

Open Quantum Systems Answers to Exercise Set 2 & 3 - JPBK

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1 Answers

1. Problem 2.1. Obtain heat in a nearly quasi-static drive in a two level system.
This question troubled me, in particular, the meaning behind "nearly quasi-static". I'm not sure what this means as a quasi-static process is one in which it is almost or nearly static or stationary so the term "nearly" here doesn't seem to have any meaning. I read the paper [4], mentioning quasistatic drive and I will assume that you wish us to study the Quasistatic limit in which W is close to the ideal limit of $-K_B T \ln(2)$ or the case where there is equality $W = Q$

There is also the case from where the Hamiltonian simply reads [10]

$$H = E_C(n - n_g)^2$$

For a single electron box where $E_C = \frac{e^2}{2C}$. The cited paper mentions 'sweeping' out the quasi-static drive to $n_g = 1/2$ allowing for the increase in the entropy of the charge system giving

$$\Delta S = K_B \ln(2)$$

And due to the process being quasi-static, ΔS is equal to the increase in the entropy of the bath so

$$\Delta Q = T \Delta S = K_B T \ln(2)$$

Maxwell's demon allows the system to be a "nearly" quasistatic driven system as the degeneracy can be 'ramped', turning the process cyclic and extracting heat of $K_B T \ln(2)$ from the bath. Experimentally, Jukka finds that the average heat dissipated is

$$Q_{avg} = -0.75 K_B T \ln(2)$$

2. Problem 2.2: Proof of Landauer's Principle
Consider a two level system with thermodynamic entropy $K_B \ln(2)$, which was proven by Von Neumann in 1949. The process of the erasure of the bit gives us a starting point of a microstate where the entropy is zero.
According to the 2nd law of thermodynamics, entropy cannot decrease - it must go somewhere-, so it is 'transferred' to the environment with temperature T giving us an energy cost of at least $K_B T \ln(2)$ so

$$E \geq K_B T \ln(2)$$

This is a very simple version where the proof as laid out by Von Neumann was not explicitly stated. Thermodynamic entropy cannot be so easily related to the information entropy. Jukka Pekola wrote a very interesting paper [3] which related the reversal to Maxwell's limit of thermodynamic efficiency and shows that it is consistent with the 2nd law of thermodynamics.

A more rigorous mathematical proof is the following:

The work extracted during isothermal quasistatic expansion is given by

$$\begin{aligned}\Delta Q &= \Delta W = \int_{\frac{V}{2}}^V \rho dV \\ &= \int_{\frac{V}{2}}^V \frac{K_B T}{V} dV \\ &= K_B T \ln(2)\end{aligned}$$

The second law of thermodynamics states that the entropy of an isolated system can never decrease with time. Applying this to the above fixes ΔQ

$$\Delta Q \geq K_B T \ln(2)$$

$$E \geq K_B T \ln(2)$$

This can also be proven from isothermal compression as well

$$\begin{aligned}\Delta Q &= -\Delta W \\ &= -\int_v^{\frac{V}{2}} \rho dV = -\int_V^{\frac{V}{2}} \frac{K_B T}{V} dV \\ &= K_B T \ln(2)\end{aligned}$$

3. Problem 3.1.

- (a) Why CBT in the universal regime where $E_C \ll K_B T$ can be considered a primary thermometer?

In a 'Primary' thermometer the measured property of matter is known so well that temperature can be calculated without any unknown quantities. Which is CBT we have

$$\begin{aligned}\frac{G}{G_T} &= 1 - \frac{2E_C}{K_B T} g\left(\frac{eV}{2K_B T}\right) \\ g(x) &= \frac{x \sinh(x) - 4 \sinh^2(x/2)}{8 \sinh^4(x/2)}\end{aligned}$$

So

$$\frac{G}{G_T} = 1 - \frac{e^2/C}{K_B T} \frac{x \sinh(x) - 4 \sinh^2(x/2)}{8 \sinh^4(x/2)}$$

Where

$$x = \frac{eV}{2K_B T}$$

And for $N = 2$ case

$$E_C = \frac{e^2}{2C}$$

In the case where $E_C \ll K_B T$ the normalized conductance does not depend on the CBT's dimensions and can be considered a primary thermometer (As $2E_c/K_B T \rightarrow 0$)

- (b) If we consider the full width half minimum of the normalized conductance so $V = V_{1/2}$ and set this to λ

$$\frac{eV_{1/2}}{NK_B T} = \lambda$$

$$V_{1/2} = \frac{\lambda NK_B T}{e}$$

The numerical number was calculated to be

$$\lambda = 5.434$$

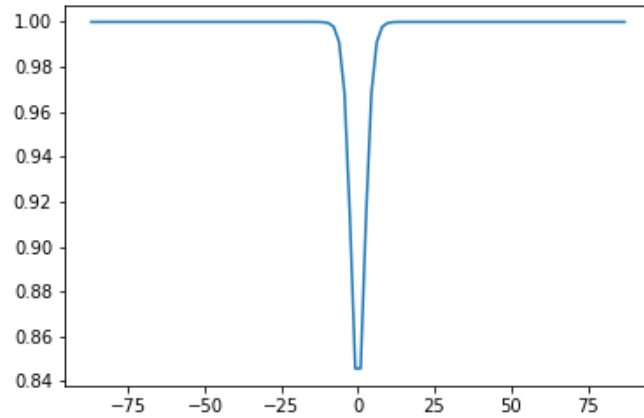


Figure 1: Plot of the normalized conductance using matplotlib. Axis values aren't quite correct but this shows an inverse bell curve. A full width half minimum approach can be used to calculate $V_{1/2}$

2 Note

I found these exercises particularly hard, mainly due to the nature of the questions, however I believe this was due to getting use to the particular thermodynamic jargon used of which I am not used to. I found an abundance of papers and research completed online (mostly by Jukka Pekola) which taught me a great deal and helped supplement what was taught in the lectures, however did not help me to complete the questions theoretically. I have attached references in my bibliography at the bottom to works which I heavily relied on.

References

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