电动力学-第十次作业

陈朝翔-2021302021121

2023年4月22日

Problem 10.16

Answer: The potential of the particle is equal to:

$$V(\vec{r},t) = \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{\left(c^2t - \vec{r} \cdot \vec{v}\right)^2 + \left(c^2 - v^2\right)\left(r^2 - c^2t^2\right)}}$$

We need to somehow "pull out" the factor of Rc from the denominator. The expression under the root is:

$$\begin{split} \left(c^2t - \vec{r} \cdot \vec{v}\right)^2 + \left(c^2 - v^2\right) \left(r^2 - c^2t^2\right) &= \\ &= c^4t^2 - 2c^2t\vec{r} \cdot \vec{v} + (\vec{r} \cdot \vec{v})^2 + c^2r^2 - c^4t^2 - v^2r^2 + v^2c^2t^2 \\ &= -2c^2t\vec{r} \cdot \vec{v} + (\vec{r} \cdot \vec{v})^2 + c^2r^2 - v^2r^2 + v^2c^2t^2 \\ \vec{R} &= \vec{r} - \vec{v}t \Longrightarrow R^2 = r^2 + v^2t^2 - 2t\vec{v} \cdot \vec{r} \\ \left(c^2t - \vec{r} \cdot \vec{v}\right)^2 + \left(c^2 - v^2\right) \left(r^2 - c^2t^2\right) &= \\ &= c^2\left(r^2 + v^2t^2 - 2t\vec{v} \cdot \vec{r}\right) + (\vec{r} \cdot \vec{v})^2 - v^2r^2 \\ &= c^2R^2 + (\vec{r} \cdot \vec{v})^2 - v^2r^2 \end{split}$$

Now, the two other two terms are equal to:

$$\begin{split} (\vec{r} \cdot \vec{v})^2 - v^2 r^2 &= ((\vec{R} + \vec{v}t) \cdot \vec{v})^2 - v^2 \left(R^2 + 2t \vec{r} \cdot \vec{v} - v^2 t^2 \right) \\ &= v^2 R^2 \cos^2 \theta + 2 (\vec{R} \cdot \vec{v}) v^2 t + v^4 t^2 - v^2 R^2 + 2 v^2 t \vec{r} \cdot \vec{v} + v^4 t^2 \\ &= v^2 R^2 \left(\cos^2 \theta - 1 \right) + 2 ((\vec{r} - \vec{v}t) \cdot \vec{v}) v^2 t + 2 v^2 t \vec{r} \cdot \vec{v} + 2 v^4 t^2 \\ &= -v^2 R^2 \sin^2 \theta + 2 v^2 t \vec{r} \cdot \vec{v} - 2 v^4 t^2 - 2 v^2 t \vec{r} \cdot \vec{v} + 2 v^4 t^2 \\ &= -v^2 R^2 \sin^2 \theta \end{split}$$

where θ is the angle the velocity makes with the vector \vec{R} . The expression under the root is then:

$$(c^{2}t - \vec{r} \cdot \vec{v})^{2} + (c^{2} - v^{2})(r^{2} - c^{2}t^{2}) = c^{2}R^{2} - v^{2}R^{2}\sin^{2}\theta$$
$$= c^{2}R^{2}(1 - \beta^{2}\sin^{2}\theta)$$

So, the potential is:

$$V = \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{(c^2t - \vec{r} \cdot \vec{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)}}$$
$$= \frac{1}{4\pi\epsilon_0} \frac{qc}{Rc\sqrt{1 - \beta^2 \sin^2 \theta}} = \sqrt{\frac{1}{4\pi\epsilon_0}} \frac{q}{R\sqrt{1 - \beta^2 \sin^2 \theta}}$$

Q.E.D

Problem 10.19

Answer:

Due to the limitations of this site I will be using letter l as the distance from the charge to the point of interest.

The Lienard-Wiechert vector potential of the charge in motion is:

$$\vec{A} = \frac{\mu_0}{4\pi} \frac{qc\vec{v}}{l_c - \vec{l} \cdot \vec{v}} = \frac{\mu_0}{4\pi} \frac{qc\vec{v}}{\vec{l} \cdot \vec{u}}$$

all quantities evaluated at retarded time. We want the partial time derivative of this, but first:

$$\begin{split} l &= c \left(t - t_r \right) \\ l^2 &= c^2 \left(t - t_r \right)^2 \\ 2 \vec{l} \cdot \frac{\partial \vec{l}}{\partial t} &= 2c^2 \left(t - t_r \right) \left(1 - \frac{\partial t_r}{\partial t} \right) \Rightarrow \frac{\partial t_r}{\partial t} = 1 - \frac{\hat{l}}{c} \cdot \frac{\partial \vec{l}}{\partial t} \end{split}$$

But:

$$\begin{split} \vec{l} &= \vec{r} - \vec{w} \left(t_r \right) \\ \frac{\partial \vec{l}}{\partial t} &= -\frac{\partial \vec{w}}{\partial t_r} \frac{\partial t_r}{\partial t} = -\vec{v} \left(t_r \right) \frac{\partial t_r}{\partial t} \\ \frac{\partial t_r}{\partial t} &= 1 + \frac{\hat{l}}{c} \cdot \vec{v} \left(t_r \right) \frac{\partial t_r}{\partial t} \end{split}$$

So, as we need to prove first:

$$\frac{\partial t_r}{\partial t} = \frac{cl}{\vec{l} \cdot (c\hat{l} - \vec{v})} = \frac{cl}{\vec{l} \cdot \vec{u}}$$

With this we can get the partial time derivative of the vector potential:

$$\begin{split} \frac{\partial \vec{A}}{\partial t} &= \frac{\mu_0 qc}{4\pi} \left[\frac{1}{\vec{l} \cdot \vec{u}} \frac{\partial \vec{v}}{\partial t} - \frac{\vec{v}}{(\vec{l} \cdot \vec{u})^2} \frac{\partial}{\partial t} (\vec{l} \cdot \vec{u}) \right] \\ \frac{\partial \vec{v}}{\partial t} &= \frac{\partial \vec{v}}{\partial t_r} \frac{\partial t_r}{\partial t} = \vec{a} \frac{cl}{\vec{l} \cdot \vec{u}} \\ \frac{\partial}{\partial t} (\vec{l} \cdot \vec{u}) &= c \frac{\partial l}{\partial t} - \frac{\partial \vec{l}}{\partial t} \cdot \vec{v} - \frac{\partial \vec{v}}{\partial t} \cdot \vec{l} = c \frac{\partial l}{\partial t} + v^2 \frac{cl}{\vec{l} \cdot \vec{u}} - (\vec{a} \cdot \vec{l}) \frac{cl}{\vec{l} \cdot \vec{u}} \end{split}$$

Hence:

$$\begin{split} \frac{\partial \vec{A}}{\partial t} &= \frac{\mu_0 q c}{4\pi} \frac{1}{(\vec{l} \cdot \vec{u})^3} \left[lc \vec{a} (\vec{l} \cdot \vec{u}) + \vec{v} \left(cl \left(c^2 - v^2 + \vec{l} \cdot \vec{a} \right) + c^2 (\vec{l} \cdot \vec{u}) \right) \right] \\ &= \frac{q c}{4\pi \epsilon_0} \left[(\vec{l} \cdot \vec{u}) \left(\frac{l}{c} \vec{a} - \vec{v} \right) + \frac{l}{c} \vec{v} \left(c^2 - v^2 + \vec{r} \cdot \vec{a} \right) \right] \end{split}$$

Q.E.D

Problem 10.22

Answer:

The electric field due to the leght of the wire of lenght dx is the field of the uniformly moving point charge:

$$d\vec{E} = \frac{\lambda dx}{4\pi\epsilon_0} \frac{1 - \beta^2}{\left(1 - \beta^2 \sin^2 \theta\right)^{3/2}} \frac{\hat{R}}{R^2}$$

Only the vertical component of the field will survive the integration, so:

$$dE = \frac{\lambda dx}{4\pi\epsilon_0} \frac{1 - \beta^2}{\left(1 - \beta^2 \sin^2 \theta\right)^{3/2}} \frac{\sin \theta}{R^2}$$

$$R = \frac{d}{\sin \theta} - x = R \cos \theta = d \cot \theta \quad dx = d \frac{d\theta}{\sin^2 \theta}$$

$$= \frac{\lambda d}{4\pi\epsilon_0} \frac{1 - \beta^2}{\left(1 - \beta^2 \sin^2 \theta\right)^{3/2}} \frac{\sin^3 \theta}{d^2} \frac{d\theta}{\sin^2 \theta}$$

$$= \frac{\lambda}{4\pi\epsilon_0 d} \left(1 - \beta^2\right) \frac{\sin \theta d\theta}{\left(1 - \beta^2 \sin^2 \theta\right)^{3/2}}$$

The total electric field is then the integral:

$$E = \frac{\lambda}{4\pi\epsilon_0 d} \left(1 - \beta^2\right) \int_0^{\pi} \frac{\sin\theta d\theta}{\left(1 - \beta^2 \sin^2\theta\right)^{3/2}}$$
$$= \frac{\lambda}{2\pi\epsilon_0 d} \left(1 - \beta^2\right) \int_0^{\pi/2} \frac{\sin\theta d\theta}{\left(1 - \beta^2 \sin^2\theta\right)^{3/2}}$$

since the integrand is symmetrical about $\theta = \pi/2$. Make the substitution:

$$u = \cos \theta \quad du = -\sin \theta d\theta$$
$$E = \frac{\lambda}{2\pi\epsilon_0 d} (1 - \beta^2) \int_0^1 \frac{du}{(1 - \beta^2 (1 - u^2))^{3/2}}$$

The integral is of the table variety:

$$\int_{0}^{1} \frac{du}{(1-\beta^{2}(1-u^{2}))^{3/2}} = \frac{1}{\beta^{3}} \int_{0}^{1} \frac{du}{(\beta^{-2}-1+u^{2}))^{3/2}}$$

$$= \frac{1}{\beta^{3}} \int_{0}^{1} \frac{du}{(A^{2}+u^{2})^{3/2}}$$

$$A^{2} = \beta^{-2} - 1$$

$$\frac{1}{\beta^{3}} \int_{0}^{1} \frac{du}{(A^{2}+u^{2}))^{3/2}} = \frac{1}{\beta^{3}} \frac{u}{A^{2}\sqrt{A^{2}+u^{2}}} \Big|_{0}^{1}$$

$$= \frac{1}{\beta^{3}} \frac{1}{A^{2}\sqrt{A^{2}+1}} = \frac{1}{\beta^{3}} \frac{1}{(\beta^{-2}-1)\beta-1} = \frac{1}{1-\beta^{2}}$$

The electric field is thus:

$$E = \frac{\lambda}{2\pi\epsilon_n d} \left(1 - \beta^2\right) \frac{1}{1 - \beta^2} = \frac{\lambda}{2\pi\epsilon_0 d}$$

$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 d} \hat{s}$$