

# MA5233 Computational Mathematics

## Lecture 19: Theory of PDEs

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# Theory of PDEs

## Differential equation

Equation in terms of an unknown function  $u : \Omega \rightarrow \mathbb{R}^n$  and its derivatives which is to hold at every point in a connected domain  $\Omega \subset \mathbb{R}^m$ .

Solution is typically only unique if we also impose values of  $u$  and its derivatives on  $\partial\Omega$ .

- ▶ Differential equation is called *ordinary* or *initial value problem* if  $\Omega \subset \mathbb{R}$  and we impose values of  $u$  and its derivatives at a single point.

Example:  $\dot{y}(t) = f(y(t))$  for all  $t \in [0, T]$ ,  $y(0) = y_0$ .

- ▶ Differential equation is called *partial* or *boundary value problem* in all other cases.

Example:  $-\Delta u(x) = f(x)$  for all  $x \in \Omega$ ,  
 $u(x) = 0$  for all  $x \in \partial\Omega$ .

Focus for the next few lectures: partial differential equations.

Focus for today: developing a mathematically sound theory of PDEs.

Motivating examples:

<https://youtu.be/ureGelZPi3o>, <https://youtu.be/00kyDKu8K-k>

# Theory of PDEs

## Introductory example

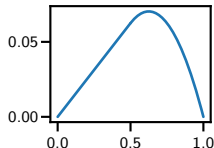
Consider the Poisson equation with Dirichlet boundary conditions,

$$-u'' = f \text{ on } [0, 1], \quad u(0) = u(1) = 0.$$

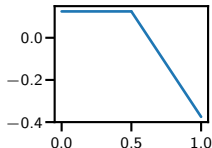
A pair  $f, u$  which “solves” this equation is given by

$$f(x) = \begin{cases} 0 & x < 0.5, \\ 1 & x > 0.5, \end{cases} \quad u(x) = \begin{cases} \frac{x}{8} & x < 0.5, \\ \frac{(x-1/4)(1-x)}{2} & x > 0.5, \end{cases}$$

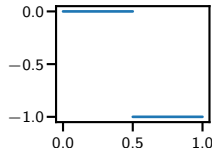
$u(x)$



$u'(x)$



$u''(x)$



Observation:  $u''(x)$  solution is discontinuous at  $x = \frac{1}{2}$ !

In what sense is  $-u''(\frac{1}{2}) = f(\frac{1}{2})$  satisfied?

# Theory of PDEs

## Introductory example (continued)

In example, we have  $-u''(x) = f(x)$  for all  $x \in [0, 1]$  except at a single point  $x = \frac{1}{2}$ .

In particular, for any continuous function  $\phi(x)$  we have

$$-\int_0^1 u''(x) \phi(x) dx = \int_0^1 f(x) \phi(x) dx. \quad (1)$$

Conversely, if (1) is satisfied for all continuous functions  $\phi(x)$  and  $f(x)$ ,  $u''(x)$  are continuous, then we can set  $\phi(x) = u''(x) + f(x)$  and obtain

$$0 = \int_0^1 \left( f(x) + u''(x) \right) \phi(x) dx = \int_0^1 \left( f(x) + u''(x) \right)^2 dx$$

which shows that  $-u''(x) = f(x)$  for all  $x \in [0, 1]$ .

# Theory of PDEs

## Introductory example (continued)

Repeated from previous slide for convenience:

$$-\int_0^1 u''(x) \phi(x) dx = \int_0^1 f(x) \phi(x) dx. \quad (1)$$

Summary:

- ▶ If  $-u''(x) = f(x)$  holds for all  $x \in [0, 1]$ , then (1) holds for all continuous  $\phi(x)$ .
- ▶ If (1) holds for all continuous  $\phi(x)$  and  $u''(x), f(x)$  are continuous, then  $-u''(x) = f(x)$  holds for all  $x \in [0, 1]$ .
- ▶ For the concrete pair  $f, u$  given above, (1) holds for all continuous  $\phi(x)$  while  $-u''(x) = f(x)$  does not make sense.

Conclusion: we may *define* that  $u(x)$  solves  $-u''(x) = f(x)$  if (1) holds for all continuous  $\phi(x)$ .

Unlike  $-\Delta u = f$ , which we derived from a flux law and conservation of mass (see Lecture 5), the above reinterpretation of this PDE has no physical meaning.

It is just a mathematical trick to arrive at a sound theory.

Also, note that this is not the final definition of what we mean by  $-u''(x) = f(x)$ .

# Theory of PDEs

## Def: Function spaces

Consider function  $f : [a, b] \rightarrow \mathbb{R}$ . We define:

- ▶  $f \in C^k([a, b])$  if  $f$  has  $k$  continuous derivatives.
- ▶  $f \in C_0^k([a, b])$  if  $f \in C^k([a, b])$  and  $f(a) = f(b) = 0$ .
- ▶  $f \in L^2([a, b])$  if  $\int_a^b f(x)^2 dx$  is well-defined and finite.

## Def: $L^2$ inner product and norm

$$\langle f, g \rangle_{L^2([a,b])} := \int_a^b f(x) g(x) dx,$$
$$\|f\|_{L^2([a,b])} := \sqrt{\langle f, f \rangle_{L^2([a,b])}} = \sqrt{\int_a^b f(x)^2 dx}.$$

By Cauchy-Schwarz inequality, we have

$$\langle f, g \rangle_{L^2([a,b])} \leq \|f\|_{L^2([a,b])} \|g\|_{L^2([a,b])};$$

hence  $\langle f, g \rangle_{L^2([a,b])}$  is bounded for all  $f, g \in L^2([a, b])$ .

# Theory of PDEs

## Remark on $L^2$ norm

Strictly speaking,  $\|f\|_{L^2([a,b])}$  is not a norm since for e.g.

$$f(x) = \begin{cases} 1 & \text{if } x = 0.5, \\ 0 & \text{otherwise} \end{cases}$$

we get  $\|f\|_{L^2([0,1])} = 0$  but  $f \neq 0$ .

We fix this by interpreting a function  $f \in L^2([a, b])$  as a *representative* of the set of all functions  $g : [a, b] \rightarrow \mathbb{R}$  such that  $\|f - g\|_{L^2([a,b])} = 0$ .

Simply put, we define that  $f = g$  for  $f, g \in L^2([a, b])$  if  $\|f - g\|_{L^2([a,b])} = 0$ .

Important consequence: the value of function  $f \in L^2([a, b])$  at a single point  $x$  is not well defined.

# Theory of PDEs

## Def: Weak derivative

Let  $f, g \in L^2([a, b])$ . We say  $g$  is a *weak derivative* of  $f$  if for all  $\phi \in C_0^\infty([a, b])$  we have

$$\int_a^b f(x) \frac{\partial \phi}{\partial x_i}(x) dx = - \int_a^b g(x) \phi(x) dx.$$

Such weak derivatives  $g$  may not exist for a given  $f \in L^2([a, b])$ , but if they do they are unique and we write  $\frac{\partial f}{\partial x} := g$ .

If  $f \in C^1$ , then the weak derivative exists and it agrees with the classical derivative.

*Rationale.* We obtain using integration by parts and

$\phi \in C_0^\infty([a, b]) \iff \phi(a) = \phi(b) = 0$  that

$$\begin{aligned} \int_a^b \underbrace{f(x)}_{\downarrow} \underbrace{\frac{\partial \phi}{\partial x}(x)}_{\uparrow} dx &= f(b) \underbrace{\phi(b)}_0 - f(a) \underbrace{\phi(a)}_0 - \int_a^b \frac{\partial f}{\partial x}(x) \phi(x) dx \\ &= - \int_a^b \frac{\partial f}{\partial x}(x) \phi(x) dx. \end{aligned}$$



# Theory of PDEs

## Thm: Weak derivatives of piecewise $C^1$ functions

Assume  $f : [a, b] \rightarrow \mathbb{R}$  is continuous on  $[a, b]$  and continuously differentiable except on a countable set  $X \subset [a, b]$ . Then, the weak derivative  $\frac{\partial f}{\partial x}$  exists and it agrees with the classical derivative on  $[a, b] \setminus X$ .

*Proof.* Assume  $X = \{\hat{x}\}$  for simplicity (result for more points in  $X$  can be shown analogously). Then, we have for every  $\phi \in C_0^\infty([a, b])$  that

$$\begin{aligned}\int_a^b f(x) \frac{\partial \phi}{\partial x}(x) dx &= \int_a^{\hat{x}} f(x) \frac{\partial \phi}{\partial x}(x) dx + \int_{\hat{x}}^b f(x) \frac{\partial \phi}{\partial x}(x) dx \\&= f(\hat{x}) \phi(\hat{x}) - f(a) \phi(a) - \int_a^{\hat{x}} \frac{\partial f}{\partial x}(x) \phi(x) dx + \dots \\&\quad f(b) \phi(b) - f(\hat{x}) \phi(\hat{x}) - \int_{\hat{x}}^b \frac{\partial f}{\partial x}(x) \phi(x) dx \\&= - \int_a^b \frac{\partial f}{\partial x}(x) \phi(x) dx.\end{aligned}$$

# Theory of PDEs

## Introductory example (continued)

We have seen that there are pairs  $f, u$  which “solve”  $-u''(x) = f(x)$  in some sense even though  $u \notin C^2$ .

We have seen that a reasonable reformulation of the PDE is as follows:

We say  $u : [0, 1] \rightarrow \mathbb{R}$  “solves”  $-u''(x) = f(x)$  if  $u$  has two weak derivatives and for all  $\phi \in C^0([0, 1])$  we have that

$$-\int_0^1 u''(x) \phi(x) dx = \int_0^1 f(x) \phi(x) dx.$$

However, it turns out that this still is not the most convenient formulation of the problem. If we restrict  $\phi$  to  $\phi \in C_0^1([0, 1])$ , then we obtain by integration by parts that

$$\int_0^1 u'(x) \phi'(x) dx = \int_0^1 f(x) \phi(x) dx.$$

This is closer to our final definition of what we mean by  $-u''(x) = f(x)$ , but it is not quite the final definition yet. In order to get there, we need some more notation.

# Theory of PDEs

## Def: Sobolev spaces

Consider  $f \in L^2([a, b])$ . We define:

- ▶  $f \in H^k([a, b])$  if  $f$  has  $k$  weak derivatives.
- ▶  $f \in H_0^k([a, b])$  if  $f \in H^k([a, b])$  and  $\frac{\partial^\ell f}{\partial x^\ell}(a) = \frac{\partial^\ell f}{\partial x^\ell}(b) = 0$  for all  $\ell \in \{0, \dots, k-1\}$ .

## Remark: boundary conditions in Sobolev spaces

According to definition, we have  $f \in H^1([a, b])$  if  $f \in L^2(\Omega)$ ,  $f$  has one weak derivative and  $f(a) = f(b) = 0$ .

It is not obvious that the last requirement makes sense: since  $f$  is only in  $L^2([a, b])$ , it is only defined up to modification with a function  $\delta f \in L^2([a, b])$  such that  $\|\delta f\|_{L^2([a, b])} = 0$ , see “Remark on  $L^2$  norm”.

However, in one dimension it can be shown that if  $f \in H^1([a, b])$ , then there exists exactly one  $\delta f$  of the above form such that  $f + \delta f \in C^0([a, b])$ .

The intended meaning of  $f(a) = f(b) = 0$  is that this condition is satisfied for the representative of  $f \in L^2([a, b])$  which is also in  $C^0([a, b])$ .

# Theory of PDEs

## Def: Weak solution to Poisson's equation

We say  $u : [a, b] \rightarrow \mathbb{R}$  is a weak solution to the Poisson equation

$$-u'' = f \text{ on } [a, b], \quad u(a) = u(b) = 0, \quad (2)$$

if  $u \in H_0^1([a, b])$  and for all  $v \in H_0^1([a, b])$  it holds

$$\int_a^b u'(x) v'(x) dx = \int_a^b f(x) v(x) dx.$$

## Remarks

- ▶ If  $u \in C^2([a, b])$  satisfies (2), then  $u$  is a weak solution.
- ▶ If  $u \in C^1([a, b])$ ,  $u \in C^2([a, b] \setminus X)$  for some countable set  $X \subset [a, b]$  and  $u$  satisfies (2) except on  $x$ , then  $u$  is a weak solution.

It follows from the second point that the solution from the introductory example is a weak solution.

# Theory of PDEs

[to be continued]