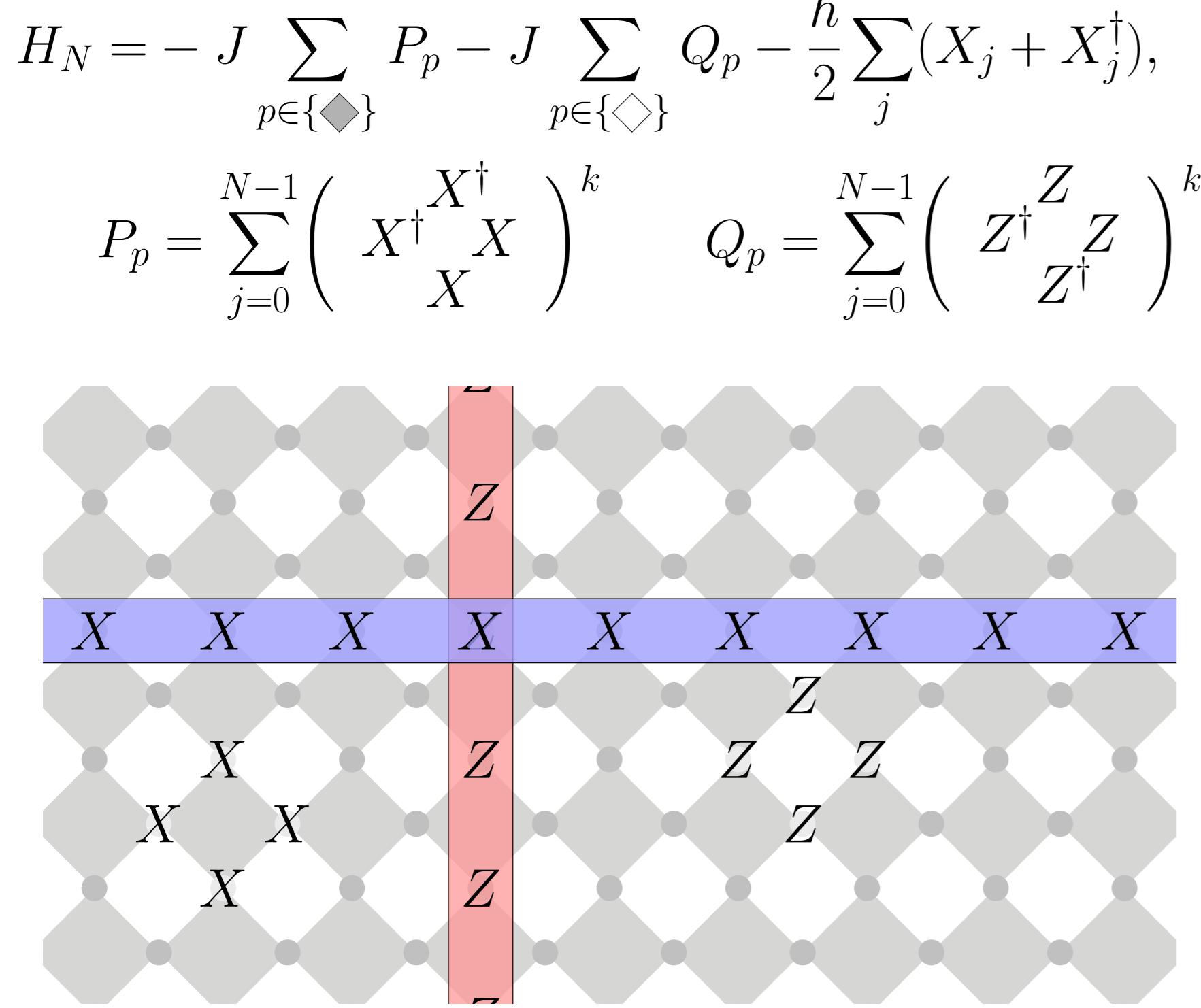


Introduction

- We introduce a numerical method for identifying topological order in two-dimensional spin models.
 - Topological order away from integrable models.
- Topologically ordered systems interesting for both condensed matter and quantum computing applications.
- Topological entanglement entropy (sometimes) identifies topological order but does not identify which order (which anyon model)[1].
 - Topological entanglement entropy sometimes mistakenly identifies topological order[2]
- By studying 1D operators associated to fundamental excitations (anyons), can efficiently search for topological order.

2D Topologically Ordered Systems

- Long range ordered but no local order parameter.
- Ground space can be used to encode quantum information which is stable under local noise.
- Ground state degeneracy depends on topology of lattice.
- Low lying excitations (anyons) created by string-like operators.
- Phases are robust to local perturbations. Anyon braiding, fusion etc. should be preserved[3].
- Ribbon operators provide a signal of a topological phase, telling us more than just the total quantum dimension.



The prototypical ribbon operators are the familiar string operators in the \mathbb{Z}_N quantum double models[5]. These commute with the Hamiltonian but have a twisted commutator with η an N th root of unity.

Properties of Anyons

- Created by string-like operator.
- Localised excitations
 - Undetectable away from particle location.
- Free to move around without changing energy.
 - String operators are deformable.
- Nontrivial braid statistics.

$$\mathcal{S}_{ab} = \frac{1}{\mathcal{D}} \bar{a} \circlearrowleft b$$

- String operators are 'slightly entangling'.
 - Anyons have few degrees of freedom, strings act to entangle them.

From Anyons to Ribbon Operators

- Operators supported on long strips with bounded width.
 - Proven in LCPC codes[4]
- Away from end points, the ribbon R should commute with the Hamiltonian.
- Existence of a ribbon should only weakly depend on the specific support chosen.
- Commutation of ribbons L and R should capture the topological data (from the \mathcal{R} or \mathcal{S} matrix)
 - We define the twisted commutator $[R, L]_\eta := RL - \eta LR$
- Restrict to the class of matrix product operators.

$$C(R; \eta) := \| [R, H] \|^2 + \| [R, L]_\eta \|^2$$

Efficiently Optimising Ribbons

- Translationally invariant matrix product operators are an efficiently representable and manipulable subclass of operators.
- By vectorising the MPO tensors, we can apply standard DMRG to optimise the ribbons.
- Cost function can be encoded as an MPO.

$$\| [R, O] \|^2 = \begin{array}{c} R \\ O \\ O^\dagger \\ R^\dagger \end{array} - \begin{array}{c} R \\ O \\ O^\dagger \\ R^\dagger \end{array} - \begin{array}{c} R \\ O \\ O^\dagger \\ R^\dagger \end{array} + \begin{array}{c} R \\ O \\ O^\dagger \\ R^\dagger \end{array} .$$

Analytic Results?

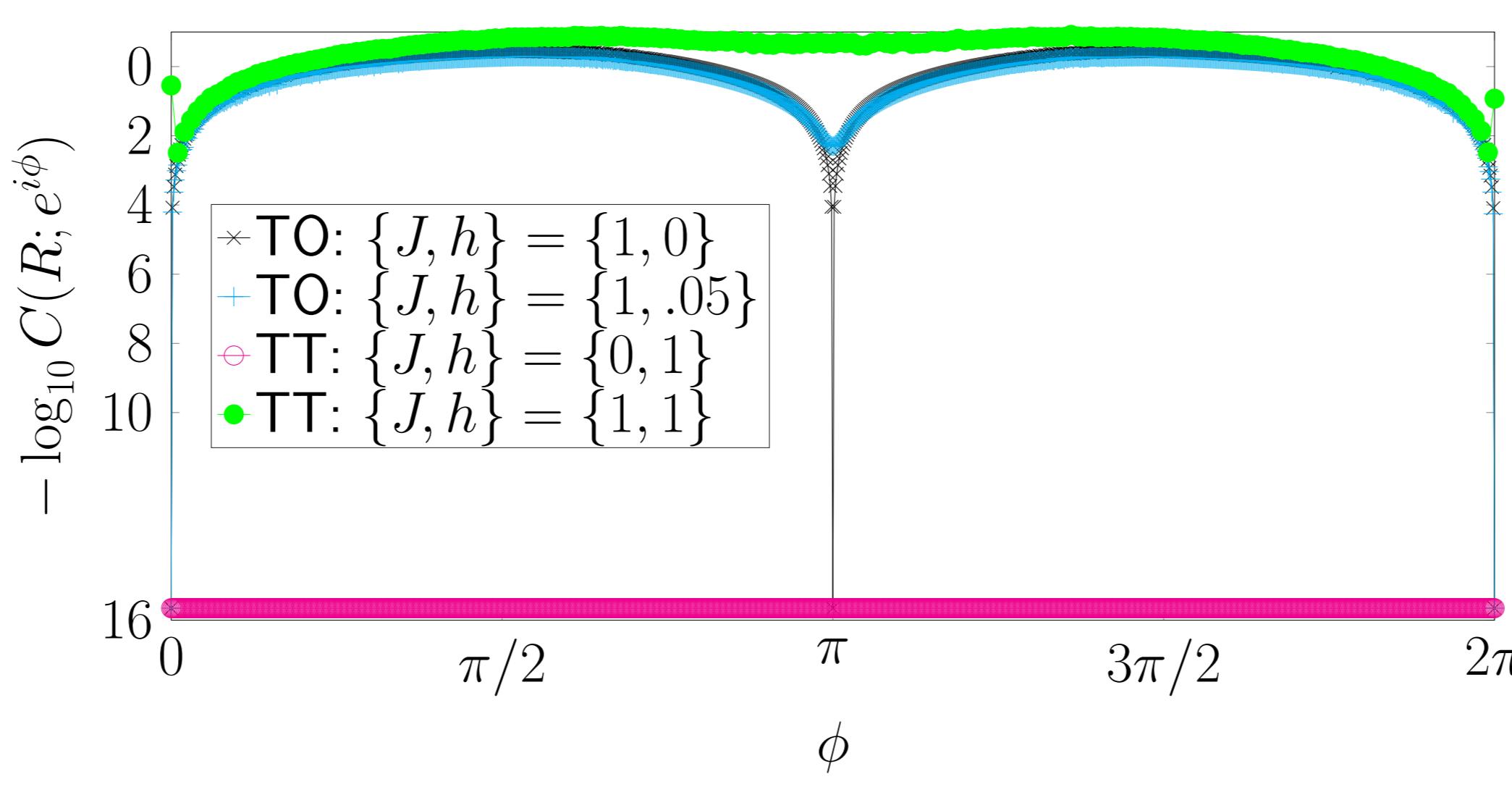
- Method presented here is heuristic.
- Analytic link between approximate symmetries (i.e. operators with small $\| [U, H] \|^2$) and ground state degeneracy discussed in poster of Chubb and Flammia.
- Can an analytic extension of this work be used to prove topological order without recourse to ground states?

References

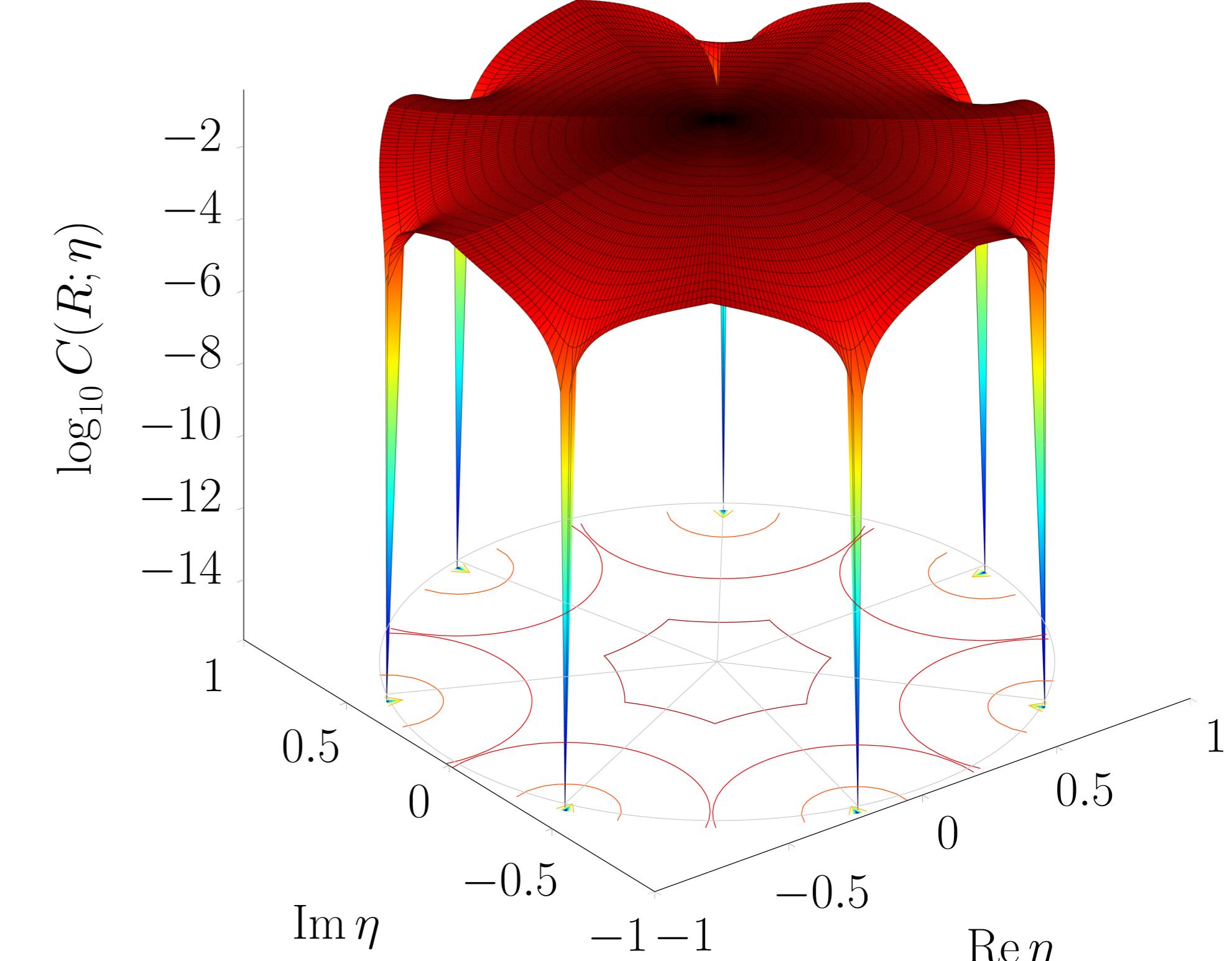
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|---------------------------------------|----------------------------------|---|
| [1] A. Kitaev et al., hep-th/0510092. | [3] S. Bravyi et al., 1001.0344. | [6] S. Trebst et al., cond-mat/0609048. |
| [2] L. Zou et al., 1604.06101. | [4] J. Haah et al., 1011.3529. | [7] A. Kitaev, cond-mat/0506438. |
| [5] A. Kitaev et al., 0904.2771. | | |

Quantum Double Models

- Exactly solvable when $h = 0$
 - Good for benchmarking numerical methods.
- Topological phase.
 - Robust to local perturbations [6]



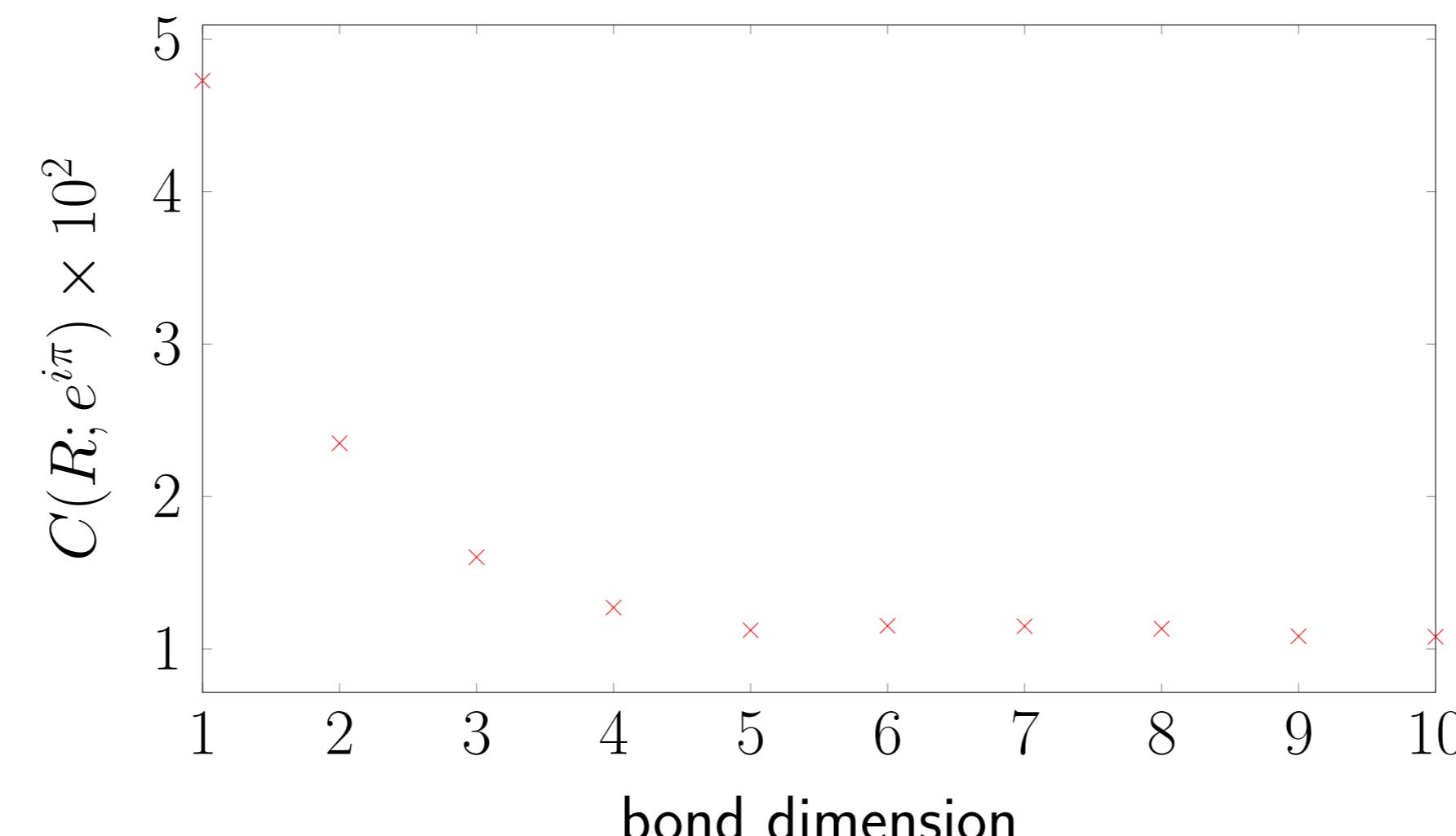
Cost for pair of ribbons in the \mathbb{Z}_2 quantum double model with a Z field. A signal of topological order (TO) is seen in the appropriate phase, no such signal is present in the paramagnetic phase. Similar behaviour is observed in the presence of an Ising-like interaction. (Here $\eta = \exp(i\phi)$)



Cost of a numerically obtained ribbon operator in the \mathbb{Z}_7 topological phase. The presence of \mathbb{Z}_7 anyons is indicated by the strong signal whenever η is a 7th root of unity.

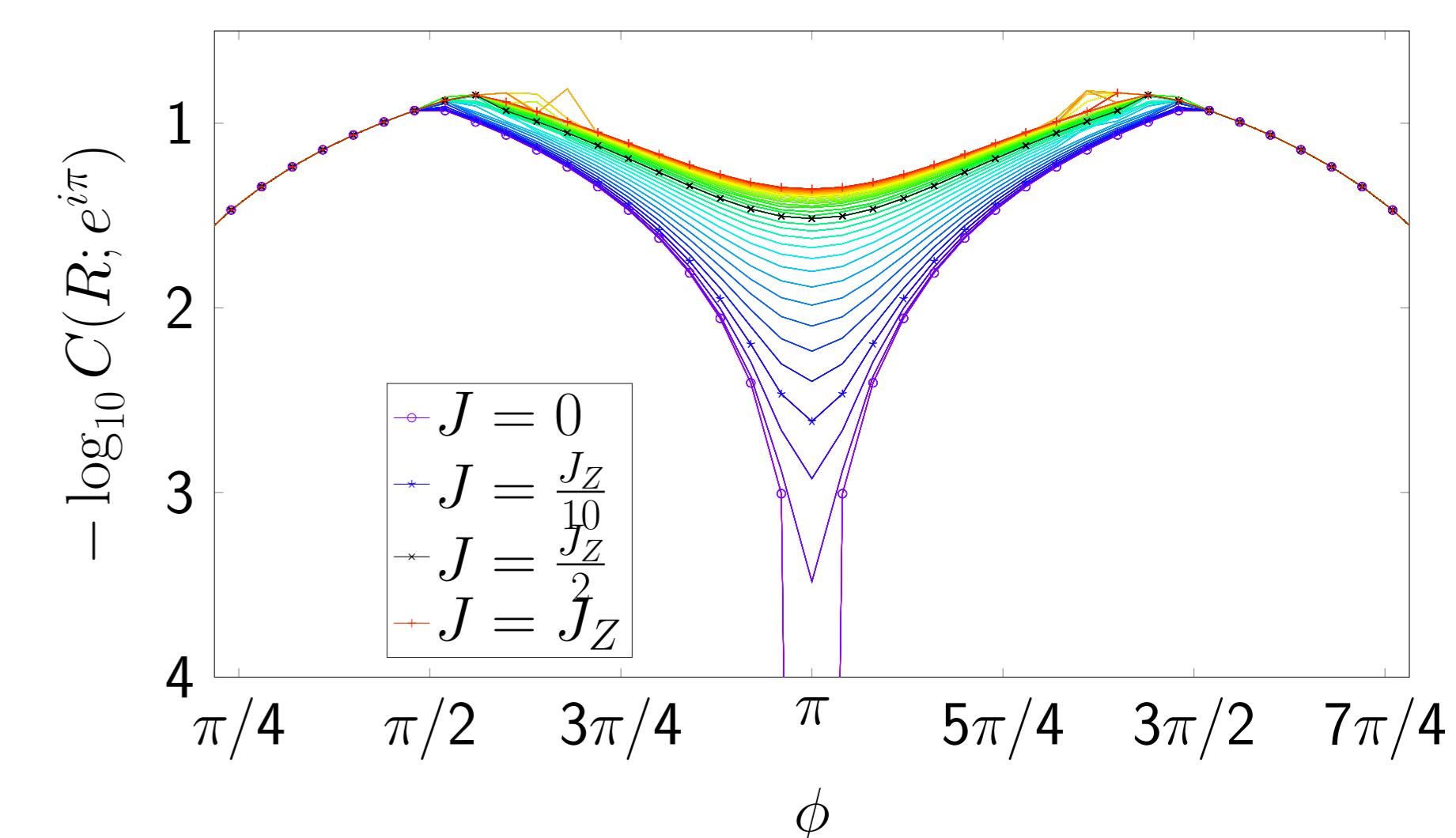
Kitaev's Honeycomb

$$H = -J_X \sum_{i,j \in X \text{ links}} X_i X_j - J_Y \sum_{i,j \in Y \text{ links}} Y_i Y_j - J_Z \sum_{i,j \in Z \text{ links}} Z_i Z_j$$

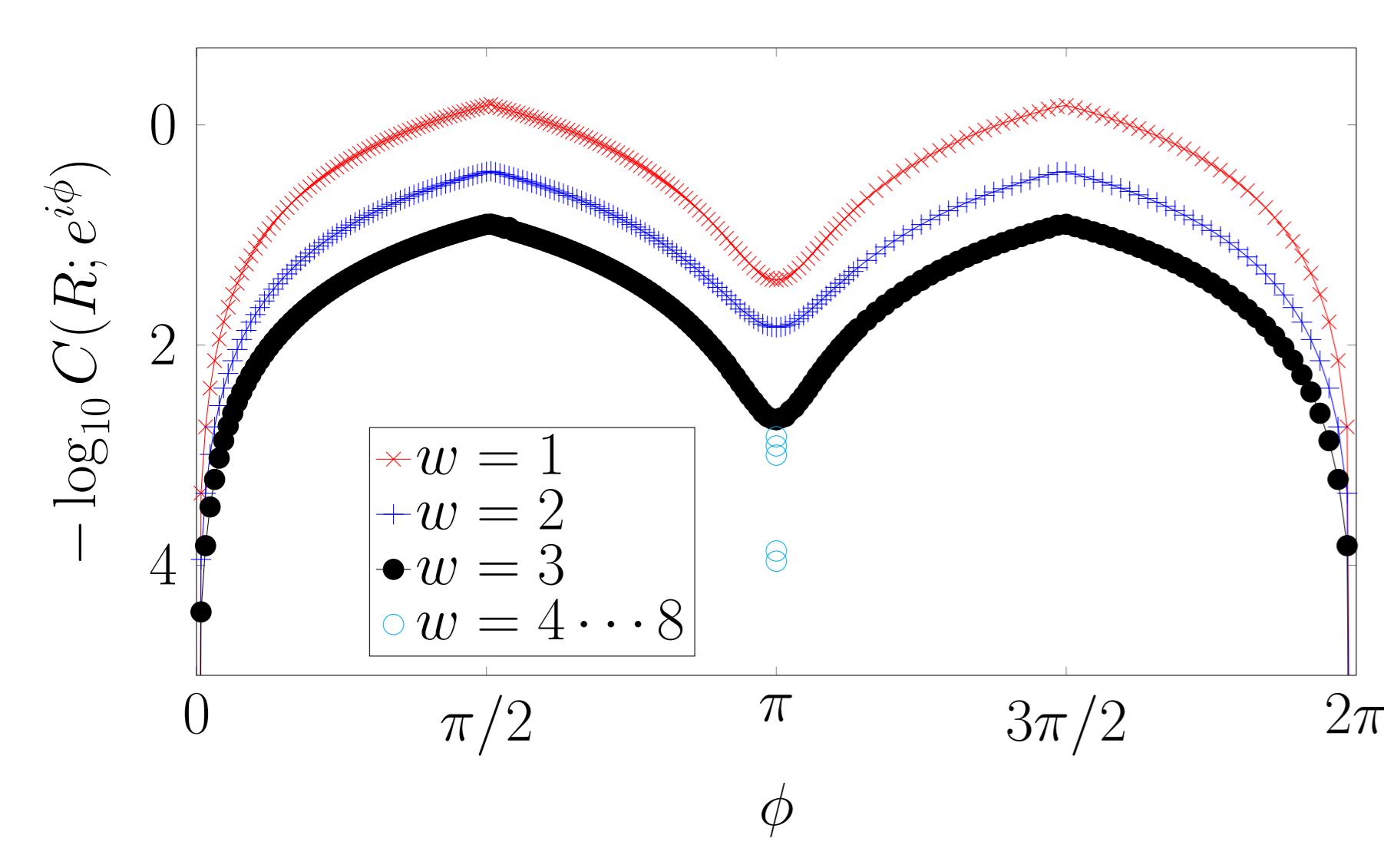


When increasing the bond dimension, we see the cost quickly saturate. This indicates that we are correct in our assumption that the ribbon operators will be 'slightly entangling'

- Perturbatively (around small J_Z) equivalent to the \mathbb{Z}_2 quantum double model[7]



When the strength of the Z term is increased, the model moves towards a phase transition. We observe the strength of the signal decreases, although does not disappear.



As the width of the support is increased, we see the cost decrease.