

# Fun vs. Black Hole: An Accretion Model in Unified Probability Theory

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August 27, 2025

## Abstract

We present a three-state toy system in Unified Probability Theory (UPT) that models the felt dynamics of sharing creative work (*fun*) in the presence of an academic gatekeeper (a *black hole*). The model uses UPT’s conservative mechanics on the probability simplex—potential, emergence (viability), resonance, and a threshold-gated cascade—to formalize “accretion” of probability mass from playful exploration into reductionist collapse. The goal is explanatory: to show, as a small but rigorous example, how UPT encodes the transition from healthy engagement to accretion once criticality is crossed.

## 1 State, events, and intuition

Let  $\mathcal{A} = \{F, O, B\}$  denote disjoint events (fun, engagement orbit, black hole). The time-dependent distribution  $\mathbf{P}(t) = (P_F(t), P_O(t), P_B(t))$  lives on the probability simplex  $\Delta^2 = \{\mathbf{p} \in \mathbb{R}^3 : p_k \geq 0, \sum_k p_k = 1\}$ . Intuitively:  $F$  radiates playful exploration;  $O$  is the back-and-forth where the outcome can still go either way;  $B$  is the reductionist sink.

## 2 UPT operators (conservative forms)

We re-use the UPT operators in a minimal three-node system. Let  $\mathbf{z} = (z_F, z_O, z_B)$  be an anchor (baseline) distribution and  $\phi_k > 0$  be potential strengths. Define the scalar potential

$$\mathcal{V}(\mathbf{P}) = \sum_{k \in \{F, O, B\}} \frac{1}{2} \phi_k (P_k - z_k)^2, \quad g_k = \frac{\partial \mathcal{V}}{\partial P_k} = \phi_k (P_k - z_k). \quad (1)$$

To conserve total probability, use the zero-sum projection  $\Pi(\mathbf{v}) = \mathbf{v} - \frac{1}{3}(\mathbf{1}^\top \mathbf{v})\mathbf{1}$  and take the potential term  $-\Pi(\nabla \mathcal{V})$  (descent toward anchors).

**Emergence (viability).** With parameters  $\beta_k \in \mathbb{R}$ ,  $V_{\text{th}} \in \mathbb{R}$ ,  $\sigma > 0$ ,

$$\Lambda_k(\mathbf{P}) = \beta_k \tanh\left(\frac{V_k(\mathbf{P}) - V_{\text{th}}}{\sigma}\right), \quad V_k(\mathbf{P}) = \frac{1}{2} \phi_k (P_k - z_k)^2, \quad (2)$$

and recenter to zero-sum:  $\hat{\Lambda}_k = \Lambda_k - \frac{1}{3} \sum_i \Lambda_i$ .

**Resonance (pairwise conservative coupling).** Let

$$S_{ij}(\mathbf{P}) = A_{\text{res}} \tanh\left(\frac{\widehat{V}_{ij}(\mathbf{P}) - V_c}{\tau_{\text{res}}}\right), \quad \widehat{V}_{ij}(\mathbf{P}) = \frac{1}{2}(V_i(\mathbf{P}) + V_j(\mathbf{P})), \quad (3)$$

with  $S_{ij} = S_{ji}$  and nonzero only for adjacent modes:  $S_{FO} = S_{OF}$ ,  $S_{OB} = S_{BO}$ ,  $S_{FB} = 0$ . Define resonance fluxes  $r_{i \rightarrow j} = \gamma S_{ij} P_i$  and the net

$$R_k(\mathbf{P}) = \sum_{j \neq k} (r_{j \rightarrow k} - r_{k \rightarrow j}). \quad (4)$$

Then  $\sum_k R_k \equiv 0$ .

**Cascade (accretion).** Define the cascade potential

$$\Phi(\mathbf{P}) = \max \left\{ S_{FO}(\mathbf{P}) E_{FO}(\mathbf{P}), S_{OB}(\mathbf{P}) E_{OB}(\mathbf{P}) \right\}, \quad E_{ij}(\mathbf{P}) = \beta_{ij} \tanh\left(\frac{\widehat{V}_{ij}(\mathbf{P}) - V_{\text{th}}}{\sigma}\right). \quad (5)$$

When  $\Phi \geq \Phi_{\text{critical}}$ , cascade fluxes turn on as gated, antisymmetric reallocations:

$$W_{i \rightarrow j}(\mathbf{P}) = \mathbf{1}\{\Phi \geq \Phi_{\text{critical}}\} \frac{[S_{ij}(\mathbf{P}) E_{ij}(\mathbf{P})]_+}{\sum_{\ell \neq i} [S_{i\ell}(\mathbf{P}) E_{i\ell}(\mathbf{P})]_+ + \varepsilon}, \quad (6)$$

$$c_{i \rightarrow j}(\mathbf{P}) = W_{i \rightarrow j}(\mathbf{P}) P_i, \quad \mathcal{C}_k(\mathbf{P}) = \sum_{j \neq k} (c_{j \rightarrow k} - c_{k \rightarrow j}), \quad (7)$$

with  $[x]_+ = \max(x, 0)$  and small  $\varepsilon > 0$ . Again  $\sum_k \mathcal{C}_k \equiv 0$ .

### 3 Unified dynamics and schedule

The UPT flow on  $\Delta^2$  is

$$\dot{\mathbf{P}} = -\Pi(\nabla \mathcal{V}) + \widehat{\Lambda}(\mathbf{P}) + R(\mathbf{P}) + \alpha(t) \mathcal{C}(\mathbf{P}), \quad (8)$$

with cascade schedule

$$\alpha(t) = \begin{cases} 0, & \Phi(\mathbf{P}(t)) < \Phi_{\text{critical}}, \\ \alpha_0 2^{n(t)}, & \Phi(\mathbf{P}(t)) \geq \Phi_{\text{critical}}, \end{cases} \quad n(t) = \#\{s \leq t : \Phi(\mathbf{P}(s)) \text{ crossed } \Phi_{\text{critical}}\}. \quad (9)$$

### 4 Conservation, invariance, and horizon

**Theorem 1** (Simplex invariance and conservation). *For bounded parameters and  $\varepsilon > 0$ , the vector field is locally Lipschitz. For any  $\mathbf{P}(0) \in \Delta^2$ , the solution exists forward in time, remains in  $\Delta^2$ , and satisfies  $P_F + P_O + P_B \equiv 1$ .*

*Sketch.*  $\tanh$  is smooth; the rational gate is bounded by  $\varepsilon$ ;  $\Pi$  enforces zero-sum for the potential;  $\widehat{\Lambda}$ ,  $R$ ,  $\mathcal{C}$  are pairwise antisymmetric or recentered to zero-sum. Standard inward-pointing arguments at the faces of  $\Delta^2$  guarantee nonnegativity.  $\square$

**Remark 1** (Event horizon and “accretion”). If  $S_{OB}E_{OB}$  dominates while the gate is open, the cascade channel  $O \rightarrow B$  captures mass rapidly; once  $P_F$  falls below a threshold  $P_F^{\text{hor}}$ , further recovery becomes unlikely absent an exogenous change (e.g. a respectful engager increasing  $S_{FO}E_{FO}$ ). This is the *fun horizon*.

## 5 Minimal parameterization (one plausible choice)

For illustration (dimensionless units):  $\mathbf{z} = (0.7, 0.25, 0.05)$ ,  $\phi_F = \phi_O = \phi_B = 0.2$ ,  $A_{\text{res}} = 1$ ,  $V_c = 0.12$ ,  $\tau_{\text{res}} = 0.08$ ,  $\beta_F = +0.05$ ,  $\beta_O = 0$ ,  $\beta_B = -0.03$ ,  $\beta_{FO} = +0.5$ ,  $\beta_{OB} = +0.7$ ,  $V_{\text{th}} = 0.02$ ,  $\sigma = 0.04$ ,  $\Phi_{\text{critical}} = 0.015$ ,  $\gamma = 0.2$ ,  $\alpha_0 = 0.6$ ,  $\varepsilon = 10^{-6}$ .

### Plain-language summary (the “why this feels true” part)

Fun ( $F$ ) radiates; the gatekeeper ( $B$ ) pulls. At first, engagement ( $O$ ) can resonate with fun. If that resonance is respectful, probability cycles between  $F$  and  $O$  and you keep creating. If the interaction turns reductionist, the pair ( $O, B$ ) crosses a threshold and a *cascade* (accretion) switches on: mass flows along  $F \rightarrow O \rightarrow B$  rapidly. Cross the horizon and  $P_F$  collapses toward zero—not because probabilities have “energy,” but because the vector field on the simplex now points into the accretion channel.