Neuro-symbolic hybridization of finite discretization methods for solving partial differential equations

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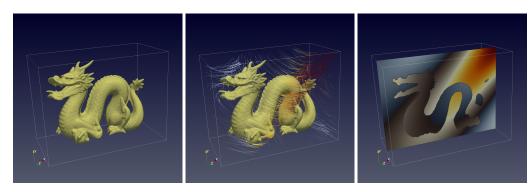


Figure 1: The proposed Neural Bootstrapping Method applied to computing the electrostatic field around a complex charged geometry in a complex dielectric. The neural approach readily enables a 8192^3 effective resolution on a GPU cluster and captures the physically correct discontinuities. Geometry (left), streamlines (center), cross-section showing jump in solution and/or gradient.

Abstract

We present a highly scalable strategy for developing mesh-free hybrid neuro-symbolic partial differential equation (PDE) solvers based on existing mesh-based numerical discretization methods. Particularly, this strategy can be used to efficiently train neural network surrogate models for the solution functions and operators of PDEs while retaining the accuracy and convergence properties of the state-of-the-art numerical solvers. The presented neural bootstrapping method (hereby dubbed NBM) is based on minimizing residuals of discretized PDE system on an array of implicit Cartesian cells at different levels of resolution centered on a set of random collocation points with respect to trainable parameters of the neural network. NBM leverages neural networks to achieve unprecedented resolution and flexibility for solving PDEs that describe complex physical systems.

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1 Introduction

Most modern physical and engineering systems are described by partial differential equations on irregular, often moving, boundaries. The difficulties in solving those problems stem from how to approximate the equations, while respecting the physically correct discontinuous nature of the solution across the boundaries. Smoothing strategies are easy to design, but unfortunately introduce unphysical characteristics in the solution and lead to systemic errors.

Since early 1990s, artificial neural networks have been used for solving partial differential equations by (i) mapping the algebraic operations of the discretized PDE systems onto specialized neural network architectures and minimizing the network energy, or (ii) treating the whole neural network as the basic approximation unit whose parameters are adjusted to minimize a specialized error function that includes the differential equation itself with its boundary/initial conditions.

In the first category, neurons output the discretized solution values over a set number of grid points and minimizing the network energy drives the neuronal values towards the solution of the linear system at the mesh points. In this case, the neural network energy is the residual of the finite discretization method summed over all neurons of the network [20]. Although the convergence properties of the finite discretization methods guarantee and control quality of the obtained solutions, the computational costs grow by increasing resolution and dimensionality. Some early examples include [15, 10, 9].

The second strategy proposed by Lagaris *et al.* [19] relies on the function approximation capabilities of the neural networks. Encoding the solution everywhere in the domain within a neural network offers a mesh-free, compact, and memory efficient surrogate model for the solution function that can be utilized in subsequent inference tasks. This method has recently re-emerged as the physics-informed neural networks (PINNs) [31] and is widely used. Despite their advantages, these methods lack controllable accuracy and convergence properties of finite discretization methods, and are biased towards lower frequency features of the solutions [37, 30].

Pursuit of hybrid solvers aims at leveraging the performance gains of neural network inference on modern accelerated hardware with the guaranteed accuracy of finite discretization methods. The hybridization efforts are algorithmic or architectural.

One important algorithmic method is the deep Galerkin method (DGM) [34] that is a neural network extension of the mesh-free Galerkin method where the solution is represented as a deep neural network rather than a linear combination of basis functions. The mesh-free nature of DGM, that stems from the underlying mesh-free Galerkin method, enables solving problems in higher dimensions by training the neural network model to satisfy the PDE operator and its initial and boundary conditions on a randomly sampled set of points rather than on an exponentially large grid. Although the number of points is huge in higher dimensions, the algorithm can process training on smaller batches of data points sequentially. Besides, second order derivatives in PDEs are calculated by a Monte Carlo method that retain scaling to higher dimensions. Another important algorithmic method is the deep Ritz method for solving variational problems [38] that implements a deep nerual network approximation of the trial function that is constrained by numerical quadrature rule for the variational functional, followed by stochastic gradient descent.

Architectural hybridization methods are based on differentiable numerical linear algebra. One emerging class involves implementing differentiable finite discretization solvers and embedding them in the neural network architectures that enable application of end-to-end differentiable gradient based optimization methods. Recently, differentiable solvers have been developed in JAX [7] for fluid dynamic problems, such as Phi-Flow [17], JAX-CFD [18], and JAX-FLUIDS [2]. These methods are suitable for inverse problems where an unknown field is modeled by the neural network, while the model influence is propagated by the differentiable solver into a measurable residual [29, 12, 24]. We also note the classic strategy for solving inverse problems is the adjoint method to obtain the gradient of the loss without differentiation across the solver [1]; however, deriving analytic expression for the adjoint equations is tedious, should be repeated after modification of the problem or its loss function, and can become impractical for multiphysics problems. Other important utilities of differentiable solvers are to model and correct for the solution errors of finite discretization methods [36], learning and controling PDE systems [13, 16].

Neural networks are not only universal approximators of continuous functions, but also of nonlinear operators [8]. Although this fact has been leveraged using data-driven strategies for learning differential operators by many authors [23, 3, 22, 21], current authors have demonstrated utility of differentiable solvers to effectively train nonlinear operators without any data in a completely physics-driven fashion, see section on learning the inverse transforms in [29].

In this work we propose a novel framework for solving PDEs based on deep neural networks by lifting any existing mesh-based finite discretization method off of its underlying grid and extend it into a mesh-free method that can be applied to high dimensional problems on unstructured random points in an embarrasingly parallel fashion. In addition, discontinuous solutions can be readily considered.

2 Problem statement

In order to illustrate our approach, we consider a closed irregular interface (Γ) that partitions the computational domain (Ω) into interior (Ω^-) and exterior (Ω^+) subdomains; *i.e.*, $\Omega = \Omega^- \cup \Gamma \cup \Omega^+$. We are interested in the solutions $u^{\pm} \in \Omega^{\pm}$ to the following class of linear elliptic problems in $\mathbf{x} \in \Omega^{\pm}$:

$$k^{\pm}u^{\pm} - \nabla \cdot (\mu^{\pm}\nabla u^{\pm}) = f^{\pm}, \qquad \mathbf{x} \in \Omega^{\pm}$$
$$[u] = \alpha, \qquad \mathbf{x} \in \Gamma$$
$$[\mu \partial_{\mathbf{n}} u] = \beta, \qquad \mathbf{x} \in \Gamma$$

Here $f^\pm=f(\mathbf{x}\in\Omega^\pm)$, $\mu^\pm=\mu(\mathbf{x}\in\Omega^\pm)$ and k^\pm are the spatially varying source term, diffusion coefficients, and coefficients, respectively in the two domains. For simplicity, we consider Dirichlet boundary conditions at the boundary of a cubic domain $\Omega=[-L/2,L/2]^3$, noting that other boundary conditions can be readily considered.

This class of problems not only captures the difficulties of solving partial differential equations on irregular domain with discontinuous solutions, but also illustrate their multiscale nature that requires a nonuniform sampling in space. In addition, this system of equations is important on its own as it arises ubiquitously in describing diffusion dominated processes in physical systems and in the life sciences, where sharp and irregular interfaces regulate transport across regions with different properties. Examples include Poisson-Boltzmann equation for describing electrostatic properties of membranes, colloids and solvated biomolecules with jump in dielectric permitivities [33, 25], in electroporation of cell aggregates with nonlinear membrane jump conditions [27], or epitaxial growth in fabrication of opto-electronic devices [26]. Other important applications are found in solidification of multicomponent alloys [35, 6], directed self-assembly of diblock copolymers for next generation lithography [14, 28, 5], multiphase flows with and without phase change.

3 Scalable & Mesh-Freeing Neuro-Symbolic PDE Solver

Neural networks are used as surrogates for the solution function that are iteratively adjusted to minimize discretization residuals at a set of randomly sampled points and at arbitrary resolutions. The key idea is that neural networks can be evaluated over vertices of any discretization stencils centered at any point in the domain effectively emulating the effect of any structured mesh without ever materializing the mesh. Therefore, we use neural networks to bootstrap mesh based finite discretization (FD) methods to compile mesh free numerical methods. We call this the Neural Bootstrapping Method (NBM), and illustrate it in figure 2.

FD methods offer guaranteed accuracy and controllable convergence properties for the training of neural network surrogate models of physical systems. NBM offers a straightforward path for applying mesh-based FD methods on unstructured random points. This is an important ability for augmenting observational data in the training pipelines. Beside its mesh-freeing advantage, NBM is a highly parallelizable strategy for FD methods. Pointwise nature of its kernels is ideally suited for GPU-accelerated computing paradigm. Multi-GPU parallel solution of PDE systems is reduced to the much simpler problem of data-parallel training using existing machine learning frameworks. Data parallelism involves distributing collocation points across multiple processors to compute gradient updates and then aggregating these locally computed updates [32].

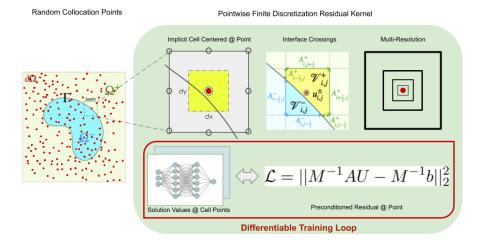


Figure 2: Neural Bootstrapping Method (NBM). NBM kernels compute residual contribution by each collocation point per GPU thread. Kernel operations involve implicit cells at different resolutions according to the bootstrapped finite discretization method. Geometric information for interface-cell crossings is computed by a level-set interpolant defined on a much lower resolution grid. Pointwise residuals are locally preconditioned based on the geometry of the interface crossing the implicit cells.

Here we bootstrap the numerical scheme proposed by [4], hence the loss function is

$$\mathcal{L} = \left| \left| \sum_{s=-,+} k_{i,j}^s u_{i,j}^s | \mathcal{V}_{i,j}^s | - \sum_{s=-,+} \left(\mu_{i-\frac{1}{2},j}^s A_{i-\frac{1}{2},j}^s \frac{u_{i-1,j}^s - u_{i,j}^s}{\Delta x} + \mu_{i+\frac{1}{2},j}^s A_{i+\frac{1}{2},j}^s \frac{u_{i+1,j}^s - u_{i,j}^s}{\Delta x} + \mu_{i+\frac{1}{2},j}^s A_{i+\frac{1}{2},j}^s \frac{u_{i+1,j}^s - u_{i,j}^s}{\Delta x} + \mu_{i,j-\frac{1}{2},j}^s A_{i+\frac{1}{2},j}^s A_{i+\frac{1}{2},j}^s \frac{u_{i+1,j}^s - u_{i,j}^s}{\Delta x} + \mu_{i+\frac{1}{2},j}^s A_{i+\frac{1}{2},j}^s A_{i+\frac{1}{2}$$

Operations in differentiable NBM kernels are strictly local. NBM starts by placing implicit compute cells of a specified resolutions at the collocation point. At the presence of discontinuities a separate coarse mesh encapsulates an interpolant for the level-set function whose intersection with the NBM cell is calculated to obtain geometric information for the FD kernel and preconditioner (denoted M^{-1}). FD kernel is applied on the compute cell where the solution values are evaluated by the neural network. Each kernel contributes a local L^2 -norm residual $r_p = ||M^{-1}AU - M^{-1}b||$ at one point p. Preconditioning helps to balance relative magnitude of contributions from all points before aggregating residuals to form a global loss value. Finally, gradient based optimization methods used in machine learning are applied to adjust neural network parameters. The automatic differentiation loop passes across the NBM kernels.

4 Numerical results

4.1 Convergence

We consider a sphere centered at the origin with radius 1/2 in a domain $[0,1]^3$, an exact solution $u^-(x,y,z)=e^z$ inside and $u^+(x,y,z)=\cos(x)\sin(y)$ outside. In addition to the jump in solution, we also consider a jump in the variable diffusion coefficient to be $\mu^-(x,y,z)=y^2\ln(x+2)+4$ inside and $\mu^+(x,y,z)=e^{-z}$ outside the sphere. Table 1 reports convergence results for the solution, which illustrates the convergence in the L^∞ -norm.

4.2 Complex geometry and physics

We simulate a Poisson problem with discontinuities on the Dragon problem presented in [11]. In this case we used the signed-distance function produced by SDFGen, and initiated an interpolant based on its values. An exact solution $u^-(x, y, z) = \sin(2x)\cos(2y)e^z$ inside and $u^+(x, y, z) = \sin(2x)\cos(2y)e^z$

Table 1: We report L^{∞} -norm error and root-mean-squared-error (RMSE) of the solution field evaluated in the domain. Rightmost column reports the overall time to solution for our JAX implementation which constitutes 10,000 epochs in each case and the initial compilation time of jaxpressions. A pair of neural networks, each with 5 hidden layers of 10 neurons wide, have 982 total trainable parameters. In each case GPU compute occupancy is at 100% on a single NVIDIA RTX A6000 GPU.

regress ∂_n	RMSE		L^{∞}		GPU Statistics	
$N_{x,y,z}$	Solution	Order	Solution	Order	t (sec/epoch)	VRAM (GB)
2^{3}	3.7×10^{-2}	-	3.25×10^{-1}	-	0.0306	1.05
$ 2^4 $	7.1×10^{-3}	2.38	1.10×10^{-1}	1.56	0.056	1.72
2^{5}	5.9×10^{-3}	0.27	8.36×10^{-2}	0.4	0.053	2.15
2^{6}	4.1×10^{-3}	0.53	6.44×10^{-2}	0.38	0.287	5.57
27	2.64×10^{-3}	0.64	3.53×10^{-2}	0.87	2.125	32.1

$$\left[16\left(\frac{y-x}{3}\right)^5 - 20\left(\frac{y-x}{3}\right)^3 + 5\left(\frac{y-x}{3}\right)\right] \ln(x+y+3)\cos(z) \text{ outside with variable diffusion coefficients} \\ \mu^-(x,y,z) = 10\left[1 + 0.2\cos(2\pi(x+y))\sin(2\pi(x-y))\cos(z)\right] \text{ inside and } \mu^+(x,y,z) = 1 \text{ outside.}$$
 The results are shown in figure 1, with L°-norm of 0.5 and RMSE of 0.06 after 1000 epochs on multi-resolutions $64^3, 128^3, 256^3, 512^3.$

4.3 Parallel scaling on network of GPUs

Put the weak and strong scaling.

5 Discussion and conclusions

We developed a differentiable GPU-based framework for solving partial differential equations with jump conditions across irregular interfaces in three spatial dimensions. Solutions in each domain are represented by a simple multi-layer perceptron (MLP) and Cartesian grid points of the underlying numerical discretization scheme are treated as collocation points for optimizing the unknown parameters of the MLPs.

More here.

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```
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```

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```
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```

produces

```
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```

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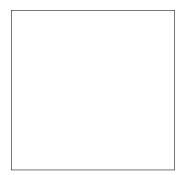


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³Sample of the first footnote.

⁴As in this example.

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```
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```

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```
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```

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- [1] Alexander, J.A. & Mozer, M.C. (1995) Template-based algorithms for connectionist rule extraction. In G. Tesauro, D.S. Touretzky and T.K. Leen (eds.), *Advances in Neural Information Processing Systems 7*, pp. 609–616. Cambridge, MA: MIT Press.
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- (c) Did you report error bars (e.g., with respect to the random seed after running experiments multiple times)? [TODO]
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A Appendix

Optionally include extra information (complete proofs, additional experiments and plots) in the appendix. This section will often be part of the supplemental material.

References

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