

FYMM/MMP IIIb 2020 Solutions to Problem Set 6

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1. Calculating the Riemann tensor, the Ricci tensor and the scalar curvature for unit sphere S^2

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

$$ds^2 = R^2 d\theta^2 + R^2 \sin^2 \theta d\phi^2$$

$$g = d\theta \otimes d\theta + \sin^2 \theta d\phi \otimes d\phi$$

The metric from the previous exercise is

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 \\ 0 & \sin^2 \theta \end{pmatrix}$$

With inverse

$$g^{\mu\nu} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{1}{\sin^2 \theta} \end{pmatrix}$$

The connection coefficients are calculated from

$$\Gamma_{\alpha\beta}^\lambda = g^{\lambda\mu} \Gamma_{\mu\beta\alpha}$$

Where

$$\Gamma_{\mu\beta\alpha} = \frac{1}{2} (\partial_\alpha g_{\beta\mu} + \partial_\beta g_{\mu\alpha} - \partial_\mu g_{\alpha\beta})$$

Therefore we have

$$\Gamma_{\phi\phi\theta} = \Gamma_{\theta\phi\phi} = \frac{1}{2} (\partial_\theta g_{\phi\phi} + \partial_\phi g_{\phi\theta} - \partial_\phi g_{\phi\theta})$$

$$= \frac{1}{2} \partial_\theta g_{\phi\phi}$$

$$\sin \theta \cos \theta$$

$$\Gamma_{\theta\phi\phi} = \frac{1}{2} (\partial_\phi g_{\theta\phi} + \partial_\phi g_{\theta\phi} - \partial_\theta g_{\phi\phi})$$

$$= -\sin \theta \cos \theta$$

I have neglected to show 0 terms. And

$$\Gamma_{\phi\theta}^\phi = \Gamma_{\theta\phi}^\phi = g^{\phi\phi} \Gamma_{\phi\phi\theta}$$

$$= \frac{1}{\sin^2 \theta} \sin \theta \cos \theta$$

$$= \frac{\cos \theta}{\sin \theta}$$

$$\begin{aligned}\Gamma_{\phi\phi}^\theta &= g^{\theta\theta}\Gamma_{\theta\phi\phi} \\ &= -\sin\theta\cos\theta\end{aligned}$$

The Riemann tensor is given by

$$R_{\lambda\mu\nu}^\kappa = \partial_\mu\Gamma_{\nu\lambda}^\kappa - \partial_\nu\Gamma_{\mu\lambda}^\kappa + \Gamma_{\mu\eta}^\eta\Gamma_{\nu\eta}^\kappa - \Gamma_{\mu\lambda}^\eta\Gamma_{\nu\eta}^\kappa$$

In the case of this question this is

$$\begin{aligned}R_{\phi\theta\phi}^\theta &= \partial_\theta\Gamma_{\phi\phi}^\theta - \partial_\phi\Gamma_{\phi\theta}^\theta + \Gamma_{\phi\phi}^\eta\Gamma_{\eta\theta}^\theta - \Gamma_{\phi\theta}^\eta\Gamma_{\eta\phi}^\theta \\ &= \partial_\theta(-\sin\theta\cos\theta) - 0 + 0 \cdot \Gamma_{\phi\phi}^\eta - \Gamma_{\phi\phi}^\theta\Gamma_{\phi\theta}^\phi \\ &= (-\cos^2\theta + \sin^2\theta) - (-\sin\theta\cos\theta)\left(\frac{\cos\theta}{\sin\theta}\right) \\ &= \sin^2\theta\end{aligned}$$

The Ricci Tensor

$$\begin{aligned}(\text{Ric})_{\mu\nu} &= R_{\mu\lambda\nu}^\lambda \\ &g^{ab}R_{a\mu b\nu}\end{aligned}$$

So we have

$$\begin{aligned}R_{\theta\theta} &= g^{ab}R_{a\theta b\theta} = g^{\theta\theta}R_{\theta\theta\theta\theta} + g^{\phi\phi}R_{\phi\theta\phi\theta} \\ &= (1 \cdot 0) + \left(\frac{1}{\sin^2\theta}\sin^2\theta\right) \\ &= 1\end{aligned}$$

And

$$\begin{aligned}R_{\phi\phi} &= g^{ab}R_{a\phi b\phi} = g^{\theta\theta}R_{\theta\phi\theta\phi} + g^{\phi\phi}R_{\phi\phi\phi\phi} \\ &= (1 \cdot \sin^2\theta) + \left(\frac{1}{\sin^2\theta} \cdot 0\right) \\ &= \sin^2\theta \\ R_{\theta\phi} &= g^{ab}R_{a\theta b\phi} = g^{\theta\theta}R_{\theta\theta\theta\phi} + g^{\phi\phi}R_{\phi\theta\phi\phi} \\ &= (1 \cdot 0) + \left(\frac{1}{\sin^2\theta} \cdot 0\right) \\ &= 0\end{aligned}$$

This is also the same for $R_{\phi\theta} = 0$. The Scalar curvature denoted by R_S is

$$\begin{aligned}R_S &= g^{\mu\nu}(\text{Ric})_{\mu\nu} = g^{\mu\nu}R_{\mu\nu} \\ &R_S = g^{\theta\theta}R_{\theta\theta} + g^{\phi\phi}R_{\phi\phi} \\ &= (1 \cdot 1) + \left(\frac{1}{\sin^2\theta} \cdot \sin^2\theta\right) \\ &= 2\end{aligned}$$

2. Symmetry of a sphere

$$g = d\theta \otimes d\theta + \sin^2 \theta d\phi \otimes d\phi$$

Killing vector fields given by

$$X^\xi \partial_\xi g_{\mu\nu} + \partial_\mu X^\alpha g_{\alpha\nu} + \partial_\nu X^\beta g_{\mu\beta} = 0$$

The metric is diagonal so we can write

$$X^\xi \partial_\xi g_{\mu\nu} + \partial_\mu X^\nu g_{\nu\nu} + \partial_\nu X^\mu g_{\mu\mu} = 0$$

The L_i killing vectors can replace X^j in the above equation for different μ, ν as either θ or ϕ . So we have Beginning with $\mu, \nu = \theta, \theta$

$$L_1^\xi \partial_\xi g_{\theta\theta} + \partial_\theta L_1^\theta g_{\theta\theta} + \partial_\theta L_1^\theta g_{\theta\theta} = 0$$

$$L_1^\xi \cdot 0 + 0 \cdot g_{\theta\theta} = 0$$

$$L_2^\xi \partial_\xi g_{\theta\theta} + \partial_\theta L_2^\theta g_{\theta\theta} + \partial_\theta L_2^\theta g_{\theta\theta} = 0$$

$$L_2^\xi \cdot 0 + 0 \cdot g_{\theta\theta} = 0$$

$$L_3^\xi \partial_\xi g_{\theta\theta} + \partial_\theta L_3^\theta g_{\theta\theta} + \partial_\theta L_3^\theta g_{\theta\theta} = 0$$

$$L_3^\xi \cdot 0 + 0 \cdot g_{\theta\theta} = 0$$

Now with $\mu, \nu = \phi, \phi$

$$L_1^\xi \partial_\xi g_{\phi\phi} + \partial_\phi L_1^\phi g_{\phi\phi} + \partial_\phi L_1^\phi g_{\phi\phi} = 0$$

$$-\cos \phi \partial_\theta \sin^2 \theta + (\partial_\phi \sin \phi \cot \theta) \sin^2 \theta + (\partial_\phi \sin \phi \cot \theta) \sin^2 \theta = 0$$

$$-2 \cos \phi \cos \theta \sin \theta + 2 \cos \phi \cot \theta \sin^2 \theta = 0$$

$$L_2^\xi \partial_\xi g_{\phi\phi} + \partial_\phi L_2^\phi g_{\phi\phi} + \partial_\phi L_2^\phi g_{\phi\phi} = 0$$

$$\sin \phi \partial_\theta \sin^2 \theta + (\partial_\phi \cos \phi \cot \theta) \sin^2 \theta + (\partial_\phi \cos \phi \cot \theta) \sin^2 \theta = 0$$

$$2 \sin \phi \cos \theta \sin \theta - 2 \sin \phi \cot \theta \sin^2 \theta = 0$$

$$L_3^\xi \partial_\xi g_{\phi\phi} + \partial_\phi L_3^\phi g_{\phi\phi} + \partial_\phi L_3^\phi g_{\phi\phi} = 0$$

$$0 \cdot \partial_\theta g_{\phi\phi} + 0 \cdot g_{\phi\phi} = 0$$

Now with $\mu, \nu = \phi\theta$

$$\underbrace{L_1^\xi \partial_\xi g_{\phi\theta}}_{=0} + \partial_\theta L_1^\phi g_{\phi\phi} + \partial_\phi L_1^\theta g_{\theta\theta} = 0$$

$$-\frac{\sin \phi g_{\phi\phi}}{\sin^2 \theta} + \cos \phi = 0$$

$$\underbrace{L_2^\xi \partial_\xi g_{\phi\theta}}_{=0} + \partial_\theta L_2^\phi g_{\phi\phi} + \partial_\phi L_2^\theta g_{\theta\theta} = 0$$

$$-\frac{\cos \phi g_{\phi\phi}}{\sin^2 \theta} + \sin \phi = 0$$

$$\underbrace{L_3^\xi \partial_\xi g_{\phi\theta}}_{=0} + \partial_\theta L_3^\phi g_{\phi\phi} + \partial_\phi L_3^\theta g_{\theta\theta} = 0$$

$$(0 \cdot 0) + (0 \cdot 0) = 0$$

And so, $L_{1,2,3}$ are the killing vectors of g . Now to calculate the commutators $[L_a, L_b]$

$$\begin{aligned} [L_1, L_2] &= [-\cos \phi \partial_\theta + \cot \theta \sin \phi \partial_\phi, \sin \phi \partial_\theta + \cot \theta \cos \phi \partial_\phi] \\ &= \cos \phi \partial_\theta \left(-\sin \phi \partial_\theta - \cot \theta \cos \phi \partial_\phi \right) - \cot \theta \sin \phi \partial_\phi \left(-\sin \phi \partial_\theta - \cot \theta \cos \phi \partial_\phi \right) \\ &\quad + \sin \phi \partial_\theta \left(\cos \phi \partial_\theta - \cot \theta \sin \phi \partial_\phi \right) + \cot \theta \cos \phi \partial_\phi \left(\cos \phi \partial_\theta - \cot \theta \sin \phi \partial_\phi \right) \\ &= -\cos \phi \cos \phi \partial_\theta^2 - \cos^2 \phi (\partial_\theta \cot \theta) \partial_\phi - \cot \theta \cos^2 \phi \partial_\theta \partial_\phi + \cot \theta \sin \phi \cos \phi \partial_\theta + \cot \theta \sin^2 \phi \partial_\phi \partial_\theta \\ &\quad - \cot^2 \theta \sin^2 \phi \partial_\phi + \cot^2 \theta \sin \phi \cos \phi \partial_\phi^2 + \sin \phi \cos \phi \partial_\theta^2 - \sin^2 \phi (\partial_\theta \cot \theta) \partial_\phi - \cot \theta \sin^2 \phi \partial_\theta \partial_\phi \\ &\quad - \cot \theta \sin \phi \cos \phi \partial_\theta + \cot \theta \cos^2 \phi \partial_\phi \partial_\theta - \cos^2 \theta \cos^2 \phi \partial_\phi - \cot^2 \theta \sin \phi \cos \phi \partial_\phi^2 \\ &= -\cot^2 \theta \partial_\phi - (\partial_\theta \cot \theta) \partial_\phi \\ &= \partial_\phi = \frac{\partial}{\partial \phi} = L_3 \end{aligned}$$

This is then repeated for $[L_2, L_3]$ and $[L_3, L_1]$ so we have

$$[L_1, L_2] = L_3$$

$$[L_2, L_3] = L_1$$

$$[L_3, L_1] = L_2$$

I think this represents the angular momentum operator.

3. T_a are $N \times N$ matrices

$$\begin{aligned}
[T_a, T_b] &= if_{abc}T_c \\
\chi_a &= \sum_{i,j=1}^N (T_a)_{ij} b_i^\dagger b_j \\
\chi_b &= \sum_{k,l=1}^N (T_b)_{kl} b_k^\dagger b_l \\
[\chi_a, \chi_b] &= \left[\sum_{i,j=1}^N (T_a)_{ij} b_i^\dagger b_j \right] \left[\sum_{k,l=1}^N (T_b)_{kl} b_k^\dagger b_l \right] \\
&\quad - \left[\sum_{k,l=1}^N (T_b)_{kl} b_k^\dagger b_l \right] \left[\sum_{i,j=1}^N (T_a)_{ij} b_i^\dagger b_j \right] \\
&= (T_a)_{ij} (T_b)_{kl} [b_i^\dagger b_j, b_k^\dagger b_l] \\
&= (T_a)_{ij} (T_b)_{kl} (b_i^\dagger b_j b_k^\dagger b_l - b_k^\dagger b_l b_i^\dagger b_j) \\
&= (T_a)_{ij} (T_b)_{kl} (b_i^\dagger [b_j, b_k^\dagger b_l] + [b_i^\dagger, b_k^\dagger b_l] b_j) \\
&= (T_a)_{ij} (T_b)_{kl} (b_i^\dagger [b_j, b_k^\dagger] b_l + b_k^\dagger [b_i^\dagger, b_l] b_j) \\
&\quad (T_a)_{ij} (T_b)_{kl} (b_i^\dagger \delta_{jk} b_l - b_k^\dagger \delta_{il} b_j) \\
&\quad (T_a T_b)_{il} b_i^\dagger b_l - (T_a T_b)_{kj} b_k^\dagger b_j \\
&= [T_a, T_b]_{ij} b_i^\dagger b_j = if_{abc} \chi_c
\end{aligned}$$

4. (a) The Gell-Mann matrices are

$$\begin{aligned}
\lambda_1 &= \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \lambda_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \\
\lambda_5 &= \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} \quad \lambda_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad \lambda_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} \quad \lambda_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}
\end{aligned}$$

The generators of $SU(3)$ are

$$T_1 = \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad T_2 = \frac{1}{2} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad T_3 = \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad T_4 = \frac{1}{2} \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$$

$$T_5 = \frac{1}{2} \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} \quad T_6 = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad T_7 = \frac{1}{2} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} \quad T_8 = \frac{1}{2\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

The generators T_1, T_2, T_3 form a $SU(2)$ subgroup of $SU(3)$ and since they are orthogonal for $a, b = \{1, 2, 3\}$, $c = \{4, 5, 6, 7\}$, $f_{abc} = 0$ and the only other values are

$$f_{147} = \frac{1}{2}$$

and

$$f_{458} = \frac{\sqrt{3}}{2}$$

Note that this also applies for other structure constants which aren't 0. The structure constants can also be calculated by the commutators of the generators as shown

$$[T_1, T_4] = \frac{1}{2}T_7, \quad \text{Tr}([T_1, T_4]T_c) = \delta^{c7} \rightarrow f_{147} = \frac{1}{2}$$

- (b) As with page 81 of the lecture notes we just need to show that $[J_a, J_b] = i\epsilon_{abc}J_c$, so we have

$$[\lambda_2, \lambda_5] = i\lambda_7$$

$$[\lambda_5, \lambda_2] = i\lambda_2$$

$$[\lambda_7, \lambda_2] = i\lambda_5$$

Thus, $\lambda_2 = J_1, \lambda_5 = J_2, \lambda_7 = J_3$ and clearly

$$[\lambda_a, \lambda_b] = i\epsilon_{abc}\lambda_c$$