Numerical Methods for Hyperbolic PDEs Homework 2

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Exercise 1

Let $U_i^n := u(t_n, x_i)$ be the exact solution at (t_n, x_i) , assuming to exist, and let \tilde{U}_i^n the numerical approximation. By Taylor expansion we know that

$$U_i^{n+1} = U_i^n + \Delta t(U_i^n)_t + \frac{\Delta t^2}{2}(U_i^n)_{tt} + \mathcal{O}(\Delta t^3)$$

$$U_{i+1}^n = U_i^n + \Delta x(U_i^n)_x + \frac{\Delta x^2}{2}(U_i^n)_{xx} + \frac{\Delta x^3}{6}(U_i^n)_{xxx} + \mathcal{O}(\Delta x^4)$$

$$U_{i-1}^n = U_i^n - \Delta x(U_i^n)_x + \frac{\Delta x^2}{2}(U_i^n)_{xx} - \frac{\Delta x^3}{6}(U_i^n)_{xxx} + \mathcal{O}(\Delta x^4).$$

We want to calculate the local error, so we assume $U_i^n = \tilde{U}_i^n$ for all i. Then the one-step error is

$$\begin{split} \tilde{U}_{i}^{n+1} - U_{i}^{n+1} \\ &= U_{i}^{n} - \frac{\lambda}{2}(U_{i+1}^{n} - U_{i-1}^{n}) + \frac{\lambda^{2}}{2}(U_{i+1}^{n} - 2U_{i}^{n} + U_{i-1}^{n}) - U_{i}^{n+1} \\ &= U_{i}^{n} - \frac{a\Delta t}{2\Delta x} \Big(2\Delta x (U_{i}^{n})_{x} + \mathcal{O}(\Delta x^{3}) \Big) + \frac{a^{2}\Delta t^{2}}{2\Delta x^{2}} \Big(\Delta x^{2} (U_{i}^{n})_{xx} + \mathcal{O}(\Delta x^{4}) \Big) \\ &- U_{i}^{n} - \Delta t (U_{i}^{n})_{t} - \frac{\Delta t^{2}}{2} (U_{i}^{n})_{tt} + \mathcal{O}(\Delta t^{3}) \\ &= \underbrace{-a\Delta t (U_{i}^{n})_{x} - \Delta t (U_{i}^{n})_{t}}_{=-\Delta t \Big((U_{i}^{n})_{t} + a(U_{i}^{n})_{x}) \Big) = 0} \\ &= \underbrace{-\frac{a\Delta t}{2} \Big(a^{2} (U_{i}^{n})_{xx} - \frac{\Delta t^{2}}{2} (U_{i}^{n})_{tt} + \mathcal{O}(\Delta t^{2}\Delta x^{2} + \Delta t^{3}) \\ = \underbrace{-\frac{a\Delta t}{2} \Big(a^{2} (U_{i}^{n})_{xx} - (U_{i}^{n})_{tt} \Big) = 0} \\ &= \mathcal{O}(\Delta t^{3}), \end{split}$$

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where we used in the last equality that $\mathcal{O}(\Delta t \Delta x^2 + \Delta t^2 \Delta x^2 + \Delta t^3) = \mathcal{O}(\Delta t^3)$, since $\lambda = a \frac{\Delta t}{\Delta x} = const.$ and that

$$u_{tt} = -au_{xt} = (-au_t)_x = (a^2u)_{xx}. (1)$$

Thus, it follows for the local truncation error τ_n

$$\tau_n = \frac{1}{\Delta t} (\tilde{U}_i^n - U_i^n) = \mathcal{O}(\Delta t^2).$$

Therefore, the method is of second order.

We perform the von-Neumann stability analysis: We assume periodic boundary conditions and expand U^n_j as Fourier series. Write $u(x,t) = \sum_{k=-\infty}^{\infty} \hat{u}_k(t) e^{-ikx}$ and let $\hat{U}^n_k = \hat{u}_k(t_n)$. Then

$$U_j^n = u(x_j, t_n) = \sum_k \hat{U}_k^n e^{-ikj\Delta x}$$

$$U_{j+1}^n = u(x_{j+1}, t_n) = \sum_k \hat{U}_k^n e^{-ik(j+1)\Delta x}$$

$$U_{j-1}^n = u(x_{j-1}, t_n) = \sum_k \hat{U}_k^n e^{-ik(j-1)\Delta x}.$$

Plugging this into the Lax-Wendroff scheme, we get

$$\begin{split} U_j^{n+1} &= \sum_k \hat{U}_k^{n+1} e^{-ikj\Delta x} \\ &= \sum_k \hat{U}_k^n e^{-ikj\Delta x} - \frac{\lambda}{2} \Big(\sum_k \hat{U}_k^n (e^{-ik(j+1)\Delta x} - e^{-ik(j-1)\Delta x}) \Big) \\ &+ \frac{\lambda^2}{2} \Big(\sum_k \hat{U}_k^n (e^{-ik(j+1)\Delta x} - 2e^{-ikj\Delta x} + e^{-ik(j-1)\Delta x}) \Big) \\ &= \sum_k e^{-ikj\Delta x} \Big(\hat{U}_k^n - \frac{\lambda}{2} \hat{U}_k^n (e^{-ik\Delta x} - e^{ik\Delta x}) + \frac{\lambda^2}{2} \hat{U}_k^n (e^{-ik\Delta x} - 2 + e^{ik\Delta x}) \Big). \end{split}$$

By projecting onto Fourier basis functions we obtain, for all $k \in \mathbb{Z}$,

$$\hat{U}_k^{n+1} = \hat{U}_k^n \underbrace{\left(1 - \frac{\lambda}{2}(e^{-ik\Delta x} - e^{ik\Delta x}) + \frac{\lambda^2}{2}(e^{-ik\Delta x} + e^{ik\Delta x} - 2)\right)}_{:=A_k}.$$

For stability, we need $|A_k| \leq 1$ for all k. Using

$$\sin(k\Delta x) = \frac{1}{2i}(e^{-ik\Delta x} - e^{ik\Delta x})$$

and

$$\cos(k\Delta x) = \frac{1}{2}(e^{-ik\Delta x} + e^{ik\Delta x}),$$

we get with substituting $\omega_k := k\Delta x$:

$$A_k = 1 - i\lambda\sin(\omega_k) + \lambda^2(\cos(\omega_k) - 1) = 1 - 2i\lambda\sin(\frac{\omega_k}{2})\cos(\frac{\omega_k}{2}),$$

where we used the identities

$$\sin(2x) = 2\sin(x)\cos(x)$$

and

$$\cos(2x) = 1 - \sin^2(x).$$

Therefore,

$$\begin{aligned} |A_k|^2 &= (1 - 2\lambda^2 \sin^2(\frac{\omega_k}{2}))^2 + 4\lambda^2 \sin^2(\frac{\omega_k}{2}) \cos^2(\frac{\omega_k}{2}) \\ &= 1 + 4\lambda^4 \sin^4(\frac{\omega_k}{2}) - 4\lambda^2 \sin^2(\frac{\omega_k}{2}) + 4\lambda^2 \sin^2(\frac{\omega_k}{2}) \cos^2(\frac{\omega_k}{2}) \\ &\stackrel{!}{\leq} 1 \end{aligned}$$

Rewriting this gives

$$0 \ge 4\lambda^2 \sin^2(\frac{\omega_k}{2}) \left(\lambda^2 \sin^2(\frac{\omega_k}{2}) \underbrace{-1 + \cos^2(\frac{\omega_k}{2})}_{=-\sin^2(\frac{\omega_k}{2})}\right)$$
$$= 4\lambda^2 \sin^4(\frac{\omega_k}{2})(\lambda^2 - 1).$$

This is equivalent to $|\lambda| \le 1$. Thus, $|\lambda| \le 1$ is sufficient as this implies $|A_k| \le 1$ for all k.

Exercise 2

We modify the order calculation from Exercise 1 to include higher order terms:

$$\begin{split} \tilde{U}_{i}^{n+1} - U_{i}^{n+1} \\ &= U_{i}^{n} - \frac{\lambda}{2}(U_{i+1}^{n} - U_{i-1}^{n}) + \frac{\lambda^{2}}{2}(U_{i+1}^{n} - 2U_{i}^{n} + U_{i-1}^{n}) - U_{i}^{n+1} \\ &= U_{i}^{n} - \frac{a\Delta t}{2\Delta x} \Big(2\Delta x (U_{i}^{n})_{x} + \frac{\Delta x^{2}}{3}(U_{i}^{n})_{xxx} + \mathcal{O}(\Delta x^{5}) \Big) + \frac{a^{2}\Delta t^{2}}{2\Delta x^{2}} \Big(\Delta x^{2}(U_{i}^{n})_{xx} + \mathcal{O}(\Delta x^{4}) \Big) \\ &- U_{i}^{n} - \Delta t (U_{i}^{n})_{t} - \frac{\Delta t^{2}}{2}(U_{i}^{n})_{tt} - \frac{\Delta t^{3}}{6}(U_{i}^{n})_{ttt} + \mathcal{O}(\Delta t^{4}) \\ &= \underbrace{-a\Delta t (U_{i}^{n})_{x} - \Delta t (U_{i}^{n})_{t}}_{= -\Delta t} - \frac{a}{6}\Delta t \Delta x^{2}(U_{i}^{n})_{xxx} + \underbrace{\frac{a^{2}\Delta t^{2}}{2}(U_{i}^{n})_{xx} - \frac{\Delta t^{2}}{2}(U_{i}^{n})_{tt}}_{= -\Delta t \Big((U_{i}^{n})_{t} + a(U_{i}^{n})_{x}\Big) = 0} \\ &- \frac{\Delta t^{3}}{6}(U_{i}^{n})_{ttt} + \mathcal{O}(\Delta t^{2}\Delta x^{2} + \Delta t^{4}) \\ &= -\frac{a}{6}\Delta t \Delta x^{2}(U_{i}^{n})_{xxx} - \frac{\Delta t^{3}}{6}(U_{i}^{n})_{ttt} + \mathcal{O}(\Delta t^{2}\Delta x^{2} + \Delta t^{4}) \\ &= \Delta t^{3}\frac{a^{3}}{6}(1 - \frac{1}{\lambda^{2}})(U_{i}^{n})_{xxx} + \mathcal{O}(\Delta t^{4}), \end{split}$$

where we used in the last equality that $u_{ttt} = -a^3 u_{xxx}$ which we can derive by Eq. (1), using

$$u_{ttt} = (a^2 u_t)_{xx} = -a^3 u_{xxx}$$

We see that the leading error term depends on $(U_i^n)_{xxx}$. This means the scheme is dispersive as discussed in the lecture.

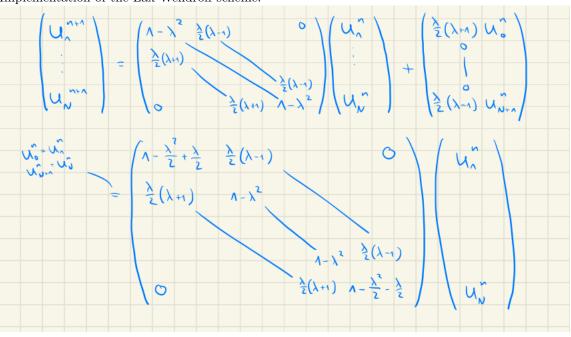
Exercise 3

As discussed in the lecture, the solution of the linear advection equation is

$$u(x,t) = u_0(x-2t) = \begin{cases} -1, & x < 2t \\ 1, & x \ge 2t. \end{cases}$$

As we see in the animation, the solution Lax-Wendroff scheme spreads out over time in a wavy pattern at the discontinuity of the exact solution, whereas in the upwind scheme the numerical solution smoothes out over time. This makes sense since the LW-scheme is dispersion dominant and the upwind scheme is diffusion dominant.

Implementation of the Lax-Wendroff scheme:



Implementation of the upwind scheme:

