MA5233 Computational Mathematics

Lecture 19: Theory of PDEs

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Differential equation

Equation in terms of an unknown function $u: \Omega \to \mathbb{R}^n$ and its derivatives which is to hold at every point in a connected domain $\Omega \subset \mathbb{R}^m$.

Solution is typically only unique if we also impose values of u and its derivatives on $\partial\Omega$.

▶ Differential equation is called *ordinary* or *initial value problem* if $\Omega \subset \mathbb{R}$ and we impose values of u and its derivatives at a single point.

Example:
$$\dot{y}(t) = f(y(t))$$
 for all $t \in [0, T]$, $y(0) = y_0$.

▶ Differential equation is called *partial* or *boundary value problem* in all other cases.

Example:
$$-\Delta u(x) = f(x)$$
 for all $x \in \Omega$, $u(x) = 0$ for all $x \in \partial \Omega$.

Focus for the next few lectures: partial differential equations. Focus for today: developing a mathematically sound theory of PDEs. Motivating examples:

https://youtu.be/ureGelZPi3o, https://youtu.be/00kyDKu8K-k

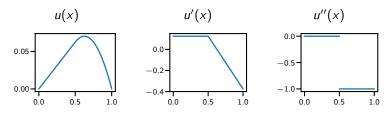
Introductory example

Consider the Poisson equation with Dirichlet boundary conditions,

$$-u'' = f$$
 on $[0,1],$ $u(0) = u(1) = 0.$

A pair f, u which "solves" this equation is given by

$$f(x) = \begin{cases} 0 & x < 0.5, \\ 1 & x > 0.5, \end{cases} \qquad u(x) = \begin{cases} \frac{x}{8} & x < 0.5, \\ \frac{(x-1/4)(1-x)}{2} & x > 0.5, \end{cases}$$



Observation: u''(x) does not exist at $x = \frac{1}{2}$! In what sense is $-u''(\frac{1}{2}) = f(\frac{1}{2})$ satisfied?

Introductory example (continued)

In example, we have -u''(x) = f(x) for all $x \in [0,1]$ except possibly at a single point $x = \frac{1}{2}$ depending on how we define u''(x).

Since $\int f(x) dx = \int g(x) dx$ if f(x) = g(x) except at a single point x, we have for any continuous function v(x) and independently of how we define $u''(\frac{1}{2})$ that

$$-\int_0^1 u''(x) v(x) dx = \int_0^1 f(x) v(x) dx.$$
 (1)

Conversely, if u satisfies (1) for all continuous functions v(x) and f(x), u''(x) are continuous, then we can set v(x) = u''(x) + f(x) and obtain

$$0 = \int_0^1 (f(x) + u''(x)) v(x) dx = \int_0^1 (f(x) + u''(x))^2 dx$$

which shows that -u''(x) = f(x) for all $x \in [0, 1]$.

Thus, equation (1) is in some sense equivalent to -u''(x) = f(x). More precise statements on next slide.

Introductory example (continued)

Repeated from previous slide for convenience:

$$-\int_0^1 u''(x) v(x) dx = \int_0^1 f(x) v(x) dx.$$
 (1)

Summary:

- ▶ If -u''(x) = f(x) holds for all $x \in [0, 1]$, then (1) holds for all continuous v(x).
- ▶ If (1) holds for all continuous v(x) and u''(x), f(x) are continuous, then -u''(x) = f(x) holds for all $x \in [0, 1]$.
- For the concrete pair f, u given on slide 3, (1) holds for all continuous v(x) while -u''(x) = f(x) does not make sense.

Conclusion: mathematical issues regarding how to interpret -u''(x) = f(x) for discontinuous f(x) disappear if we define that u(x) solves -u''(x) = f(x) if and only if (1) is satisfied for all continuous v(x).

Topics to discuss next

- Introduce a generalised notion of derivatives for functions which are not differentiable in the classical sense.
- ► Continue "reinterpreting" -u''(x) = f(x) until we arrive at a form for which we can prove existence and uniqueness of solutions.

Def: C^k and L^2 function spaces

Consider $f:[a,b]\to\mathbb{R}$. We define:

- $ightharpoonup f \in C^k([a,b])$ if f has k continuous derivatives.
- ▶ $f \in C_0^k([a, b])$ if $f \in C^k([a, b])$ and f(a) = f(b) = 0.
- ▶ $f \in L^2([a, b])$ if $\int_a^b f(x)^2 dx$ is well-defined and finite.

Def: L^2 inner product and norm

$$\langle f, g \rangle_{L^2([a,b])} := \int_a^b f(x) g(x) dx,$$

 $\|f\|_{L^2([a,b])} := \sqrt{\langle f, f \rangle_{L^2([a,b])}} = \sqrt{\int_a^b f(x)^2 dx}.$

By Cauchy-Schwarz inequality, we have

$$\langle f, g \rangle_{L^2([a,b])} \le ||f||_{L^2([a,b])} ||g||_{L^2([a,b])};$$

hence $\langle f, g \rangle_{L^2([a,b])}$ is bounded for all $f, g \in L^2([a,b])$.

Remark on L^2 norm

Strictly speaking, $||f||_{L^2([a,b])}$ is not a norm since for e.g.

$$f(x) = \begin{cases} 1 & \text{if } x = 0.5, \\ 0 & \text{otherwise} \end{cases}$$

we get $||f||_{L^2([0,1])} = 0$ but $f \neq 0$.

We fix this by interpreting a function $f \in L^2([a,b])$ as a *representative* of the set of all functions $g:[a,b] \to \mathbb{R}$ such that $\|f-g\|_{L^2([a,b])}=0$.

Simply put, we define that f = g for $f, g \in L^2([a, b])$ if $||f - g||_{L^2([a, b])} = 0$.

Important consequence: the value f(x) of $f \in L^2([a, b])$ at a single point x is not well defined.

Def: Weak derivative

Let $f, g \in L^2([a, b])$. We say g is a weak derivative of f if for all $v \in C_0^1([a, b])$ we have

$$\int_{a}^{b} f(x) v'(x) dx = - \int_{a}^{b} g(x) v(x) dx.$$

Such weak derivatives g may not exist for a given $f \in L^2([a,b])$, but if they do they are unique and we write f' := g.

If $f \in C^1$, then the weak derivatives exists and it agrees with the classical derivative.

Rationale. We obtain using integration by parts and $v \in C_0^1([a,b]) \implies v(a) = v(b) = 0$ that

$$\int_{a}^{b} \underbrace{f(x)}_{\downarrow} \underbrace{v'(x)}_{\uparrow} dx = f(b) \underbrace{v(b)}_{0} - f(a) \underbrace{v(a)}_{0} - \int_{a}^{b} f'(x) v(x) dx$$
$$= -\int_{a}^{b} f'(x) v(x) dx.$$

Thm: Weak derivatives of piecewise C^1 functions

Assume $f:[a,b]\to\mathbb{R}$ is continuous on [a,b] and continuously differentiable except on a countable set $X\subset [a,b]$. Then, the weak derivative f'(x) exists and it agrees with the classical derivative on $[a,b]\setminus X$.

Proof. Assume $X = \{\hat{x}\}$ for simplicity (result for more points in X can be shown analogously). Then, we have for every $v \in C_0^1([a,b])$ that

$$\int_{a}^{b} f(x) v'(x) dx = \int_{a}^{\hat{x}} f(x) v'(x) dx + \int_{\hat{x}}^{b} f(x) v'(x) dx$$

$$= f(\hat{x}) v(\hat{x}) - f(a) v(a) - \int_{a}^{\hat{x}} f'(x) v(x) dx + \dots$$

$$f(b) v(b) - f(\hat{x}) v(\hat{x}) - \int_{\hat{x}}^{b} f'(x) v(x) dx$$

$$= -\int_{a}^{b} f'(x) v(x) dx.$$

Introductory example (continued)

With the notation developed on the last few slides, we may formulate our reinterpretation of -u''(x) = f(x) from slide 5 in a precise way as follows.

We say $u:[0,1]\to\mathbb{R}$ "solves" -u''(x)=f(x) if u has two weak derivatives and for all $v\in C^0([0,1])$ we have that

$$-\int_0^1 u''(x) v(x) dx = \int_0^1 f(x) v(x) dx.$$
 (1)

However, it turns out that this is still not the most mathematically convenient formulation of the problem. In order to get to the final formulation, we need some more notation.

Def: Sobolev spaces

Consider $f \in L^2([a, b])$. We define:

- ▶ $f \in H^k([a, b])$ if f has k weak derivatives.
- ▶ $f \in H_0^k([a,b])$ if $f \in H^k([a,b])$ and $f^{(\ell)}(a) = f^{(\ell)}(b) = 0$ for all $\ell \in \{0, ..., k-1\}$.

Remark: boundary conditions in Sobolev spaces

According to definition, we have $f \in H^1([a, b])$ if $f \in L^2(\Omega)$, f has one weak derivative and f(a) = f(b) = 0.

It is not obvious that the last requirement makes sense: since f is in $L^2([a,b])$, it is only defined up to modification with $\delta f \in L^2([a,b])$ such that $\|\delta f\|_{L^2([a,b])} = 0$, see "Remark on L^2 norm".

This issue is resolved as follows:

It can be shown that if $f \in H^1([a,b])$, then there exists exactly one δf of the above form such that $f + \delta f \in C^0([a,b])$.

The intended meaning of f(a) = f(b) = 0 is that this condition is satisfied for the representative of $f \in L^2([a,b])$ which is also in $C^0([a,b])$.

Def: Sobolev inner product and spaces

Given $f, g \in H^k([a, b])$, we define

$$\langle f, g \rangle_{H^k([a,b])} := \sum_{\ell=0}^k \langle f^{(\ell)}, g^{(\ell)} \rangle_{L^2([a,b])},$$

$$||f||_{H^k([a,b])} := \sqrt{\langle f, f \rangle_{H^k([a,b])}}.$$

Since the weak derivatives $f^{(\ell)}, g^{(\ell)}$ are in $L^2([a, b])$, it follows that $\langle f, g \rangle_{H^k([a,b])}$ and $||f||_{H^k([a,b])}$ are bounded for $f, g \in H^k([a,b])$.

Example

$$\langle f,g\rangle_{H^1([a,b])}=\int_a^b f(x)\,g(x)\,dx+\int_a^b f'(x)\,g'(x)\,dx.$$

Hilbert space

A vector space V equipped with an inner product $\langle f,g\rangle_V$ is called a *Hilbert space* if V is complete under the norm $\|f\|_V:=\sqrt{\langle f,f\rangle_V}$.

Examples

- $ightharpoonup \mathbb{R}^n$ with inner product $\langle a,b\rangle:=a^Tb$.
- $ightharpoonup L^2([a,b])$ with L^2 inner product.
- \vdash $H^1([a,b])$ and $H^1_0([a,b])$ with H^1 inner product.

Def: Weak solution to Poisson's equation

We say $u:[0,1] \to \mathbb{R}$ is a weak solution to the Poisson equation

$$-u'' = f \text{ on } [0,1], \qquad u(0) = u(1) = 0,$$
 (2)

if $u \in H^1_0([0,1])$ and for all $v \in H^1_0([0,1])$ it holds

$$\int_0^1 u'(x) \, v'(x) \, dx = \int_0^1 f(x) \, v(x) \, dx.$$

Remarks

- ▶ If $u \in H_0^1([0,1]) \cap C^2([0,1])$ is a weak solution, then u satisfies (2).
- ▶ If $u \in C^2([0,1])$ satisfies (2), then u is a weak solution.
- ▶ If $u \in C^1([a, b])$, $u \in C^2([0, 1] \setminus X)$ for some countable set $X \subset [0, 1]$ and u satisfies (2) except on X, then u is a weak solution.

It follows from the third point that the solution from the introductory example is a weak solution.

Abstract weak formulation of Poisson's equation

The key point of reinterpreting -u'' = f according to the definition on previous slide is that the problem can now be formulated as follows:

Given Hilbert space V, bilinear $a: V \times V \to \mathbb{R}$ and linear $b: V \to \mathbb{R}$, find $u \in V$ such that

$$a(u, v) = b(v)$$
 for all $v \in V$.

For Poisson's equation, we have $V = H_0^1([0,1])$,

$$a(u,v) := \int_0^1 u'(x) v'(x) dx, \qquad b(v) := \int_0^1 f(x) v(x) dx.$$

Terminology

- ▶ Bilinear $a: V \times V \rightarrow \mathbb{R}$ is called *bilinear form*.
- ▶ Linear $b: V \to \mathbb{R}$ is called *functional*.

Lax-Milgram theorem

The abstract problem from previous slide has a unique solution and it holds $||u||_V \leq \frac{B}{A}$ if all of the following conditions are satisfied:

- ▶ a(u, v) is bounded: $\exists A > 0$ such that $|a(u, v)| \leq A ||u||_V ||v||_V$.
- ▶ a(u, v) is coercive: $\exists c > 0$ such that $a(v, v) \ge c ||v||_V^2$.
- ▶ b(v) is bounded: $\exists B > 0$ such that $|b(v)| \leq B ||v||_V$.

If a(u, v) = a(v, u) is symmetric, then LM is also called Riesz representation theorem.

Application of Lax-Milgram to Poisson equation

Boundedness of a(u, v) and b(v) is straightforward:

$$|a(u,v)| = |\langle u',v'\rangle_{L^{2}([0,1])}| \qquad |b(v)| = |\langle f,v\rangle_{L^{2}([0,1])}|$$

$$\leq ||u'||_{L^{2}([0,1])} ||v'||_{L^{2}([0,1])} \qquad \leq ||f||_{L^{2}([0,1])} ||v||_{L^{2}([0,1])}$$

$$\leq ||u||_{H^{1}([0,1])} ||v||_{H^{1}([0,1])} \qquad \leq ||f||_{L^{2}([0,1])} ||v||_{H^{1}([0,1])}$$

Hence A = 1 and $B = ||f||_{L^2([0,1])}$.

Showing coercivity of a(u, v) requires Poincaré's inequality, see next slide.

Coercivity of Poisson's equation

For $V=H^1_0([0,1])$, showing coercivity amounts to finding c>0 such that

$$0a(v,v) \ge c \left(\|v\|_{L^2([0,1])}^2 + \|v'\|_{L^2([0,1])}^2 \right). \tag{3}$$

By definition, we have that $a(v,v) = \langle v',v' \rangle_{L^2([0,1])} = ||v'||_{L^2([0,1])}^2$; hence if we can find C > 0 such that

$$||v||_{L^2([0,1])} \le C ||v'||_{L^2([0,1])},$$
 (4)

we get

$$\|v\|_{L^2([0,1])}^2 + \|v'\|_{L^2([0,1])}^2 \le (C^2 + 1) \|v'\|_{L^2([0,1])}^2 = (C^2 + 1) a(v, v)$$

which is (3) with $c = (C^2 + 1)^{-1}$.

The bound (4) indeed holds, and it is known as Poincaré inequality.

Poincaré inequality

We have $||v||_{L^2([0,1])} \le C ||v'||_{L^2([0,1])}$ for all $v \in H^1_0([0,1])$ and some C > 0 independent of v.

Discussion

- ▶ Loosely speaking, Poincaré's inequality says v' small $\implies v$ small.
- ▶ This is not true in general: v(x) := C with $C \in \mathbb{R}$ has derivative v'(x) = 0 but v(x) can be arbitrarily large.
- Poincaré's inequality can hold because we restrict $v \in H_0^1([0,1])$, which limits the above counterexample to C = 0.

What you should remember from this lecture

- ► Idea of weak derivatives.
- ► Function spaces C^k , C_0^k , L^2 , H^1 , H_0^1 , their inner products and norms.
- ▶ Weak formulation of Poisson's equation.
- ► Lax-Milgram theorem.
- Poincaré inequality.

Final remarks

This module (MA5233) is about solving PDEs numerically, and this lecture is intended to provide a minimal background for the discussion of such numerical methods.

If you want to know more about the theory of PDEs, take a module which is fully dedicated to this topic, e.g. MA4211 or MA5206.