

MA3227 Numerical Analysis II

Lecture 3: Finite Differences in Two Dimensions

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Finite Differences in Two Dimensions

Introduction

Consider Poisson's equation on $\Omega = (0, 1)^2$, i.e. the problem of finding $u : [0, 1]^2 \rightarrow \mathbb{R}$ such that for all $(x_1, x_2) \in (0, 1)^2$ we have

$$-\Delta u(x_1, x_2) = f(x_1, x_2), \quad u(0, x_2) = u(1, x_2) = u(x_1, 0) = u(x_1, 1) = 0.$$

Finite difference discretisation works just like in 1D:

- Functions: $u(x_1, x_2) \rightarrow u[i_1, i_2] = u\left(\frac{i_1}{n+1}, \frac{i_2}{n+1}\right)$
- Derivatives:

$$\frac{\partial^2 u}{\partial x_1^2}\left(\frac{i_1}{n+1}, \frac{i_2}{n+1}\right) \approx (n+1)^2 \left(u[i_1+1, i_2] - 2u[i_1, i_2] + u[i_1-1, i_2] \right),$$

$$\frac{\partial^2 u}{\partial x_2^2}\left(\frac{i_1}{n+1}, \frac{i_2}{n+1}\right) \approx (n+1)^2 \left(u[i_1, i_2+1] - 2u[i_1, i_2] + u[i_1, i_2-1] \right)$$

and thus

$$\Delta u\left(\frac{i_1}{n+1}, \frac{i_2}{n+1}\right) \approx (n+1)^2 \left(u[i_1+1, i_2] + u[i_1, i_2+1] - 4u[i_1, i_2] + u[i_1-1, i_2] + u[i_1, i_2-1] \right).$$

New issue: natural “layout” for point values $u[i_1, i_2]$ is a matrix, but for linear-algebra-purposes we would like $u[i_1, i_2]$ to be a vector.

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Vectorisation of a matrix

Given $u \in \mathbb{R}^{n \times n}$, we define $\text{vec}(u) \in \mathbb{R}^{n^2}$ through

$$\text{vec}(u)[i_1 + n(i_2 - 1)] := u[i_1, i_2].$$

Example

$\text{vec}(u)$ enumerates the entries of a 4×4 matrix u as follows:

| | | | |
|---|---|----|----|
| 1 | 5 | 9 | 13 |
| 2 | 6 | 10 | 14 |
| 3 | 7 | 11 | 15 |
| 4 | 8 | 12 | 16 |

.

Remarks

- ▶ We write $\Delta_n^{(1)}$ and $\Delta_n^{(2)}$ to indicate the dimension of Δ_n .
- ▶ Vectorisation of u leads to a more complicated sparsity pattern for the finite-difference-discretised Laplacian matrix $\Delta_n^{(2)}$. See next slide.

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Two-dimensional Laplacian matrix

$$\Delta_n^{(2)} = (n+1)^2 \begin{pmatrix} \begin{bmatrix} -4 & 1 & & \\ 1 & \ddots & \ddots & \\ & \ddots & \ddots & 1 \\ & & 1 & -4 \end{bmatrix} & \begin{bmatrix} 1 & & & \\ & \ddots & & \\ & & \ddots & \\ & & & 1 \end{bmatrix} & & \\ \begin{bmatrix} 1 & & & \\ & \ddots & & \\ & & \ddots & \\ & & & 1 \end{bmatrix} & \begin{matrix} \ddots & & \\ & \ddots & \\ & & \ddots \end{matrix} & & \\ & \begin{matrix} \ddots & & \\ & \ddots & \\ & & \ddots \end{matrix} & & \begin{bmatrix} 1 & & & \\ & \ddots & & \\ & & \ddots & \\ & & & 1 \end{bmatrix} \\ & & \begin{bmatrix} 1 & & & \\ & \ddots & & \\ & & \ddots & \\ & & & 1 \end{bmatrix} & \begin{bmatrix} -4 & 1 & & \\ 1 & \ddots & \ddots & \\ & \ddots & \ddots & 1 \\ & & 1 & -4 \end{bmatrix} \end{pmatrix}$$

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Discussion

Solving Poisson's equation in 2D in principle works just like in 1D: assemble the linear system and solve.

However, there are a few technical complications:

- ▶ $\Delta_n^{(2)}$ has n^4 entries, but only $\mathcal{O}(n^2)$ of these entries are nonzero. Exploiting this special sparsity property is crucial to stay within the memory limits of a typical laptop. Julia's `SparseArrays` package provides a data type which allows us to represent sparse matrices efficiently.
- ▶ The structure of $\Delta_n^{(2)}$ is quite complicated. Assembling this matrix is tricky unless we use the Kronecker product trick introduced on the next slide.

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Kronecker product

Let a, b be vectors and A, B be matrices. We define

$$a \otimes b := \begin{pmatrix} a[1] b \\ \vdots \\ a[n] b \end{pmatrix}, \quad A \otimes B := \begin{pmatrix} A[1, 1] B & \cdots & A[1, n] B \\ \vdots & \ddots & \vdots \\ A[n, 1] B & \cdots & A[n, n] B \end{pmatrix}.$$

Remarks

- ▶ $h = f \otimes g$ corresponds to the product $h(x_1, x_2) = g(x_1) f(x_2)$.
- ▶ Let I be the identity matrix. We have the correspondences

$$\frac{\partial^2}{\partial x_1^2} \longrightarrow I \otimes \Delta_n^{(1)}, \quad \frac{\partial^2}{\partial x_2^2} \longrightarrow \Delta_n^{(1)} \otimes I.$$

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Two-dimensional Laplacian

It follows from previous slide that the finite-difference discretisation of

$$\Delta = \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2}$$

is given by

$$\Delta_n^{(2)} = \Delta_n^{(1)} \otimes I + I \otimes \Delta_n^{(1)}$$

Implementing $\Delta_n^{(2)}$ is straightforward with this formula.

See `laplacian_2d()`.

Convergence analysis for two-dimensional finite differences

The “stability & consistency \Rightarrow convergence” trick also works in 2D.

- ▶ Consistency can be shown using Taylor expansions just like in 1D.
- ▶ Stability of $\Delta_n^{(2)}$ requires us to study the eigenvalues of $\Delta_n^{(2)}$.
These eigenvalues can be derived from the results on the next slide.

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Lemma

$$(A \otimes B)(a \otimes b) = (Aa) \otimes (Bb)$$

Proof. Straightforward but tedious computations.

Eigenvalues and -vectors of 2D Laplacian

Let λ_k, u_k be eigenpairs of $\Delta_n^{(1)}$. Then, eigenpairs of $\Delta_n^{(2)}$ are

$$\lambda_{k_1, k_2} = \lambda_{k_1} + \lambda_{k_2}, \quad u_{k_1, k_2} = u_{k_1} \otimes u_{k_2}.$$

Proof.

$$\begin{aligned} \Delta_n^{(2)} u_{k_1, k_2} &= \left(\Delta_n^{(1)} \otimes I + I \otimes \Delta_n^{(1)} \right) (u_{k_1} \otimes u_{k_2}) \\ &= (\Delta_n^{(1)} u_{k_1}) \otimes u_{k_2} + u_{k_1} \otimes (\Delta_n^{(1)} u_{k_2}) \\ &= \lambda_{k_1} u_{k_1} \otimes u_{k_2} + \lambda_{k_2} u_{k_1} \otimes u_{k_2} \\ &= (\lambda_{k_1} + \lambda_{k_2}) u_{k_1} \otimes u_{k_2} \\ &= (\lambda_{k_1} + \lambda_{k_2}) u_{k_1, k_2}. \end{aligned}$$

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Stability of 2D Laplacian

It follows from results on previous slide that

$$|\lambda_{\min}| = |\lambda_{1,1}| = 2 |\lambda_1| = 2\pi^2 + \mathcal{O}(n^{-2}).$$

Hence, we have $\|(\Delta_n^{(2)})^{-1}\|_{2,n} = \mathcal{O}(1)$, i.e. 2D Laplacian is stable.

Convergence of 2D finite differences

Combining the consistency and stability result, we obtain

$$\|u - u_n\|_{n,2} = \mathcal{O}(n^{-2})$$

as in 1D.

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Summary

- ▶ The structure of $\Delta_n^{(2)}$ is “tridiagonal + two far off-diagonals”.
- ▶ Eigenpairs of $\Delta_n^{(2)}$ are given by

$$\lambda_{k_1, k_2} = \lambda_{k_1} + \lambda_{k_2}, \quad u_{k_1, k_2} = u_{k_1} \otimes u_{k_2}$$

with λ_k, u_k the eigenpairs of $\Delta_n^{(1)}$.