. 533]:

$$\mathcal{L} = \frac{N_1 N_2 f}{A}. \quad (4.1)$$

¹⁵⁰ The LHC is designed to reach the nominal luminosity of $\mathcal{L} = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. See Fig. 4.2. Each
¹⁵¹ bunch contains about 110 billion protons [33, 35].

¹⁵² Figure 4.4 shows the distribution of the number of interactions per bunch crossing in all

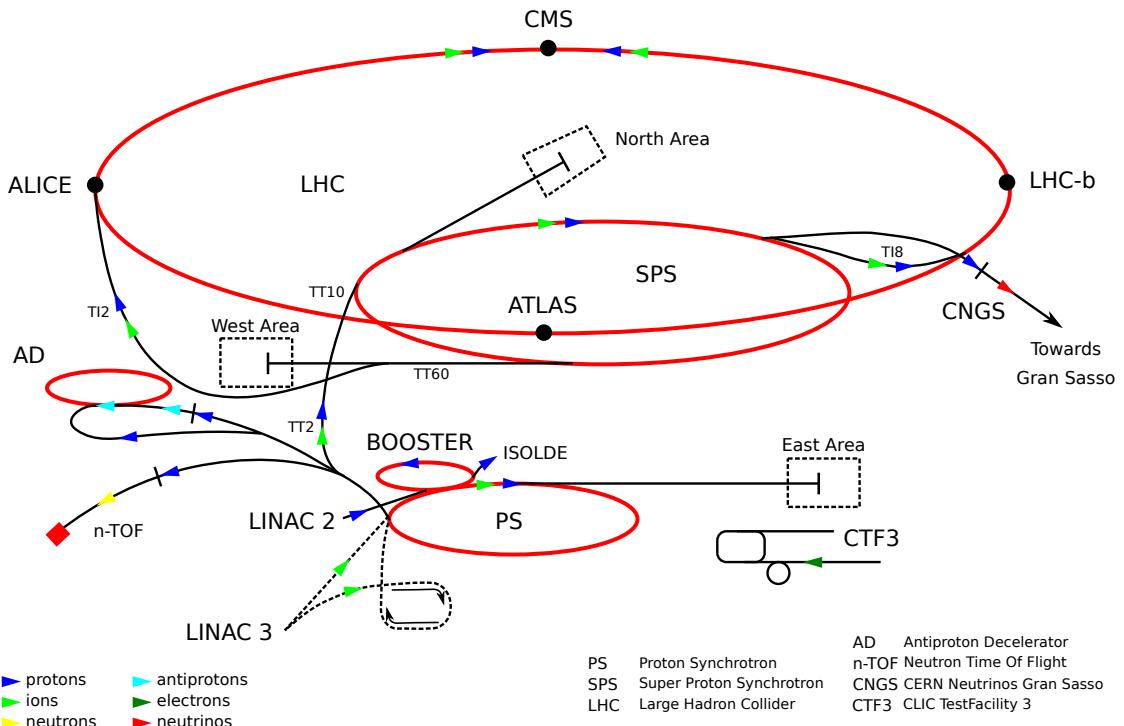


Figure 4.1: CERN's accelerator complex. Taken from [31].

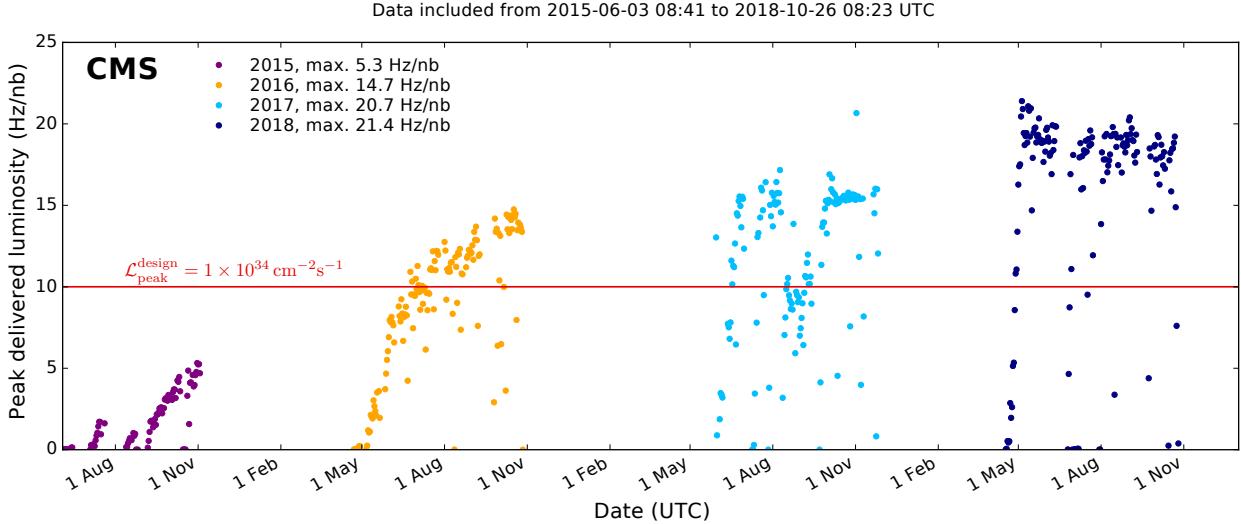


Figure 4.2: The instantaneous luminosity \mathcal{L} . The design luminosity is $\mathcal{L}_{\text{peak}}^{\text{design}} = 10^{34} \text{ cm}^{-2} \text{s}^{-1}$. Adapted from [34].

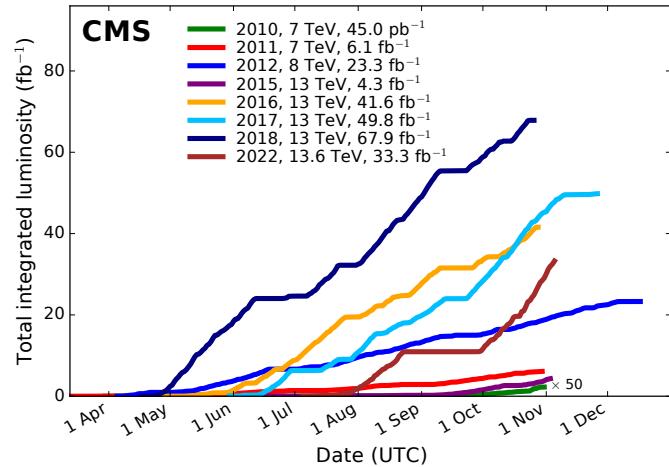


Figure 4.3: Integrated luminosity collected by CMS at $\sqrt{s} = 13 \text{ TeV}$. Figure taken from [34].

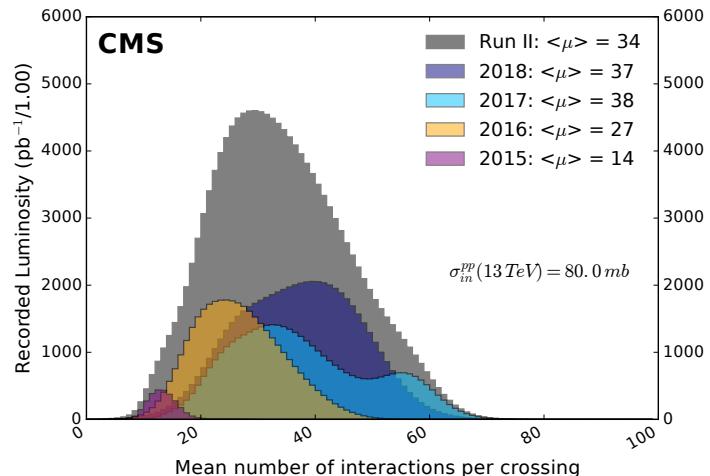


Figure 4.4: Distribution of number of pp interactions, assuming an inelastic cross section of 80 mb. Figure taken from [34].

153 data-taking years of Run 2.

154 Figure 4.3 shows the data of each data-taking year with *integrated luminosity* L , which is

$$L := \int_{\text{data taking}} \mathcal{L}(t) dt, \quad (4.2)$$

155 given in units of inverse area, such as the inverse barn (b^{-1}). Between 2016 and 2018, CMS
156 collected about 138 fb^{-1} of pp collision data for physics analysis.

157 4.2 The CMS detector

158 The Compact Muon Solenoid (CMS) detector at one of the interaction points at the LHC. An
159 illustration is shown in Fig. 4.5. A detailed description can be found in Ref. [36].

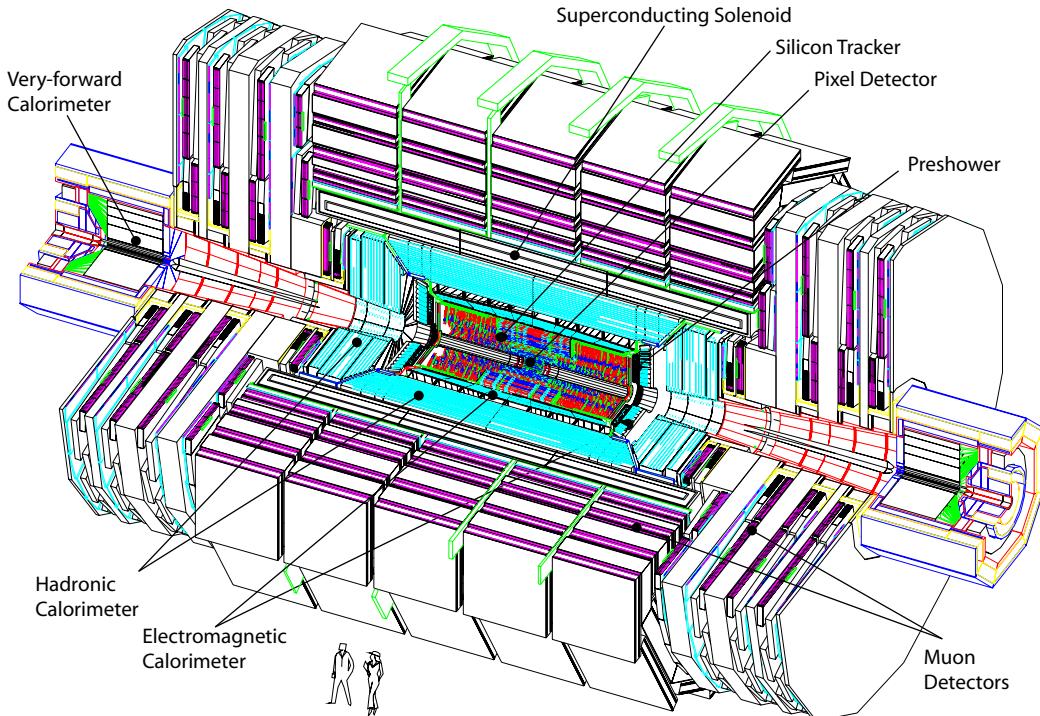


Figure 4.5: The CMS detector. Taken from [36].

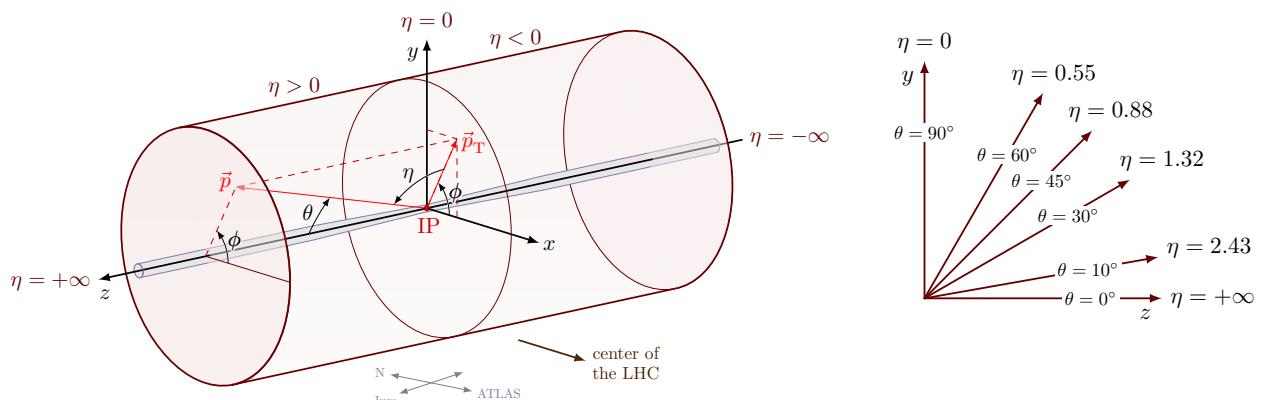


Figure 4.6: Left: The conventional coordinate system of CMS with momentum vector \vec{p} . Taken from [37]. Right: Pseudorapidity. Taken from [38].

160 **4.2.1 Coordinate system**

161 The conventional coordinate system of CMS is defined in Fig. 4.6.

162 **4.2.2 The solenoid magnet & flux-return yoke**

163 The *superconducting solenoid magnet* is a 13 m long cylinder with an inner diameter of 6 m. The
164 yoke [42].

165 **4.2.3 The pixel tracker**

166 The inner tracker system has a *silicon pixel tracker* and a *silicon microstrip tracker*. Figure 4.8
167 presents a closer look of the full tracker layout. The pixel tracker from 2008 up to 2016 was
168 composed of 1440 sensor modules with a total of 66 million silicon pixels [43]. During the
169 technical shutdown it was upgraded with 1856 modules with 124 million pixels in total [41, 44].
170 Their layouts are compared in Fig. 4.9. See Refs. [43, 45] for more information on the resolution.

171 **4.2.4 The silicon strip tracker**

172 The silicon strip tracker is larger. A detailed description of the tracking and vertexing software
173 is given in Ref. [43].

174 **4.2.5 The electromagnetic calorimeter**

175 The *electromagnetic calorimeter* (ECAL) is built around the tracker. The energy resolution can
176 be parametrized as a function of energy [36]:

$$\left(\frac{\sigma_E}{E}\right)^2 = \left(\frac{S}{\sqrt{E}}\right)^2 \oplus \left(\frac{N}{E}\right)^2 \oplus C^2, \quad (4.3)$$

177 with the stochastic term S , the noise N , and a constant term C . The technical design report [46]
178 describes the ECAL in more technical detail.

179 **4.2.6 The hadronic calorimeter**

180 The last subdetector inside the solenoid is the *hadron calorimeter* (HCAL). The jet energy
181 resolution can be found in Ref. [47]

182 The *very forward HCAL* (HF) is positioned around the beamline outside the detector.

183 **4.2.7 The muon system**

184 The muon system is outside the solenoid magnet. The final momentum resolution can be roughly
185 parametrized as follows:

$$\left(\frac{\sigma_{p_T}}{p_T}\right)^2 = (A \cdot p_T)^2 \oplus C^2, \quad (4.4)$$

186 where A , and C are constants determined by the hit resolution and multiple scattering, respec-
187 tively [48]. The resolution grows with momentum, because the track becomes more straight,
188 which increases the uncertainty in its curvature. The technical design report can be found in
189 Ref. [49].

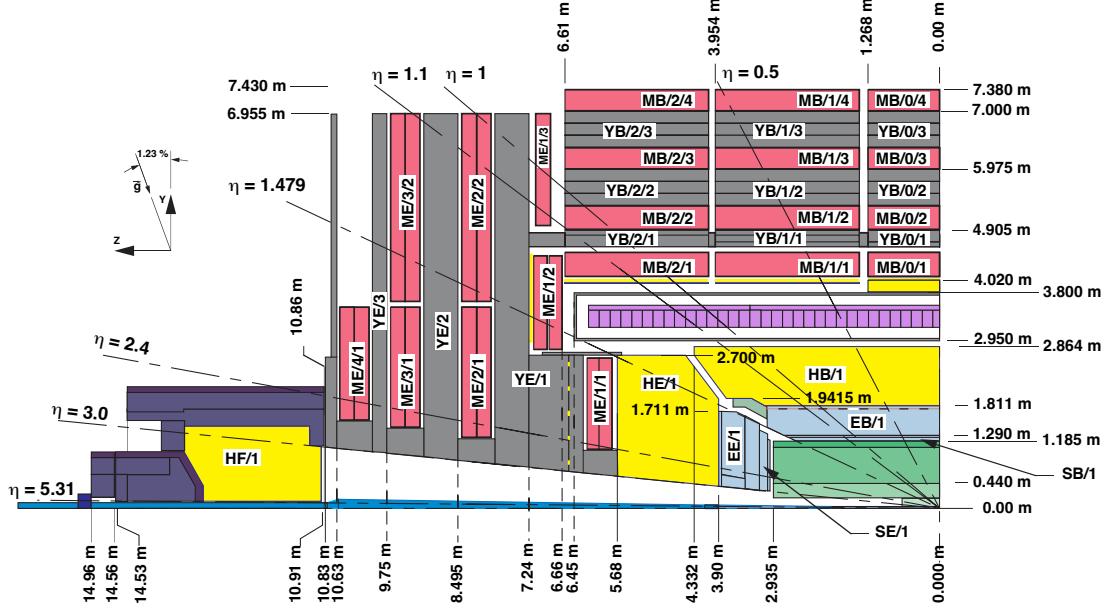


Figure 4.7: Schematic of one quadrant of the CMS detector in the positive zy -plane. Adapted from [39].

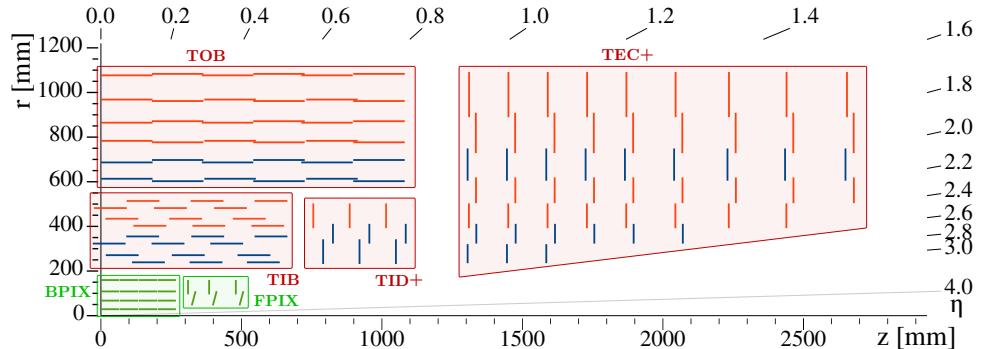


Figure 4.8: Schematic view of the CMS tracking system. Adapted from [40].

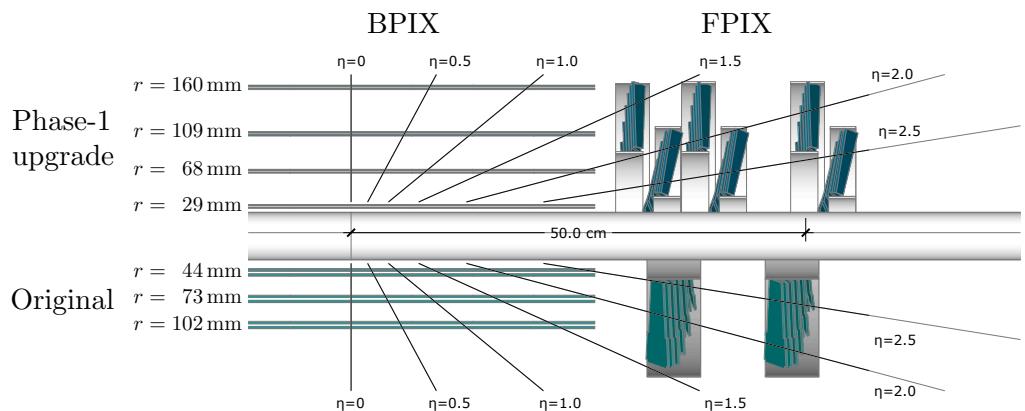


Figure 4.9: Layout of the CMS pixel detector. Adapted from [41].

190 **4.2.8 The trigger system**

191 The collision rate of 40 MHz is higher than is possible to record offline. The *L1 trigger* has a
192 rate of about 100 kHz. The next level, is the *high-level trigger* (HLT). See Ref. [50].

5 Object & Event Reconstruction

193 This chapter explains object and event reconstruction in the CMS experiment.

5.1 Principles of particle identification

196 Figure 5.1 compares the lifetime, dividing them in different regions of stability in the context of
197 the CMS experiment.

198 Each of these particles has a unique set of properties that leave a characteristic signal in one
199 or more of the subdetector, as illustrated in Fig. 5.2.

5.2 Particle-flow algorithm

201 The particle-flow (PF) algorithm [53, 54] fully reconstructs the event of a pp collision with an
202 optimized combination of all measurements of the CMS subsystems.

203 Tracks in the inner tracker or muon system are iteratively built from hits, using the *Kalman-*
204 *filter (KF) technique* [55, 56].

5.3 Primary vertex

206 The algorithm locates primary vertices (PVs) of pp collisions. The tracks are clustered using the
207 *deterministic annealing algorithm* [57]. Finally, candidate vertices are fitted using an *adaptive*
208 *vertex fitter* [58]. The PV with the largest sum of momenta is assumed to be the vertex of the
209 hard scattering process [43, 54, 59].

5.4 Electrons

211 Electrons are reconstructed by associating a track in the inner tracker to clusters of energy
212 deposits in the ECAL. Candidate tracks are refitted with the *Gaussian sum filter* (GSF) [60].
213 The ECAL clusters are recombined into a so-called *supercluster*. The resolution is documented
214 in Ref. [61, 62].

215 The reconstruction and identification of electrons, as well as high-energy photons, are dis-
216 cussed in more detail in Refs. [63] and [61].

217 The *relative isolation* for electrons is computed with the so-called *rho-effective-area method* [64]:

$$I_{\text{rel}}^e := \frac{\sum_{\text{charged}} p_T + \max(0, \sum_{\text{neutral}} E_T - \rho A_{\text{eff}})}{p_T^e}. \quad (5.1)$$

5.5 Muons

219 Muon candidates in CMS are reconstructed as *standalone muons*, *tracker muons*, and/or *global*
220 *muons* [65].

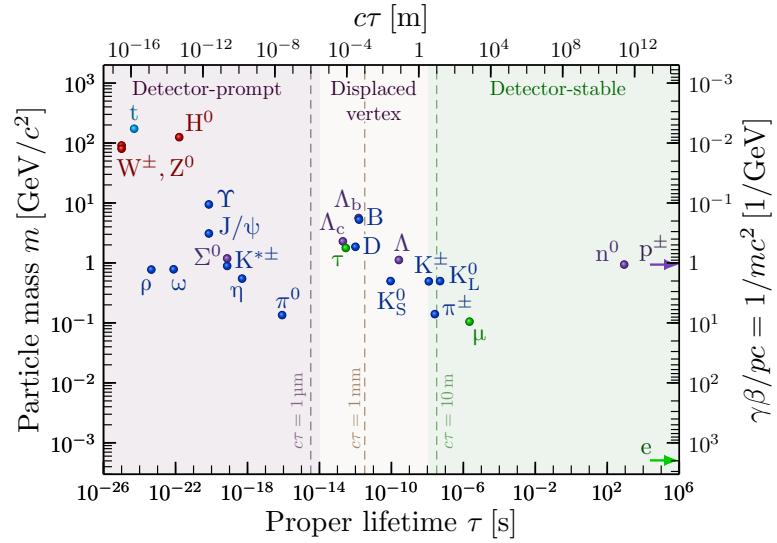


Figure 5.1: Plot of the mass versus lifetime τ of many composite and fundamental SM particles. The decay length is given by $L = \gamma\beta c\tau$, with $\gamma\beta = p/mc$. Adapted from [51].

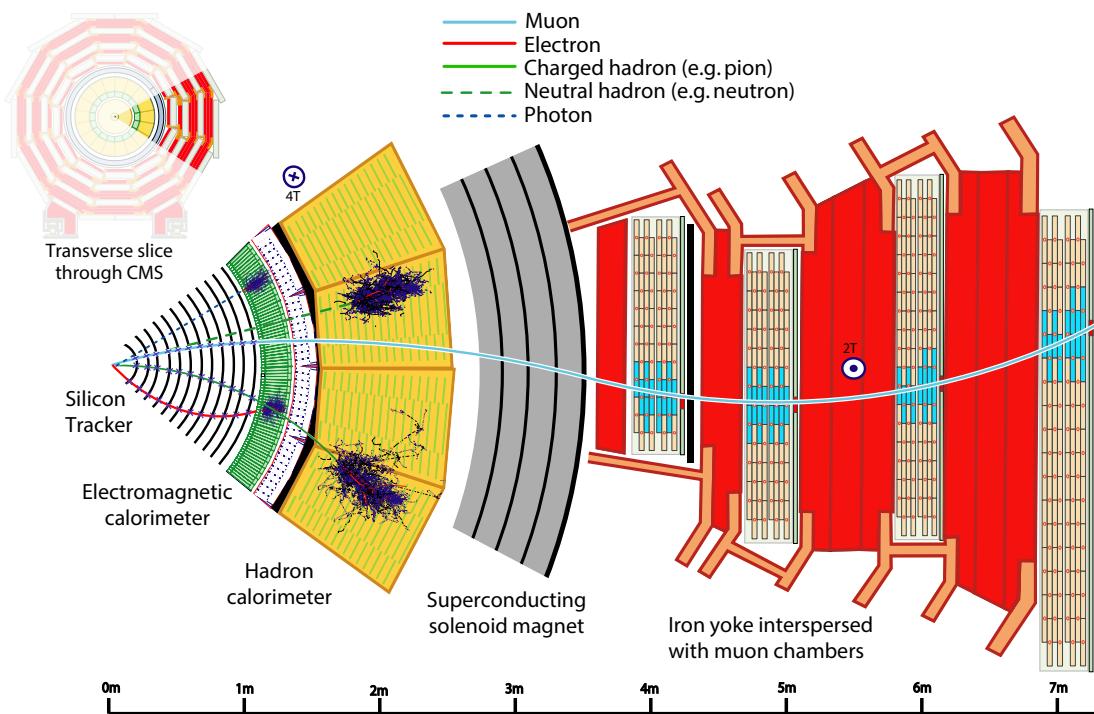


Figure 5.2: Particles in the CMS detector. Adapted from from [52].

221 The relative isolation of a muon is defined with the so-called $\Delta\beta$ -corrections:

$$I_{\text{rel}}^{\mu} := \frac{\sum_{\text{charged}} p_T + \max(0, \sum_{\text{neutral}} E_T - \Delta\beta \sum_{\text{charged, PU}} p_T)}{p_T^{\mu}}. \quad (5.2)$$

222 5.6 Hadronically Decayed τ Leptons

223 Hadronic decays of τ leptons (τ_h) are reconstructed with the *hadron-plus-strips* (HPS) algorithm
224 [66–69]. About 65% of hadronic τ decays involve neutral pions, which decay promptly ($\tau =$
225 8.4×10^{-17} s) to two photons 98.8% of the time [10, p. 33].

226 The τ_h identification algorithm DEEPTAU [69] was used to discriminate against jets that are
227 initiated by quarks and gluons, as well as electrons and muons, see Fig. 5.5.

228 5.7 Jets

229 Jets are reconstructed with a clustering algorithm. The jet properties must be *infrared safe* and
230 *collinear safe*. The *anti- k_T* (AK) algorithm utilizes the metric

$$d_{ij} = \min(p_{T,i}^{-2}, p_{T,j}^{-2}) \frac{\Delta R_{ij}^2}{\Delta R_y^2}, \quad (5.3)$$

231 where $\Delta R_y = \sqrt{\Delta y^2 + \Delta\phi^2}$ is the distance in (y, ϕ) -space with the *rapidity*

$$y := \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right). \quad (5.4)$$

232 If a cluster i of combined objects satisfies

$$d_{ij} > d_{iB} = p_{T,i}^{-2}, \quad (5.5)$$

233 More info is given in Ref. [54, 71]. The AK algorithm is implemented in the FASTJET library [72,
234 73], and *charged-hadron subtraction* (CHS) is used.

235 The CMS Collaboration determines *jet energy corrections* (JECs) in several steps, which are
236 explained in detail in Ref. [74]. Additionally, the jet energy and *jet energy resolution* (JER) of
237 jets in simulation are corrected to obtain better agreement between simulation and data.

238 Genuine jets are distinguished from pileup jets with the identification algorithm [75]. Iden-
239 tification [76] removes spurious jet-like features that originate from isolated noise patterns in
240 certain HCAL regions.

241 5.8 Bottom quark identification

242 The DEEPCSV algorithm [77, 78] for b tagging jets. DEEPCSV is a deep neural network
243 (DNN) with four hidden layers that is an extension of the combined secondary vertex (CSV)
244 algorithm [77, 78].

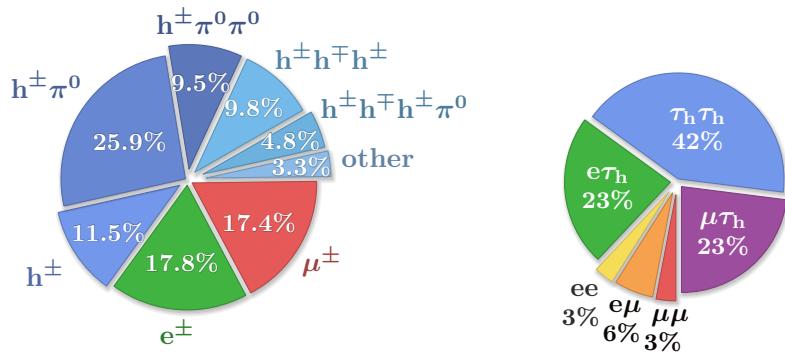


Figure 5.3: Pie charts of branching fractions. **Left:** τ lepton decay. **Right:** Decay channels of a pair of τ leptons. Numbers from PDG [10, p. 28].

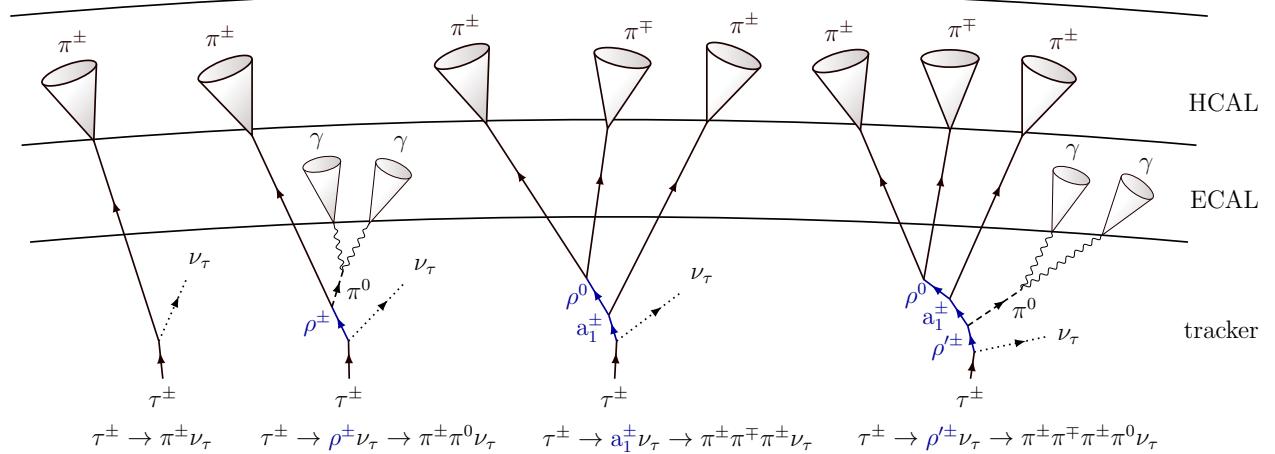


Figure 5.4: An illustration τ_h signatures. Inspired by Ref. [70, p. 37].

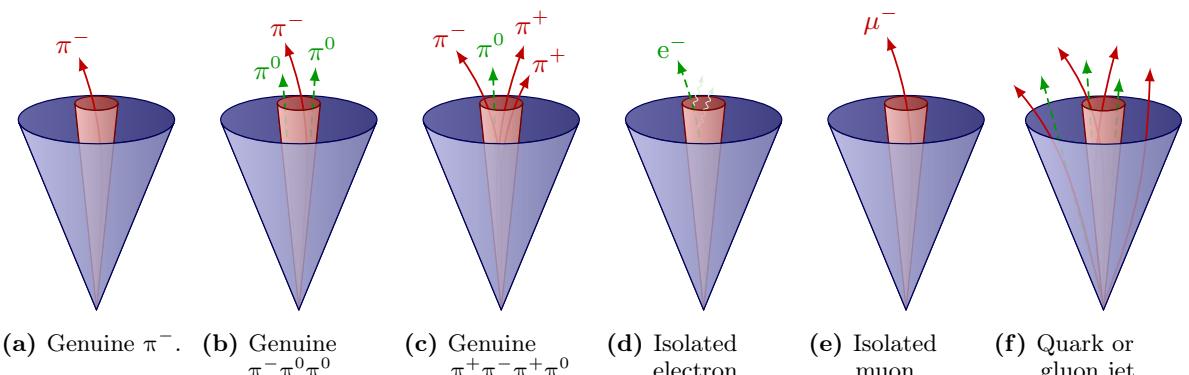


Figure 5.5: Illustration of several hadronic decay modes of the τ leptons (a-c) and their backgrounds (d-f).

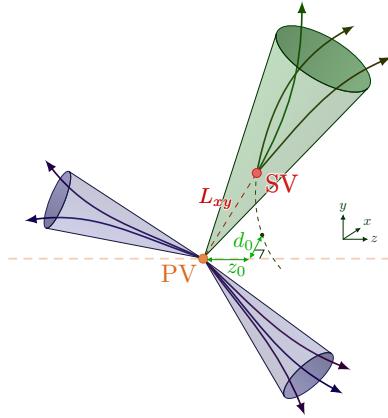


Figure 5.6: B-tagging.

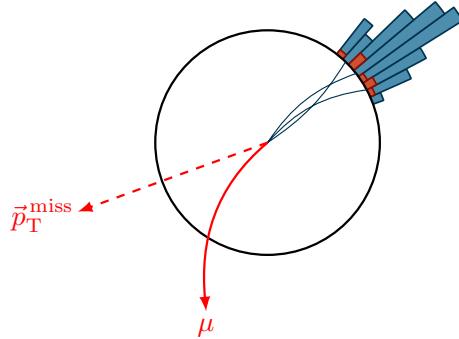


Figure 5.7: Illustration of missing transverse energy (MET) with a muon track (red), and energy deposits in the calorimeters.

245 **5.9 Missing transverse energy**

246 The vectorially sum of all particles gives the missing momentum

$$\vec{p}_T^{\text{miss}} = - \sum_i \vec{p}_T^i. \quad (5.6)$$

247 This vector, or its length, is often referred to as the *missing transverse momentum*, or *missing
248 transverse energy* (MET). Several corrections are applied [54]. More details can be found in
249 Refs. [79] and [80].

250 **6 Data sets & Simulated Samples**

251 This chapter will discuss in more detail the data sets of pp collision events used, as well as the
252 samples of simulated background events.

253 **6.1 Data sets**

254 The total data set recorded between 2016 and 2018 corresponds to an integrated luminosity of
255 approximately 138 fb^{-1} at $\sqrt{s} = 13 \text{ TeV}$. The run ranges and integrated luminosity per year are
256 listed in Table 6.1.

Table 6.1: LHC run number ranges and integrated luminosity L .

Run number range	$L [\text{fb}^{-1}]$
272007–284044	36.3
297020–306462	41.5
315252–325175	59.7

257 **6.2 Event simulation**

258 The PYTHIA parameters affecting the description of the underlying event are set to the CUETP8M1
259 (CP5) tune for all 2016 (2017 and 2018) samples [82, 83], except for the 2016 $t\bar{t}$ sample, for
260 which CUETP8M2T4 [84] is used. The interaction with the CMS detector is simulated by
261 GEANT4 [85]. Pileup is generated with PYTHIA. Event generation is described in more detail in
262 Refs. [10, p. 717], [86] and [87].

263 **6.3 Backgrounds**

264 **6.3.1 Simulated backgrounds**

265 All the simulated background samples used in this thesis are listed in Table 6.2. The $W + \text{jets}$ and
266 $Z + \text{jets}$ processes are generated with MADGRAPH [88] at LO precision. The MLM jet matching
267 and merging scheme [89] is used to match partons between MADGRAPH and PYTHIA and prevent
268 overcounting. The $Z + \text{jets}$ samples generated at NLO by MADGRAPH5_aMC@NLO [90]. The
269 production of $t\bar{t}$ and singly produced top quarks is simulated with the POWHEG [91–93] 2.0 and
270 1.0 generators, respectively, at NLO precision [94–97]. Diboson production is generated at LO
271 with PYTHIA 8 [98, 99]. The NNPDF3.0 PDF sets [100] are used for 2016. The NNPDF3.1
272 PDF [24] sets are used for 2017 and 2018 samples.

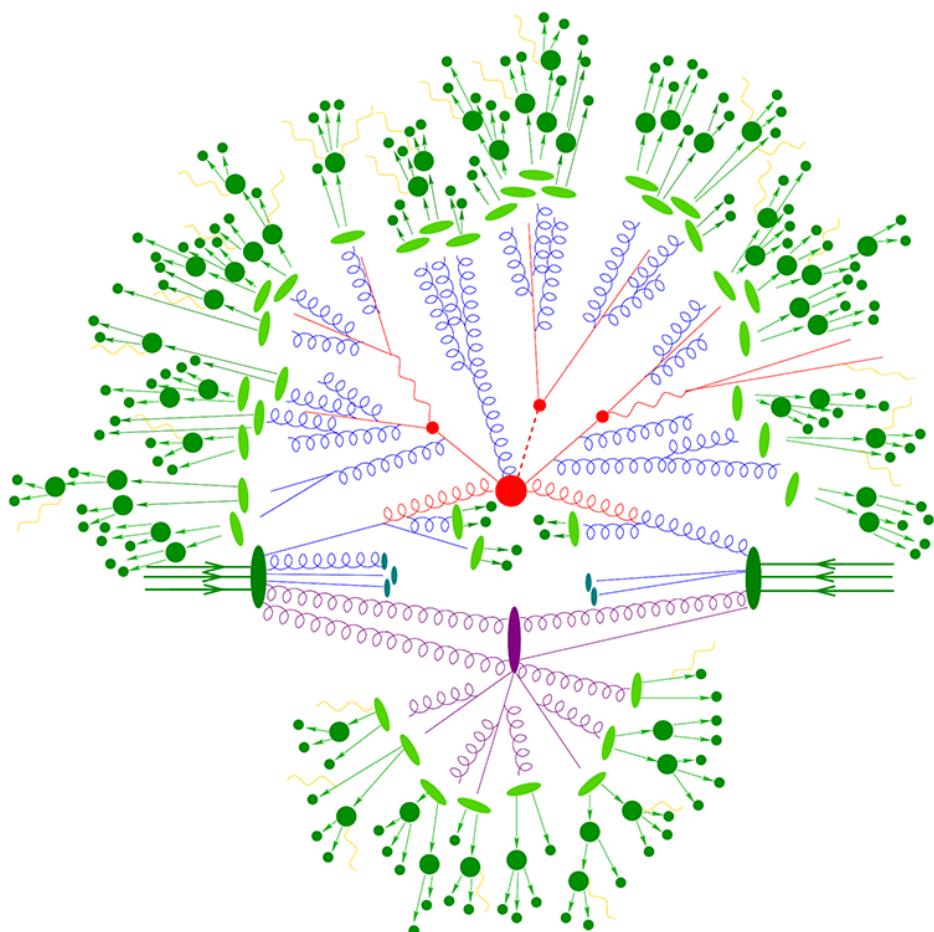


Figure 6.1: Schematic diagram explaining event generation of two colliding protons. Retrieved from [81].

Table 6.2: Summary of simulated SM backgrounds

Process	Generators	Cross section σ [pb]
Drell-Yan, $Z/\gamma^* \rightarrow \ell^+\ell^-$, LO		
+ jets, $10 < m_{\ell\ell} < 50$ GeV	MADGRAPH, PYTHIA	15810.0 (LO), 18610.0 (NLO)
+ jets, $m_{\ell\ell} > 50$ GeV	MADGRAPH, PYTHIA	5343.0 (LO), 6077.2 (NNLO)
+ 1 jets, $m_{\ell\ell} > 50$ GeV	MADGRAPH, PYTHIA	877.8 (LO)
+ 2 jets, $m_{\ell\ell} > 50$ GeV	MADGRAPH, PYTHIA	304.4 (LO)
+ 3 jets, $m_{\ell\ell} > 50$ GeV	MADGRAPH, PYTHIA	111.5 (LO)
+ 4 jets, $m_{\ell\ell} > 50$ GeV	MADGRAPH, PYTHIA	44.05 (LO)
Drell-Yan, $Z/\gamma^* \rightarrow \ell^+\ell^-$, NLO		
+ jets, $m_{\ell\ell} \in [100, 200]$ GeV	aMC@NLO, PYTHIA	247.8 (NLO)
+ jets, $m_{\ell\ell} \in [200, 400]$ GeV	aMC@NLO, PYTHIA	8.502 (NLO)
+ jets, $m_{\ell\ell} \in [400, 500]$ GeV	aMC@NLO, PYTHIA	0.4514 (NLO)
+ jets, $m_{\ell\ell} \in [500, 700]$ GeV	aMC@NLO, PYTHIA	0.2558 (NLO)
+ jets, $m_{\ell\ell} \in [700, 800]$ GeV	aMC@NLO, PYTHIA	0.04023 (NLO)
+ jets, $m_{\ell\ell} \in [800, 1000]$ GeV	aMC@NLO, PYTHIA	0.03406 (NLO)
+ jets, $m_{\ell\ell} \in [1000, 1500]$ GeV	aMC@NLO, PYTHIA	0.01828 (NLO)
+ jets, $m_{\ell\ell} \in [1500, 2000]$ GeV	aMC@NLO, PYTHIA	0.002367 (NLO)
+ jets, $m_{\ell\ell} \in [2000, 3000]$ GeV	aMC@NLO, PYTHIA	5.409×10^{-4} (NLO)
+ jets, $m_{\ell\ell} \in [3000, \infty]$ GeV	aMC@NLO, PYTHIA	3.048×10^{-5} (NLO)
W + jets, $W \rightarrow \ell\nu$		
+ jets	MADGRAPH, PYTHIA	52940.0 (LO), 61526.7 (NLO)
+ 1 jets	MADGRAPH, PYTHIA	8104.0 (LO)
+ 2 jets	MADGRAPH, PYTHIA	2793.0 (LO)
+ 3 jets	MADGRAPH, PYTHIA	992.5 (LO)
+ 4 jets	MADGRAPH, PYTHIA	544.3 (LO)
t̄t + jets		833.9 (NNLO)
Fully leptonic	POWHEG, PYTHIA	88.29 (NNLO), $\mathcal{B} = 10.6\%$
Semi-leptonic	POWHEG, PYTHIA	365.35 (NNLO), $\mathcal{B} = 43.9\%$
Fully Hadronic	POWHEG, PYTHIA	377.96 (NNLO), $\mathcal{B} = 45.4\%$
Single top		
t + W ⁻	POWHEG, PYTHIA	35.85 (NNLO)
̄t + W ⁺	POWHEG, PYTHIA	35.85 (NNLO)
Single t, t channel	POWHEG, PYTHIA	136.02 (NNLO)
Single ̄t, t channel	POWHEG, PYTHIA	80.95 (NNLO)
Diboson		
WW	PYTHIA	75.88 (LO)
WZ	PYTHIA	27.60 (LO)
ZZ	PYTHIA	12.14 (LO)

273 **6.3.2 Cross sections**

274 Table 6.2 also lists the theoretical cross sections σ for each sample. The events are weighted by

$$Z = \frac{L\sigma}{N_{\text{tot}}}, \quad (6.1)$$

275 or with generator weight w_{gen} ,

$$Z = \frac{L\sigma}{\sum w_{\text{gen}}} w_{\text{gen}}. \quad (6.2)$$

276 The LO cross sections of the Z + jets and W + jets samples in Table 6.2 are computed with
277 MADGRAPH. The Z + jets cross section is computed with FEWZ [101] program at NNLO in per-
278 turbative QCD, and with NLO electroweak corrections. The tt}+jets cross section was computed
279 with the TOP++v2.0 program [102, 103] at NNLO and at next-to-next-to-leading logarithmic
280 (NNLL) accuracy. The W + jets production is normalized with cross sections computed at NLO
281 accuracy.

²⁸² 7 Event Selection

²⁸³ This chapter will discuss in more detail the data sets of pp collision events used, as well as the
²⁸⁴ samples of simulated background events.

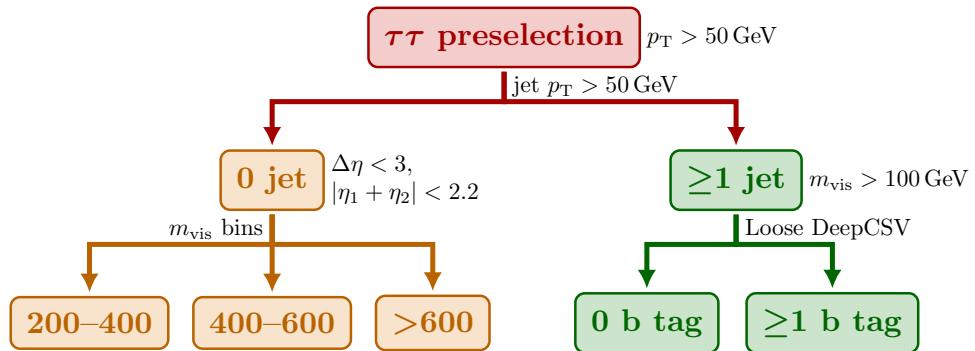


Figure 7.1: Flow chart of event selection.

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