

N-Body Simulations of the Solar System – First Semester Report

Jack Dymond - 150194462

December 6, 2018

Abstract

A literature search was performed, and progress made on creating a 4th order predictor corrector algorithm. The first steps were creating an efficient, and accurate second order N-body code, to help initialise the non-self-starting predictor corrector. A plan was also created, to help organise my time in the second semester.

1 Literature Review

In order to undertake the project it was necessary to first inform myself of the N-Body problem, and its pertinence in Astrophysical problems. To do this, a literature review was undertaken.

1.1 The N-Body Problem

Consider a system of particles, or Bodies, in a dynamic system. In order to understand how the system will behave, the forces acting on each particle need to be ascertained, as well as how this force varies with time. In very simple systems, namely those involving only two particles, there is often an analytical solution to understand how it will behave. Such as the Bohr model of an atom. However, when multiple particles are introduced, more than three to be specific, the problem is impossible to solve analytically, without making simplifying assumptions¹. Hence, the N-Body Problem.

To overcome this, the equations of motion governing the system are instead solved numerically. When the problem was first identified In the early 17th century, solving these equations to resolve long term dynamical behaviour was completely unfeasible. Consequently, it wasn't until the advent of computers that N-Body simulations began to proliferate in the Astrophysics community.

1.2 The N-Body Problem in Astrophysics

In Astrophysics the distance scales being considered are often much greater than those at which the electrostatic forces act, let alone the nuclear. As a result, the gravitational force is often the one that needs to be considered. Hence, astrophysical problems that analyse the dynamical effects of gravity dominated systems, are the ideal candidate for N-Body simulations.

The latter stages of star formation is perhaps the smallest scale at which N-Body simulations become useful. In the former stages, particle sizes are very small, leaving millions of particles to account for, hence, fluid mechanics are a more computationally feasible way of modelling them. However, this changes once in the runaway and oligarchic growth stages: The onset of more significant gravitational forces, and decrease in number of bodies sees N-Body simulations become a much more effective method.

E. Kokubo and S. Ida were among the first to utilise showcase this, in their paper *On Runaway Growth of Planetesimals*[1]. They performed 3-D N-body simulations of 3000 equally massive planetesimals, orbiting a central body. The simulations provided conclusive evidence of naturally occurring runaway growth, something only previously inferred by statistical analysis of the coagulation condition[2, 3]. The results confirmed the validity of a statistical approach to demonstrate runaway growth. It also highlighted

¹Such a Rydberg atom in atomic Physics that makes the assumption of a 'Hydrogenic' atom, simplifying the problem to the Bohr model.

the disadvantages of the statistical approach in later stages of runaway growth, where the planetesimals are grouped spatially: A statistical approach assumes a homogenous distribution of planetesimals.

Once planetary systems are formed they are far from stable, understanding these instabilities, and how they evolve dynamically are another use for N-body simulations.

One example of such instabilities is our early solar system, and studying the dynamism of this system is an active area of research. Namely, unravelling the mysteries of its structure, such as the existence of the Kuiper Belt and Oort cloud. But also explaining its history, in events such as the Late heavy Bombardment. The Nice model solves all of these, and is an example of how N-body simulations can provide answers to problems requiring more specificity to solve.

The nice model, first hypothesised in 2005[4], describes an event in which the outer gas giants, Uranus and Neptune, undergo considerable migrations. The inner gas giants get into a 1:2 resonance, causing dynamical instabilities in the outer orbits. As a result Neptune is propelled past Uranus, as both planets move to their current positions. These migrations disband an early Kuiper belt, pushing asteroids and planetesimals both into the inner system, and out to its far reaches. Thereby explaining both the Oort Cloud, and the Late Heavy Bombardment.

The principal investigative method used in the 2005 paper was the N-Body simulations of an early solar System. K. Tsiganis et al. performed N-Body simulations of the system, including with it a circumstellar disk of planetesimals, truncated at 30 AU. They found orbital resonances in Jupiter and Saturn would cause excitations in their own orbits, and instabilities in those of Uranus and Neptune. Once relaxed, in 50% of the simulations the system would diverge to its current architecture. Perhaps most importantly, the simulations also reproduced the current eccentricities found in the orbits of the Gas Giants, something previously unexplained by planet formation theories.

This was a very high impact paper; since being published it's been cited 717 times. One of which was the extension of the model, namely, its initial conditions.

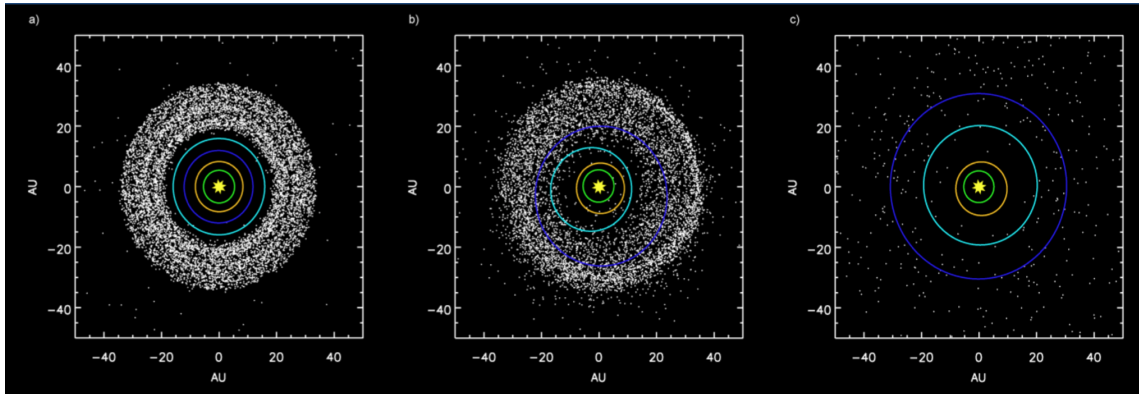


Figure 1: An illustration of the Nice model, with (a) through (c) showing the time evolution of the solar system. It is clear that Uranus (Light Blue) and Neptune (Dark Blue) switch positions in the simulation.

The early solar system is inherently unobservable, hence its initial conditions will always remain a mystery. However, a paper by Batygin et al. aimed to put constraints on the early state of the solar system[5]. They aimed to find *fully resonant* initial conditions, that resulted in the architecture of our current system, post-evolution. Identifying which of these systems evolved in ways that corroborated with the Nice model, was also one of their aims. In their findings they found 4 fully resonant initial conditions that were compatible with the Nice model, and 4 that inferred different evolutions. Their paper explained the importance of multi-resonant systems in the migration of the planets, and made predictions of the eccentricity jumps and evolutionary paths that given resonances would infer. These hypotheses were then supported using N-Body simulations of such fully resonant systems.

With the Nice model being so well supported by subsequent work, it is now important for future developments to concur with this model.

One such development is the advent of ‘Planet 9’, not to be confused with Pluto, Planet 9 refers to a planet several earth sizes in mass, orbiting in the far reaches of the solar system. There are many theories that allow for a 9th planet, such as an ejection, early in the dynamical timescale of the solar system. A much simpler solution however, would be the capture of a 9th planet from the Sun’s birth cluster.

The capture of a 9th planet would be congruent with current theories of the solar system, so long as certain criteria are met, these criteria are detailed at the beginning of a paper by Mustill et al.[6]: The star–star encounter must be at a distance that would not disrupt the Kuiper belt (~ 150 AU); The other star must have a planet on a very wide orbit (≥ 100 AU); Lastly, the planet must be captured onto a trajectory that allows for the current orbital distributions of the gas giants. The paper shows, through N-body simulations of star-star encounters, that all of these criteria could be satisfied in a capture event, albeit at a small probability.

In order for a planet to be captured from a foreign star, the sun would likely need to be in a star cluster at the time of capture. These are once again gravitationally dominated systems, and hence the study of their dynamics has been encapsulated by N-body simulations.

The dynamics between stars in clusters, and that of the planets encompassing said stars, are often intertwined physically. This is true to such an extent, that the structure of a protoplanetary system can infer the dynamics of the cluster in which it was formed.

Perturbations will often effect the orbits of young planetary systems, and these perturbations will persist long after the parent cluster has evaporated. The outer orbits of these systems will relax very slowly, and can hence have the potential to carry signatures of their cluster environment. In a paper published by Cai and his colleagues, this hypotheses is supported in detail[7]. They carry out N-body simulations,

N-body simulations were performed of star clusters of varying stellar mass, between $0.04 M_{\odot}$ and $25 M_{\odot}$. These stars hosted multi-planet systems using two different models, the distribution of semi-major axes remained consistent between these models, starting at 5.2 AU, and extending outwards to ~ 200 AU. The planetary masses were different in each model, one used Jupiter mass planets, the other $\frac{1}{3}$ Earth mass planets. The dynamical simulations aimed to observe a departure from these initial conditions, these departures arising due to both perturbations from stars, and through planet-planet scattering.

They found that the birth environment of a planetary system can be understood from observed orbital elements. They also suggested that the star clusters cause ‘natural selection’ between the planetary systems formed within them, where only the most resilient systems would survive. Furthermore, the more dense stellar environments, such as deep inside the cluster, resulted in stronger signatures in the planetary orbits. With these statements they also provided quantitative relationships between the perturbation method, and the inclination and eccentricity of planetary orbits: They found that the inclination and eccentricity would rise systematically when subjected to external perturbations; Or if internal evolution dominated, the inclination would fluctuate in anti-phase with the eccentricity, this mechanism being dominated by the conservation of angular momentum. All of these findings were true regardless of which mass planet was used. These results also allowed the group to hypothesise the Sun’s birth environment: The outskirts of a star cluster ~ 1000 stars in size.

It follows that the reverse is also true, perhaps by observing stars in clusters, you could infer properties of planetary systems encompassing them, by analysing their location in the cluster.

Photometric properties of clusters can be measured from observations, these properties are often used to predict physical parameters such as mass and radii. However, these are often estimations, and to obtain values with certainty, many parameters need to be measured to a high degree of accuracy.

The N-body method sees another use in supporting values estimated using observations, this is shown by H. Baumgardt, in a paper published in 2016[8]. Baumgardt performs 900 N-body simulations of differently initialised star clusters, their concentration, size, and black hole mass fraction were varied between simulations.

By comparing data obtained from observations, against that provided through N-body simulations, the best fitting model was obtained for each cluster. With these models absolute masses could be derived, and hence their mass-light ratios. The presence of black holes in their simulations allowed predictions to be made, regarding the presence of black holes in the observed clusters. Furthermore, cluster distances could be ascertained for galaxies with more accurate velocity dispersion measurements.

They found theoretically predicted mass-light ratios were in strong agreement with those predicted using N-body methods. For the majority of clusters, comparisons with N-body simulations didn’t support the existence of black holes. One cluster however, ω Cen, strongly disagreed with N-body models that lacked a large black hole mass fraction, suggesting the cluster hosts a black hole at its centre, the paper suggesting a mass of $\sim 40,000 M_{\odot}$.

While being very useful to support findings from observations, the N-body method is most powerful when being used to model the evolution of these environments.

An astrophysical body often under the spotlight is the binary star system, this interest is independent of the environment in which they are in. Clusters by definition have a dense star population, and hence carry the potential to host many binary systems. The binary systems seen in clusters are often in the form of *intermediate* binaries, these are binaries with separations of 10s–100s AU.

The initial mass function in star formation is universal, but it is unclear as to whether this carries to stars found in binaries, as they are frequently separated in clusters due to their often high separations. Parker and Goodwin investigated this[9].

N-body simulations were performed of clusters with 100% binary population, each with separations in the range of 62–620 AU, and with primary and secondary mass functions being derived from previous work in the field. They allowed these clusters to evolve dynamically over a period of 10^6 years, and then reassessed the binary populations. Two cluster morphologies were used, the more sub-structured fractal clusters, and the centrally dominated Plummer spheres. In each cluster environment ten simulations would be performed for a given initial binary population, whilst varying the positions and velocities of the individual binary systems.

After the simulations they found the binary population could decrease by a factor of 2 after 1 Myr, separation distances could also change to “statistically significant” values over this time: In the field these separation distances are sometimes used to infer dynamical history and initial conditions. This paper shows caution should be taken when doing this as even with the full knowledge of initial conditions, the dynamical evolution of clusters often created statistical outliers. That is to say, when observing binary separation distances, you cannot be sure if this is due to initial conditions, or due to a particularly unusual dynamical evolution.

In particularly active environments, not only are stars seen in binaries, but also the clusters themselves. In fact $\sim 10\%$ of clusters are found in binary pairs, these are known as *binary clusters*.

In a paper by B. Arnold et al. the formation of these systems is investigated[10]. The group simulated systems of 1000 stars, to emulate star forming regions, over a period 20 Myr and analysed the structure of the cluster after its dynamical evolution. They varied the initial degree of sub-structuring, as to capture both smooth distributions, and the more hierarchical. Their virial ratios² were also varied, from relatively ‘cool’ values of 0.3, to ‘hot’ values of 0.7. The velocity distribution was such that stars were given velocities similar to that of stars spatially close to them, known as ‘coherence’. As a method of analysing the resultant morphology, three different classifications were given: Binary Clusters, and Single clusters. The final class is less self evident: Binary Merger. In this paper, a binary merger specifies a cluster that was once a binary at some point throughout its evolution, but merged over the 20 Myr period, hence a merger.

They found that binary clusters could form from these star forming regions, and that this would be the case much more frequently if the cluster was initially ‘hot’ dynamically. However, they also noted this was a very stochastic mechanism, and that it was impossible to determine whether a binary would form, and indeed which stars occupy the two clusters, from the initial conditions alone. When binaries did form, they found that they’d often drift apart from one another, and that initial substructure would significantly influence the mass ratio of the binaries. They found high amounts of substructure gave high/low mass ratios, no substructure gave more equal mass ratios, and the most equal mass ratios came from regions with a small amount of initial substructure. In the event that clusters remained bound to one another, they’d recombine later in time, this was mostly the case in systems that began in virial equilibrium.

Clusters are very interesting astrophysical objects, that can provide information about astrophysical problems, from the formation of planets, to the formation and evolution of galaxies. However these areas are often complicated by the dominance of turbulence, something N-body models haven’t consider in the papers discussed.

Galaxies are often dominated by dark matter, something which traditional N-body models can’t simulate, due to its diffuse fluid like behaviour. Being so dominant in many large scale processes, dark matter cannot be ignored at the large scale. Hence, a model was developed to assist in simulations of galaxies, as well as those at the cosmological scale.

²This is the ratio of Kinetic to potential energy

GADGET (**GA**laxies with **D**ark matter and **G**as int**E**rac**T**), is a tool, first developed in 2001[11], to allow N-body simulations to be performed while still considering dark matter and gas distributions. This was achieved by implementing smoothed particle hydrodynamics, a lagrangian technique which is effective in solving hydrodynamical problems, in conjunction with the simulation of individual particles. This could be used with collisionless particles, therefore capturing dark matter in the scope of the code's capabilities.

The code has seen prolific use in the field, quite recently a more advanced version of the code, GADGET-3, was used in simulations of the Milky Way.

The Milky way's disk shows vertical warps, this is highlighted at the beginning of a paper by M. Chequers et al.[12], and the paper presented a method of analysing wave like perturbations in the disk. They also hypothesised that these warps may be caused by inhomogeneities in the dark matter distribution in the halo.

They perform simulations of smoothly distributed halos, as well as those with more substructure, in form of dark matter satellites, 'sub-halos'. They found the bending of the disk was quantitatively similar in both halo distributions, albeit with amplified amplitudes produced when satellites were present in the halos. This is shown in fig. 2, a figure lifted directly from the paper. It is clear the effect is more pronounced in the 'satellites' simulation.

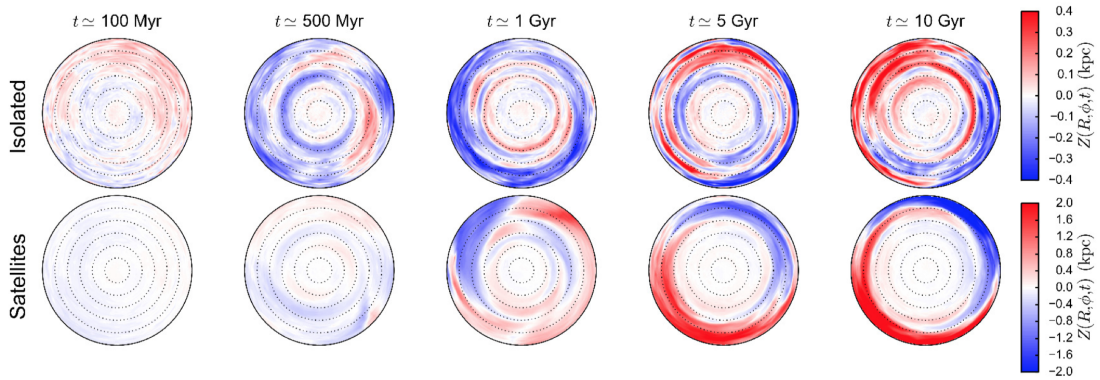


Figure 2: The mean vertical displacement in each simulation is highlighted using a colour coded map of the galaxy.

To better understand the linkage between the halo and the disk, they propose the simulations be used in conjunction with observations, as is so often the case in N-body research in astrophysics.

It is clear from this literature survey that the N-body method is a very powerful tool, being prevalent from the planetary scale, to the cosmological, its scope is unparalleled in this field. The algorithmic complexity necessary to undertake these simulations efficiently is also very high, especially when hydrodynamic systems are implemented in parallel. Hence using these methods in this project will prove invaluable in the pursuit of a career in the academic and industrial environment alike.

2 Progress on Project

There are many ways to approach the N-Body problem to the 4th order. More accurate methods are predictor-correctors, in particular the Hermite Polynomial, and the Adams-Bashforth-Moulton model. The former tries to predict the higher order kinematic components, the aptly named Jerk, Snap, Crack, and Pop. This is computationally expensive, as crossproducts between velocity and acceleration vectors are needed, something relatively inefficient for computer processors to calculate. The Adams-Bashforth-Moulton method predicts future values by fitting a polynomial to previous values, this requires storing previous values of the quantity that needs predicting.

At low values of N, storing these values is more computationally feasible hence, the method is more efficient than the Hermite Method, hence, this project is going to implement the Adams-Bashforth-Moulton model. This method is not self starting, meaning in order to function it requires time-history of the velocity and acceleration vectors. In order to do this, a second order method is necessary to initialise

the values used for a time history, in a method known as ‘Bootstrapping’. Hence, the first part of the project required implementing an accurate 2nd order N-Body.

2.1 Second order Kinematics

In order to calculate a body’s changes in velocity and position over a given time-step, the acceleration vector acting on it must first be calculated. In a system dominated by gravitational forces, this requires knowledge of every body’s position in a consistent coordinate system. For a given body, n , its interaction with another body m , can be calculated using Newton’s Law of gravitation, provided the distance between them, r_{nm} , is known. By summing over all interactions the body will have, its total acceleration vector can be calculated as shown in eq. (1).

$$\vec{a}_n = \sum_{m \neq n}^N \frac{Gm_m}{|r_{nm}|^3} r_{nm} \quad (1)$$

Given a discrete time-step over which to apply this acceleration, the changes in distance and velocity can be determined using the kinematic equations. These are shown below, where the subscripts i and f denote the initial and final positions respectively, δt represents the time-step, and \vec{r}, \vec{v} and \vec{a} , the position, velocity and acceleration.

$$\vec{r}_{n,f} = \vec{r}_{n,i} + \vec{v}_n \delta t \quad \vec{v}_{n,f} = \vec{v}_{n,i} + \vec{a}_n \delta t$$

For the position, this can easily be extended to the second order, by including the current acceleration. This is shown in eq. (2).

$$\vec{r}_{n,f} = \vec{r}_{n,i} + \vec{v}_n \delta t + \frac{1}{2} \vec{a}_n \delta t^2 \quad (2)$$

To calculate the velocity to the second order, the rate of change of the acceleration is required, the jerk. This is not a simple calculation, and is computationally expensive, hence, an approximation should be made: $\frac{da}{dt} \approx \frac{a_i - a_f}{\delta t}$. Given a small change in acceleration this is a valid assumption. Programatically, this requires storing the previous value of acceleration, a_p . Then by using it in conjunction with the current acceleration, a_c , the second order equation for the velocity can now be taken as eq. (3).

$$\vec{v}_{n,f} = \vec{v}_{n,i} + \vec{a}_n \delta t + \frac{1}{2} \frac{d\vec{a}_n}{dt} \delta t^2 = \vec{v}_{n,i} + \vec{a}_{n,c} \delta t + \frac{1}{2} (\vec{a}_{n,p} - \vec{a}_{n,c}) \delta t \quad (3)$$

2.2 Recreating the Solar System accurately

To undertake N-Body simulations, these are the only equations necessary. However, to start the process, initial conditions should be given to each body, namely a mass, position, and velocity. These were taken from a NASA database[13], see appendix A for the values pertinent in this project.

To initialise the simulation, the planets were positioned on the x axis at their orbital radii, and at $y = 0$. Their respective orbital velocities were then applied along the y axis, with $v_x = 0$. Hence, when the simulation was run they would orbit on the x - y plane.

At the beginning of each time-step, the acceleration acting on each body was calculated using eq. (1), this value was then applied to the kinematic equations, eq. (2), and eq. (3). The planets were moved to these positions, and given the corresponding velocities. New accelerations were calculated, and the process was repeated. Since no time history was given, initially the change in acceleration couldn’t be calculated in eq. (3). Hence, only the first derivative of velocity was used in the calculation of $\vec{v}_{n,2}$, this was done by manually setting the change in acceleration to zero.

To demonstrate the accuracy of this second order code, a 1000 year simulation of the solar system was carried out with 10 second time-steps. The resulting simulation produced a system with stable orbits, this is reflected in fig. 3. Figure 3 (a) shows the log of orbital distance against time, it is clear the orbits remain at a consistent distance with only Earth drifting very slightly. The oscillations in the distances of Saturn,

Uranus, and Neptune, are due to the elliptical nature of their orbits, the eccentricities were reasonable: ~ 0.02 , ~ 0.04 , and ~ 0.03 respectively. Figure 3 (b) shows the orbits in the x - y plane, demonstrating their circular nature.

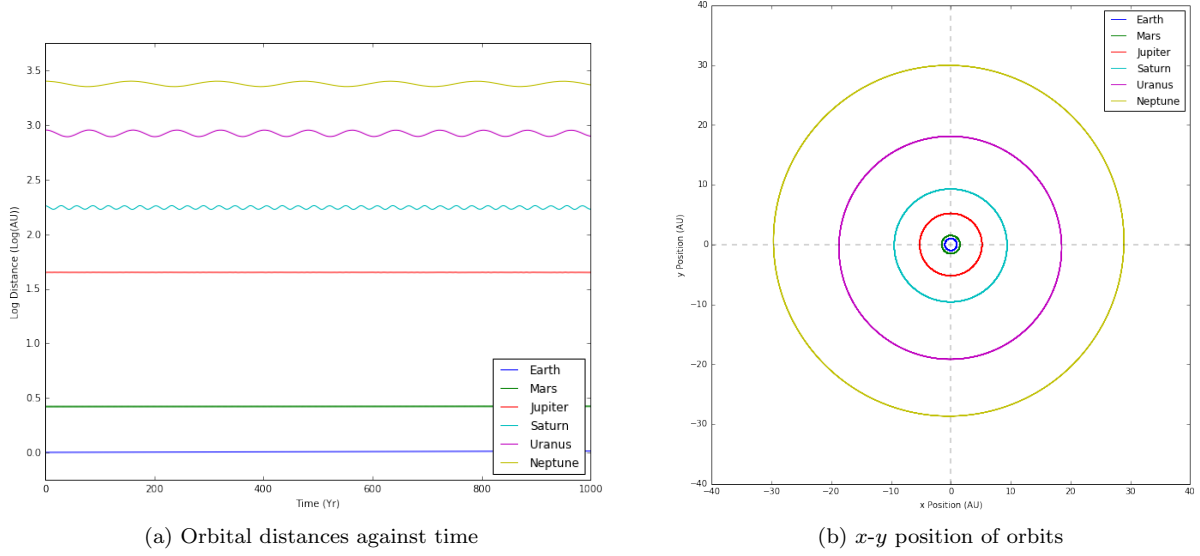


Figure 3: Plots of orbital position throughout the simulation.

To further ensure the model's accuracy, energy conservation was considered. The total energy was calculated when the system was initialised to obtain E_i , this was simply the sum of the kinetic and potential energies, as shown in eq. (4). A factor of a half was applied to the potential energy, as in this summation it would be calculated twice for a given interaction $n \leftrightarrow m$.

$$E_{tot} = \sum_{n=1}^N \left[\frac{1}{2} m_n |v_n|^2 - \sum_{m \neq n}^N \frac{G m_m m_n}{2 |r_{nm}|} \right] = \frac{1}{2} \sum_{n=1}^N m_n \left[|v_n|^2 - \sum_{m \neq n}^N \frac{G m_m}{|r_{nm}|} \right] \quad (4)$$

This value was recalculated throughout the simulation, and compared against the initial value, using eq. (5). Where ΔE , represents the fractional energy loss, and the subscript t denoting the time-step at which the value is calculated.

$$\Delta E_t = \frac{E_i - E_t}{E_i} \quad (5)$$

This fractional energy loss was plotted against time, to see how it varied throughout the simulation. This is shown in fig. 4. Figure 4 (a) shows the variation across the entire simulation, whereas fig. 4 (b) shows the value over a 100 year time period.

Over the whole simulation, the average energy loss remains close to zero, only lowering slightly, mainly due to the drift of Earth. However, there are clearly oscillatory fluctuations in the value, these oscillations are due to interactions of Jupiter with Saturn. Upon increasing the time-resolution of this graph, fig. 4 (b), smaller scale oscillations can be seen. These are presumably a result of interactions between the other planets in the system. Nonetheless, the energy losses (and its fluctuation) proved to be minor in this simulation; the large scale oscillation having an amplitude of ~ 0.0003 , and only lowering by ~ 0.0002 .

2.3 Suitability for Bootstrapping

The previous section demonstrates the code's ability to accurately recreate the orbits of the Solar System, to a reasonable degree of accuracy. However, if a simulation is to run for millions of years, the discrepancies seen in the 2nd order system will manifest as large inaccuracies. Hence, the 4th order Adams-Bashforth-Moulton will be preferred.

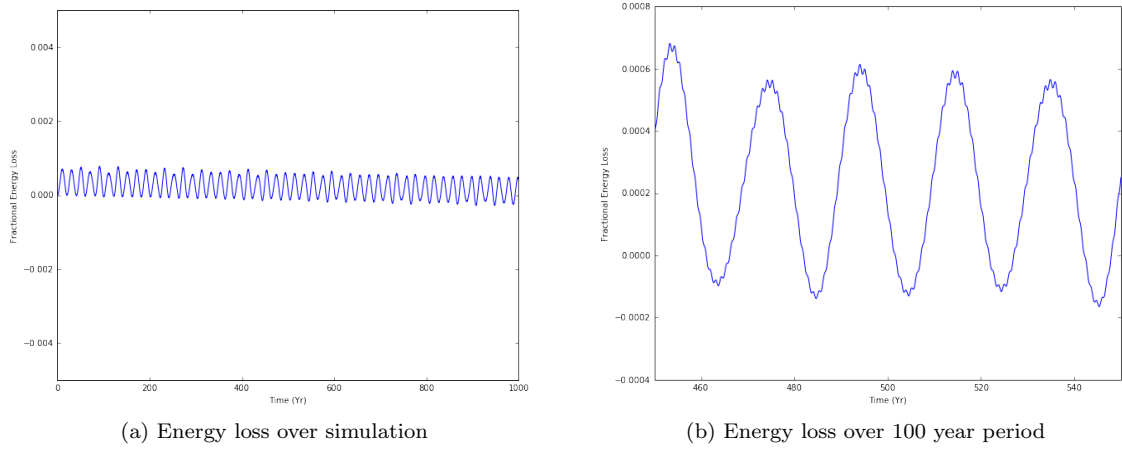


Figure 4: Plots showing the Energy loss throughout the simulation.

As a method of producing the initial time history, the 2nd order code is more than sufficient; over a small timescale, < 80 seconds, the errors seen at long timescales will be infinitesimal. Accordingly, this code will be used initially in the ‘bootstrapping’ phase.

2.4 Randomising Initial Positions Along Orbit

If the simulation is to be run more than once, it will be beneficial to randomise the initial positions of the planets along their orbit. Should the code be applied to astrophysical problems, this would prevent any deterministic results from arising, thereby producing more natural behaviour.

There are many ways to approach this, some requiring more knowledge of the orbital paths than others. However, with knowledge of the mean orbital radius and velocity, the planet can be placed at any position at a radius equal to the orbital radius. Provided the velocity vector is both perpendicular to the radial vector, and equal to the mean orbital velocity.

Consider the positive quadrant of a planet’s orbit (see fig. 3 (b)). On a given axis, the position will take a value between 0 and R , and the velocity will take a value between 0 and v_T , where R and v_T denote the orbital radius and velocity respectively. The component on the other axis can then be calculated using trigonometry. This is represented mathematically as shown in eq. (6), where X_n represents the random number between 0 and 1, assigned to the body, n .

$$r_{x,n} = X_n R_n \quad \text{and} \quad v_{y,n} = X_n v_{T,n} \quad (6)$$

Then by using Trigonometry, the components along the other axes can be calculated, as shown in eq. (7). Note the opposite axis on r_n , and v_n respectively in eq. (6), this is due to the velocity vector being perpendicular to the radius vector.

$$r_{y,n} = R_n^2 - r_{x,n}^2 \quad \text{and} \quad v_{x,n} = v_{T,n}^2 - v_{y,n}^2 \quad (7)$$

To allow for the other three quadrants, a random factor of ± 1 should be applied to each component of the position. The value applied to the velocity will depend on the value applied to the radius vector. That is to say, the quadrant allocated by the random factors of ± 1 will determine the value of ± 1 applied to the velocity components.

This was achieved programatically by first randomly assigning a value of ± 1 to each position vector component, a parity. Then, using a simple series of *IF* statements, the velocity components were allocated parities, thus creating two 1×2 arrays for each quantity. The conditions for the algorithm are illustrated in table 1.

Regardless of which quadrant they’re in, the same trigonometric relations in eq. (7) exist between the quantities, and their components. Hence, with a random number allocated to each position component, the planet can be randomly positioned in any quadrant, and in any position along its orbit.

| Position [x,y] | Velocity [v_x, v_y] |
|----------------|-------------------------|
| [1,1] | [-1,1] |
| [-1,1] | [-1,-1] |
| [-1,-1] | [1,-1] |
| [1,-1] | [1,1] |

Table 1: Position parity vectors, with corresponding velocity parity vectors.

This functionality was added to the code, and the resultant behaviour was identical to that seen in fig. 3 and fig. 4.

2.5 Moving Forward

As section 2.3 states, this code will be more than sufficient for initialising the first 8 time-steps of a predictor-corrector. Furthermore, with planets being positioned at completely random positions along their orbit, there is no chance of a deterministic effects arising from the initial setup of the system.

Hence, the project can continue towards the goal of creating a 4th Order Predictor Corrector. The logistics of this are detailed in the next section.

3 Project Plan

I'll produce a gantt chart, and explain it in this section. Although, I've not got round to doing this yet!

References

- [1] Kokubo, E., and S. Ida
On Runaway growth of Planetesimals: N-body simulation
Icarus 1996, 123, 180–191.
- [2] Barge, P., and R. Pellat
Mass spectrum and velocity dispersions during planetesimal accumulation. I. Accretion
Icarus 1991 93, 270–287
- [3] Philip J. Armitage
Astrophysics of Planet Formation
Sec. 4.5.1, Pg. 131
10 Dec. 2009
- [4] K. Tsiganis, R. Gomes, A. Morbidelli and H. F. Levison
Origin of the orbital architecture of the giant planets of the Solar System
Nature, Vol 435, P 469–451, 26 May 2005
- [5] Konstantin Batygin and Michael E. Brown
Early Dynamical Evolution of the Solar System: Pinning Down the Initial Conditions of the Nice Model
ApJ, 716:1323–1331, 2010 June 20
- [6] Alexander J. Mustill, Sean N. Raymond, and Melvyn B. Davies
Is there an exoplanet in the Solar System?
MNRAS Issue 1, p.109–113
- [7] Maxwell Xu Cai, Simon Portegies Zwart, and Arjen van Elteren
The signatures of the parental cluster on field planetary systems
MNRAS, V. 474, Issue 4, 11 March 2018, Pages 5114–5121
- [8] H. Baumgardt
N-body modelling of globular clusters: masses, mass-to-light ratios and intermediate-mass black holes
MNRAS, V. 464, Issue 2, 11 January 2017, P 2174–2202.
- [9] Richard J. Parker. and Simon P. Goodwin
The same, but different: Stochasticity in binary destruction
MNRAS, V. 424, Issue 1, 21 July 2012, P. 272–281

- [10] Becky Arnold, Simon P. Goodwin, D. W. Griffiths and Richard. J. Parker
How do binary clusters form?
MNRAS, V. 471, Issue 2, 21 October 2017, P. 2498–2507
- [11] Volker Springel, Naoki Yoshida, and Simon D. M. White
GADGET: a code for collisionless and gasdynamical cosmological simulations
New Astronomy, V. 6, Issue 2, April 2001, P. 79–117
- [12] Matthew H Chequers, Lawrence M Widrow, and Keir Darling
Bending waves in the Milky Way’s disc from halo substructure
MNRAS, V. 480, Issue 3, 1 November 2018, P. 4244–4258
- [13] Williams, David
TPlanetary Fact Sheet - Metric
<https://nssdc.gsfc.nasa.gov/planetary/factsheet/>
As of: December 6, 2018

A Initial Conditions for Solar System

| Planet | Mass (10^{24} kg) | Initial Orbital Velocity (km s^{-1}) | Initial Orbital Radius (AU) |
|---------|----------------------|-------------------------------------------------|-----------------------------|
| Earth | 5.97 | 29.8 | 1.00 |
| Mars | 0.642 | 24.1 | 1.41 |
| Jupiter | 1898 | 13.1 | 5.03 |
| Saturn | 568 | 9.7 | 9.20 |
| Uranus | 86.8 | 6.8 | 18.64 |
| Neptune | 102 | 5.4 | 30.22 |

Table 2: Initial Conditions taken from [13]

B Relationship to previous projects

In my second semester 3rd year project, I used the N-body method to analyse the fragility of the solar system to external perturbations. This used code very similar to that in the first half of this project. Hence, I focussed on streamlining the process, making it more efficient, and increasing its functionality. I also focussed on creating a framework to implement it with a 4th order corrector, rather than collecting and collating results from simulations.

Hence, while the projects overlap in the fundamental method used, my focus has been different in each project. With this 4th year project concentrating much more on the code itself, rather than implementing it to an astrophysical problem.